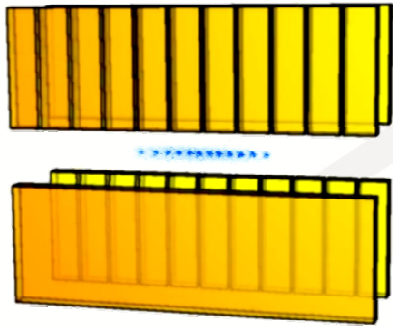




Tailoring Coupling Constants to Simulate Spin Models in a Segmented Trap



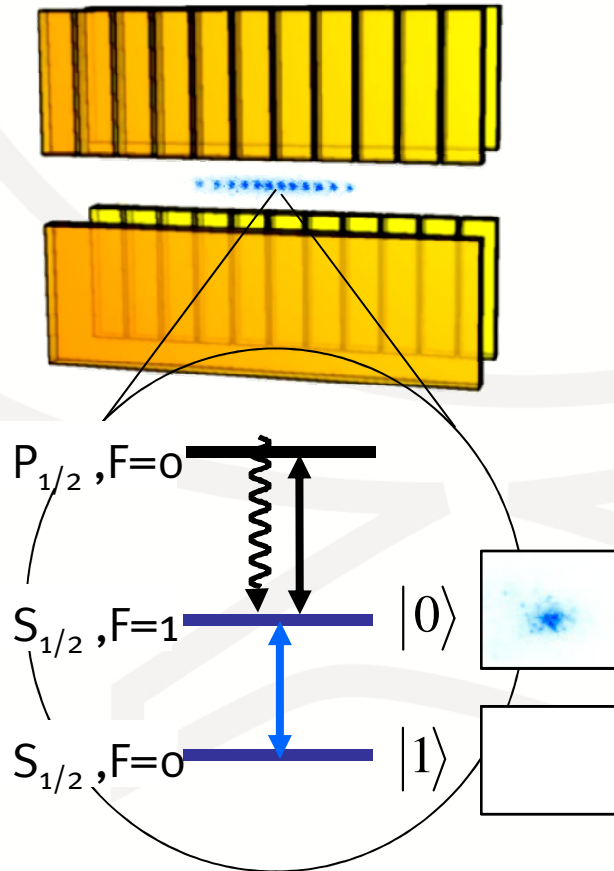
M. Johanning¹, S. Zipilli², M. Giampaolo²,
F. Illuminati², and Ch. Wunderlich¹

¹Faculty of Science and Technology, Dep. of Physics, University of Siegen, Germany

²Dep. of Mathematics and Informatics, University of Salerno, Italia



Microwaves & Hyperfine Qubit



- use microwaves and rf for coherent state manipulation of $^{171}\text{Yb}^+$
- the magnetic gradient makes Zeeman splitting position dependent
- the resonance frequency is unique for each ion

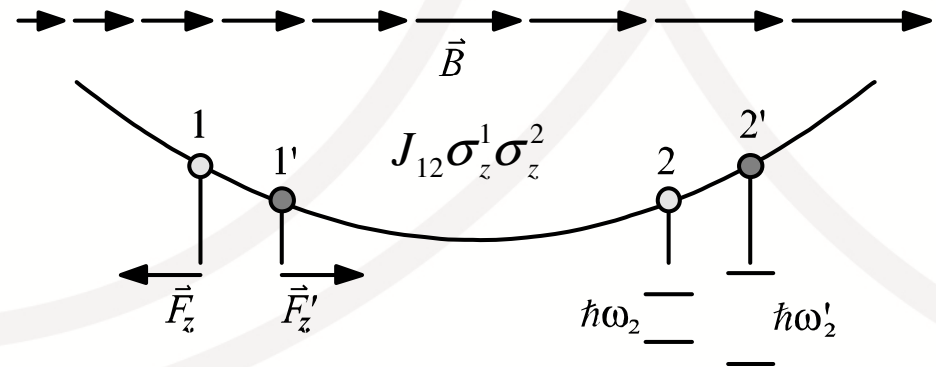
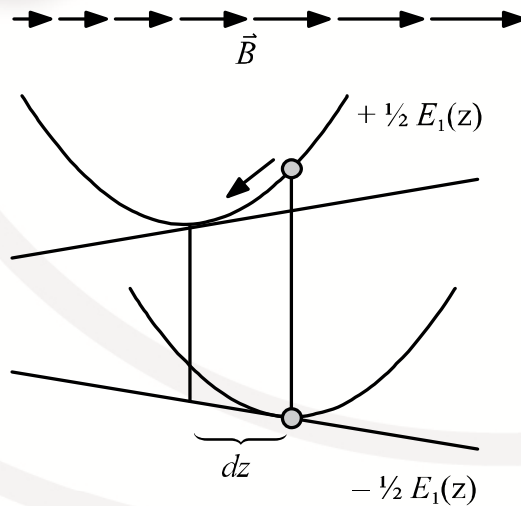
F. Mintert & Ch. Wunderlich, *Phys. Rev. Lett.* 87, 4 (2001); 91, 029902 (2003).

