GRAVITATIONAL WAVES FROM/AFTER REHEATING

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GARCÍA-BELLIDO & DGF, **PRL 2007** GARCÍA-BELLIDO, DGF & SASTRE, **PRD 2008** FENU, DGF, DURRER & GARCÍA-BELLIDO, **JCAP 2009** DUFAUX, DGF & GARCÍA-BELLIDO, **PRD 2010**

BENASQUE 2010, August 17th 2010, SPAIN

GRAVITATIONAL WAVES (GW): PROBING the EARLY UNIVERSE ($t \leq 1$ s)

WEAKNESS of GRAVITY:

ADVANTAGE: GW DECOUPLE upon Production DISADVANTAGE: DIFFICULT DETECTION

O ADVANTAGE: GW → Probe for Early Universe

 $\rightarrow \left\{ \begin{array}{l} \mathbf{Decouple} \rightarrow \mathrm{Spectral} \ \mathrm{Form} \ \mathrm{Retained} \\ \mathbf{Specific} \ \mathbf{HEP} \ \Leftrightarrow \ \mathrm{Specific} \ \mathrm{GW} \end{array} \right.$

Physical Processes:

Inflation Reheating Phase Transitions Turbulence

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 Inflation Reheating Phase Transitions Turbulence

 $\textbf{INFLATION} \longrightarrow \textbf{REHEATING} \longrightarrow \textbf{BIG} \text{ BANG THEORY}$



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INFLATION \longrightarrow **REHEATING** \longrightarrow **BIG BANG THEORY**



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SCALAR REHEATING: SIMPLE EXAMPLES

$$V(\phi, \chi) = \frac{1}{4}\lambda\phi^{4} + \frac{1}{2}m_{\chi}^{2}\chi^{2} + \frac{1}{2}g^{2}\phi^{2}\chi^{2}$$
 (Chaotic)
$$V(\phi, \chi) = \frac{1}{2}\mu^{2}\phi^{2} + \frac{\lambda}{4}(\chi^{2} - v^{2})^{2} + \frac{1}{2}g^{2}\phi^{2}\chi^{2}$$
 (Hybrid)

 $\begin{cases} \ddot{\phi}(t) + 3H\dot{\phi} + V'(\phi) = 0 \quad (\text{Inflaton Zero-Mode}: \text{Damped Oscillator}) \\ \Box \phi_k + F(\int dq \phi_q \chi_{|k-q|})\phi_k + \dots = 0 \quad (\text{Inflaton Fluctuations}) \\ \Box \chi_k + F(\int dq \chi_q, \phi_{|k-q|})\chi_k + \dots = 0 \quad (\text{Matter Fluctuations}) \end{cases}$

DYNAMICS:

Non-Linear, Non-Perturbative and Far-From-Equilibrium

Reheating (Hybrid Scenarios): SPINODAL INSTABILITY

$$\ddot{\phi}(t) + (\mu^2 + g^2 |\chi|^2) \phi(t) = 0$$

$$\ddot{\chi}_k + (k^2 + m^2 \left(\frac{\phi^2}{\phi_c^2} - 1\right) + \lambda |\chi|^2) \chi_k = 0$$

$$\left. \begin{pmatrix} k < m = \sqrt{\lambda}v \end{pmatrix} \right.$$

$$\chi_k, n_k \sim e^{\sqrt{m^2 - k^2}t}$$



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Scalar Reheating: Observable Phenomena

PHENOMENA FROM (SCALAR) REHEATING:

SUB-HORIZON GRAVITATIONAL WAVES

SUPER-HORIZON GRAVITATIONAL WAVES

EXPECTATIONs of (p)REHEATING: SubH-GW

Physics of (p)REHEATING: $\ddot{\varphi}_k + \omega^2(k,t)\varphi_k = 0$

 $\begin{cases} \text{Hybrid Preheating}: \quad \omega^2 = k^2 + m^2(1 - V t) < 0 \quad (\text{Tachyonic}) \\ \text{Chaotic Preheating}: \quad \omega^2 = k^2 + \Phi^2(t) \sin^2 \mu t \quad (\text{Periodic}) \end{cases}$

At
$$\mathbf{k}_i$$
: $\varphi_{k_i}, n_{k_i} \sim e^{\mu(k,t)t} \Rightarrow$ Inhomogeneities:
$$\begin{cases} L_i \sim 1/k_i \\ \delta \rho / \rho \gtrsim 1 \\ v \approx c \end{cases}$$

(p)REHEATING: VERY EFFECTIVE GW GENERATOR

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GW from RELATIVISTIC WAVES of MATTER





