Operational interpretation of quantum discord

Leandro Aolita (ICFO, Barcelona)

In collaboration with D. Cavalcanti and K. Modi (Singapore), S. Boixo (CalTech), M. Piani (Waterloo), and A. Winter (Singapore & Bristol)





Quantum Discord (QD)

Quantum Discord (QD)

• was introduced [Zurek (00), Ollivier & Zurek (01)] to quantify all quantum correlations.

• Since its definition, it has received LOTS of attention. (36 arXiv titles over the last two years.)

• Interpretations in terms of the gain (in work extraction) a Maxwell's demon obtains when operating quantumly with respect to classically have been provided. See Aharon's talk!!!

• However, astonishingly, up to now QD lacked an (information-theoretic) operational interpretation.

• Quantum information community not happy about this :-(

• Curiosity: discord is an asymmetric correlation quantifier...

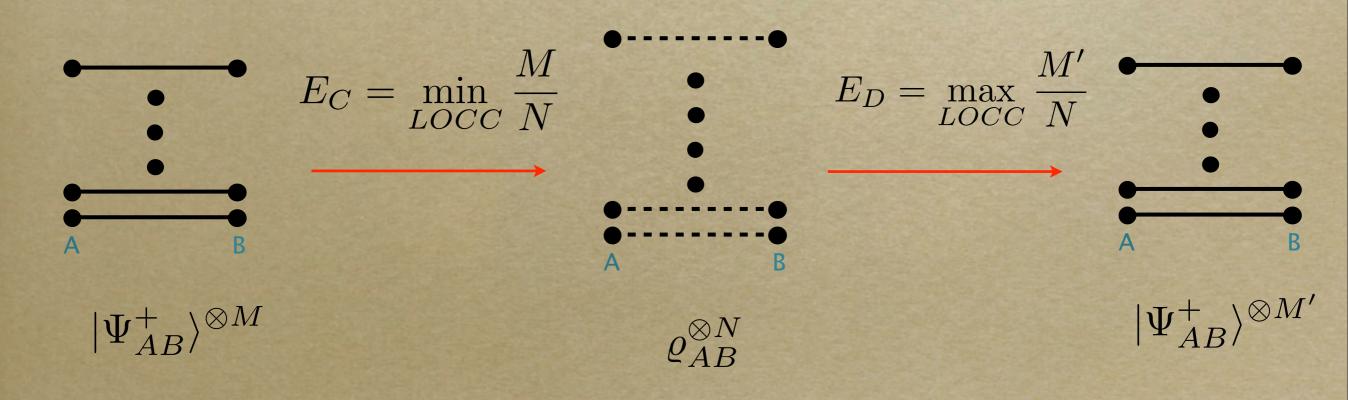
Information-theoretic quantities in the spirit of Shannon: asymptoticly many copies.

We say that a quantifier of correlations has an operational meaning if it measures the performance or efficiency of a given physical (information processing) protocol.

Information-theoretic quantities in the spirit of Shannon: asymptoticly many copies.

We say that a quantifier of correlations has an operational meaning if it measures the performance or efficiency of a given physical (information processing) protocol.

• Paradigmatic examples:



QD has a clear operational interpretation: it quantifies the total singlet consumption in state merging!!!!

The intrinsic asymmetry in QD plays a natural role in this scenario!!!

Discord imbalances quantify the efficiencies in different strategies of state merging and dense coding!!!

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

See also [V. Madhok & A. Datta, arXiv:1008.4135] for related results!!!

Outline of the talk

and a second for water me and instant a second to the second with the

JAMANA A. MANARL

Outline of the talk

- Conditional entropy and coherent information.
- Definition of QD.
- State merging and its total entanglement consumption.
- Operational interpretation of QD.
- Asymmetry of QD.
- QD, state merging and the quantum advantage of dense coding.
- Asymptotic regularization and concluding remarks.

Quantum conditional entropy and coherent information

Low did to the table to

the tor was some second to the

78422

stand & Advent

Quantum conditional entropy and coherent information

• The classical (Shannon) entropy measures the (average) uncertainty in the value of a classical random variable a:

$$H(a) \equiv H(\{p_i^a\}) := -\sum_i p_i^a \log_2 p_i^a$$

Quantum conditional entropy and coherent information

• The classical (Shannon) entropy measures the (average) uncertainty in the value of a classical random variable a:

$$H(a) \equiv H(\{p_i^a\}) := -\sum_i p_i^a \log_2 p_i^a$$

• The von Neumman entropy is the quantum counterpart: $S(\rho) := -\text{Tr}[\rho \log_2 \rho]$

• Notation (for the reduced state of part X): $S(X) = S(\varrho_X)$



• The classical conditional entropy measures the uncertainty left - on average - for the value of a given that the value of b has been discovered:

$$H(a|b) := H(a,b) - H(b)$$

• Classical info theory: H(a|b) is the average amount of (partial) classical information that A must give to B (who already knows the value of b) so that the latter gains full knowledge of (a,b) [Slepian & Wolf (71)].

