

QUANTUM CORRELATIONS ACROSS THE DE SITTER HORIZON

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INTRODUCTION

➤ **Our aim is to study the possible quantum influence across cosmological horizons**

- **We construct a model of a spacetime with two regions, separated by a cosmological horizon**

- **We will find an exact quantization of the this spacetime**

➤ **The chosen model derives from de Sitter spacetime**

➤ **The result is closely related with black hole models and quantum wormholes**

THE MODEL

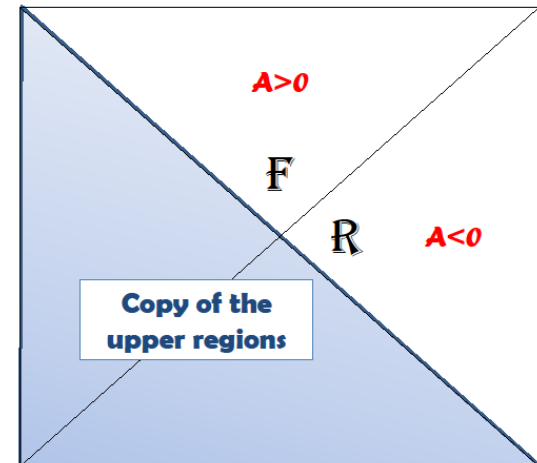
➤ We consider a model provided with a cosmological constant and described by this metric

$$ds^2 = -\frac{N^2(T)}{A(T)}dT^2 + A(T)dR^2 + b^2(T)d\Omega_2^2$$

Classically

• $A > 0 \rightarrow b^2 > \Lambda/3$ Static de Sitter spacetime

• $A < 0 \rightarrow b^2 < \Lambda/3$ Kantowski-Sachs spacetime



Penrose Diagram

CONFIGURATION SPACE

➤ **The Lagrangian in each region $i=R,F$ is given by**

$$L_i(N_i, A_i, b_i, \dot{A}_i, \dot{b}_i) = \frac{1}{2} \left(\frac{b_i \dot{b}_i \dot{A}_i}{N_i} + \frac{A_i \dot{b}_i^2}{N_i} - N_i + \Lambda b^2 N_i \right)$$

➤ **Identifying an appropriately rescaled lapse function N' , in each region, we can write the action in the Hamiltonian formalism as**

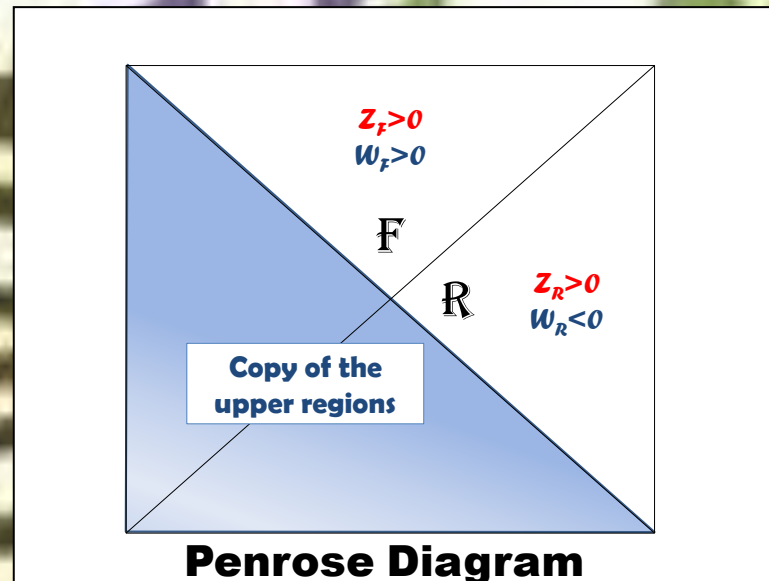
$$S = \sum_i \int dt \left[\Pi_{A_i} \dot{A}_i + \Pi_{b_i} \dot{b}_i - N' H_i \right]$$

➤ **New configuration variables**

$$w_i = A_i b_i^2 \left| \frac{\Lambda b_i^2}{3} - 1 \right|, \quad z_i = \frac{A_i}{\left(\frac{\Lambda b_i^2}{3} - 1 \right)}$$

▪ **The classical solutions are**

$$z_i = 1$$



HAMILTONIAN CONSTRAINT

$$H = [w_R \partial_{w_R} w_R \partial_{w_R} + \frac{|w_R|}{4} + w_F \partial_{w_F} w_F \partial_{w_F} + \frac{|w_F|}{4} - z_R \partial_{z_R} z_R \partial_{z_R} - z_F \partial_{z_F} z_F \partial_{z_F}]$$

