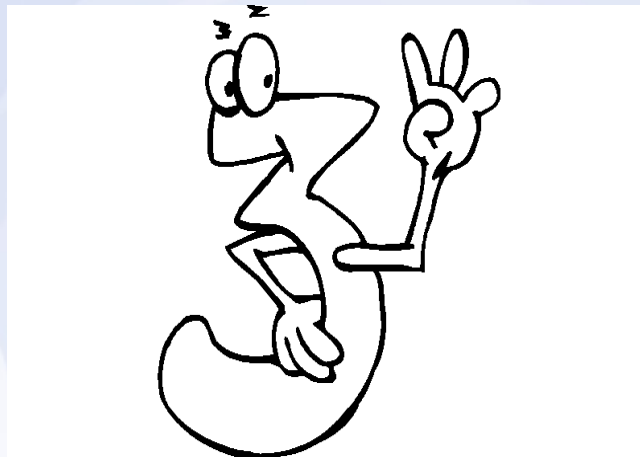


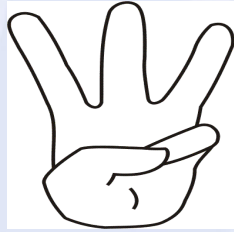
Three-form cosmology

in $< 3^3$ minutes

- Tomi Koivisto: Three-formal dark energy
- Nelson Nunes: Three-inflation and nongaussianity
- Federico Urban: Three-magnetic fields

BENASQUE, 2012: **Modern Cosmology**: **Early Universe**, **CMB** and **LSS**





Three-formalities

- Canonical action

$$S_A = \int d^4x \sqrt{-g} \left(\frac{1}{2\kappa^2} R - \frac{1}{48} F^2 - V(A^2) \right)$$

$$F_{\alpha\beta\gamma\delta} = 4\nabla_{[\alpha} A_{\beta\gamma\delta]}$$

- Always dual to a vector,

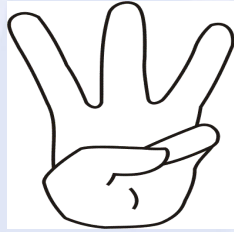
$$\star A_\mu = \frac{1}{6} \epsilon_{\mu\alpha\beta\gamma} A^{\alpha\beta\gamma}$$

in some cases reducible to a scalar:

$$(dA, dA) + m^2(A, A) = (\delta \star A, \delta \star A) + m^2(\star A, \star A) \rightarrow m^2(\phi, \phi) + (d\phi, d\phi)$$

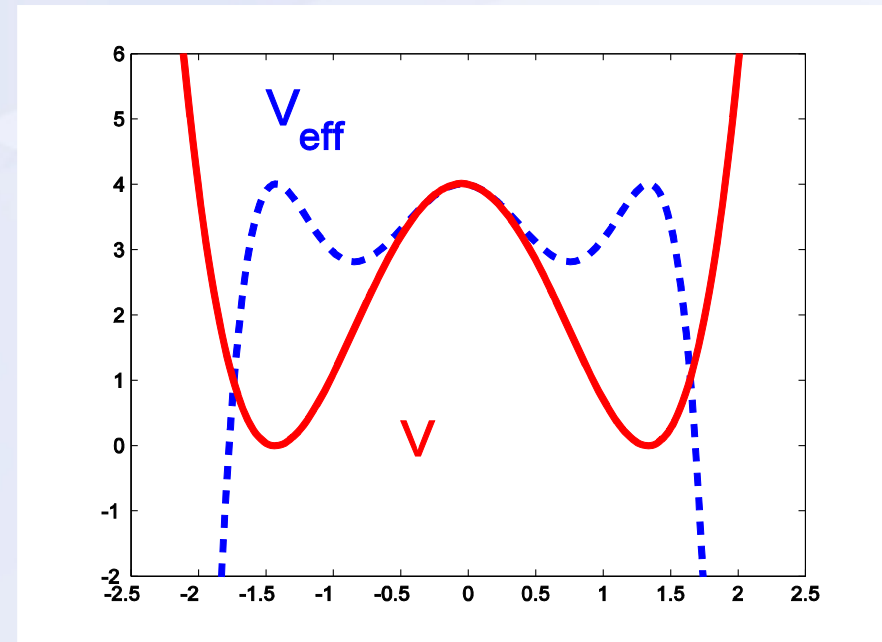
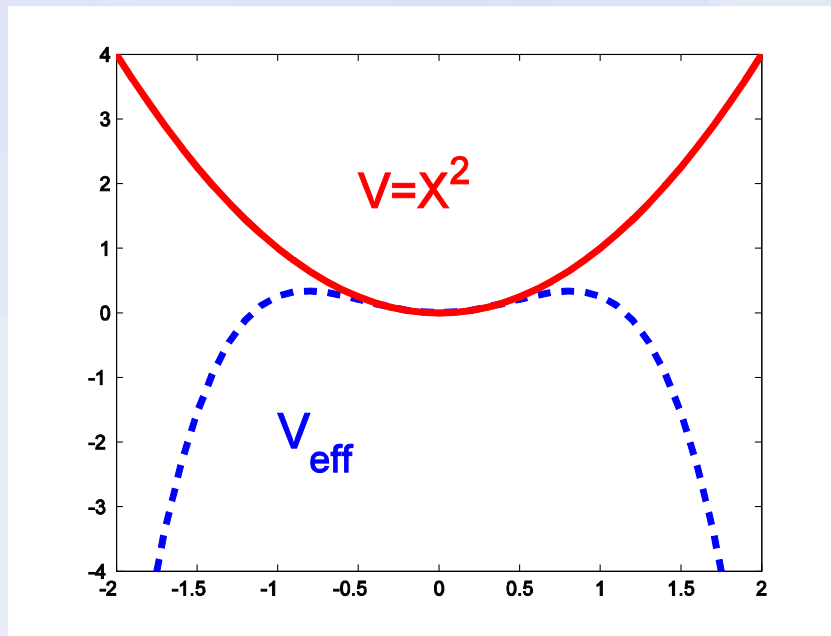
- The simplest case:

$$F^2 = \Lambda!$$

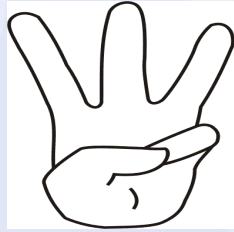


Effective potential

$$V_{effective}(X) = V(X) - \frac{3}{2} \int V'(X) X dX$$



$$V(X^2) = (X^2 - C)^2$$



Three-cosmology

- The FLRW symmetry allows only

$$A = a^3(t) X(t) dx \wedge dy \wedge dz$$

- The field $X(t)$ then behaves as:

$$\rho_X = \frac{1}{2} (\dot{X} + 3HX)^2 + V(X),$$

$$p_X = -\frac{1}{2} (\dot{X} + 3HX)^2 + V'(X)X - V.$$

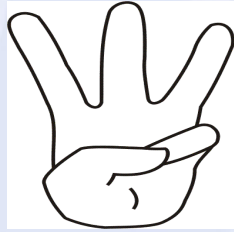
$$w_X = -1 + \frac{V_{,X} X}{\rho_X}$$

$$\ddot{X} = -3H\dot{X} - V_{,X} - 3\dot{H}X,$$

$$c_S^2 = \frac{V_{,XX} X}{V_{,X}}$$

- A critical point: $X^2 = 2/3$

$$H^2 = \frac{1}{3} \frac{V + \rho_B}{1 - \kappa^2 (X' + 3X)^2 / 6}$$

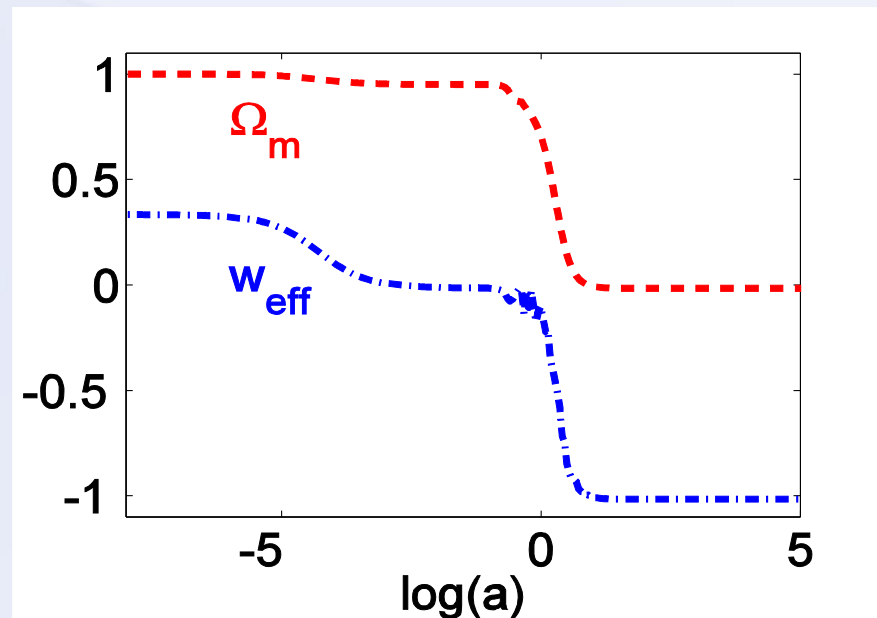
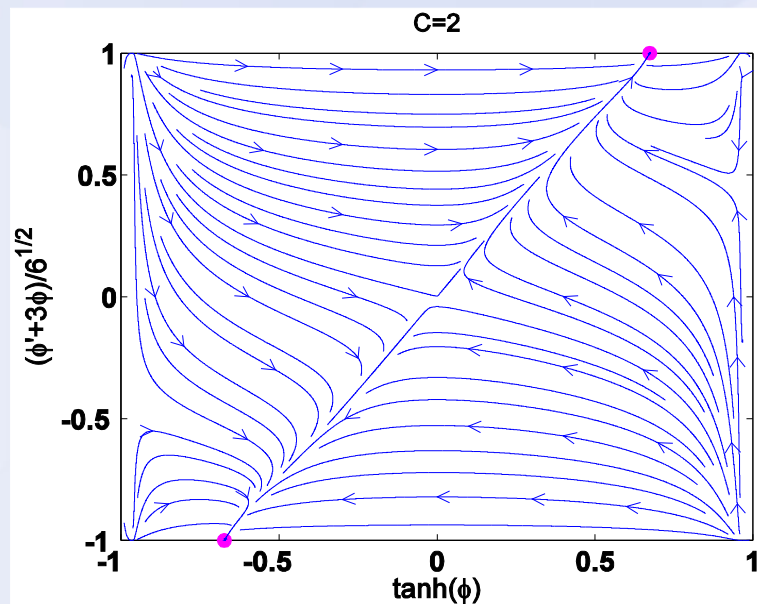


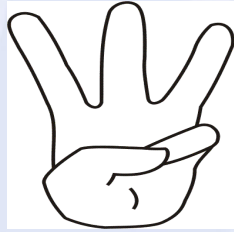