Ch. Wunderlich, *Laser Physics at the Limit*, Springer Berlin;

C. Ospelkaus et al., *Phys. Rev. Lett.* 101, 090502 (2008)



Spin-Spin Coupling



- total energy: sum of Zeeman shift and the trapping potential
- equilibrium positions becomes state dependent, mimicks momentum transfer; **M**agnetic **G**radient **I**nduced **C**oupling: **MAGIC**
- via Coulomb repulsion each ion affects the equilibrium positions of all other ions, changing their eigen values

F. Mintert & Ch. Wunderlich, *Phys. Rev. Lett.* 87, 4 (2001); 91, 029902 (2003).

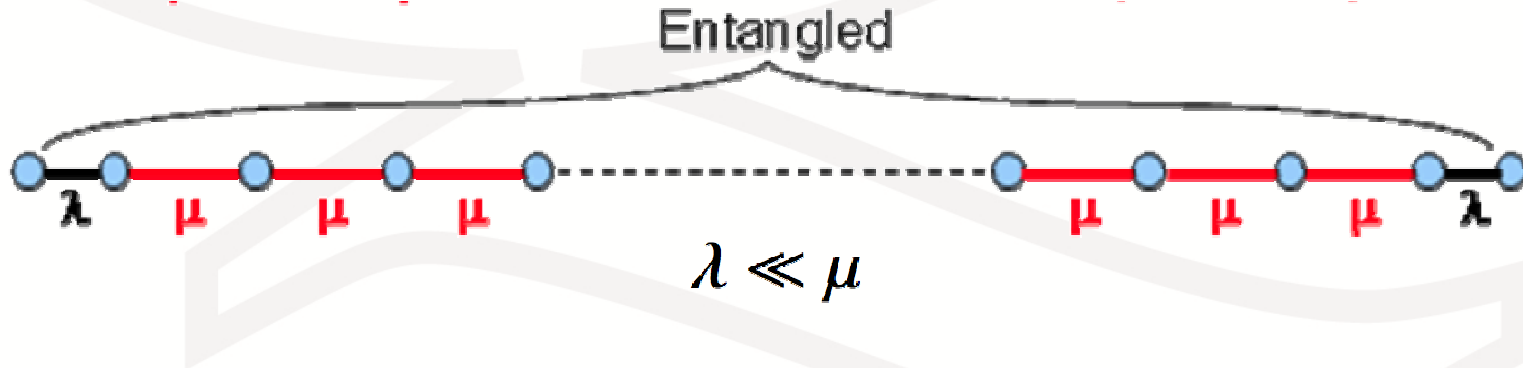
Ch. Wunderlich, *Laser Physics at the Limit*, Springer Berlin;

Ch. Wunderlich & Ch. Balzer, *Adv. in At. Mol. Opt. Phys.* 49, 293–372 (2003).



Long distance entanglement (LDE)

- ground-state, indirect, end-to-end entanglement in spin-chains, useful for quantum bus
- loosely coupled end/messenger ions



L. Campos Venuti et al., Phys. Rev. Lett., 96, 247206 (2006);

L. Campos Venuti et al., Phys. Rev. A 76, 052328 (2007);

S. M. Giampaolo et al., New J. of Phys. 12, 025019 (2010).



LDE Conditions

- weak coupling between the two external spins and the remaining spins
- non-degenerate ground state, no B field
- examples: models with competing interactions along orthogonal axes as XY, XYZ, Heisenberg...

$$H_{XY} = \sum_{j,k} (J_x \sigma_j^x \sigma_k^x + J_y \sigma_j^y \sigma_k^y)$$

$$H_{XYZ} = \sum_{j,k} (J_x \sigma_j^x \sigma_k^x + J_y \sigma_j^y \sigma_k^y + J_z \sigma_j^z \sigma_k^z)$$

$$H_{Heisenberg} = J \sum_{j,k} (\sigma_j^x \sigma_k^x + \sigma_j^y \sigma_k^y + \sigma_j^z \sigma_k^z)$$



