

(PRIMORDIAL) NON-GAUSSIANITY & N-BODY SIMULATIONS

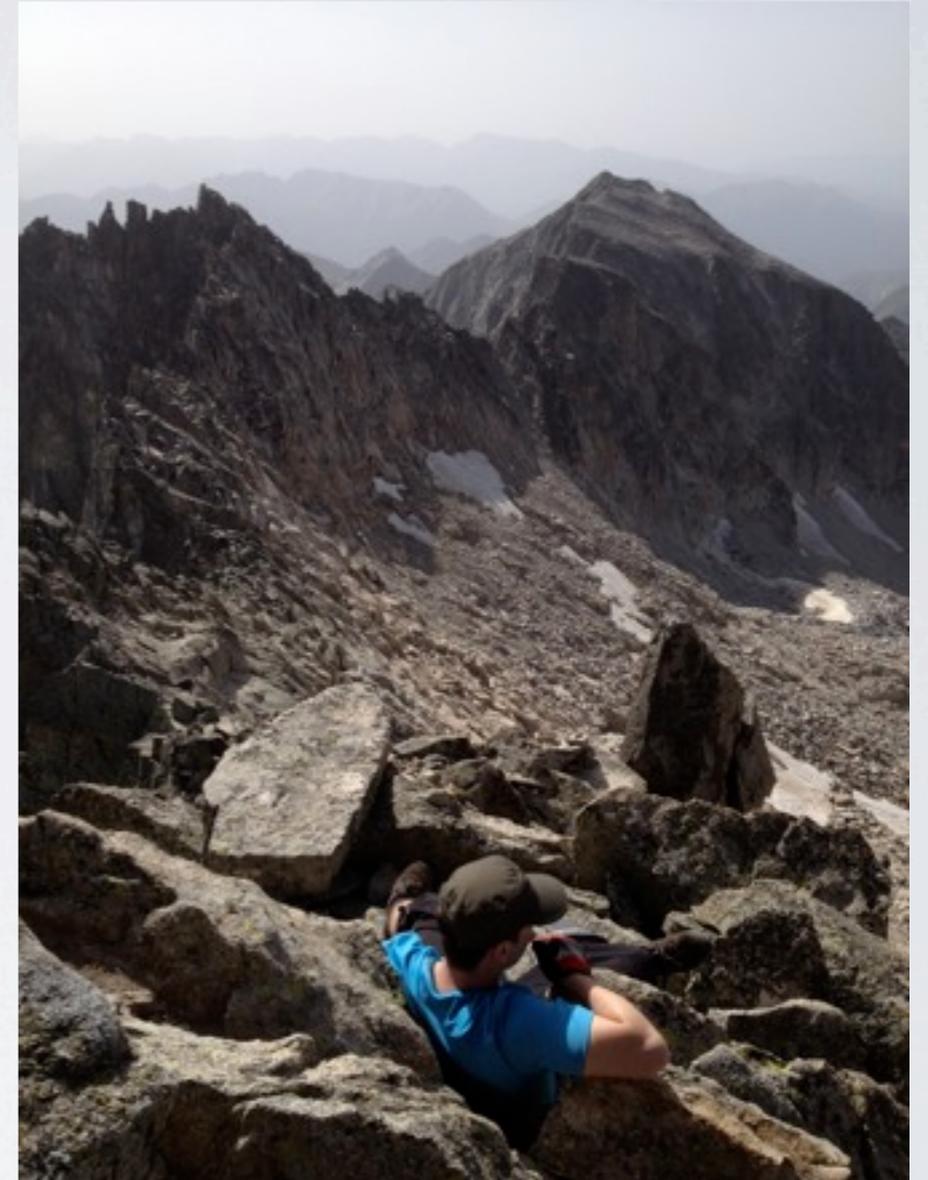
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arXiv: 1008.1730, 1108.3813, 1208.5678, 1209.xxxx

Centro de Ciencias de Benasque Pedro Pascual
22nd August 2012

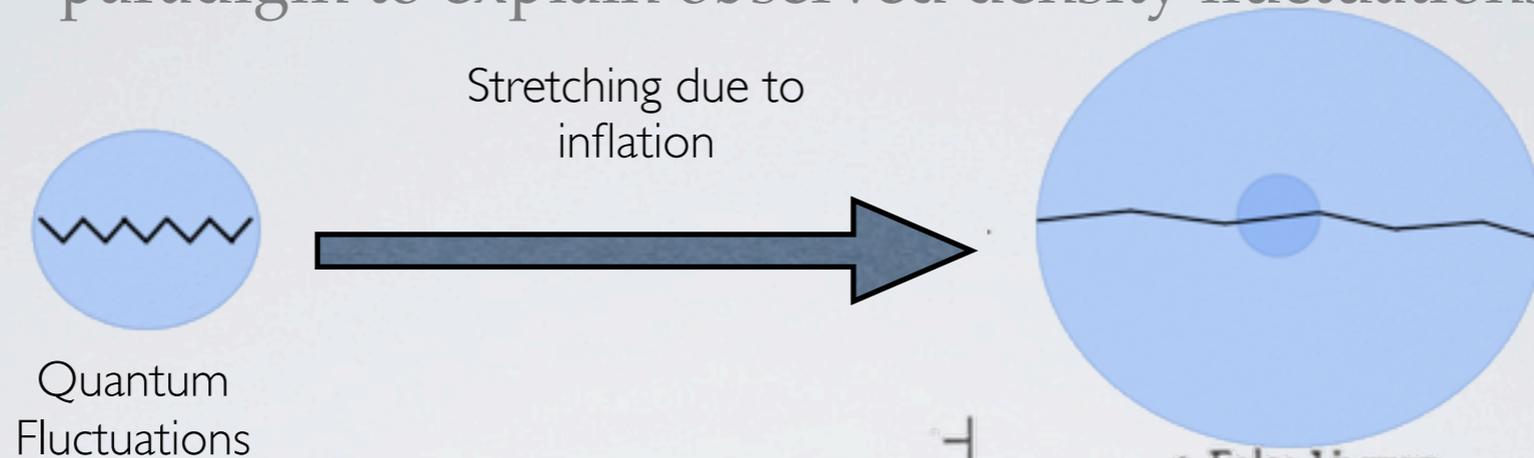
OUTLINE

- Non-Gaussianity primer
- Initial Conditions for N-body Simulations
- Analysing results of N-body Simulations
- Comparison to Halo model and 1-loop Corrections (+ simple fitting formula)

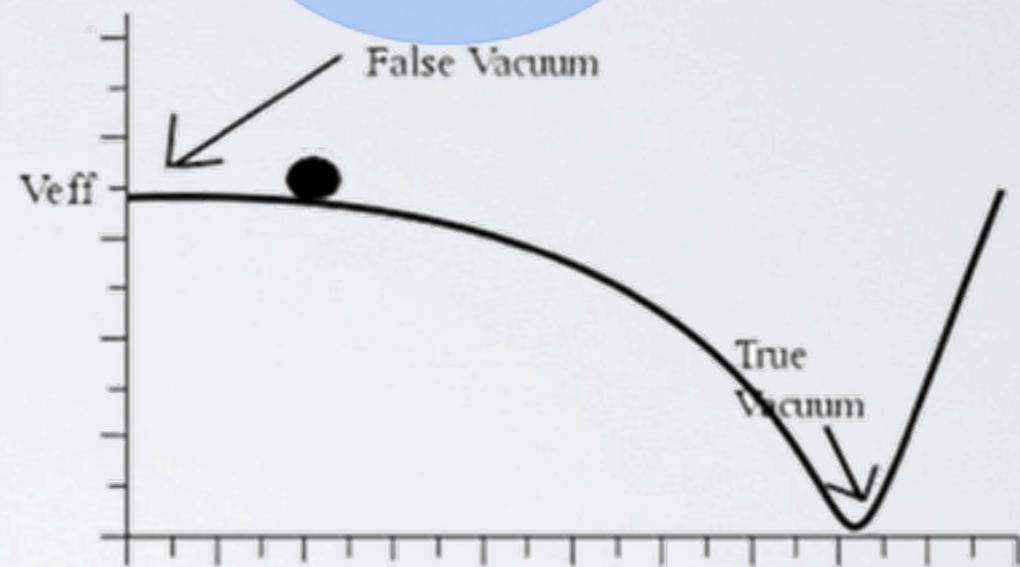


@ the primordial level

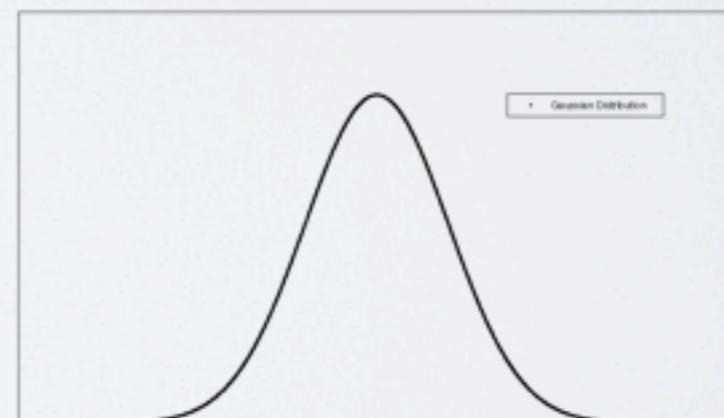
- Inflation = paradigm to explain observed density fluctuations



- Standard Inflation ... Require **slow-roll** to satisfy the horizon and flatness problems.



- Prediction... **Gaussian** primordial fluctuations



- Conditions of **standard** case:

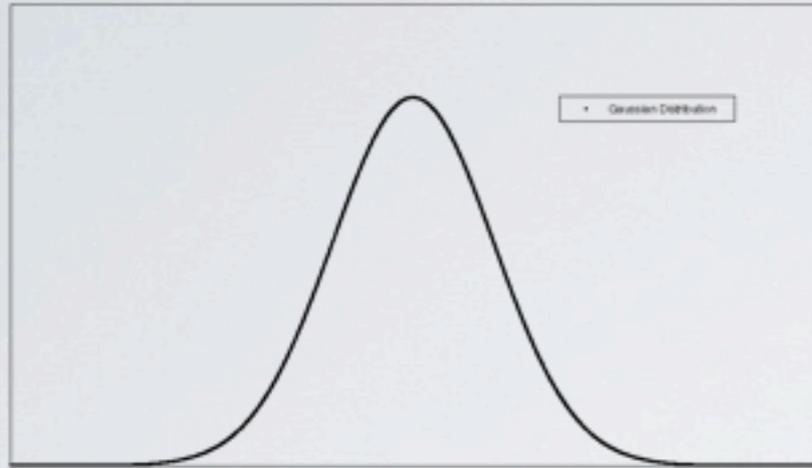
- (a) Single scalar field
- (b) Canonical kinetic energy
- (c) Slow roll
- (d) Standard vacuum state

Break any of these



Non-Gaussianities

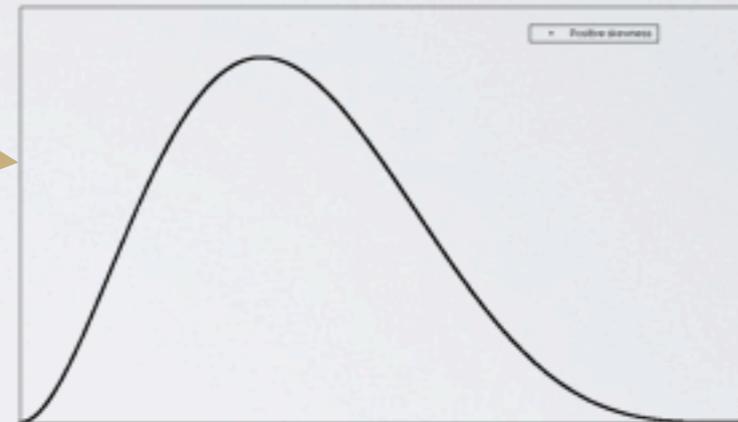
Skewness ... Bispectrum



Gaussian distribution

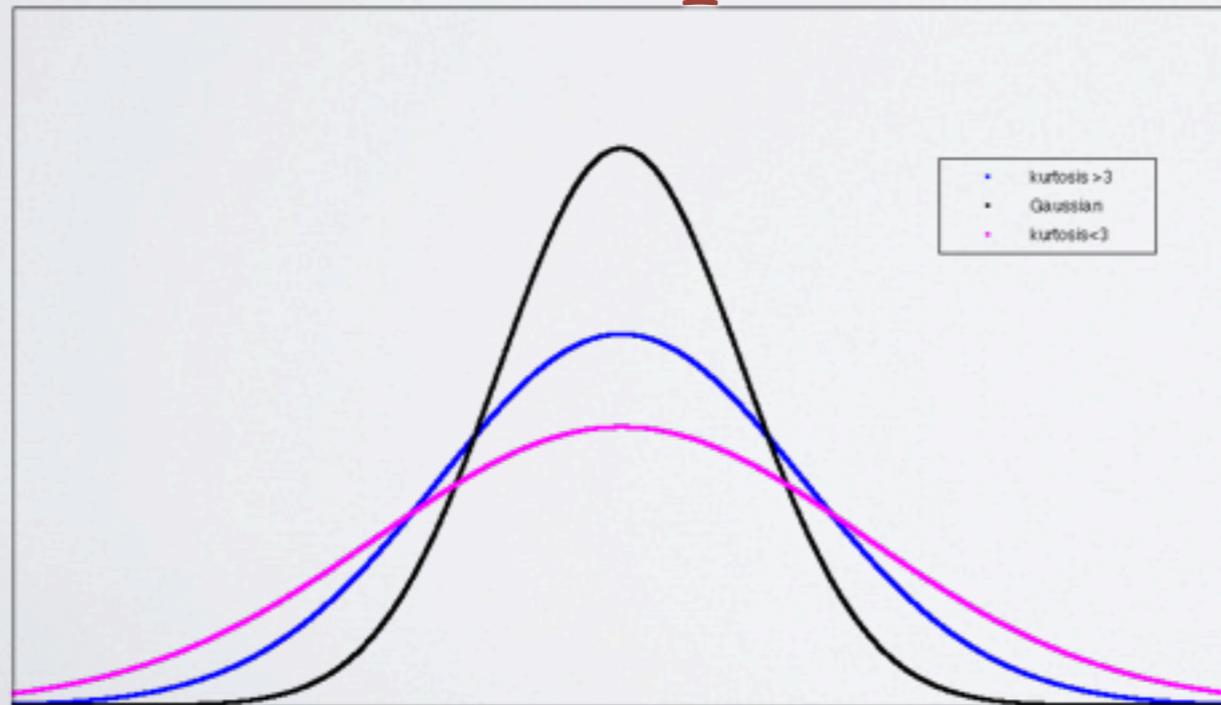


Negatively skewed



Positively skewed

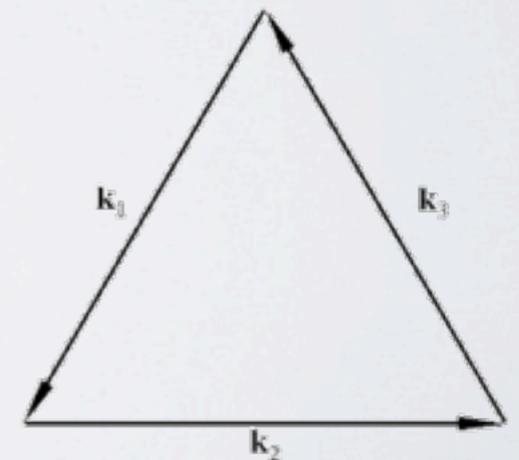
Kurtosis ... Trispectrum



Trispectrum = (connected) 4-point correlator

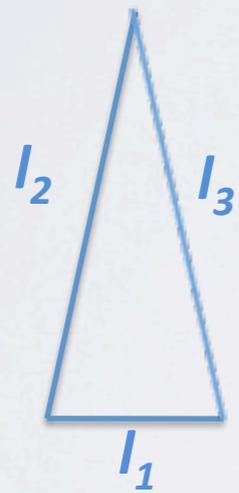
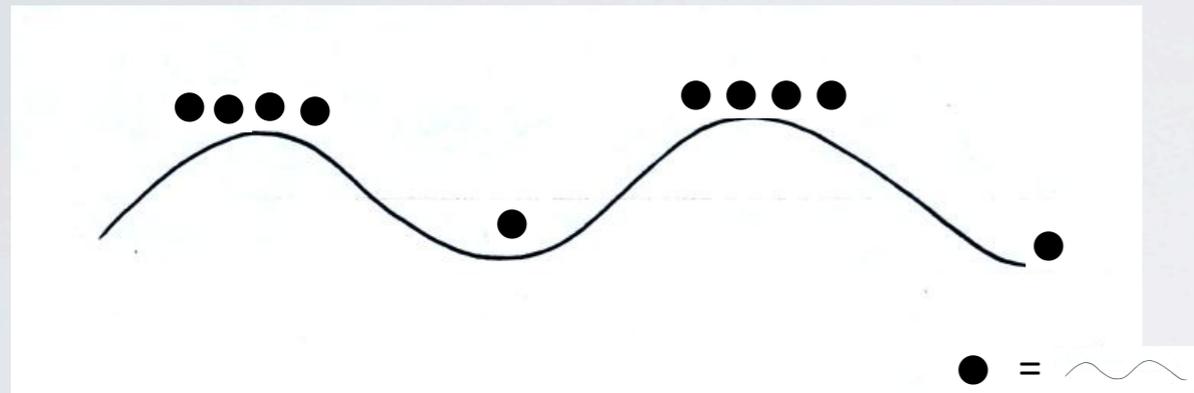
$$\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2)\Phi(\mathbf{k}_3)\Phi(\mathbf{k}_4) \rangle_c = (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4) T_\Phi(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4)$$

$$\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2)\Phi(\mathbf{k}_3) \rangle = (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\Phi(k_1, k_2, k_3)$$

