

Observation of Fractional Quantum Hall effect in suspended graphene



Xu Du, Ivan Skachko,



Fabian Duerr, Adina Luican



Eva Y. Andrei



- 1'st report: May, 2009 (Science brevia submission)
- reported June 25, 09, EPQHS3 Lucca, Italy
- reported July 1, 09, "Recent Progress in Graphene", Seoul, Korea
- reported July 28, 09 "Graphene 2009", Benasque, Spain

Interaction and correlation effects in suspended graphene using STM and Transport

- Transport in Suspended graphene
 - Reduced potential fluctuations

- B finite

- Quantum Hall effect
 - $\nu=0$ insulating phase

IQHE $\nu=1,3,4$
FQHE $\nu=1/3$

- STM in suspended graphene

- Structure
- Density of States

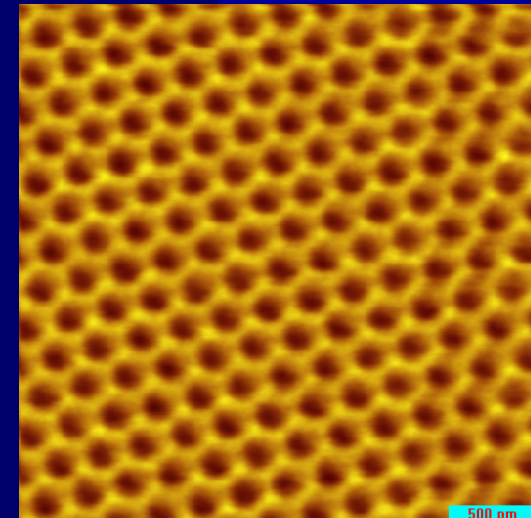
- B finite

- Landau levels
- Fermi Velocity
- e-ph interactions
- e-e interactions

- Rotated graphene

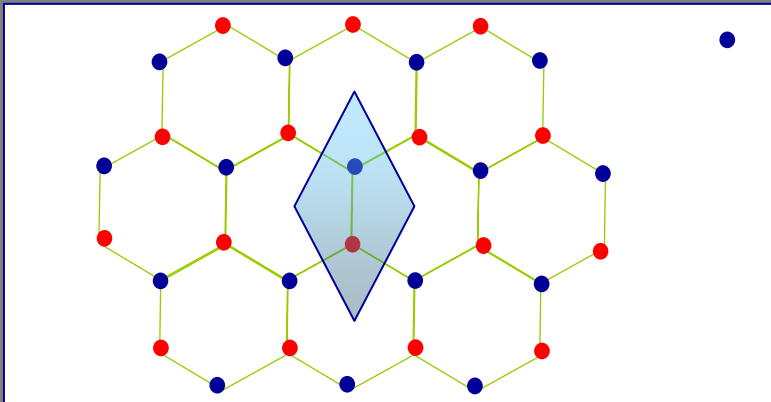
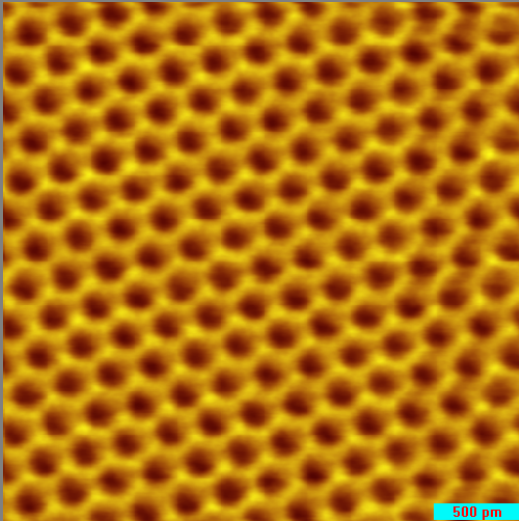
- Tunable Van-Hove singularities

see talk by J. Lopes dos Santos



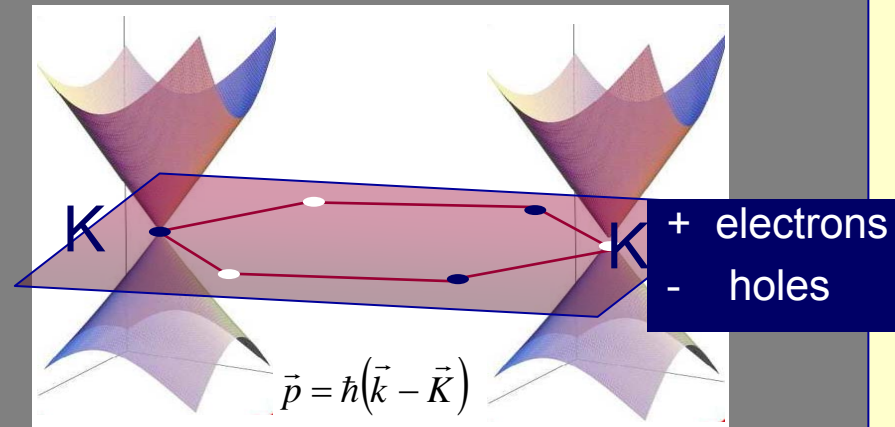
Graphene – the hallmarks

Structure: honeycomb



$$\Psi_{\vec{k}}(\vec{r}) = c_A |\text{blue circle}_A\rangle + c_B |\text{red circle}_B\rangle = \begin{pmatrix} c_A \\ c_B \end{pmatrix}$$

Band structure



Degeneracy 2-spin x 2-valley

$$H_K^{eff} = v_F \vec{\sigma} \cdot \vec{p}$$

Dirac-Weyl

$$E = \pm v_F |\vec{p}|$$

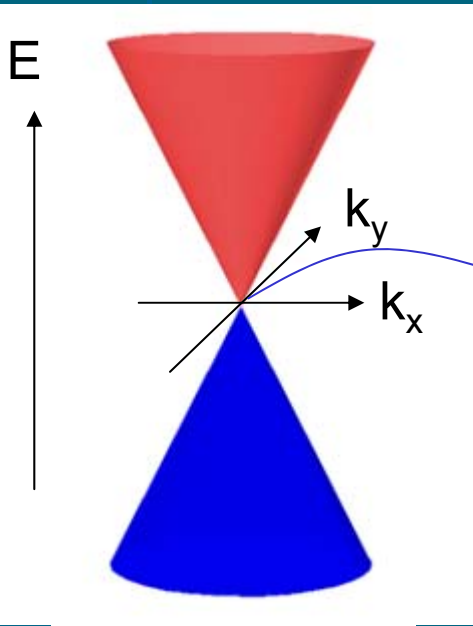
Zero mass
 $v_f \sim 10^6$ m/s

$$\psi = \begin{pmatrix} c_A \\ c_B \end{pmatrix} = \frac{1}{\sqrt{2}} e^{i\vec{k}\vec{r}} \begin{pmatrix} e^{-i\phi/2} \\ e^{i\phi/2} \end{pmatrix}$$

spinor



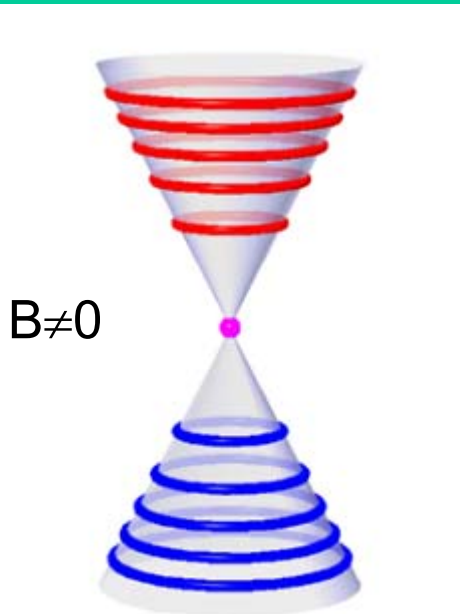
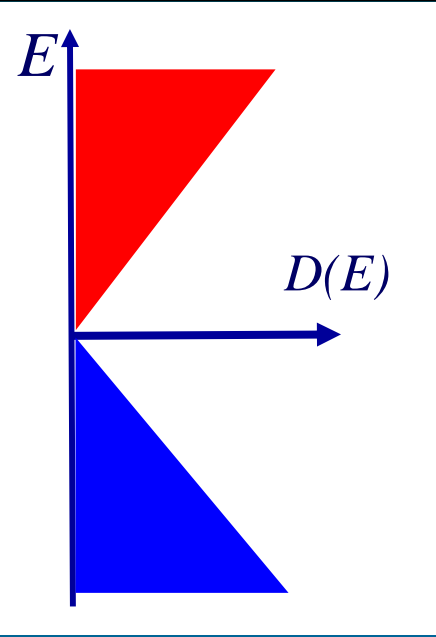
Density of states



$$E(k) = v_F \hbar k$$

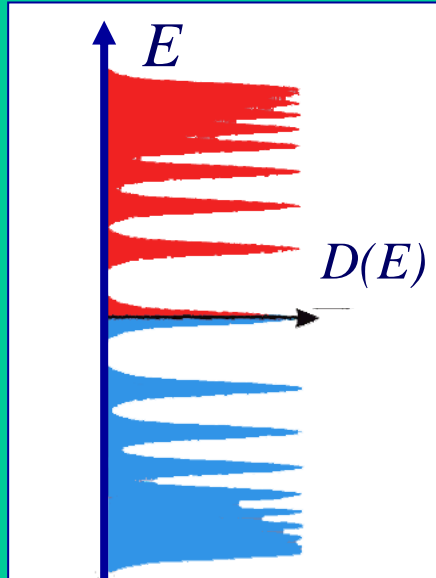
Dirac Point

$$D(E) = \frac{2}{\pi \hbar^2} \frac{1}{v_F^2} |E|$$



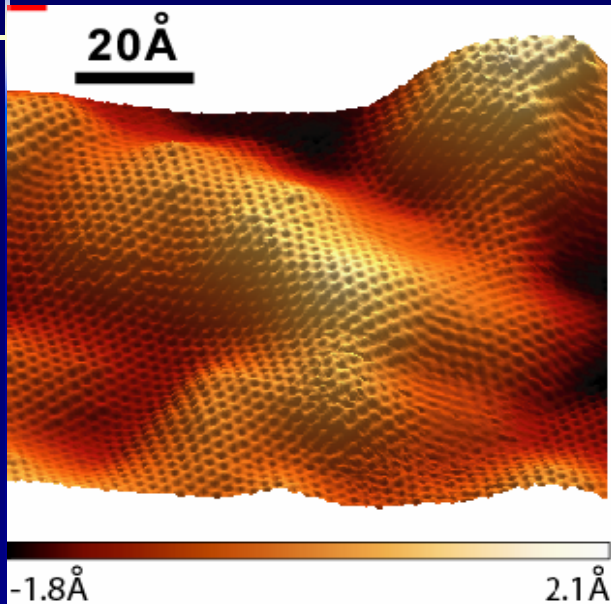
Landau Levels

$$E_n = \text{sgn}(n) \sqrt{2e\hbar v_F^2 |n| B}$$



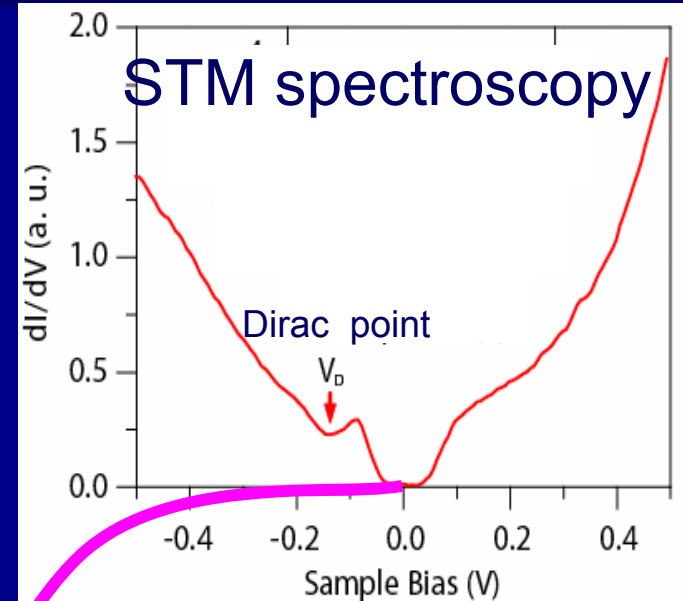
STM - Graphene on insulating substrates (SiO_2 , SiC)