➤ We choose the kinematical Hilbert space to be the tensor product spanned by

$$\phi_{k_R}(z_R) \phi_{k_F}(z_F) \varphi_{K_R}(w_R) \varphi_{K_F}(w_F)$$

with

$$k_R, k_F, K_R, K_F > 0$$

where

$$\begin{aligned} \phi_{k_i}(z_i) &= e^{\pm i k_i \ln z_i} \\ \varphi_{K_i}(w_i) &= C_{K_i} J_{2iK_i}(\sqrt{|w_i|}) + C'_{K_i} J_{-2iK_i}(\sqrt{|w_i|}) \end{aligned}$$

PHYSICAL HILBERT SPACE

INDEPENDENT QUANTIZATION OF THE REGIONS

➤ We consider the Hilbert space as $\mathcal{H} = \mathcal{H}^R \oplus \mathcal{H}^F$

➤ Defining $\alpha_i = \ln |w_i|$, $\beta_i = \ln z_i$

The Wheeler-De Witt equation for each region reads

$$\left[-\partial_{\alpha_i}^2 + \partial_{\beta_i}^2 - \frac{e^{\alpha_i}}{4} \right] \phi_{k_i}(z_i) \varphi_{K_i}(w_i) = 0$$

where $k_i = K_i$

➤ The physical states are given by

$$\Psi = \prod_{i=R,F} \int dk_i C_{k_i} \phi_{k_i}(z_i) \varphi_{k_i}(w_i)$$

QUANTIZATION WITH CORRELATIONS BETWEEN THE REGIONS

➤ We consider the Hilbert space as

$$\mathcal{H} = \mathcal{H}^R \otimes \mathcal{H}^F$$

➤ The Wheeler-De Witt equation takes the form

$$\left[-\partial_{\alpha_R}^2 + \partial_{\beta_R}^2 - \frac{e^{\alpha_R}}{4} - \partial_{\alpha_F}^2 + \partial_{\beta_F}^2 - \frac{e^{\alpha_F}}{4} \right] \Psi = 0$$

➤ $\ln z_i$ behaves like an internal time in each region

➤ We define a common internal time ($\ln z$) with the identification

$$\ln z_R = \ln z_F$$

So now the states are spanned by

$$\phi_k(z) \varphi_{K_R}(w_R) \varphi_{K_F}(w_F)$$

- **The physical Hilbert space is characterized by the dispersion relation**

$$2k^2 = K_R^2 + K_F^2$$

- **The physical states are given by**

$$\Psi = \int dK_R dK_F C_{K_R K_F} \phi_k(z) \varphi_{K_R}(w_R) \varphi_{K_F}(w_F)$$

- **The solutions $k=0$ that imply $K_i=0$ have the asymptotic form**

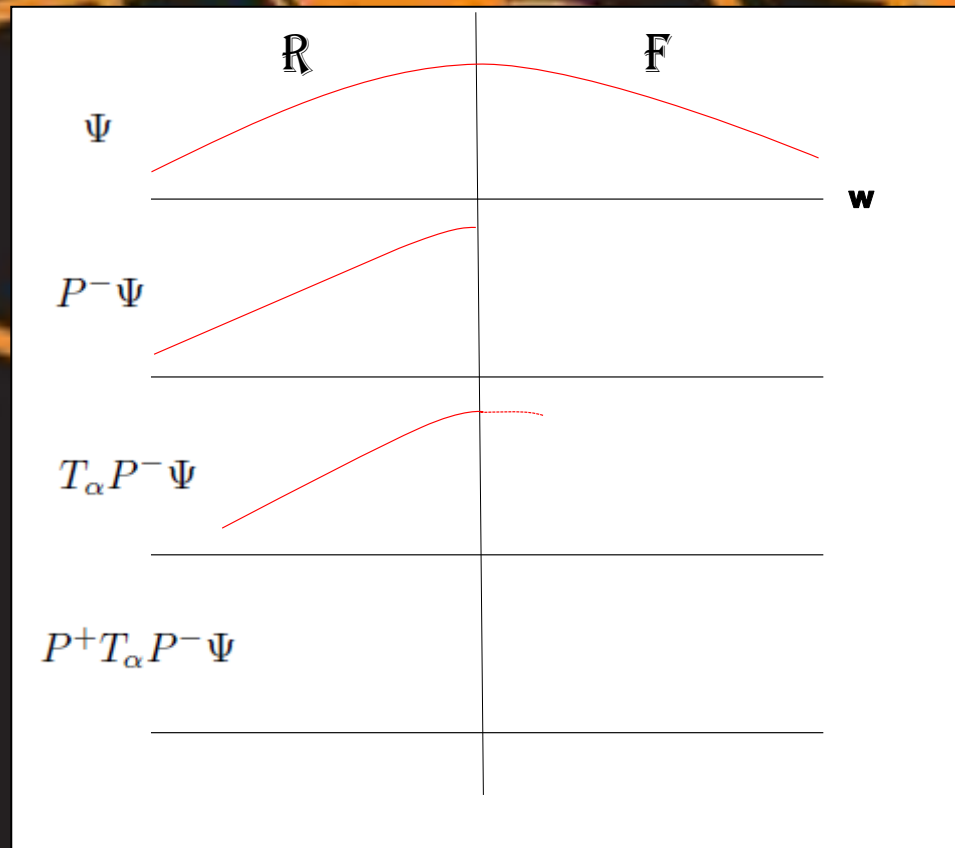
$$\Psi \sim (|w_R| |w_F|)^{-1/4} \cos \sqrt{|w_R|} \cos \sqrt{|w_F|}$$