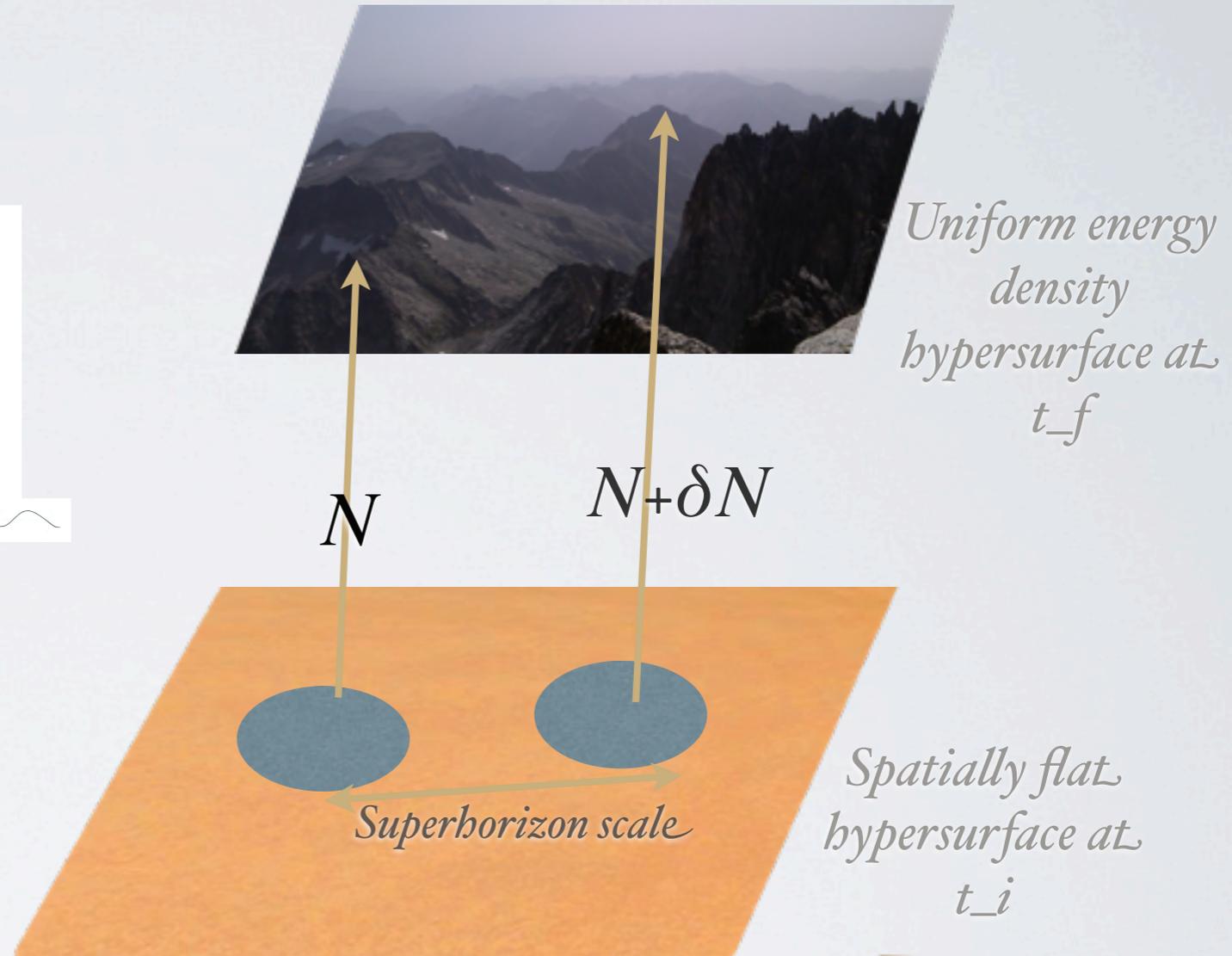


Local Bispectrum ... superhorizon evolution

$$B(k_1, k_2, k_3) = 2f_{\text{NL}} [P(k_1)P(k_2) + P(k_2)P(k_3) + P(k_3)P(k_1)]$$



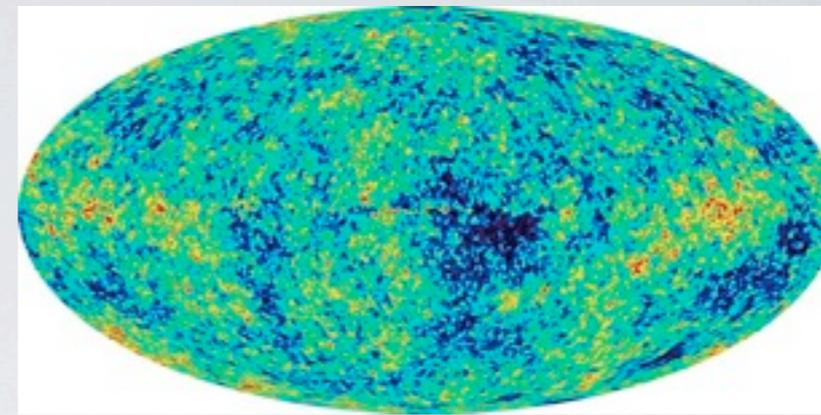
NOTE SQUEEZED SHAPE



Many more examples...

- *DBI inflation*
- *Ghost inflation*
- *Second order corrections to single field inflation*
- *Warm inflation*
- *Inflation with non-standard vacua (&/or higher order kinetic terms)*
- *Inflation with a step in the potential*
- *Inflation with oscillatory modifications to the potential*
- *Non-local inflation*
- ...

@ the CMB level



$$a_{lm} \longleftarrow \xrightarrow{\Delta_l(k)} \Phi(\mathbf{k})$$

Primordial to CMB via
transfer function

$$a_{lm} = 4\pi(-i)^l \int \frac{d^3k}{(2\pi)^3} \Delta_l(k) \Phi(\mathbf{k}) Y_{lm}^*(\hat{\mathbf{k}})$$

CMB Bispectrum



Primordial Bispectrum

$$B_{m_1 m_2 m_3}^{l_1 l_2 l_3} = \langle a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3} \rangle = \mathcal{G}_{m_1 m_2 m_3}^{l_1 l_2 l_3} b_{l_1 l_2 l_3}$$

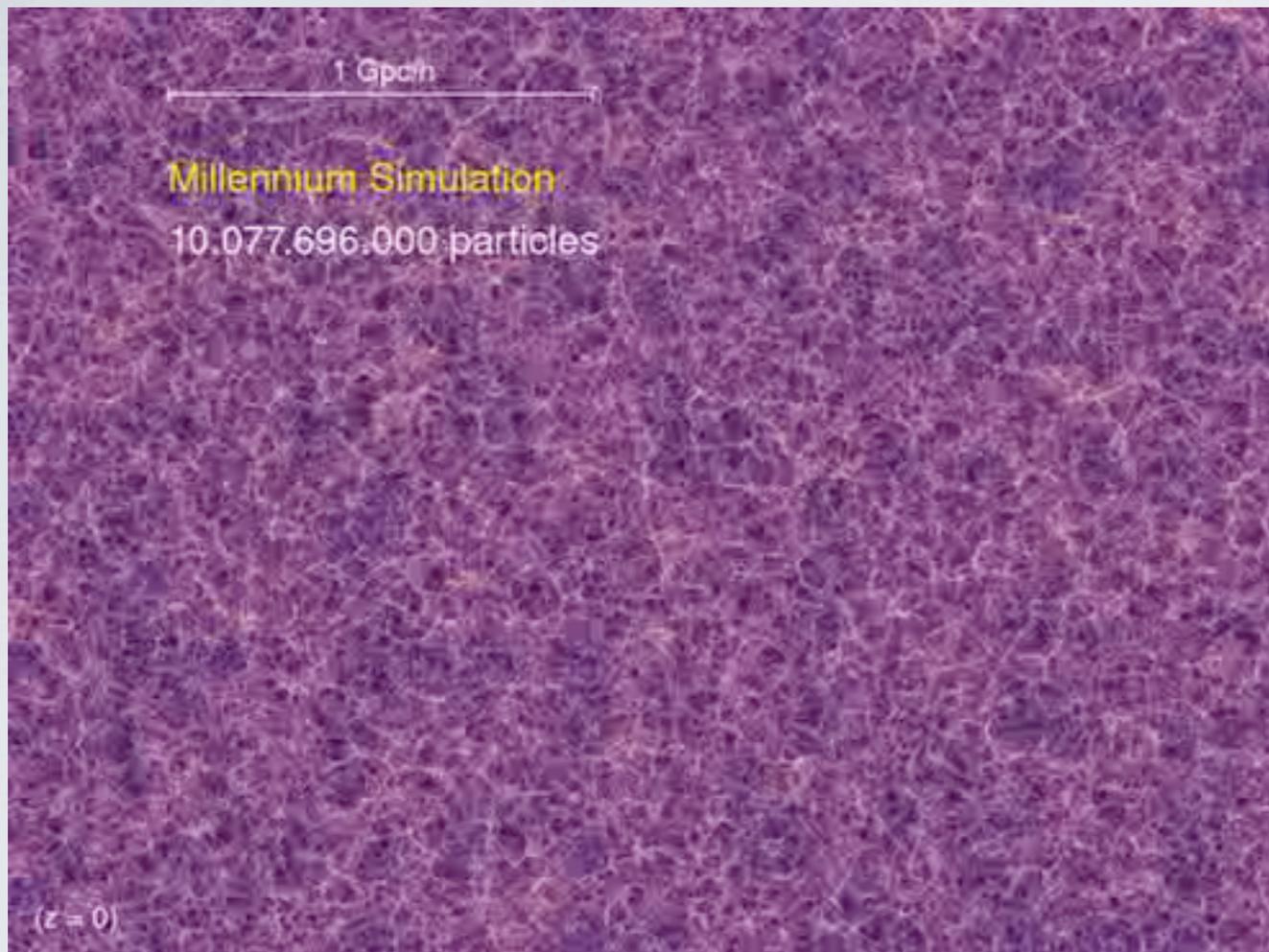
$$\langle \Phi(\mathbf{k}_1) \Phi(\mathbf{k}_2) \Phi(\mathbf{k}_3) \rangle = (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B_\Phi(k_1, k_2, k_3)$$

$$b_{l_1 l_2 l_3} = \Delta_\Phi^2 \left(\frac{2}{\pi} \right)^3 \int dk_1 dk_2 dk_3 \Delta_{l_1}(k_1) \Delta_{l_2}(k_2) \Delta_{l_3}(k_3) S(k_1, k_2, k_3) \\ \times \int dx x^2 j_{l_1}(k_1 x) j_{l_2}(k_2 x) j_{l_3}(k_3 x),$$

$$S(k_1, k_2, k_3) = \frac{(k_1 k_2 k_3)^2}{\Delta_\Phi^2} B_\Phi(k_1, k_2, k_3)$$

@ the *LSS* level

At late times fluctuations grow & non-Gaussianity is induced by gravitational evolution



- ◆ Need N-body simulations, e.g. Gadget3

Video: <http://www.mpa-garching.mpg.de/galform/virgo/millennium>

Advantage: More information (z dependence)

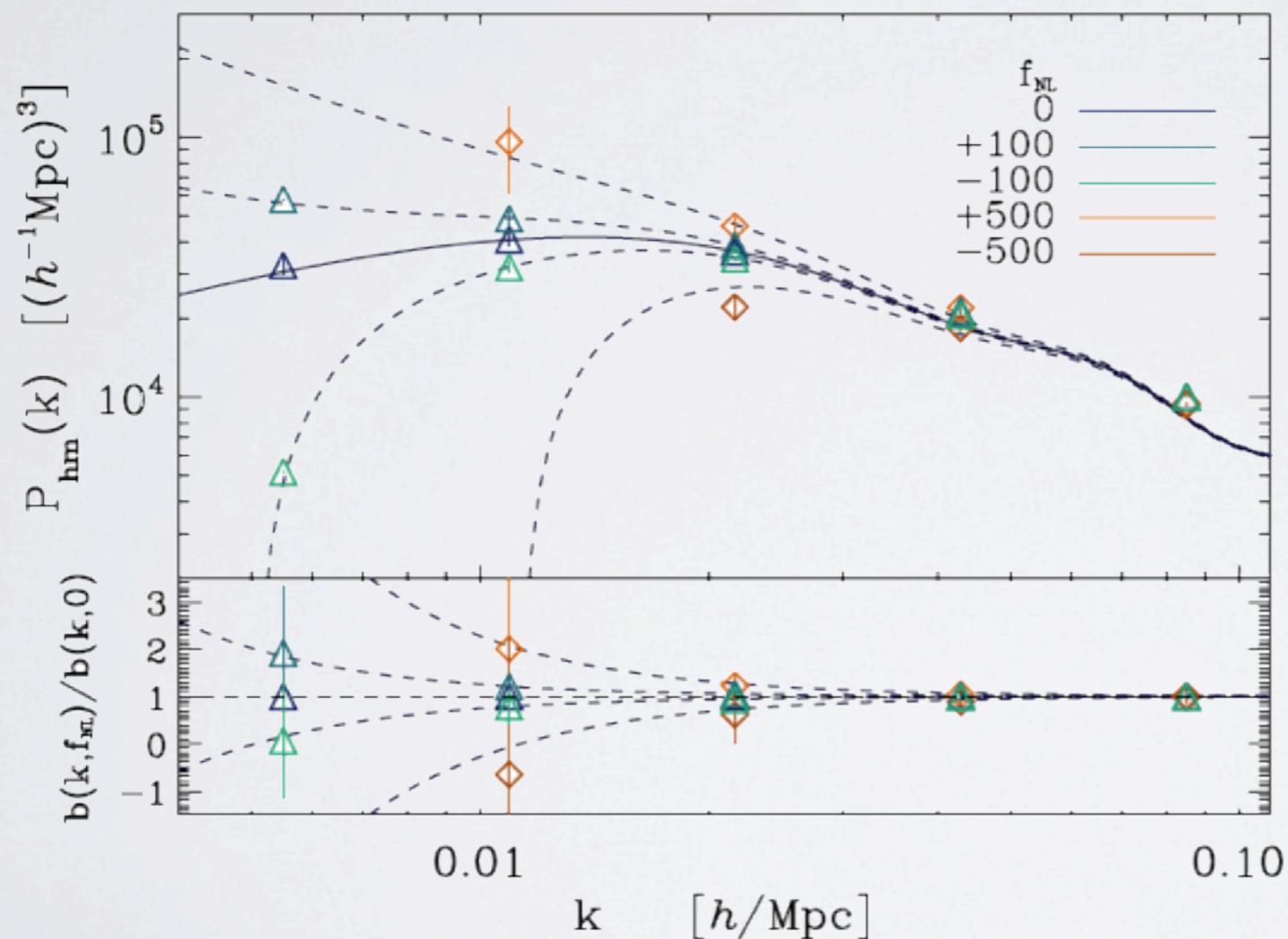
Drawback: Non-linear evolution...difficult to distinguish primordial signal

Primordial non-Gaussianity and N-body Simulations

Searching for non-Gaussianity in large scale structure

e.g. Use clustering of rare peaks to get a constraint on f_{NL}

(0710.4560 Dalal et al)



The primordial bispectrum gives a contribution to the galaxy power spectrum

State of play of non-Gaussianity in N-body simulations

Until recently only the **local** type was simulated

Then... Verde, Wagner, (Boubekkeur) [arXiv:1006.5793](#), [arXiv:1102.3229](#)

Scoccimarro, Hui, Manera, Chan [arXiv:1108.5512](#)

However, they still require the underlying shape to be **separable**

Also, analysis still is still done on **slices** at particular redshifts.

Difficult to be sure you're seeing everything!

Initial Conditions for N-body Simulations

$$\Phi = \Phi^G + \frac{1}{6}F_{\text{NL}}\Phi^B + \frac{1}{24}\tau_{\text{NL}}\Phi^T$$

$$\Phi^B(\mathbf{k}) = \int \frac{d^3\mathbf{k}'d^3\mathbf{k}''}{(2\pi)^3} \delta_D(\mathbf{k} - \mathbf{k}' - \mathbf{k}'') \Phi^G(\mathbf{k}')\Phi^G(\mathbf{k}'') \frac{B(k, k', k'')}{P(k)P(k') + P(k)P(k'') + P(k')P(k'')}$$

Complexity of N^6 ($N=512!$)

Problem is non-separability...

Hint: Try the modal approach arXiv:1008.1730 Fergusson,DMR,Shellard

$$\frac{B(k, k', k'')}{P(k)P(k') + P(k)P(k'') + P(k')P(k'')} = \sum_{rst} \alpha_{rst}^Q q_r(k)q_s(k')q_t(k''),$$

$$\Phi^B(\mathbf{k}) = \sum_n \alpha_n^Q q_{\{r}(k) \int d^3\mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} M_s(\mathbf{x}) M_{t\}}(\mathbf{x})$$

$$M_s(\mathbf{x}) = \int \frac{d^3\mathbf{k}}{(2\pi)^3} \Phi^G(\mathbf{k}) q_s(k) e^{-i\mathbf{k}\cdot\mathbf{x}} \text{ FFT only!!}$$

Application to Estimation

$$\mathcal{E} = \int \frac{d^3k_1}{(2\pi)^3} \frac{d^3k_2}{(2\pi)^3} \frac{d^3k_3}{(2\pi)^3} \frac{(2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) B(k_1, k_2, k_3)}{P(k_1)P(k_2)P(k_3)} [\delta_{\mathbf{k}_1}^{obs} \delta_{\mathbf{k}_2}^{obs} \delta_{\mathbf{k}_3}^{obs} - 3 \langle \delta_{\mathbf{k}_1}^{sim} \delta_{\mathbf{k}_2}^{sim} \rangle \delta_{\mathbf{k}_3}^{obs}]$$

Modal approach to the rescue!