Cosmological dynamics

$$x \equiv \kappa X, \quad y \equiv \frac{\kappa}{\sqrt{6}}(X' + 3X), \quad z^2 = \frac{\kappa^2 V}{3H^2}$$

$$w^2 \equiv \frac{\kappa^2 \rho_B}{3H^2}, \quad \lambda(x) \equiv -\frac{1}{\kappa} \frac{V_{,X}}{V}$$

	x	y	w	\dot{H}/H^2	λ	description
A	0	0	± 1	$-3\gamma/2$	any	matter domination
B_{\pm}	$\pm\sqrt{2/3}$	± 1	0	0	any	maximal point
C	x_{ext}	$\sqrt{3/2} x_{\text{ext}}$	0	0	0	potential extremum





Coupled three-form

- Consider coupling to CDM particles

$$\mathcal{L} = -\frac{1}{48}F^2 - V(A^2) - \Sigma_a m_a(A^2)\delta(x - x(\lambda))\sqrt{\dot{x}^2/g}$$

- Contributes to the background pressure

$$2\dot{H} + 3H^2 = 8\pi G \left[\frac{1}{2} (\dot{X} + 3HX)^2 + V - (V_{,X} + \kappa\rho_m f) X \right] \quad f = \frac{1}{\kappa} \log m^2_{,X}$$

- The matter overdensity δ evolves as:

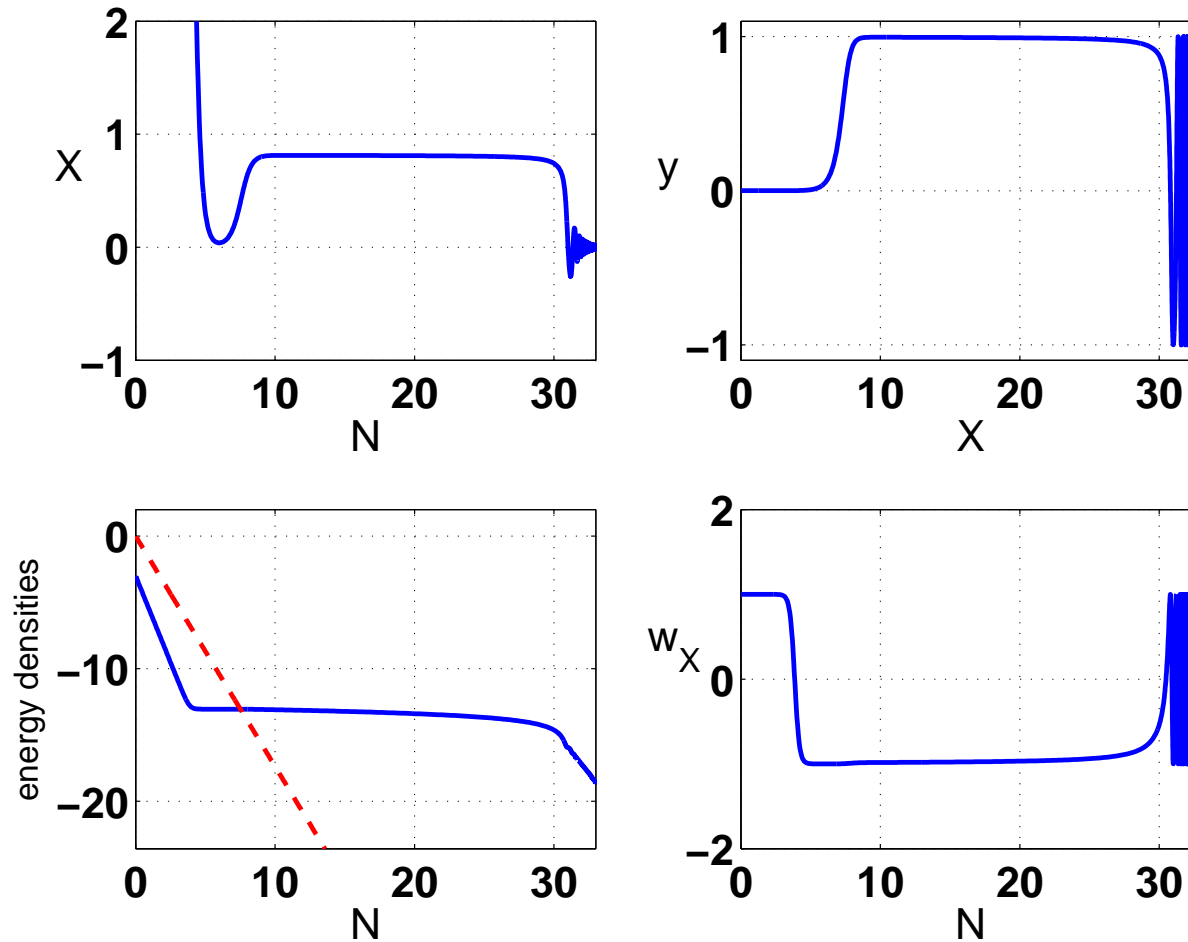
$$\ddot{\delta} + \left[2H + \frac{f\kappa\dot{X} - XF - 2\dot{F}}{1 - F} \right] \dot{\delta} = 4\pi G_{eff} f \rho \delta - \frac{k^2}{a^2} c_{eff}^2 \delta$$

$$\frac{G_{eff}}{G} = \frac{1 + \frac{\dot{F} + (2H + X)\dot{F}}{4\pi G \rho_m}}{1 - F}$$

$$c_{eff}^2 = \frac{F}{1 - F}$$

$$\frac{-f^2 \kappa^2 \rho}{V_{,XX} + f_{,X} \kappa \rho} \delta \equiv F \delta$$

Inflation



$$H^2 = \frac{1}{3} \frac{V + \rho}{1 - (X_{,N} + 3X)^2/6}$$

$$V_{\text{eff},X} = V_{,X} \left(1 - \frac{3}{2} X^2 \right) - \frac{3}{2} \gamma \rho_B X$$

How much inflation?

$$\epsilon \equiv -\frac{\dot{H}}{H^2} \approx \frac{3V_{,X}X}{2V} \left(1 - \frac{3}{2}X^2\right)$$

For $V = X^n$

$$\Delta N \approx \frac{2}{9n} \frac{1}{2/3 - X_{\text{init}}^2}$$

More inflation for $X^2 \rightarrow 2/3$

Power spectrum and consistency relation

$$\mathcal{P}_\zeta = \left(\frac{k}{aH}\right)^{3-2\nu} c_S \left(\frac{c_S}{1+\sigma}\right)^{1-2\nu} \frac{1}{32\pi^2} 2^{2\nu-1} (1-\epsilon)^{2\nu-1} \frac{\Gamma(\nu)^2}{\Gamma(3/2)^2} \frac{H^2}{|\epsilon|}$$

$$r \equiv \frac{A_T^2}{A_\zeta^2} = 16 c_S |\epsilon|$$

For $V = X^n$ and 60 e -folds

$$n_s \approx 0.97$$

independent of the value of power n !