STM topography



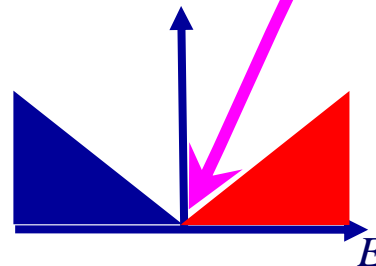
W. Brar, et al arxiv:0802.4315

STM spectroscopy

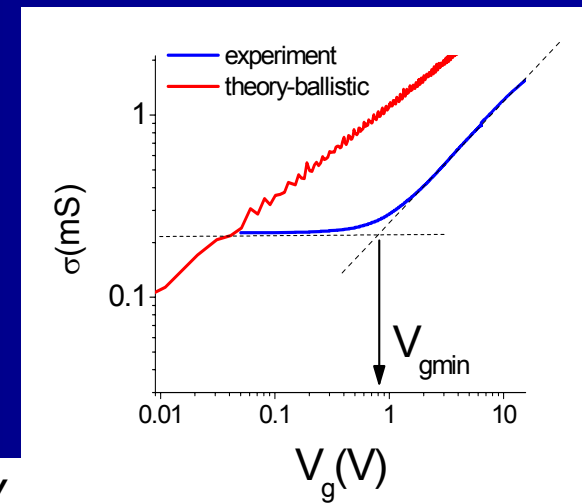
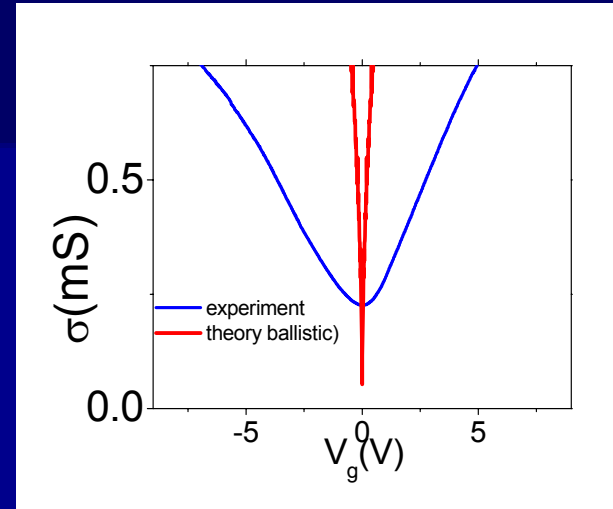
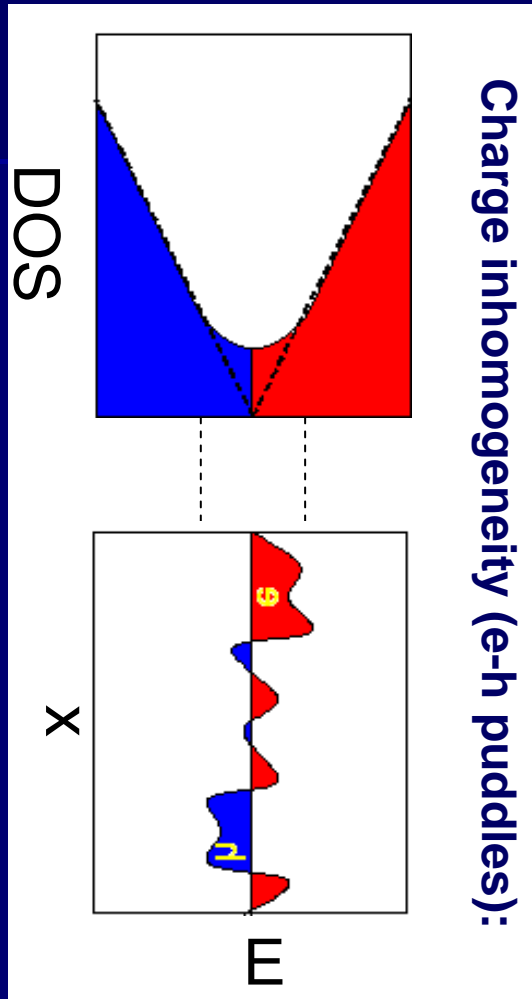
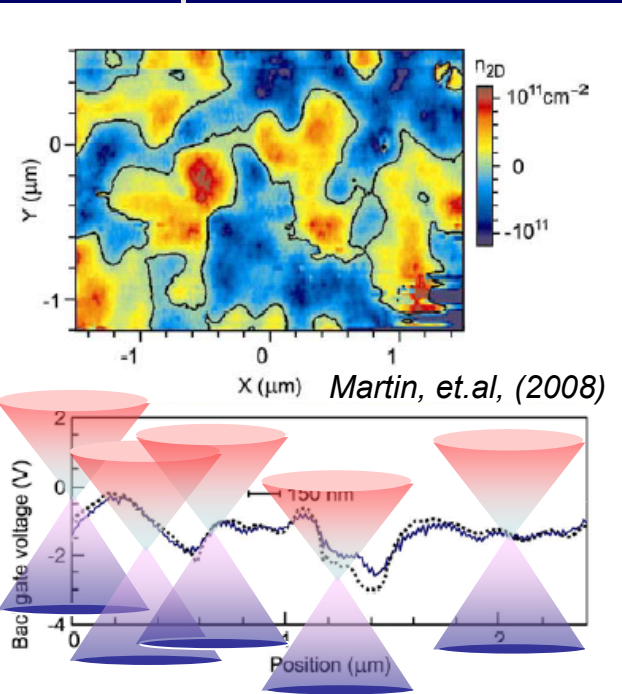


No Dirac cone

$D(E)$



Graphene on SiO_2 : e-h puddles and



e-h puddles : smeared Dirac point
 $\rightarrow n_{min}$ minimum carrier density

$\Delta V_g \sim 1-10\text{V}$
 $n_{min} \sim 10^{11} - 10^{12} \text{cm}^{-2}$
 $(\Delta E_F \sim 30-100 \text{meV})$



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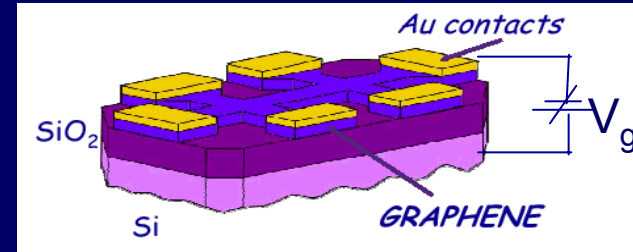
- Rotated graphene

Benasque-09

- Tunable Van Hove singularities



Transport: graphene on SiO₂



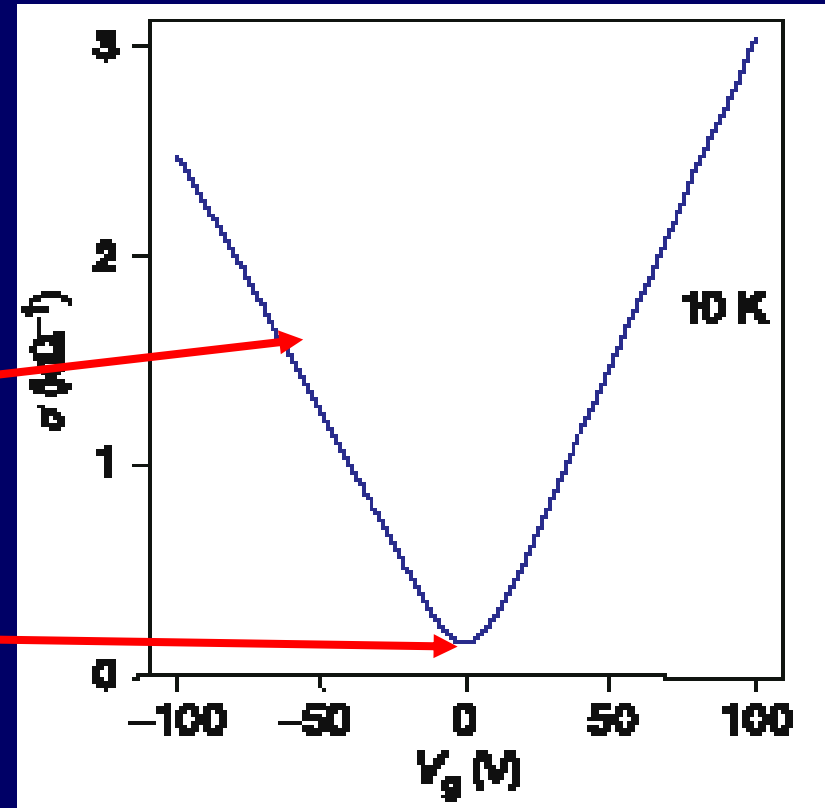
Experiment

Far from Dirac point

$$\sigma \propto n$$

At Dirac point

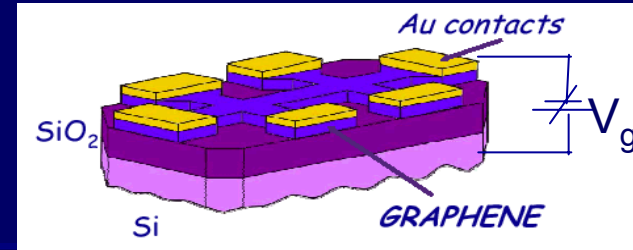
$$\sigma_{\min} = \frac{4e^2}{h}$$



Novoselov et al, Nature 2005
Y.W. Tan et al 2007



Transport: graphene on SiO₂



	Experiment	Theory: Ballistic transport
		$G(\mu) = \frac{4e^2}{h} \sum_n t_n(\mu)$
Far from Dirac point	$\sigma \propto n$	$\sigma \propto n^{1/2} \quad n \gg 10^8 \text{ cm}^{-2}$
At Dirac point	$\sigma_{\min} = \frac{4e^2}{h}$	$\sigma_{\min} = \frac{4e^2}{\pi h} \quad (L \ll W)$

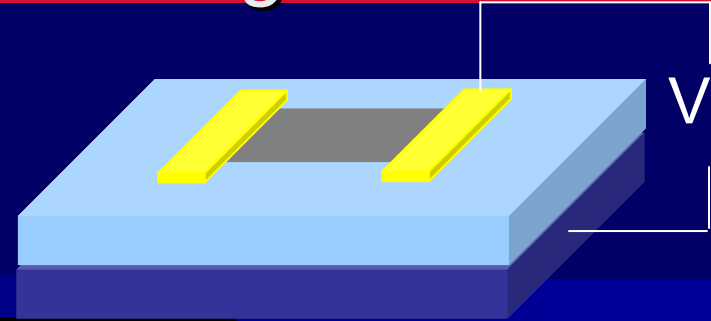
Novoselov et al, Nature 2005
Y.W. Tan et al 2007

Tworzydło et al PRL 2006
Katsnelson et.al, 2006



Conductivity and scattering

$$\sigma = \frac{e^2 v_F^2 N(E_F) \tau(k_F)}{2}$$



scattering	Source	Contributions
Long range	Charged impurities	$\sigma \sim n$ $\mu \sim \text{const}$ $l_m \sim n^{1/2}$
Short-range	atomic roughness Neutral impurities	$\sigma \sim \text{const}$ $\mu \sim n^{-1}$ $l_m \sim n^{-1/2}$
Midgap states	Vacancies boundaries corrugations	$\sigma \sim n [\ln(n^{1/2}R_0)]^2$ $\mu \sim [\ln(n^{1/2}R_0)]^2$ $l_m \sim n^{1/2} [\ln(n^{1/2}R_0)]^2$
Ballistic		$\sigma \sim n^{1/2}$ $\mu \sim n^{-1/2}$ $l_m = L/2$

S. Adam et al 2007

$$E_F = \hbar v_F k_F \sim n^{1/2}$$

$$\sigma = ne\mu$$

$$l = v_F \tau$$

Peres et al PRB 2006
T. Stauber, et al (2007).



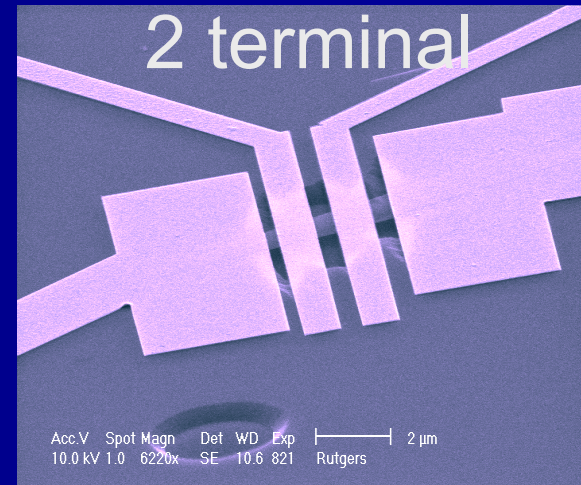
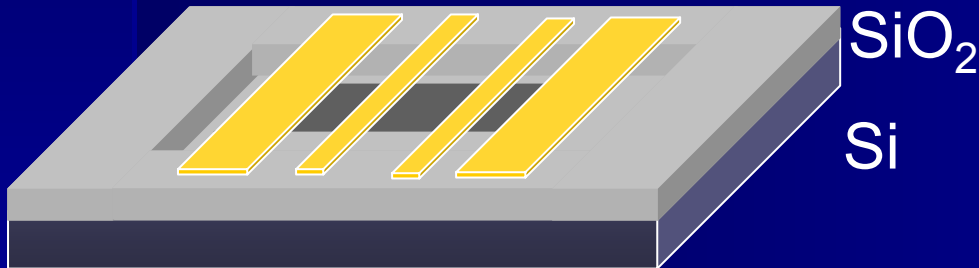
Suspended Graphene

2-terminal technique :

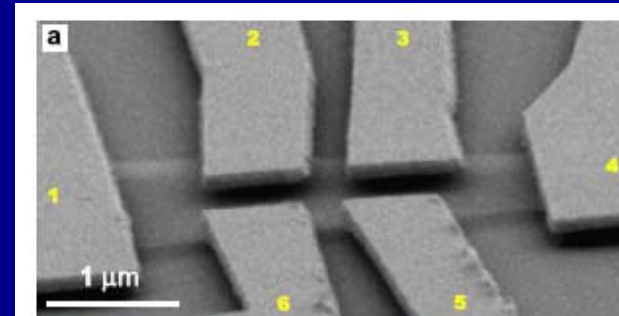
X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)

- Substrate roughness
- Trapped charges
- Quench condensed ripples

Get rid of the substrate!