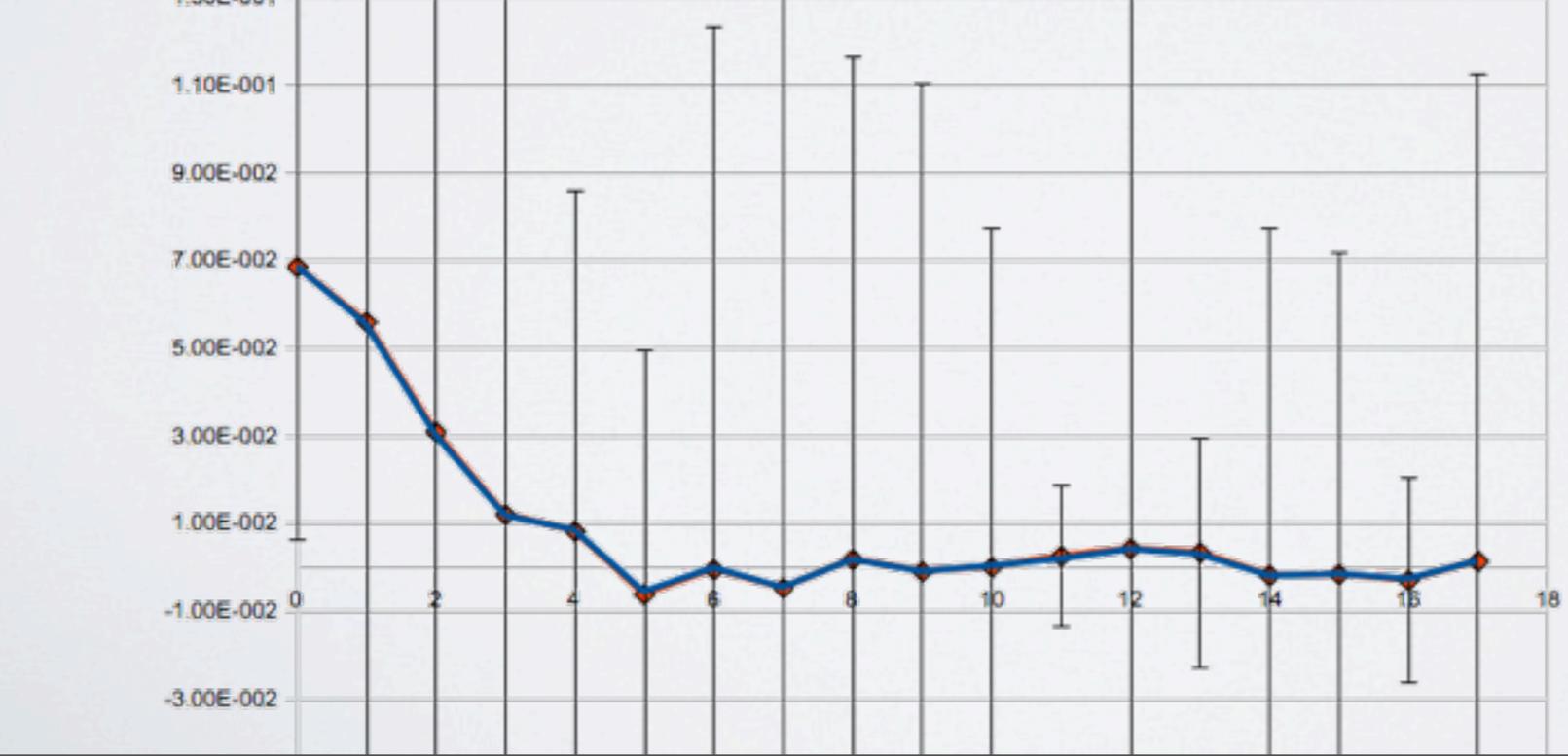
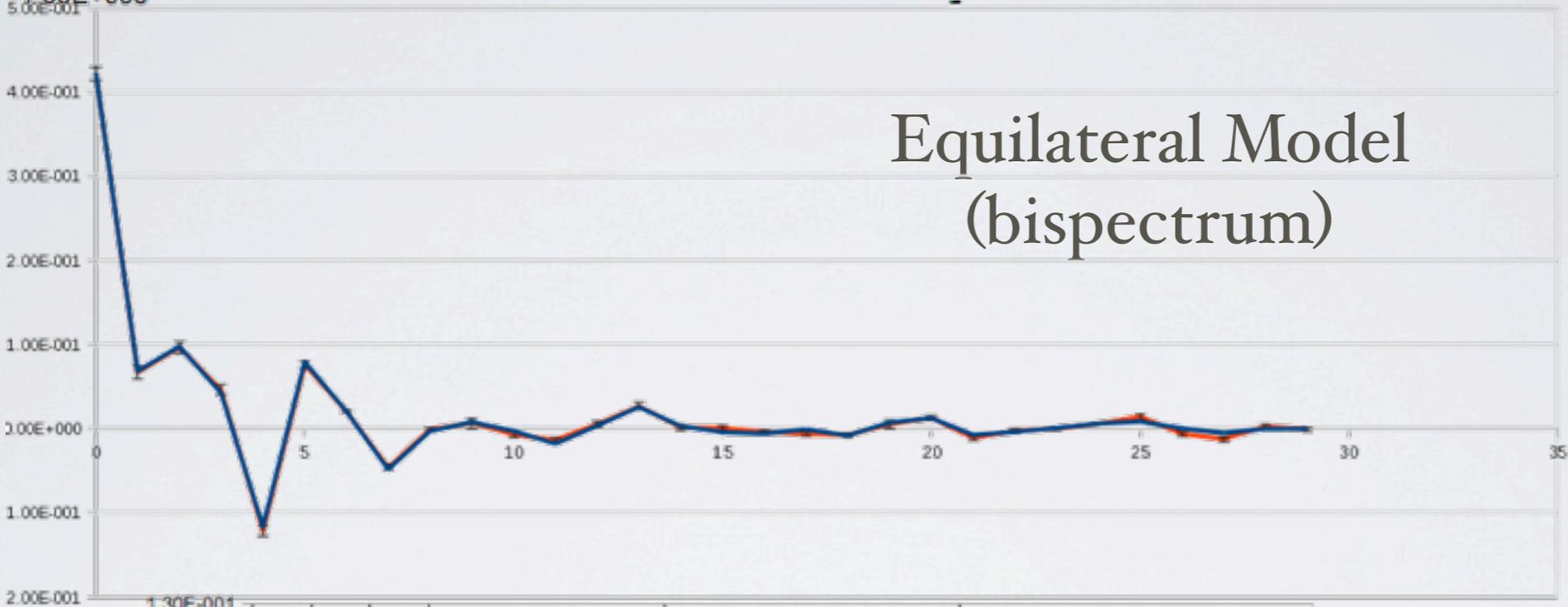
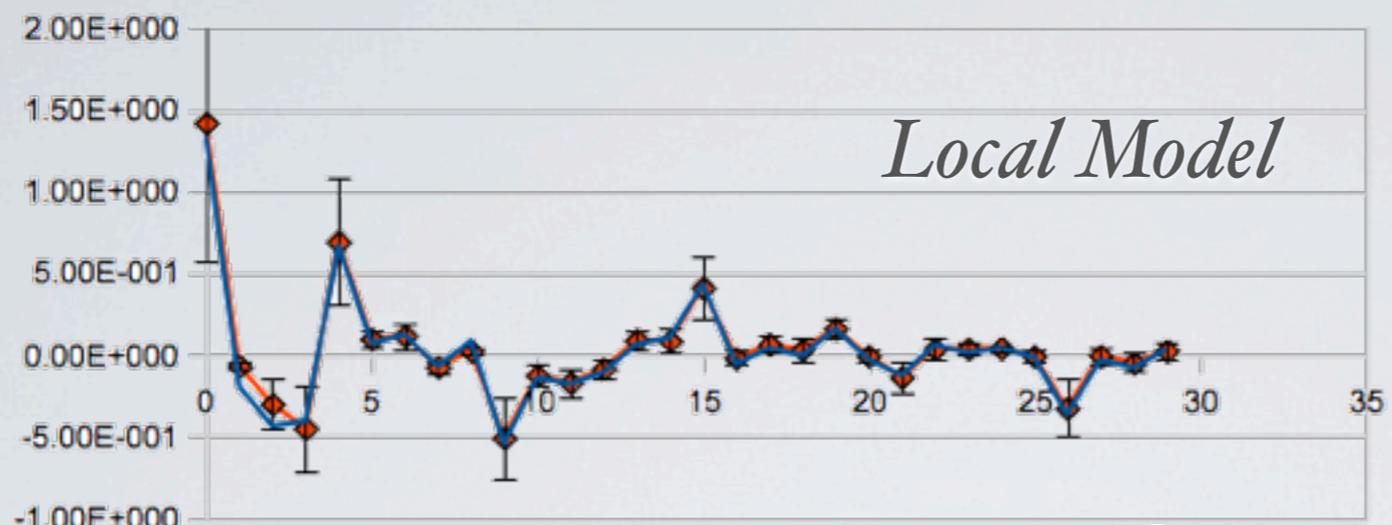
$$\frac{B(k_1, k_2, k_3) v(k_1)v(k_2)v(k_3)}{\sqrt{P(k_1)P(k_2)P(k_3)}} = \sum \alpha_n^{\mathcal{Q}} \mathcal{Q}_n(k_1, k_2, k_3)$$


 $\mathcal{E} = \sum_n \alpha_n^{\mathcal{Q}} \beta_n^{\mathcal{Q}}$
 where $\beta_n^{\mathcal{Q}} = \int d^3x M_r(\mathbf{x}) M_s(\mathbf{x}) M_t(\mathbf{x})$

$M_p(\mathbf{x}) = \int d^3k \frac{\delta_{\mathbf{k}}^{obs} q_r(k) e^{i\mathbf{k}\cdot\mathbf{x}}}{\sqrt{kP(k)}}$

$\left. \begin{array}{l} \text{FFT \&} \\ \text{a single 3D} \\ \text{integral} \end{array} \right\}$

Similarly can define a correlator between shapes

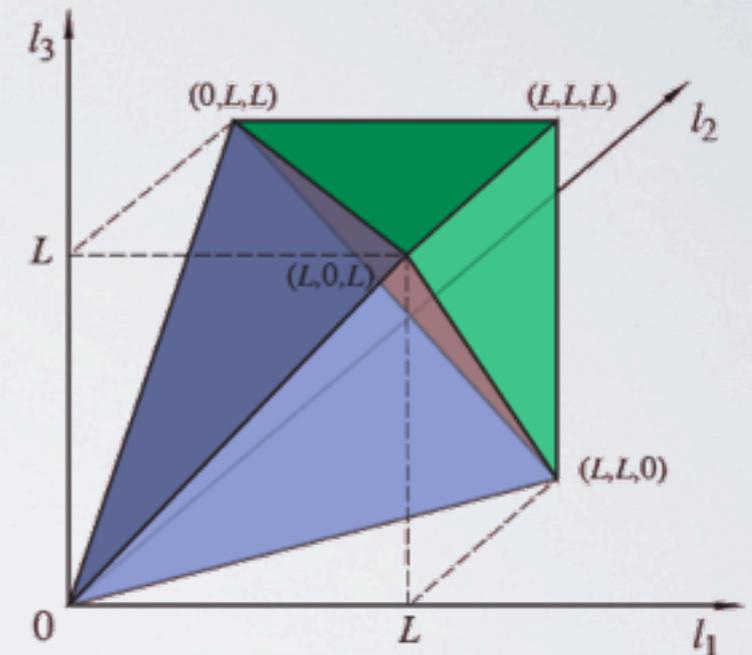
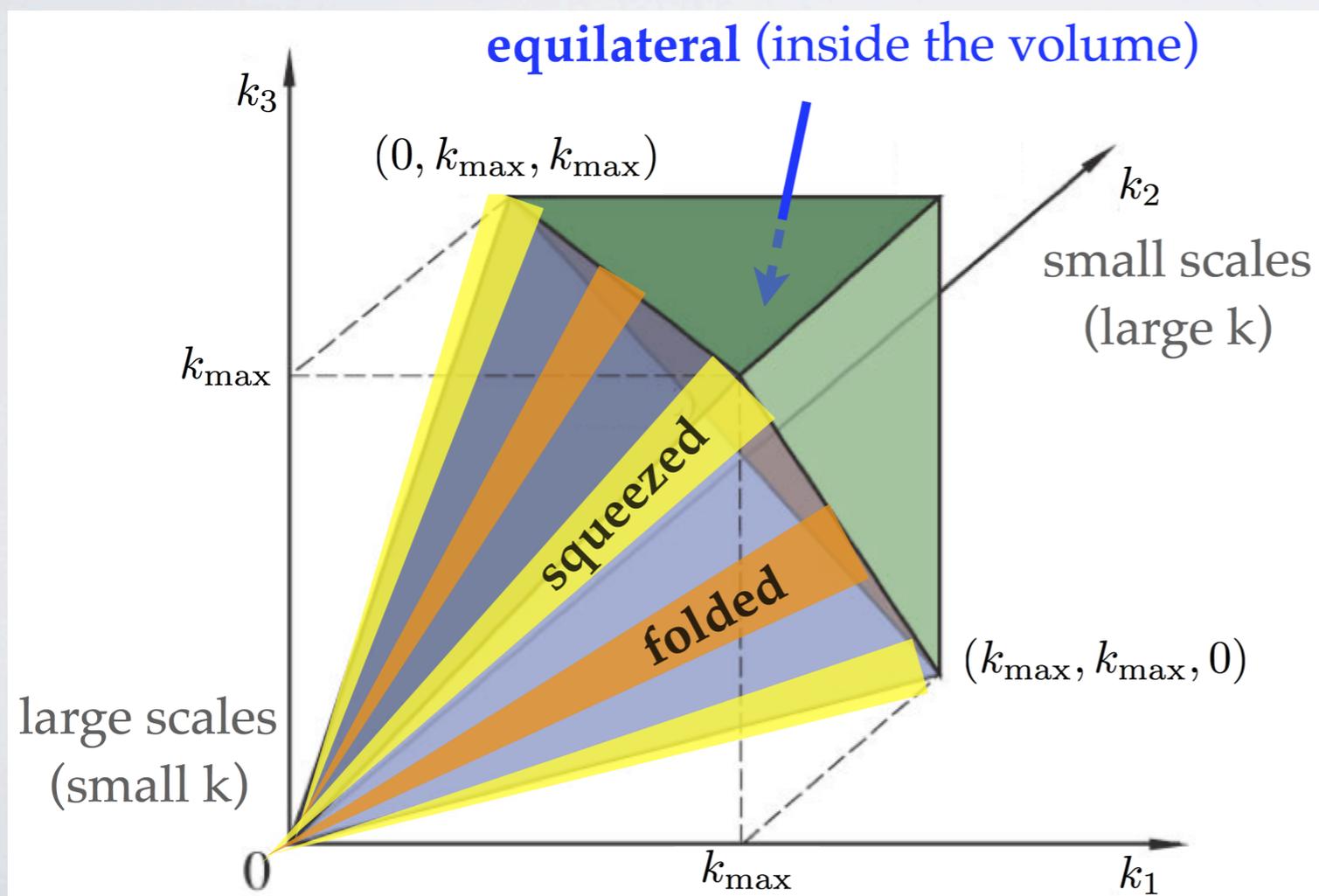


Local Model g_{n1}
(trispectrum)

