

Onsager relations in a two-dimensional electron gas with spin-orbit coupling

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Spin-orbit in a 2DEG

- What for?
- Spin-Hall effect(s)
- Role of disorder

G. Vignale, J. Supercond. Nov. Magn. (2010)

T. Jungwirth, J. Wunderlich and K. Olejník, Nat. Mat. (2012)

Charge vs. spin

- Spin currents
- Onsager relations

L. Y. Wang, A. G. Mal'shukov, C. S. Chu, PRB (2012)

C. Gorini, R. Raimondi, P. Schwab, arXiv (2012)

Electronics (charge) \longrightarrow Spintronics (charge&spin)

Why the spin?

- Long lifetime, low power consumption
- Multi-functional (semiconductor) devices
- Quantum computing

Nat. Mat. Insight (2012)

The main question&some possible answers

How to manipulate spins in solid state systems?

- Magnetic fields/materials
- Optical methods
- Spin-orbit coupling

Spin-orbit in a 2DEG

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Nat. Mat. Insight (2012)

The main question&some possible answers

Spin-orbit coupling \Rightarrow

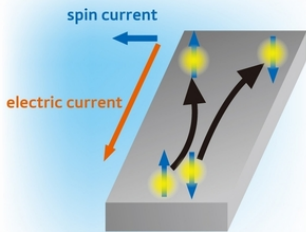
- 1 **Electrical** handle on spin
- 2 Spin **not** conserved

Spin-orbit in a 2DEG

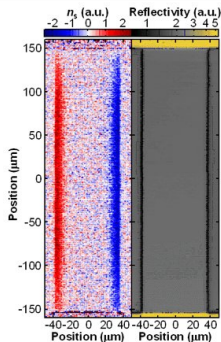
The spin-Hall effect

An electric field \mathbf{E} generates a transverse (pure) spin current

M. I. Dyakonov and V. I. Perel, Phys. Lett. A (1971)



Physicsworld (2006)



Y. K. Kato et al., Science (2004)

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Y. K. Kato et al., Science (2004)

J. Wunderlich et al., PRL (2005)

S. O. Valenzuela and M. Tinkham, Nature (2006)

E. Saitoh et al., Appl. Phys. Lett. (2006)

T. Seki et al., Nat. Mat. (2008)

K. Ando et al., J. Appl. Phys. (2009)

E. S. Garlid et al., PRL (2010)

K. Olejník et al., PRL (2012)

...

The spin-Hall effect

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M. I. Dyakonov and V. I. Perel, Phys. Lett. A (1971)

Intrinsic Spin Hall effect in the Rashba model

$$\mathbf{J}_y^{S_z} = \sigma_{SH} E_x$$

Sinova et al., PRL (2004)

- No disorder $\Rightarrow \sigma_{SH} = \frac{e}{8\pi}$ “universal” result (interesting)
- Any disorder $\Rightarrow \sigma_{SH} = 0$ “universal” result (boring)

Disruptive role of impurities

Spin-orbit in a 2DEG

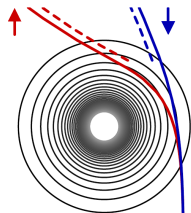
Intrinsic origin

- Crystal structure (broken bulk inversion symmetry)
- Device structure (broken structure inversion symmetry)

Extrinsic origin

- Mott skew-scattering
- Side-jump

Constructive role of impurities



Adapted from Engel et al., PRL (2005)

Intrinsic vs. extrinsic

Nontrivial dependence on physical parameters

Spin-orbit in a 2DEG

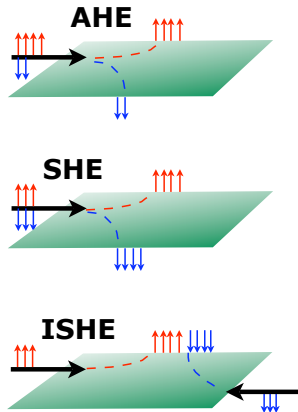
- a **spin-up current** generates a current in the transverse direction

$$\delta j_{y\uparrow} = 2\gamma_{\uparrow} j_{x\uparrow}$$

- a **spin-down current** generates a current in the transverse **opposite** direction

$$\delta j_{y\downarrow} = -2\gamma_{\downarrow} j_{x\downarrow}$$

$\gamma_{\uparrow}, \gamma_{\downarrow} \leftrightarrow$ Intrinsic and/or extrinsic



The main points

- Spin non-conservation
- Fundamental role of disorder
- Nontrivial interplay intrinsic-extrinsic

Two (related) questions

- Definition of the spin current?
- Reciprocity spin-Hall \leftrightarrow inverse spin-Hall?