4-terminal (Hall bar)



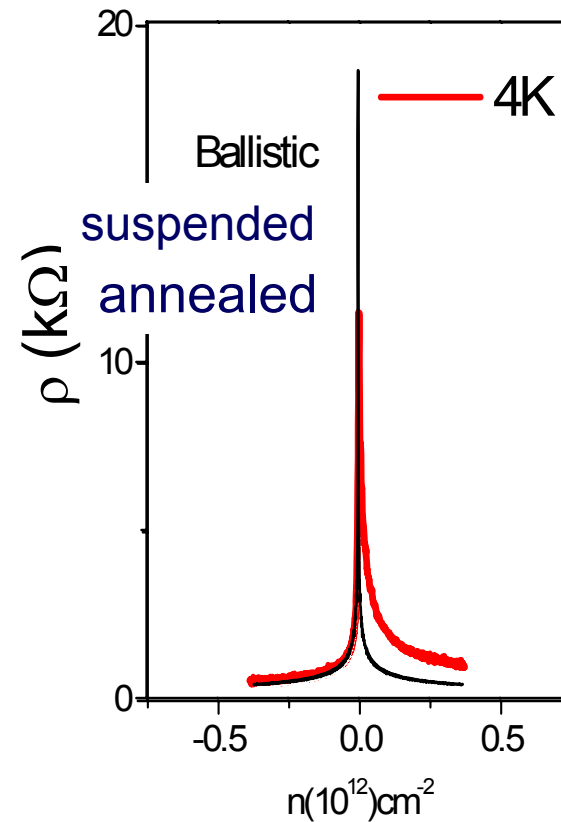
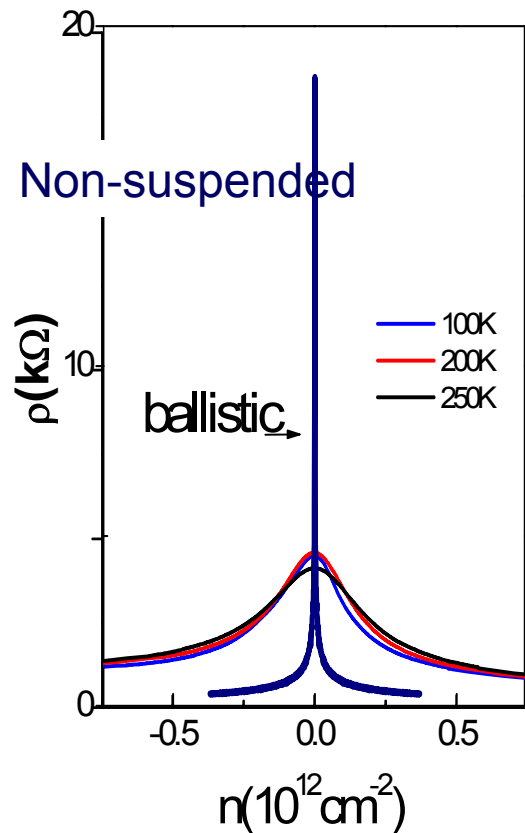
4-terminal

Bolotin et al , Solid State Communications (2008)

Suspended Graphene: T dependence

2-terminal technique :

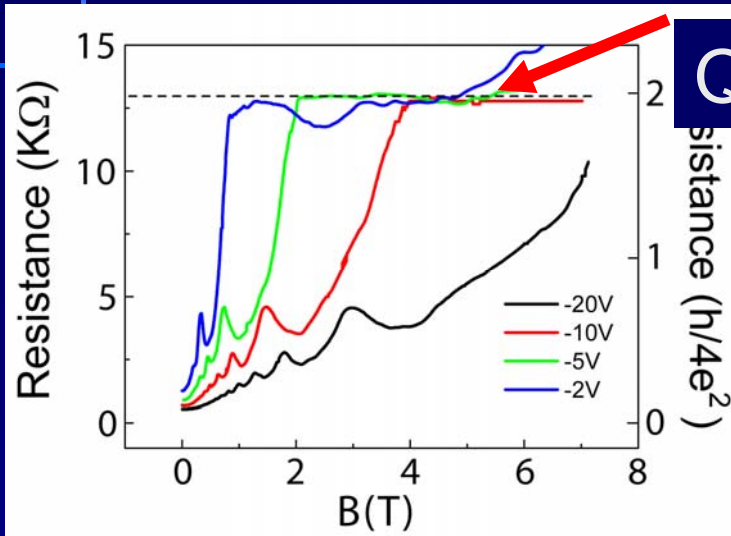
X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)



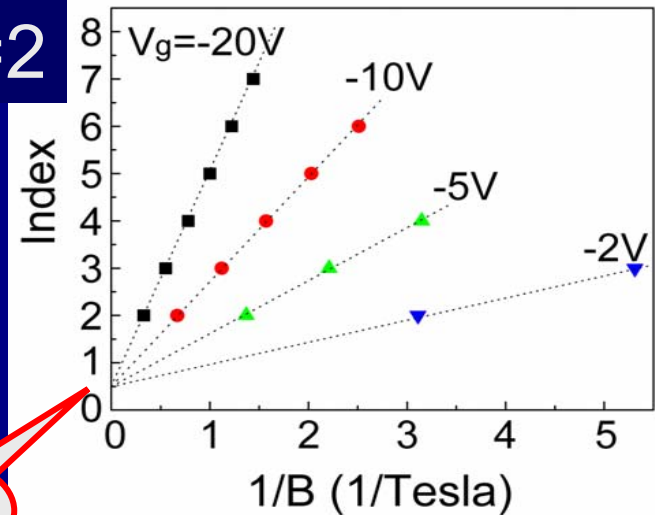
$$\alpha_{SG} = n/V_g$$

ShdH and QHE – density calibration

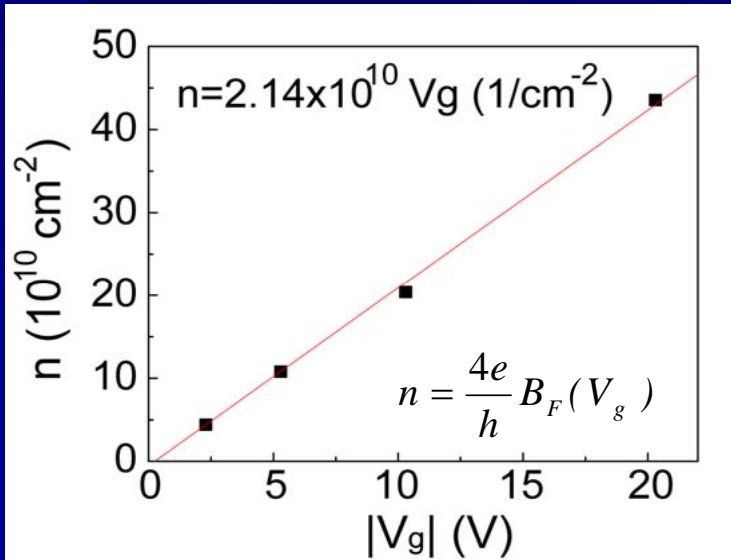
2-terminal technique X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)



QHE: $\nu=2$



1/2



➤ Single layer

$$\alpha_{SG} = n/V_g = 2.14 \times 10^{10} \text{ cm}^{-2} \text{ V}^{-1}$$

$$\alpha_{SG} / \alpha_{NSG} \approx \epsilon_{SiO_2}$$

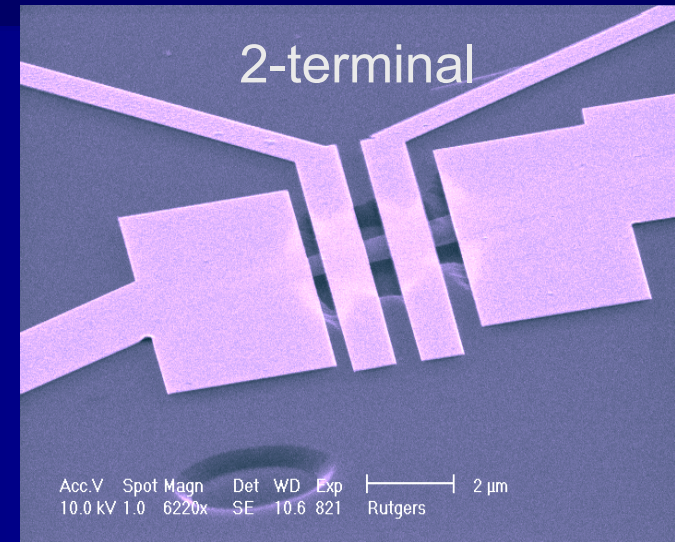
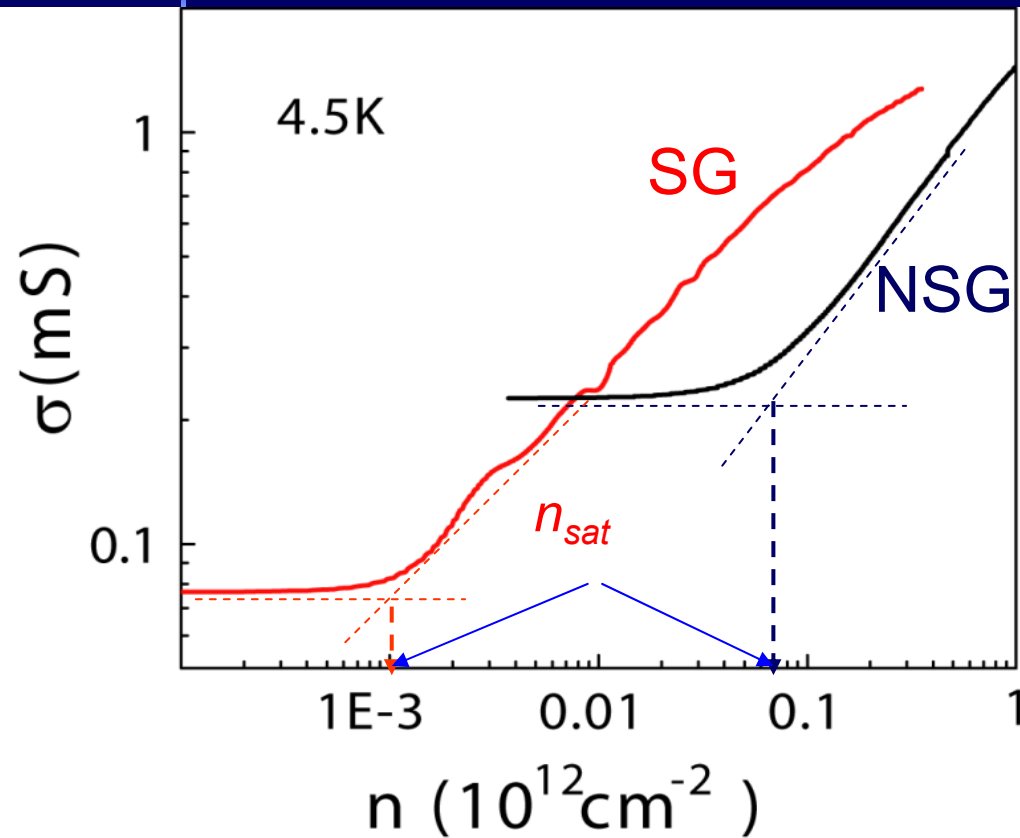
➤ Suspended



Suspended Graphene: reduced residual carriers

2-terminal technique :

X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)



NSG: $n_{\text{sat}} \sim 10^{10} - 10^{11} \text{cm}^{-2}$

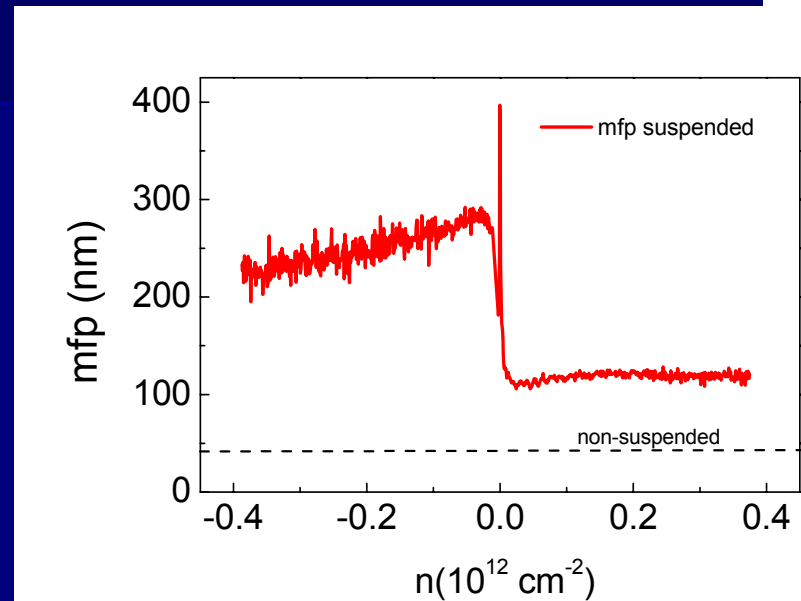
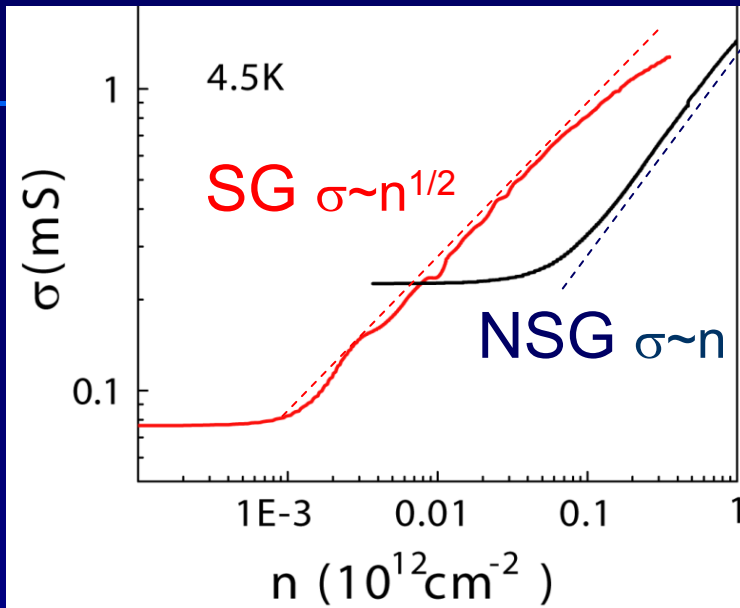
SG: $n_{\text{sat}} \sim 10^9 \text{cm}^{-2}$



Suspended Graphene: approaching ballistic transport

2-terminal technique :