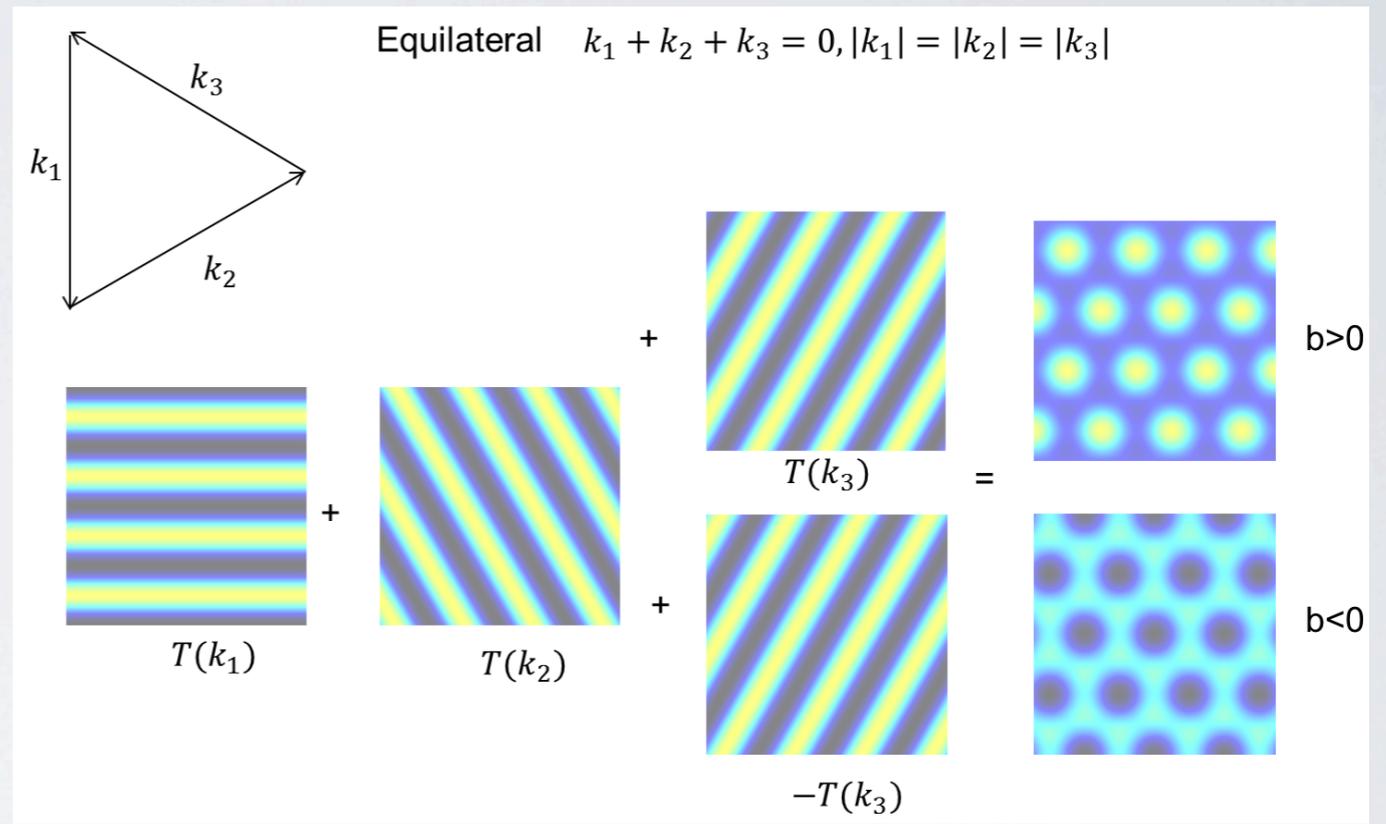
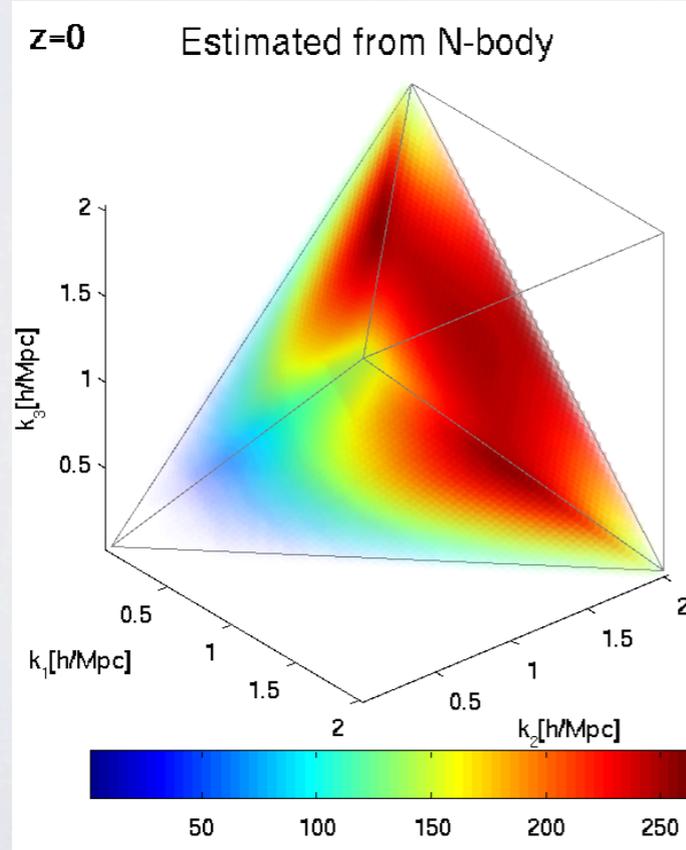
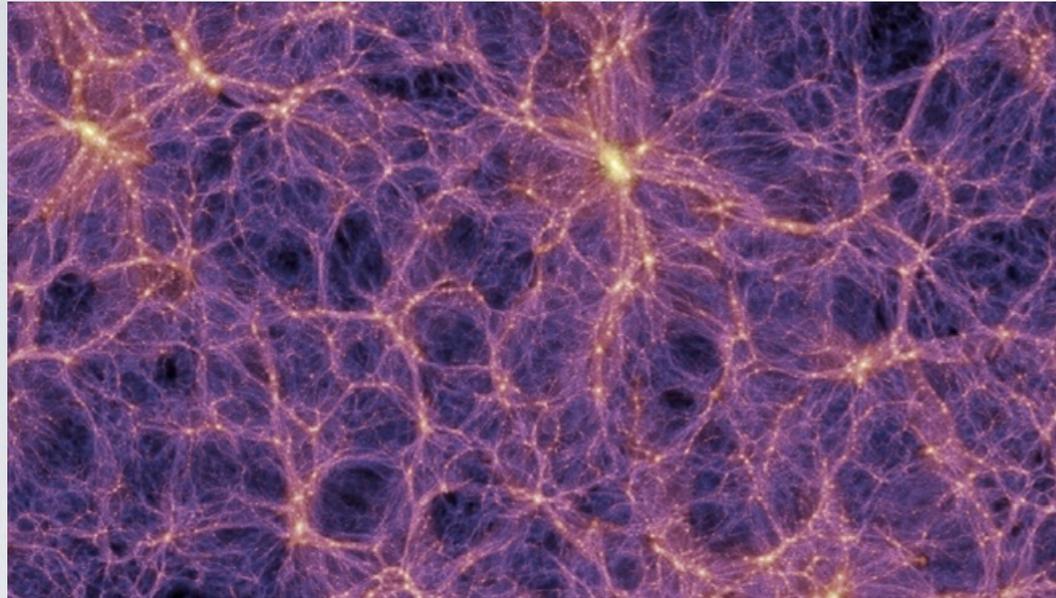
Column BG
Column BH

A BIT OF VISUALISATION...

$$S(k_1, k_2, k_3) = \sum_{prs} \alpha_{prs} q_p(k_1) q_r(k_2) q_s(k_3)$$

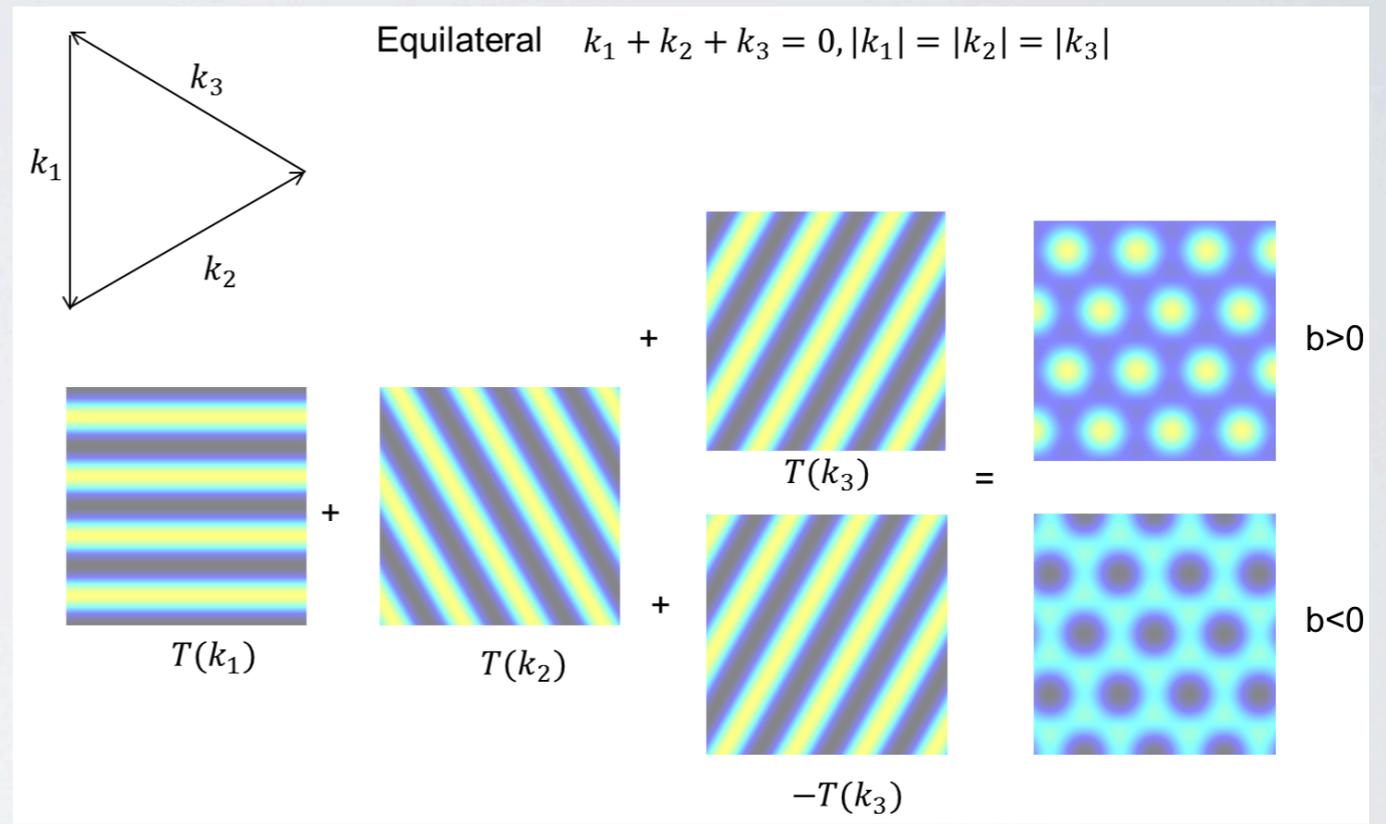
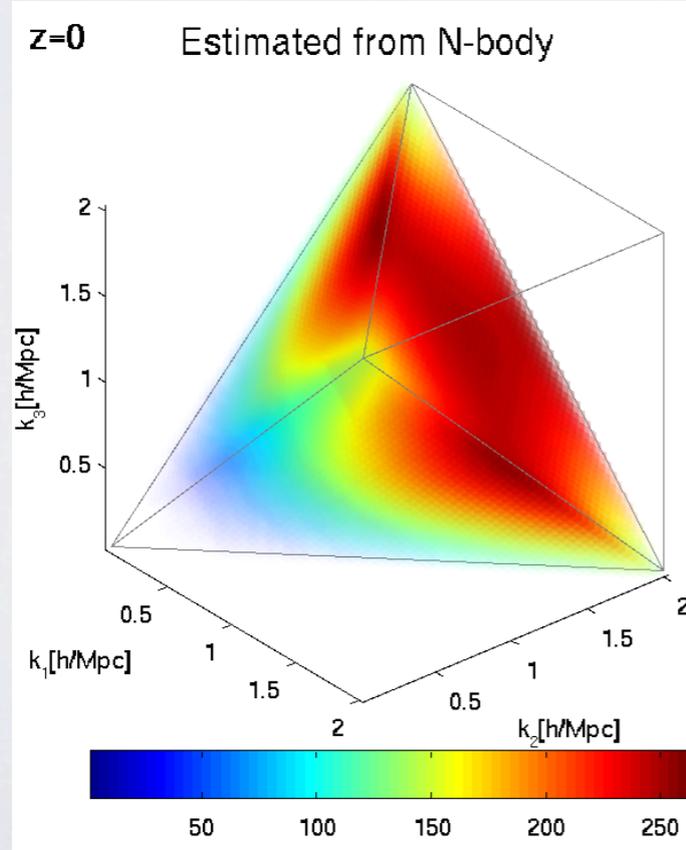
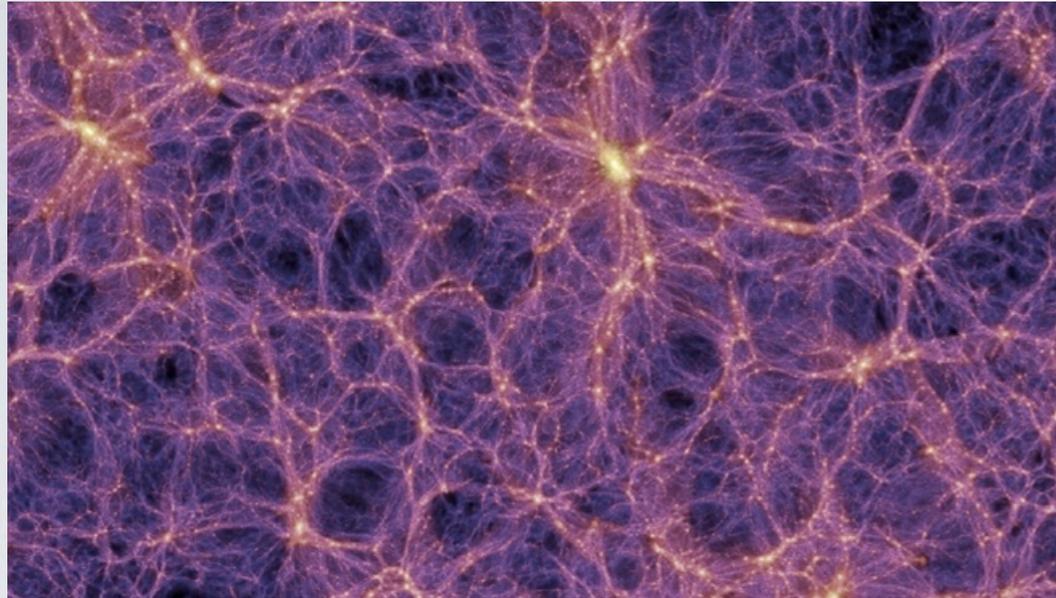


- What does the Gravity Bispectrum look like?



arXiv:1107.5431 Lewis

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arXiv:1107.5431 Lewis

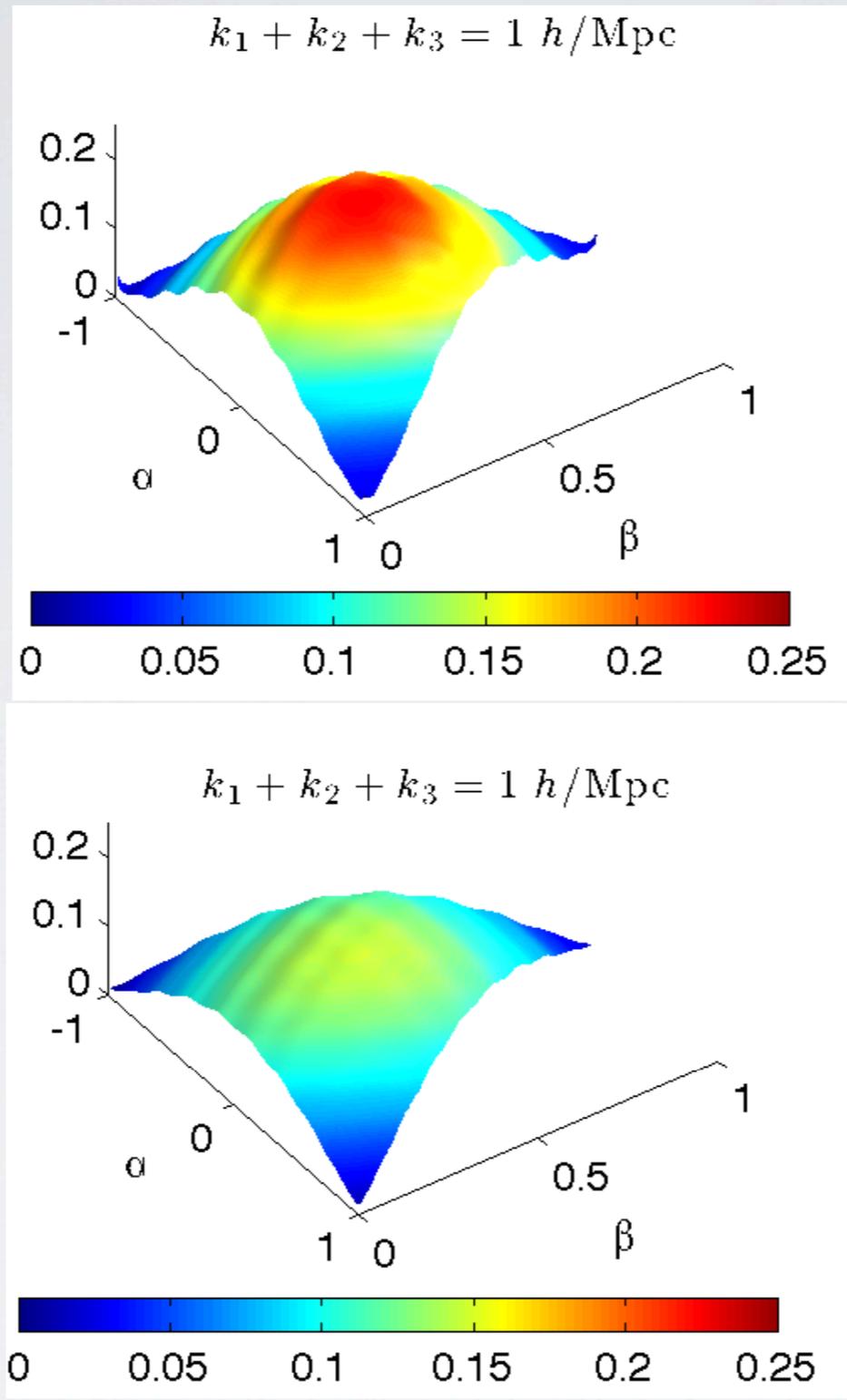


Figure 4. Bispectrum weights $[P_\delta(k_1)P_\delta(k_2) + 2 \text{ perms}]^{-1}$ (top) and $\sqrt{k_1 k_2 k_3 / [P_\delta(k_1)P_\delta(k_2)P_\delta(k_3)]}$ (bottom) evaluated with CAMB [30] at redshift $z = 30$ on slices with $k_1 + k_2 + k_3 = 1 \text{ h/Mpc}$.

Lots of results - *the evolving bispectrum*

Getting the N-body initial conditions

(A) Poisson equation to convert to density perturbations

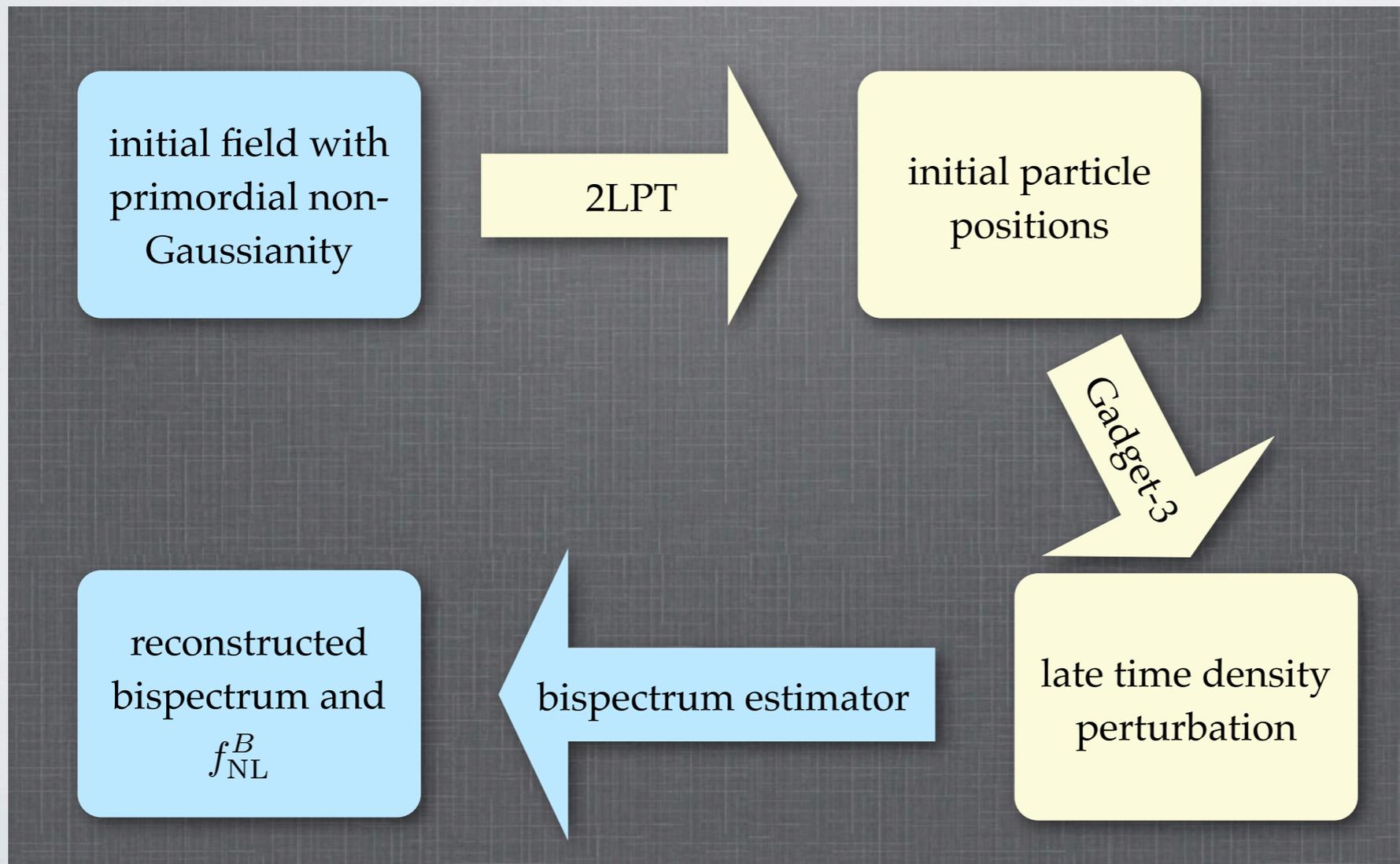
$$\delta_{\mathbf{k}}(a) = M(k; a)\Phi_{\mathbf{k}} \qquad M(k; a) = -\frac{3 k^2 T(k)}{5 \Omega_m H_0^2} D_+(a)$$

(B) 2nd order Lagrangian perturbation theory to get the **positions and velocities** of the particles... must decide whether to use a glass or a grid configuration.

