



Coherent and broadband enhanced optical absorption in graphene

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Acknowledgments

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Funding:



PHILIPS

Outline

Motivation and introduction

- Coherent absorption
- Measurements and modeling
- Optimization
- Conclusions

Motivation



Problem: the low absolute value of the absorption of graphene (2.3% of the incident light) limits the photocurrent efficiency

F. Xia, T. Mueller, Yu. Li, A. Valdes-Garcia and P. Avouris, Nat. Nanotech. 2009, 4, 839. F. Bonaccorso, Z. Sun, T. Hasan and A. C. Ferrari, Nat. Photon., 2010, 4, 611.

Enhanced photodetection



M. Furchi, et al., Nano Lett., 2012, 12, 2773 M. Engel, M. Steiner, A. Lombardo, A. C. Ferrari, H. v. Lohneysen, P. Avouris and R. Krupke, Nat. Commun., 2012, 3, 906.

Hybrid devices



G. Konstantatos, M. Badioli, L. Gaudreau, , J. Osmond, M. Bernechea, F. P. Garcia de Arquer, F. Gatti and F. H. L. Koppens, Nat. Nanotech., 2012, 7, 363.
Z. Fang, Y. Wang, Z. Liu, A. Schlather, P. M. Ajayan, F. H. L. Koppens, P. Nordlander and N. J. Halas, ACS Nano, 2012, 6, 10222

Plasmonics to enhance absorption



S. Thongrattanasiri, F. H. L. Koppens and J. Garcia de Abajo, Phys. Rev. Lett., 2012, 108, 047401. Y. Nikitin, F. Guinea and L. Martin Moreno, Appl. Phys. Lett., 2012, 101, 151119.

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Scattering problem

Perfect energy transfer between incoming channels (optical modes) and the structure.



Building a destructive interference pattern in the far-field will cause the light to be trapped in the structure. The material losses will eventually dissipate the energy.



Weakly absorbing material



Coherent Perfect Absorption

2 channels CPA





Weakly absorbing material

W. Wan, Y. Chong, L.Ge, H. Noh, A. Douglas Stone and H. Cao, Science, 2011, 331, 889.

From 2 to 1 channel

1 channel CPA (single port reflector)



W. Wan, Y. Chong, L. Ge, H. Noh, A. Douglas Stone and H. Cao, Science, 2011, 331, 889.

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Graphene characterization



CVD deposition of graphene on copper foils. PMMA transfer onto silica substrates (turbostratic deposition)

G. Pirruccio, G. Lozano, L. Martin Moreno and J. Gomez Rivas, submitted. V. G. Kravets, A.N. Grigorenko, R. R. Nair, P. Blake, S. Anissimova, K. S. Novoselov and A.K. Geim, Phys. Rev. B, 2010, 81, 155413

Sample





Absorptance measurements

Absorptance (5 layers)



When $\theta^* < \theta < \theta_{c,1}$ the incident wave is transmitted through the prism-silica layer interface and undergoes total internal reflection at the silica-substrate interface. \longrightarrow T=0

Broadband absorption





Polarization independence (5 layers)



Transfer matrix calculation



V. G. Kravets, A.N. Grigorenko, R. R. Nair, P. Blake, S. Anissimova, K. S. Novoselov and A.K. Geim, Phys. Rev. B, 2010, 81, 155413

Interpretation

$$S = \begin{pmatrix} t_{1234} & r_{4321} \\ r_{1234} & t_{4321} \end{pmatrix}$$

$$r_{12}$$

$$r_{1234} = r_{12} + r_{scatt}$$

$$R = \left| \left| r_{12} \right| e^{i\varphi_{12}} + \left| r_{scatt} \right| e^{i\varphi_{scatt}} \right|^2$$

$$R = 0 \iff \begin{cases} |r_{12}| = |r_{1234}|^2 = R \\ |S_{11}|^2 = |t_{1234}|^2 \propto T \\ A = 1 - R - T \end{cases}$$

$$R = 0 \iff \begin{cases} |r_{12}| = |r_{scatt}| \\ \varphi_{12} - \varphi_{scatt} = \pi \\ T = 0 \end{cases}$$

$$T = 0$$

$$T = 0$$

$$T = 0$$

$$T = 0$$

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 \wedge

$$\min_{d_{silica}} (R) = \lim_{d_{silica}} \left(\left| \frac{r_{12}^{p}(\lambda_{0}, \theta) + r_{234}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}}{1 + r_{12}^{p}(\lambda_{0}, \theta)r_{234}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}} \right|^{2} + \frac{r_{12}^{p}(\lambda_{0}, \theta)r_{234}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}}{1 + r_{12}^{p}(\lambda_{0}, \theta)r_{234}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}} + \frac{r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}}{1 + r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}} + \frac{r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}}{1 + r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}} + \frac{r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}}{1 + r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}, \theta)e^{-2ik_{x2}d_{silica}}} + \frac{r_{12}^{p}(\lambda_{0}, \theta)r_{23}^{p}(\lambda_{0}$$





 λ_0 =540 nm











Conclusions

- We have experimentally demonstrated broadband and enhanced absorption in graphene
- This enhancement is explained in terms of coherent absorption arising from interference and dissipation in a multilayer structure.
- For 10 layers of graphene it is possible to enhance the absorptance over 91%. For the 5 layer and monolayer sample a similar analysis leads to a maximum absorption of 76% and 15%.

Controlling the absorption





Microcavity control



illumination

M. Engel, M. Steiner, A. Lombardo, A. C. Ferrari, H. v. Lohneysen, P. Avouris and R. Krupke, Nat. Commun., 2012, 3, 906.

Plasmonics to enhance absorption



Y. Nikitin, F. Guinea and L. Martin Moreno, Appl. Phys. Lett., 2012, 101, 151119. T. J. Echtermeyer, L. Britnell, P. K. Jasnos, A. Lombardo, R. V. Gorbachev, A. N. Grigorenko, A. K. Geim, A. C. Ferrari and K. S. Novoselov, Nat. Commun., 2001, 2, 458.

Plasmonics to enhance absorption



T. J. Echtermeyer, L. Britnell, P. K. Jasnos, A. Lombardo, R. V. Gorbachev, A. N. Grigorenko, A. K. Geim, A. C. Ferrari and K. S. Novoselov, Nat. Commun., 2001, 2, 458.

Total absorption



S. Thongrattanasiri, F. H. L. Koppens and J. Garcia de Abajo, Phys. Rev. Lett., 2012, 108, 047401.



Raman shift / cm-1











1 layer G-line position







1 layer intensity of D-line









Coherent absorption

$$S = \begin{pmatrix} t_{1234}^{p} & r_{4321}^{p} \\ r_{1234}^{p} & t_{4321}^{p} \end{pmatrix}$$

$$\begin{cases} \left| S_{21} \right|^2 = \left| r_{1234}^p \right|^2 = R = 0 \\ \left| S_{11} \right|^2 = \left| t_{1234}^p \right|^2 \propto T = 0 \end{cases}$$

$$A = 1 - R - T$$