2D effective Hamiltonian

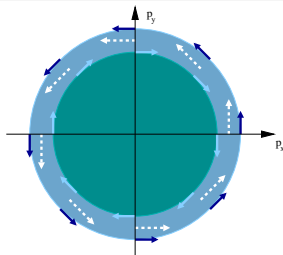
$$H = \underbrace{p^2/2m + H_{\text{intr}}}_{\text{band/device structure}} + \underbrace{V_{\text{imp}} + H_{\text{extr}}}_{\text{impurities}}$$

$$H_{\text{intr}} = -\mathbf{b}(\mathbf{p}) \cdot \boldsymbol{\sigma}, \quad H_{\text{extr}} = -\frac{\lambda^2}{4\hbar} [\mathbf{p} \times \nabla V_{\text{imp}}(\mathbf{r})] \cdot \boldsymbol{\sigma}$$

$$\epsilon_{\pm}(\mathbf{p}) = \frac{p^2}{2m} \pm |\mathbf{b}(\mathbf{p})|$$

$$|\mathbf{b}| \sim \text{meV} \sim 10^{-1} \epsilon_F$$

$$(\lambda/\lambda_c)^2 \sim 10^6$$



2D effective Hamiltonian

$$H = \underbrace{p^2/2m + H_{\text{intr}}}_{\text{band/device structure}} + \underbrace{V_{\text{imp}} + H_{\text{extr}}}_{\text{impurities}}$$

$$H_{\text{intr}} = -\mathbf{b}(\mathbf{p}) \cdot \boldsymbol{\sigma}, \quad H_{\text{extr}} = -\frac{\lambda^2}{4\hbar} [\mathbf{p} \times \nabla V_{\text{imp}}(\mathbf{r})] \cdot \boldsymbol{\sigma}$$

Spin current

$$J_i^a \equiv \frac{1}{2} \{s^a, v_i\} \quad \Rightarrow \quad \partial_t s^a + \nabla \cdot \mathbf{J}^a = T^a$$

$$v_i = \nabla_{p_i} H, \quad a(\text{spin}), i(\text{space}) = x, y, z$$

Spin non-conservation \Rightarrow “spin-torque” T^a

Spin-orbit \leftrightarrow non-Abelian gauge fields

$$H = \frac{p^2}{2m} + \alpha (p_y \sigma^x - p_x \sigma^y) = \frac{(p + \eta \mathcal{A}^a \sigma^a / 2)^2}{2m} + \text{const.}$$

$$\eta \mathcal{A}_y^x = -\eta \mathcal{A}_x^y = 2m\alpha, \quad \text{else } \mathcal{A}_i^a = 0$$

Advantages

- Unambiguous definition of \mathbf{J}_s : “colour” current
- Spatially/time-modulated spin-orbit or Zeeman fields
- Interplay intrinsic-extrinsic

But...

Limited to linear-in-momentum spin-orbit (?)

N. J. Fröhlich and U. M. Studer, Rev. Mod. Phys. (1993), I. V. Tokatly, PRL (2008)

C. Gorini et al., PRB (2010)

Charge vs. spin

Charge

$$\Phi, \mathbf{A} \Rightarrow \mathbf{E}, \mathbf{B}, \mathbf{J}_c$$

Spin

$$\Psi, \mathcal{A} \Rightarrow \mathcal{E}, \mathcal{B}, \mathbf{J}_s$$

$$\tilde{\partial}_t \mathbf{J}_s + \tilde{\nabla} \cdot \mathbf{J}_s = 0, \quad \mathbf{J}_s = \mathbf{J}_{s,\text{intr}} + \mathbf{J}_{s,\text{extr}}$$

$$\tilde{\partial}_t = \partial_t \underbrace{-i\eta[\Psi, \dots]}_{\text{Zeeman}} \quad \tilde{\nabla} = \nabla \underbrace{+i\eta[\mathcal{A}, \dots]}_{\text{spin-orbit}}$$

Hall effects

- $\mathbf{J}_c = -\frac{e\tau}{m} \mathbf{J}_c \times \mathbf{B}$ Hall
- $\mathbf{J}_s = -\frac{\eta\tau}{4m} \mathbf{J}_c \times \mathcal{B}$ Spin-Hall
- $\mathbf{J}_c = -\frac{\eta\tau}{m} \mathbf{J}_s \times \mathcal{B}$ Inverse spin-Hall

Gorini et al., PRB (2010), Raimondi et al., Ann. Phys. (Berlin) (2012)

Onsager relations

- Original formulation: fluxes and forces

L. Onsager, Phys. Rev. (1931)

- Microscopic approach: Kubo formula

N. Nagaosa et al., Rev. Mod. Phys. (2010)

$$\delta H = \mathbf{J}_\beta \cdot \mathbf{A}_\beta \quad \Rightarrow \quad \delta \mathbf{J}_\alpha = \underbrace{\langle \mathbf{J}_\alpha \mathbf{J}_\beta \rangle}_{K_{\alpha\beta}} \mathbf{A}_\beta$$

