

# An Inhomogeneous World View

Modern Cosmology – Benasque August 2010

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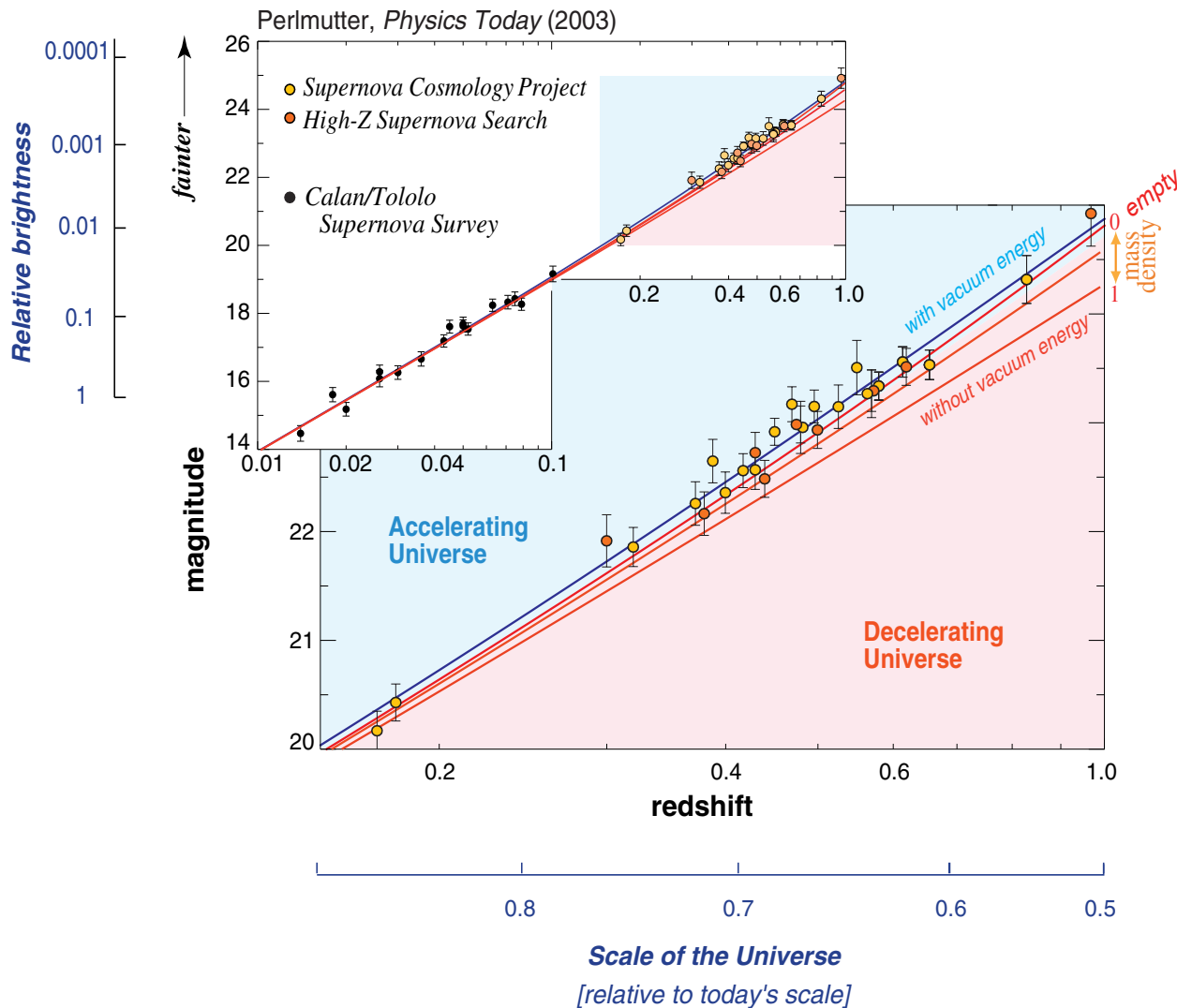
**Juan Garcia-Bellido**

**Benjamin Sinclair**

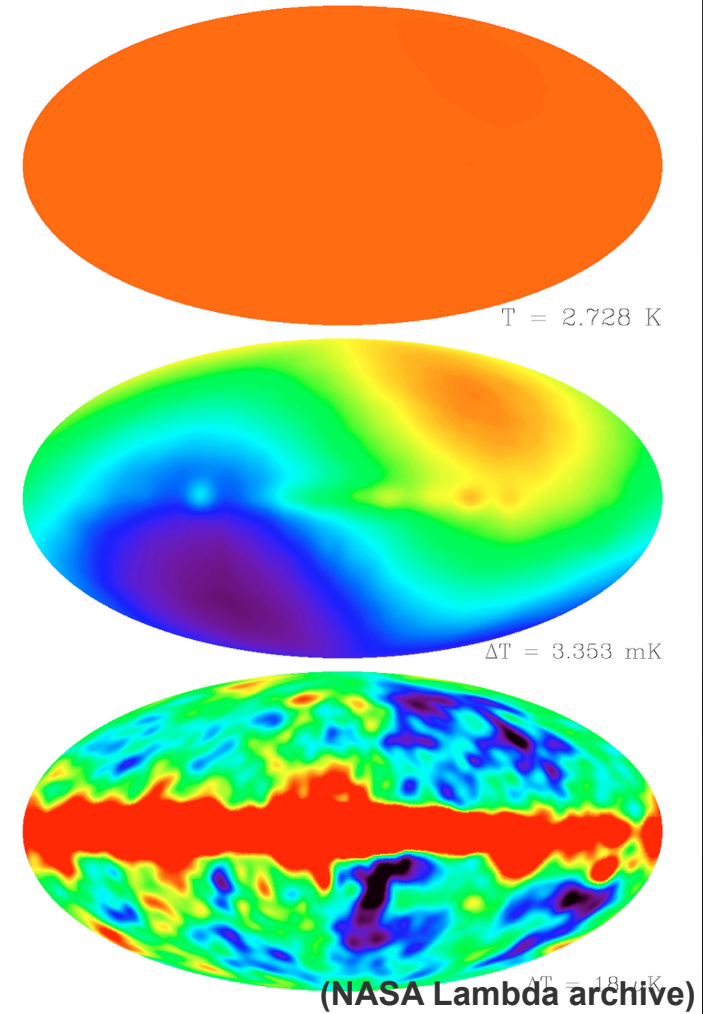
**Julian Vicente**

# We live in an Accelerating & Isotropic Universe

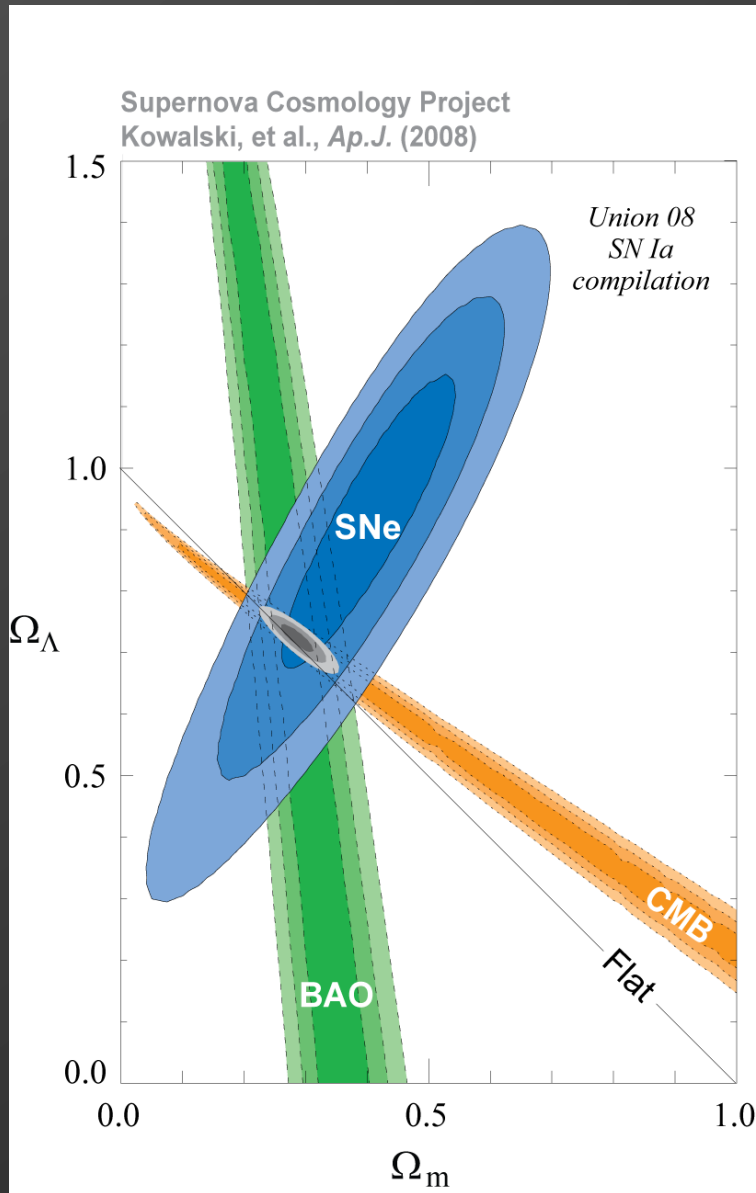
Type Ia Supernovae



DMR 53 GHz Maps



# The Standard $\Lambda$ CDM model



The numbers

are impressive,  
but can we trust  
the model ???

$t_0 = 13.73 \pm 0.12$  Gyr Age of Universe

$H_0 = 70.1 \pm 1.3$  km s<sup>-1</sup> Mpc<sup>-1</sup> Expansion rate

$\Omega_b = 0.0462 \pm 0.0015$  Baryon density

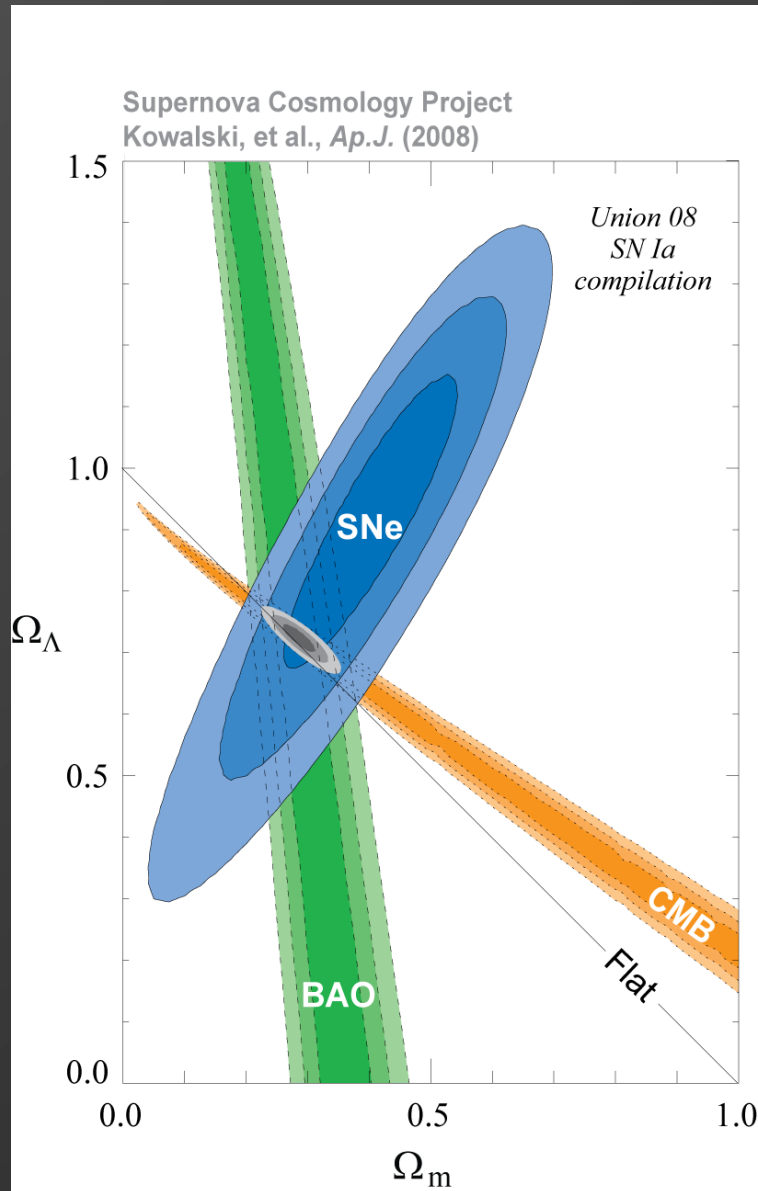
$\Omega_c = 0.233 \pm 0.013$  Dark Matter density

$\Omega_\Lambda = 0.721 \pm 0.015$  Dark energy density

$n_s = 0.960 \pm 0.014$  Scalar spectral index

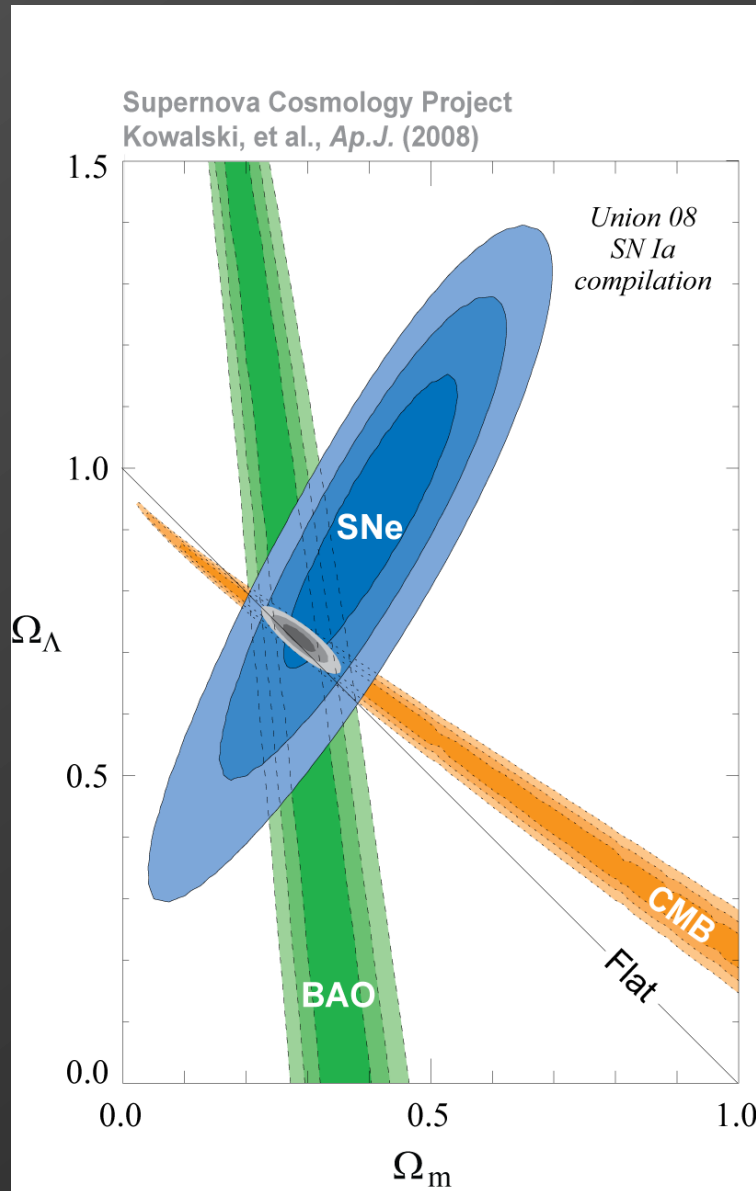
$\Delta_R^2 = 2.46 \pm 0.09 \cdot 10^{-9}$  Fluctuation amplitude

# The Pillars of the $\Lambda$ CDM model



- General Relativity
- The FRW model of a homogeneous universe
- Dark Matter & fields with strange equations of state

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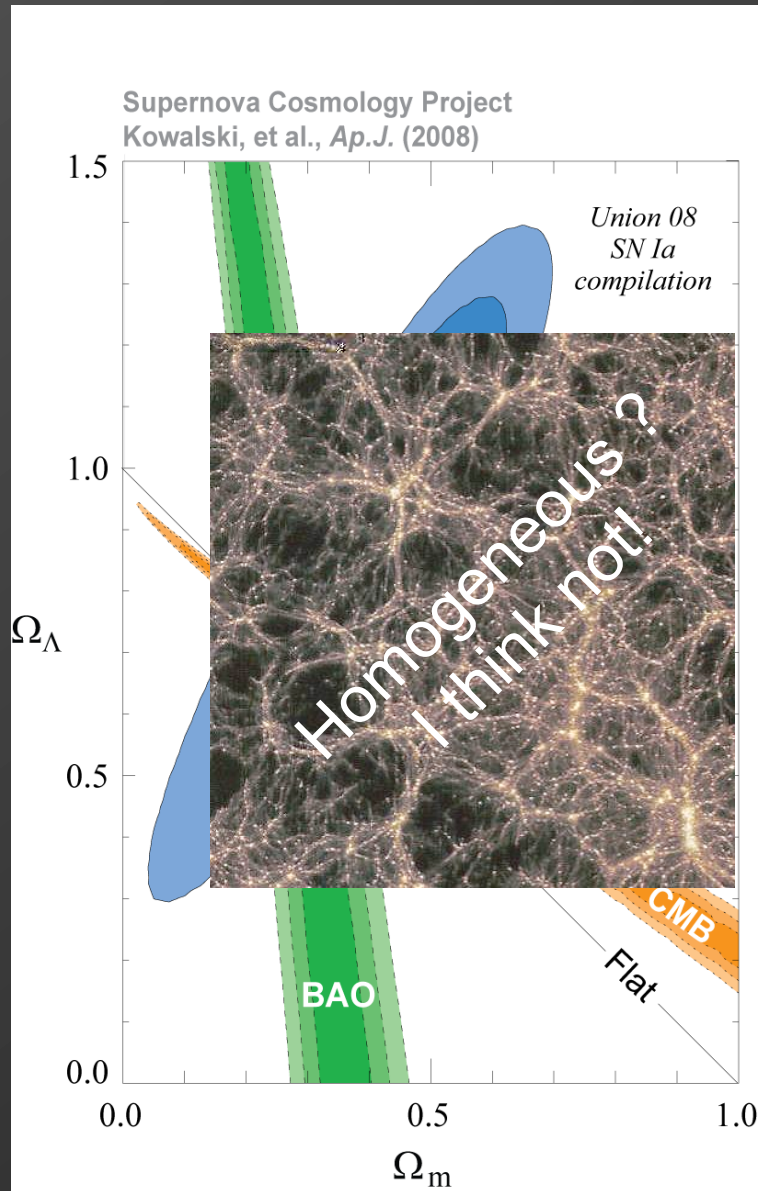