G^a-Bellido et al '02 (Hybrid Scenario: $\lambda \approx 0.1$, $VEV = 10^{-3}M_p$)

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• Scalars $(n_k \gg 1)$: $\Box \phi + V_{,\phi} = 0, \ \Box \chi_a + V_{,\chi_a} = 0$

Semi-classical regime $\pi_k \approx \kappa \phi_k + \dots$ (Squeezed States)

• FRW:
$$H^2 = \frac{8\pi G}{3}\rho$$
, $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p)$, $\begin{cases} \rho = \langle \rho_{\phi} + \rho_{\chi} + ... \rangle \\ p = \langle p_{\phi} + p_{\chi} + ... \rangle \end{cases}$

• GW:
$$h_{ij}'' + 2\mathcal{H}h_{ij}' - \nabla^2 h_{ij} = 16\pi G \Pi_{ij}^{\text{TT}}, \quad \Pi_{ij} = T_{ij} - \langle T_{ij} \rangle_{_{\text{FRW}}}$$

$$ds^2 = a^2 (-d\eta^2 + (\delta_{ij} + h_{ij}) dx^i dx^j), \quad \text{TT} : \begin{cases} h_{ii} = 0 \\ h_{ij,j} = 0 \end{cases}$$

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TT dof carry energy out of the source!!!

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$$g^2 = 2\lambda = 0.25, v = 10^{-3}M_p, V_I = 0.024$$

Bubble Nucleation and Collisions (Animation by Alfonso Sastre)



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3 stages: Exp. Instabilities \rightarrow Bubble Collisions \rightarrow Turbulence



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Today's Signal (GW RedShifted)



$$\Omega_{
m GW}(f) \propto (v/M_p)^2\,, ~~f_* \propto \lambda^{1/4} imes k_*$$

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SUPER-HORIZON GRAVITATIONAL WAVES



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(k η_* < 1, Super-Horizon Scales)

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$$\begin{array}{c} \phi_{a}(\mathbf{k},\eta) & \longrightarrow & T_{\mu\nu}(\phi_{a}) & \longrightarrow & \Pi^{TT}_{\mu\nu}(\phi_{a}) & \longrightarrow & \Box h_{\mu\nu} = 16\pi \mathrm{G} \ \Pi^{TT}_{\mu\nu} \\ \hline \\ FIELD & STRESS & ANISOTROPIC (TT) & GW EQUATIONS \\ FLUCTUATIONS & tensor & STRESS tensor & (TT metric perturb.) \\ \hline \end{array}$$

$$\rho_{\rm GW} = \frac{\langle \dot{\mathbf{h}}_{\mu\nu} \dot{\mathbf{h}}^{\mu\nu} \rangle}{16\pi \mathrm{G}} = \int \!\! \frac{d\rho_{\rm GW}(k,\eta)}{d\log k} d\log k \implies \Omega_{\rm GW}(k,\eta) \equiv \frac{1}{\rho_c} \frac{d\rho_{\rm GW}(k,\eta)}{d\log k}$$

TECHNICALLY

$$\left\langle \phi_{a}(\mathbf{k},\eta) \phi_{a}(\mathbf{k},\eta) \right\rangle \longrightarrow \left\langle \Pi_{\mu\nu}^{\mathrm{TT}}(\phi_{a}) \Pi_{\mu\nu}^{\mathrm{TT}}(\phi_{a}) \right\rangle \longrightarrow \left\langle \dot{\mathbf{h}}_{\mu\nu} \dot{\mathbf{h}}_{\mu\nu} \right\rangle$$

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Jones-Smith et al, 2008

Fenu, DGF, Durrer, Garcia-Bellido 2009

Aftermath of Hybrid Reheating: Scale Inv SubH GW





GW are Early Universe Ideal Probe: decoupled upon production ⇒ spectral signature retained till today ⇒ GWB: "photo" of very Early Universe

④ GW from Reheating: Form, freq. Peak and Amplitude → Specific Model of Inflation (Disadvantage/Advantage)

Scalar Reheating Models: GW (high amplitude, too high frequency)

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- Scalar Reheating Models: GW (high amplitude, too high frequency)

Summary: Scale-Inv GW from Global PhT.

NLSM + large-N limit: Self-Ordering Scalar Fields after SSB (Global PhT).

②
$$k\eta_* \ll 1 \rightarrow k\eta \gg 1$$
: $\Omega_{GW}(k,\eta) = const.$
Observable at LIGO, LISA, BBO,...