Time-reversal: $\mathbf{J}_\alpha \rightarrow \epsilon_\alpha \mathbf{J}_\alpha$, $\epsilon_\alpha = \pm 1$

$$K_{\alpha\beta} = \epsilon_\alpha \epsilon_\beta K_{\beta\alpha}$$

α, β = charge, spin, heat...

$$\alpha = \mathbf{s} \text{ (spin)}, \beta = \mathbf{c} \text{ (charge)}$$

$$\text{Time-reversal: } \mathbf{J}_S \rightarrow \mathbf{J}_S, \mathbf{J}_C \rightarrow -\mathbf{J}_C \Rightarrow K_{SC} = -K_{CS}$$

Spin-Hall response

$$\delta \mathbf{J}_S = \underbrace{\langle \mathbf{J}_S \mathbf{J}_C \rangle}_{K_{SC}} \mathbf{A}, \quad \delta \mathbf{J}_C = \underbrace{\langle \mathbf{J}_C \mathbf{J}_S \rangle}_{-K_{SC}} \mathcal{A}$$

“Spin gauge” dependent reciprocal relations

- Charge conservation: $\mathbf{J}_C \cdot \mathbf{A} = -e\Phi_C$
- Spin non-conservation: $\mathbf{J}_S \cdot \mathcal{A} \neq -\eta\Psi \Rightarrow \sigma_{SH,A} \neq \sigma_{SH,\Phi}$

“Spin-gauge” \leftrightarrow experimental setup

C. Gorini, R. Raimondi and P. Schwab, arXiv (2012)

$$H = \frac{p^2}{2m} + H_{\text{intr}} + V_{\text{imp}} + H_{\text{extr}}$$

Optical injection (homogeneous)

$$\delta H_1 = \sum_i [J_i e A_i + J_i^Z \eta \mathcal{A}_i^Z]$$

$$J_y = -\sigma_{sH} \eta \mathcal{E}_x^Z \Leftrightarrow J_x^Z = \sigma_{sH} e E_y$$

Scalar potentials (inhomogeneous)

$$\delta H_2 = -e\Phi - \eta\Psi$$

$$J_y = -\bar{\sigma}_{sH} \eta \mathcal{E}_x^Z \Leftrightarrow \bar{J}_x^Z = \bar{\sigma}_{sH} e E_y$$

J_x^Z = standard (non-conserved), \bar{J}_x^Z = conserved

To sum up

- Spin-orbit \Rightarrow electrical handle on spin
- Fundamental role of impurities (intrinsic vs. extrinsic)
- Spin-charge Onsager: dependence on experimental setup

Something for the future

- Extension of non-Abelian formalism
- Spin-thermoelectric effects

G. Vignale, J. Supercond. Nov. Magn. (2010)

T. Jungwirth, J. Wunderlich and K. Olejník, Nat. Mat. (2012)

C. Gorini, R. Raimondi and P. Schwab, arXiv (2012)

$$\mathbf{J}^0 = -D(\nabla n + 2eN_0\mathbf{E}) - D^a([\tilde{\nabla}s]^a + \frac{\eta N_0}{2}\boldsymbol{\varepsilon}^a) +$$

$$-\frac{e\tau}{m}\mathbf{J}^0 \wedge \mathbf{B} - \frac{\eta\tau}{m}\mathbf{J}^a \wedge \mathcal{B}^a$$

$$\mathbf{J}^a = -D^a(\nabla n + 2eN_0\mathbf{E}) - D([\tilde{\nabla}s]^a + \frac{\eta N_0}{2}\boldsymbol{\varepsilon}^a) +$$

$$-\frac{e\tau}{m}\mathbf{J}^a \wedge \mathbf{B} - \frac{\eta\tau}{4m}\mathbf{J}^0 \wedge \mathcal{B}^a,$$

$$D = \text{charge diffusion constant}, \quad D^a = \frac{D'}{n}s^a$$

Rashba+disorder

$$\sigma_{sH}(\omega) = -\frac{\gamma\sigma}{e} \left[\frac{-i\omega + 1/\tau_s}{-i\omega + 1/\tau_{DP} + 1/\tau_s} \right]$$

$$\bar{\sigma}^{sH}(\omega) = -\frac{\gamma\sigma}{e} \left[\frac{-i\omega}{-i\omega + 2/\tau_{DP}} \right] \left[\frac{-i\omega - 1/\tau_{DP} + 1/\tau_s}{-i\omega + 1/\tau_{DP} + 1/\tau_s} \right]$$

$$\gamma = \gamma_{\text{intr}} + \gamma_{\text{extr}}, \quad 1/\tau_{DP} = (2m\alpha)^2 D, \quad 1/\tau_s = 1/\tau (\lambda p_F/4)^4$$

$$\mathbf{J}_x^z = \mathbf{s}^z \left(\frac{1}{i\hbar} [x, H] \right), \quad \bar{\mathbf{J}}_x^z = \frac{1}{i\hbar} [s^z x, H]$$