## Modify Gravity

- General ~~Relativity~~
- The FRW model of a homogeneous universe
- Dark Matter & ~~fields~~ with strange equations of state

# The Pillars of the $\Lambda$ CDM model

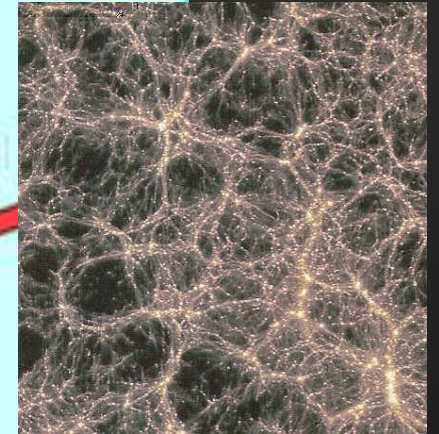
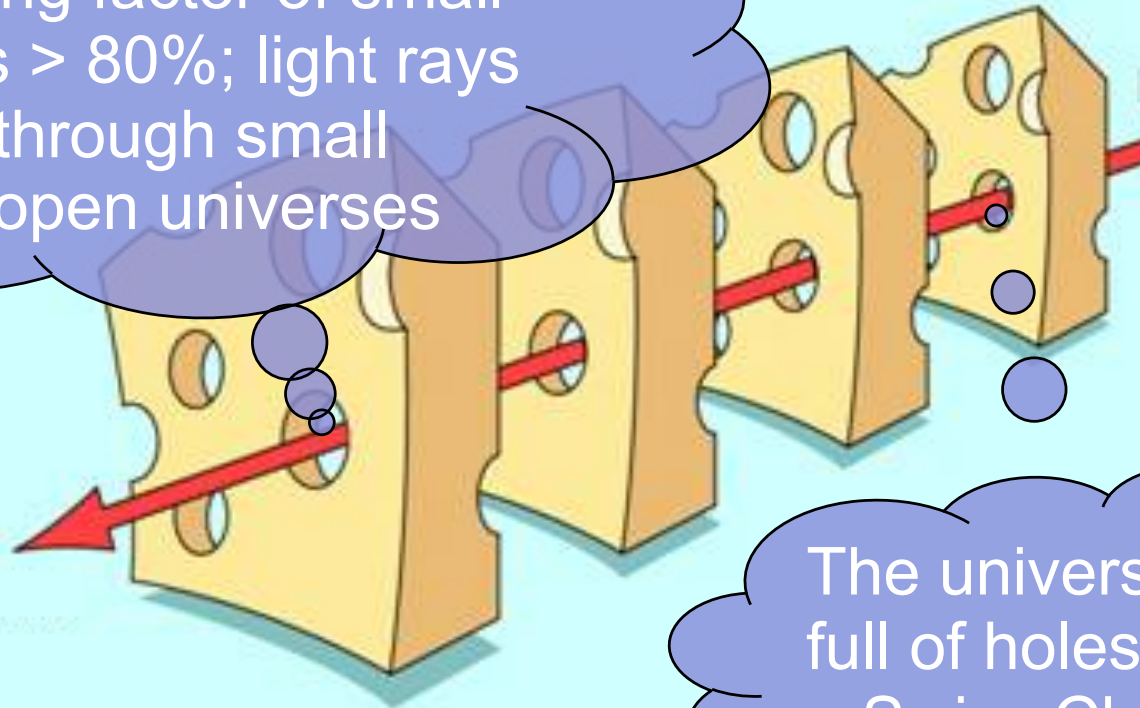
## Alternative Scenario



- General Relativity
- **Modify the Geometry**
- The FRW model of a homogeneous universe
- Dark Matter & fields with strange equations of state

# Alternative scenario I Its Swiss Cheese

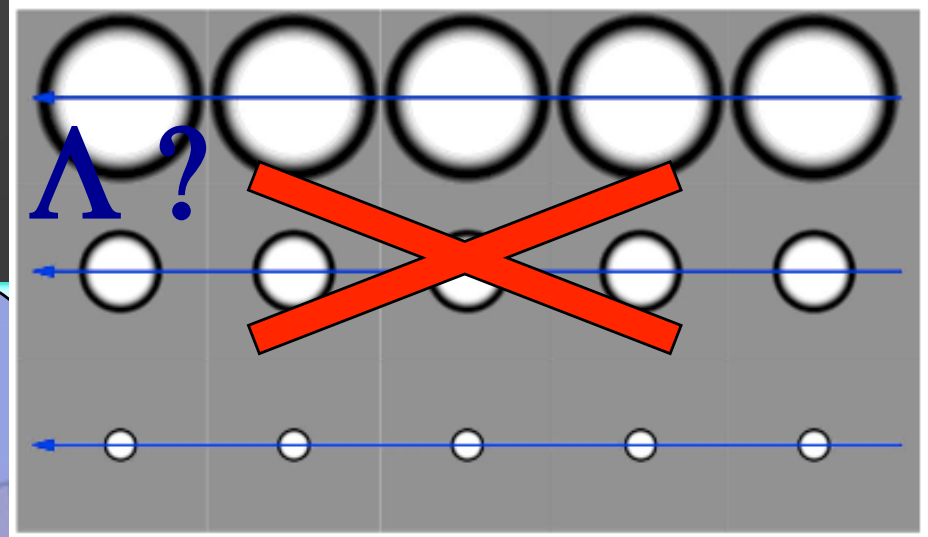
The filling factor of small voids is  $> 80\%$ ; light rays will go through small locally open universes



The universe is full of holes, like a Swiss Cheese

# Alternative scenario I Its Swiss

Since voids becomes bigger and bigger we get apparent acceleration



But randomising the positions of the voids removes the effect.

Swiss Cheese Stinks!

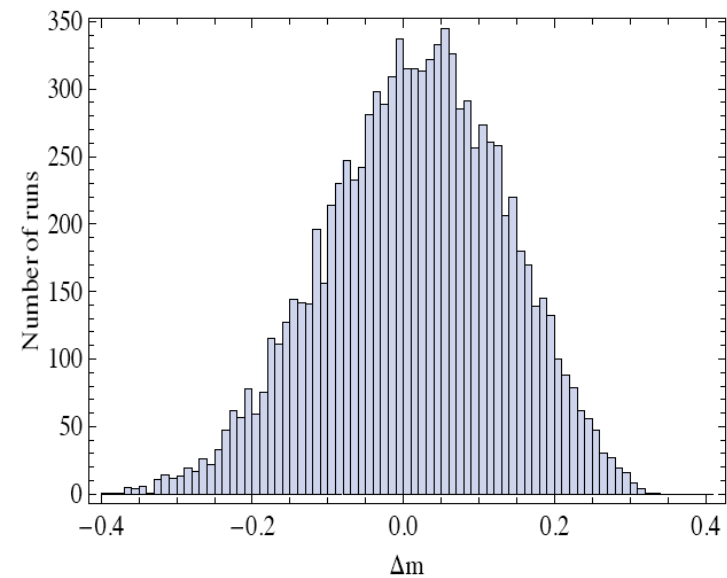


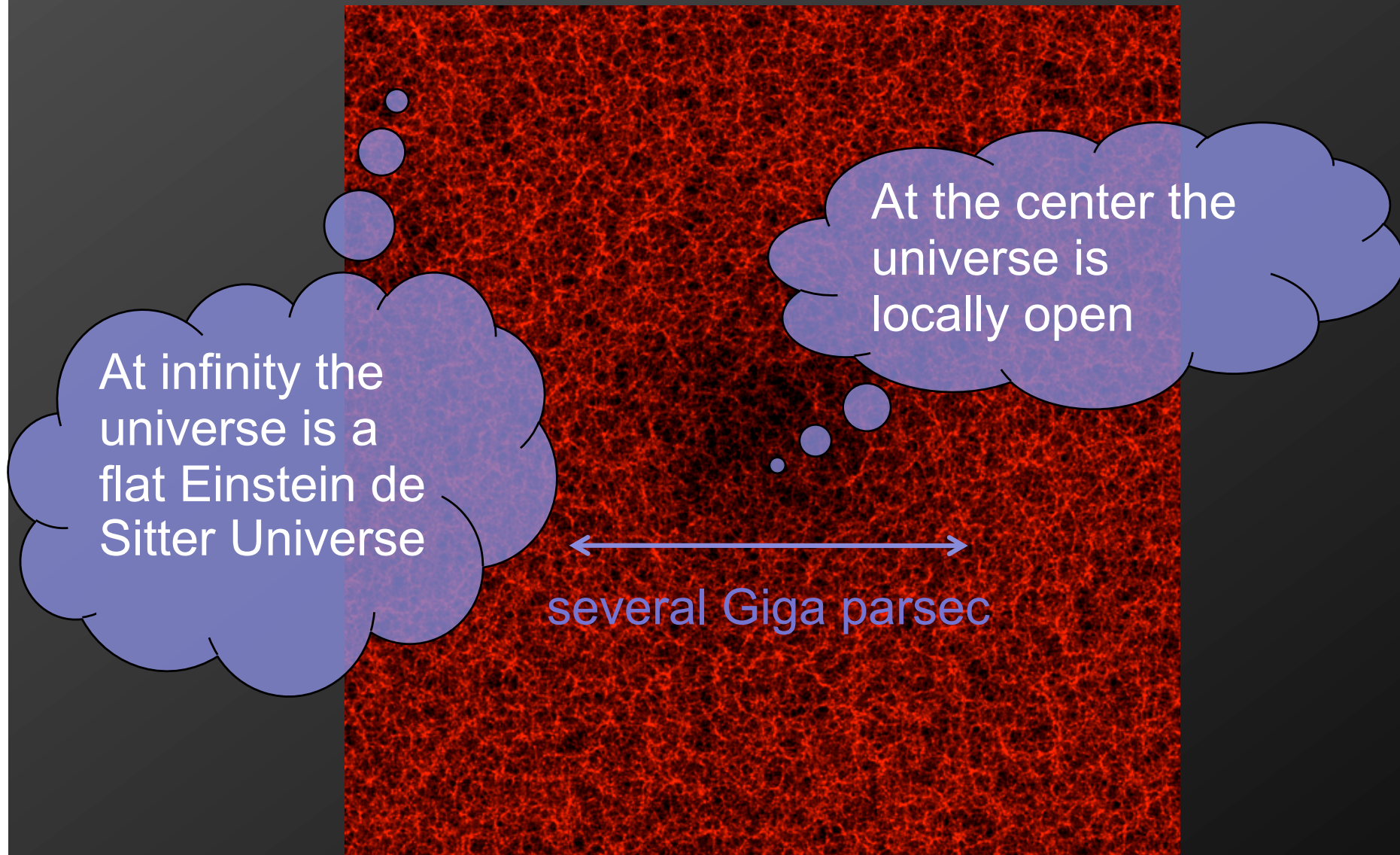
FIG. 2: A histogram of the distance modulus shifts  $\Delta m$  found for 10,000 realizations of our randomized void scenario

Vanderveld et al 2008



# Alternative scenario II

## We are living inside a giant void



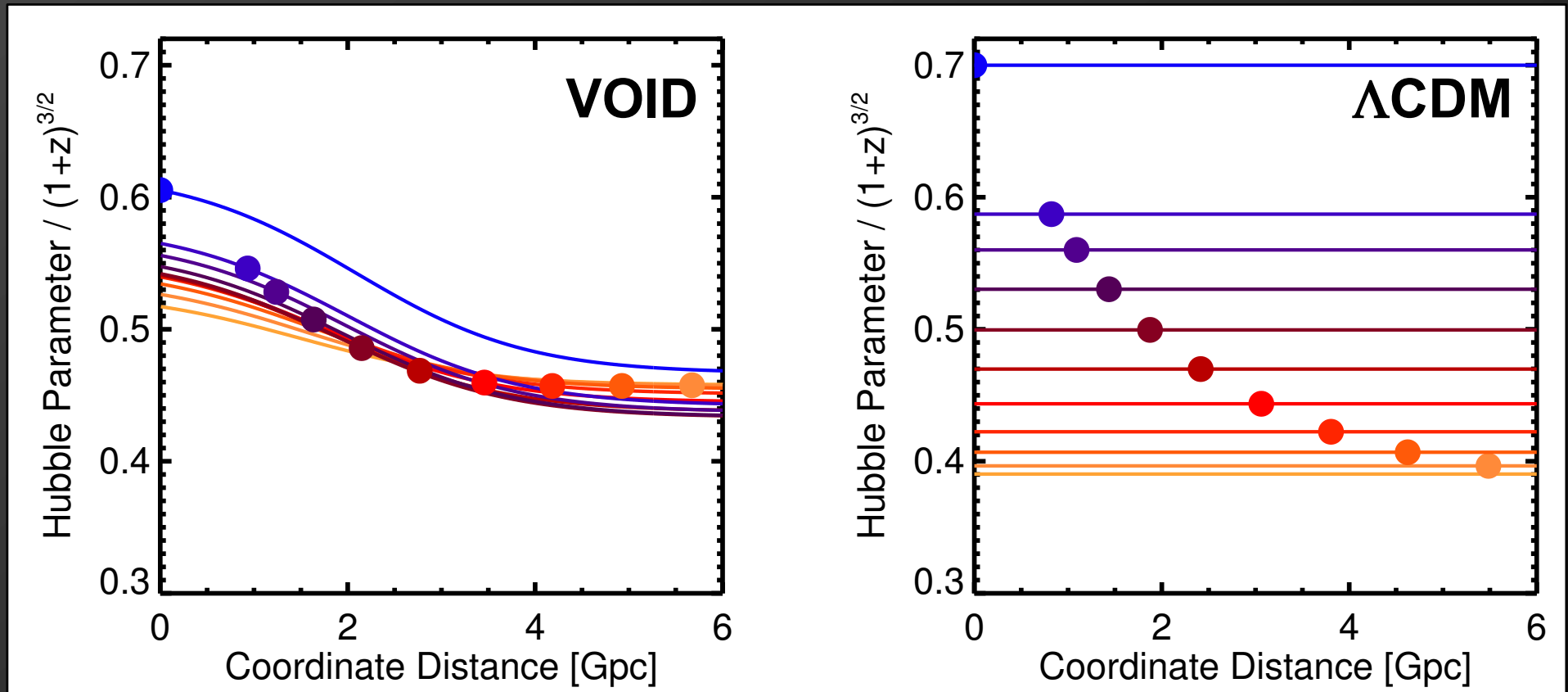
# How does a Central Void Help Explain Dark Energy?