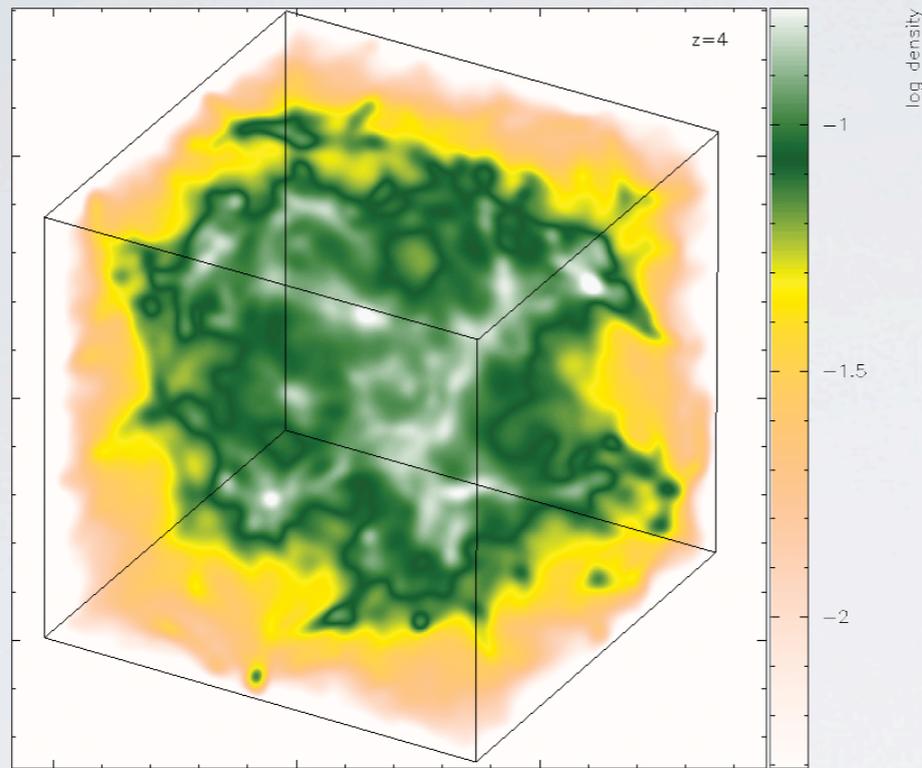
(C) Other things to choose: Boxsize
 Number of particles
 Initial redshift
 plus the 'gas softening length'

Gravity
only

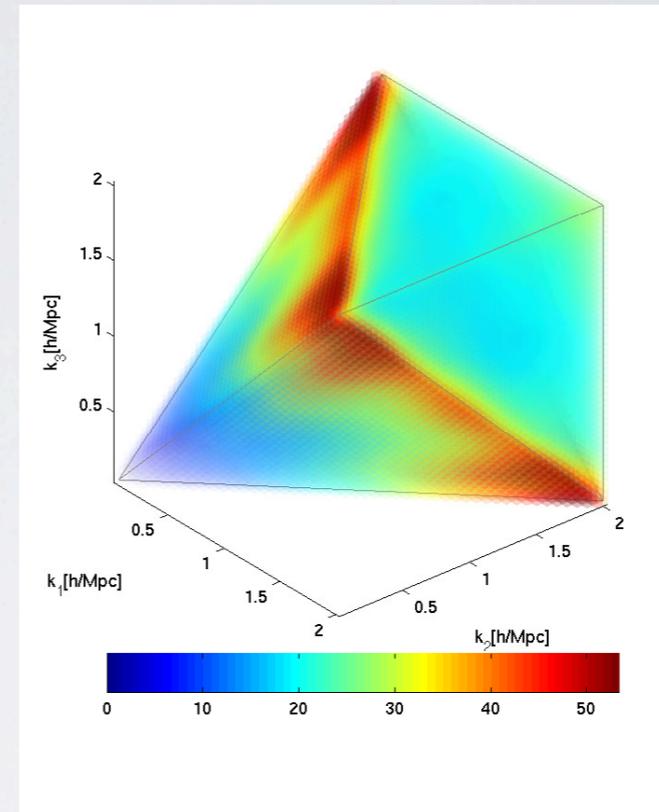
Name	NG shape	f_{NL}	$L[\frac{\text{Mpc}}{h}]$	N_p	z_i	$L_s[\frac{\text{kpc}}{h}]$	N_r	glass
G512	—	—	1600	512	49	156	3	no
G512g	—	—	1600	512	49	156	3	yes
G768	—	—	2400	768	19	90	3	no
G1024	—	—	1875	1024	19	40	2	no
Loc10	local	10	1600	512	49	156	3	no
Loc10g	local	10	1600	512	49	156	3	yes
Eq100	equil	100	1600	512	49	156	3	no
Eq100g	equil	100	1600	512	49	156	3	yes
Orth100	orth	100	1600	512	49	156	3	no
Orth100g	orth	100	1600	512	49	156	3	yes
Orth100 ⁻	orth	-100	1600	512	49	156	3	no
Flat10	flat	10	1600	512	49	156	3	no



Evolving 'Gaussian' Bispectrum

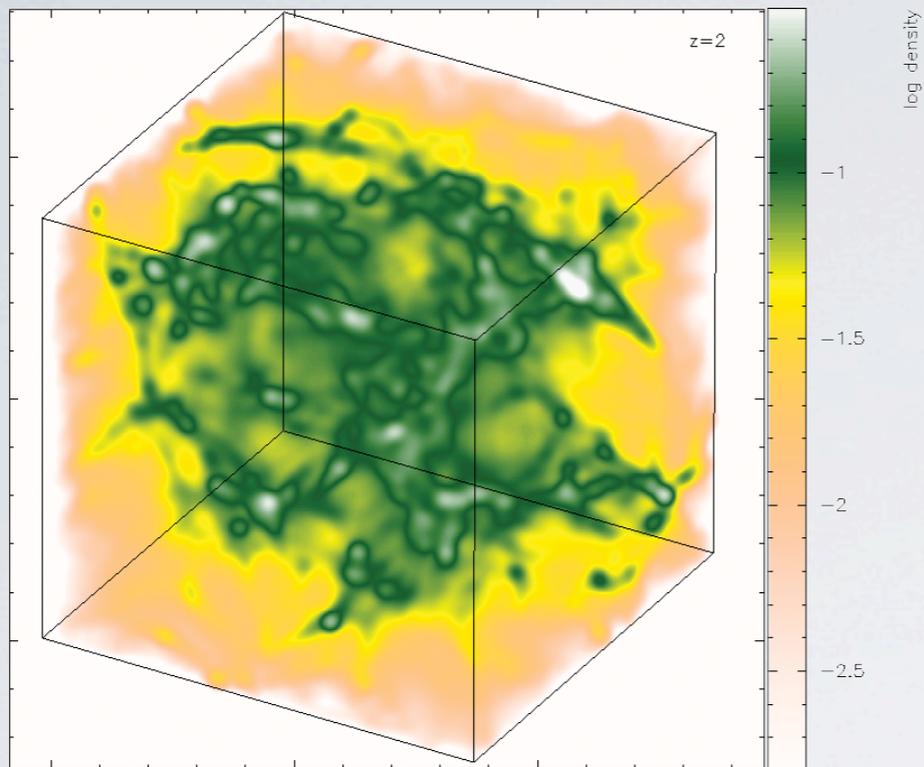


(a) Dark matter, $z = 4$

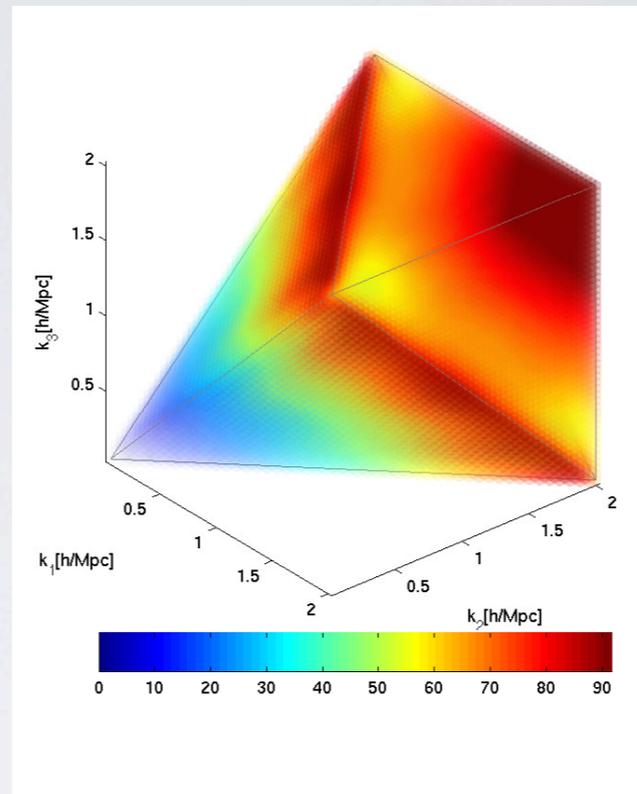


(b) Bispectrum signal, $z = 4$

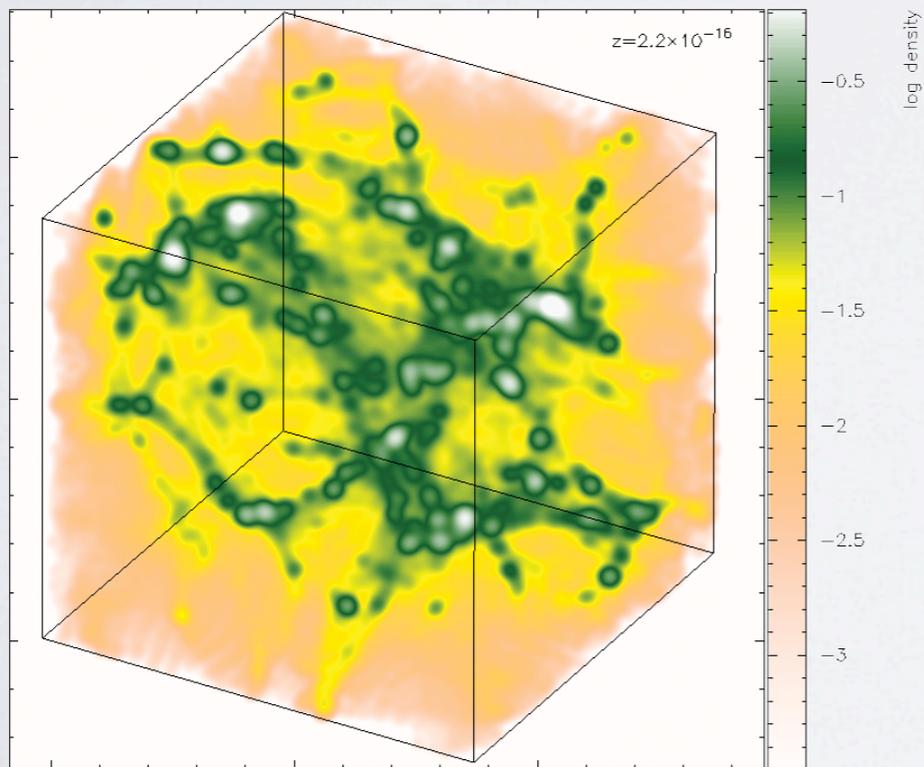
Early blob-like structure \longrightarrow Flattened bispectrum