X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)



suspended

$$n_{\text{sat}} \sim 10^9 \text{cm}^{-2}$$

$$\sigma \sim n^{1/2} \text{ Ballistic}$$

$$\mu \sim n^{-1/2}$$

$$\mu_{\text{sat}} \sim 2 \times 10^5 - 10^6 \text{cm}^2/\text{V s}$$

Non suspended

$$n_{\text{sat}} \sim 10^{11} \text{cm}^{-2}$$

$$\sigma \sim n - \text{Long range scatterers}$$

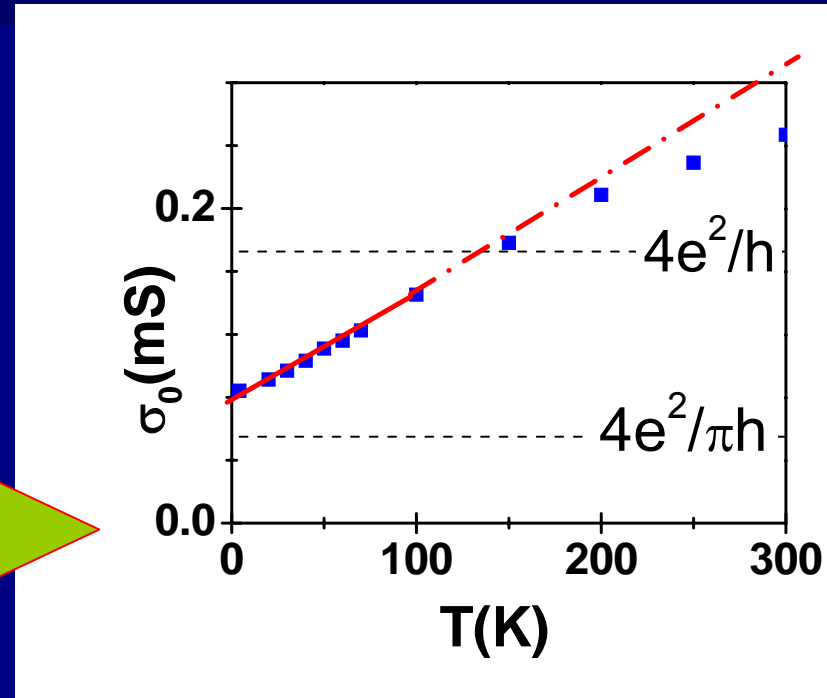
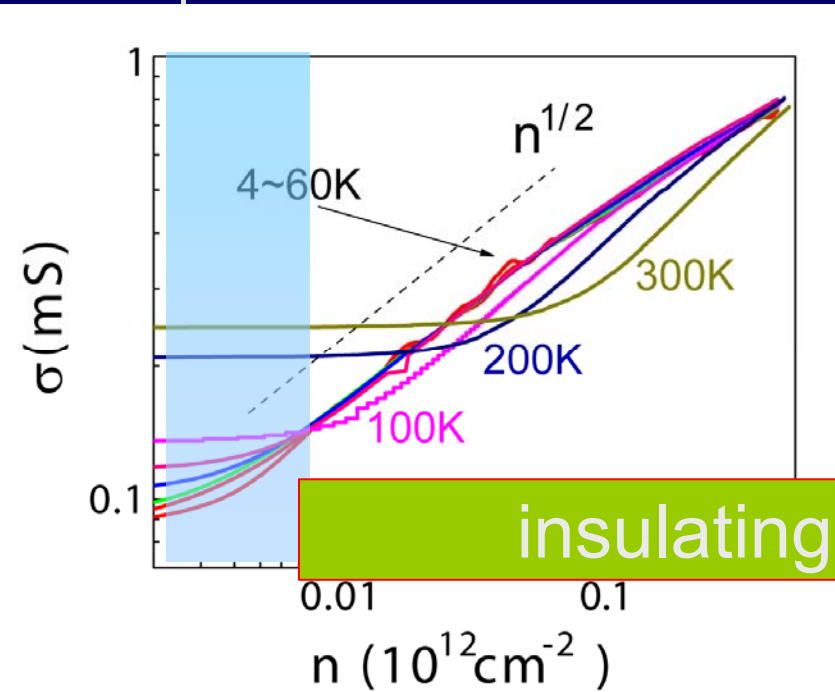
$$\mu \sim 1 - 2 \times 10^4 \text{cm}^2/\text{V s}$$



Suspended Graphene: T dependence

2-terminal technique :

X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)



$T < 100\text{K}$: Approaching Ballistic
 $\sigma \sim n^{1/2} \sim E_F$

$T > 100\text{K}$:
Long-range scattering, ripples?

- $\sigma_{\min} \sim T$ for $T < 150\text{K}$



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 - Fractional Quantum Hall effect
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Quantum Hall Effect

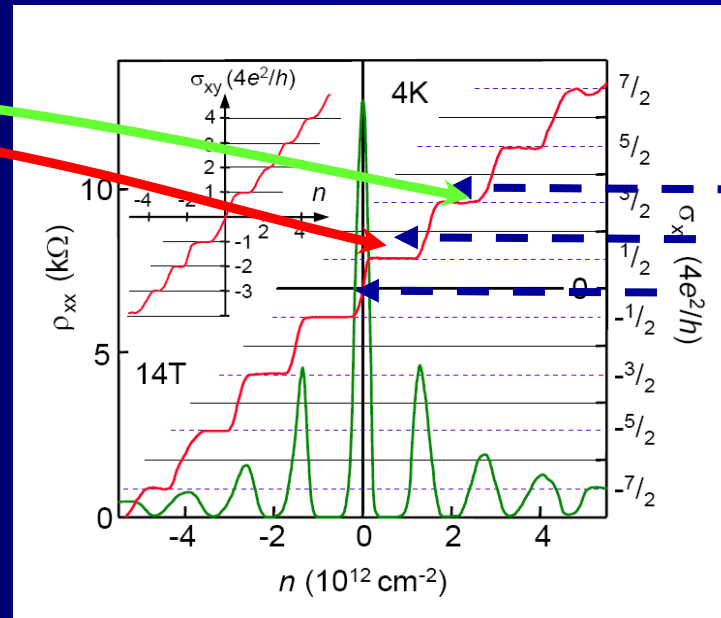
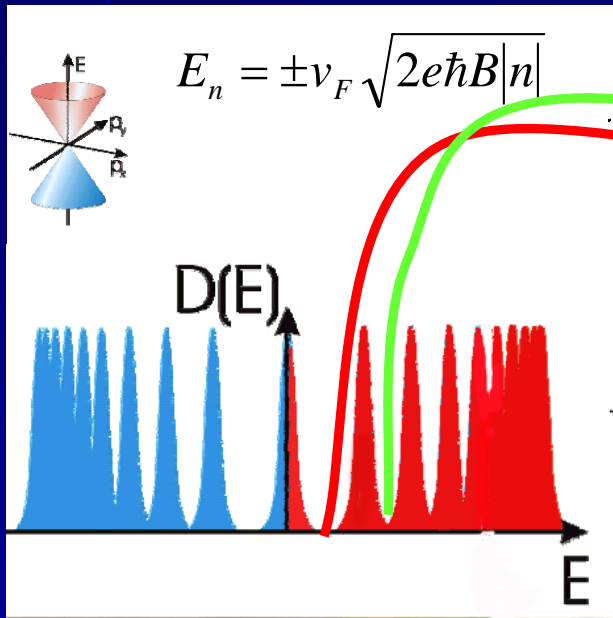
Each filled Landau level contributes g quanta of Hall conductance
(g = degeneracy)

$$\sigma_{xy} = \nu \frac{e^2}{h}$$

$$\nu = g_s g_v (n + 1/2) = \pm 2, \pm 6$$

$$g = 4$$

non-interacting sequence



non-relativistic
 $\nu = 2n \quad n = 0, 1, \dots$
 $g = 2$

Landau level at $E=0$:
no QHE plateau at 0.

Novoselev et al Nature 2005
Zhang et al Nature 2005

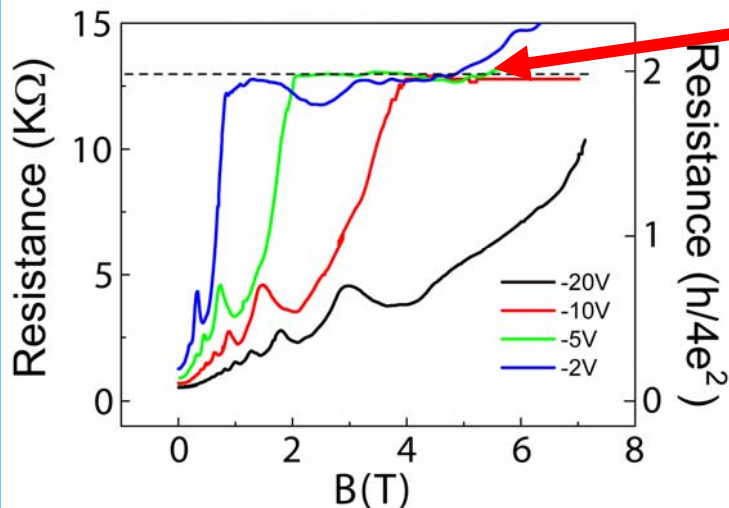


$$\alpha_{SG} = n/V_g$$

QHE: 2-terminal versus 4-terminal

suspended graphene with 2-terminal technique

X. Du, I. Skachako, A. Barker, E. Y. A. Nature Nanotech. 3, 491 (2008)

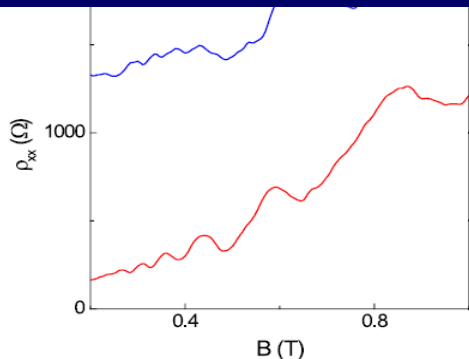


QHE: $\nu=2$

- well defined QHE plateaus clearly visible at low density $n \sim 2 \times 10^{10} \text{ cm}^{-2}$.

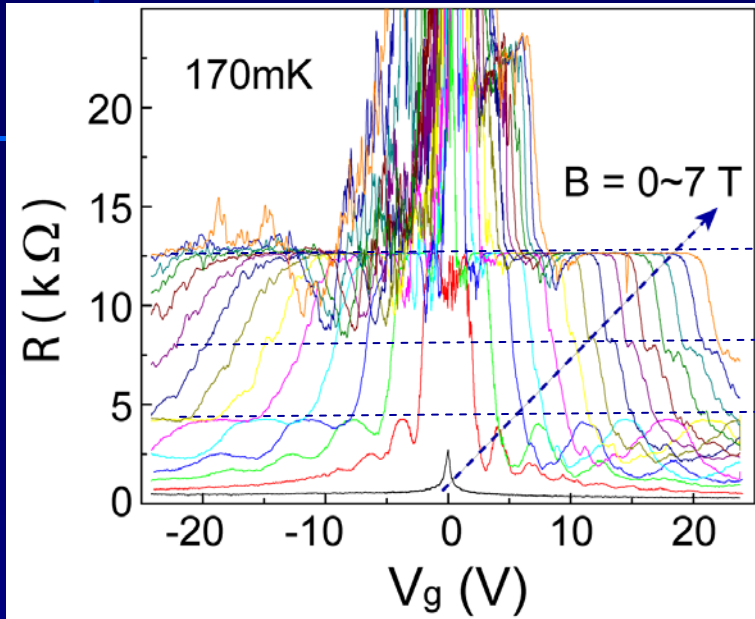
compare:

suspended graphene with 4-terminal technique Bolotin et al ,
Solid State Communications (2008)



- No evidence of QHE plateaus!
- even at higher density $n \sim 10^{11} \text{ cm}^{-2}$. !

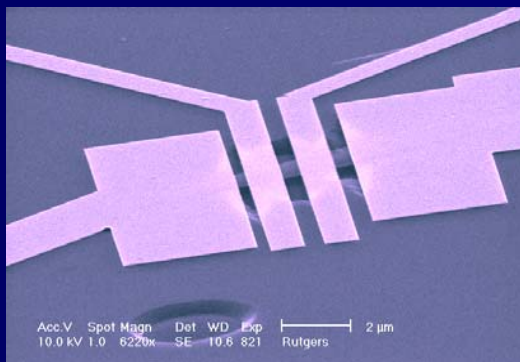
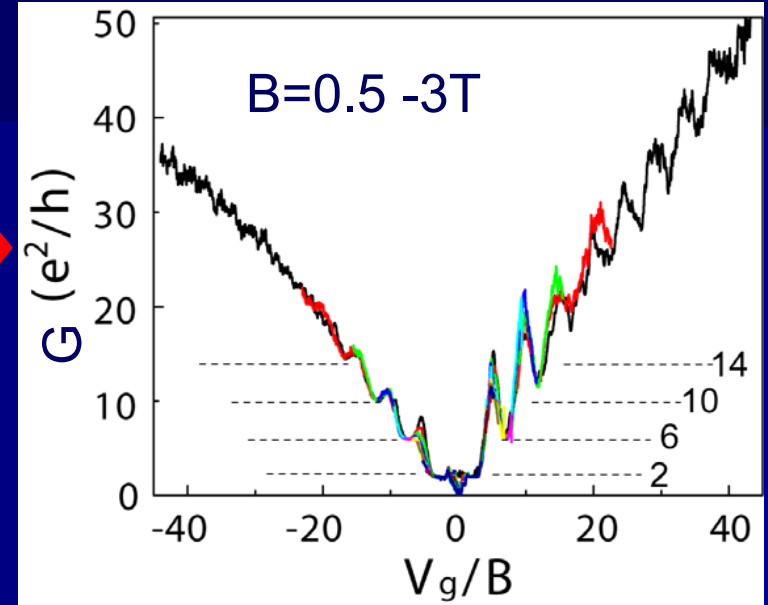
Suspended Graphene: QHE 2-terminal measurement