- We observe in the redshift cone: Acceleration can be due to both spatial and temporal changes in the expansion rate. Maybe the Universe is tricking us?

$$\frac{D}{Dt} \approx \frac{\partial}{\partial t} - c \frac{\partial}{\partial r}$$

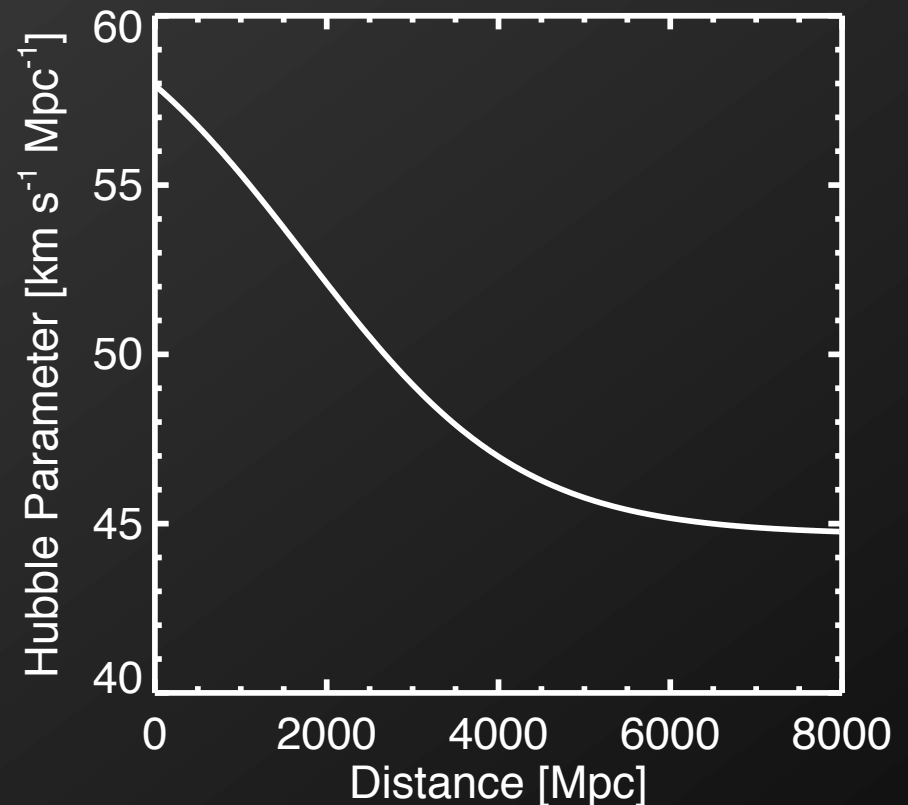
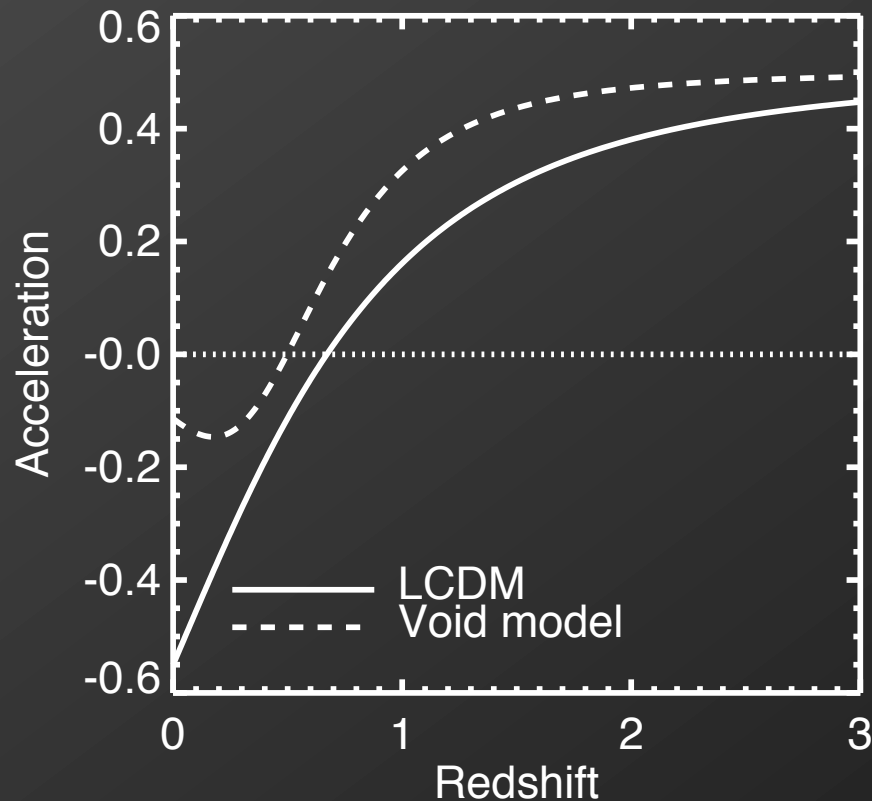
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# Void Construction 101

## The Lemaître-Tolman-Bondi Model

- Describes a space-time, which has spherical symmetry in the spatial dimensions, but with time and radial dependence

$$ds^2 = -dt^2 + X^2(r, t) dr^2 + A^2(r, t) d\Omega^2$$

- Defining an effective matter density and the Hubble rate as

$$H(r, t) = \frac{\dot{A}(r, t)}{A(r, t)},$$

$$F(r) = H_0^2(r) \Omega_M(r) A_0^3(r),$$

$$k(r) = H_0^2(r) \left( \Omega_M(r) - 1 \right) A_0^2(r),$$

# Void Construction 101

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$$ds^2 = -dt^2 + X^2(r, t) dr^2 + A^2(r, t) d\Omega^2$$

- A “local analogy” to the Friedman equation can be derived

$$H^2(r, t) = H_0^2(r) \left[ \Omega_M(r) \left( \frac{A_0(r)}{A(r, t)} \right)^3 + (1 - \Omega_M(r)) \left( \frac{A_0(r)}{A(r, t)} \right)^2 \right]$$

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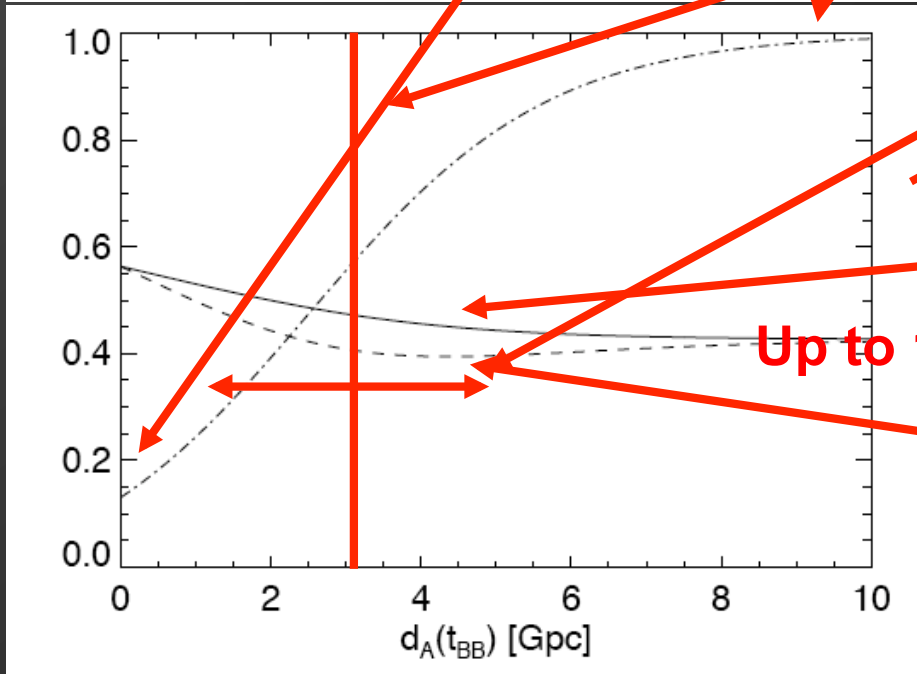
where  $H_0(r)$ , the expansion rate at  $t_0$ , and  $\Omega_M(r)$ , the total matter density inside  $r$  at  $t_0$ , uniquely describes our model.

Functional form fixed due to gauge choice

# A simple model for the void

$$\Omega_M(r) = \Omega_{\text{out}} + (\Omega_{\text{in}} - \Omega_{\text{out}}) \left( \frac{1 - \tanh[(r - r_0)/2\Delta r]}{1 + \tanh[r_0/2\Delta r]} \right)$$

$$H_0(r) = H_0 \left[ \frac{1}{\Omega_K(r)} - \frac{\Omega_M(r)}{\sqrt{\Omega_K^3(r)}} \sinh^{-1} \sqrt{\frac{\Omega_K(r)}{\Omega_M(r)}} \right]$$



Void Size

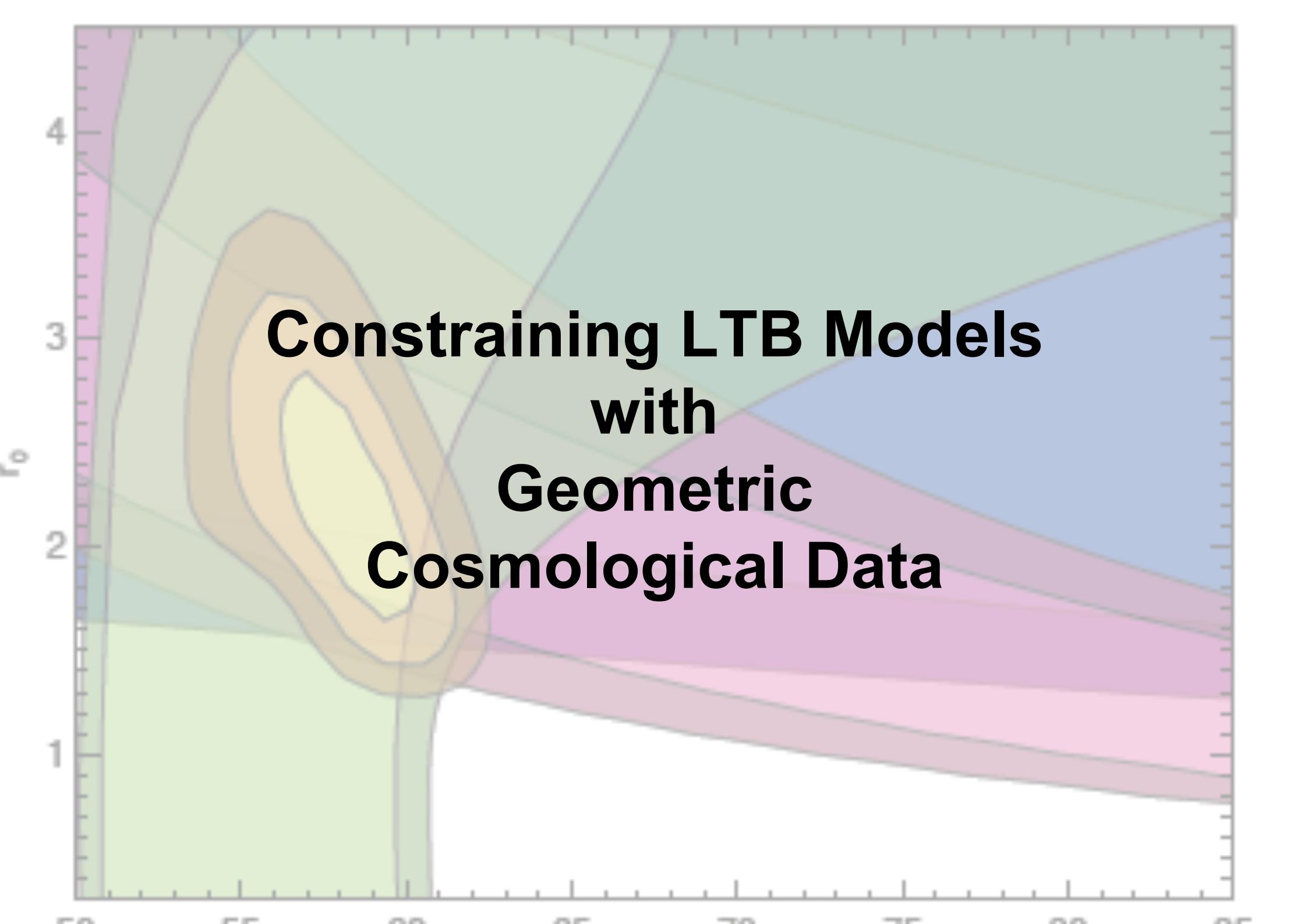
Transition Width

Up to 10 % difference!

$$H_T(r, t) \equiv \frac{\dot{A}(r, t)}{A(r, t)},$$

$$H_L(r, t) \equiv \frac{\dot{A}'(r, t)}{A'(r, t)}.$$





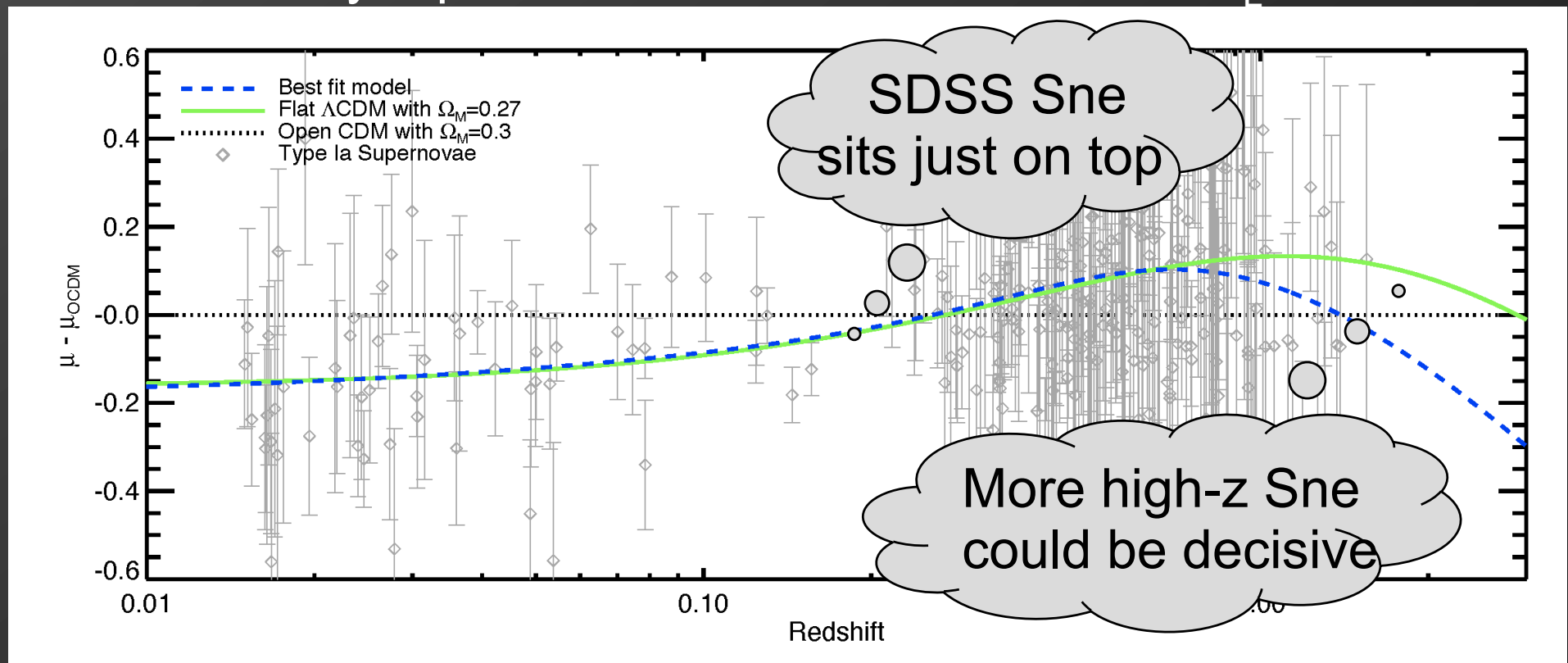
**Constraining LTB Models  
with  
Geometric  
Cosmological Data**

# Constraining Cosmological Data

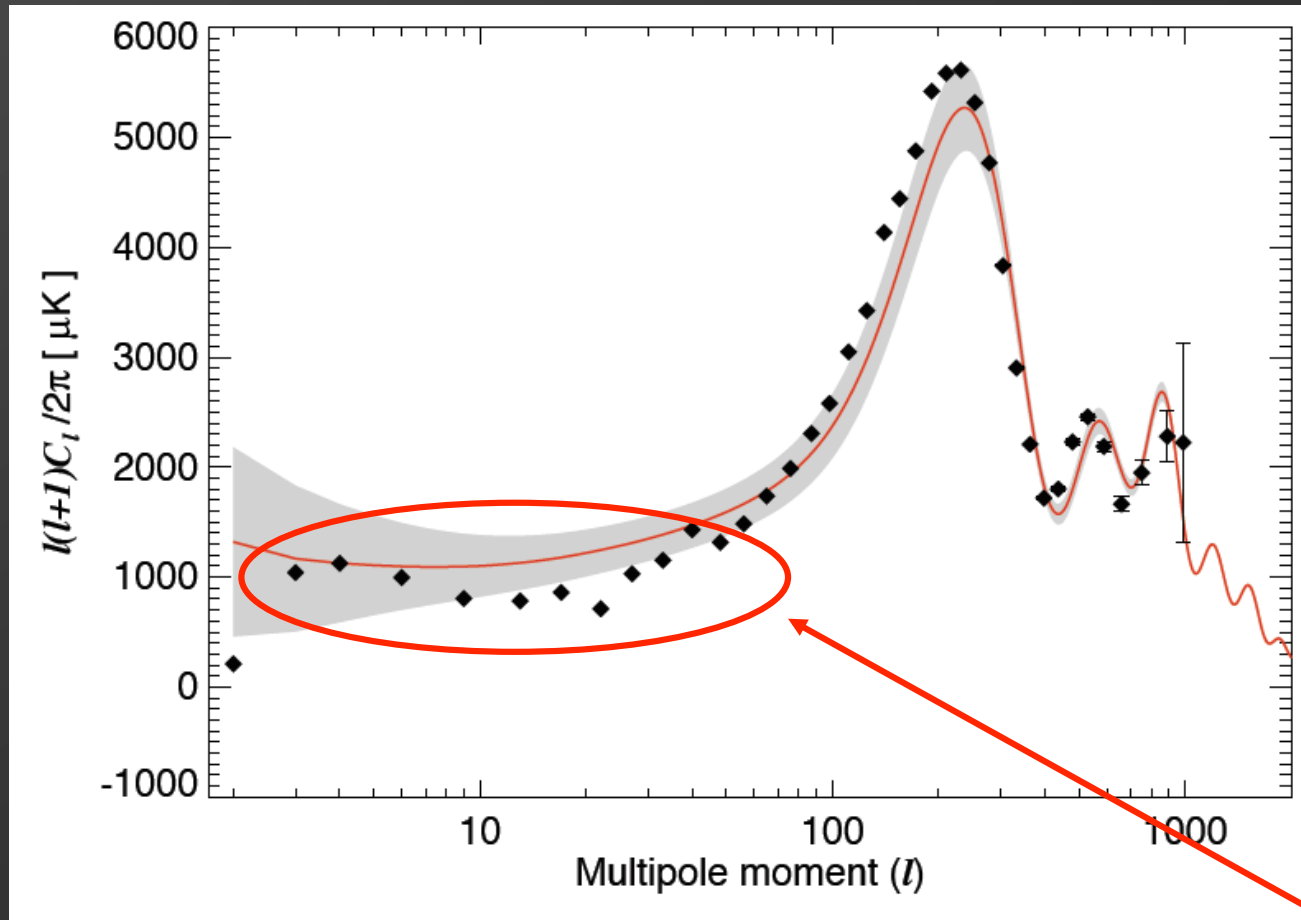
- **Type Ia Supernovae:** 307 SNe compiled by many different sources, and published as the UNION data set.
  - Simple to do since we just fit against  $d_L(z)$
- **1<sup>st</sup> acoustic peak in the CMB:**  $d_C(z_{\text{rec}})$ , sound horizon
- **Radial Baryon Acoustic Oscillations:**
  - Sound horizon, and the longitudinal Hubble rate  $H_L$

# Fitting the Type Ia Supernovae

The best fit model has no problems with Type Ia SNe,  
It basically reproduces the standard  $\Lambda$ CDM  $z$ - $d_L$  relation.