• Given this interpretation, H(a|b) is of course non-negative. And, in fact, it can also be expressed as

$$H(a|b) = \sum_{j} p_j^b H(a|b=j),$$

where H(a|b=j) is the entropy of the conditional probability $p_{i|b=j}^{a} := p_{ij}^{ab}/p_{j}^{b}$



• The quantum conditional entropy is defined analogously:

S(A|B) := S(AB) - S(B),

but, in contrast, it can take negative values!!!

• The quantum conditional entropy is defined analogously:

$$S(A|B) := S(AB) - S(B),$$

but, in contrast, it can take negative values!!!

• The possible negativity was for a long time a hard obstacle to an operational interpretation for S(A|B)

• As a matter of fact, its opposite was even given a name of its own. The coherent information $I(A \ge B) \coloneqq -S(A | B)$.

• Originally introduced in quantum info as purely-quantum quantity to measure the amount of quantum info conveyable by a quantum channel [Schumacher & Nielsen (96)].

Quantum Discord

And the state of the second of

Quantum Discord

• A "remedy" to negative quantum conditional entropy is [Henderson & Vedral (01), Ollivier & Zurek (01)]:

$$S(A|B_c) := \min_{\{N_j\}} \sum_{j} p_j^B S(A|B=j),$$

over a positive rank-1 decomposition of the identity:

 $\sum_{j} N_j = \mathbf{1}_B.$

Quantum Discord

• A "remedy" to negative quantum conditional entropy is [Henderson & Vedral (01), Ollivier & Zurek (01)]:

$$S(A|B_c) := \min_{\{N_j\}} \sum_{j} p_j^B S(A|B=j),$$

 $\sum_{j} N_j = \mathbf{1}_B.$

over a positive rank-1 decomposition of the identity:

• The quantum discord of AB with measurements on B [Ollivier & Zurek (01)]:

$$D(A|B) := S(A|B_c) - S(A|B).$$

Quantum Discord

• A "remedy" to negative quantum conditional entropy is [Henderson & Vedral (01), Ollivier & Zurek (01)]:

$$S(A|B_c) := \min_{\{N_j\}} \sum_{j} p_j^B S(A|B=j),$$

 $\sum_{j} N_j = \mathbf{1}_B.$

over a positive rank-1 decomposition of the identity:

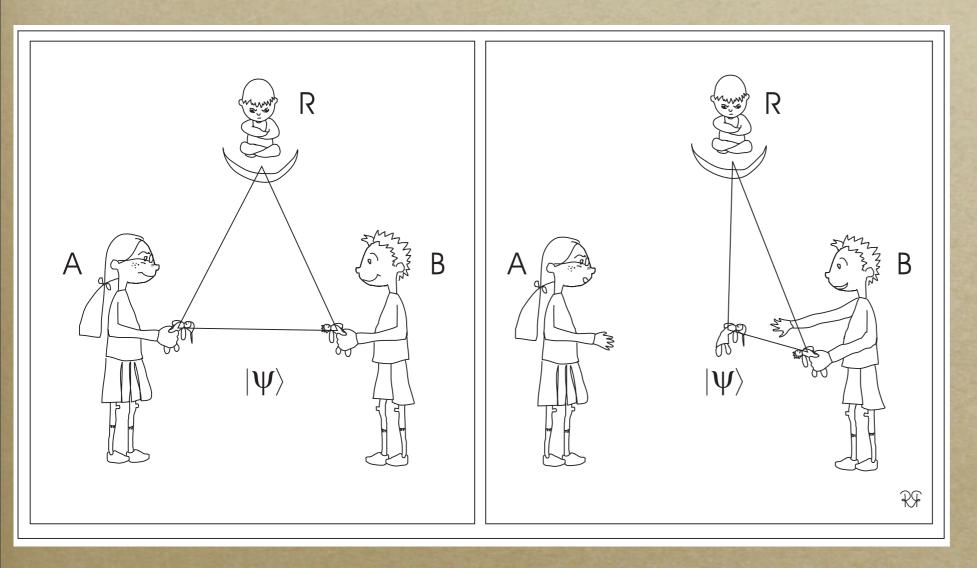
• The quantum discord of AB with measurements on B [Ollivier & Zurek (01)]:

$$D(A|B) := S(A|B_c) - S(A|B).$$

[A Ferraro, L. Aolita, D. Cavalcanti, F. M. Cucchietti, & A. Acín, PRA (2010)]

