

# Lectures on Quark Flavor Physics: Exercises

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Benasque Summer School on Flavor Physics, July 2008

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### 1. Yukawa couplings, CKM matrix, and unitarity triangles:

a) Show that flavor non-diagonal kinetic terms in the Standard Model Lagrangian can always be diagonalized and brought into standard form by field redefinitions. To this end, study the Lagrangian

$$\mathcal{L}_{\text{kinetic}} = \bar{Q}_L Z_Q i\not{D} Q_L + \bar{u}_R Z_u i\not{D} u_R + \bar{d}_R Z_d i\not{D} d_R,$$

where all fields are 3-component vectors in generation space, and  $Z_A$  are non-negative, hermitian  $3 \times 3$  matrices.

b) Show that an arbitrary complex matrix  $Y$  can be diagonalized by a biunitary transformation:

$$W^\dagger Y U = \lambda,$$

where  $U, W$  are unitary matrices, and  $\lambda$  is a real, diagonal matrix with non-negative eigenvalues. (*Hint:* Consider the matrices  $YY^\dagger$  and  $Y^\dagger Y$ .)

c) Derive the number of mixing angles and physical (i.e., observable) phases of the CKM matrix for the Standard Model with  $N$  fermion generations.

d) Show that the Jarlskog determinant  $J$  defined as

$$\text{Im} \left( V_{ij} V_{kl} V_{il}^* V_{kj}^* \right) = J \sum_{m,n} \epsilon_{ikm} \epsilon_{jln} \quad (i \neq k, j \neq l)$$

is invariant under phase redefinitions of the quark fields, and calculate its value in terms of the Wolfenstein parameters to leading nontrivial order in  $\lambda$ .

e) Show that all unitarity triangles have the same area  $J/2$ .

### 2. Matching of Wilson coefficients in the effective weak Hamiltonian:

Assume that, in addition to its standard interactions, the  $Z^0$  boson has a small flavor-changing coupling to left-handed  $b$  and  $s$  quarks:

$$\mathcal{L}_Z = \frac{g_2}{\cos \theta_W} Z^\mu \left\{ \sum_f \bar{f} \gamma_\mu \left( T_f^3 \frac{1 - \gamma_5}{2} - Q_f \sin^2 \theta_W \right) f + \left( \varepsilon_{bs} \bar{s} \gamma_\mu \frac{1 - \gamma_5}{2} b + \text{h.c.} \right) \right\},$$

where  $|\varepsilon_{bs}| \ll 1$ . The sum in the first term is over all Standard Model fermions.  $T_f^3$  is the third component of weak isospin,  $Q_f$  the electric charge in units of  $e$ ,  $g_2$  the SU(2) gauge coupling, and  $\theta_W$  the weak mixing angle.

Calculate the contributions to the Wilson coefficients  $C_{3-10}$  in the effective weak Hamiltonian for  $b \rightarrow s\bar{q}q$  transitions arising from tree-level  $Z$ -boson exchange, working to first order in  $\varepsilon_{bs}$ . Recall that  $m_Z \cos \theta_W = m_W$  and  $G_F/\sqrt{2} = g_2^2/8m_W^2$ . Use the fact that  $T_f^3 = 0$  for right-handed quarks, while  $T_f^3 = Q_f - Y$  with  $Y = 1/6$  for left-handed quarks.

### 3. OPE for inclusive $B$ -meson decays:

Calculate the leading contribution (in powers of  $\alpha_s$  and  $1/m_b$ ) to the inclusive  $\bar{B} \rightarrow X_s \gamma$  decay rate using the optical theorem,

$$\Gamma(\bar{B} \rightarrow X_s \gamma) = \text{Disc} \frac{\langle \bar{B} | \mathbf{T} | \bar{B} \rangle}{2m_B}, \quad \mathbf{T} = i \int d^4x T \{ \mathcal{H}_{\text{eff}}(x) \mathcal{H}_{\text{eff}}(0) \},$$

by evaluating the discontinuity of the one-loop diagram with two insertions of the dipole operator  $Q_{\tau\gamma}$  in the effective weak Hamiltonian.

a) Show that the Feynman rule (in momentum space) for the dipole operator

$$Q_{\tau\gamma} = -\frac{em_b}{4\pi^2} \bar{s} \sigma_{\mu\nu} \frac{1 + \gamma_5}{2} F^{\mu\nu} b$$

is

$$-\frac{em_b}{4\pi^2} [\gamma_\mu, \not{q}] \frac{1 + \gamma_5}{2},$$

where  $q$  is the outgoing photon momentum and  $\mu$  its polarization index.

b) Prove the following identity for Dirac matrices in 4 dimensions:

$$[\gamma_\mu, \not{q}] \gamma^\alpha [\gamma^\mu, \not{q}] = 16q^\alpha \not{q} - 4q^2 \gamma^\alpha$$

c) Evaluate the discontinuity of the partonic  $b \rightarrow s \gamma$  forward scattering amplitude by using the Cutkosky rule to replace the cut propagators with

$$\text{Disc} \frac{1}{q^2 + i\epsilon} \frac{1}{(p - q)^2 + i\epsilon} = (2\pi) \delta(q^2) \theta(q^0) (2\pi) \delta((p - q)^2) \theta(p^0 - q^0).$$

Use Feynman gauge, set the strange-quark mass to zero, and work in the  $b$ -quark rest frame, where the on-shell  $b$ -quark momentum is  $p^\mu = (m_b, 0, 0, 0)$ . First show that

$$\int \frac{d^4q}{(2\pi)^4} (2\pi) \delta(q^2) \theta(q^0) (2\pi) \delta((p - q)^2) \theta(p^0 - q^0) = \int \frac{d\Omega_q}{32\pi^2}, \quad q^0 = |\vec{q}| = \frac{m_b}{2}$$

in that frame, where  $d\Omega_q$  is the measure for the angular integration over the direction of the vector  $\vec{q}$ . Then use the results from parts a) and b).

d) Make the replacement

$$\bar{u}(p) \dots u(p) \rightarrow \frac{\langle \bar{B} | \bar{b} \dots b | \bar{B} \rangle}{2m_B}$$

for the  $b$ -quark spinor product and evaluate the resulting matrix element to obtain the total hadronic  $\bar{B} \rightarrow X_s \gamma$  decay rate.