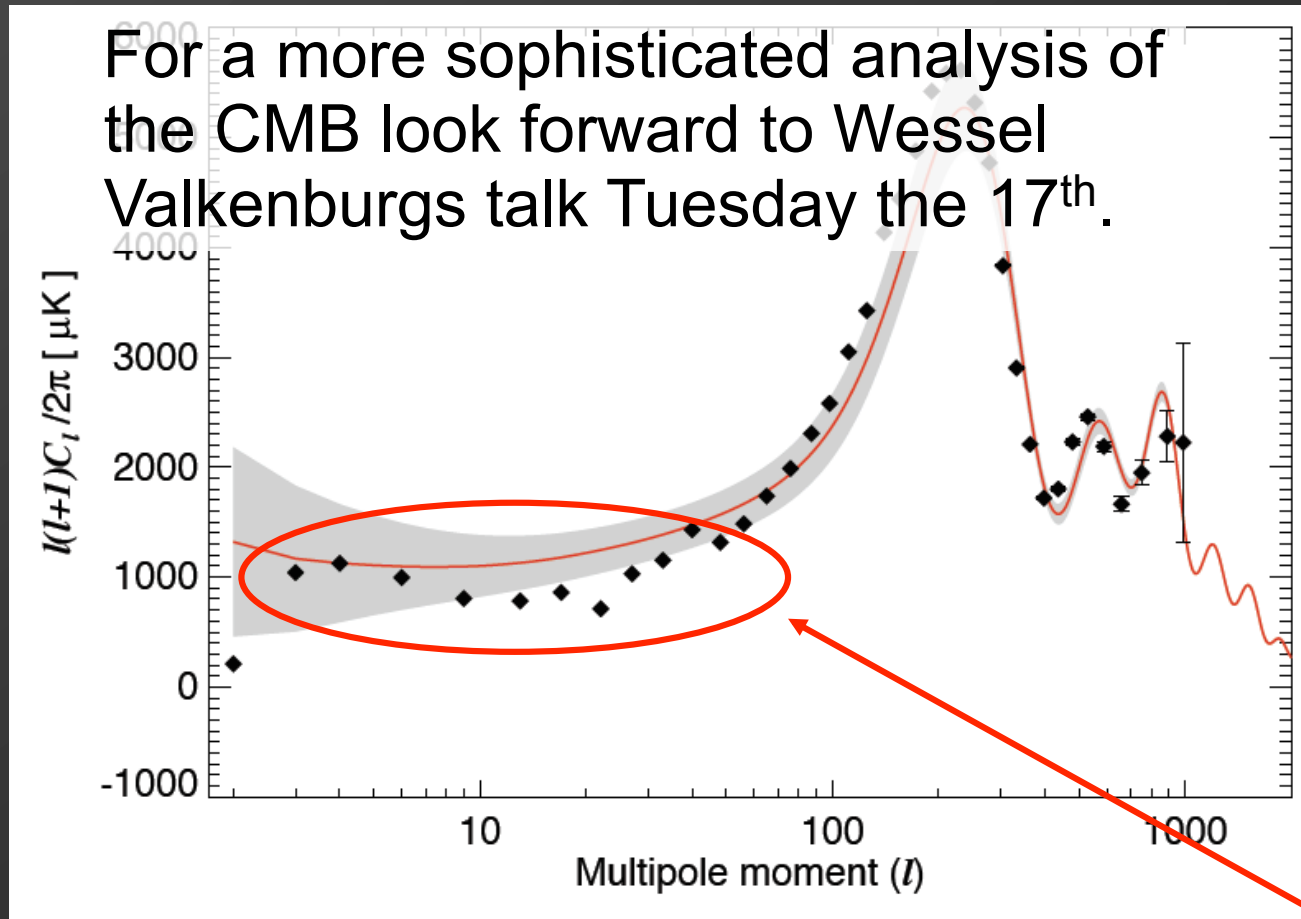


# Fitting the 1<sup>st</sup> Peak in the CMB



- The fit to the *first* peak is ok - we did not try to fit all data
- LTB perturbation theory is needed to explain low  $l$  (ISW)

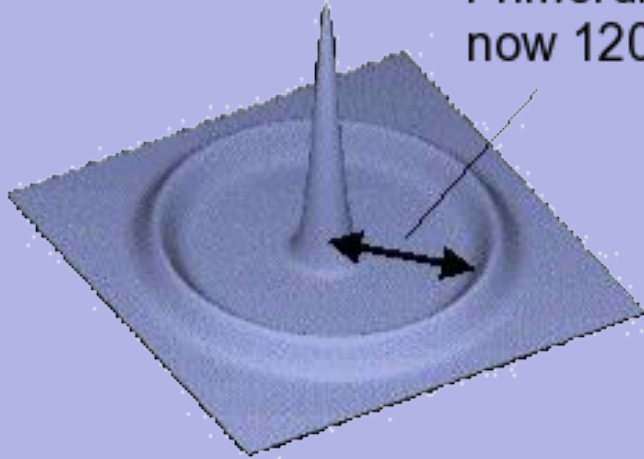
# Fitting the 1<sup>st</sup> Peak in the CMB



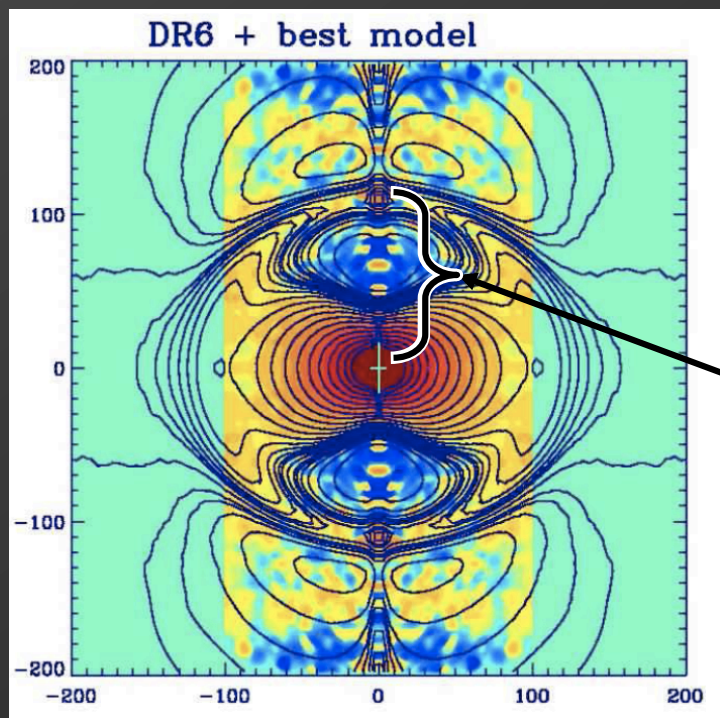
- The fit to the *first* peak is ok - we did not try to fit all data
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# Radial Baryon Acoustic Scale

Primordial sound wave,  
now 120 Mpc across



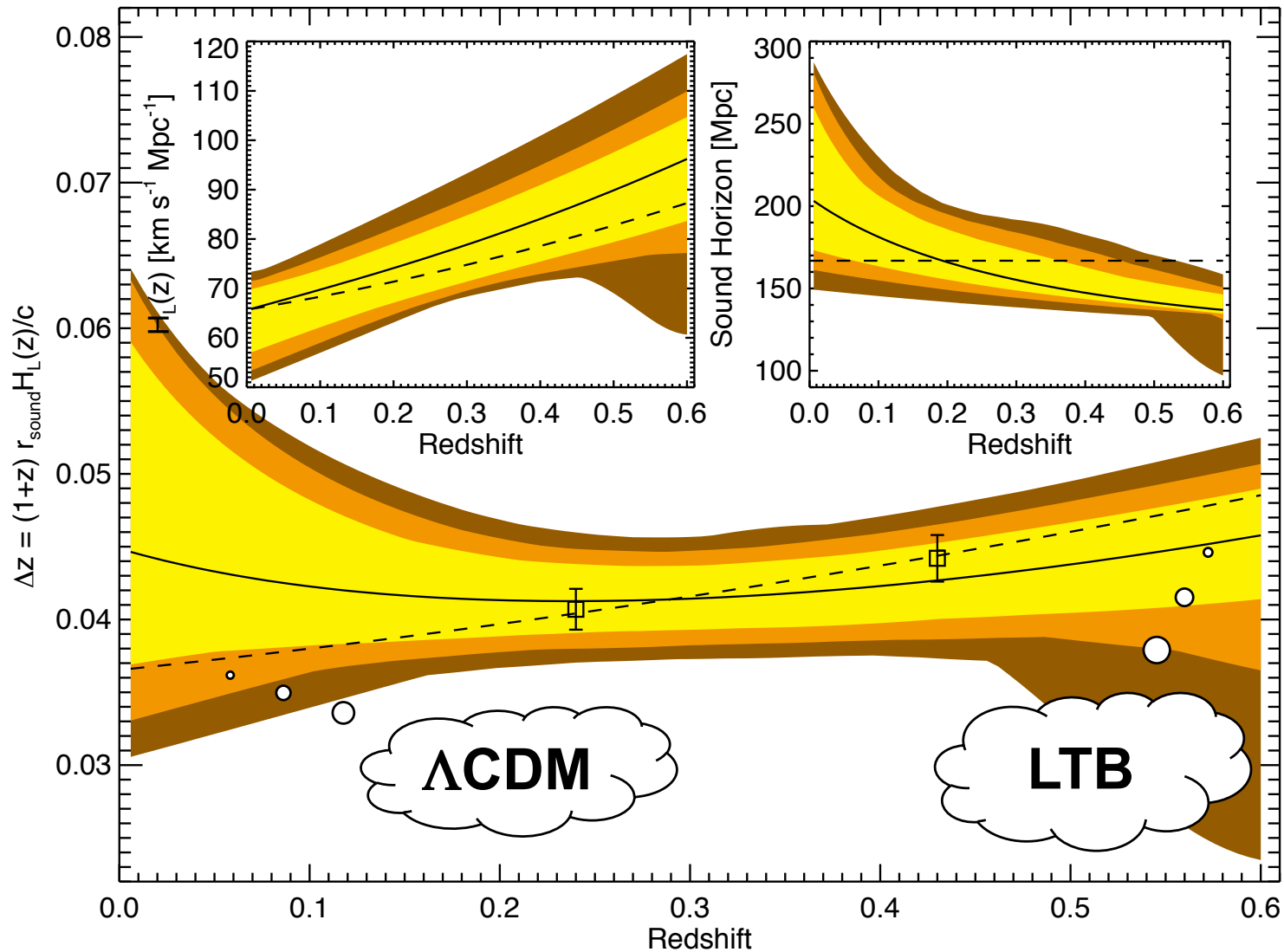
- Acoustic oscillations are frozen sound waves and the physical size is known from early universe physics



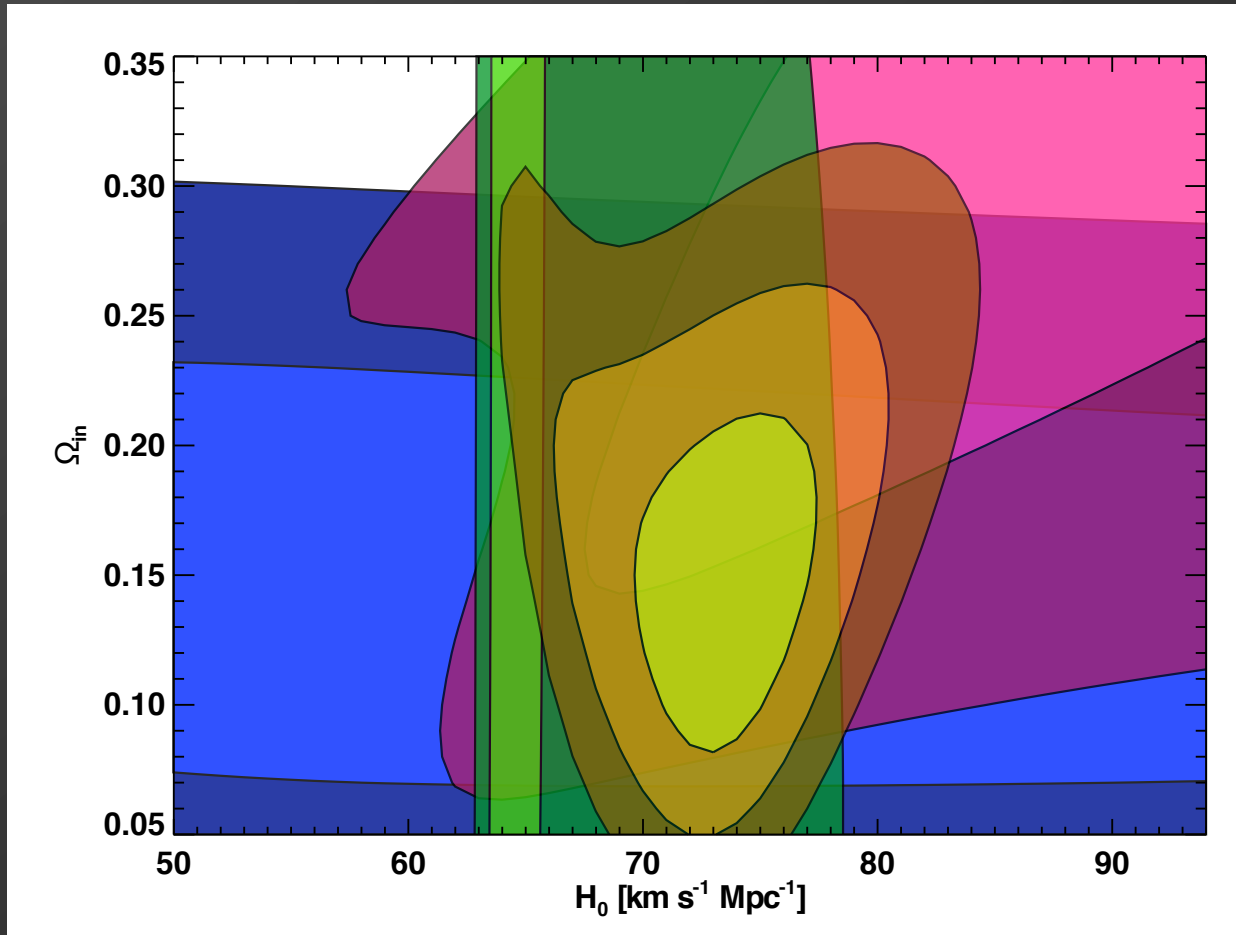
- Measuring the correlation function along the line of sight gives a clean measure of  $H_L(z)$

$$\Delta z_{LTB} = \frac{H_L(z)r_s(z)}{c}$$

# Radial Baryon Acoustic Scale



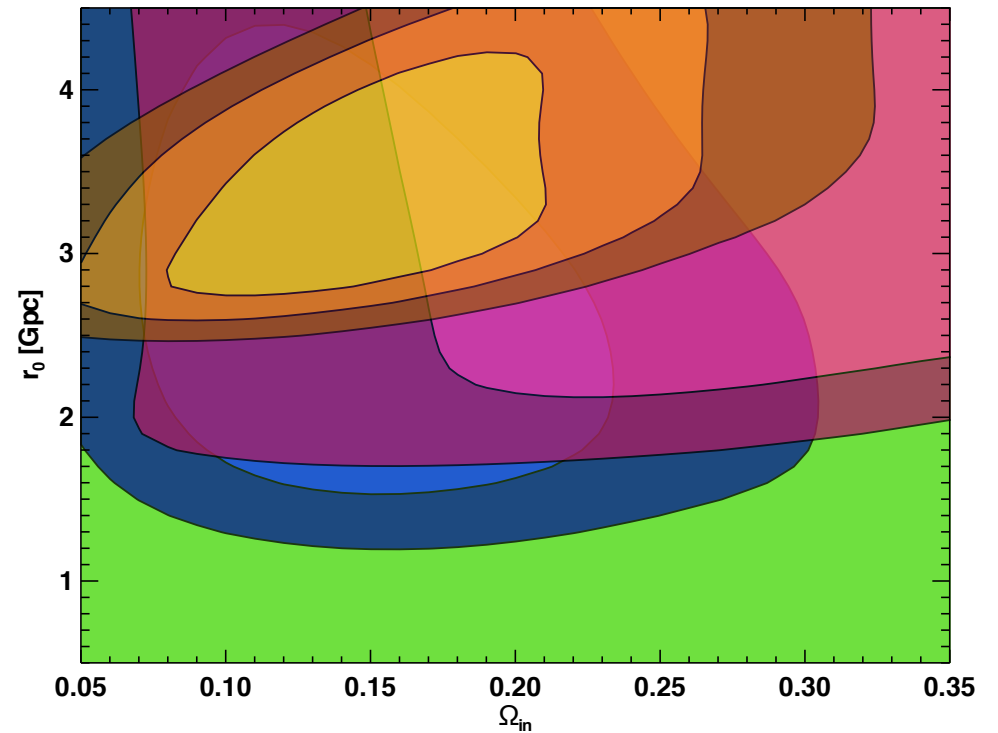
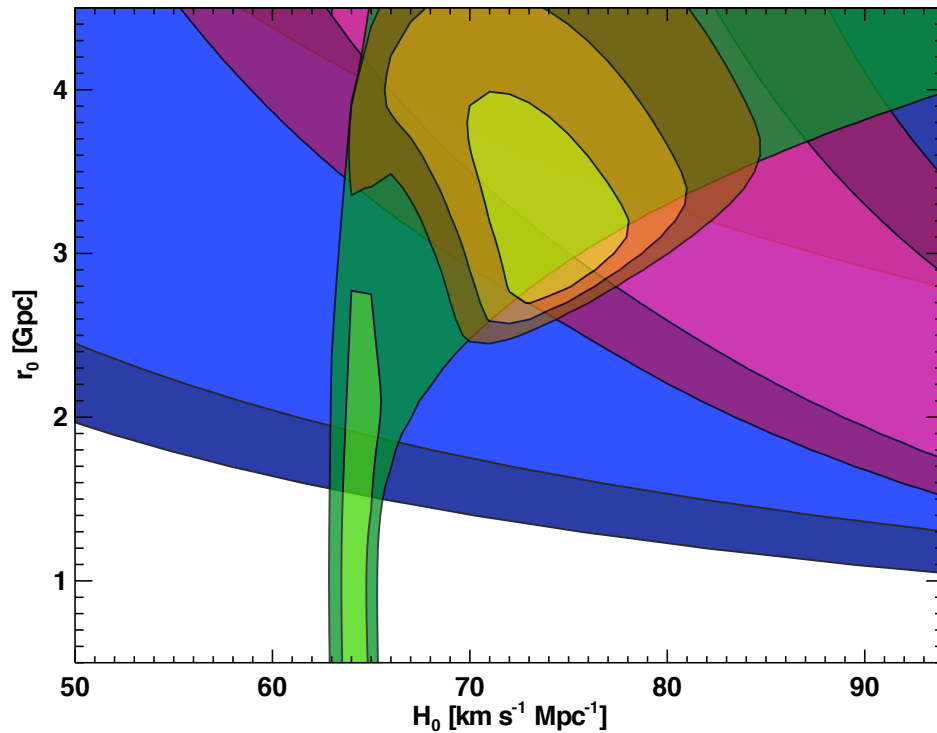
# Scanning the Model Space