State merging and entanglement consumption

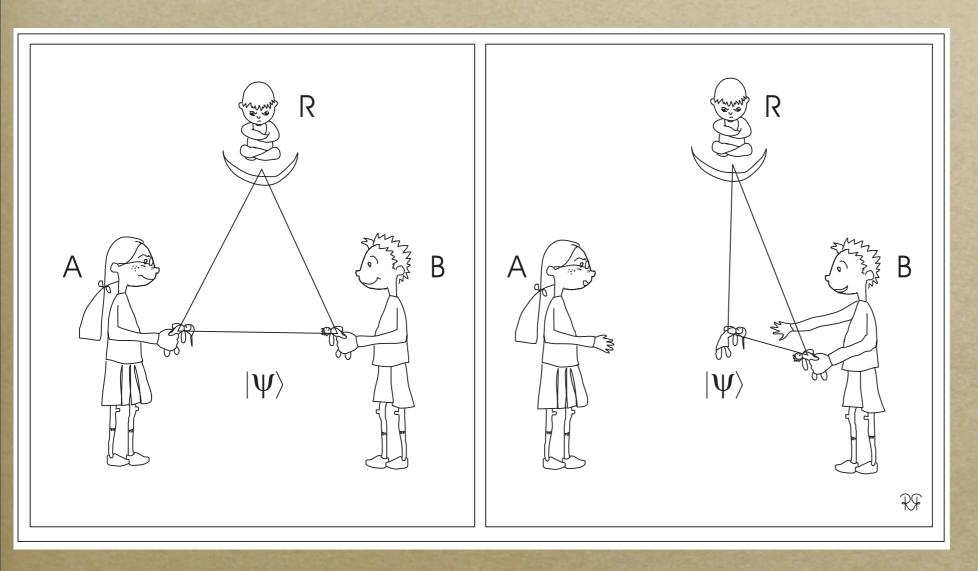
• A satisfactory operational interpretation for S(A|B) - and therefore also $I(A \ge B)$ - was found in the context of state merging:



[Horodecki, Oppenheim & Winter, Nature (05)]

State merging and entanglement consumption

• A satisfactory operational interpretation for S(A|B) - and therefore also $I(A \ge B)$ - was found in the context of state merging:

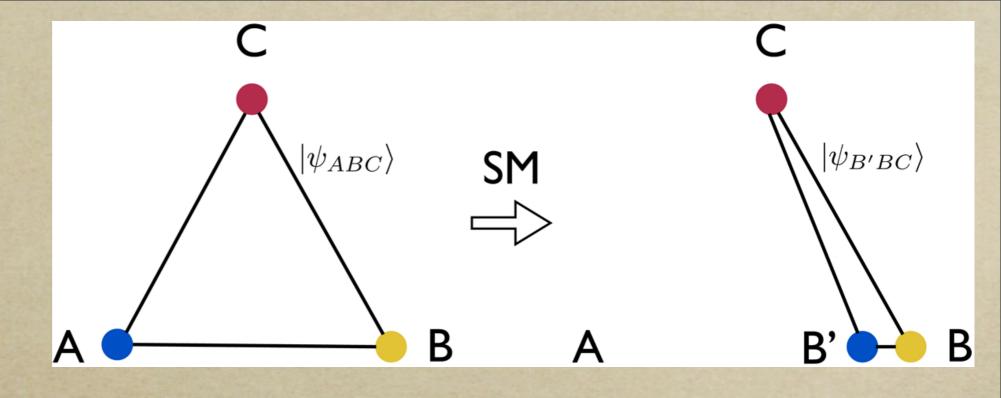


• S(A|B) quantifies exactly the optimal amount of uses of a perfect quantum channel!!!

• In a sense this is similar to what happened with H(a|b)...

[Horodecki, Oppenheim & Winter, Nature (05)]

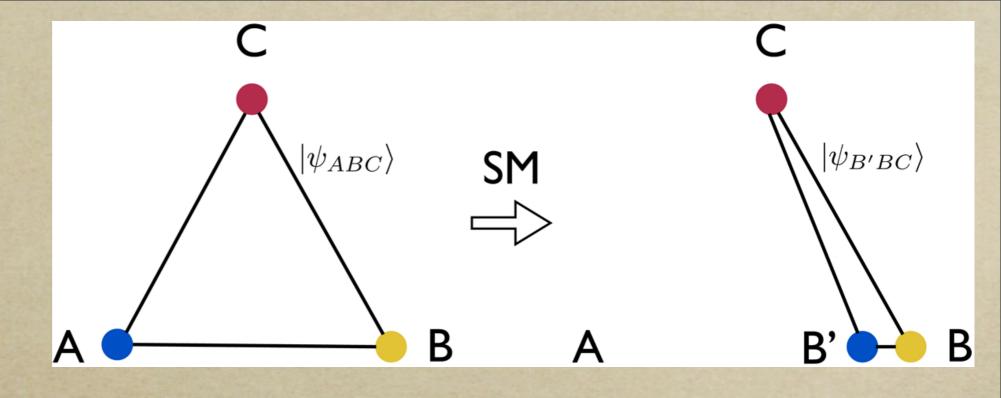
More technically,



• A and B know the state they have, QAB.

