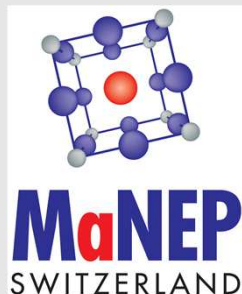
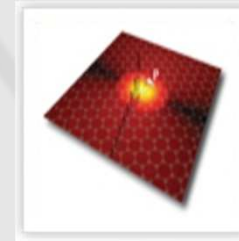


# Magneto-plasmonic effects in epitaxial graphene

Alexey Kuzmenko

University of Geneva

Graphene Nanophotonics  
Benasque, 4 March 2013



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[www.graphene-flagship.eu](http://www.graphene-flagship.eu)

GRAPHENE FLAGSHIP



# Collaborators



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**University of Geneva**



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**Univ. of Erlangen  $\Rightarrow$  Univ. Chemnitz**



J. Chen, F. Huth, R. Hillenbrand  
**CIC nanoGUNE, San Sebastian**

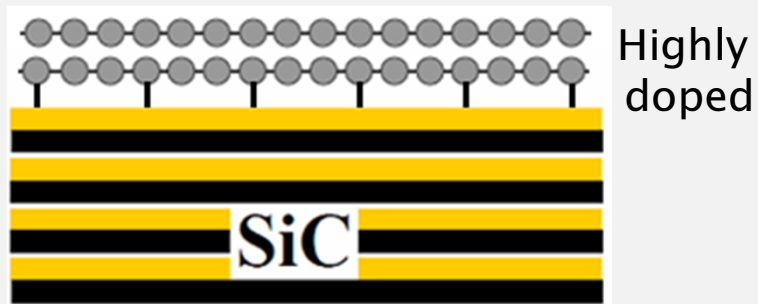


M. Orlita, M. Potemski  
**CNRS, Grenoble**

# Different types of epitaxial graphene (EG)

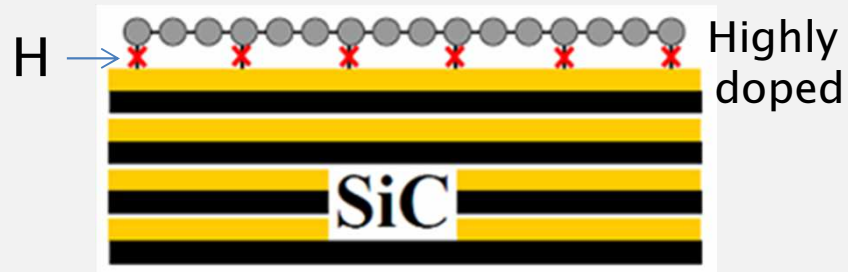
## Si-face

monolayer on buffer layer



K. Emtsev, *Nature Materials*, 8, 203 (2009)

Quasi-freestanding monolayer  
(hydrogenated dangling bonds)

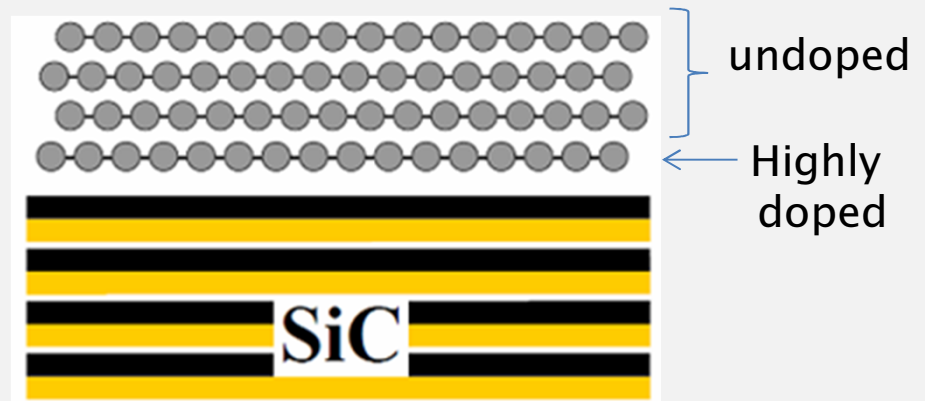


C. Riedl *et al.* PRL 103, 246804 (2009)

Growth at  $\sim 1500$  °C in Ar

## C-face

twisted multilayer



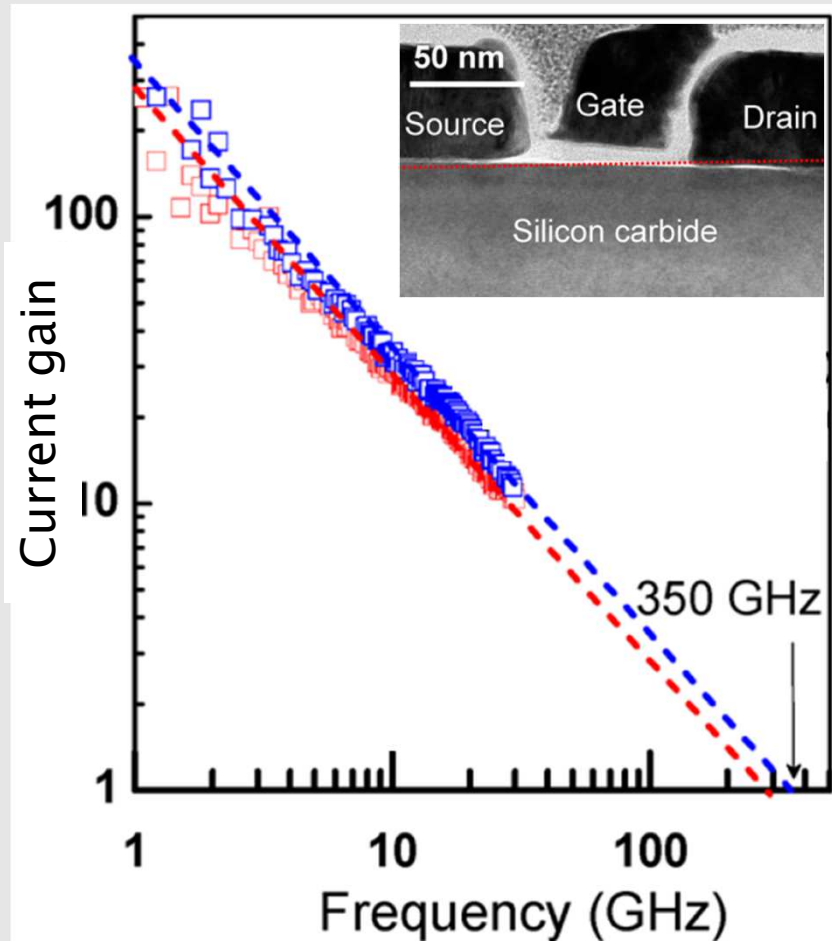
C. Berger *et al.* J. Phys. Chem. B 108, 19912 (2004)

Defect free-regions  
of several micron size

# Single layer graphene on Si-face of SiC

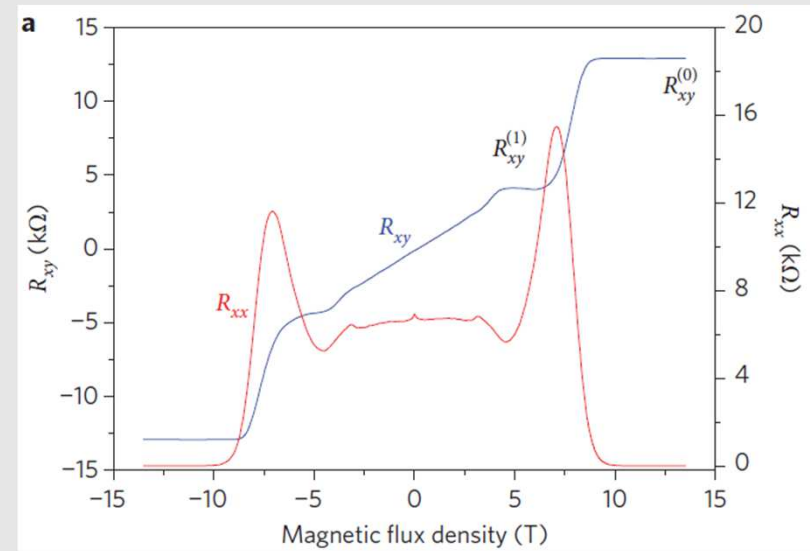
Ultrafast graphene transistors

$$\text{Mobility } \mu \approx 3000 \frac{\text{cm}^2}{\text{V s}}$$

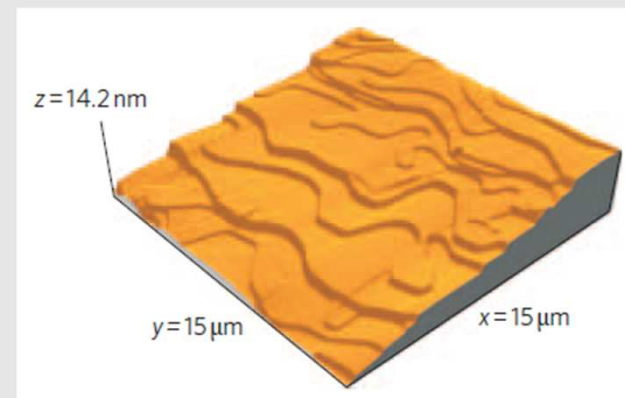


IBM group: Wu *et al.* Nano Letters (2012)

QHE resistance standard

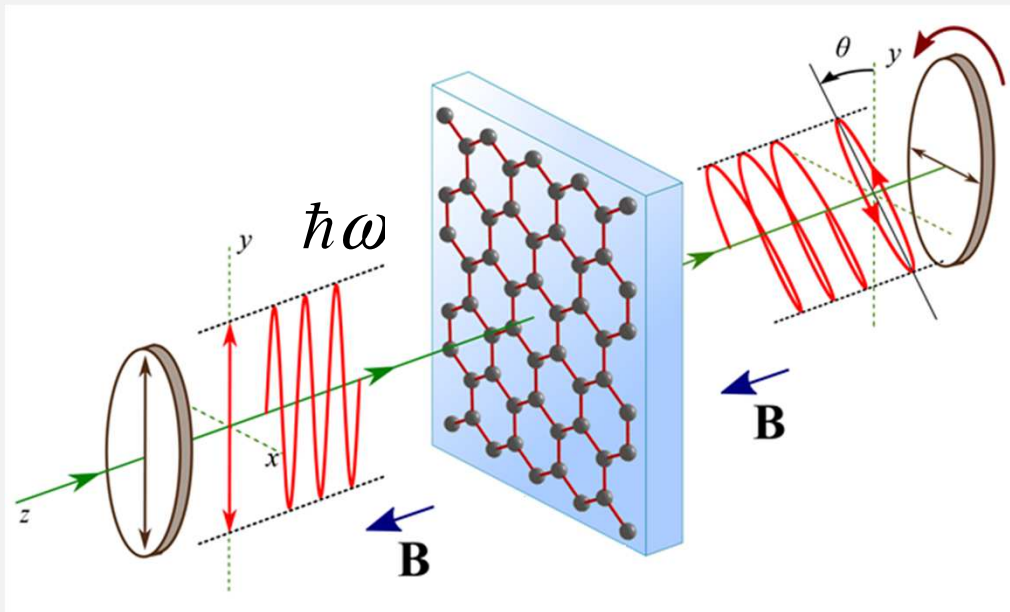
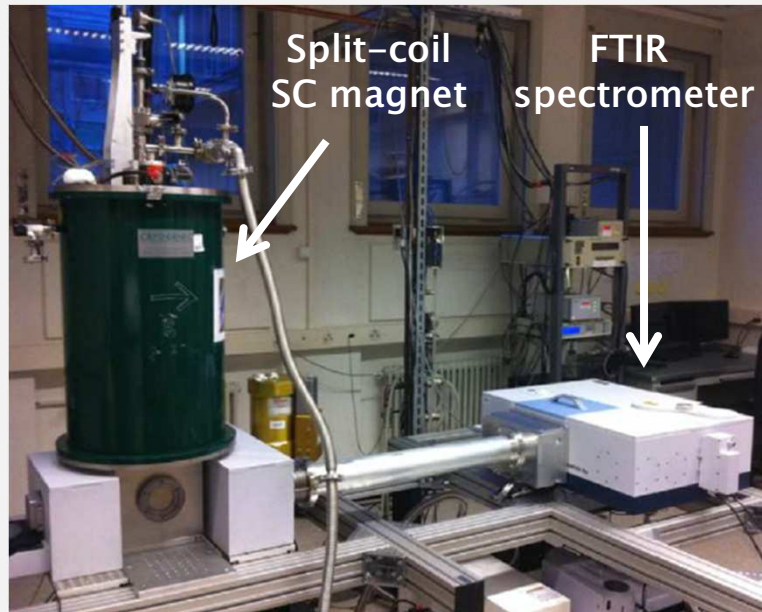