• For $VEV = M_I$, then $\Omega_{GW} / \Omega_{GW}^{inf} \sim \mathcal{O}(10) - \mathcal{O}(100)$ Scale-Inv GW is not any more a smoking gun of inflation.

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NON-GAUSSIANITY AFTER GLOBAL PhT

A. Jaffe '93:
$$\delta(\mathbf{x}, \eta) = \frac{2\pi G}{5} \eta^2 \int d\eta' \, \partial_i T_{0i}(\mathbf{x}, \eta') , \begin{pmatrix} \mathsf{MATTER} \\ \mathsf{DOMINATION} \end{pmatrix}$$

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NON-GAUSSIANITY AFTER GLOBAL PhT

$$\delta(\mathbf{x},\eta) = \frac{2\pi G}{5}\eta^2 \int d\eta' \,\partial_i(\partial_0 \phi^a)(\partial_i \phi^a)$$

$$\langle \delta(\mathbf{k})\delta(\mathbf{k}')\rangle = (2\pi)^3 \,\delta_D(\mathbf{k} + \mathbf{k}') P^{\sigma}(k), \qquad P^{\sigma}(k,\eta) \equiv \frac{C^2 \eta^4}{A^2} \bigotimes_{\mathbf{N}} q_{2},$$
2-Point Correlator
$$g_2 \equiv \int \frac{d^3 v}{(2\pi)^3} \left[\mathcal{I}(v, |\hat{\mathbf{k}} - \mathbf{v}|)\right]^2 (\hat{\mathbf{k}} \cdot \mathbf{v}) \left[2(\hat{\mathbf{k}} \cdot \mathbf{v}) - 1\right]$$

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NON-GAUSSIANITY AFTER GLOBAL PhT

$$\delta(\mathbf{x},\eta) = \frac{2\pi G}{5}\eta^2 \int d\eta' \,\partial_i T_{0i}(\mathbf{x},\eta'),$$

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P<sup>\sigma(k,\eta) \equiv \frac{C^2 \eta^4}{A^2} \begin{minipage}{2} & & \\ \hline P^{\sigma}(k,\eta) & = \frac{C^2 \eta^4}{A^2} \begin{minipage}{2} & & \\ \hline P^{\sigma}(k,\eta) & = \frac{C^2 \eta^4}{A^2} \end{minipage} \end{minipage}
2-Point Correlator
Power Spectrum</sup>

$$\begin{array}{c} \begin{array}{c} \text{Physical} \\ \hline \textbf{Constraints} \end{array} & \overbrace{N^{1/4}}^{v} = \left(\frac{p_{\sigma} A^2 \Delta_{\mathcal{R}}^2}{8G_{\text{sw}}^2 g_2} \right)^{1/4} M_{\text{Pl}} \lesssim \frac{M_{\text{Pl}}}{2000} \quad \left(\begin{array}{c} G_{\text{sw}} = 10\\ p_{\sigma} = 0.1 \end{array} \right) \end{array}$$

NON-GAUSSIANITY AFTER GLOBAL PhT

 $\langle \delta(\mathbf{k}_1) \delta(\mathbf{k}_2) \delta(\mathbf{k}_3)
angle = (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B(k_1,k_2,k_3)$

(3-Point Correlator)

(Bispectrum)

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NON-GAUSSIANITY AFTER GLOBAL PhT







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NON-GAUSSIANITY AFTER GLOBAL PhT





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Turbulence

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Results for Hybrid Reheating: Turbulence

Turbulece: $var(\phi) \sim t^{-2p}$, $n(k,t) = t^{-\gamma p} n_o(t^{-p}k)$ (M,T '04)



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Lattice Techniques



Lattice Simulations: Numerics

$$\partial_{\mu}O(x) \rightarrow (O(x+\mu) - O(x-\mu))/2a_{\mu}$$

 $\partial_{\mu}\partial_{\mu}O(x) \rightarrow (O(x+2\mu) + O(x-2\mu) - 2O(x))/4a_{\mu}^{2}$



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Scalar Source (Configuration Space):

 $\ddot{h}_{ij}(\mathbf{x},t) + 3H\dot{h}_{ij}(\mathbf{x},t) - \frac{1}{a^2}\nabla^2 h_{ij}(\mathbf{x},t) = \frac{16\pi}{a^2} \operatorname{TT}\left\{\nabla_l \phi^a \nabla_m \phi^a\right\}(\mathbf{x},t)$