- C is a neutral (inactive) reference system.
- LOCCs are for free, but quantum channels (singlets!) are expensive.
- Acting on N copies of the state, their goal is to end up with $|\psi_{B'BC}\rangle^{\otimes N}$, such that $|\psi_{B'BC}\rangle \rightarrow |\psi_{ABC}\rangle$, for $N \rightarrow \infty$.

More technically,

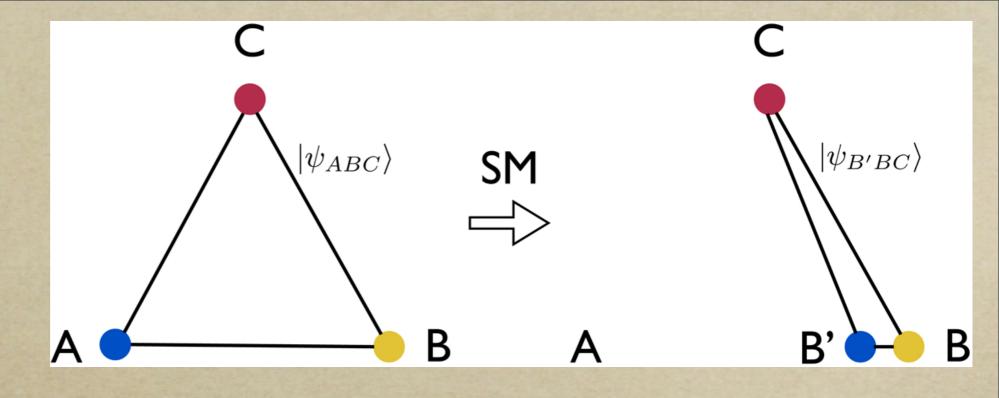


• A and B know the state they have, QAB.

- C is a neutral (inactive) reference system.
- LOCCs are for free, but quantum channels (singlets!) are expensive.
- Acting on N copies of the state, their goal is to end up with $|\psi_{B'BC}\rangle^{\otimes N}$, such that $|\psi_{B'BC}\rangle \rightarrow |\psi_{ABC}\rangle$, for $N \rightarrow \infty$.

• For $S(A|B) \ge 0$, A and B consume S(A|B) extra singlets (per copy of the state) and end up fully uncorrelated.

More technically,



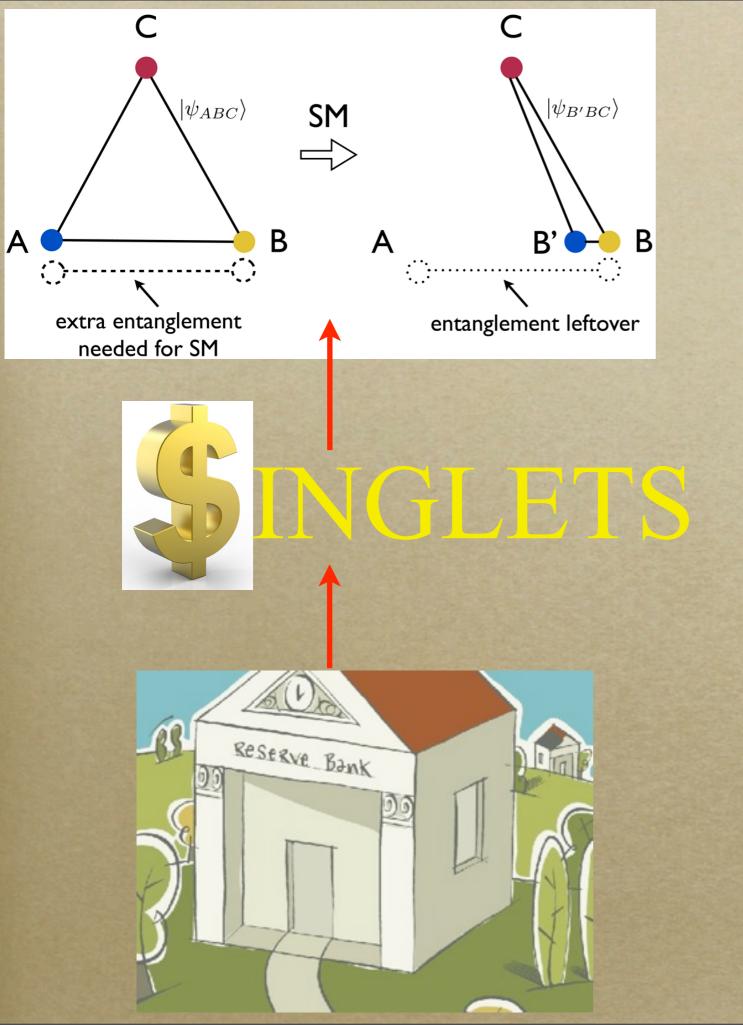
• A and B know the state they have, QAB.