Terrace steps





A. Tzalenchuk *et al.* Nature Nano. (2010)

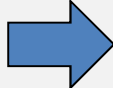
# Infrared/THz magneto-optical setup

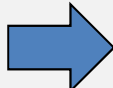


- ✚  $\hbar\omega$  from 1 meV to 0.5 eV
- ✚ B from -7 to +7 T
- ✚ Linear polarization
- ✚ Far-field / large spot

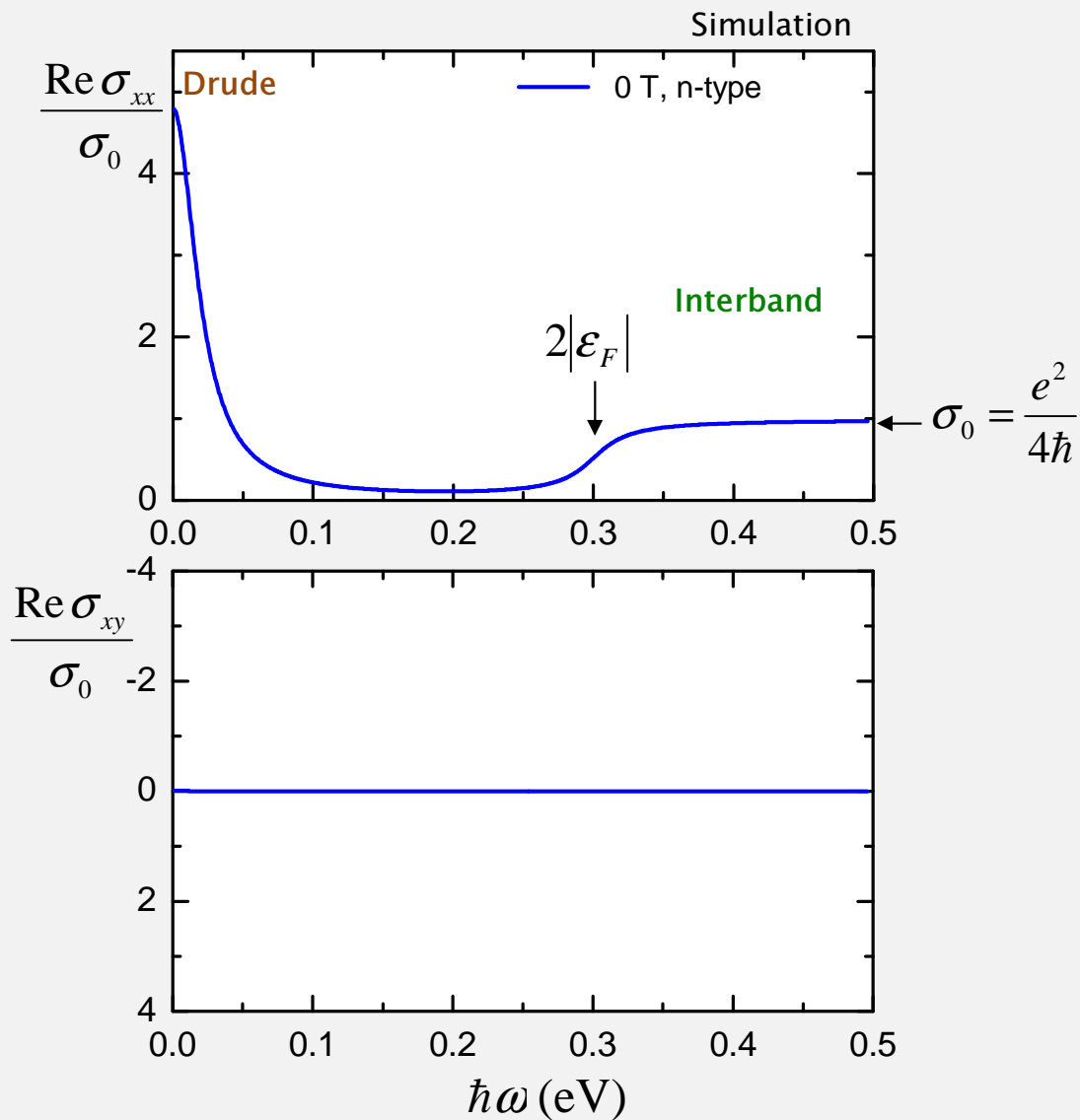
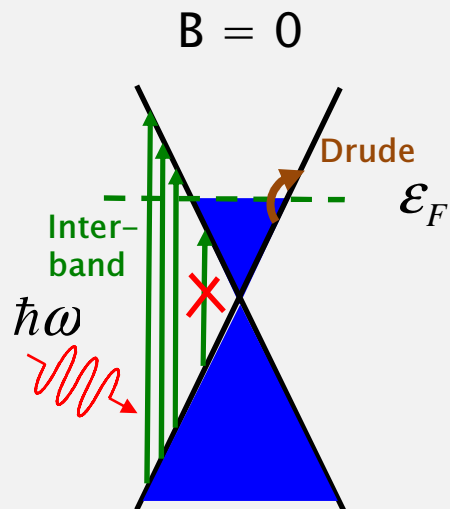
Finally we extract:  

$$\sigma_{\pm}(\omega) = \sigma_{xx}(\omega) \pm i\sigma_{xy}(\omega)$$

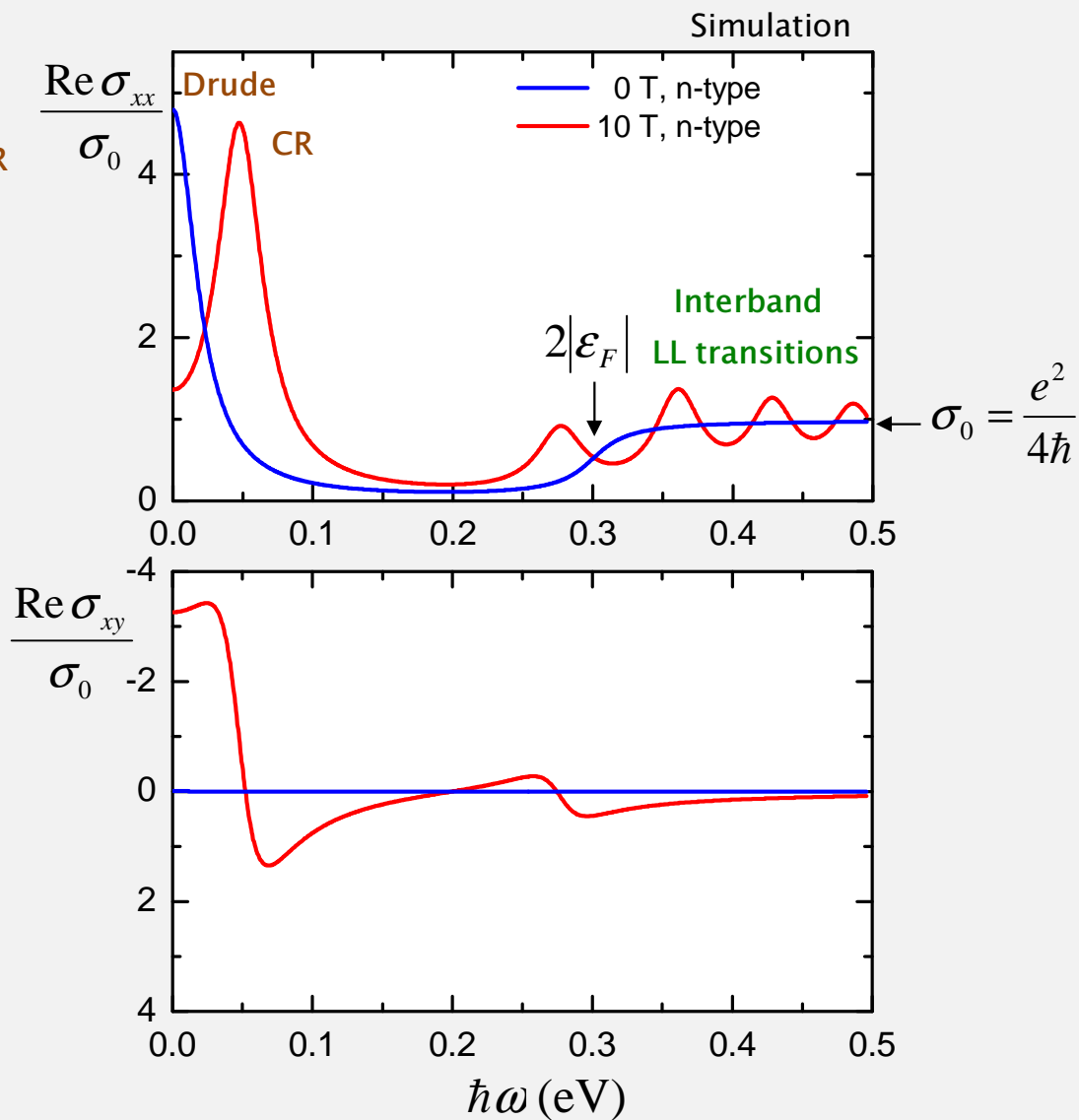
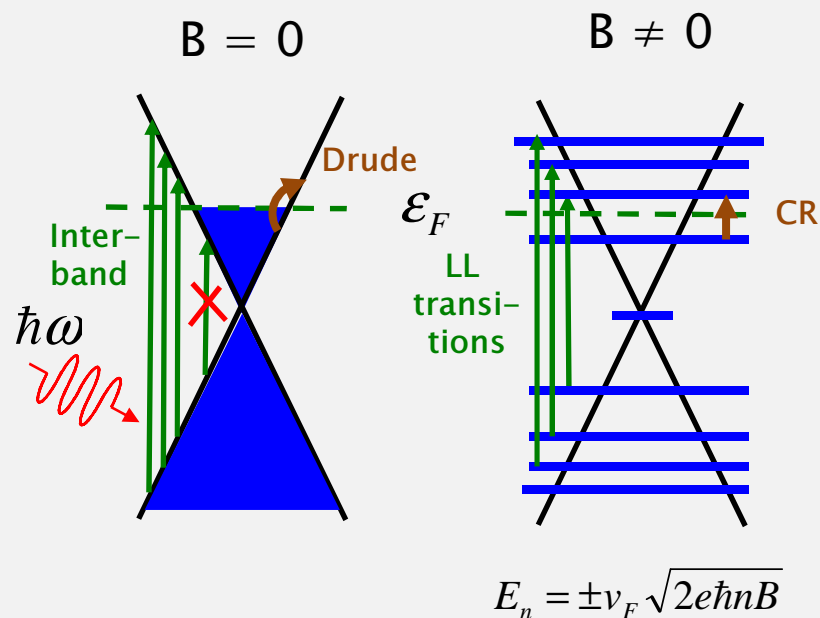
Absorption  Optical conductivity  $\sigma_{xx}(\omega)$