Non-gaussianity

Speed of sound

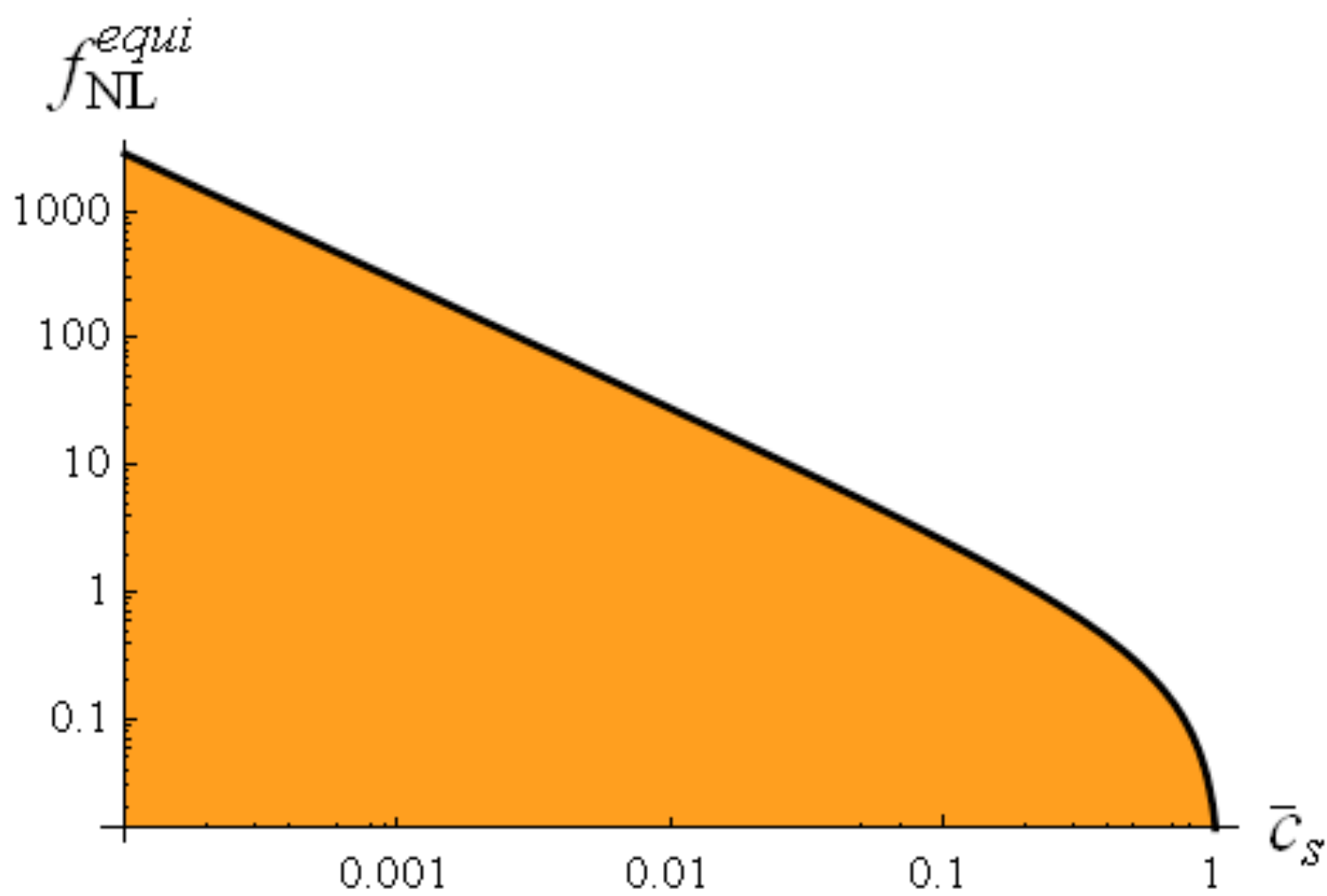
$$c_s^2 = \frac{V_{,XX}X}{V_{,X}}$$

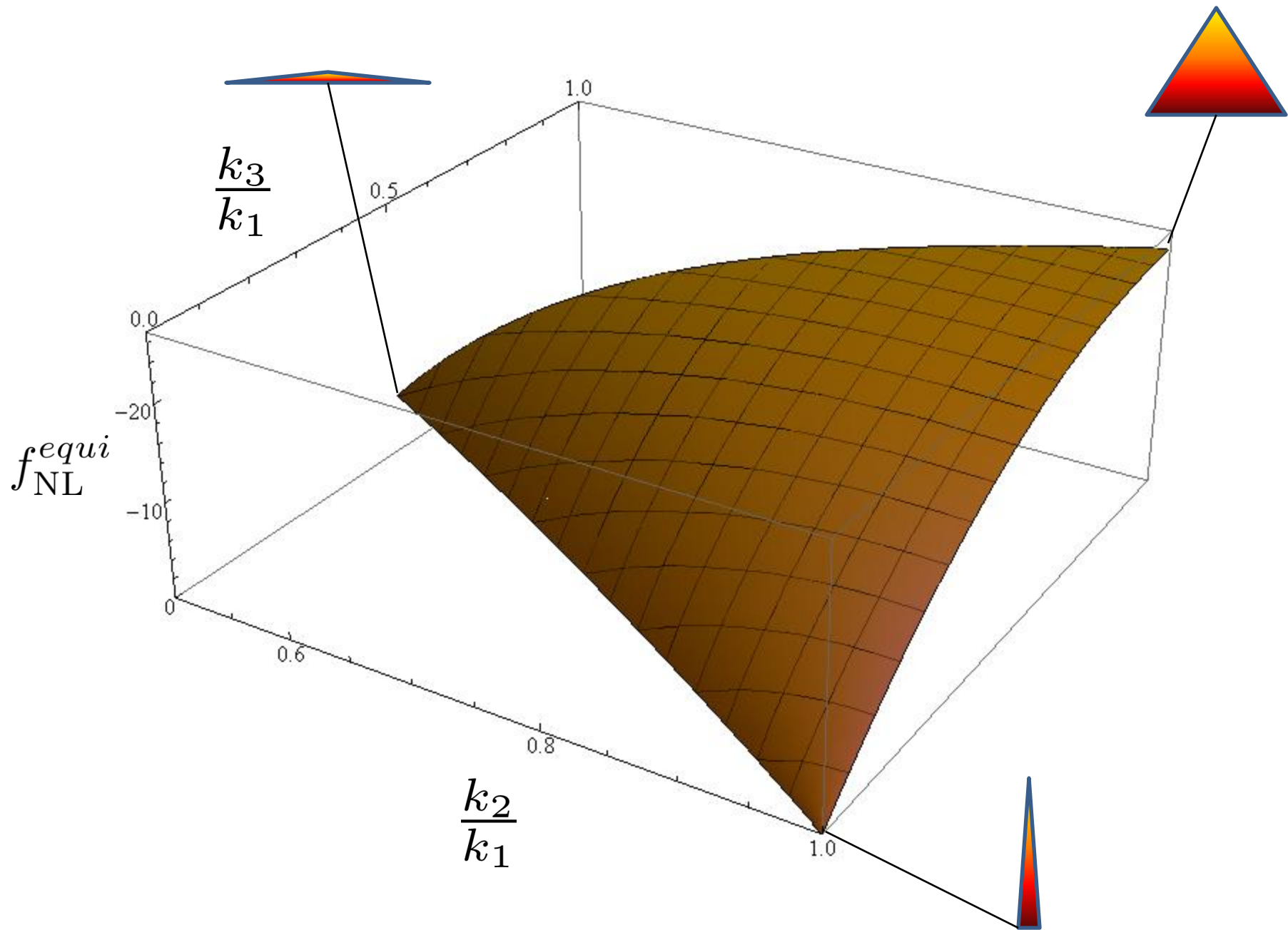