$$\frac{h}{2e^2}$$

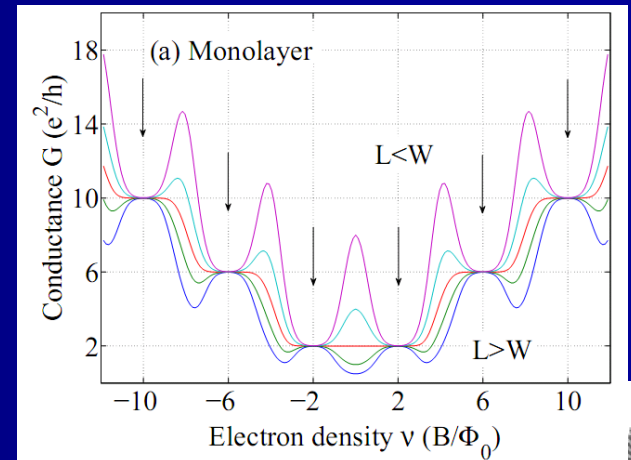
$$\frac{h}{3e^2}$$

$$\frac{h}{4e^2}$$

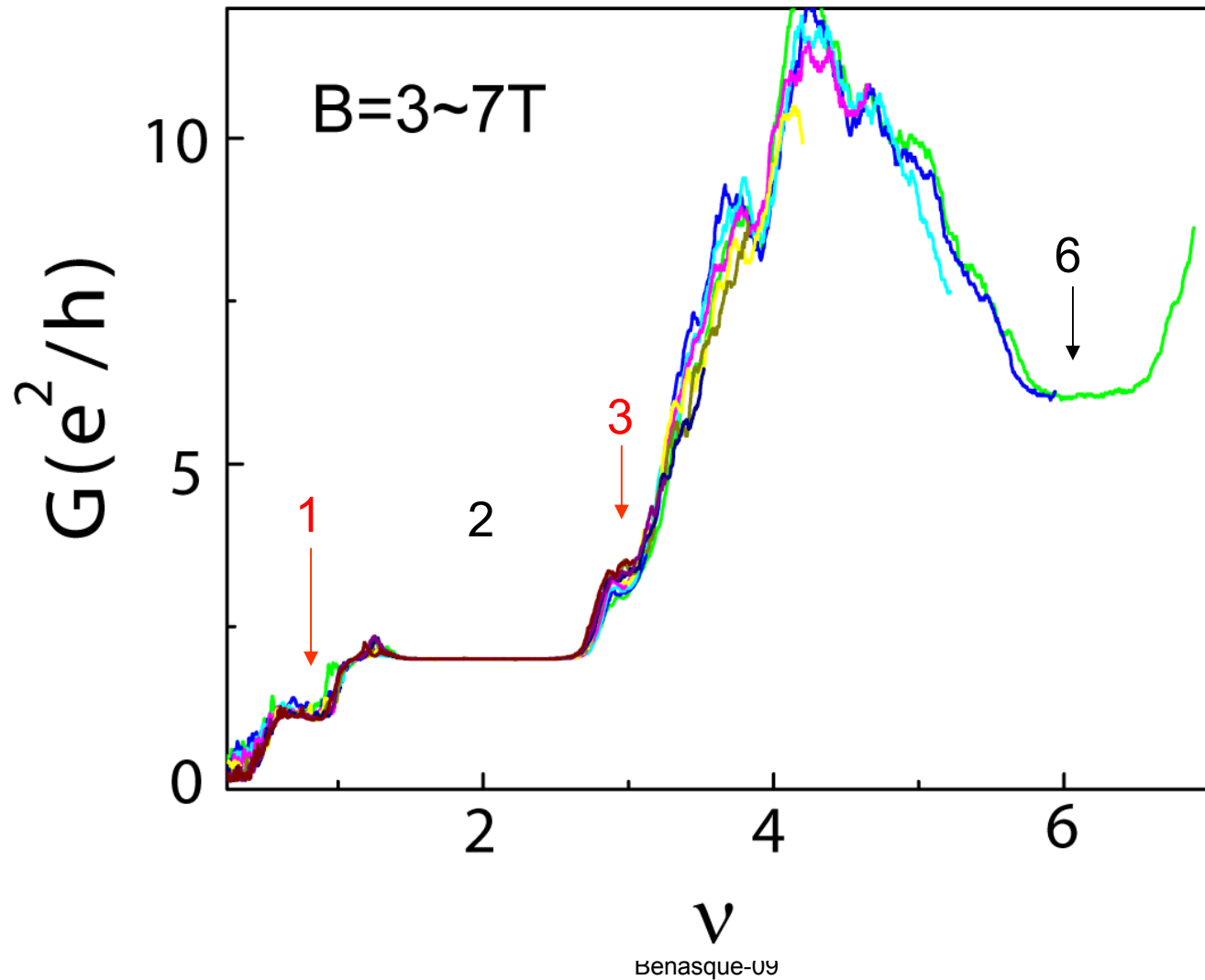


$L / W \sim 2-3$

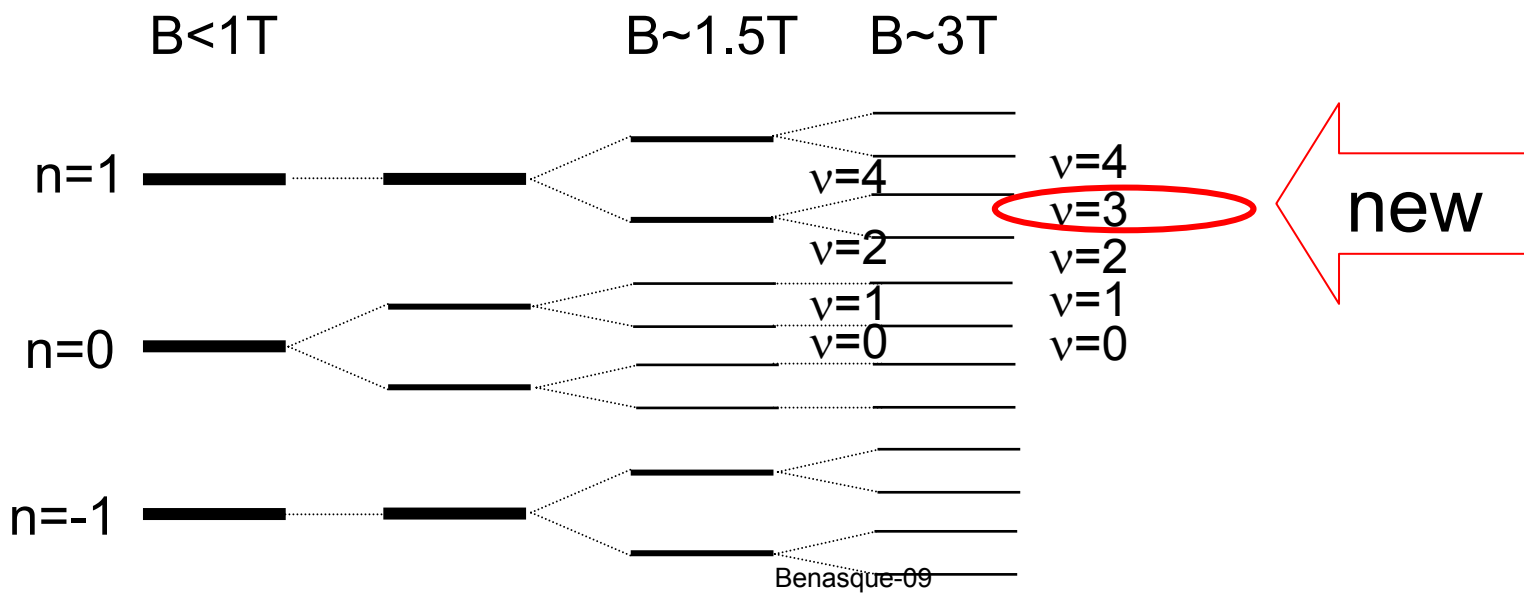
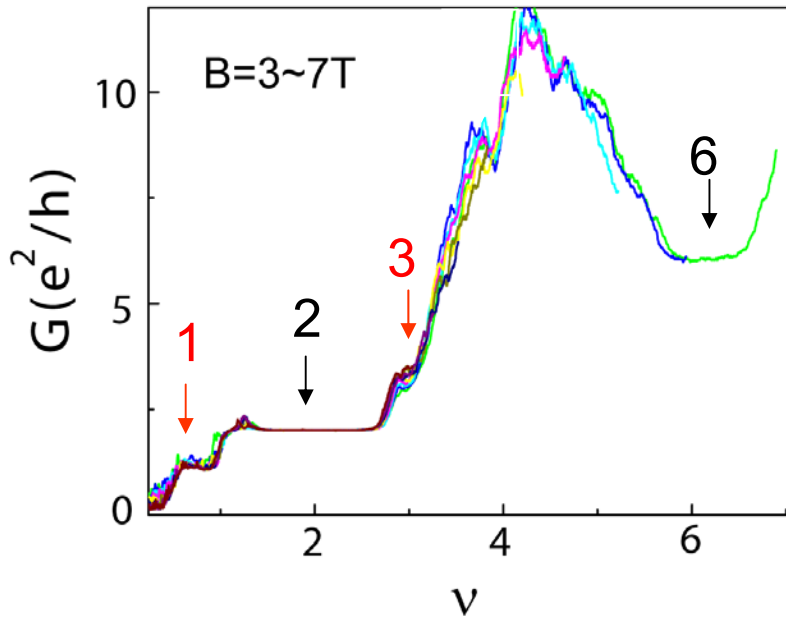
QHE in 2-terminal device:
Abanin & Levitov, PRB (2008)
 Benasque-09



Suspended Graphene: QHE 2-terminal measurement

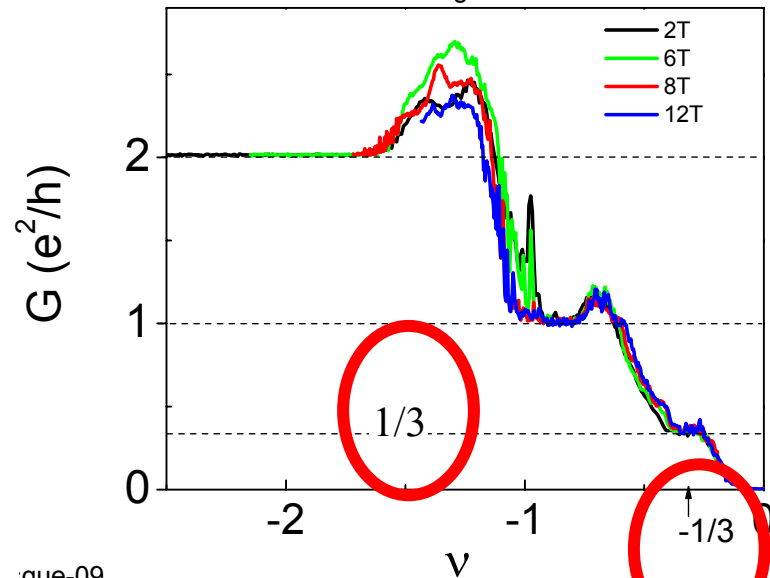
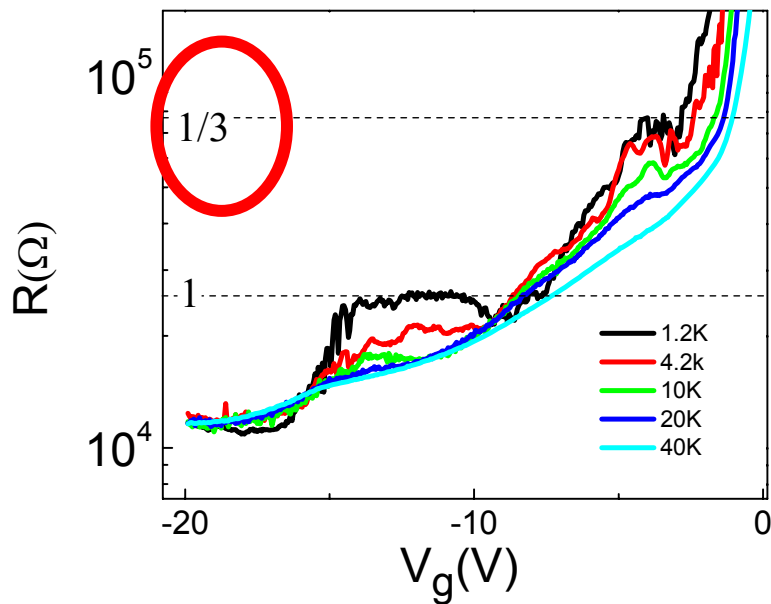
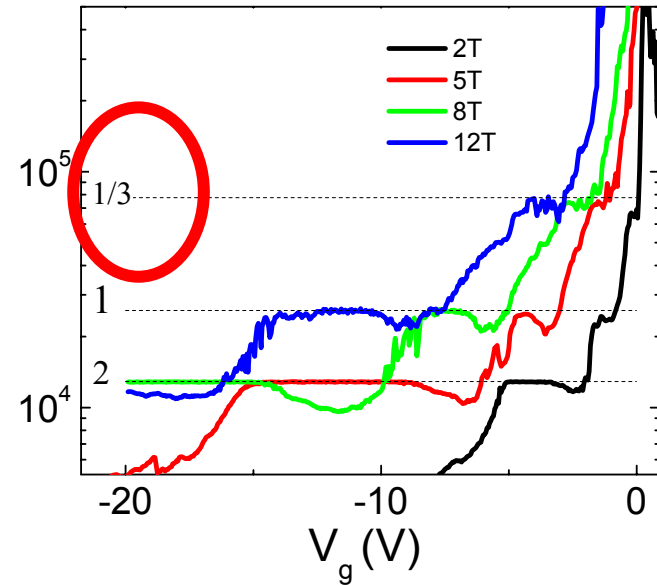
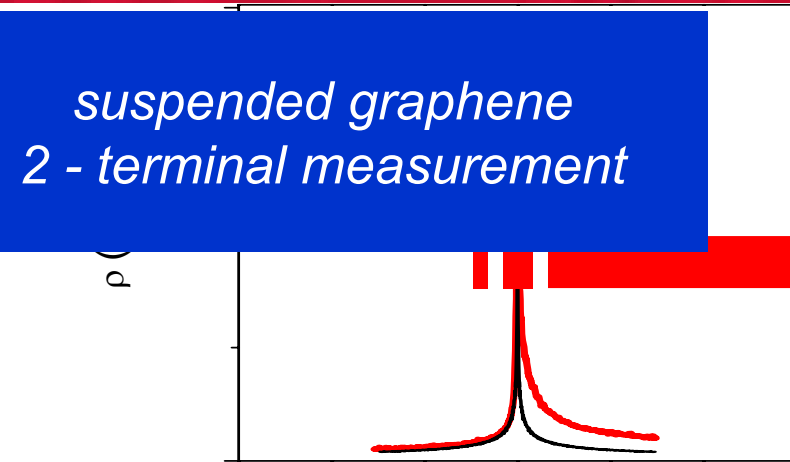