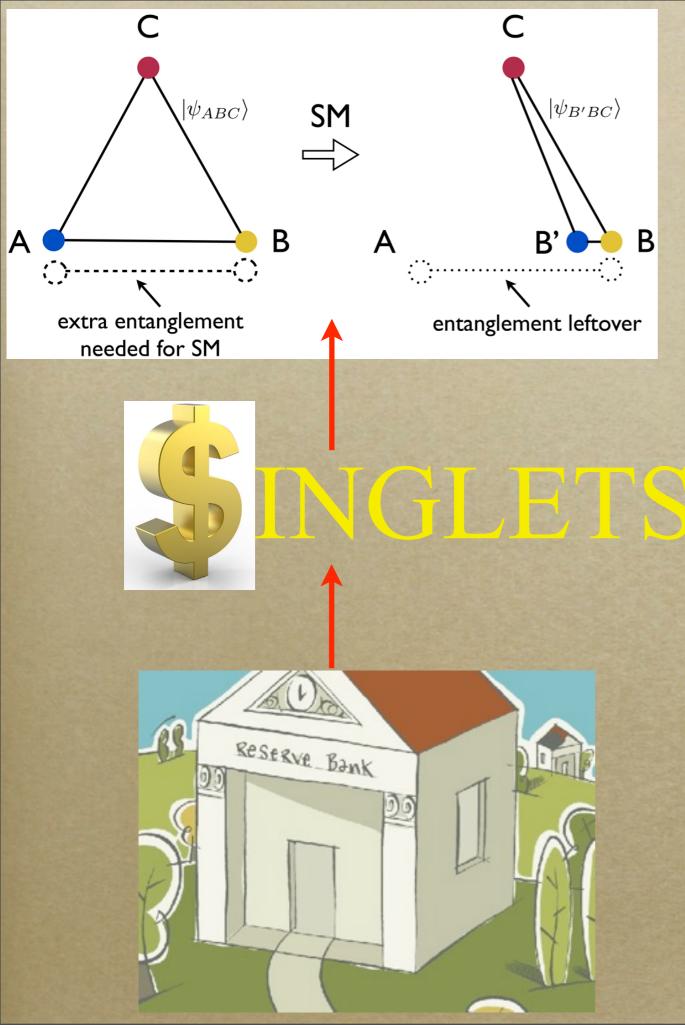
- C is a neutral (inactive) reference system.
- LOCCs are for free, but quantum channels (singlets!) are expensive.
- Acting on N copies of the state, their goal is to end up with $|\psi_{B'BC}\rangle^{\otimes N}$, such that $|\psi_{B'BC}\rangle \rightarrow |\psi_{ABC}\rangle$, for $N \rightarrow \infty$.

• For $S(A|B) \ge 0$, A and B consume S(A|B) extra singlets (per copy of the state) and end up fully uncorrelated.

• For S(A|B) < 0, not only do they perform the SM for free but they also retain -S(A|B) = I(A > B) singlets (per copy of the state), which they can use for future merings.