which is sharply peaked around the semiclassical solutions

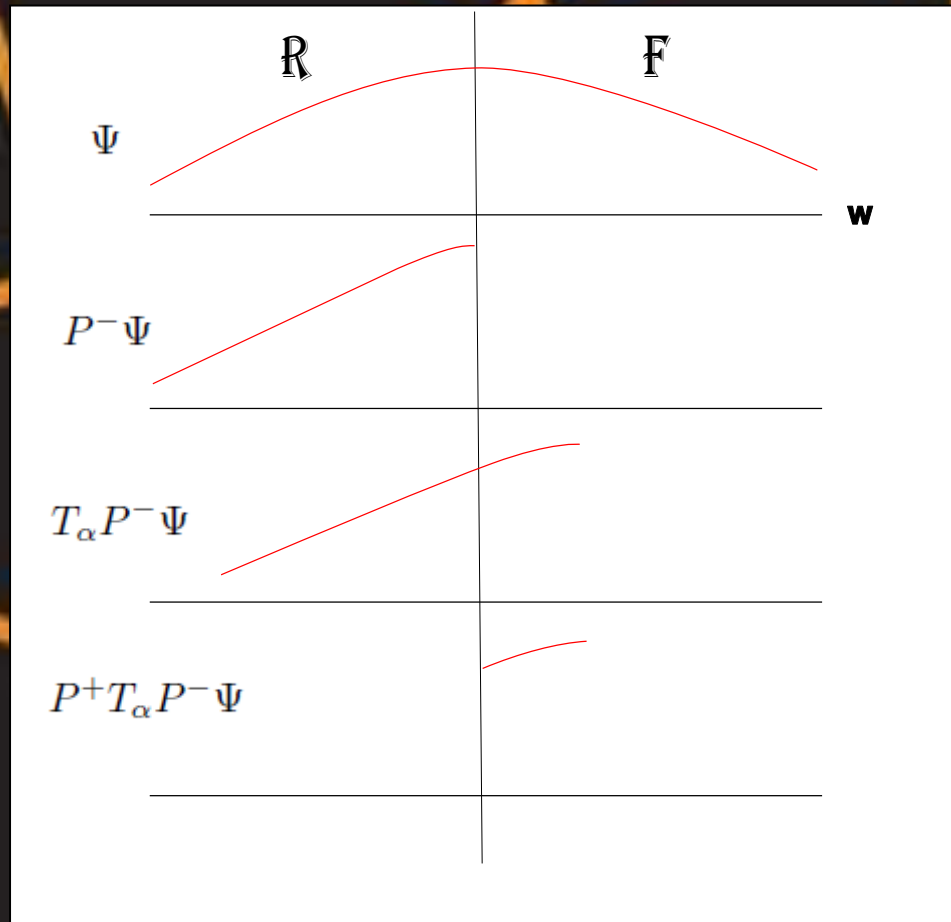
QUANTUM CORRELATIONS

➤ We define a translation operator (T_α) in the variable w

INDEPENDENT QUANTIZATION OF THE REGIONS



QUANTIZATION WITH CORRELATIONS BETWEEN THE REGIONS



➤ We have a new contribution coming from the region R in the region F

CONCLUSIONS

- **Exact quantization of the (maximal) extension of the de Sitter spacetime**
- **Recover semiclassical behavior**
- **We construct a quantization that establishes quantum correlations between both regions, mixing them**
- **We can find effects in one region by tracing out the other one, which is not accessible**
- **The analysis of the correlation in the case of the translation operator can be extended to other (more physical) operators mixing both regions**
- **Work in progress...**

In memory of Pedro F. González Díaz