Suspended Graphene: QHE 2-terminal measurement



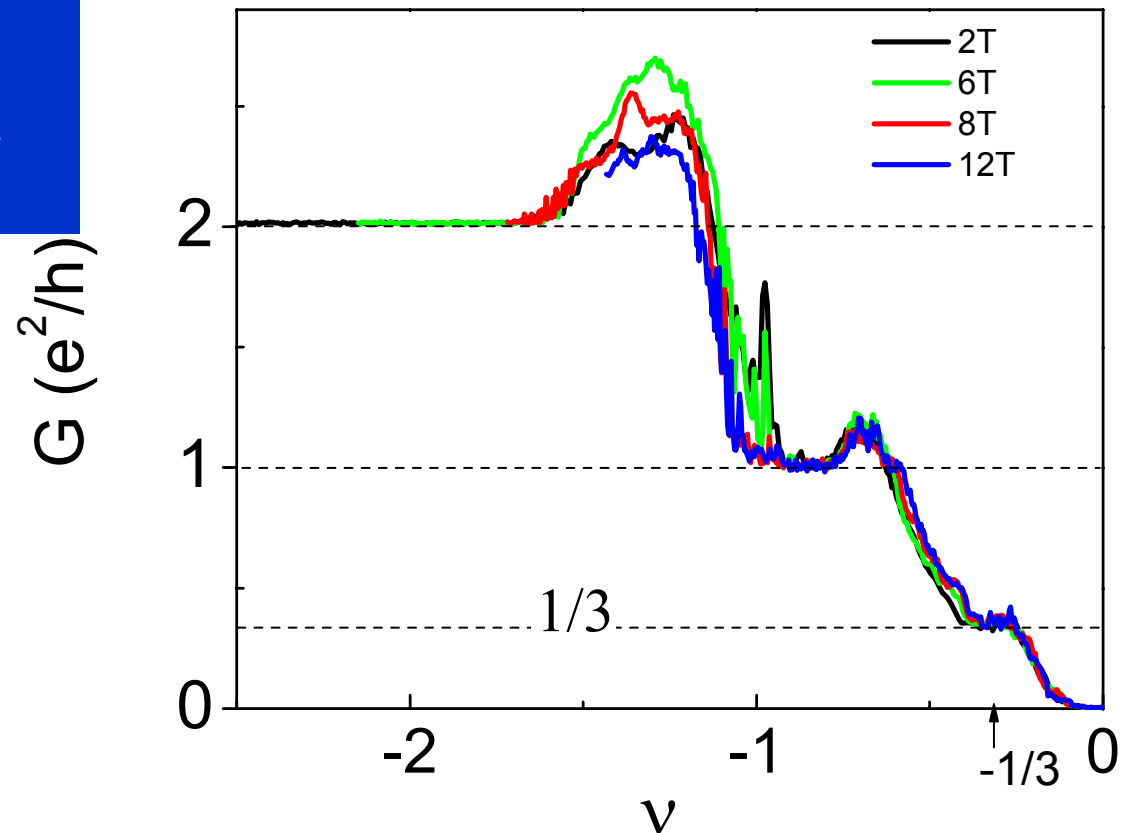
1'st observation of Fractional QHE in graphene

suspended graphene
2-terminal measurement



1'st observation of Fractional QHE in graphene

suspended graphene
2 - terminal measurement



FQHE not seen in 4-terminal measurements !!



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"New" Substrate - Graphite

G. Li, A. Luican, E. Y. A., PRL (2009)

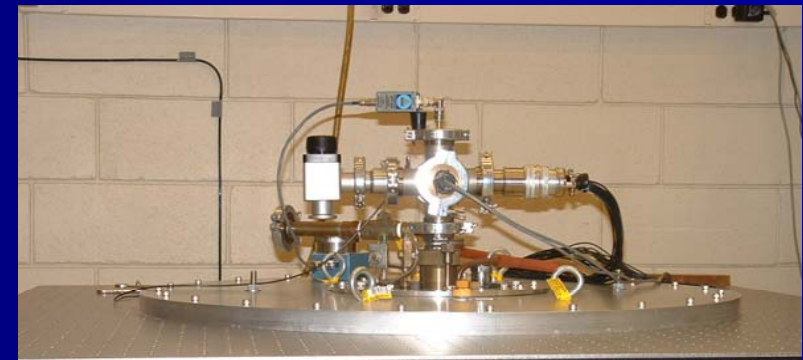
G. Li, E.Y. A - Nature Physics, 3, 623 (2007)

- Clean
- Lattice matched
- Conductor

■ STM –

- $T=4$ (2K)
- $B=13$ (15T)
- 10^{-10} - 10^{-3} m

graphene



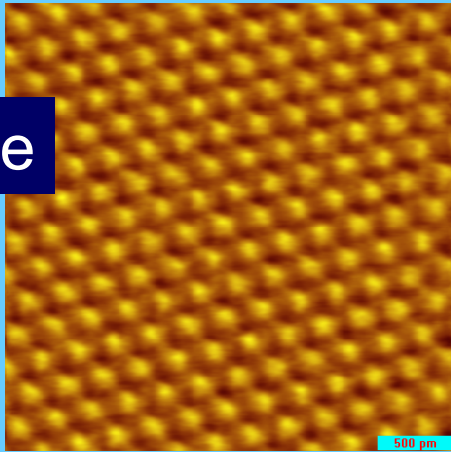
- Topography
- Spectroscopy $B=0$
- Spectroscopy $B>0$



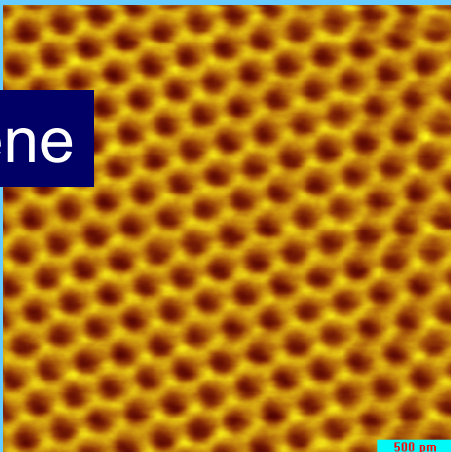
STM: Graphene on Graphite



topography

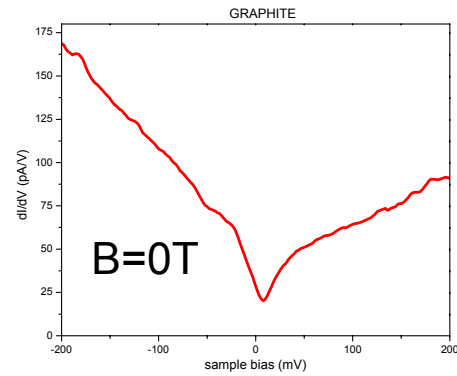


Graphite



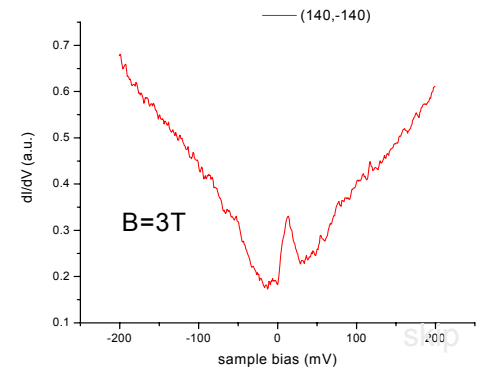
Graphene

B=0 spectroscopy

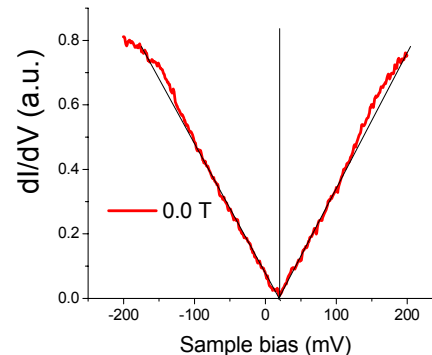


B=0T

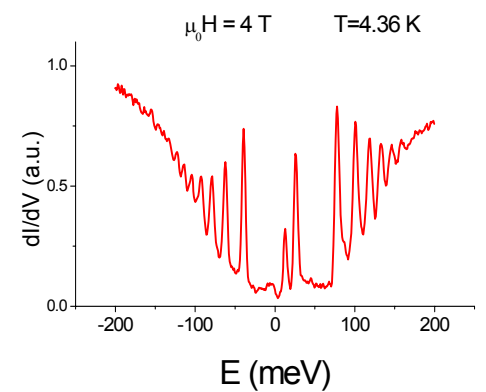
B>0 spectroscopy



B=3T



0.0 T



$\mu_0 H = 4 \text{ T}$ $T = 4.36 \text{ K}$

E (meV)

Finding graphene on graphite

➤ Search on graphite surface

Macroscopic defects – Terraces, ribbons

➤ Characterize: Landau level spectroscopy

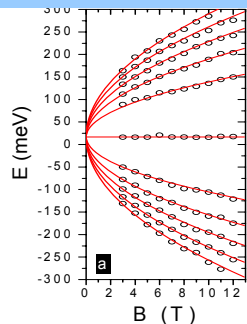
➤ Coupling strength and Landau levels



Decoupled graphene layer
Single massless sequence

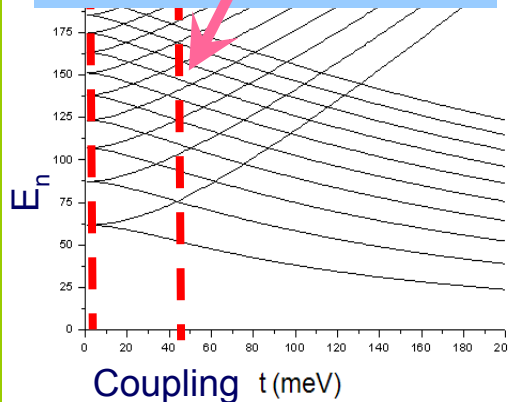
$$E_n \sim (nB)^{1/2}$$

G. Li, A. Luican, E.Y. A
PRL102, 176804 (2009)



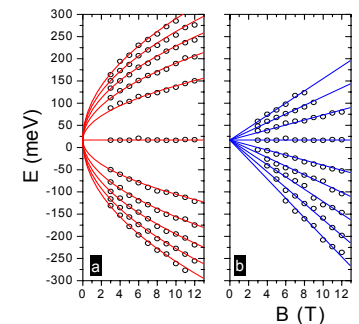
Weakly coupled layer
crossing levels

A. Luican, G. Li, E.Y. A
S. S. Com (2009)



3-10 coupled layers
2 LL sequences:
massless and massive

G. Li, E.Y. A
Nature Physics, 3, 623 (2007)

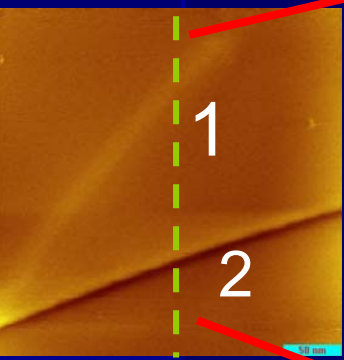


Interlayer coupling

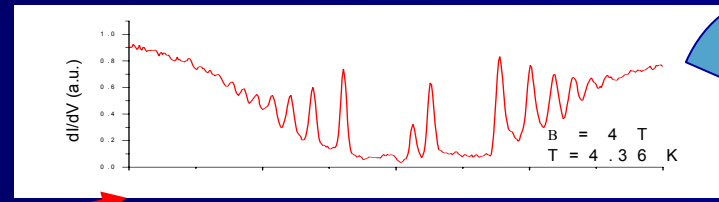


A. Luican, G. Li, E.Y. A., Solid State Com (2009)

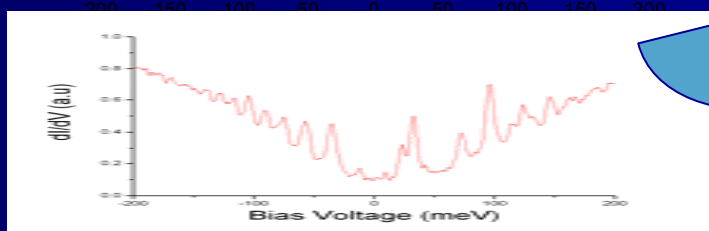
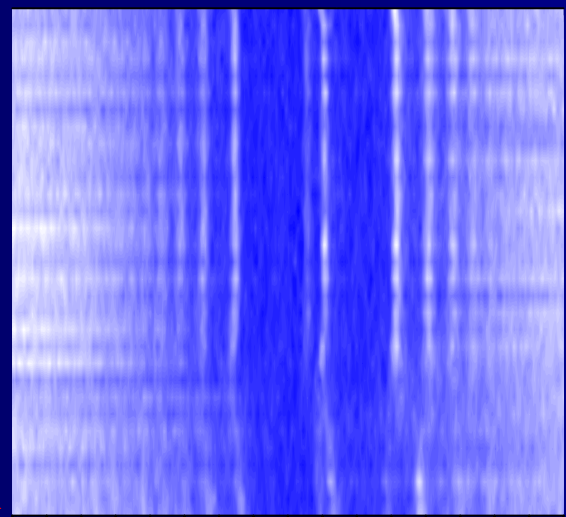
Pereira et al (PRB,2007)
LL vs interlayer coupling parameter