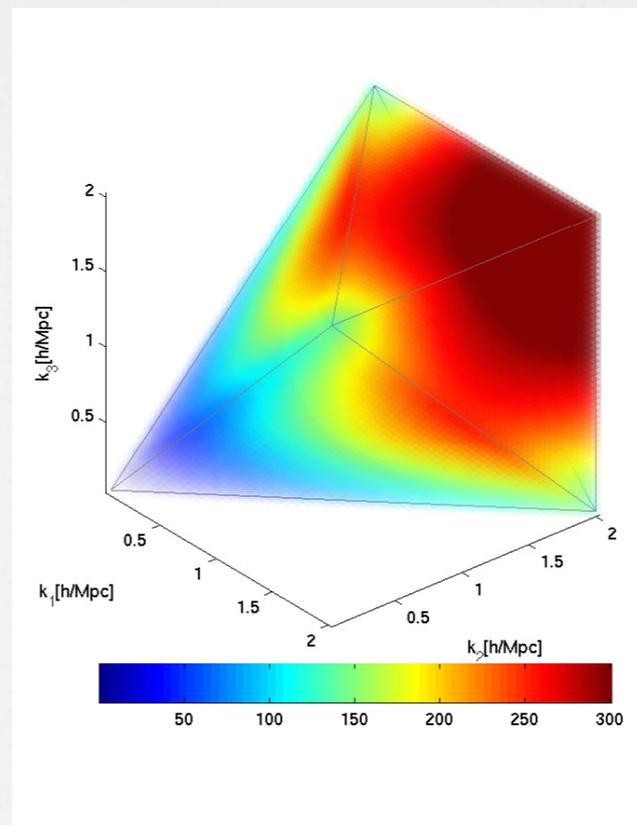
(c) Dark matter, $z = 2$



(d) Bispectrum signal, $z = 2$

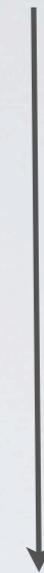


(e) Dark matter, $z = 0$

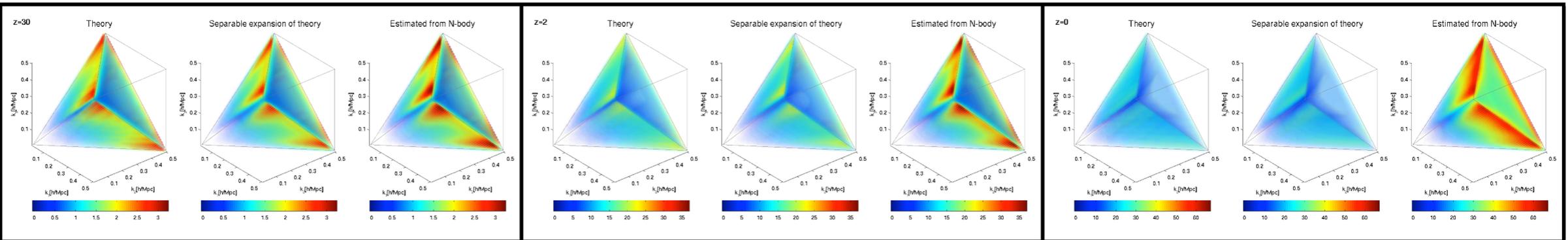
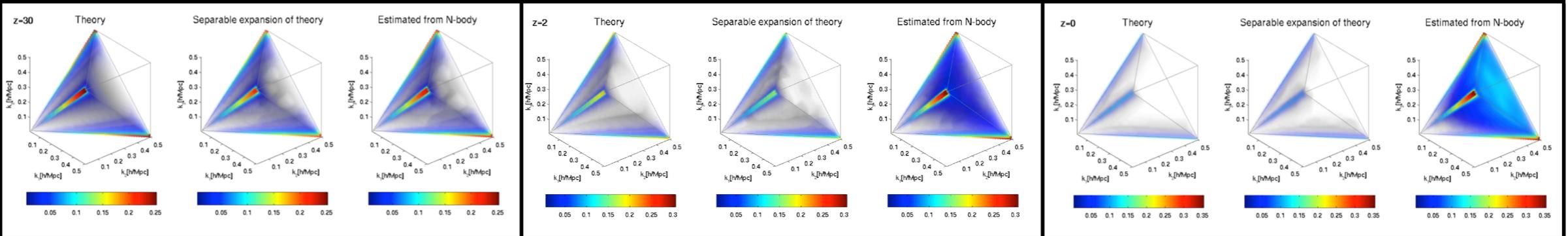
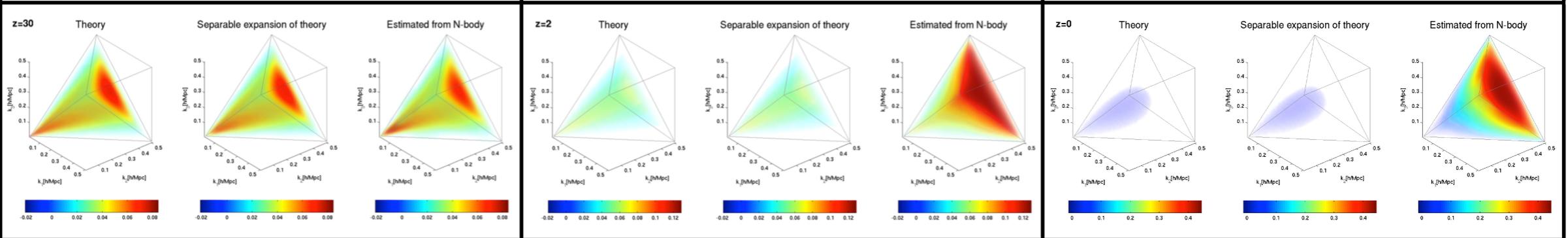
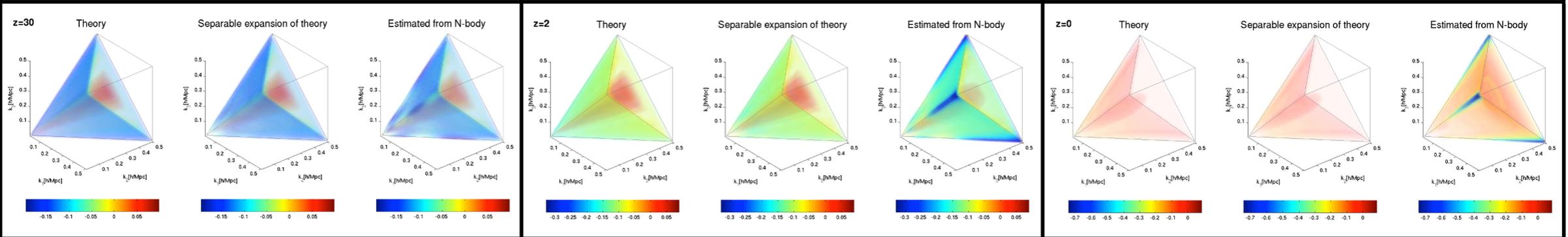
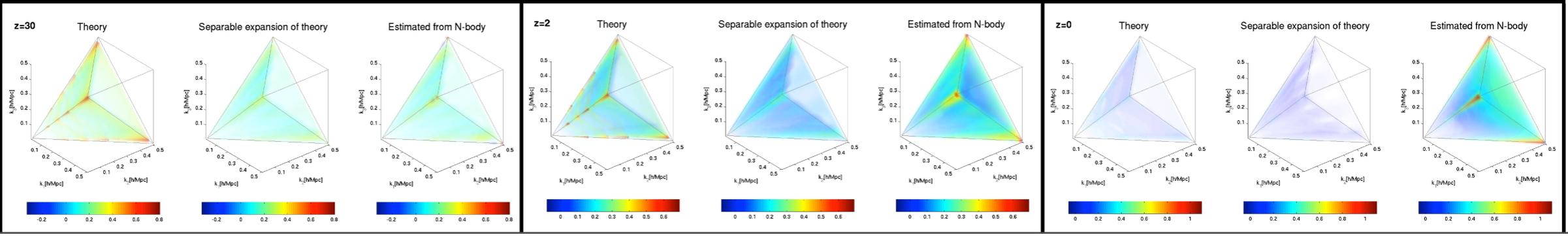


(f) Bispectrum signal, $z = 0$

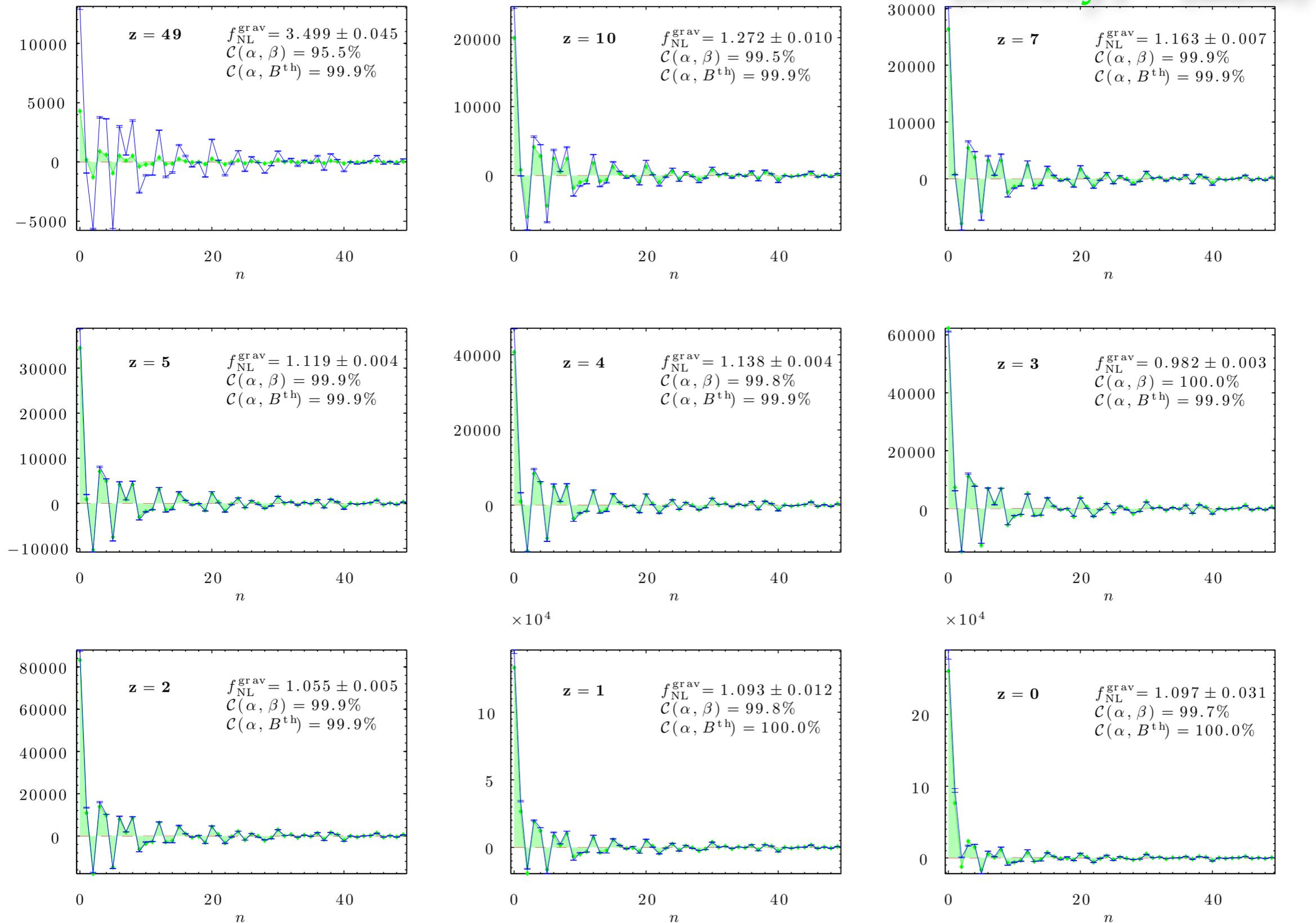
Filaments and clusters appearing



Looking more and more equilateral

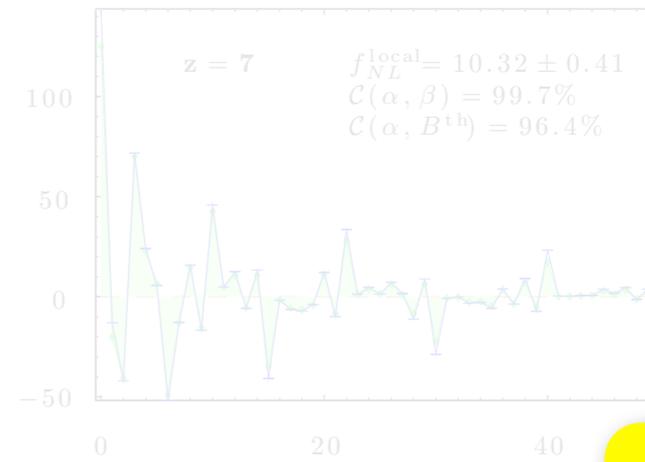
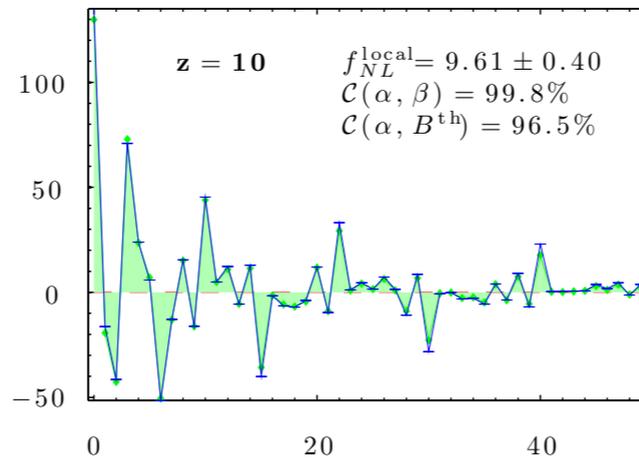
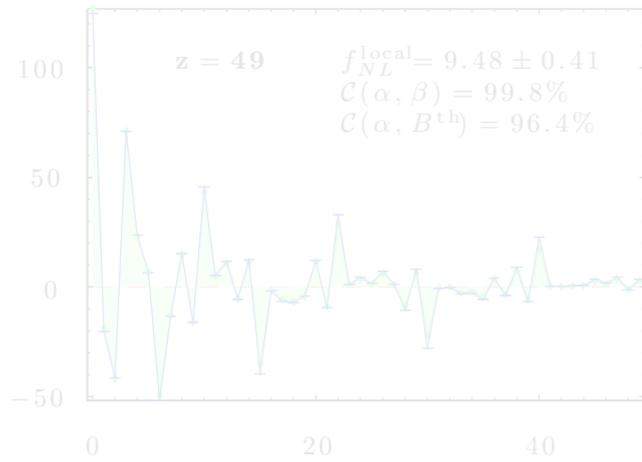
$z = 30$ $z = 2$ $z = 0$ Gravity
(Gaussian ICs)local
 $f_{NL}^{loc} = 10$ equilateral
 $f_{NL}^{eq} = 100$ orthogonal
 $f_{NL}^{orth} = 100$ flattened
 $f_{NL}^{flat} = 10$ 

Gravitational bispectrum, Gaussian initial conditions



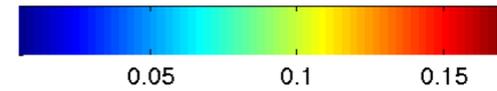
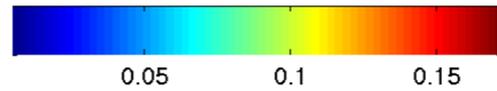
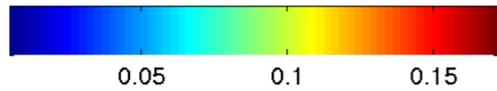
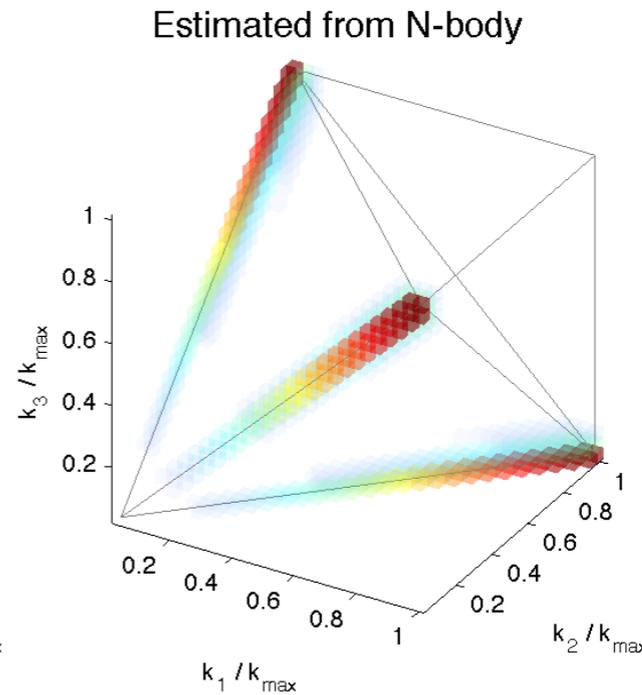
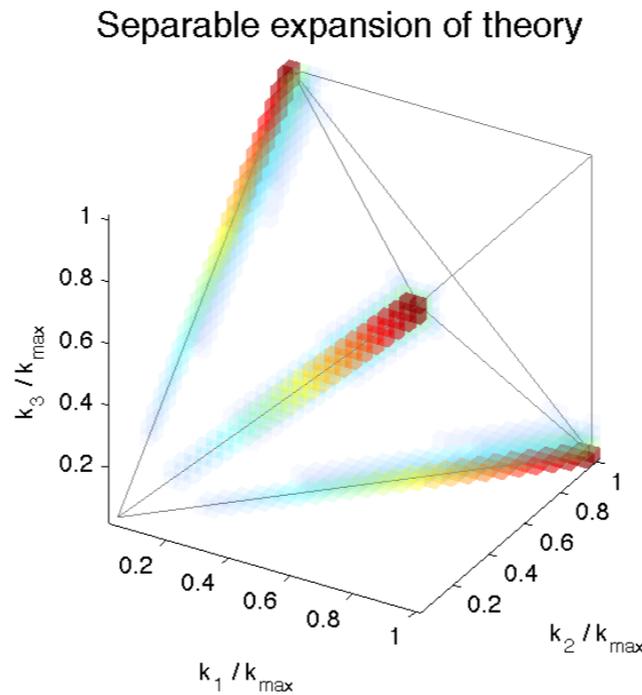
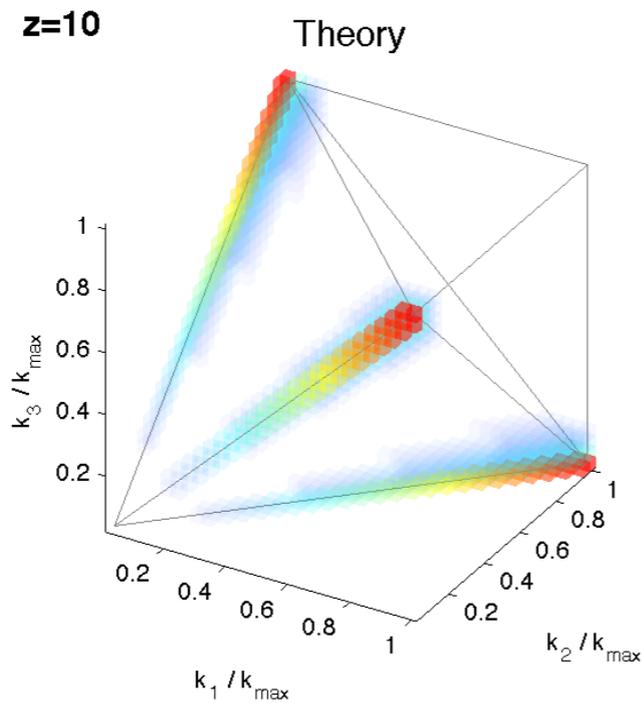
3 realisations of 512^3 particles in a $L = 1600$ Mpc/h box with $z_{\text{init}} = 49$ and $k = 0.004 - 0.5$ h/Mpc

Non-Gaussian initial conditions $f_{NL}^{local} = 10$, plot $(\alpha_n^R)_{tree}^{primordial}$ and $\langle \beta_n^R \rangle - \langle (\beta_n^R)_{Gaussian} \rangle$



z=10

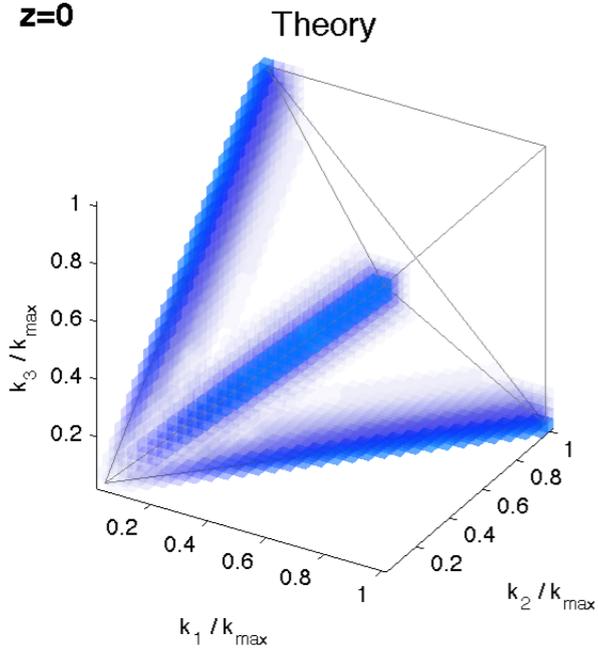
plot $\sqrt{k_1 k_2 k_3} B_\delta(k_1, k_2, k_3) / \sqrt{P_\delta(k_1) P_\delta(k_2) P_\delta(k_3)}$



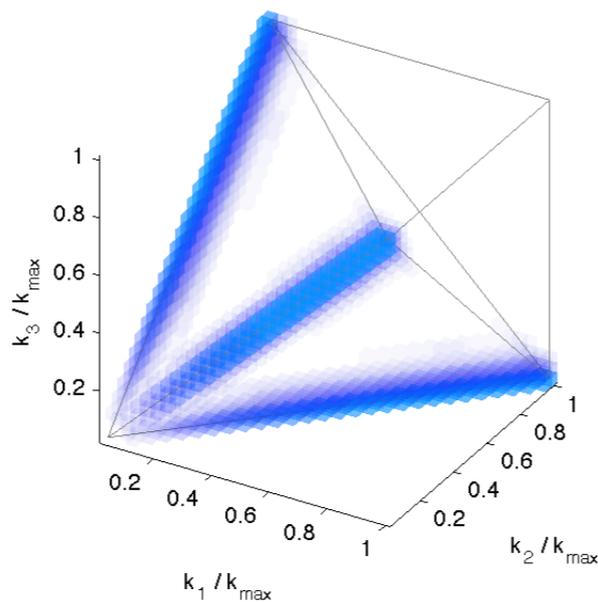
3 realisations of 512^3 particles in a $L = 1600$ Mpc/h box with $z_{init} = 49$ and $k = 0.004 - 0.5$ h/Mpc