Faraday rotation  Optical **Hall** conductivity  $\sigma_{xy}(\omega)$

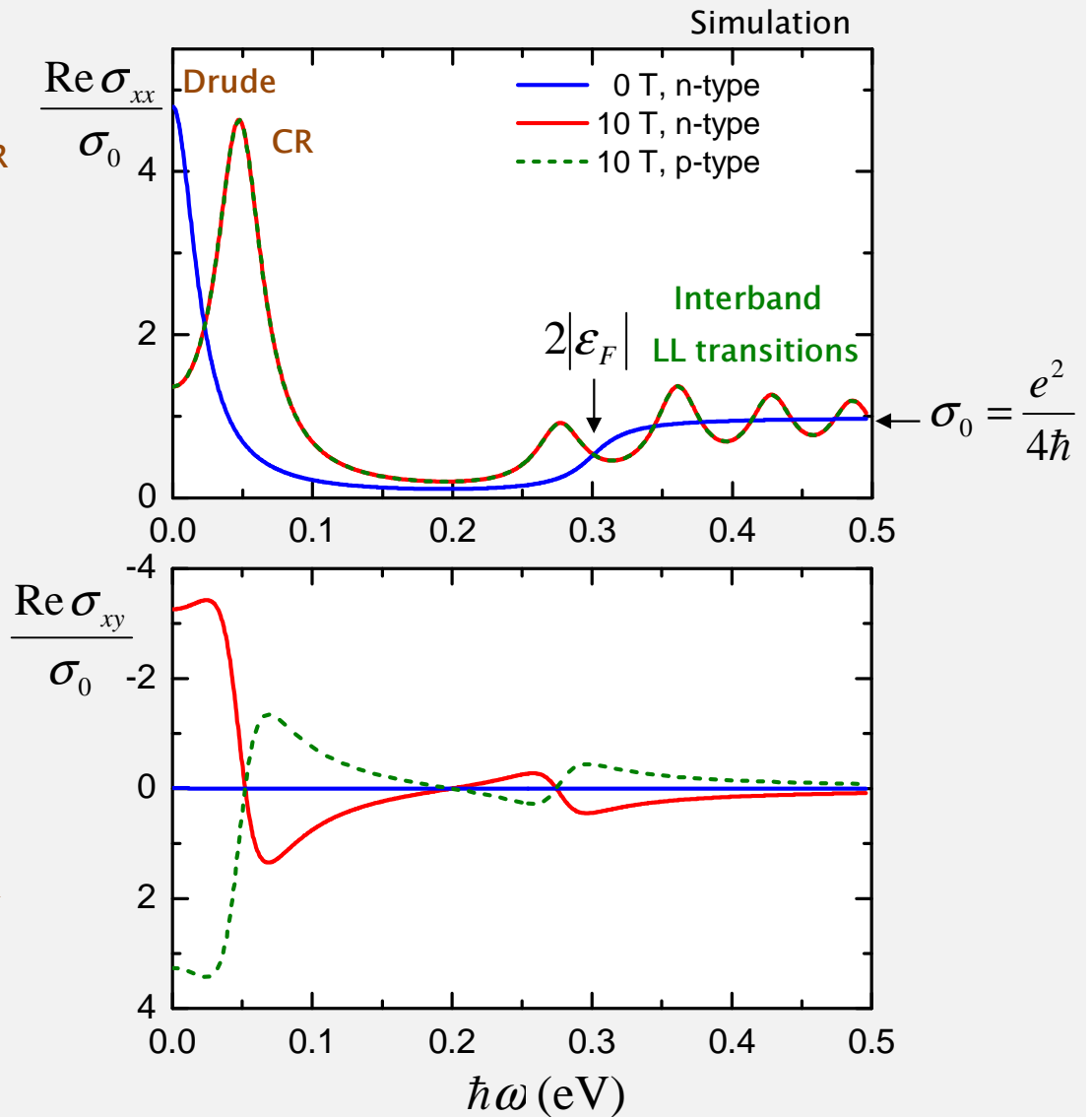
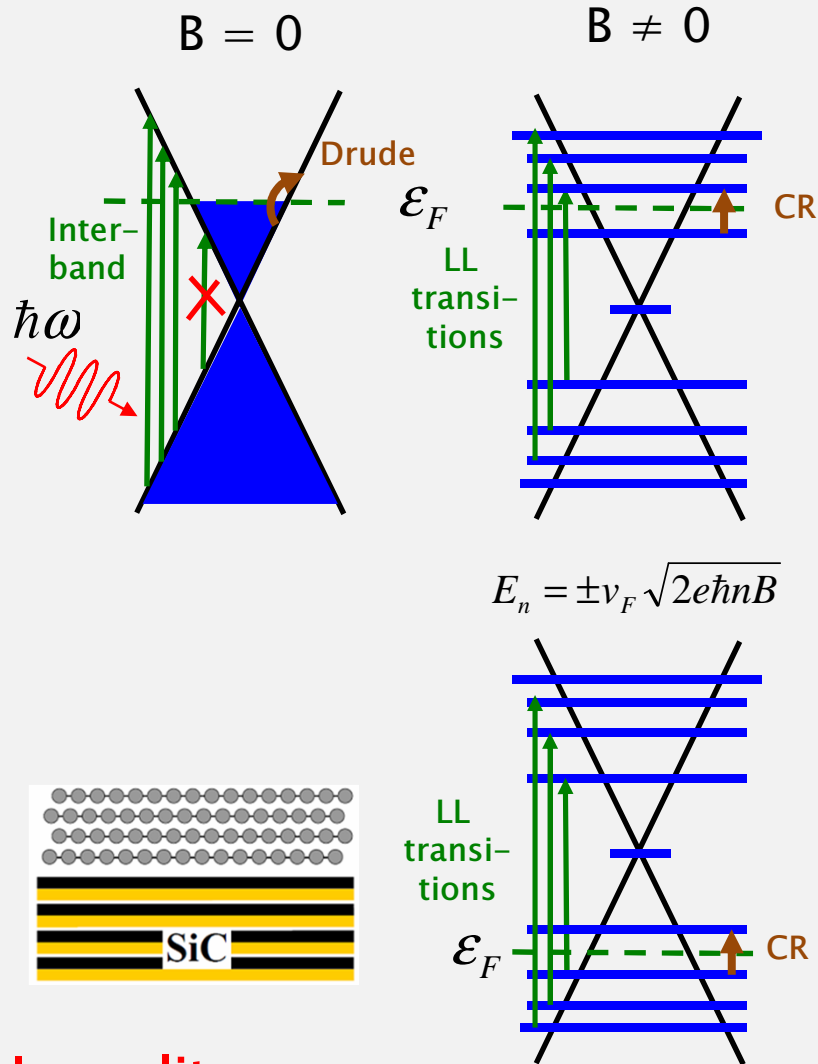
# Optical conductivity of one monolayer (theory)



# Optical conductivity of one monolayer (theory)



# Optical conductivity of one monolayer (theory)

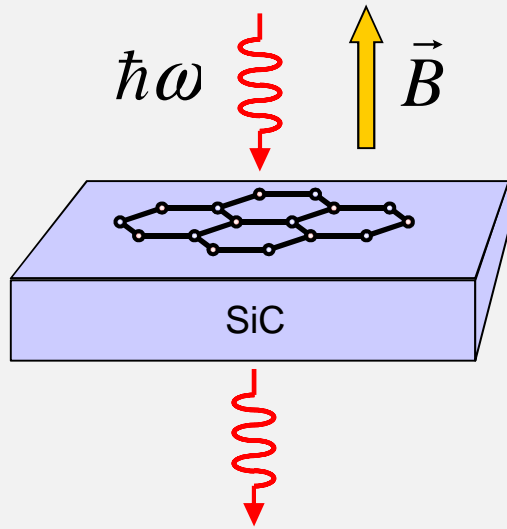


## In reality:

- ⚡ all layers add to the signal
- ⚡ doping changes across layers
- ⚡ stacking and grains may invalidate simple theory

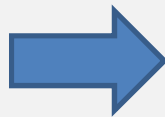


# What magneto-optics can tell about graphene ?



- ✚ Thickness / homogeneity
- ✚ Doping level
- ✚ Doping type (p or n)
- ✚ Doping homogeneity
- ✚ Mobility
- ✚ Fermi velocity
- ✚ Cyclotron mass
- ✚ Electron-hole asymmetry
- ✚ Stacking
- ✚ Grains boundaries
- ✚ ...