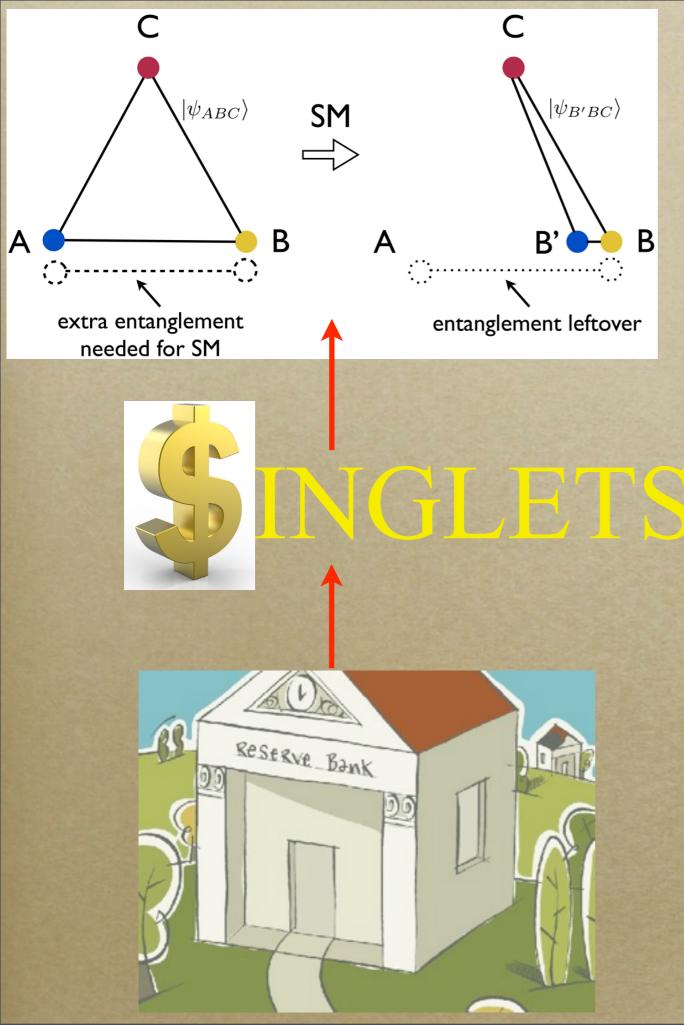






• A and B have a joint checking account in a entanglement bank from where they can borrow singlets.

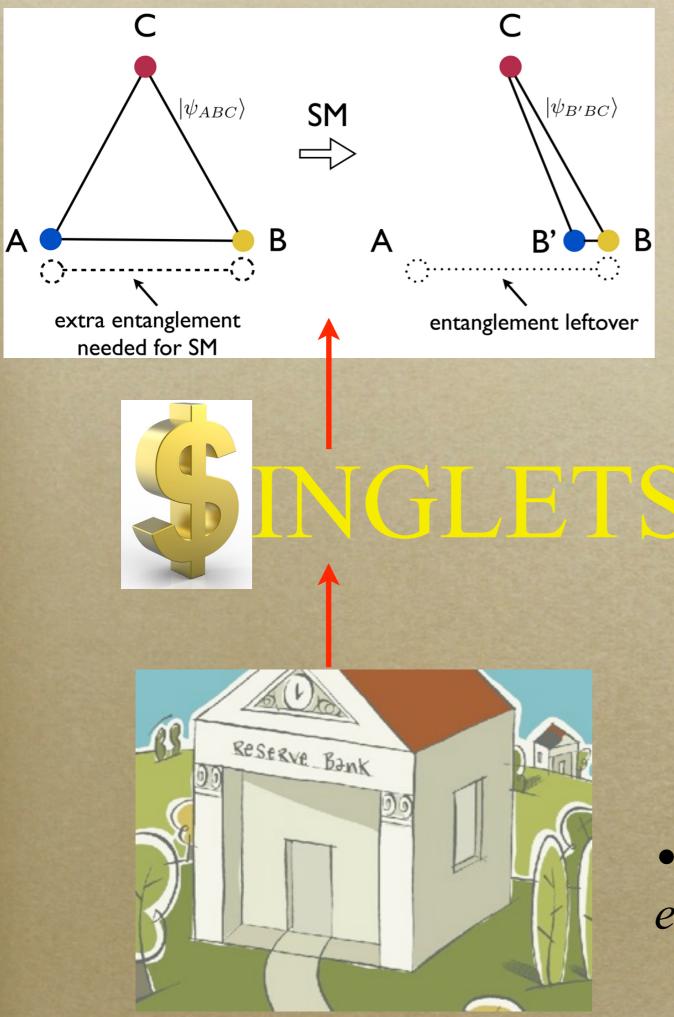
• *Then*, S(A|B) = DEBIT - CREDIT.



• A and B have a joint checking account in a entanglement bank from where they can borrow singlets.

• Then, S(A|B) = DEBIT - CREDIT.

• The balance of the SM operation is given precisely by $I(A \ge B)$.



• A and B have a joint checking account in a entanglement bank from where they can borrow singlets.

• Then, S(A|B) = DEBIT - CREDIT.

• The balance of the SM operation is given precisely by $I(A \ge B)$.

• But what about the initial entanglement between A and B????

• The initial entanglement between A and B is completely lost.

• Therefore, the total entanglement consumption during the process of SM has to take this initial amount into account:

TOTAL CONSUMPTION = INITIAL BALANCE + DEBIT - CREDIT.

• The initial entanglement between A and B is completely lost.

• Therefore, the total entanglement consumption during the process of SM has to take this initial amount into account:

TOTAL CONSUMPTION = INITIAL BALANCE + DEBIT - CREDIT.

 $\Gamma(A\rangle B) := E_F(A:B) + S(A|B),$

with $E_F(A:B) := \min_{\{p_i,\psi_i^{AB}\}} \sum_i p_i S(\operatorname{Tr}_A[\psi_i^{AB}]), \text{ for } \varrho_{AB} = \sum_i p_i |\psi_i^{AB}\rangle \langle \psi_i^{AB}|.$

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

• The initial entanglement between A and B is completely lost.