- **Yellow: Everything, Blue: SNe Ia. Green: CMB. Purple: BAO**
- The Type Ia Supernovae constrain  $\Omega_{\text{matter}}$
- CMB constrains the Hubble param, because  $\Omega_{\text{out}}=1$  &  $\omega_b=\text{const}$



# Scanning the Model Space



- **Yellow: Everything, Blue: SNe Ia. Green: CMB. Purple: BAO**
- The SNe and BAO pushes the void size to  $> 2$  Gpc
- Some tension between BAO and SNe (high-z SNe are useful)

# Best Fit and Marginalised Errors

	$H_0$	$H_{r=0}$	$H_{r=\infty}$
units	100 km s <sup>-1</sup> Mpc <sup>-1</sup>		
Priors	0.50–0.95	0.4–0.89	0.33–0.63
Best Fit $\pm 2\text{-}\sigma$	0.67 $\pm$ 0.03	0.58	0.45
	$\Omega_{\text{in}}$	$r_0$	$\Delta r$
units		Gpc	$r_0$
Priors	0.05–0.35	0.5–4.5	0.1–0.9
Best Fit $\pm 2\text{-}\sigma$	0.16 $\pm$ 0.09	2.7 $\pm$ 0.8	0.44(>0.12)

# Conclusions From First Part

- An LTB model of a giant void can convincingly fit a large set of current cosmological observations and do it as well as the  $\Lambda$ CDM model.
- The void model only contains 4 parameters: It is a simple model.
- The best fit void size is  $\sim 2.7 \pm 1$  Gpc, approximately the size of the cold spot in the CMB, if it was near the surface of last scattering.
- Combining many small voids in a “swiss-cheese universe” does not seem to work

3- $\sigma$

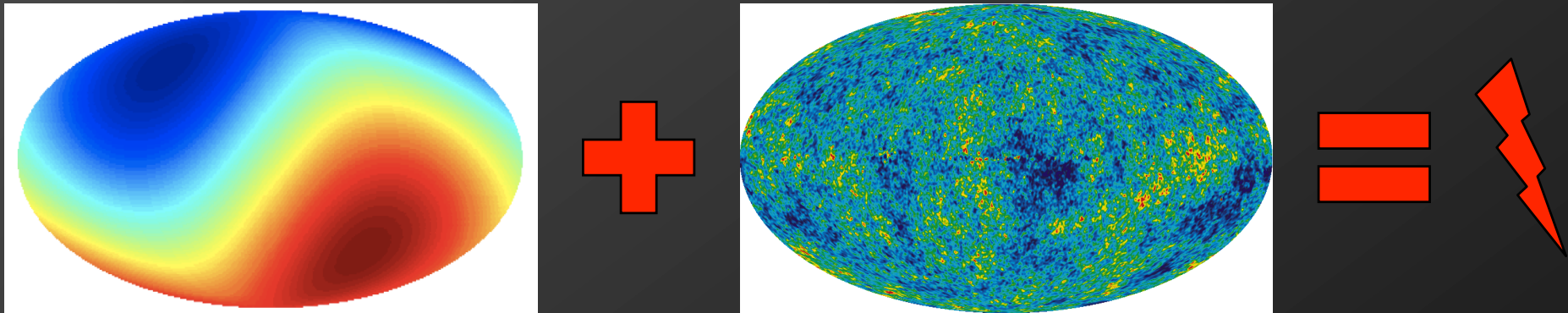
2- $\sigma$

1- $\sigma$

# Constraining Void Models with “non-geometric” measurements

$\sigma = -1600 \text{ y} \quad -750 \alpha =$

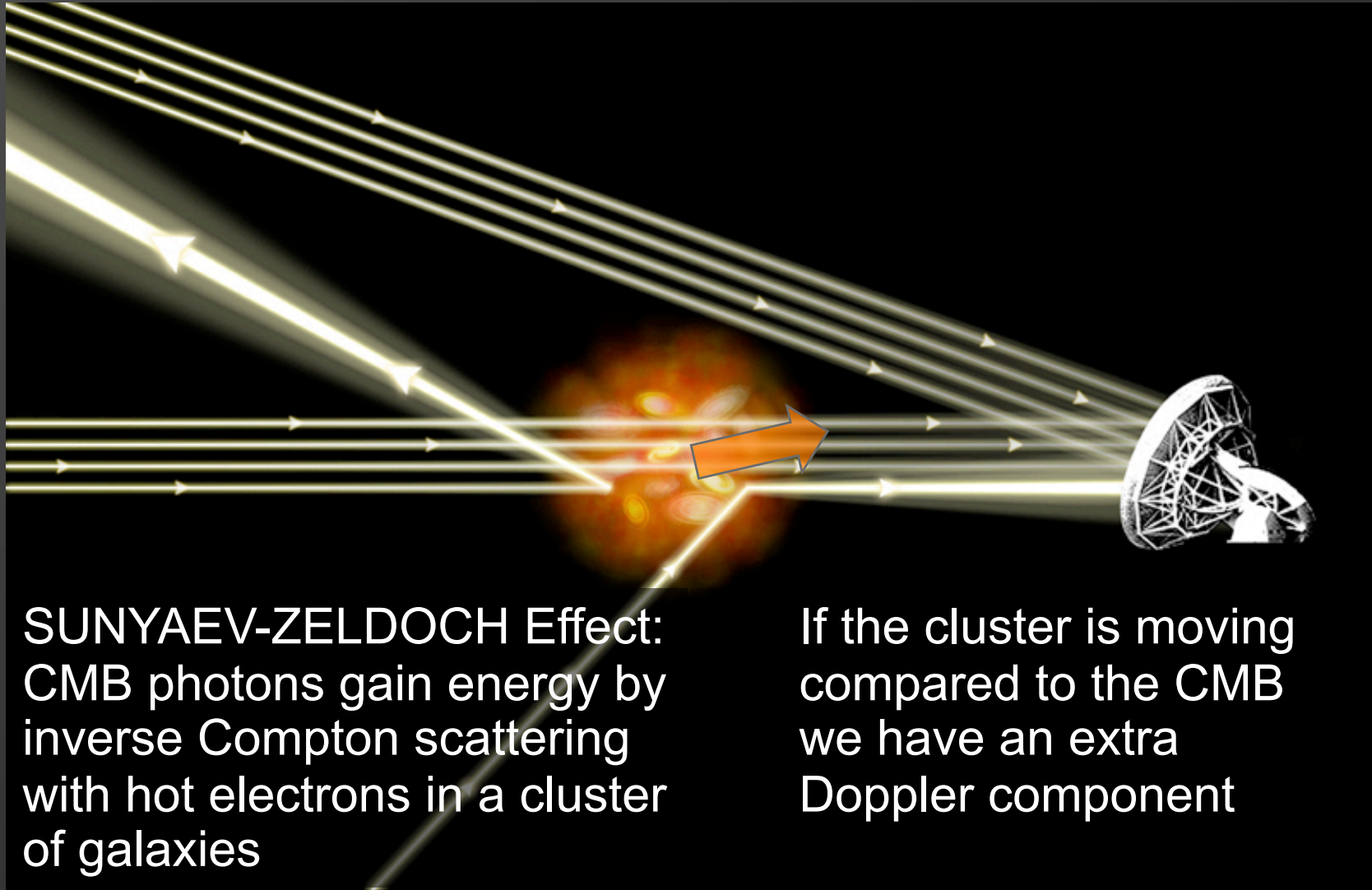
# Using the CMB dipole to constrain our position in the void



- An off-center observer will see an extra dipole in the CMB because of different *integrated* expansion rates – i.e. redshifts – in different directions.
- Same effect for Supernovae, can be combined
- For a large void the constraint is pretty tight  $\sim 40 \text{ Mpc } h^{-1}$

(Blomqvist & Mörtzell 2010)

# Using the Kinetic Sunyaev-Zeldovich Effect



# Using the Kinetic Sunyaev-Zeldovich Effect

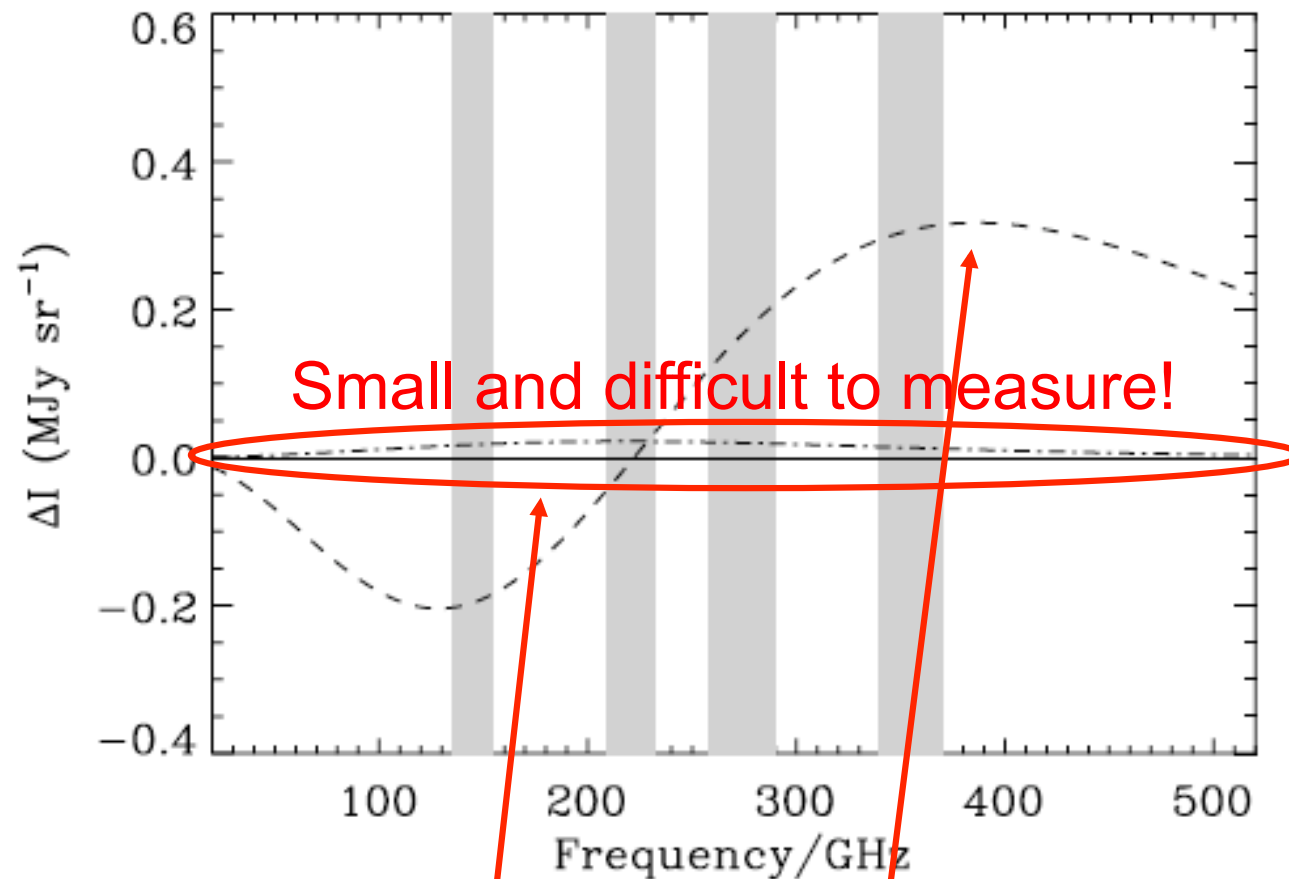
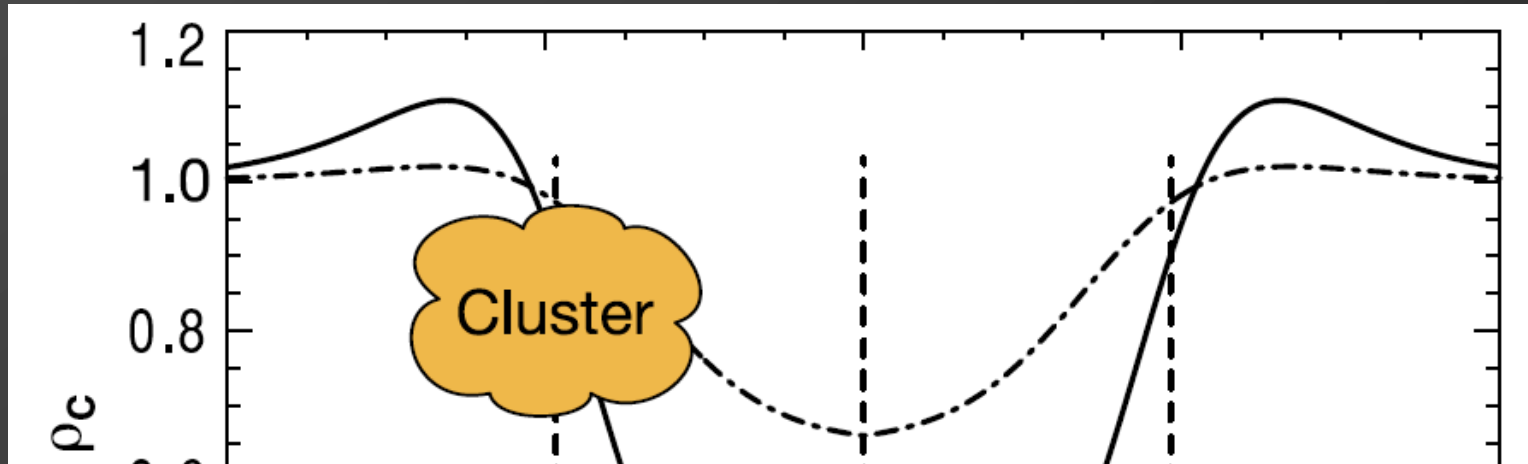
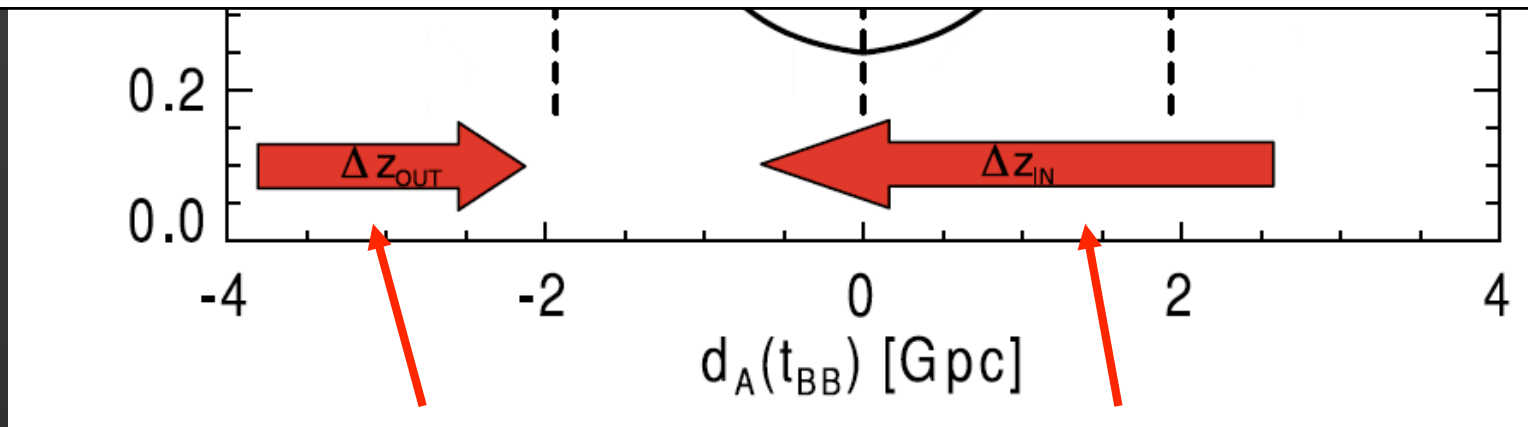


FIG. 1.—Frequency dependence of the SZ effect for a cluster with optical depth  $\tau = 0.01$ , gas temperature 10 keV, and a peculiar velocity of  $-500 \text{ km s}^{-1}$  (toward the observer). The thermal SZ spectrum is indicated by the dashed line, the kinematic effect by the dot-dashed line. The shaded regions indicate the bands in which SuZIE II observes.