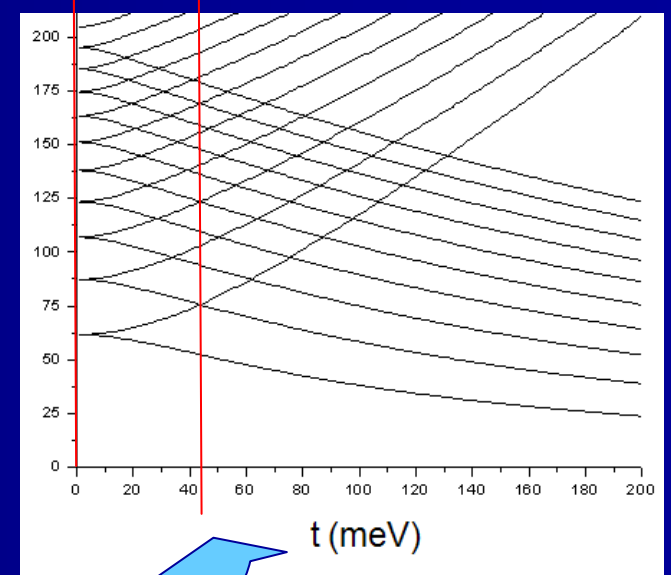
$t=0$



$t=45\text{meV}$



Region 1
No coupling – single layer graphene $t=0$



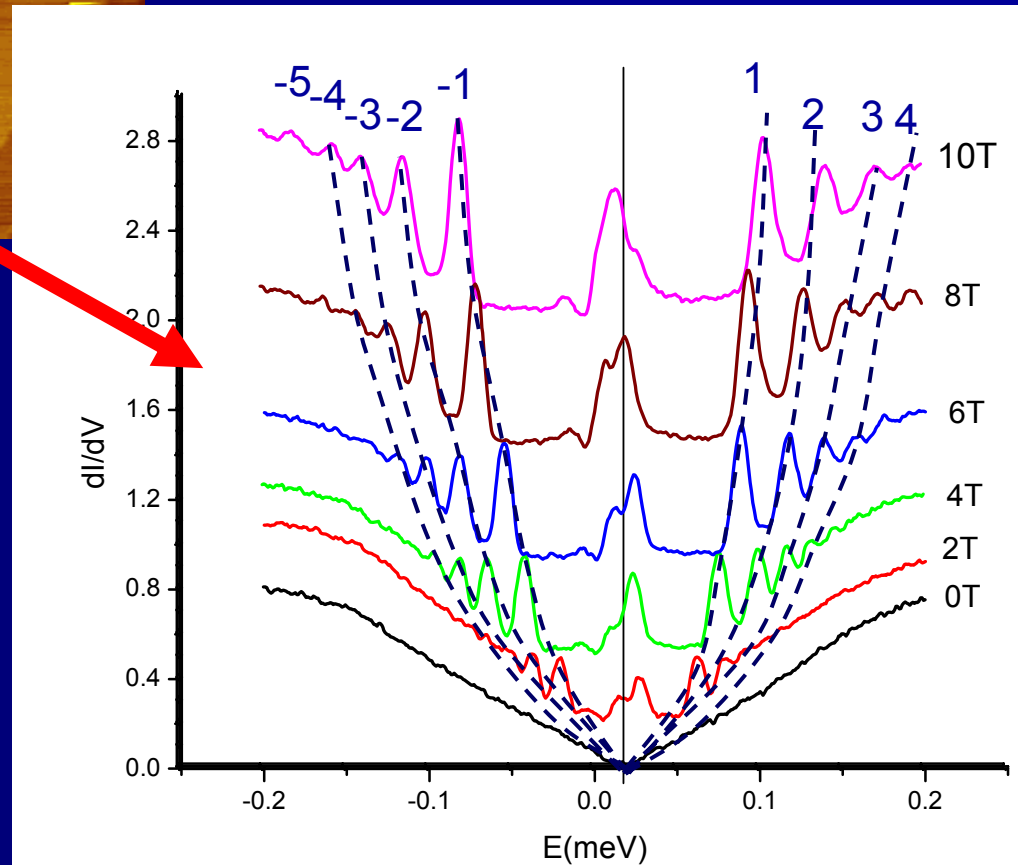
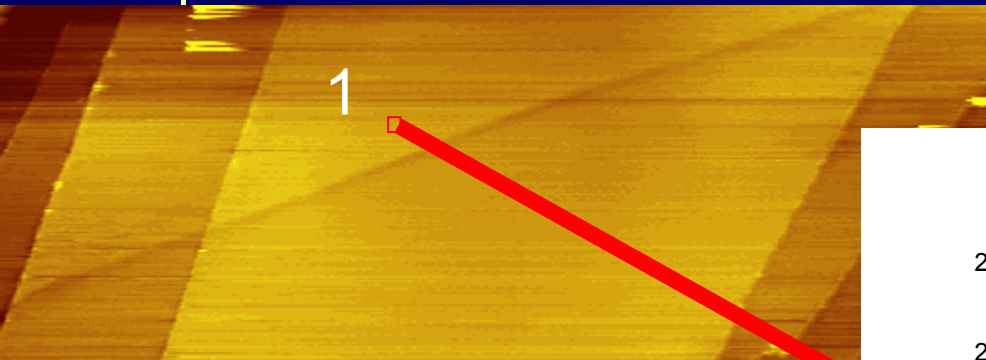
Region 2
Coupled layer $t=0.045\text{eV}$
Compare to $\gamma_1=0.3 \text{ eV}$



Landau level spectroscopy

G. Li, A. Luican, E. Y. A., *PRL* (2009)

G. Li, E.Y. A. - *Nature Physics*, 3, 623 (2007)

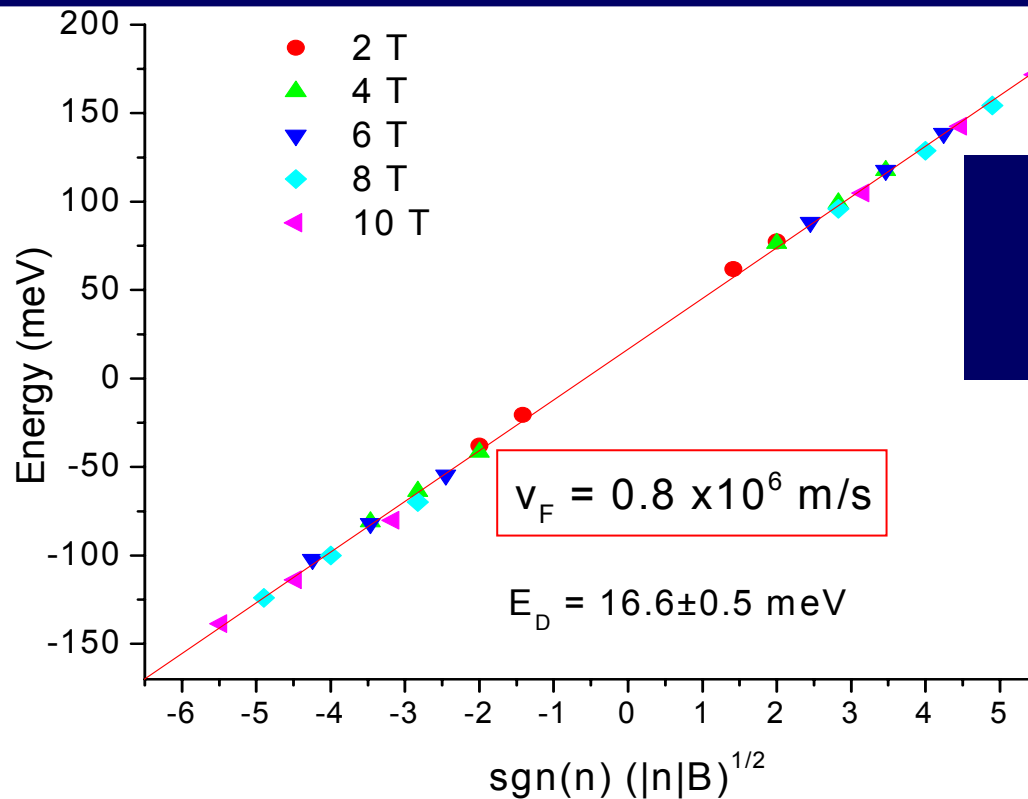


Massless Dirac Fermions

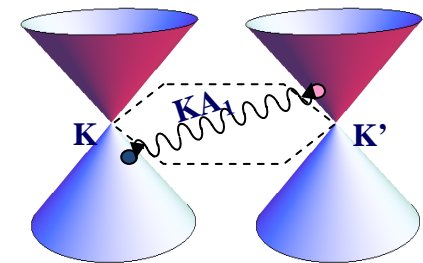
G. Li, A. Luican, E. Y. A., PRL (2009)

G. Li, E.Y. A. - Nature Physics, 3, 623 (2007)

$$E_j = \pm v_F \sqrt{2e\hbar |n| B}, \quad n = 0, \pm 1, \dots$$



Slow down due to Electron-phonon interactions



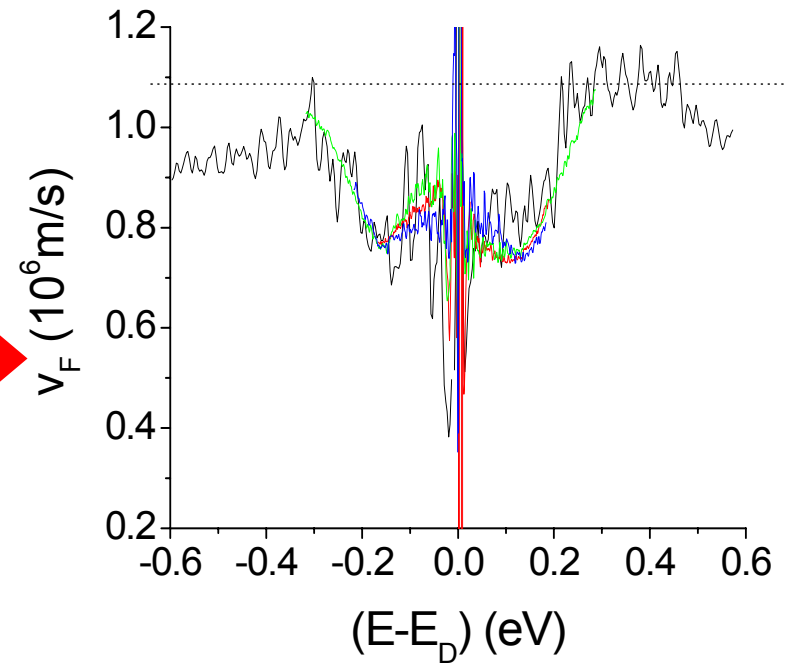
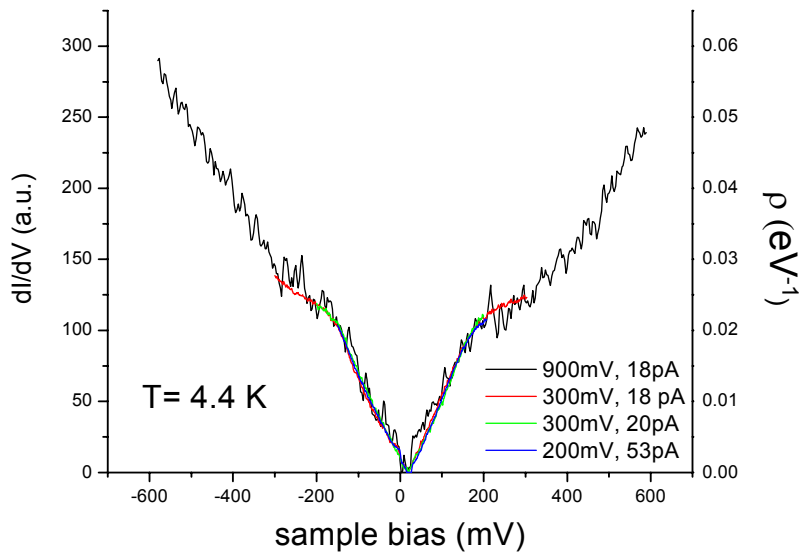
Expect: $V_F = 1 \times 10^6 \text{ m/s}$