- ✚ Sees all layers
- ✚ No contacts/resist
- ✚ No UHV needed
- ✚ Done routinely

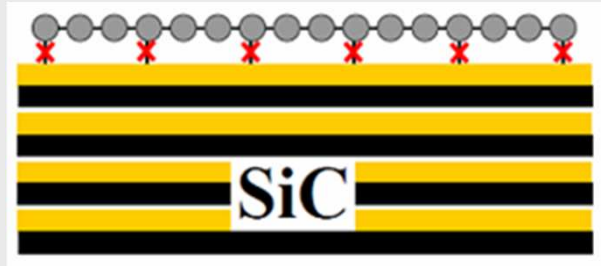


useful for routine  
characterization



# Extracting physical parameters

Quasi-free standing monolayer



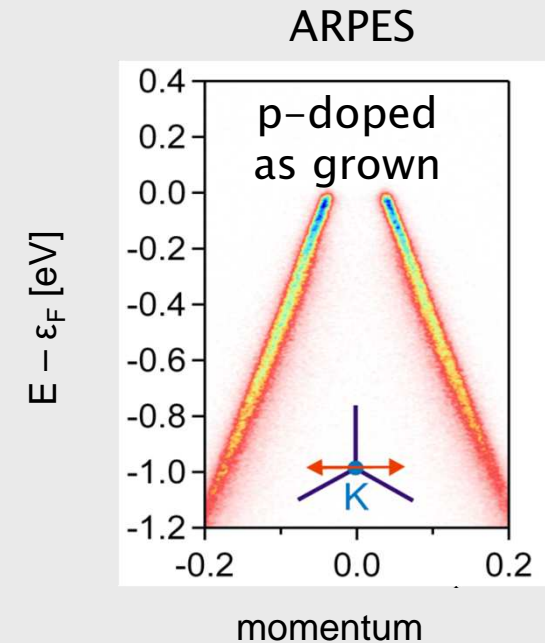
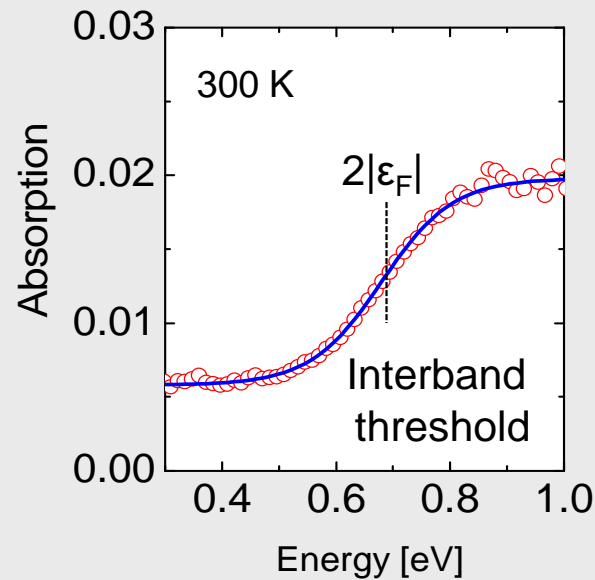
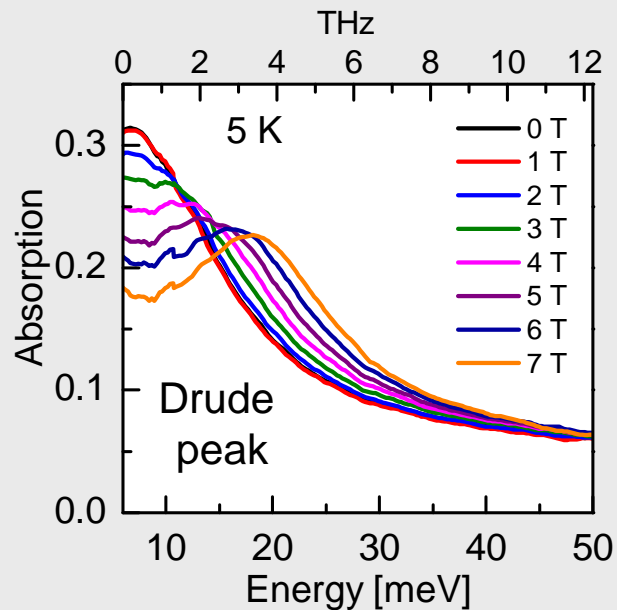
$$|\varepsilon_F| \approx 0.35 \text{ eV}$$

$$m_{CR} \approx 0.055 m_e$$

$$n \approx 8 \cdot 10^{12} \text{ cm}^{-2}$$

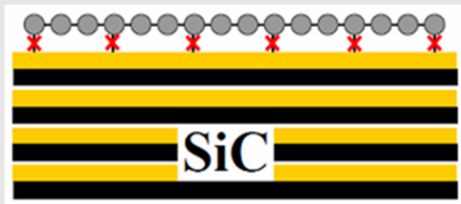
$$v_F = 1.05 \cdot 10^6 \text{ m/s}$$

$$\mu \approx 2500 \text{ cm}^2/(\text{V} \cdot \text{s})$$

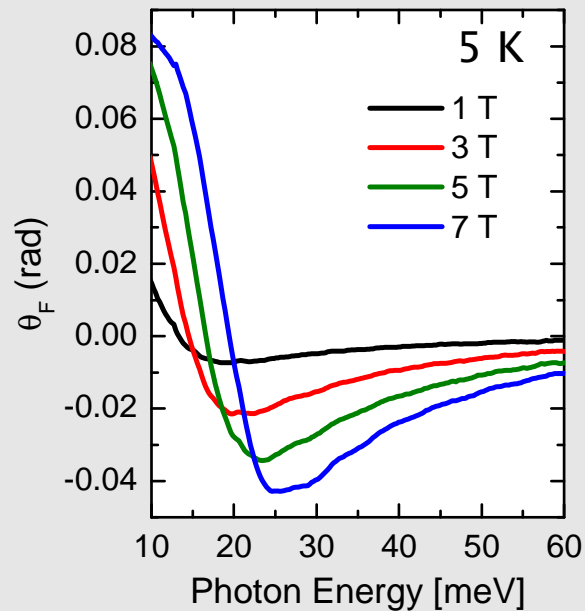


# Electrons or holes ?

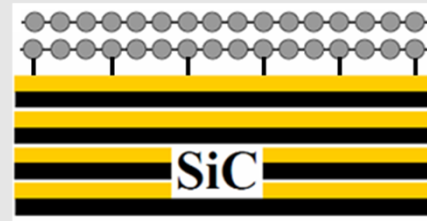
Quasi-free standing monolayer



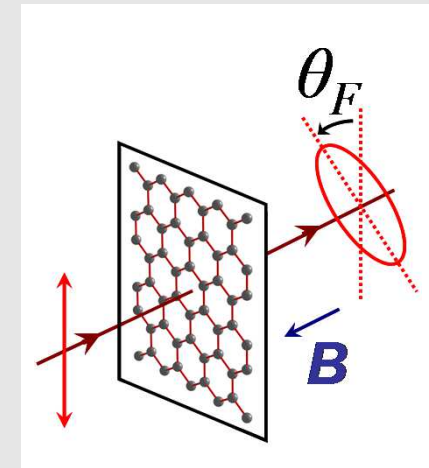
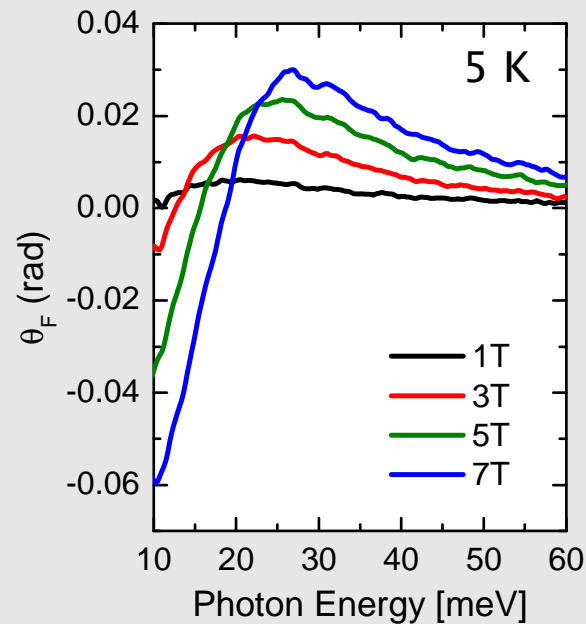
p-doped