# The Induced Dipole for an Off-centered Cluster



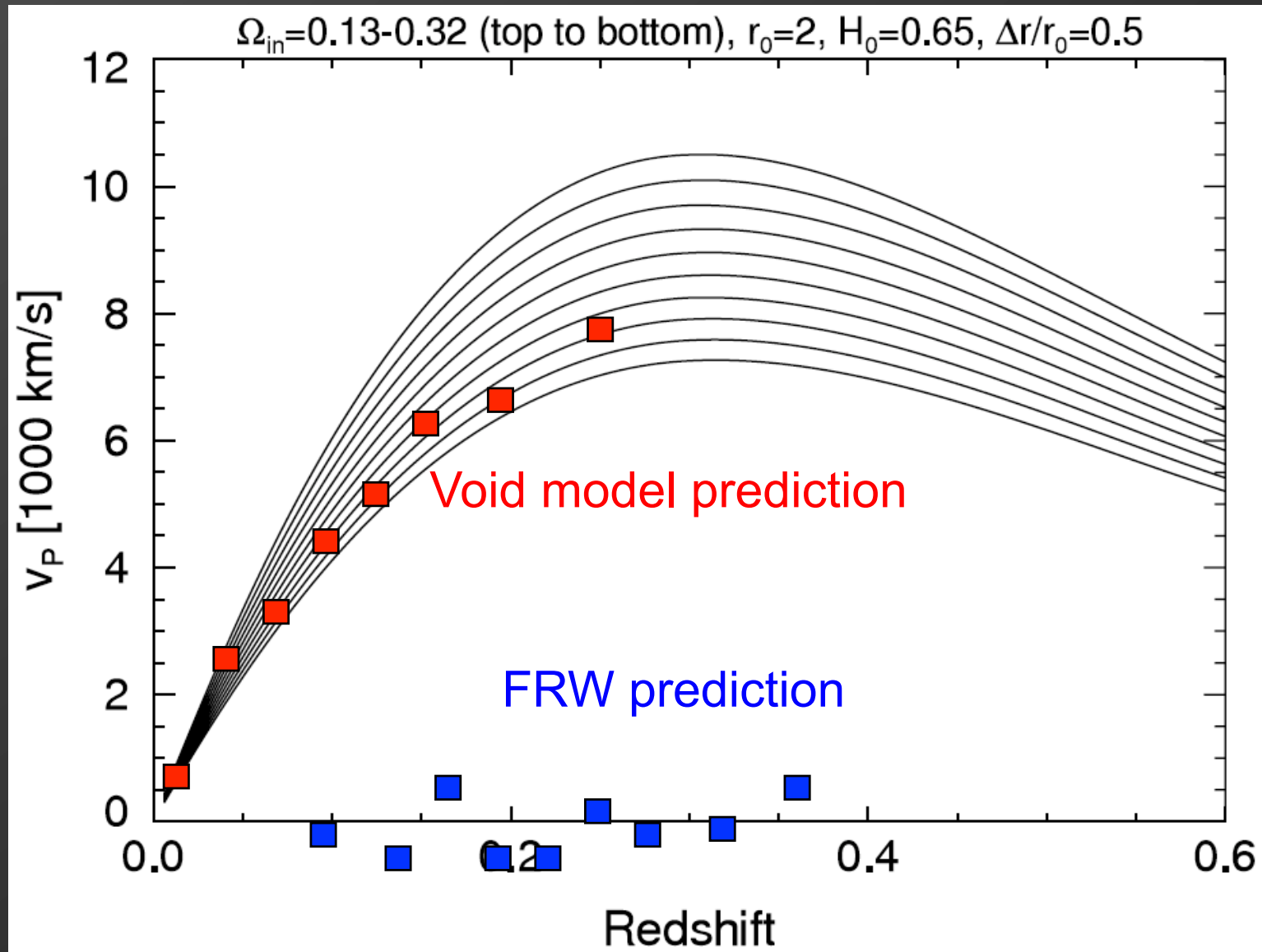
$$\frac{v_p}{c} = \frac{\Delta T}{T_{\text{dipole}}} = \frac{T(\theta) - \hat{T}}{\hat{T}} = \frac{|T_{\text{in}} - T_{\text{out}}|}{T_{\text{in}} + T_{\text{out}}} = \frac{|z_{\text{in}} - z_{\text{out}}|}{2 + z_{\text{in}} + z_{\text{out}}}$$



Different expansion histories => different redshifts!

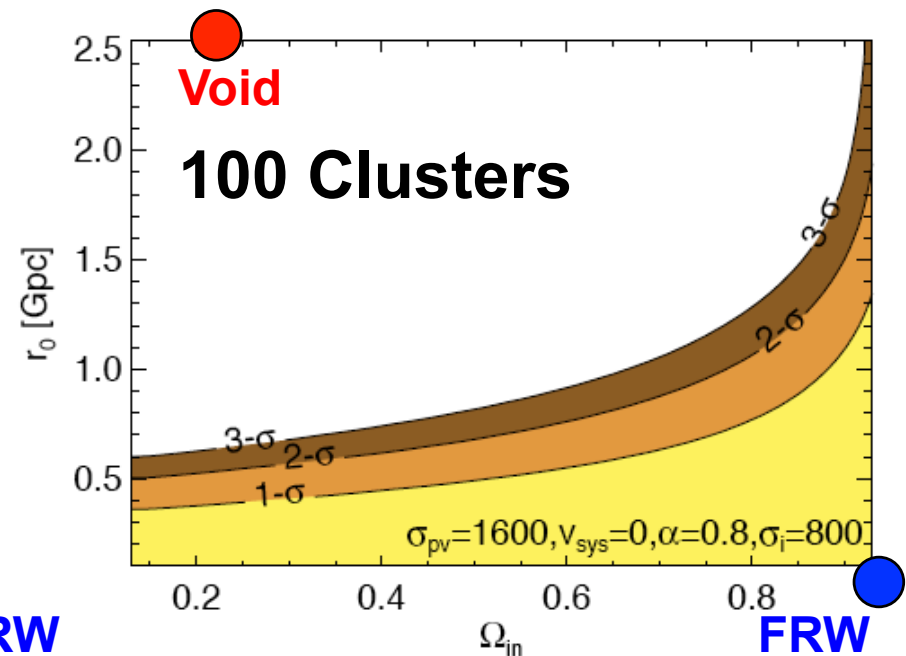
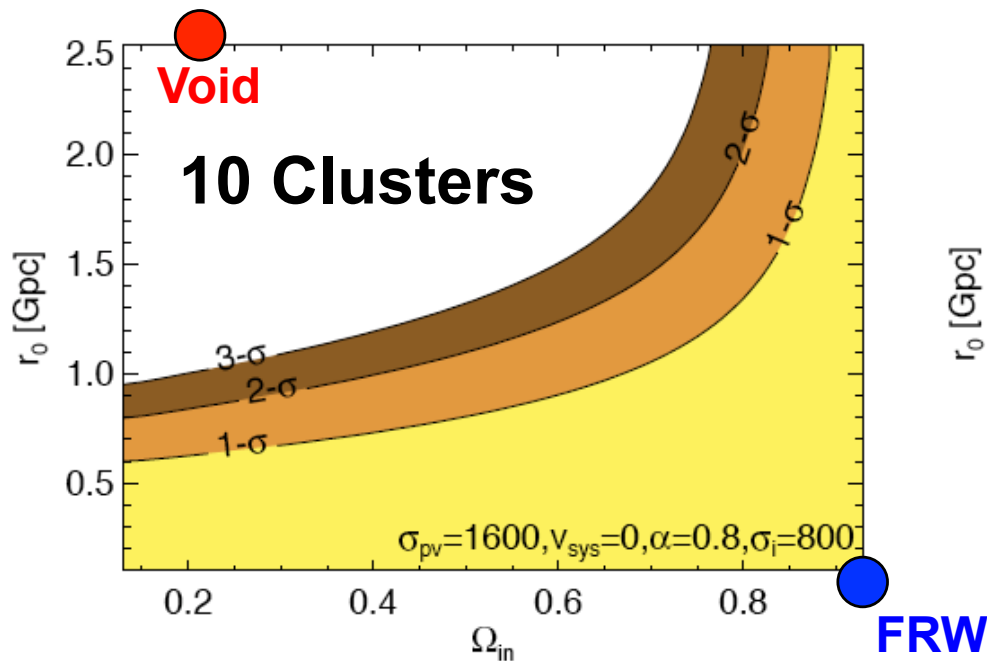


# Forecasted kSZ for void models



# Future Bounds from ACT or SPT

- While the ACT and the SPT telescopes will make thousands of thermal SZ cluster observations we need follow up in X-rays, radio and/or optical for kSZ
- In the very first kSZ data release the LTB model could be definitively ruled out



# Using Cosmic reionization as a mirror

$u$ -distortion

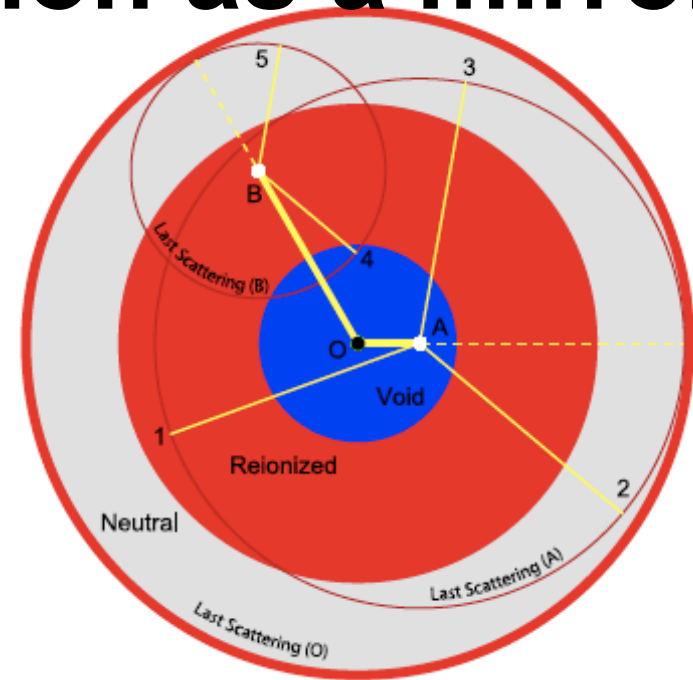
A blackbody spectrum at temperature  $T$  mixed with a blackbody at temperature  $T+\Delta T$  produces a  $u$ -distorted blackbody.

Stebbins, astro-ph/0703541

$$u[\hat{n}] = \frac{3}{16\pi} \int_0^\infty dz' \frac{d\tau}{dz'} \int d\hat{n}' (1 + (\hat{n} \cdot \hat{n}')^2) \times \left( \frac{\Delta T}{T} [\hat{n}, \hat{n}, z] - \frac{\Delta T}{T} [\hat{n}', \hat{n}, z] \right)^2$$

Degenerate with Compton  $\gamma$ -distortion parameter:  $u = 2\gamma$

FIRAS:  $\gamma < 15 \times 10^{-6}$  (95%): Fixen et al, ApJ 473, 576 (1996)



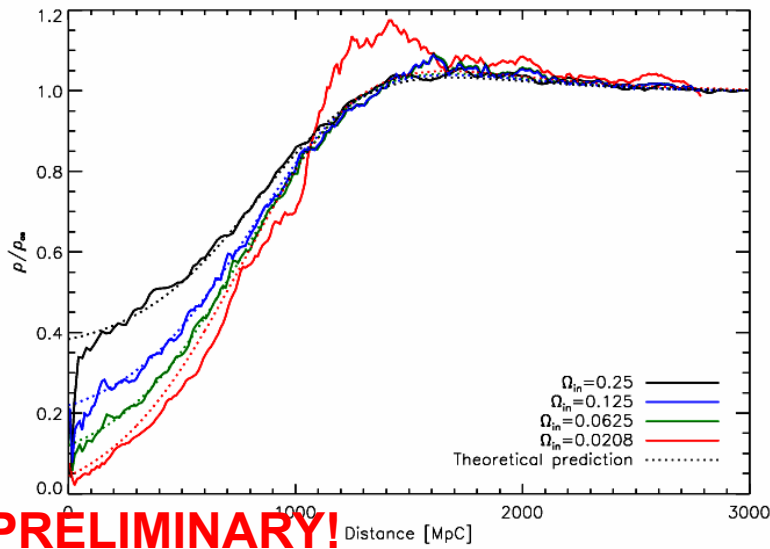
# The next frontier: Structure Formation & Perturbations

(Dunsby et al 2010)

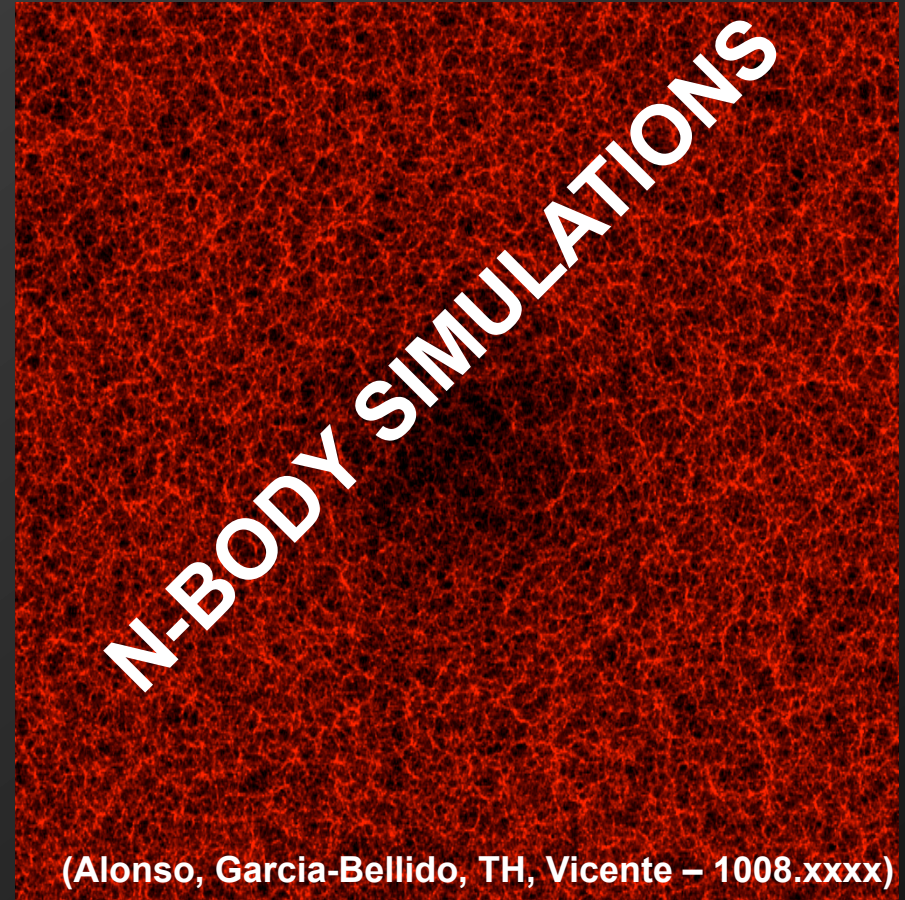
$$\begin{aligned} \frac{1}{3}\dot{\Theta} - \dot{\Sigma} &= \left(\frac{1}{3}\Theta + \Sigma\right)^2 + \frac{4}{3}\pi G\rho + \mathcal{E} \\ -\dot{\Theta} &= \frac{1}{3}\Theta^2 + \frac{3}{2}\Sigma^2 + 4\pi G\rho \\ \dot{\Sigma} &= -\left(\frac{2}{3}\Theta + \frac{1}{2}\Sigma\right)\Sigma - \mathcal{E} \\ \dot{\rho} &= -\Theta\rho \\ \dot{\mathcal{E}} &= \left(\frac{3}{2}\Sigma - \Theta\right)\mathcal{E} - 4\pi G\rho\Sigma + \varepsilon_{ab}d^a d^b, \end{aligned}$$

propagation equations along  $n^a$  are

$$\begin{aligned} \phi' &= -\frac{1}{2}\phi^2 - \left(\frac{1}{3}\Theta + \Sigma\right)\left(\frac{2}{3}\Theta - \Sigma\right) - \frac{16}{3}\pi G\rho - \mathcal{E} + d_a a^a \\ \xi' &= -\phi\xi + \frac{1}{2}\varepsilon_{ab}d^a d^b \\ \zeta'_{\{ab\}} &= -\phi\zeta_{\{ab\}} + d_{\{a}a_{b\}} + \left(\frac{1}{3}\Theta + \Sigma\right)\Sigma_{ab} - \mathcal{E}_{ab} \\ \Sigma' - \frac{2}{3}\Theta' &= -\frac{2}{3}\Sigma d_a \Sigma^a \end{aligned}$$



**PRELIMINARY!**



(Alonso, Garcia-Bellido, TH, Vicente – 1008.xxxx)