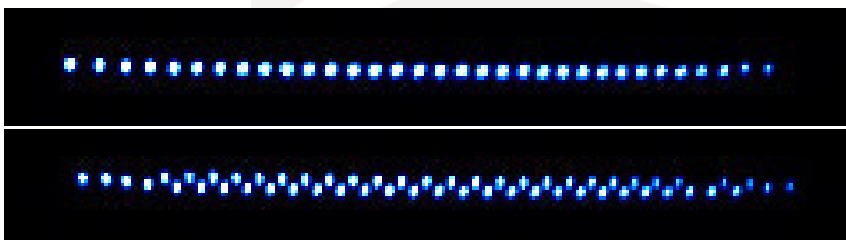
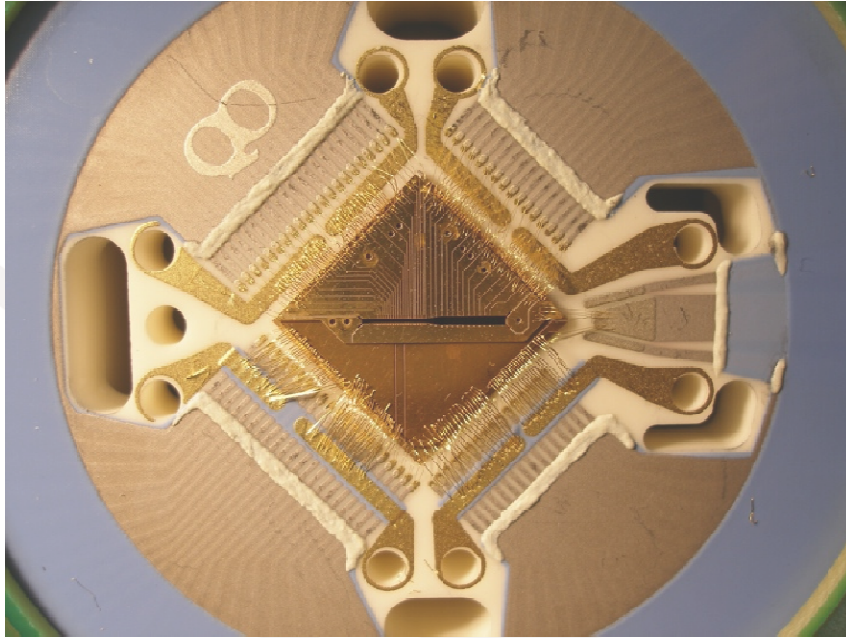
LDE Work Plan

1. Manipulation of the spin-spin couplings:
 - weakly coupled end-spins
2. Engineering of Hamiltonian:
 - interaction along z and along another axis

...
3. Preparation of the ground state:
 - adiabatic variation of some Hamiltonian parameter



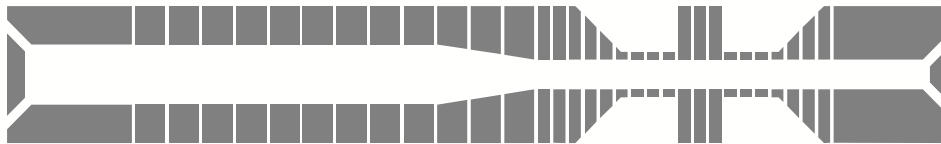
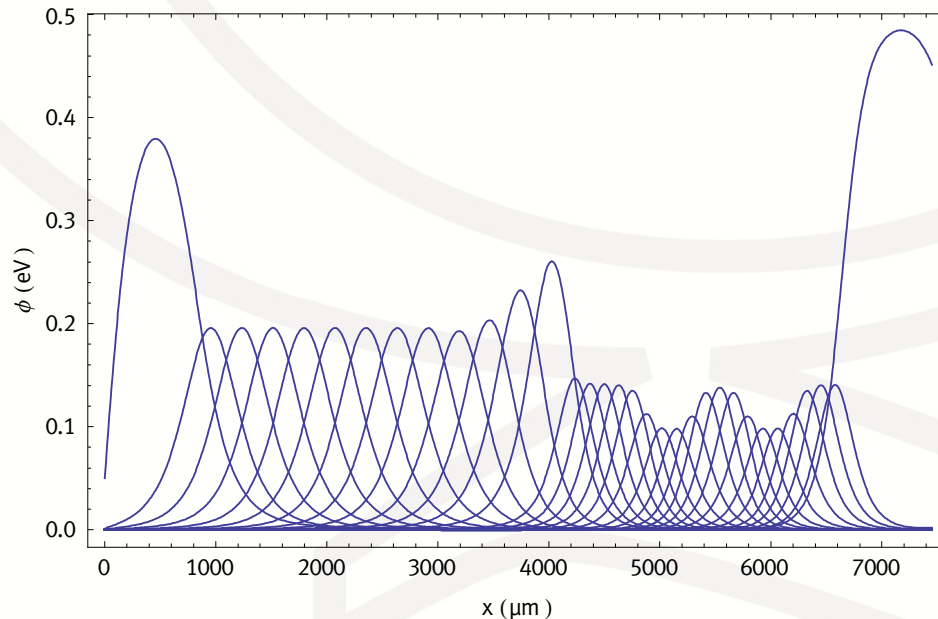
Segmented μ trap