Monolayer on buffer layer



n-doped

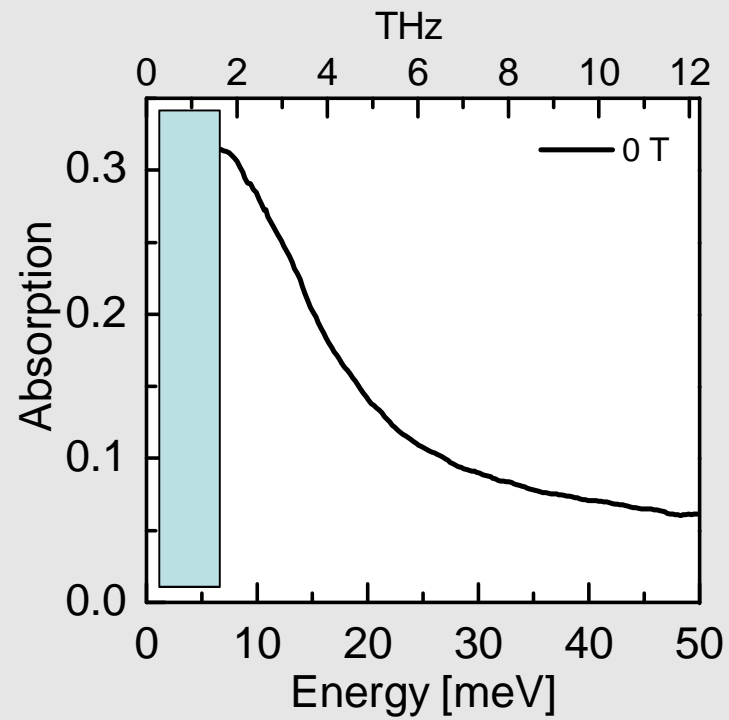


Faraday rotation is sensitive to the doping type

I. Crassee *et al.* Nature Physics 7, 48 (2011)

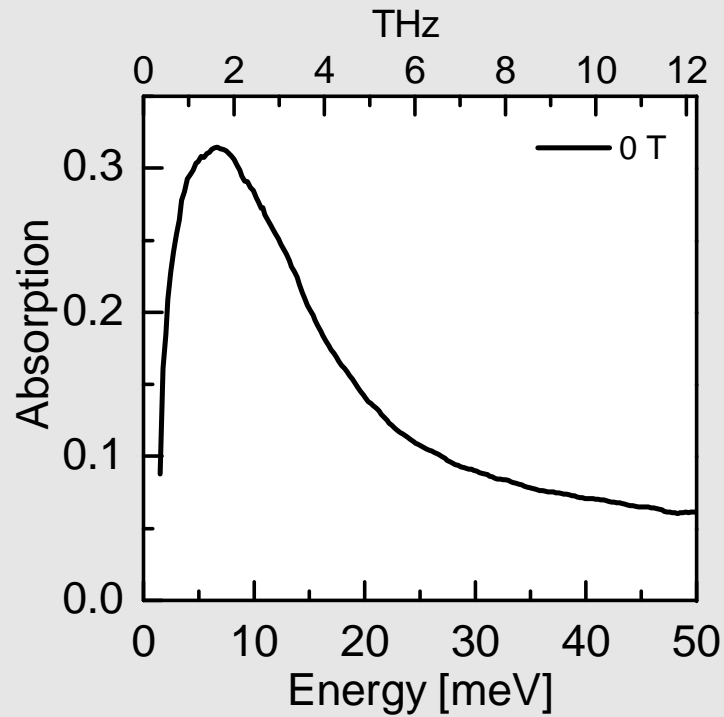
N. Ubrig *et al.*, in preparation

# Drude peak



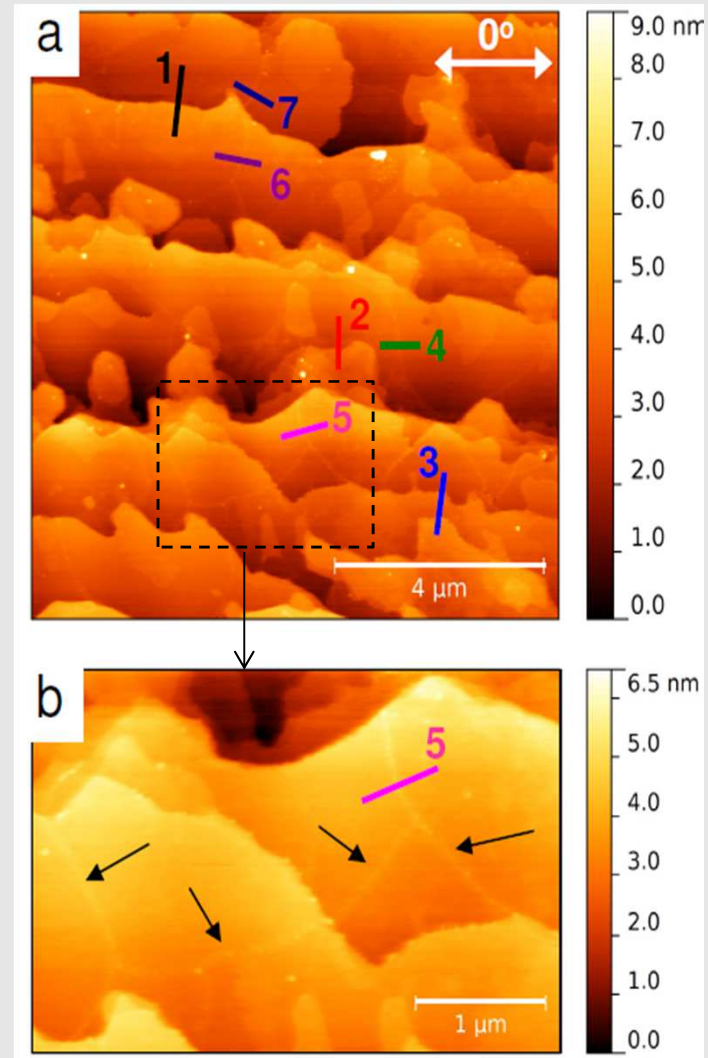
QFS monolayer

# Drude peak ... is not a Drude peak !



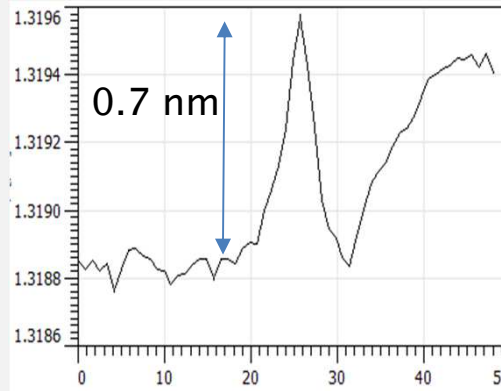
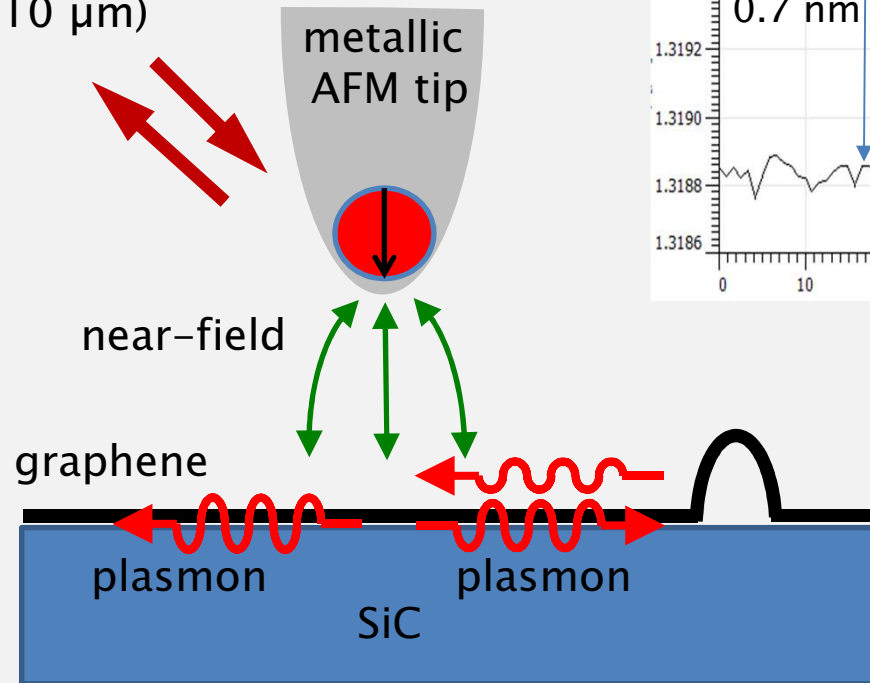
- ⚡ Plasmon excitation
- ⚡ Caused by steps (and wrinkles)
- ⚡ Edge resistance should be big

Quasi-freestanding monolayer

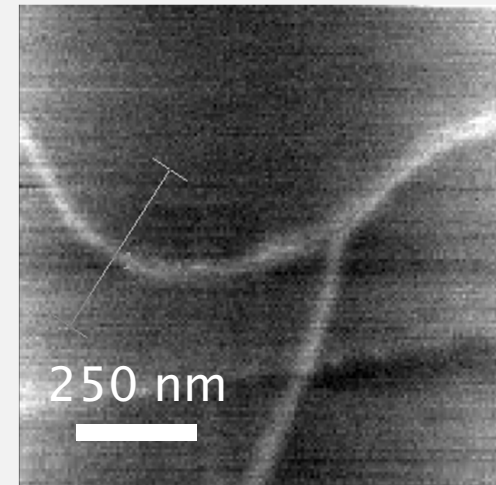


# AFM–near field optical plasmon imaging

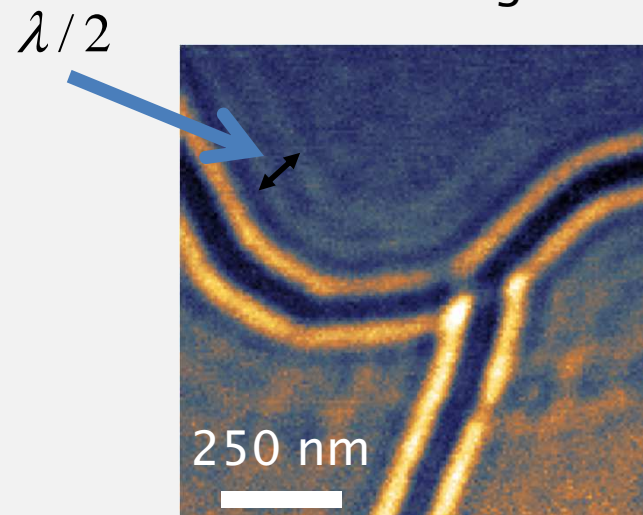
IR laser  
(10  $\mu\text{m}$ )



AFM topography



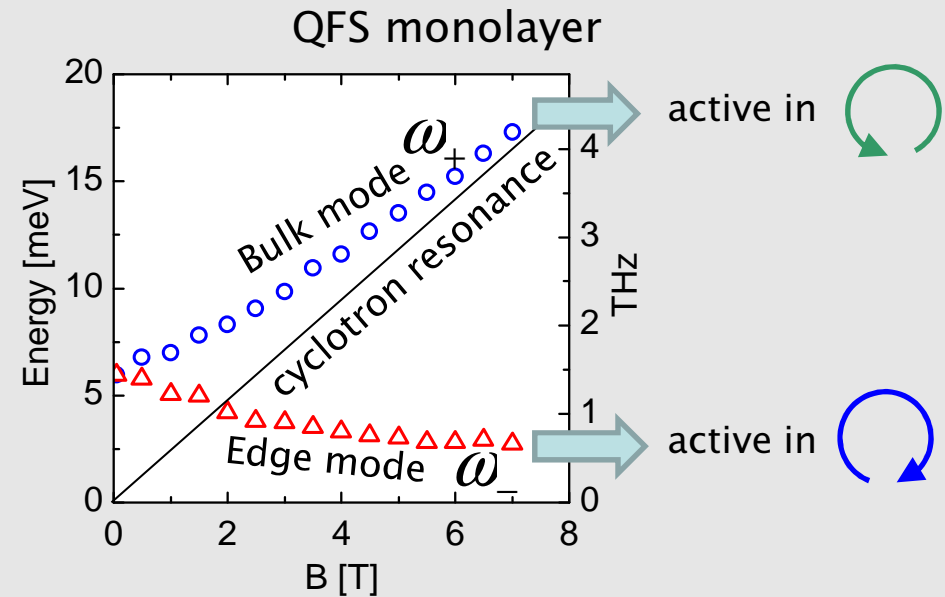
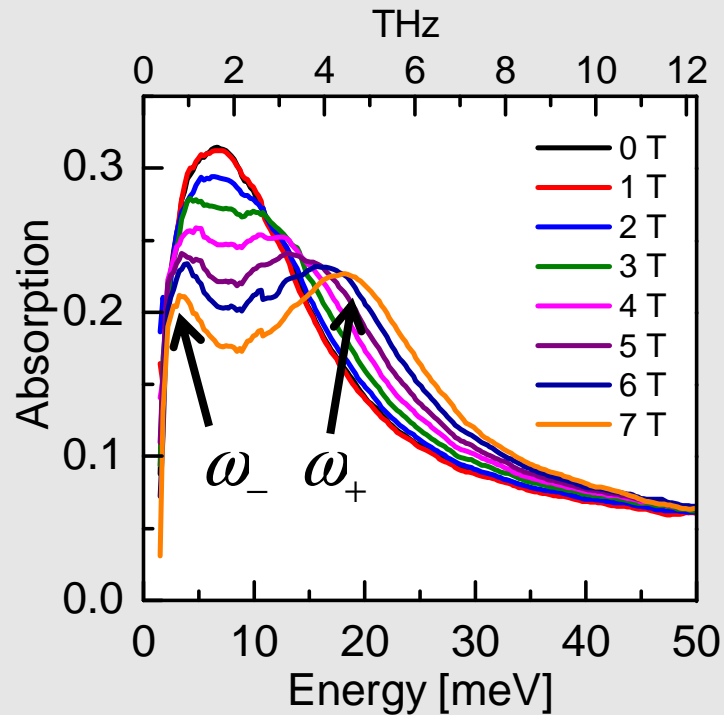
Near field signal



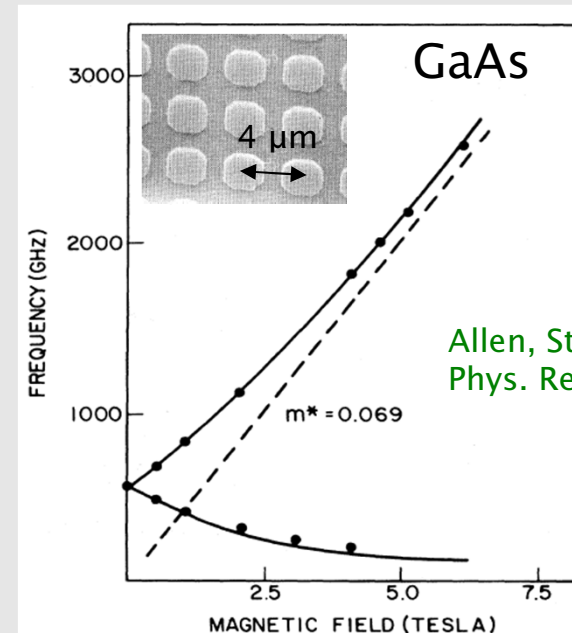
SiC terraces and graphene wrinkles are strong plasmon scatterers

Details: talk of R. Hillenbrand on Tuesday  
See also: J. Chen *et al*, Nature 487, 77 (2012); Fei *et al*, *ibid*, p. 82.