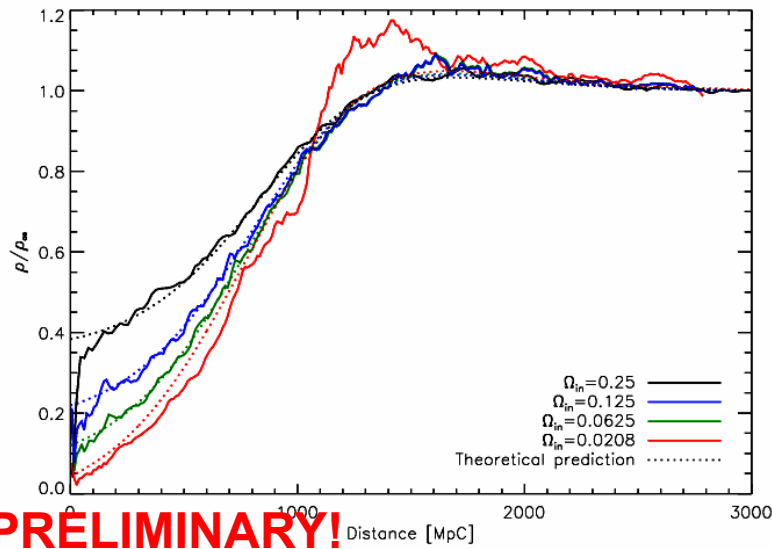
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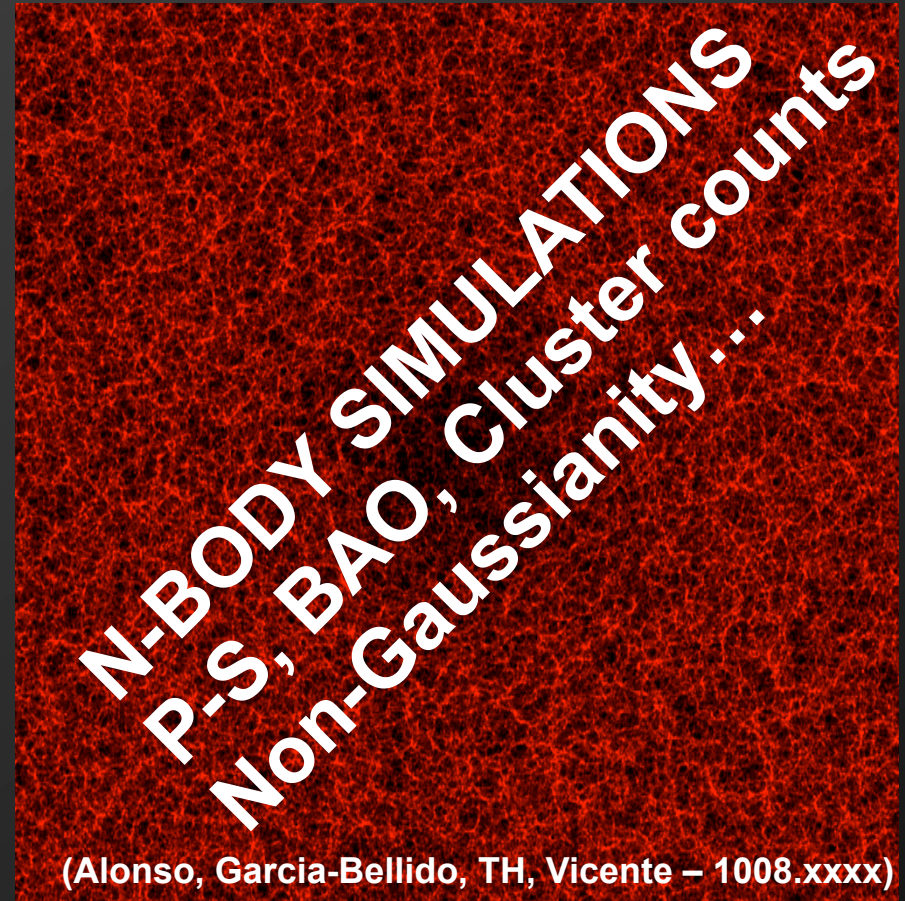
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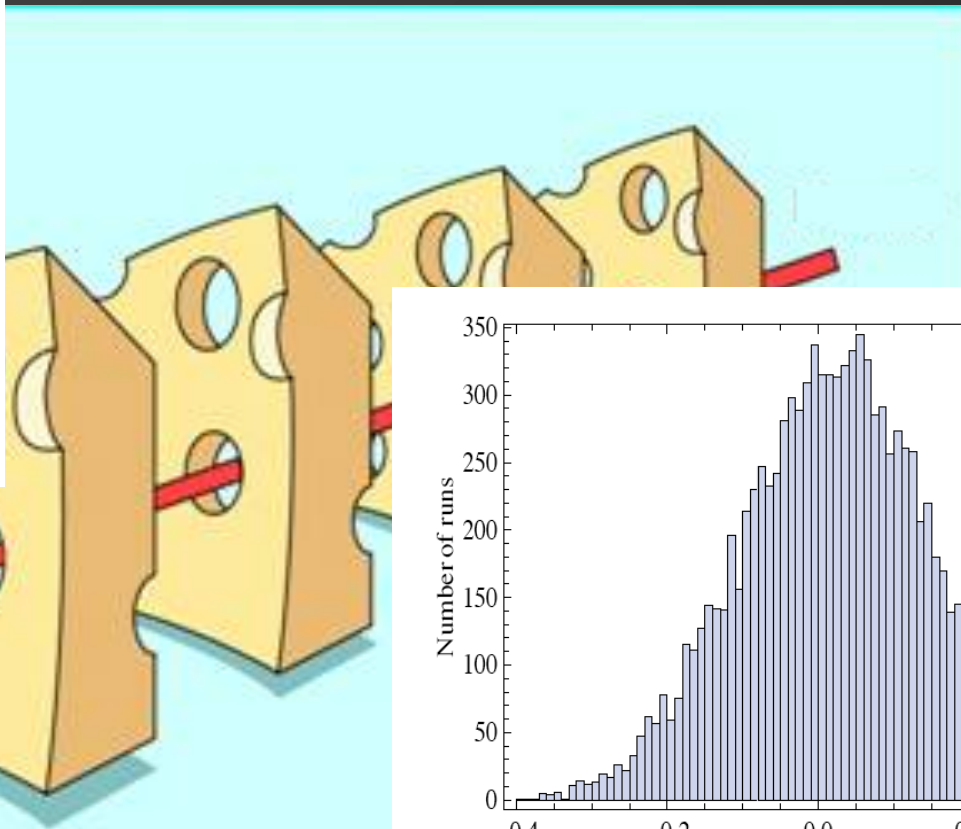
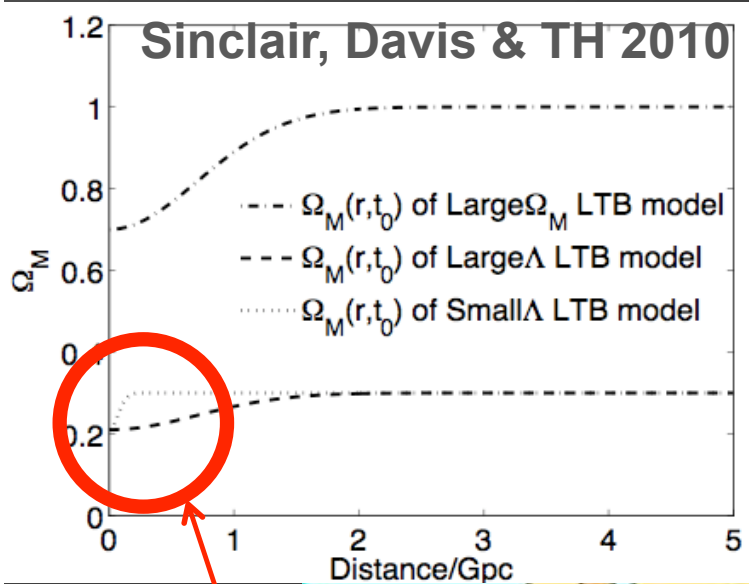
# Summary

- Void models, observationally, seem a real alternative to the standard model. While they break with the Copernican principle, they do not need dark energy.
- A void model with a size of  $\sim 2.7$  Gpc yields a very good fit to observations constraining the *geometry* of the universe.
- There are many new observations and ways to tests the model related to structure, and to distortions of the CMB:
  - CMB and large scale structure (More theory+simulations)
  - Remote measurements of the CMB: The Kinetic Sunyaev Zeldovich effect (Apex, ACT, SPT, Planck)
  - Spectral distortions in the CMB (FIRAS is almost good enough)

# Outlook

- The void model may be ruled out in the near future, just like the Swiss Cheese model was
- But given the quality of the observations today, people will keep testing inhomogeneous models – it is the next natural step after FRW
- The void model may be wrong, but it helps in getting a better understanding of the **perturbed  $\Lambda$ CDM model**
- Example:
  - How are  $\Omega_{\text{matter}}$  and  $\Omega_{\Lambda}$  biased by a local void? How do we cope with the bias ?

# Using LTB to model the impact of a small void on Supernova Cosmology



The last hole in the cheese matters

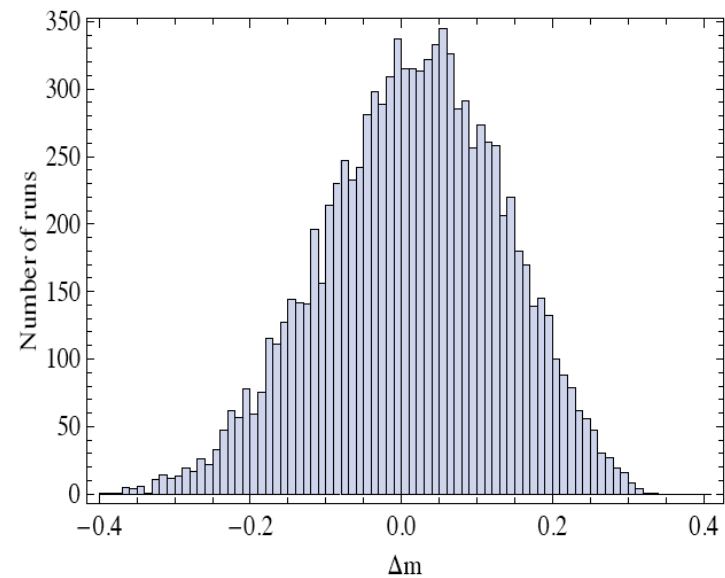
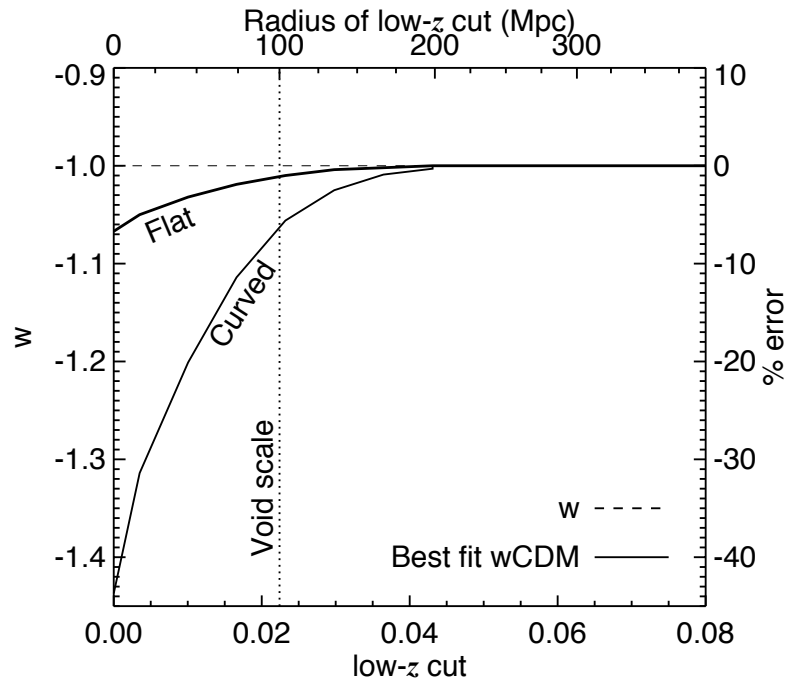


FIG. 2: A histogram of the distance modulus shifts  $\Delta m$  found for 10,000 realizations of our randomized void scenario.

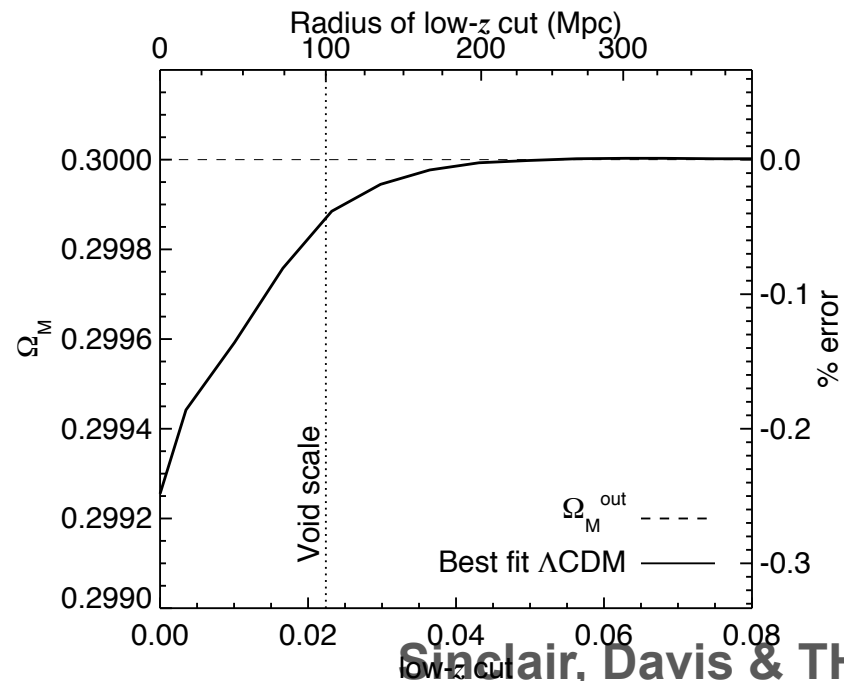
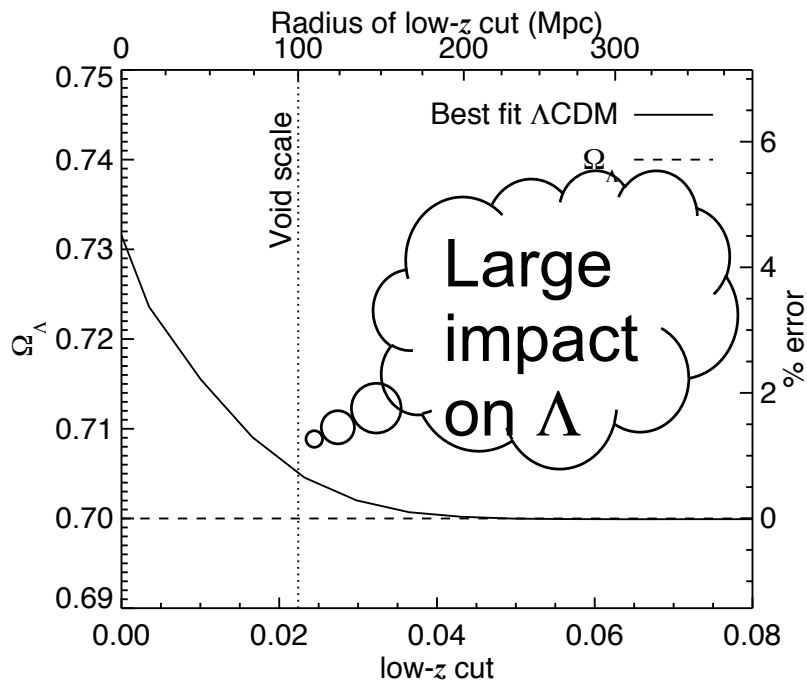
Vanderveld et al 2008



# Impact of a small void on Supernova Cosmology



- The LTB model is extended with  $\Lambda$
- We can then embed a small void – **we use  $70 \text{ Mpc } h^{-1}$**  – in a LCDM universe
- We construct a **sample of 301 supernovae with  $z < 1.7$**
- When fitting a LCDM model to the mock data, the impact can be appreciable for upcoming surveys
- It is essential to **discard low redshift supernovae with  $z < 0.02$  for a  $< 1\%$  error**



# Can we measure geometry ??

- The shear to expansion ratio gives a normalised measure of inhomogeneity, with a straight forward theoretical interpretation

$$\varepsilon = \sqrt{\frac{3}{2}} \frac{\sigma}{\Theta} = \frac{H_T - H_L}{H_L + 2H_T}$$

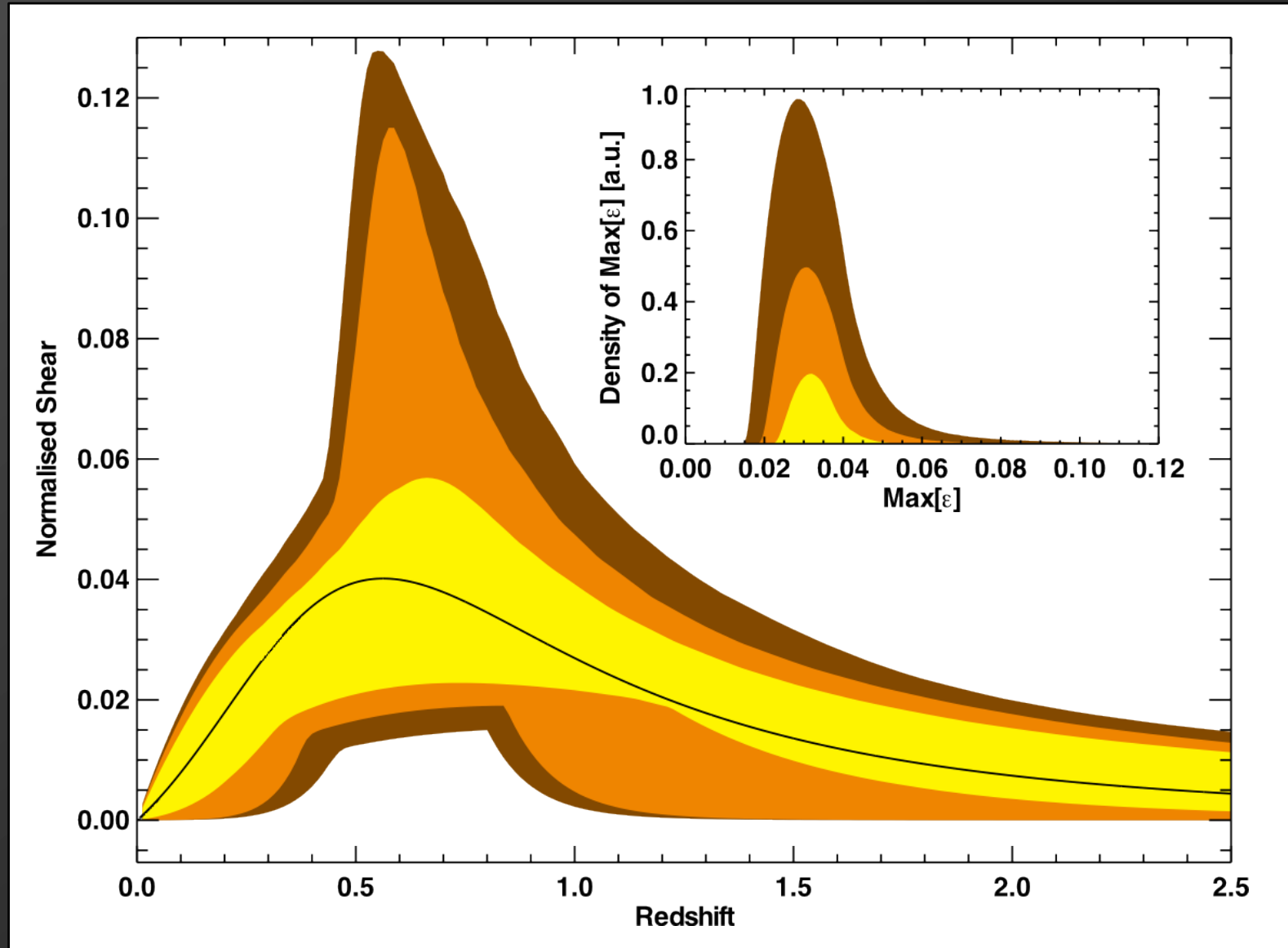
- Notice that  $H_L = H_T + \partial_{\ln A} H_T$