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 $\ddot{h}_{ij}(\mathbf{x},t) + 3H\dot{h}_{ij}(\mathbf{x},t) - \frac{1}{a^2}\nabla^2 h_{ij}(\mathbf{x},t) = \frac{16\pi}{a^2} \operatorname{TT}\left\{\nabla_l \phi^a \nabla_m \phi^a\right\}(\mathbf{x},t)$

Scalar Source (Fourier):

$$\begin{split} \ddot{h}_{ij}(\mathbf{k},t) + 3H\dot{h}_{ij}(\mathbf{k},t) + \frac{k^2}{a^2}h_{ij}(\mathbf{k},t) &= 16\pi \Lambda_{ij,lm}(\hat{\mathbf{k}}) \left\{ \nabla_l \phi^a \nabla_m \phi^a \right\} (\mathbf{k},t) \\ \Lambda_{ij,lm} &= P_{il}P_{jm} - \frac{1}{2}P_{ij}P_{lm}, \quad P_{ij} = \delta_{ij} - k_i k_j / k^2 \end{split}$$

Scalar Source (Configuration Space):

 $\ddot{h}_{ij}(\mathbf{x},t) + 3H\dot{h}_{ij}(\mathbf{x},t) - \frac{1}{a^2}\nabla^2 h_{ij}(\mathbf{x},t) = \frac{16\pi}{a^2} \operatorname{TT}\left\{\nabla_l \phi^a \nabla_m \phi^a\right\}(\mathbf{x},t)$

Scalar Source (Fourier):

$$\begin{split} \ddot{h}_{ij}(\mathbf{k},t) + 3H\dot{h}_{ij}(\mathbf{k},t) + \frac{k^2}{a^2}h_{ij}(\mathbf{k},t) &= 16\pi \Lambda_{ij,lm}(\hat{\mathbf{k}}) \left\{ \nabla_l \phi^a \nabla_m \phi^a \right\} (\mathbf{k},t) \\ \Lambda_{ij,lm} &= P_{il}P_{jm} - \frac{1}{2}P_{ij}P_{lm}, \quad P_{ij} = \delta_{ij} - k_i k_j / k^2 \end{split}$$

Solution: $(h_{ij}(t_0) = \dot{h}_{ij}(t_0) = 0)$ $h_{ij}(\mathbf{k}, t) = \Lambda_{ij,lm}(\hat{\mathbf{k}}) \int_{t_0}^t dt' G(t - t') \Pi_{lm}^{\text{eff}}(\mathbf{k}, t'), \qquad \Pi_{lm}^{\text{eff}} = \nabla_l \phi \, \nabla_m \phi$

Building the Solution:

1) Non-Physical eq.: $\ddot{u}_{ij}(\mathbf{x},t) + 3H\dot{u}_{ij}(\mathbf{x},t) - \frac{\nabla^2}{a^2}u_{ij}(\mathbf{x},t) = 16\pi \{\phi^a,_i \phi^a,_j\}(\mathbf{x},t)$

2) Fourier transform: $u_{ij}(\mathbf{x},t) \rightarrow u_{ij}(\mathbf{k},t)$

3) Proyection: $h_{ij}(\mathbf{k},t) = \Lambda_{ij,lm}(\hat{\mathbf{k}})u_{lm}(\mathbf{k},t)$

Outputs:
$$\rho_{GW} = \frac{1}{32\pi G} \frac{1}{L^3} \int d^3 \mathbf{x} \, \dot{h}_{ij} \dot{h}_{ij} = \frac{1}{32\pi G} \frac{1}{L^3} \int d^3 \mathbf{k} |\dot{h}_{ij}(t, \mathbf{k})|^2$$

1) Total GW density: $\rho_{GW} = \frac{1}{32\pi GL^3} \times \int k^2 dk \int d\Omega \Lambda_{ij,lm}(\hat{\mathbf{k}}) \dot{u}_{ij}(t,\mathbf{k}) \dot{u}_{lm}^*(t,\mathbf{k})$

2) Spectrum:
$$\frac{d\rho}{d\log k} = \frac{1}{8GL^3} k^3 \left\langle \Lambda_{ij,lm}(\hat{\mathbf{k}}) \dot{u}_{ij}(t,\mathbf{k}) \dot{u}_{lm}^*(t,\mathbf{k}) \right\rangle_{4\pi}$$

3) Snapshots: $h_{ij}(t, \mathbf{x}) = (2\pi)^{-3/2} \int d^3k e^{-i\mathbf{k}\mathbf{x}} \Lambda_{ij,lm}(\hat{\mathbf{k}}) u_{lm}(t, \mathbf{k})$