

A modified gravity from metric quantum fluctuations

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As we know gravity is highly non-linear theory. The non-linearities appear in front of derivative terms:

$$g^{\dots} g^{\dots}$$

One of the most neglected of approaches to the quantization of a non-linear theory is nonperturbative (NP) quantization by Heisenberg.

Heisenberg has applied NP
quantization for non-linear spinor field.

The essence of NP quantization is to write operator equation

$$\hat{R}_{\mu\nu} - \frac{1}{2}\hat{g}_{\mu\nu}\hat{R} = \kappa\hat{T}_{\mu\nu}, \quad (1)$$

Infinite equations set for Green functions

Heisenberg offer to write an infinite equations set for all Green functions as follows

$$\langle Q | \text{Eq. (1)} | Q \rangle = 0,$$

$$\langle Q | \hat{g}(x_1) \cdot \text{Eq. (1)} | Q \rangle = 0,$$

$$\langle Q | \hat{g}(x_1)\hat{g}(x_2) \cdot \text{Eq. (1)} | Q \rangle = 0,$$

$$\dots = 0,$$

$$\langle Q | \text{prod. of } g \text{ at different points } (x_1, \dots, x_n) \cdot \text{Eq. (1)} | Q \rangle = 0$$

How to solve ?

The first way is to cut off equations set by using some decomposition $G_{m+n} \approx G_m G_n$ and taking into account the first $p < m + n$ equations. The second one is to write some functional (for example, action) and average it using some assumptions about expectation value of Green functions.

$$\hat{g}_{\mu\nu} = g_{\mu\nu} + \widehat{\delta g}_{\mu\nu}$$



$$\left\langle \mathcal{L}(g + \widehat{\delta g}) \right\rangle \approx \mathcal{L}(g) + \mathcal{L}' \widehat{\delta g}^{\mu\nu} + \mathcal{L}'' \widehat{\delta g}^{\mu\nu} \widehat{\delta g}^{\rho\sigma}$$



$$\left\langle \widehat{\delta g}_{\mu\nu}(x_1) \cdot \widehat{\delta g}_{\rho\sigma}(x_2) \right\rangle \approx P_{\mu\nu}(x_1) P_{\rho\sigma}(x_2)$$

$$\hat{g}_{\mu\nu} = g_{\mu\nu} + \widehat{\delta g}_{\mu\nu}$$



$$\left\langle \mathcal{L}(g + \widehat{\delta g}) \right\rangle \approx$$

$$\mathcal{L}(g) + \mathcal{L}' \widehat{\delta g}^{\mu\nu} + \mathcal{L}'' \widehat{\delta g}^{\mu\nu} \widehat{\delta g}^{\rho\sigma}$$



$$\left\langle \widehat{\delta g}_{\mu\nu}(x_1) \cdot \widehat{\delta g}_{\rho\sigma}(x_2) \right\rangle \approx$$

$$P_{\mu\nu}(x_1) P_{\rho\sigma}(x_2)$$

$$\begin{aligned}\hat{g}_{\mu\nu} &= g_{\mu\nu} + \widehat{\delta g}_{\mu\nu} \\ &\downarrow \\ \langle \mathcal{L}(g + \widehat{\delta g}) \rangle &\approx \\ \mathcal{L}(g) + \mathcal{L}' \widehat{\delta g}^{\mu\nu} + \mathcal{L}'' \widehat{\delta g}^{\mu\nu} \widehat{\delta g}^{\rho\sigma} & \\ &\downarrow \\ \langle \widehat{\delta g}_{\mu\nu}(x_1) \cdot \widehat{\delta g}_{\rho\sigma}(x_2) \rangle &\approx \\ P_{\mu\nu}(x_1) P_{\rho\sigma}(x_2) &\end{aligned}$$

$$\left\langle \widehat{\delta g_{\mu\nu}}(x_1) \cdot \widehat{\delta g_{\rho\sigma}}(x_2) \right\rangle \approx P_{\mu\nu}(x_1) P_{\rho\sigma}(x_2)$$

It is necessary to emphasize that **even** for x_1, x_2 spacelike points this expression **is not zero !**

For perturbative quantization the situation is opposite.

Two possibilities for 2-point Green function

- $P_{\mu\nu}$ is proportional to the metric tensor

$$P_{\mu\nu} \propto g_{\mu\nu};$$

- $P_{\mu\nu}$ is proportional to the Ricci tensor

$$P_{\mu\nu} \propto \frac{R_{\mu\nu}}{R};$$

The proportionality coefficient should be some invariant. Consequently it has to be $\tilde{F}(R, R_{\mu\nu}R^{\mu\nu}, \dots)$.

For the ansatz $P_{\mu\nu} = g_{\mu\nu}$ we have

$$\begin{aligned} \langle \mathcal{L}(g + \delta g) \rangle &\approx \\ \sqrt{-g} \left(-\frac{c}{2\kappa} R - 2R\tilde{F}(R, \dots) \right) &= \\ \sqrt{-g} F(R, \dots) \end{aligned}$$

Thus we have obtained a *modified gravity theory*.

For the choice $\tilde{F} = \frac{\Lambda}{2R}$ we obtain

$$\langle \mathcal{L}(g + \delta g) \rangle \approx \sqrt{-g} \left(-\frac{c}{2\kappa} R - \Lambda \right)$$

Thus we have obtained Einstein gravity with Λ -term.

For the choice $\tilde{F} = \text{const}$ we obtain

$$\langle \mathcal{L}(g + \delta g) \rangle \approx \sqrt{-g}(-R) \left(\frac{c}{2\kappa} + \text{const} \right)$$

Thus we have obtained Einstein gravity with *modified gravitational constant*.

The calculation of the expectation value of $\langle \mathcal{L} + \delta^2 \mathcal{L} \rangle$ gives us

$$\langle \mathcal{L} + \delta^2 \mathcal{L} \rangle = \sqrt{-g} \left\{ \left[\frac{1}{2} - F(R, \dots) \right] \nabla^\mu \phi \nabla_\mu \phi - \left[1 + 2F(R, \dots) \right] V(\phi) \right\}. \quad (2)$$

We see that: (a) for big enough metric quantum fluctuations the function F becomes $F > 1/2$ and the scalar field can be a phantom one; (b) non-minimal coupling between scalar field and gravity appears.

Quantum fluctuations of metric give rise to $F(R)$ modified gravitational theories (with the applications to the explanation of modern Universe acceleration and so on).