For power law $V = X^n$

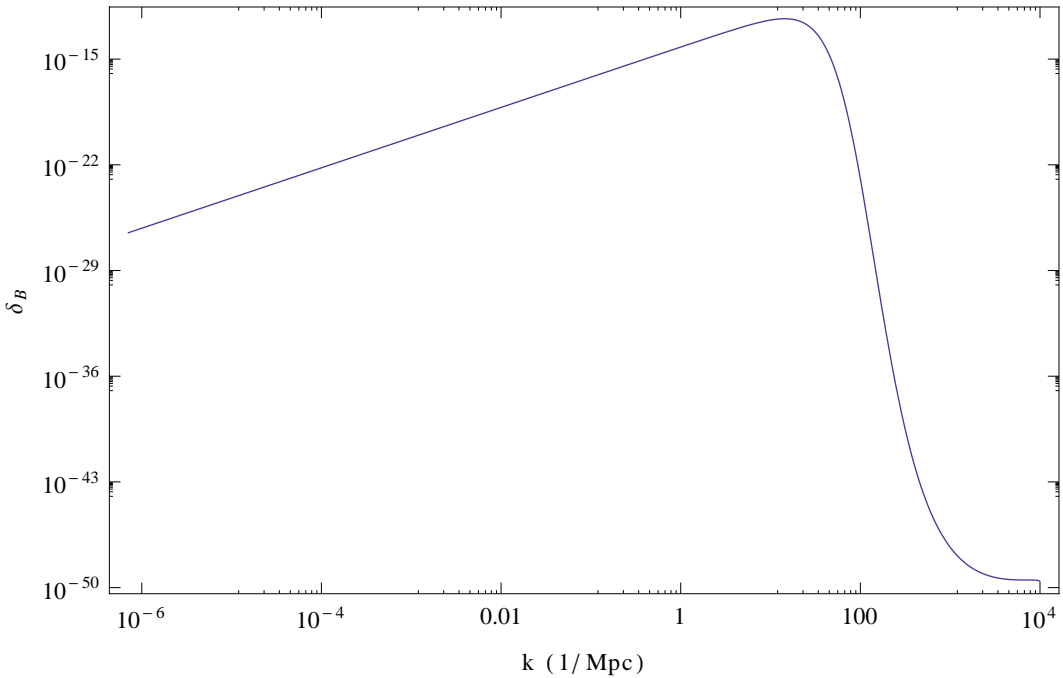
$$c_s^2 = n - 1$$

\Rightarrow non-gaussianity larger for smaller c_s^2