# Magnetoplasmons in graphene



- ⚡ Plasmon splits in two magnetoplasmons
- ⚡ Bulk and edge modes
- ⚡ similar to classical works on 2DEGs



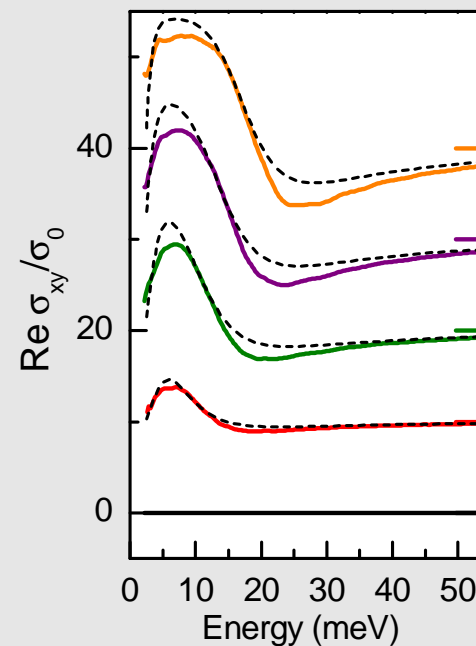
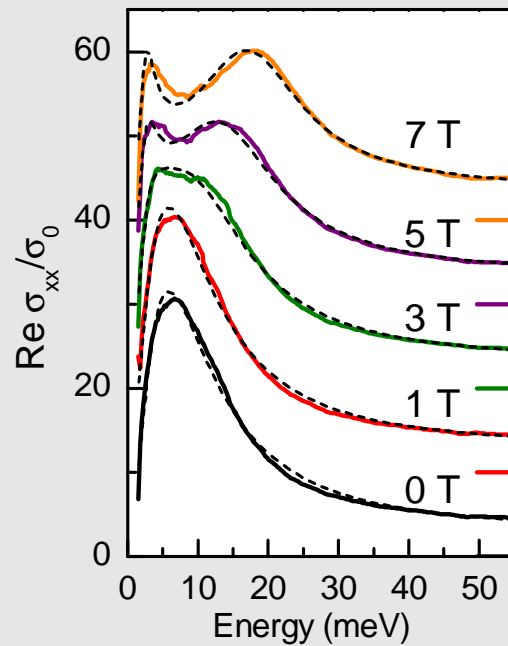
Allen, Stormer and Hwang, Phys. Rev. B 28, 4875 (1983)

# Effective medium approximation

$$\sigma_{\pm}(\omega) = \frac{ne^2}{m} \frac{i}{\omega \pm \omega_c + i\gamma - \omega_0^2/\omega}$$

S. Mikhailov, PRB 54, 10335 (1996)

$$\omega_{\pm} = \sqrt{\omega_0^2 + \frac{\omega_c^2}{4}} \pm \frac{\omega_c}{2}$$

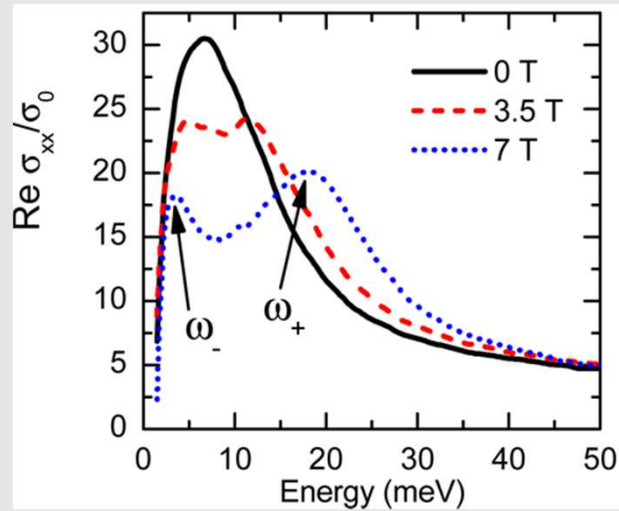
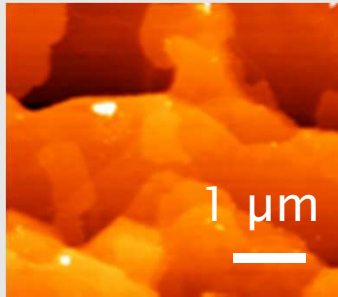


EMA works surprisingly well !



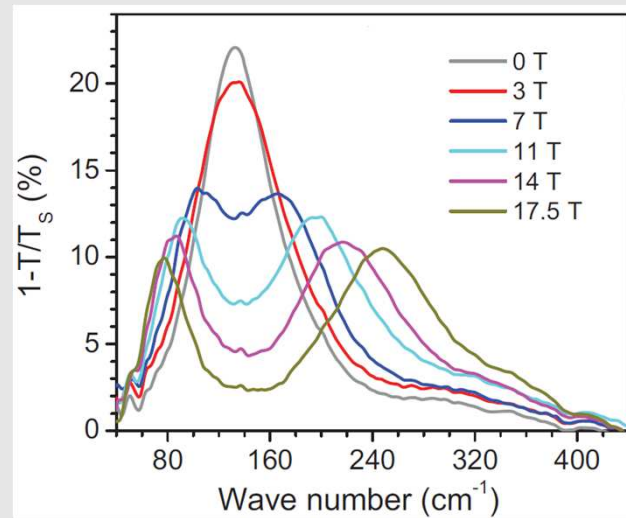
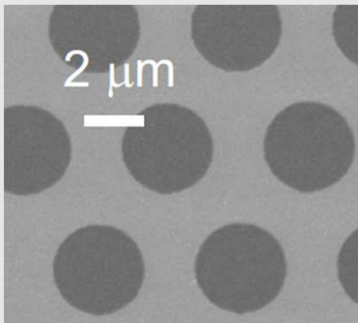
# Magnetoplasmons in graphene

epitaxial  
graphene



I. Crassee *et al.* Nano Lett. 12, 2470 (2012)

CVD graphene  
disks

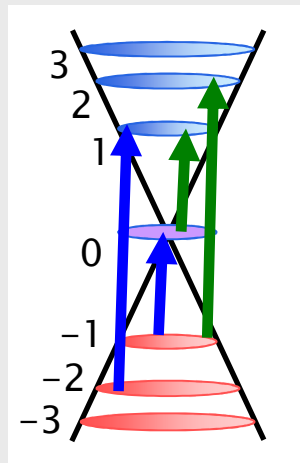
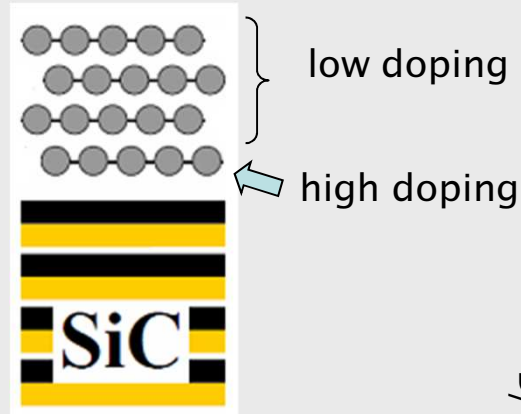


H. Yan *et al.* Nano Lett. 12, 3766 (2012)

Also magnetoplasmons  
observed in the QHE regime  
I. Petkovic *et al.* PRL 110, 016801 (2013)

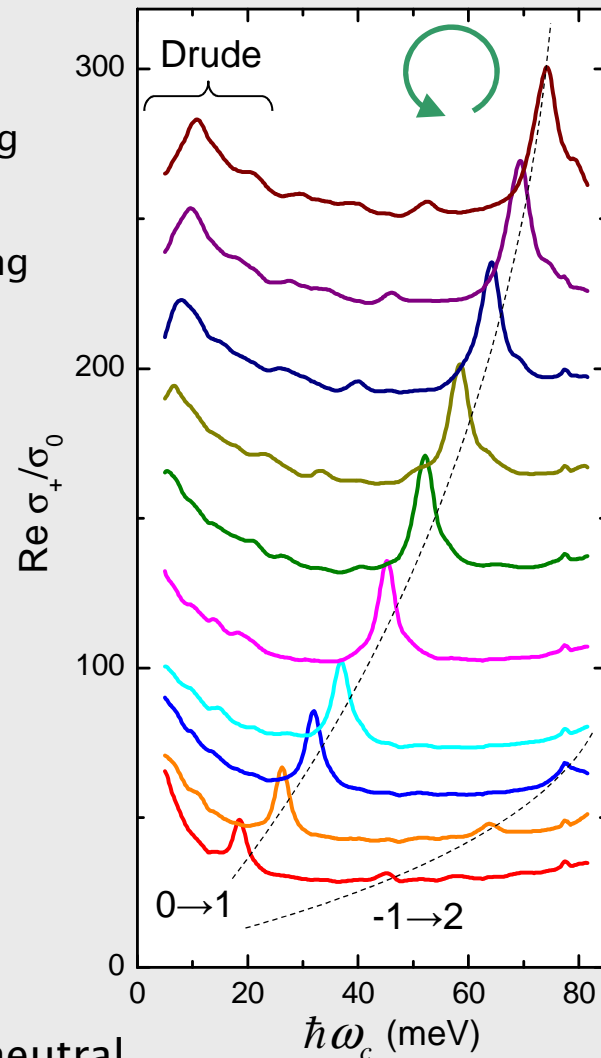
# Multilayer epitaxial graphene on C-side of SiC

twisted multilayer  
(10–15 layers)

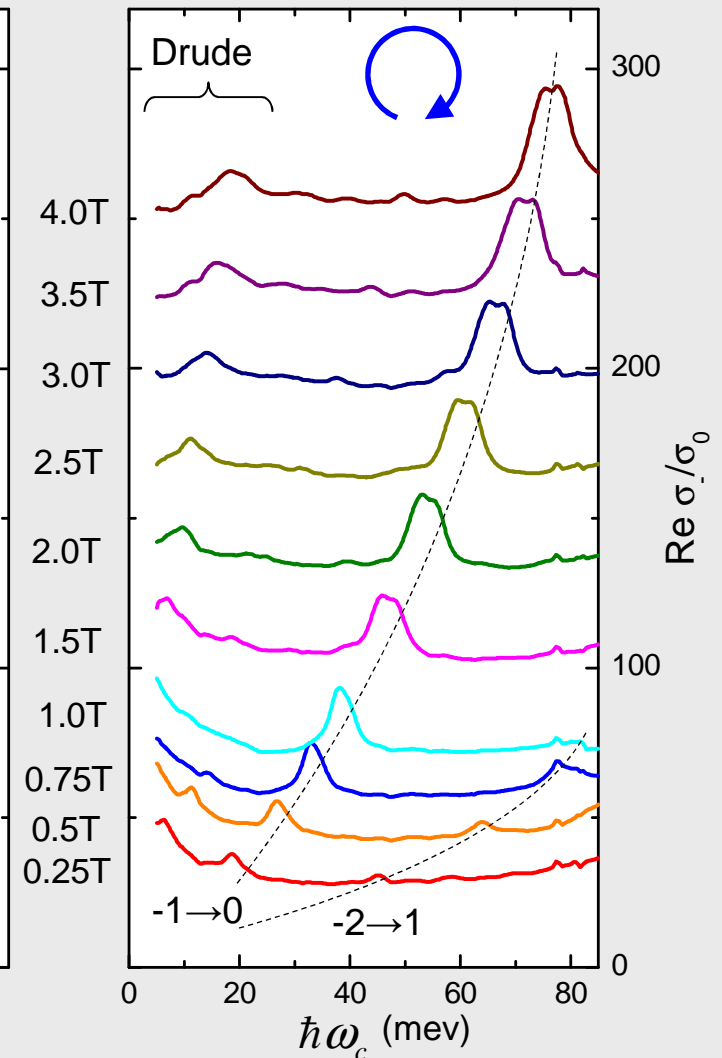


- ⊕ approximately charge-neutral
- ⊕  $\propto \sqrt{B}$
- ⊕ electron-hole asymmetry
- ⊕ stacking effects