Non-Gaussian initial conditions $f_{NL}^{local} = 10$, plot $(\alpha_n^R)_{tree}^{primordial}$ and $\langle \beta_n^R \rangle - \langle (\beta_n^R)_{Gaussian} \rangle$

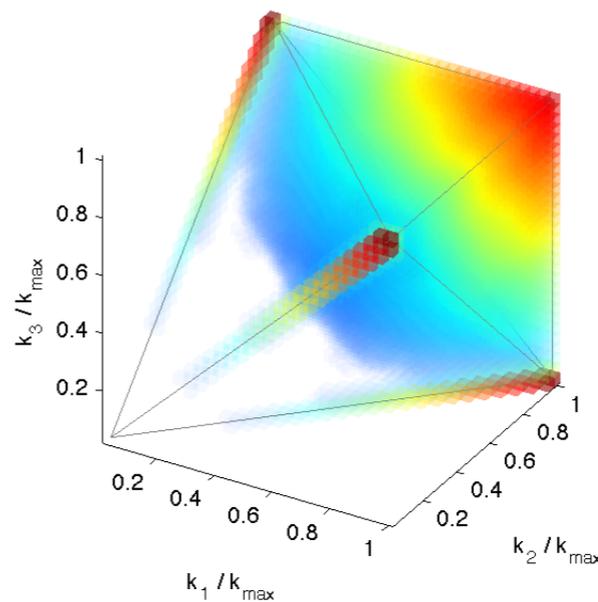
$z=0$



Separable expansion of theory

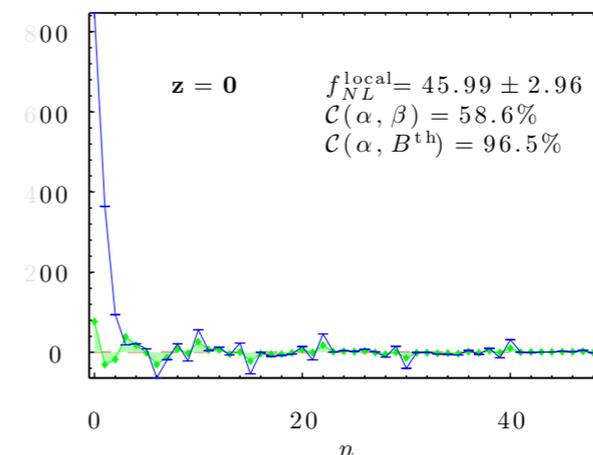
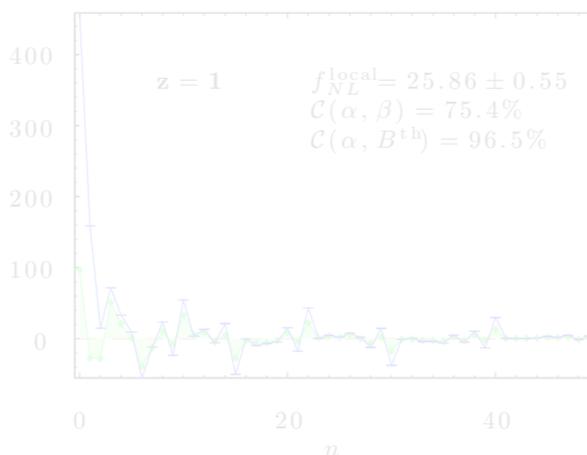
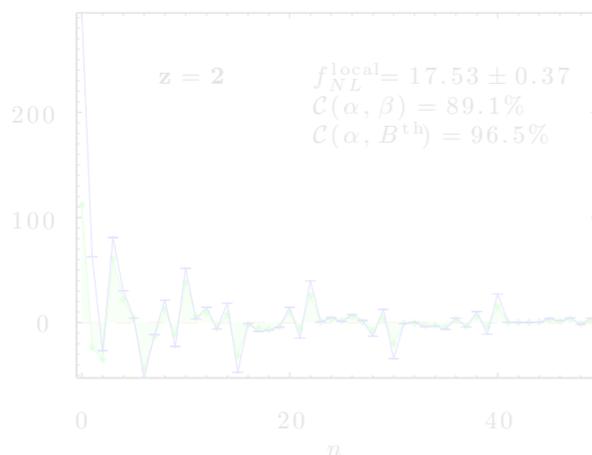


Estimated from N-body



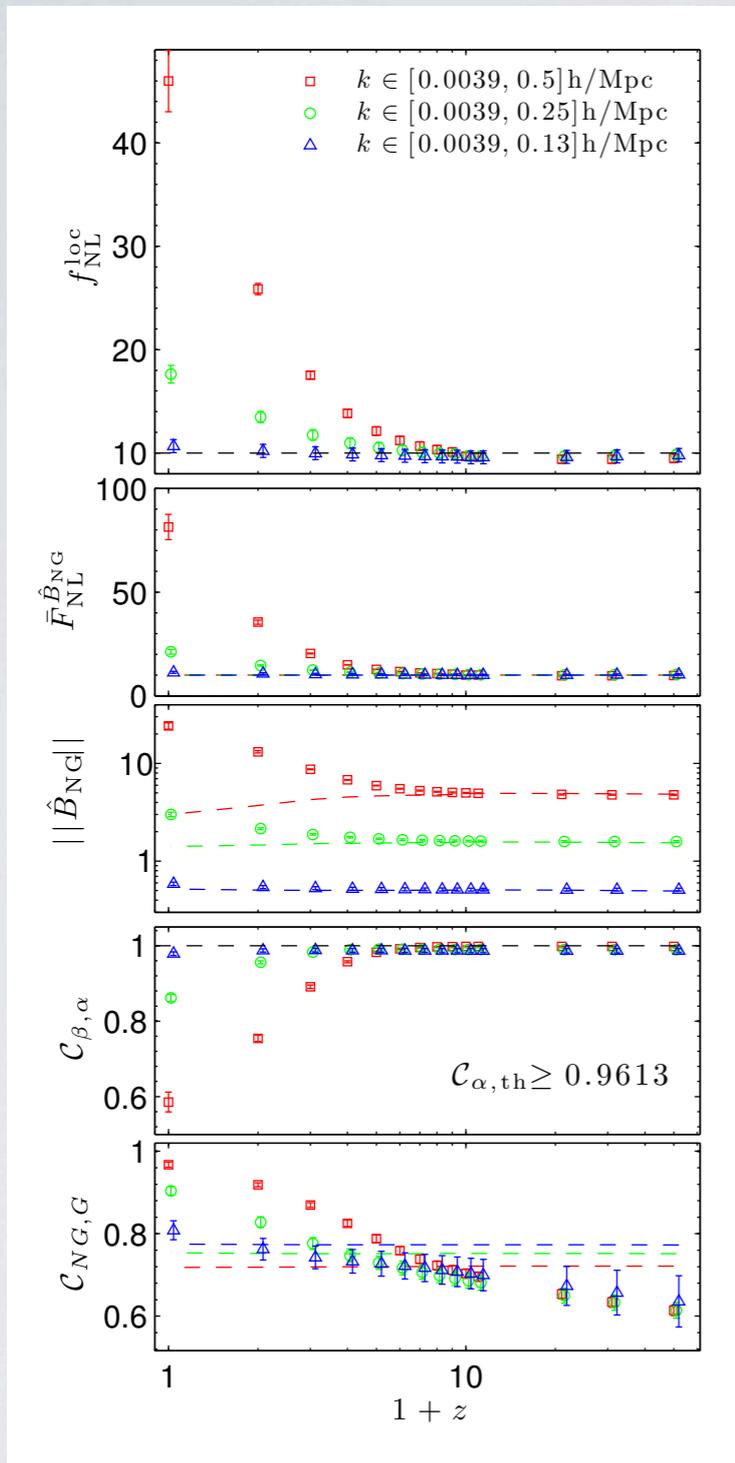
$z=0$

plot $\sqrt{k_1 k_2 k_3 B_\delta(k_1, k_2, k_3)} / \sqrt{P_\delta(k_1) P_\delta(k_2) P_\delta(k_3)}$

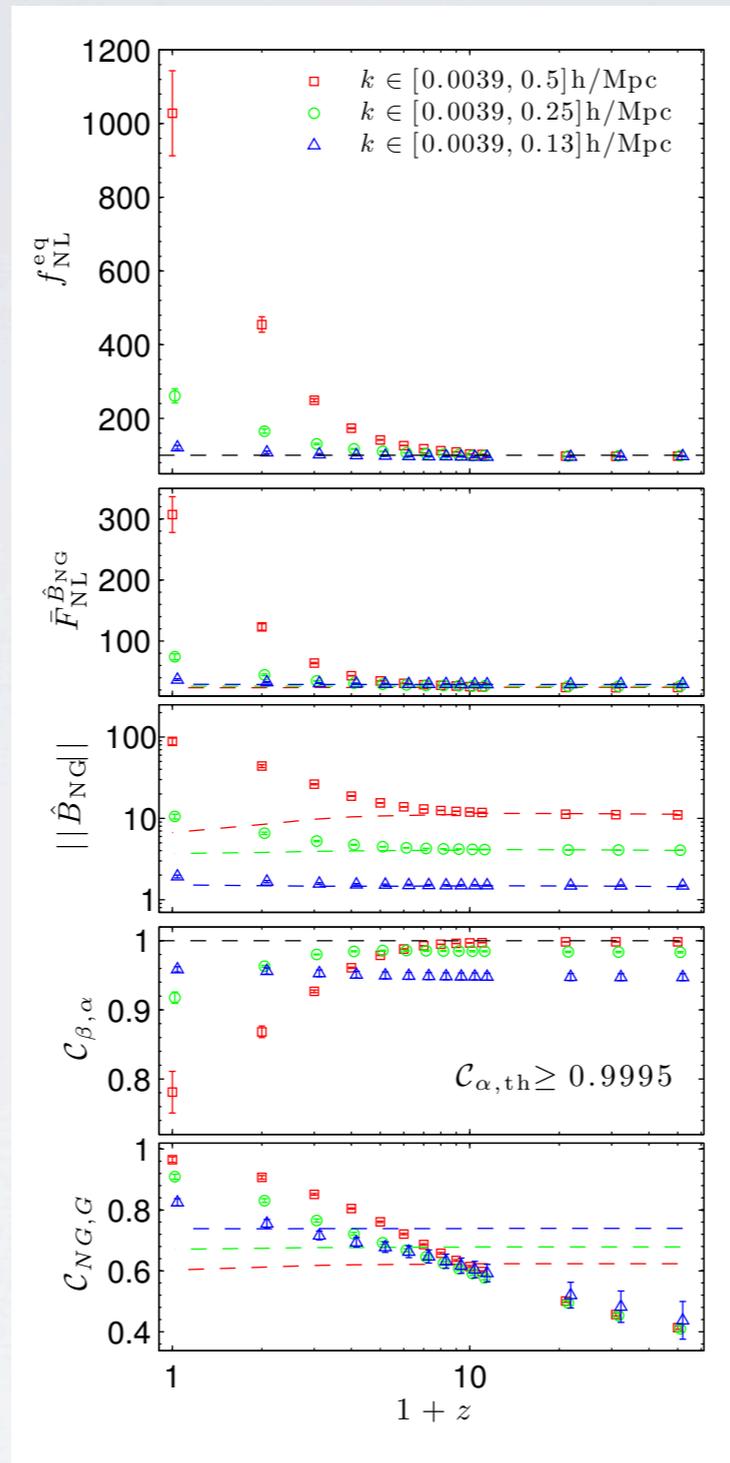


3 realisations of 512^3 particles in a $L = 1600$ Mpc/h box with $z_{init} = 49$ and $k = 0.004 - 0.5$ h/Mpc

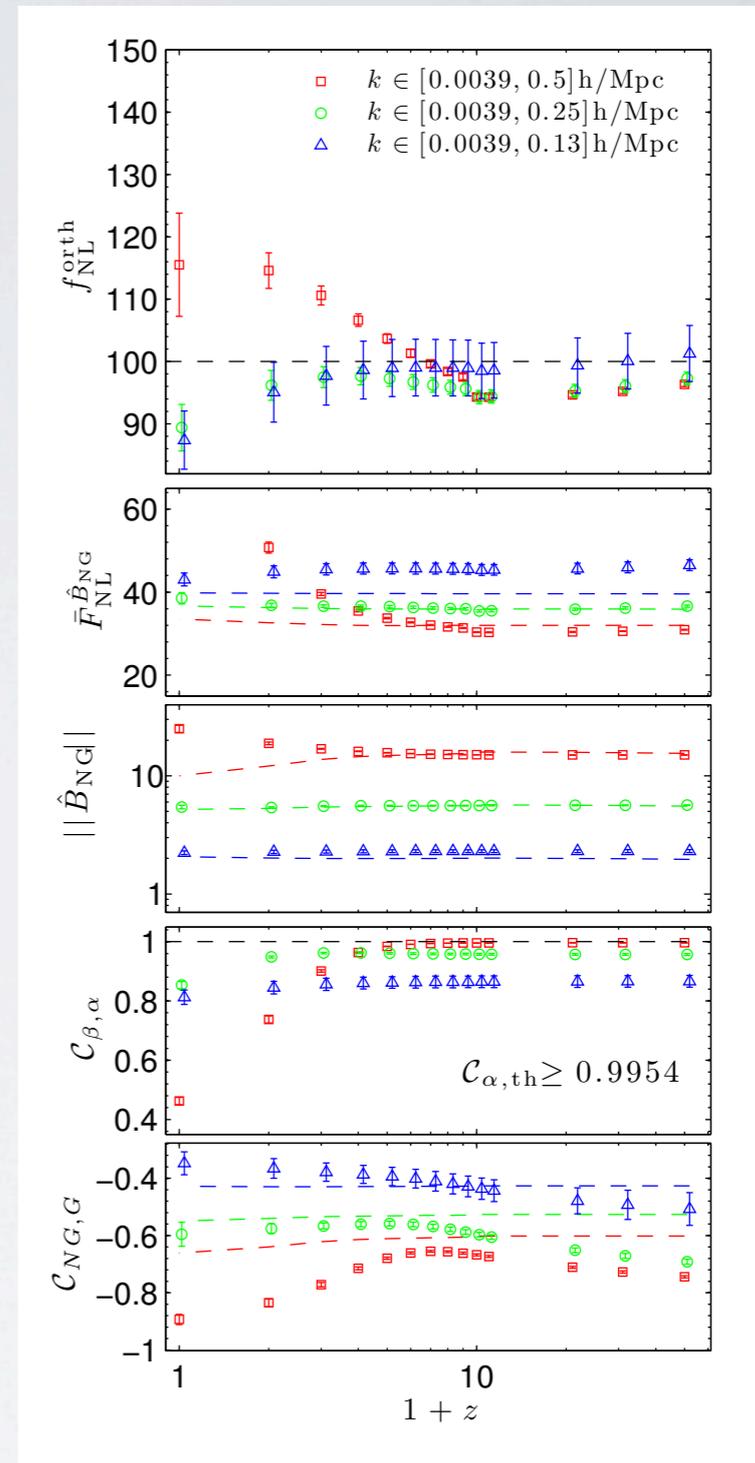
Comparing to the tree level result



(a) Local



(b) Equilateral



(c) Orthogonal

COMPARISON AGAINST LOOP LEVEL RESULTS & PHENOMENOLOGICAL MODELS

Gaussian case

Tree level

$$B_{\delta}^{\text{grav}}(k_1, k_2, k_3) = 2P_{\delta}^L(k_1)P_{\delta}^L(k_2)F_2^{(s)}(\mathbf{k}_1, \mathbf{k}_2) + 2 \text{ perms}$$

$$\text{where } F_2^{(s)}(\mathbf{k}_1, \mathbf{k}_2) = \frac{5}{7} + \frac{1}{2} \frac{\mathbf{k}_1 \cdot \mathbf{k}_2}{k_1 k_2} \left(\frac{k_1}{k_2} + \frac{k_2}{k_1} \right) + \frac{2}{7} \left(\frac{\mathbf{k}_1 \cdot \mathbf{k}_2}{k_1 k_2} \right)^2$$

Phenomenological fit (Gil-Marín et al arXiv:1111.4477)

$$F_2^{(s)\text{eff}}(\mathbf{k}_1, \mathbf{k}_2) = \frac{5}{7} a(n_1, k_1) a(n_2, k_2) + \frac{1}{2} \frac{\mathbf{k}_1 \cdot \mathbf{k}_2}{k_1 k_2} \left(\frac{k_1}{k_2} + \frac{k_2}{k_1} \right) b(n_1, k_1) b(n_2, k_2) + \frac{2}{7} \left(\frac{\mathbf{k}_1 \cdot \mathbf{k}_2}{k_1 k_2} \right)^2 c(n_1, k_1) c(n_2, k_2), \quad 99\% \text{ correlation to N-body result}$$

Non-Gaussian case

Tree level $B_{\delta}^{\text{prim}}(k_1, k_2, k_3; z) =$
 $M(k_1, z)M(k_2, z)M(k_3, z)B_{\Phi}(k_1, k_2, k_3)$

... < 60% correlation for $k > 0.1$ h/Mpc

Loop Level Result (Sefusatti et al arXiv:1111.6966) $B = B_{112}^{II} + B_{122}^I + B_{122}^{II} + B_{113}^I + B_{113}^{II}$

$$B_{112}^{II} = \int \frac{d^3q}{(2\pi)^3} F_2(\mathbf{q}, \mathbf{k}_3 - \mathbf{q}) T_{\delta}^L(\mathbf{k}_1, \mathbf{k}_2, \mathbf{q}, \mathbf{k}_3 - \mathbf{q}), \quad (17)$$

$$B_{112}^I = F_2^{(s)}(\mathbf{k}_1, \mathbf{k}_2) [P_{\delta}^L(k_1) P_{12}(k_2) + k_1 \leftrightarrow k_2] \\ + 2 \text{ perms}, \quad (18)$$

$$B_{122}^{II} = 4 \int \frac{d^3q}{(2\pi)^3} F_2^{(s)}(\mathbf{q}, \mathbf{k}_2 - \mathbf{q}) F_2^{(s)}(\mathbf{k}_1 + \mathbf{q}, \mathbf{k}_2 - \mathbf{q}) \times \\ B_{\delta}^L(k_1, q, |\mathbf{k}_1 + \mathbf{q}|) P_{\delta}^L(|\mathbf{k}_2 - \mathbf{q}|) + 2 \text{ perms}, \quad (19)$$

$$B_{113}^I = 3 B_{\delta}^L(k_1, k_2, k_3) \int \frac{d^3q}{(2\pi)^3} F_3^{(s)}(\mathbf{k}_3, \mathbf{q}, -\mathbf{q}) P_{\delta}^L(q) \\ + 2 \text{ perms}, \quad (20)$$