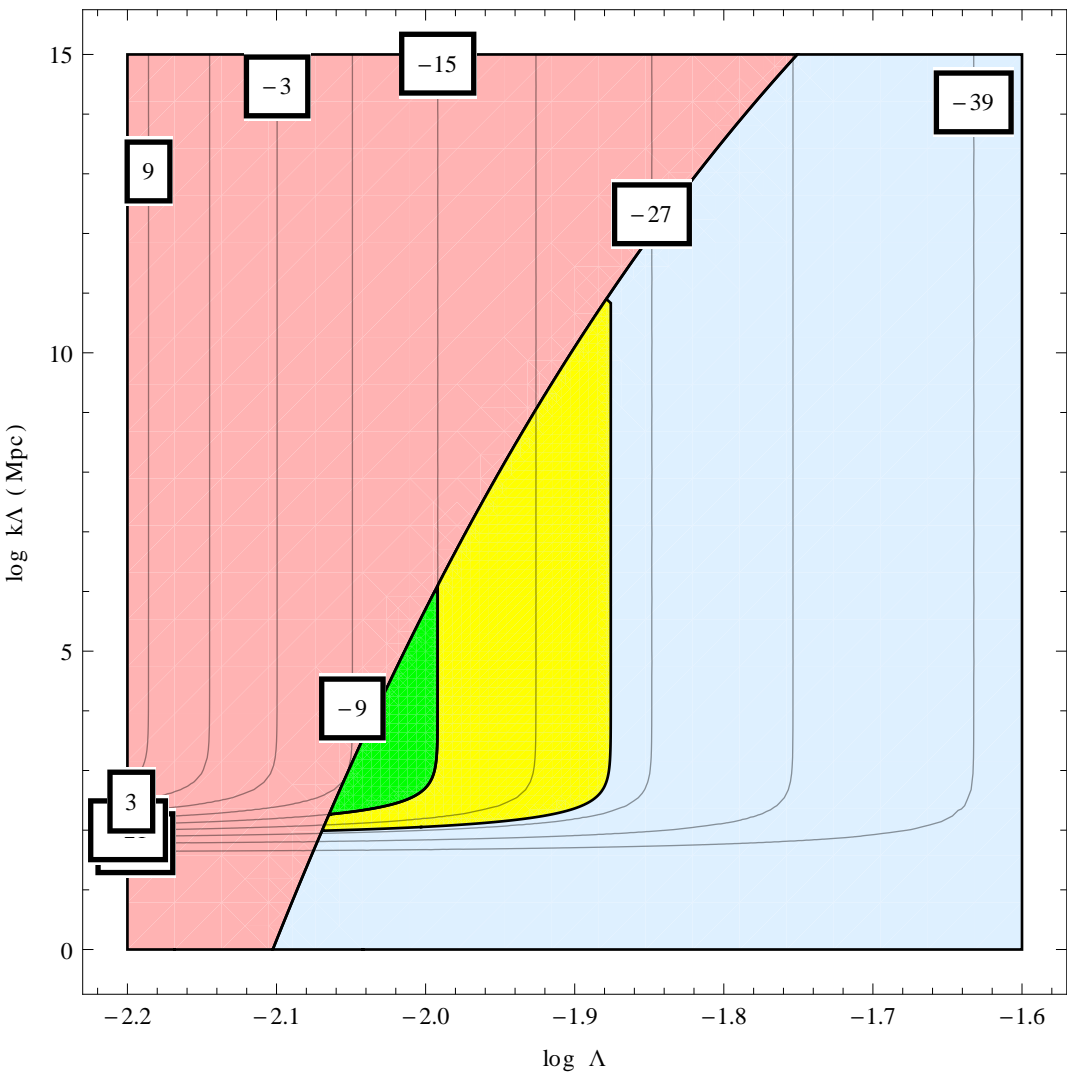
$$f_{NL} \propto \frac{1}{c_s^2}$$



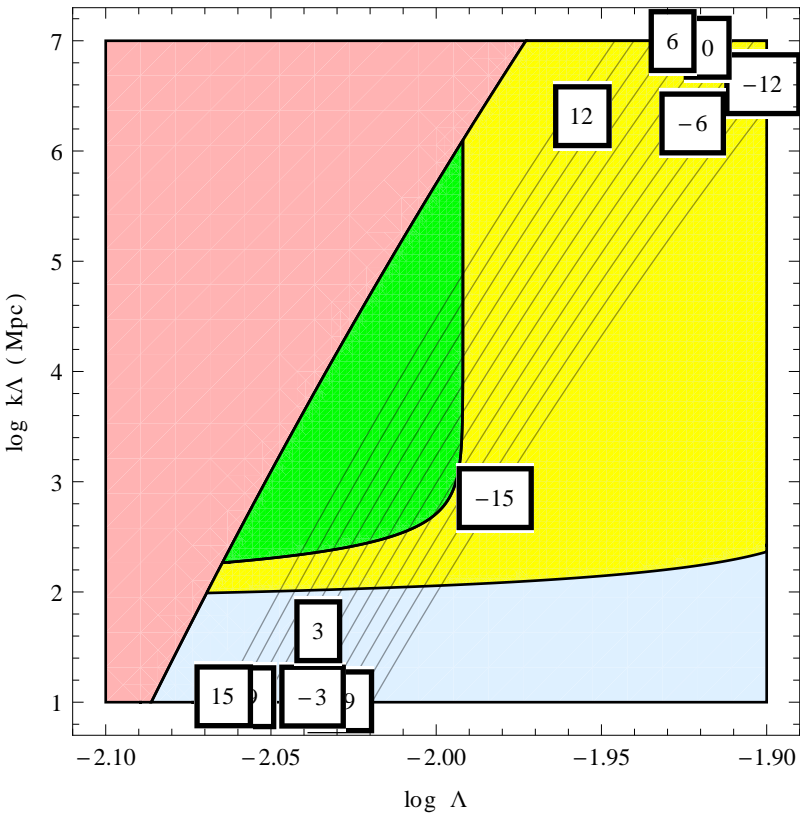