- linear Paul trap, three layer structure
- basic design from Mainz group (Schmidt-Kaler), altered by our group
- segments allow flexible axial trapping potentials and multiple trapping zones
- carrier acts as vacuum interface (!)



Potential Simulation



- boundary element method [1]
- structures reflected in potentials simulation:
 - width of electrodes
 - electrode separation
 - middle layer coils
- helpful for shuttling, tailoring of potentials



Extension: Tikhonov Regularization

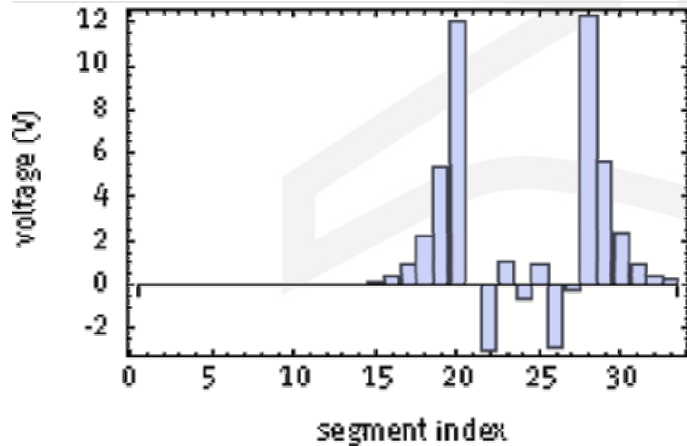
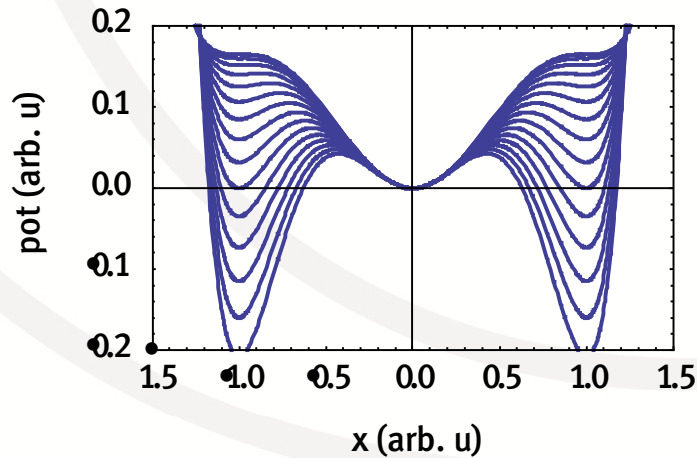
- we want to get a specific potential $\vec{\phi}$
- we need to solve for the voltages V_i :

$$\hat{A}\vec{V} = \vec{\phi}$$

- often ill conditioned, instead we minimize $\|\hat{A}\vec{V} - \vec{\phi}\|^2$
(least squares)
- boundary condition $\|\hat{A}\vec{V} - \vec{\phi}\|^2 - \|\hat{\Gamma}\vec{V}\|^2$
- if Tikhonov matrix $\hat{\Gamma}$ equals identity, algorithm prefers solution with smaller norm