• Therefore, the total entanglement consumption during the process of SM has to take this initial amount into account:

TOTAL CONSUMPTION = INITIAL BALANCE + DEBIT - CREDIT.

 $\Gamma(A\rangle B) := E_F(A:B) + S(A|B),$

with $E_F(A:B) := \min_{\{p_i,\psi_i^{AB}\}} \sum_i p_i S(\operatorname{Tr}_A[\psi_i^{AB}]), \text{ for } \varrho_{AB} = \sum_i p_i |\psi_i^{AB}\rangle \langle \psi_i^{AB}|.$

• The entanglement of formation quantifies the minimum amount of pure-state entanglement that A and B consume to create (asymptotically many copies of) QAB by LOCC with strategies where each pure-state member of the ensemble is prepared independently.

•Thus, $\Gamma(A \mid B)$ quantifies the total entanglement consumed, by taking into account the amount A and B would have needed to prepare QAB by LOCC - and ``lost" during SM - plus the amount used by the process of SM itself. >>>> EXTENDED STATE MERGING.

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

Operational interpretation of QD

The is the support of the second of the seco

Operational interpretation of QD

Call tor pasts manuscier period

Consider now $D(A|C) := S(A|C_c) - S(A|C).$

Operational interpretation of QD

a ton was some correction

Consider now
$$D(A|C) := S(A|C_c) - S(A|C).$$

• For a pure tripartite state it is: $S(B) = E_F(A:B) + I(B:C_c) := E_F(A:B) + S(B) - S(B|C_c) \Rightarrow E_F(A:B) = S(B|C_c).$ [Koashi & Winter (04)]

Operational interpretation of QD

Consider now
$$D(A|C) := S(A|C_c) - S(A|C).$$

• For a pure tripartite state it is: $S(B) = E_F(A:B) + I(B:C_c) := E_F(A:B) + S(B) - S(B|C_c) \Rightarrow E_F(A:B) = S(B|C_c).$ [Koashi & Winter (04)]

• And also S(A|C) := S(AC) - S(C) = S(B) - S(AB) := -S(B|A).

Operational interpretation of QD

Consider now
$$D(A|C) := S(A|C_c) - S(A|C).$$

• For a pure tripartite state it is: $S(B) = E_F(A:B) + I(B:C_c) := E_F(A:B) + S(B) - S(B|C_c) \Rightarrow E_F(A:B) = S(B|C_c).$ [Koashi & Winter (04)]

• And also S(A|C) := S(AC) - S(C) = S(B) - S(AB) := -S(B|A).

Then $D(A|C) = E_F(A:B) + S(A|B) := \Gamma(A > B)!!!$

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

Operational meaning of the asymmetry

7×422

stand & Adver

Tor supply mining corperations

Operational meaning of the asymmetry

• The asymmetry in discord has bothered many. From the previous result it follows immediately that

$D(A|C) - D(C|A) = \Gamma(A \rangle B) - \Gamma(C \rangle B)$

... now we can understand the asymmetry in terms differences in the cost of ESM!!!

Operational meaning of the asymmetry

• The asymmetry in discord has bothered many. From the previous result it follows immediately that

$D(A|C) - D(C|A) = \Gamma(A \rangle B) - \Gamma(C \rangle B)$

... now we can understand the asymmetry in terms differences in the cost of ESM!!!

$D(A|C) - D(A|B) = \Gamma(A \rangle B) - \Gamma(A \rangle C)$

... for the first time a physical scenario where the values of QD provide concrete quantitative info about which of two possible strategies is most convenient!!!

QD, dense coding, and extended state merging

• DC: by sending her part of QAB, A transmits classical info more efficiently than she could was the system classical [Bennett & Wiesner (92)].

• Conventional (pure-state) DC scenario: each letter in an alphabet is associated to a unitary rotation.

• Then the correction to the classical capacity (rate of information transmission per shared state used) is exactly the coherent information $I(A \ge B)$ [Horodecki et al. (01), Winter (02), Bruss et al. (04)].

QD, dense coding, and extended state merging

• DC: by sending her part of QAB, A transmits classical info more efficiently than she could was the system classical [Bennett & Wiesner (92)].

• Conventional (pure-state) DC scenario: each letter in an alphabet is associated to a unitary rotation.

• Then the correction to the classical capacity (rate of information transmission per shared state used) is exactly the coherent information $I(A \ge B)$ [Horodecki et al. (01), Winter (02), Bruss et al. (04)].

• The most general (mixed-state) DC scenario: A's optimal encoding also consists of unitary rotations, but preceded by a pre-processing general quantum operation $\Lambda_A: M_{d_A} \to M_{d'_A}$.