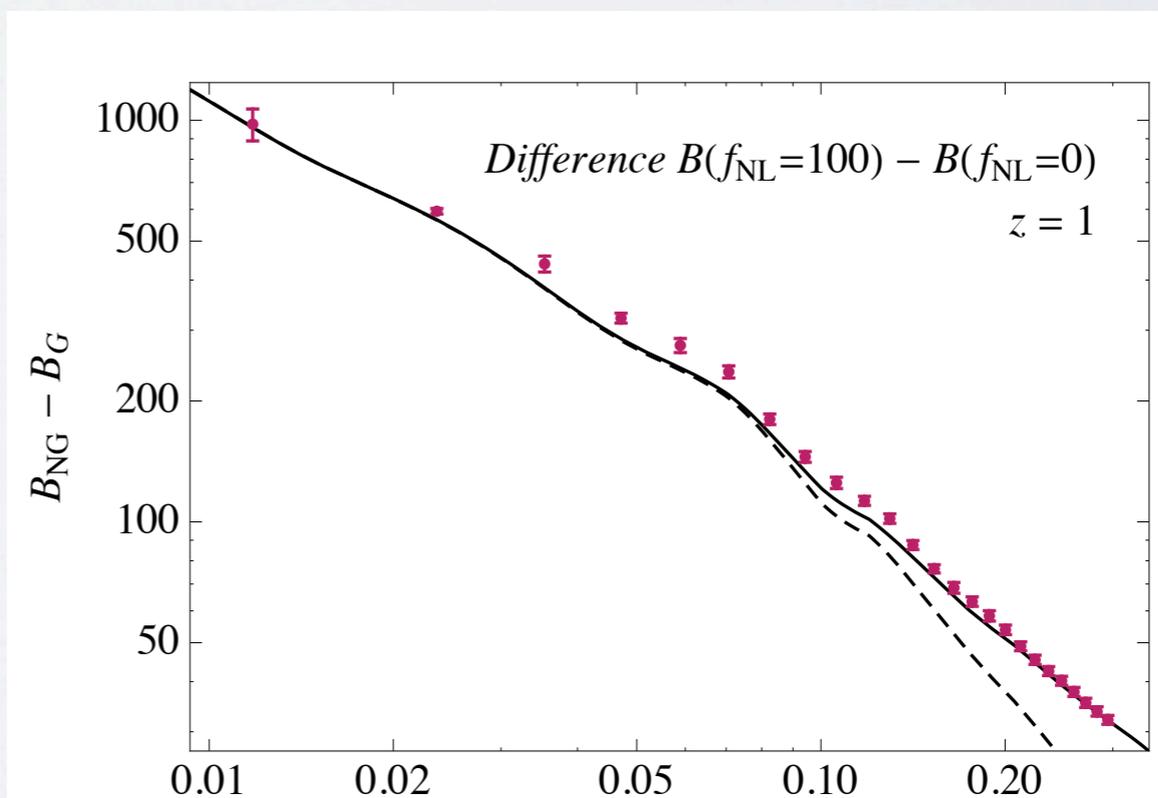
$$B_{113}^{II} = 3 P_{\delta}^L(k_1) \int \frac{d^3q}{(2\pi)^3} F_3^{(s)}(\mathbf{k}_1, \mathbf{q}, \mathbf{k}_2 - \mathbf{q}) \times \\ B_{\delta}^L(k_2, q, |\mathbf{k}_2 - \mathbf{q}|) + (k_1 \leftrightarrow k_2) + 2 \text{ perms}, \quad (21)$$

where

$$P_{12}(k) = 2 \int \frac{d^3q}{(2\pi)^3} F_2^{(s)}(\mathbf{q}, \mathbf{k} - \mathbf{q}) B_{\delta}^L(k, q, |\mathbf{k} - \mathbf{q}|).$$

Model	z	Tree $k_{\max} = 0.5$	Tree+Loop (0.5)	Tree $k_{\max} = 0.25$	Tree+Loop (0.25)	Tree $k_{\max} = 0.125$	Tree+Loop (0.125)
Local	0	58%, 66.6	79%, 12.3	84%	84%	94%, 11	98%, 10.2
	1	87%, 29.2	87%, 8.5	95%	98%	96%, 10.6	97%, 10.5
	2	89%, 25.9	92%, 13.8	98%	99%	96%, 10.5	97%, 10.4
	3	96%, 20.8	96%, 13.8	99%	100%	96%, 10.3	97%, 10.3
Equil	0	78%	84%	91%	95%	96%	96%
	1	87%	91%	96%	97%	95%	97%
	2	92%	96%	98%	98%	95%	95%
	3	96%	98%	98%	99%	95%	95%
Flat	0	48%, 47.45	88%, 11.1	76%	98%	90%, 10.3	95%, 9.8
	1	66%, 30.3	94%, 11.8	90%	98%	92%, 10.2	93%, 10.1
	2	83%, 22.9	98%, 12.1	94%	97%	94%, 10.3	94%, 10.1
	3	92%, 19.1	99%, 12.2	95%	97%	92%, 10.2	93%, 10.1
Orthog	0	47%	70%	87%	87%		
	1	74%	85%	95%	96%		
	2	90%	95%	96%	97%		
	3	96%	98%	96%	96%		

Previous analysis only
particular slices=>



- Beyond perturbation theory... The Halo Model

Assumes all dark matter resides in halos

Ingredients: (1) Halo mass function $n(m, z)$
 (2) Halo density profile $\hat{\rho}(k_1, m, z)$
 (3) Halo bias functions $b_1(m_1)$

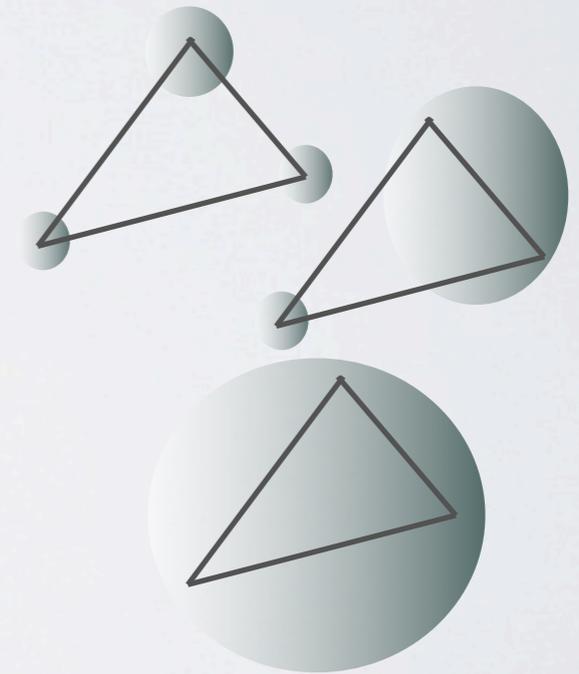
(Figueroa et al arXiv:1205.2015)

$$B(k_1, k_2, k_3) = B_{3h}(k_1, k_2, k_3) + B_{2h}(k_1, k_2, k_3) + B_{1h}(k_1, k_2, k_3)$$

$$B_{3h}(k_1, k_2, k_3, z) = \frac{1}{\bar{\rho}^3} \left[\prod_{i=1}^3 \int d m_i n(m_i, z) \hat{\rho}(m_i, z, k_i) \right] B_h(k_1, m_1; k_2, m_2; k_3, m_3; z),$$

$$B_{2h}(k_1, k_2, k_3, z) = \frac{1}{\bar{\rho}^3} \int d m n(m, z) \hat{\rho}(m, z, k_1) \int d m' n(m', z) \hat{\rho}(m', z, k_2) \hat{\rho}(m', z, k_3) \\ \times P_h(k_1, m, m', z) + \text{cyc.},$$

$$B_{1h}(k_1, k_2, k_3, z) = \frac{1}{\bar{\rho}^3} \int d m n(m, z) \hat{\rho}(k_1, m, z) \hat{\rho}(k_2, m, z) \hat{\rho}(k_3, m, z).$$



where

$$B_h(k_1, m_1; k_2, m_2; k_3, m_3; z) = b_1(m_1) b_1(m_2) b_1(m_3) B(k_1, k_2, k_3) \\ + [b_1(m_1) b_1(m_2) b_2(m_3) P(k_1) P(k_2) + \text{cyc.}]$$

Gaussian case: Correlation $> 99.2\%$ out to $k=2$ h/Mpc
($z=0,1,2$)

Local non-Gaussian case:

Correlation $> 97.5\%$ out to $k=2$ h/Mpc
($z=0,1,2$)

but ... overestimates amplitude beyond $k=0.5$ h/Mpc.
Need to divide the I-halo term by factor of 4.

Simple phenomenological fitting formulae

$$B_{\delta}^{\text{fit}}(k_1, k_2, k_3) \equiv B_{\delta, \text{NL}}^{\text{grav}} + c_1 \tilde{D}(z)^{d_1} B_{\delta}^{\text{const}}$$

where $B_{\delta}^{\text{const}}(k_1, k_2, k_3) \equiv (k_1 + k_2 + k_3)^{-1.7}$

99.3% correlation to halo model at 2h/Mpc

Simulation	c_1	d_1	$\min_{z \leq 20}(\mathcal{C}_{\beta, \alpha})$	$\mathcal{C}_{\beta, \alpha}(z = 0)$
G ₄₀₀ ⁵¹²	1.0×10^7	8	99.8%	99.8%
G512g	4.1×10^6	7	99.8%	99.8%
Loc10	2.4×10^3	6	99.7%	99.7%
Eq100	8.6×10^2	6	97.9%	99.4%
Flat10	1.1×10^4	6	98.8%	98.9%
Orth100	-2.6×10^2	6	90.5%	90.5%

Primordial Non-Gaussianity

Time shifted model? $B_{\delta, \text{const}}^{\text{grav}}(k_1, k_2, k_3)|_z \equiv c_1 \bar{D}^{n_h}(z) (k_1 + k_2 + k_3)^\nu$

Constant component of primordial bispectrum \longleftrightarrow One-Halo term

 Faster growth of structure than tree-level

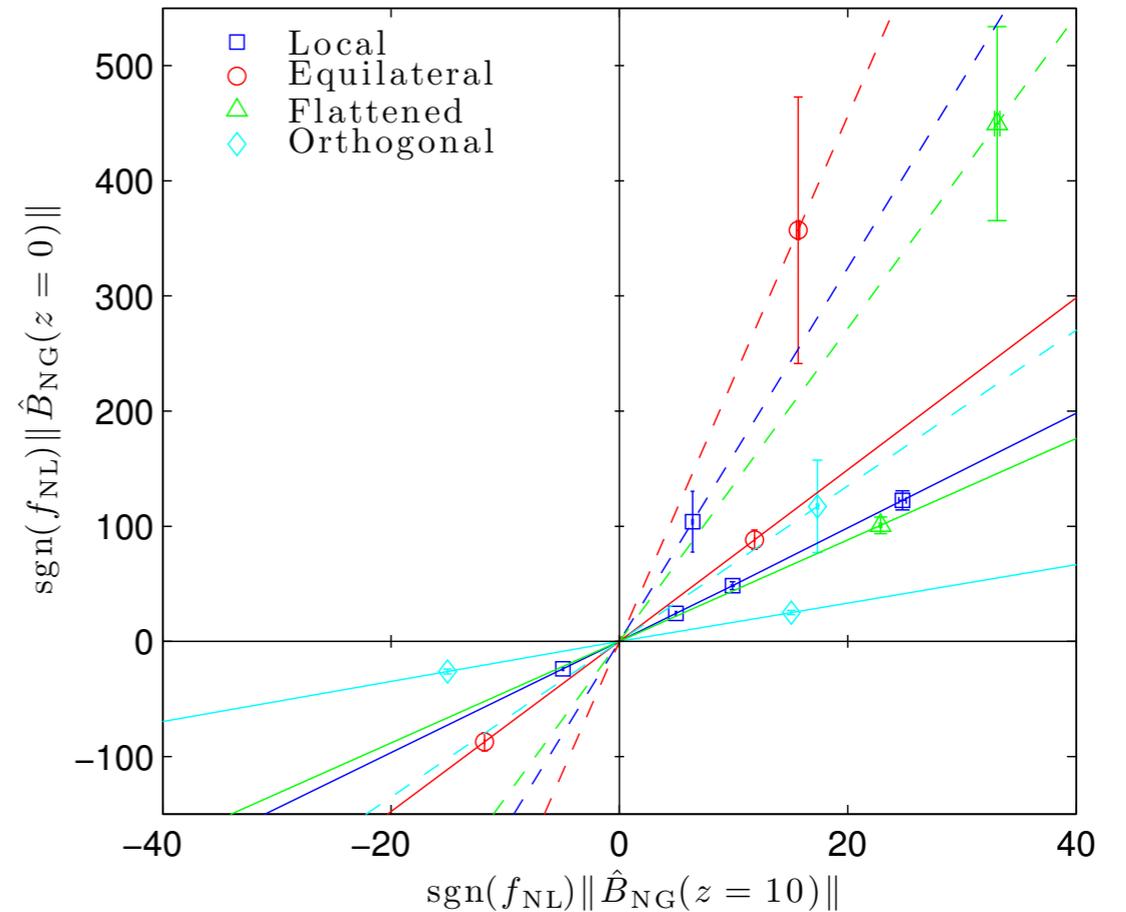
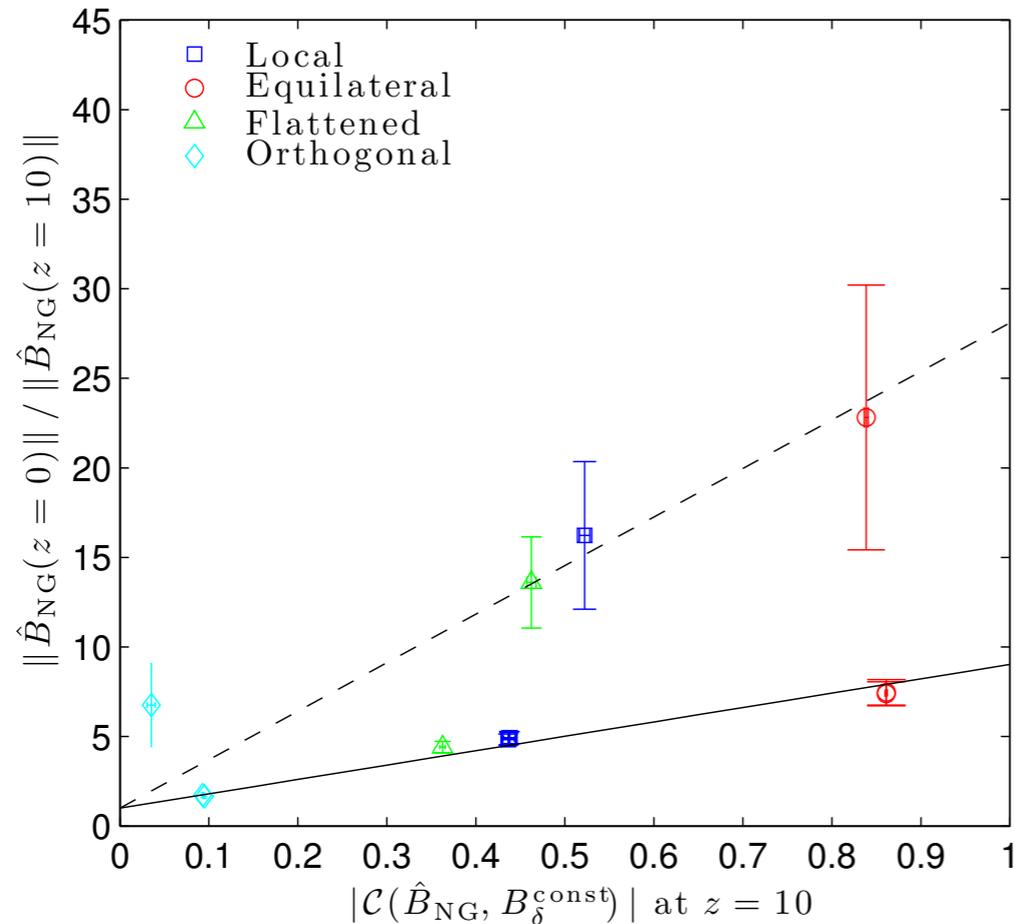
$$\bar{D}^{n_h}(z + \Delta z) \approx \bar{D}^{n_h}(z) + n_h \bar{D}^{n_h-1}(z) \frac{d\bar{D}(z)}{dz} \Delta z.$$

$$(\hat{f}_{\text{NL}}^{\text{const}})_{\text{Gauss}}(z + \Delta z) = (\hat{f}_{\text{NL}}^{\text{const}})_{\text{NG}}(z)$$

 $B_{\text{NG}}^{\text{fit}}(k_1, k_2, k_3) \equiv f_{\text{NL}} [B_{\delta, \text{NL}}^{\text{prim}} + B_{\delta, \text{const}}^{\text{prim}}]$

$$B_{\delta, \text{const}}^{\text{prim}}(k_1, k_2, k_3) \equiv c_2 \bar{D}^{n_h^{\text{prim}}}(z) (k_1 + k_2 + k_3)^\nu$$

$$\|B_{\text{NG}}(z_{\text{late}})\| \propto \mathcal{C}(B_{\text{NG}}(z_{\text{early}}), B_{\delta, \text{const}}^{\text{grav}}(z_{\text{early}})) \times \|B_{\text{NG}}(z_{\text{early}})\|.$$



Message: To predict the relative *growth* of the *bispectrum* only need to correlate to *constant* model at *early* times

CONCLUSIONS

- Can efficiently produce reliable general non-Gaussian **initial conditions** for N-body simulations
- Can more accurately analyse the results of N-body simulations
- More complete comparison to tree-level, loop-level and phenomenological predictions
- Next step... application to the galaxy bispectrum