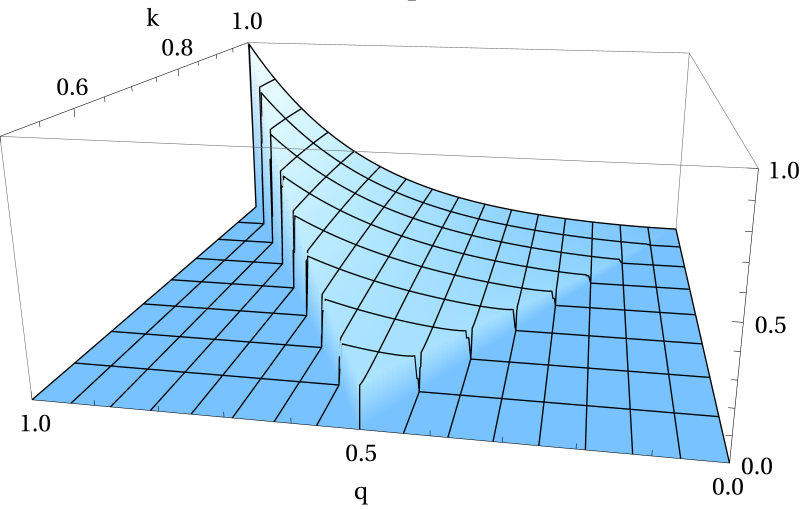
$\log \delta_B @ 1 / \text{Mpc}$



log f_{NL} prefactor



$F(k, q)$



$\log \delta_B @ 1 / \text{Mpc}$