Then the (single-shot) capacity is $\chi(A \rangle B) := \log_2 d'_A + \max_{\Lambda_A} I(A' \rangle B)$

QD, dense coding, and extended state merging

• DC: by sending her part of QAB, A transmits classical info more efficiently than she could was the system classical [Bennett & Wiesner (92)].

• Conventional (pure-state) DC scenario: each letter in an alphabet is associated to a unitary rotation.

• Then the correction to the classical capacity (rate of information transmission per shared state used) is exactly the coherent information $I(A \ge B)$ [Horodecki et al. (01), Winter (02), Bruss et al. (04)].

• The most general (mixed-state) DC scenario: A's optimal encoding also consists of unitary rotations, but preceded by a pre-processing general quantum operation $\Lambda_A: M_{d_A} \to M_{d'_A}.$

Then the (single-shot) capacity is

$$\chi(A\rangle B) := \log_2 d'_A + \max_{\Lambda_A} I(A'\rangle B)$$

And the quantum advantage [Horodecki & Piani (07)]:

$S(A) = E_P(B:A) + \Delta_{DC}(C \rangle A)$

$$S(A) = E_P(B:A) + \Delta_{DC}(C \rangle A)$$

 $\Rightarrow D(A|C) - D(B|C) = \Gamma(A \rangle C) - \Gamma(A \rangle C) = S(A) - S(B)$

 $S(A) = E_P(B:A) + \Delta_{DC}(C \rangle A)$ $S(B) = E_P(A:B) + \Delta_{DC}(C \rangle B)$

 $\Rightarrow D(A|C) - D(B|C) = \Gamma(A \rangle C) - \Gamma(A \rangle C) = S(A) - S(B)$

$\Rightarrow D(A|C) - D(B|C) = \Delta_{DC}(C \rangle A) - \Delta_{DC}(C \rangle B)!!!$

QD imbalance quantifies how much more efficient it is for C to do DC toward A as compared to toward B!!!!!

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

 $S(A) = E_P(B:A) + \Delta_{DC}(C \rangle A)$ $S(B) = E_P(A:B) + \Delta_{DC}(C \rangle B)$

 $\Rightarrow D(A|C) - D(B|C) = \Gamma(A \rangle C) - \Gamma(A \rangle C) = S(A) - S(B)$

$\Rightarrow D(A|C) - D(B|C) = \Delta_{DC}(C \rangle A) - \Delta_{DC}(C \rangle B)!!!$

QD imbalance quantifies how much more efficient it is for C to do DC toward A as compared to toward B!!!!!

• And in fact, if C sends always the same subsystem then it is:

$$\Rightarrow D(A|C) - D(B|C) = \chi_{DC}(C \rangle A) - \chi_{DC}(C \rangle B)!!!$$

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

Asymptotic regularizations

Low and the Washington War You King the King the

stand & - Advent

Attende Tor wast's me we corperated in at

Asymptotic regularizations

• All the relations we have found can be recast in their regularized version:

 $E_C(A:B) = \lim_{N \to \infty} \frac{1}{N} E_F(A:B)_{\rho_{AB}^{\otimes N}}$ $\Gamma^{\infty}(A \rangle B) := \lim_{n \to \infty} \Gamma(A \rangle B)_{\rho_{AB}^{\otimes n}} / n = E_C(A:B) + S(A|B)$

> Total pure-state entanglement consumption of ESM for general state-creation strategies (with arbitrary LOCC)

Asymptotic regularizations

• All the relations we have found can be recast in their regularized version:

 $E_C(A:B) = \lim_{N \to \infty} \frac{1}{N} E_F(A:B)_{\rho_{AB}^{\otimes N}}$ $\Gamma^{\infty}(A \rangle B) := \lim_{n \to \infty} \Gamma(A \rangle B)_{\rho_{AB}^{\otimes n}} / n = E_C(A:B) + S(A|B)$

> Total pure-state entanglement consumption of ESM for general state-creation strategies (with arbitrary LOCC)

$$D^{\infty}(A|B) = \lim_{N \to \infty} \frac{1}{N} D(A|B)_{\rho_{AB}^{\otimes N}}$$

 $\Rightarrow D^{\infty}(A|B) = \Gamma^{\infty}(A\rangle B)$



Salar a the second to a second to a second to a second to the second tot

and the second

Conclusions

- We introduced extended state merging and its total entanglement consumption.
- QD quantifies the total singlet consumption in ESM.
- The intrinsic asymmetry in QD plays a natural role in this scenario: it tells us which of two possible strategies is cheapest.
- *QD* imbalance (with the measured system as the one in common) quantifies the difference in efficiency gain between DC toward two different receivers.
- These results define for the first time a physical scenario where the values of QD provide concrete quantitative info about the efficiency or cost involved in physical protocols.

[D. Cavalcanti, L. Aolita, S. Boixo, K. Modi, M. Piani, & A. Winter, arXiv:1008.3205]

THANKS!