electrons



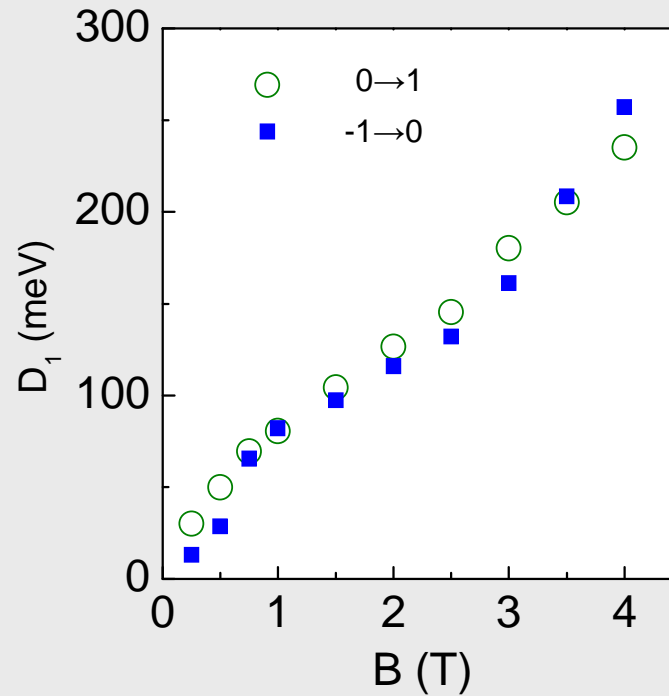
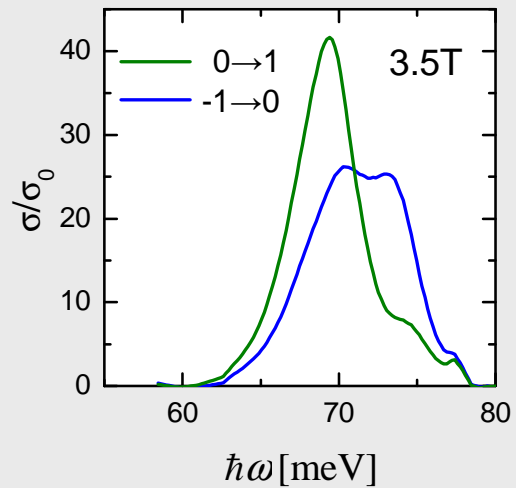
holes



$$\varepsilon_1 = v_F \sqrt{2e\hbar B}$$

# Optical LL transition intensities

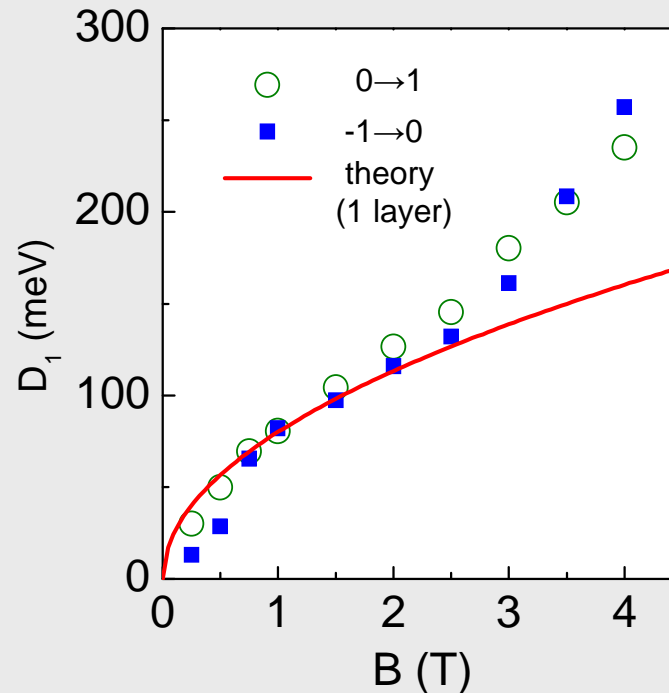
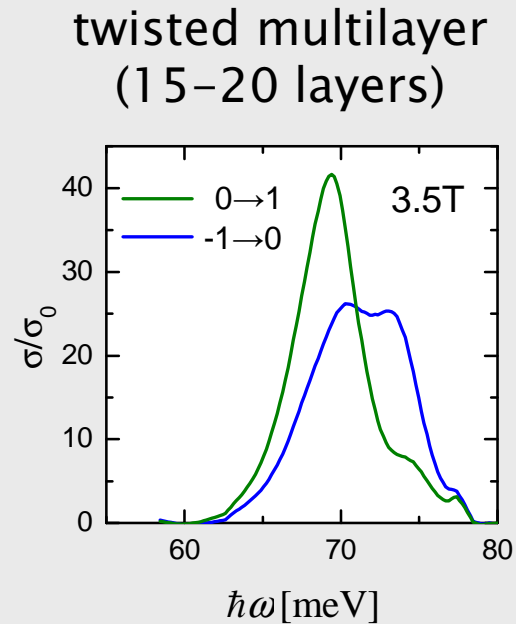
twisted multilayer  
(15–20 layers)



$$D_1 = \int \frac{\sigma(\omega)}{\sigma_0} d(\hbar\omega)$$



# Optical LL transition intensities



$$D_1 = \int \frac{\sigma(\omega)}{\sigma_0} d(\hbar\omega)$$



theory for  
Dirac fermions

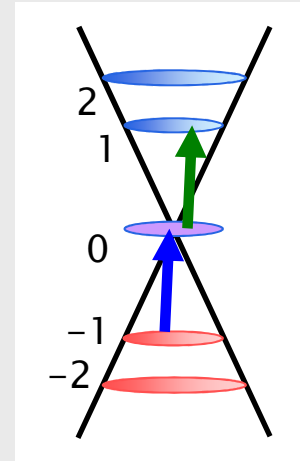
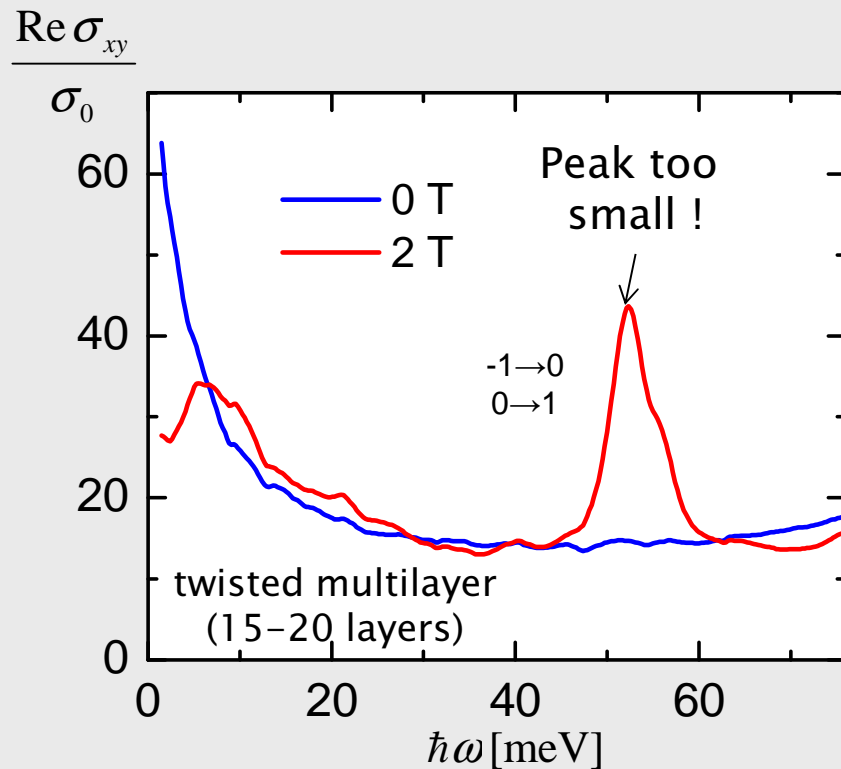
1 layer:  
 $D_1 = 2\varepsilon_1$

$N$  layers:  
 $D_1 = 2\varepsilon_1 N$

✚ **NO**  $B^{1/2}$  dependence

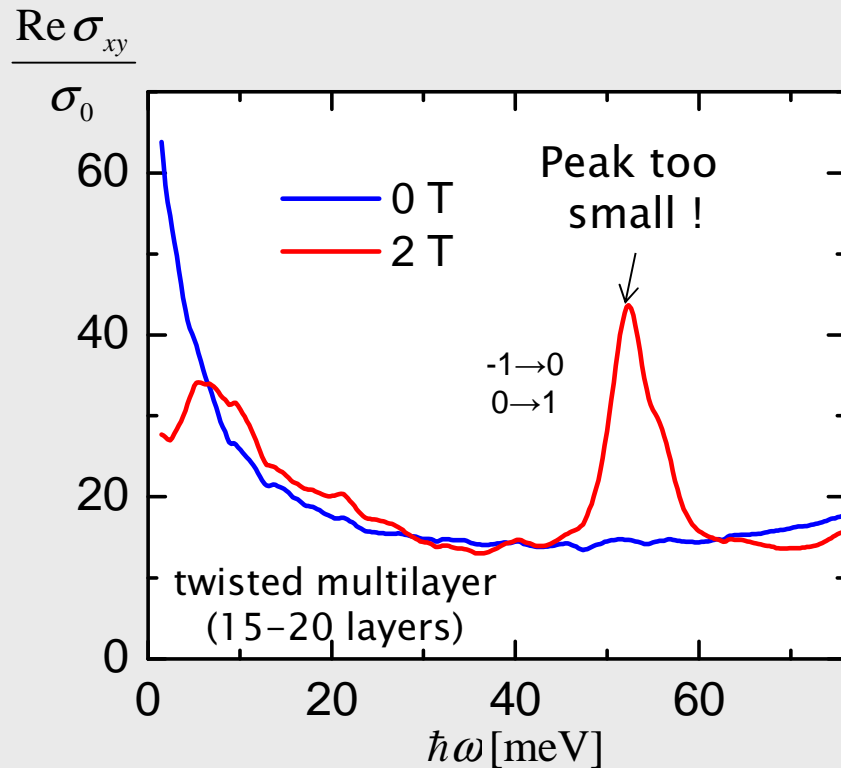
✚ Intensity **catastrophically** lower than expected !!!

# Optical LL transition intensities



- ✚ only 1–2 layers (out of 15–20) contribute to Landau peaks
- ✚ A stacking effect ?
- ✚ Most layers overdamped?