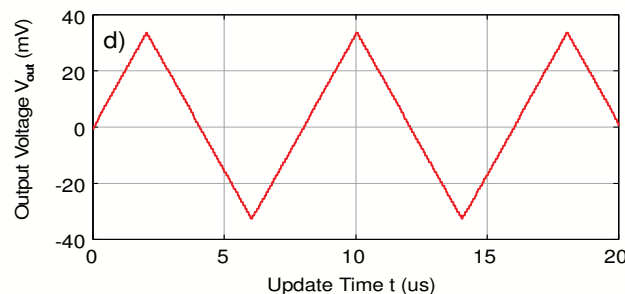
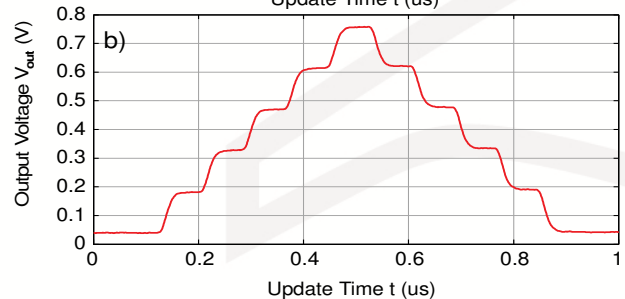
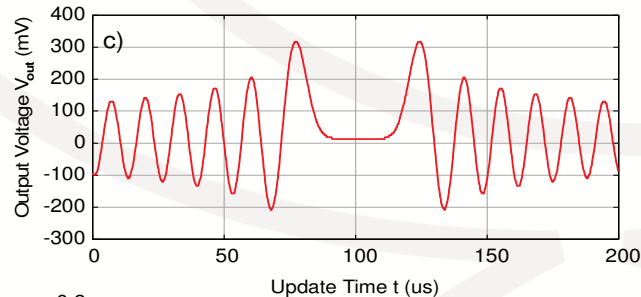
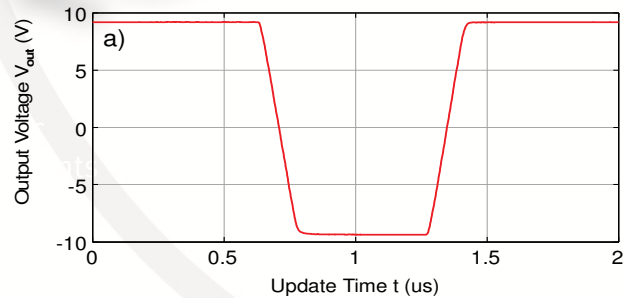
Designing the Coupling



- triple well potential defines couplings with three parameters
- find appropriate triple well
- minima can be much closer than 2 segment widths (!)
- find voltages



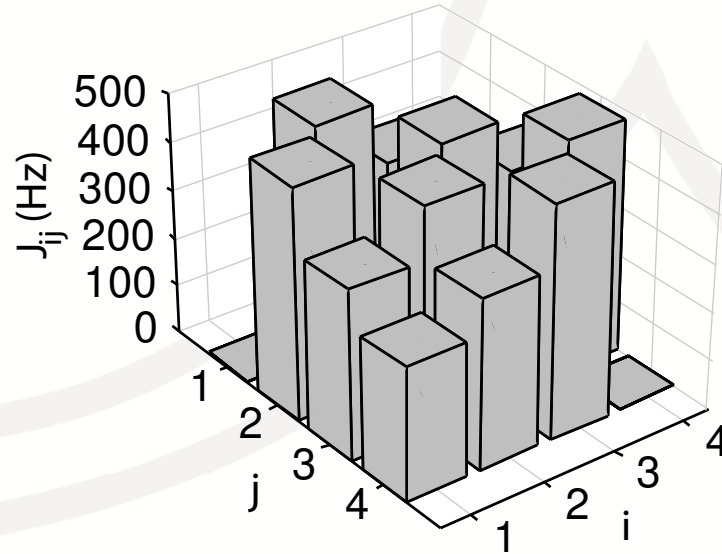
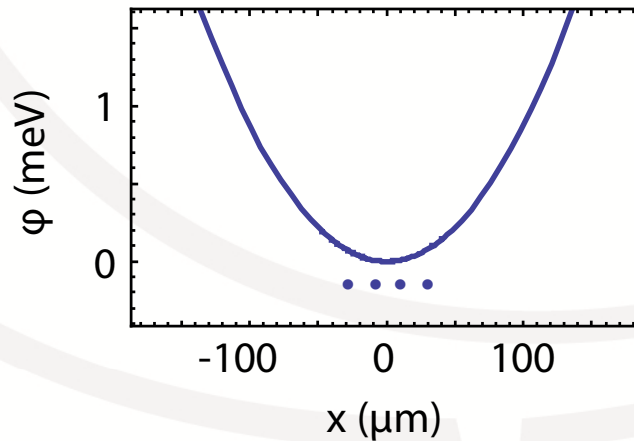
Multi Channel AWG