**Radial inhomogeneity  $\Leftrightarrow \varepsilon \neq 0$**

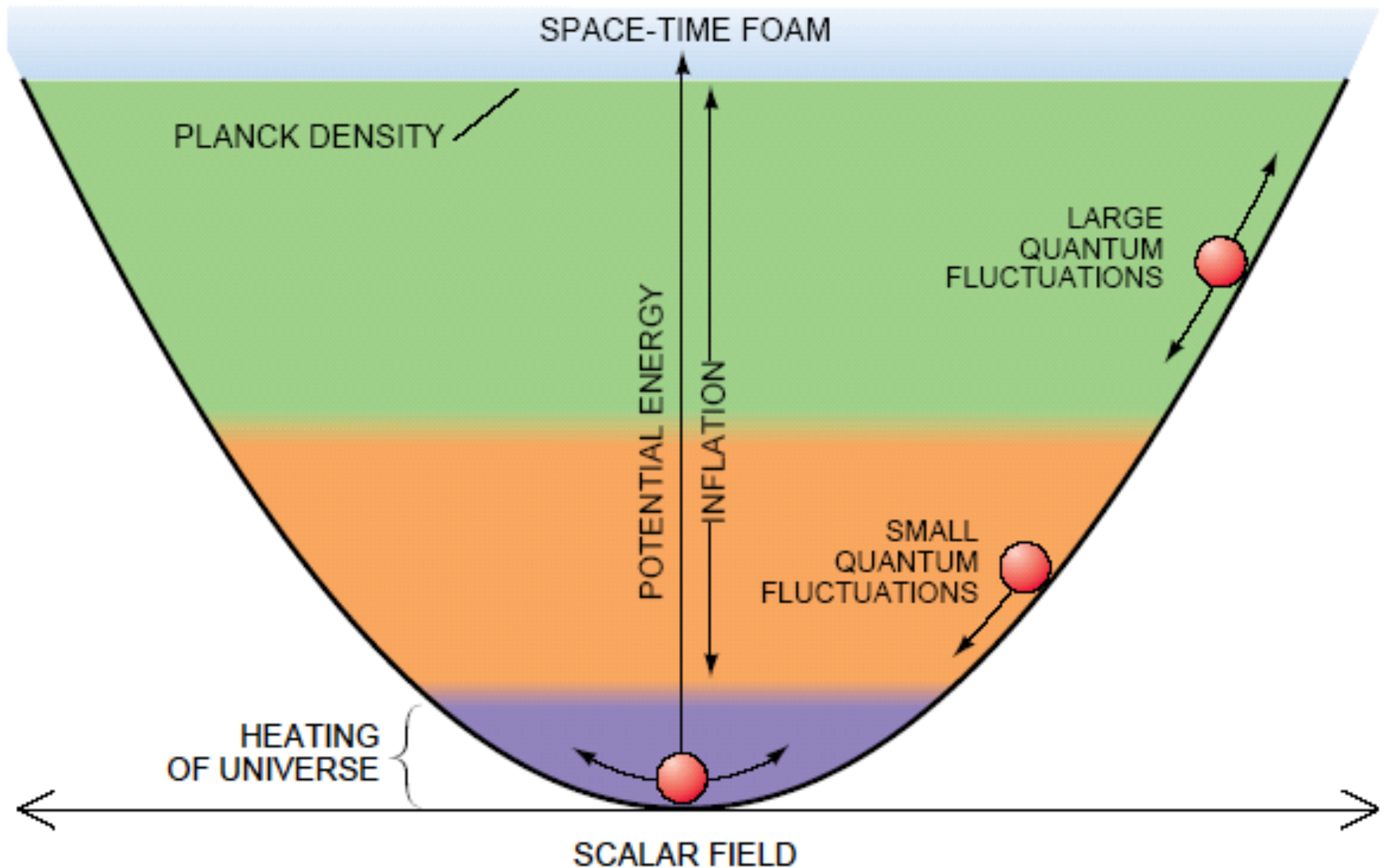
- The strength of  $\varepsilon$  is that it can be measured in the future by combining radial BAO measurements (giving  $H_L$ ) and measurements of the distance (giving  $d_A$ )

$$\varepsilon = \frac{\sqrt{1 - k(r)} - H_L(d_A + (1 + z) d'_A(z))}{2\sqrt{1 - k(r)} + H_L(d_A - 2(1 + z) d'_A(z))}$$

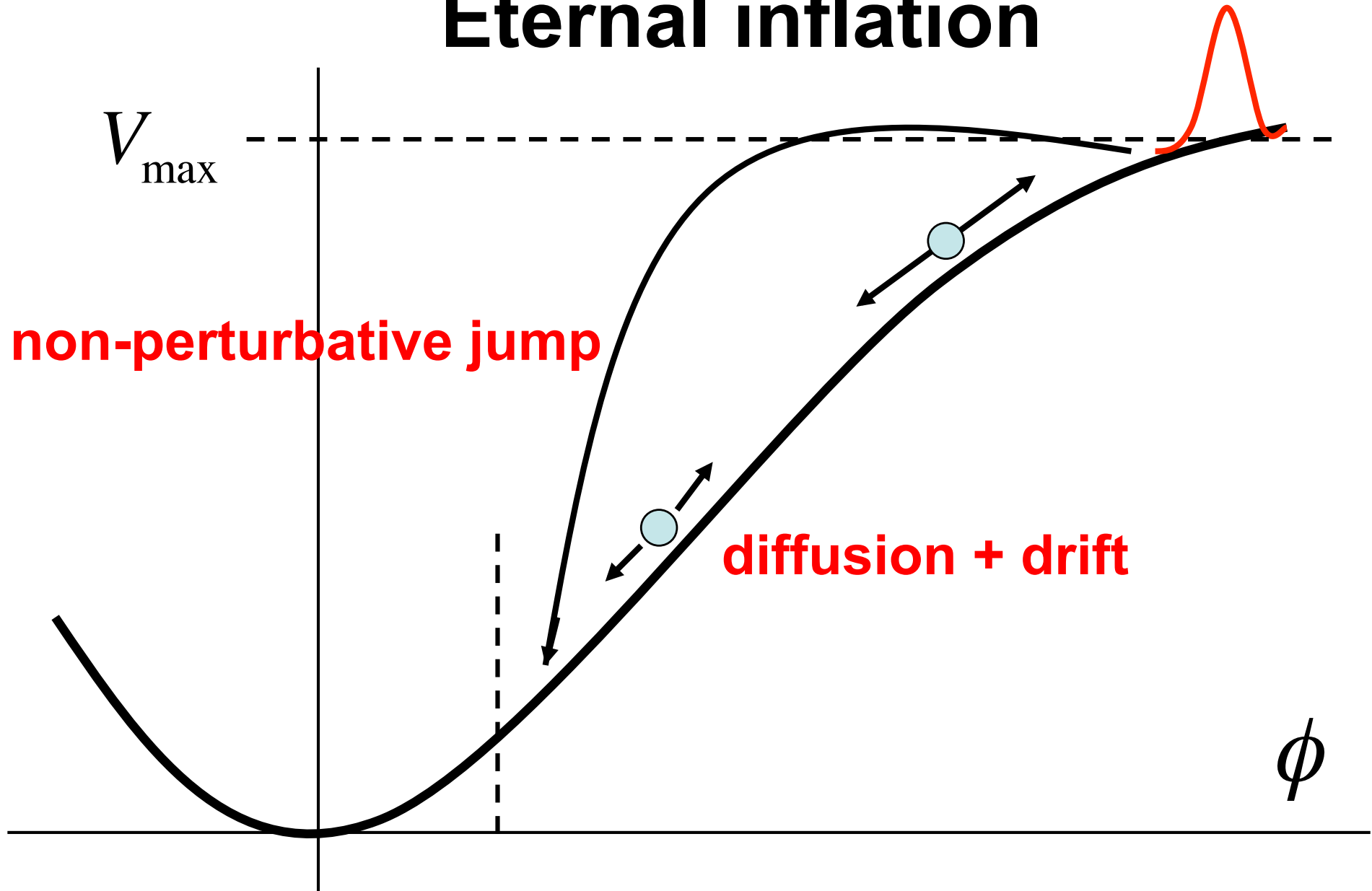
# Can we measure geometry ??



# A void model from Chaotic Inflation

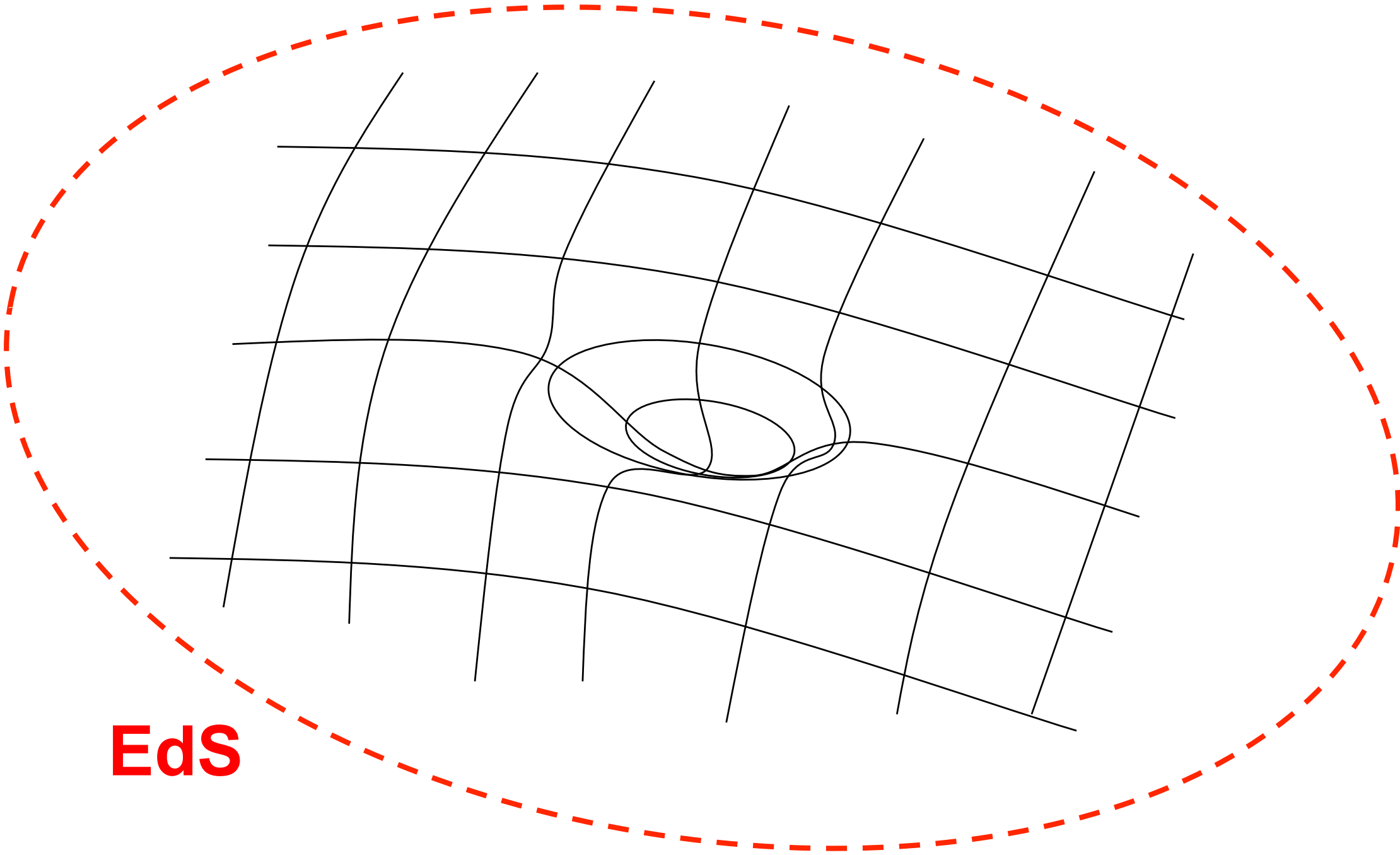


# Eternal inflation



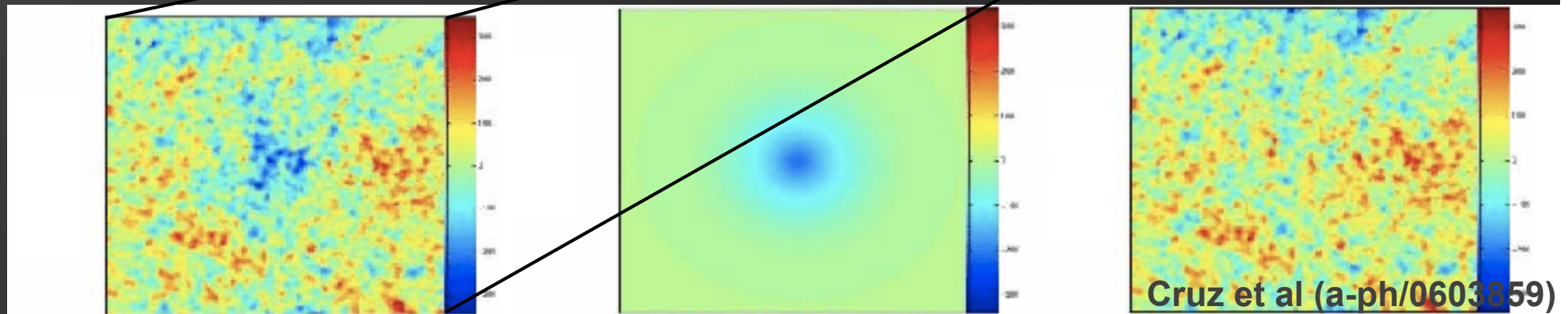
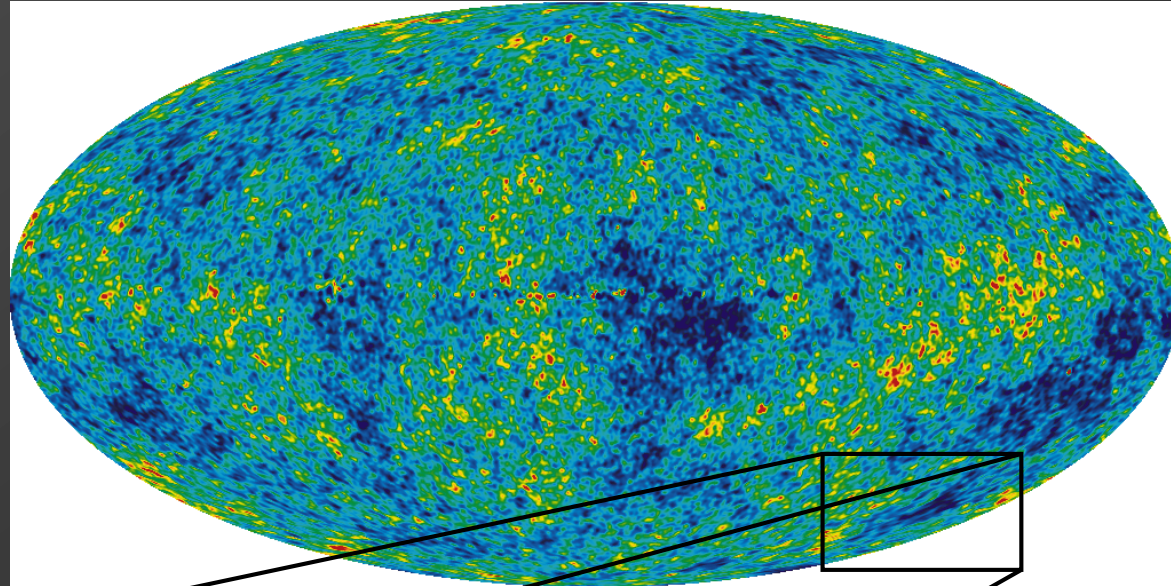
(Linde, Linde, Mezhlumian, PLB 345, 203, 1995)

# The Inflow = LTB Model



**EdS**

# Could the *Cold Spot* in the cosmic microwave background be an inflow ?



Approximately 2 Gpc comoving size