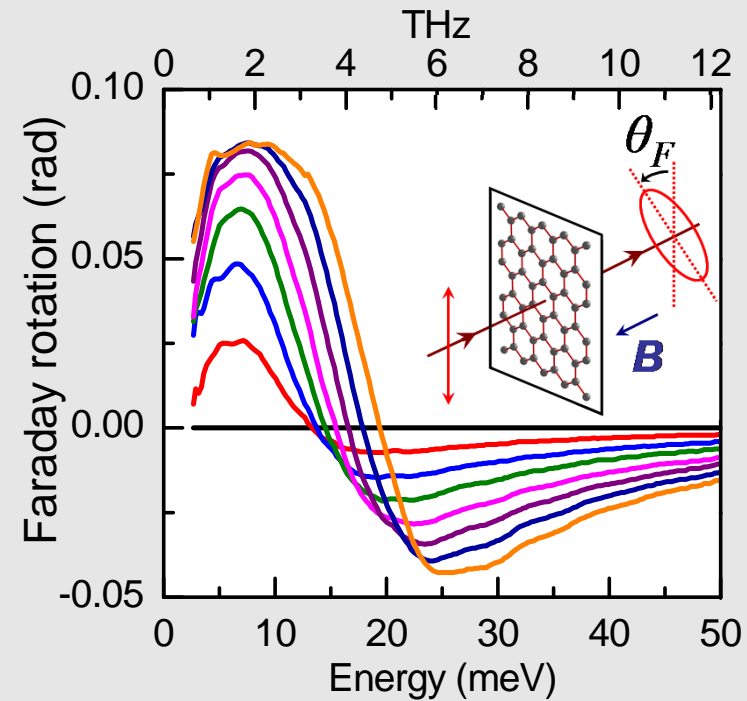
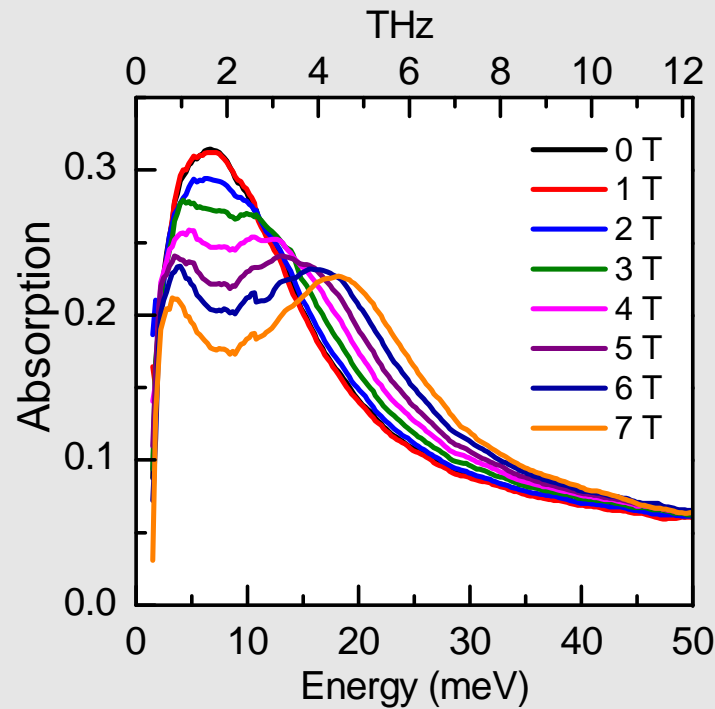
# Optical LL transition intensities



- ✚ only 1–2 layers (out of 15–20) contribute to Landau peaks
- ✚ A stacking effect ?
- ✚ Most layers overdamped?



# Graphene for magneto-optical applications

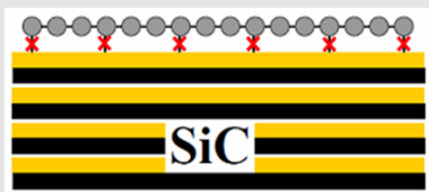


Absorption coefficient\*  $\approx 70 \frac{\%}{\text{nm}}$

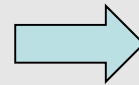
Verdet constant\*  $\approx 15 \frac{\text{degrees}}{\text{T} \cdot \text{nm}}$

\*provided that layers can be stacked without affecting their properties

QFS monolayer Potentially tunable

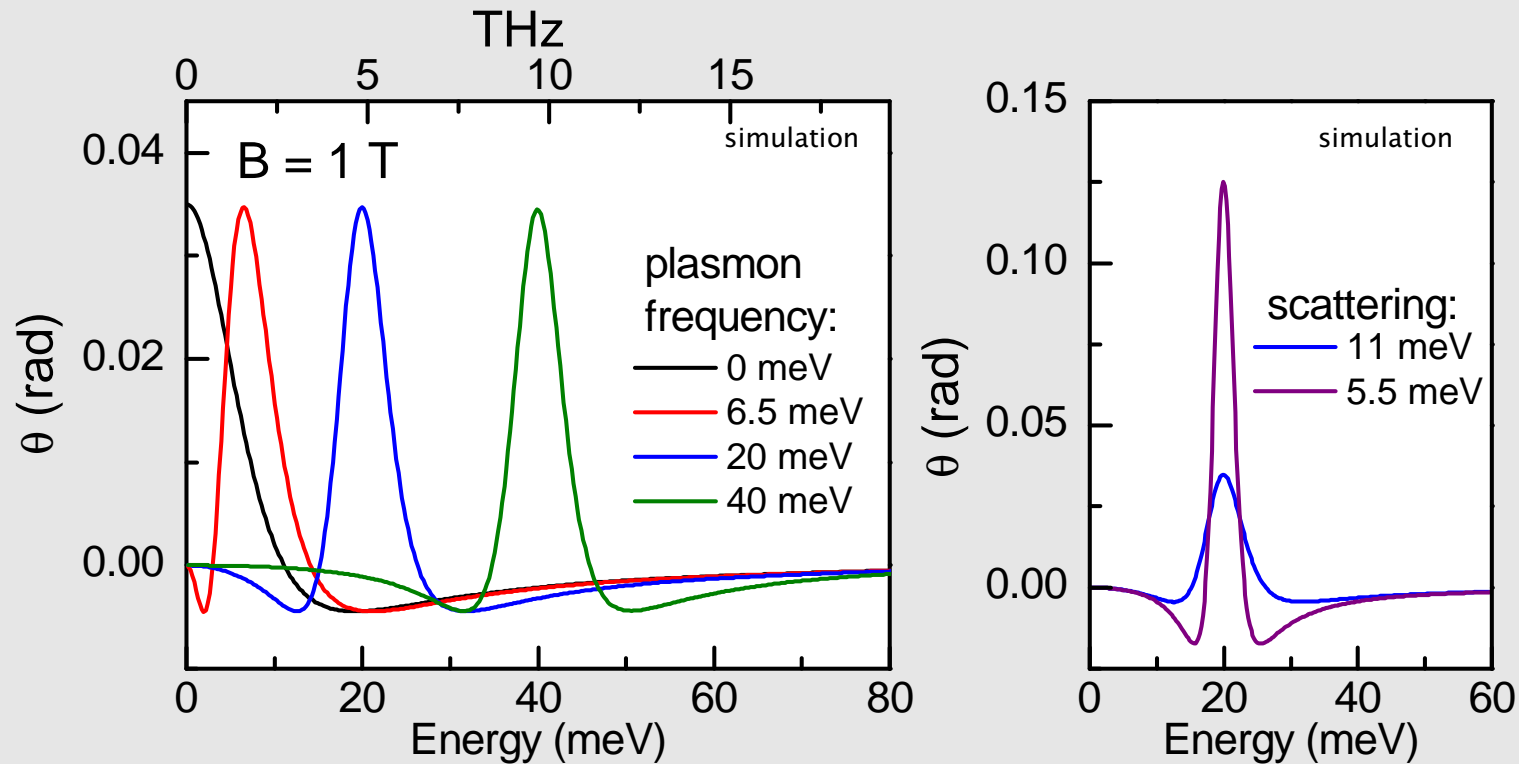


by gate !



Ultrafast control !

# Faraday rotation affected by magneto-plasmons



- rotation can be controlled by the plasmon frequency
- higher-mobility samples produce larger rotation
- defects are useful !



# Graphene-based magneto-optical applications

## Types of devices

- ✦ Absorption modulators
- ✦ Polarization modulators
- ✦ Faraday isolators (valves)...

## Spheres of application

- ✦ Telecommunications
- ✦ Biosensing
- ✦ Security
- ✦ Astronomy...

## Advantages of epitaxial graphene

- ✦ Samples very large
- ✦ Chemical potential homogeneous
- ✦ SiC is transparent

## Disadvantages of epitaxial graphene

- ✦ Gating is relatively difficult
- ✦ Controlled multilayer growth is tricky

H. Da & C.W.Qiu, APL 100, 241106 (2012)

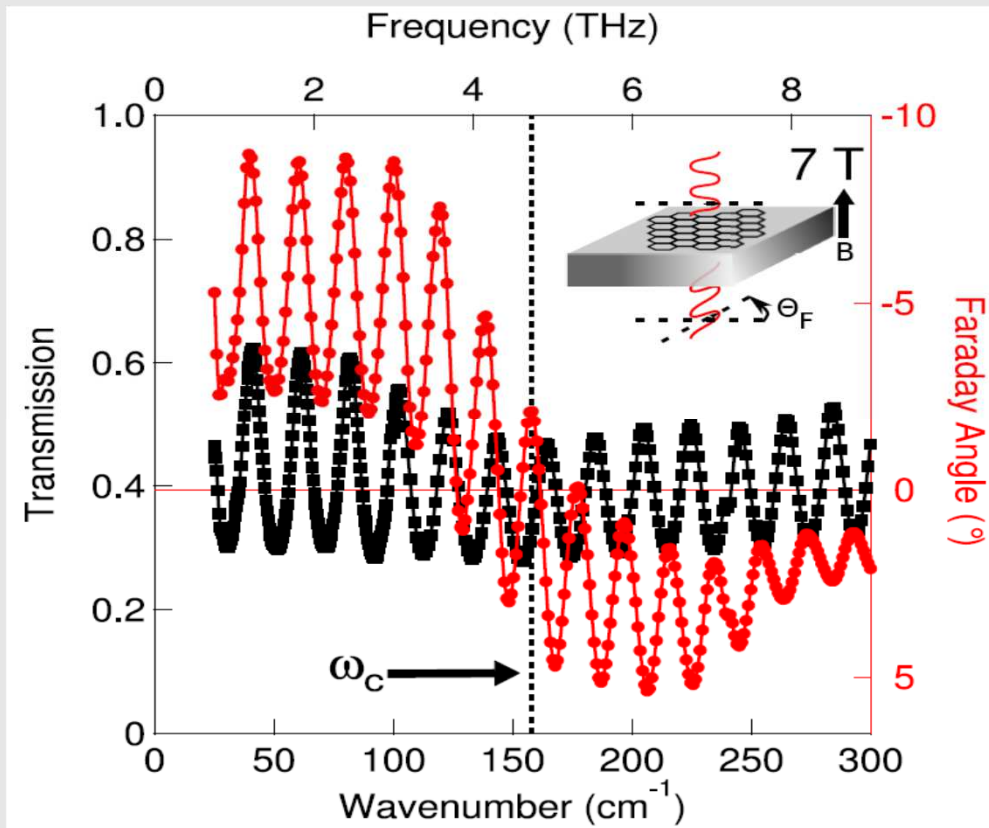
A. Fallahi & Perruisseau-Carrier, APL 101, 231605 (2012)

D. L. Sounas & C. Caloz, IEEE Trans. Microw. Theory Tech. 60, 901 (2012)

Y. Zhou *et al*, Phys.Chem. Chem. Phys. (2013)

...

# Fabry-Perot enhanced Faraday rotation in graphene



- interference in the substrate
- Rotation up to 9°
- Rotation and transmission increase simultaneously



Theory of FR in graphene in a cavity:  
A. Ferreira *et al*, PRB 84, 235410 (2011)

For details: Nicolas Ubrig, poster on Wednesday

# Summary and outlook

- ✚ Magneto-optics (MO) is useful for routine characterization
- ✚ Giant THz Faraday rotation (present record is  $9^\circ$ )
- ✚ Robust (magneto-)plasmons due to nanoscale defects (i.e. steps and wrinkles)
- ✚ Strong interplay between MO and plasmonic effects (cyclotron mass about  $10^2$  smaller than in noble metals)
- ✚ Graphene is promising for MO applications