skip



Velocity normalization

- Zero field density of states

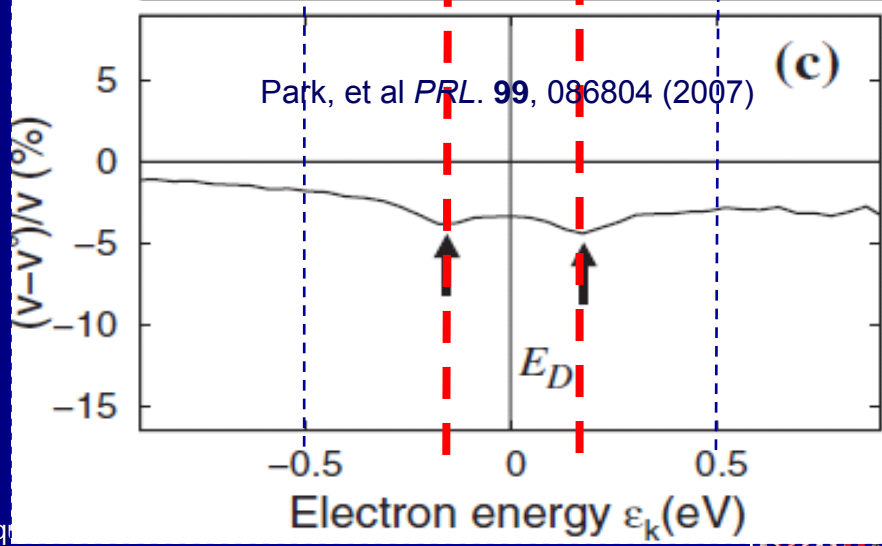
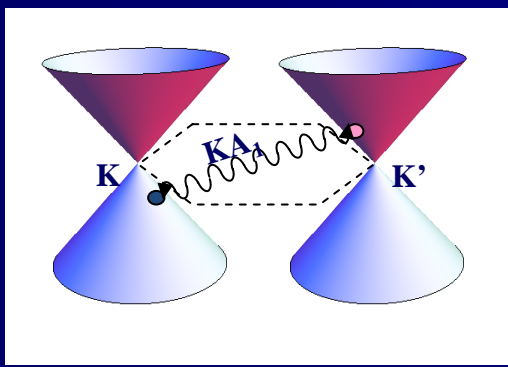
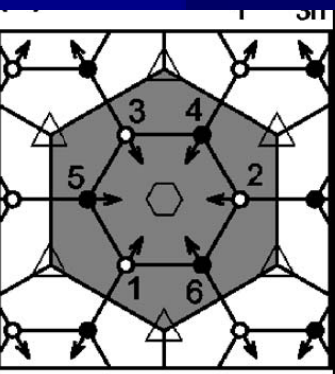
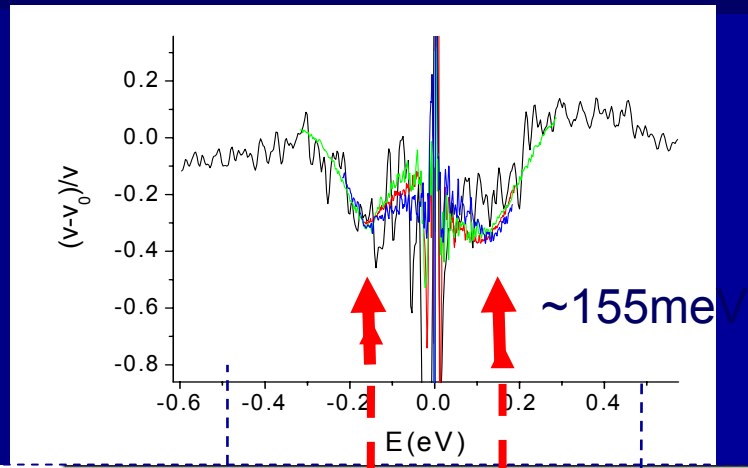
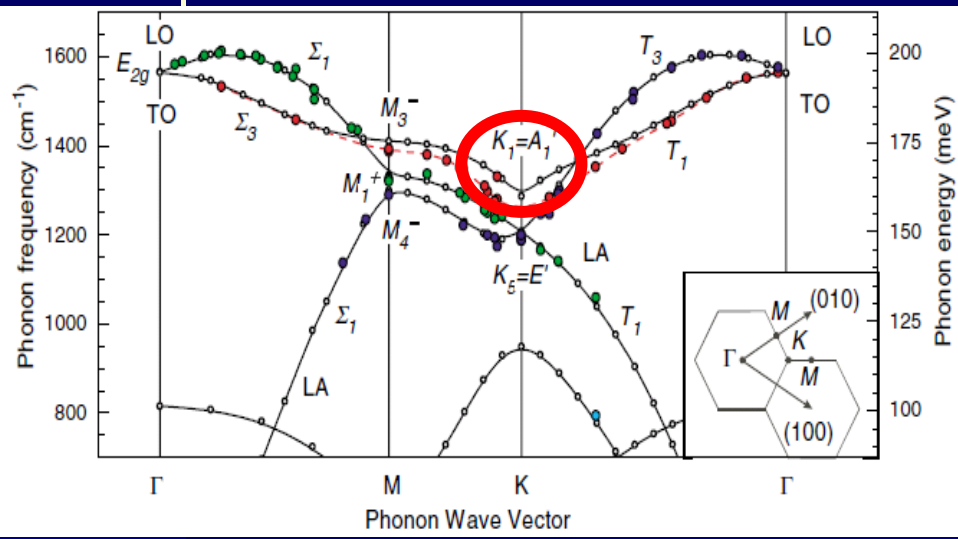


$$v_F \equiv \frac{dE}{\hbar dk} = \frac{2}{\hbar} \sqrt{\frac{A_c}{\pi}} \left(\int_0^E \rho(E') dE' \right)^{1/2} / \rho(E)$$



$$A_c = 3\sqrt{3}a^2/2$$

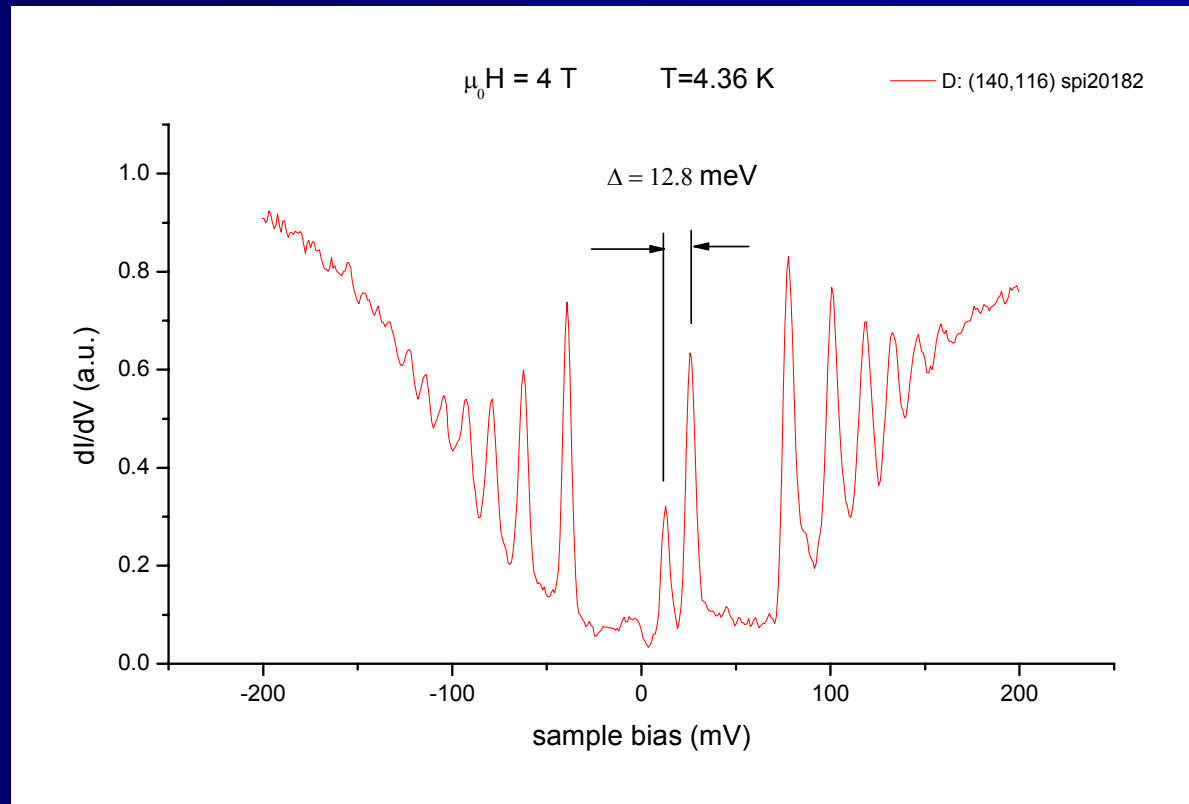
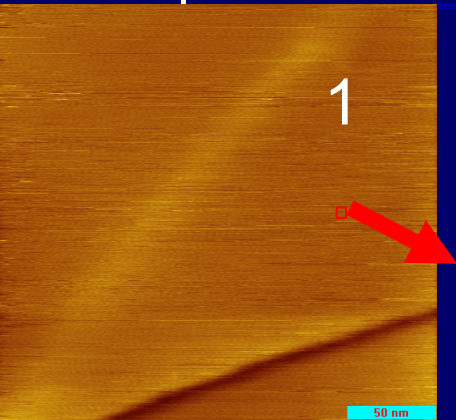
Electron phonon coupling



High resolution STS – 4T

G. Li, A. Luican, E. Y. A., PRL (2009)

G. Li, E.Y. A - Nature Physics, 3, 623 (2007)



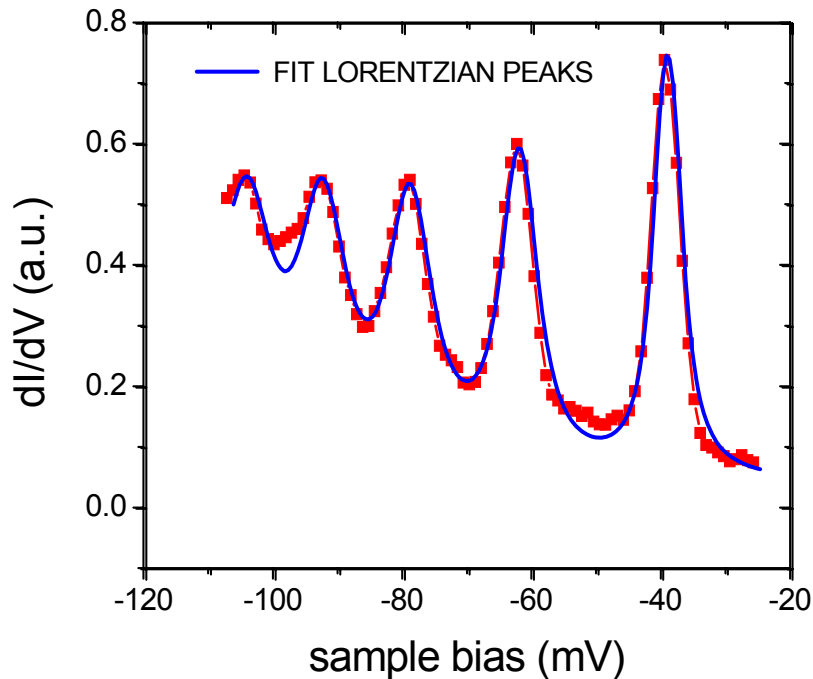
16 resolved Landau Levels



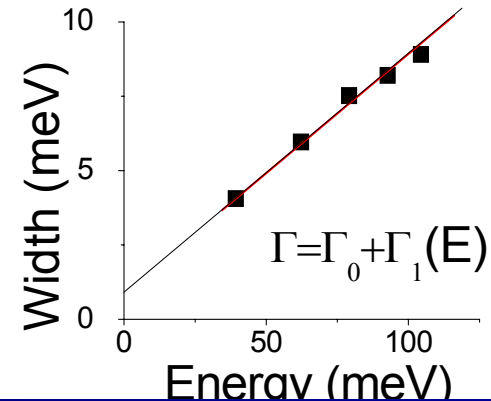
Quasiparticle lifetime

G. Li, A. Luican, E. Y. A., PRL (2009)

G. Li, E.Y. A - Nature Physics, 3, 623 (2007)



Line-width $\sim E$



$$\tau_0 = 0.7 \text{ ps} \Rightarrow l_{mfp} \sim v_F \tau_0 = 700 \text{ nm}$$

$$\tau \propto E^{-1} \approx 9 \text{ ps} / \text{meV}$$

Inelastic e-e interactions

$$\tau \sim E^{-1} \sim 18 \text{ ps} / \text{meV}$$

$$\mu = \frac{ev_F l_{mfp}}{E_F} = 220,000 \text{ cm}^2 / \text{V} \cdot \text{sec}$$

$$n = 3 \times 10^{10} \text{ cm}^{-2}$$

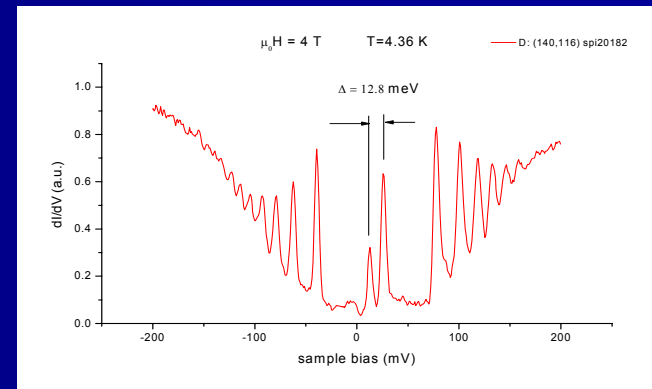
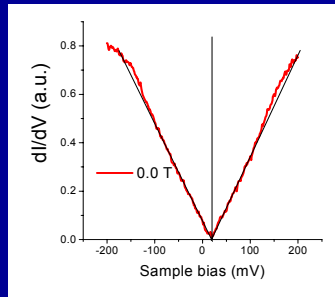
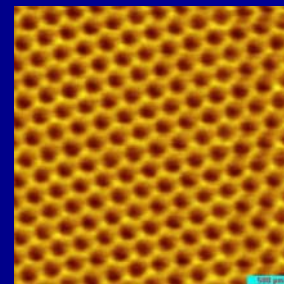
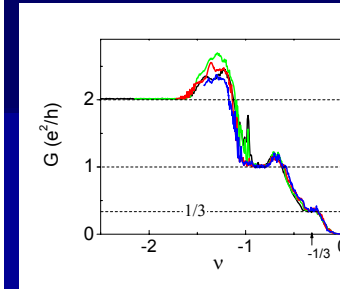
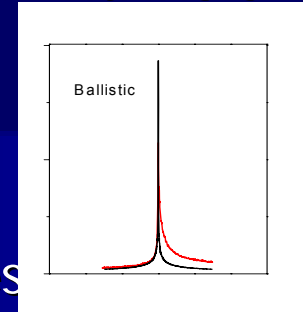
Gonzalez et al 1993
Castro Neto al PRB 2006
Fritz et al arXiv:0802.4289



Summary

Graphene without insulating substrate: Intrinsic properties

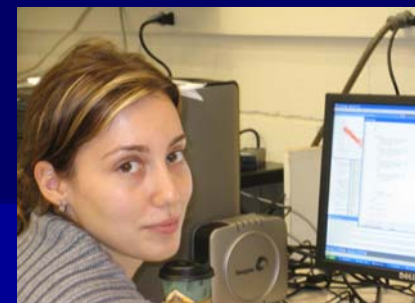
- Transport – 2-terminal
 - Ballistic transport on micron length scales
 - IQHE $\nu=1,2,3,4$
 - FQHE $\nu=1/3$
 - $\nu=0$ insulator
- STM
 - Honeycomb structure
 - Spectroscopy
 - Linear Density of states
 - Well defined Dirac point
 - Direct observation of Landau levels
 - Fermi velocity
 - e-phonon interactions
 - e-e interactions





STM -

Guohong Li
Adina Luican



Transport - Xu Du

Ivan Skachko
Anthony Barker
Fabian Duerr



Eva Y. Andrei



Collaborators:

J. Kong, A. Reina, A. Geim, R. Nahir, K. Novoselov

J. Lopes dos Santos, A. H. Castro Neto

