## Lectures on the Soft-Collinear Effective Theory

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## Problem 3) Decoupling of Soft and Collinear Gluons in $SCET_{II}$

Consider an operator with one collinear quark and one soft quark in  $\text{SCET}_{\text{II}}$ . This case differs from  $\text{SCET}_{\text{I}}$  in that soft gluons knock collinear quarks offshell and collinear gluons knock soft quarks offshell. To be definite lets make the soft quark a heavy quark from HQET,  $h_v$ , and the energetic quark a massless collinear quark,  $\xi_n$ . The SCET operator will be

$$(\bar{\xi}_n W_{\bar{n}}) \Gamma(S_n^{\dagger} h_v) \tag{1}$$

where W is the Wilson line we saw in problem 2, and  $S_n$  is an exact analog of the Y from lecture, but with  $A_s$  gluons rather than ultrasoft gluons. To find Eq. (1) requires tree level matching (to determine the direction and gluon components appearing in the Wilson lines), and gauge invariance (to ensure that loops do not spoil the structure so obtained). The first non-trivial term in perturbation theory have one soft and one collinear gluon. By computing the QCD graphs in Figure 1 expanded to LO, verify that the three-gluon interactions are responsible for putting the gluons in the right order in Eq. (1).

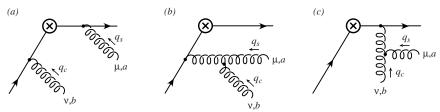


FIG. 1: QCD graphs with collinear,  $q_c$ , and soft,  $q_s$ , momenta.

## Problem 4) SCET Loops for Two-Jet Production

Consider the two-jet production process through a virtual photon in SCET, namely  $e^+e^- \rightarrow J_n J_{\bar{n}} X_{us}$  where  $J_n$  is a jet in the n = (1, 0, 0, -1) direction,  $J_{\bar{n}}$  is a jet in the  $\bar{n} = (1, 0, 0, 1)$  direction, and any remaining particles in the final state are ultrasoft, contained in  $X_{us}$ .

a) Write down two collinear quark Lagrangians, one for  $\xi_n$  fields and one for  $\xi_{\bar{n}}$  fields. Interactions between these two types of collinear fields are hard, and so do not effect your analysis. What are the Feynman rules for the ultrasoft gluon coupling to each of these collinear quarks?

b) Start with  $J^{\text{QCD}} = \bar{\psi}\gamma_{\mu}\psi$  and determine the appropriate LO SCET current  $J^{\text{SCET}} = \bar{\xi}_n \cdots \bar{\xi}_n$ , ie. fill in the dots with appropriate collinear Wilson lines and Dirac structure.

c) Draw the five one-loop Feynman diagrams that are non-zero for  $e^+e^- \rightarrow q_n\bar{q}_n$  (use Feynman gauge for all gluons when determining which graphs are zero). Here  $q_n$  has *n*-collinear momentum p, and  $\bar{q}_n$  has  $\bar{n}$ -collinear momentum  $\bar{p}$  and you should work in the CM frame. All graphs but one can be directly read off using the loop computations done in lecture (or given in the handout notes), as long as you use the same IR regulator. That is, you should keep both collinear quarks offshell,  $p^2 \neq 0$  and  $\bar{p}^2 \neq 0$ . Compute the divergent terms in the one remaining ultrasoft graph using dimensional regularization in the UV.

d) Add up the  $1/\epsilon$  terms from the graphs in c) and determine the lowest order anomalous dimension equation for C the Wilson coefficient of  $J^{\text{SCET}}$ . Solve this equation keeping only the  $\ln \mu/Q$  term and using a fixed coupling  $\alpha_s$ , and then with a running coupling  $\alpha_s(\mu)$ . (Voilá, Sudakov double logs resummed.)