- 24 independent channels
- synchronous, arbitrary sequences
- amplitude ± 10 V
- update rate 20 MHz
- 16 bit resolution
- low noise & drift
- loss free transmission up to 2 m
- freely programmable via USB



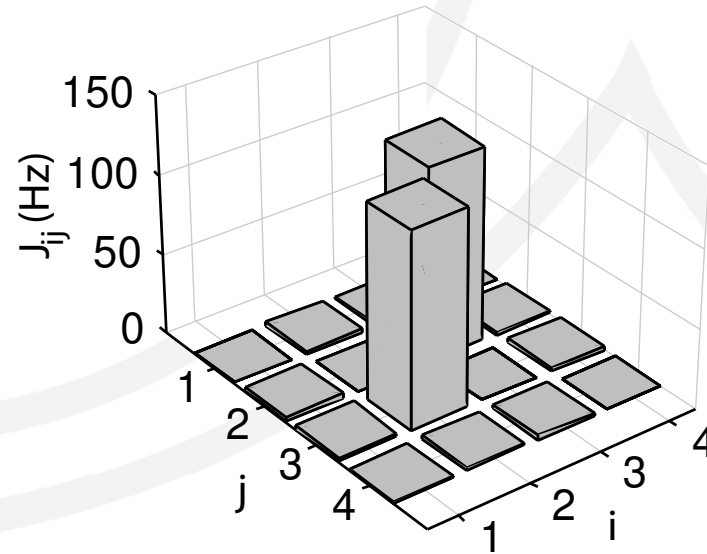
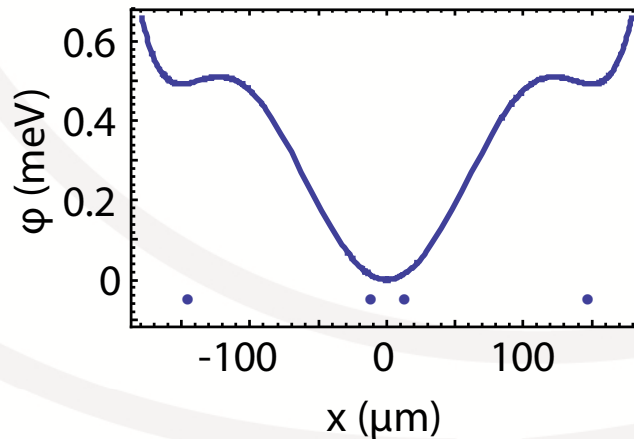
1. Tailoring Coupling Constants



- harmonic coupling:
 - delocalized normal modes
 - all ions are strongly coupled, long range
 - global scaling by trap frequency and gradient



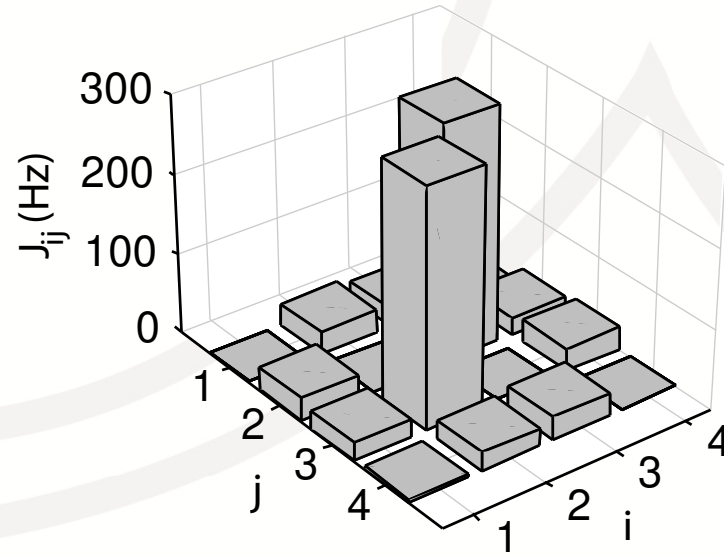
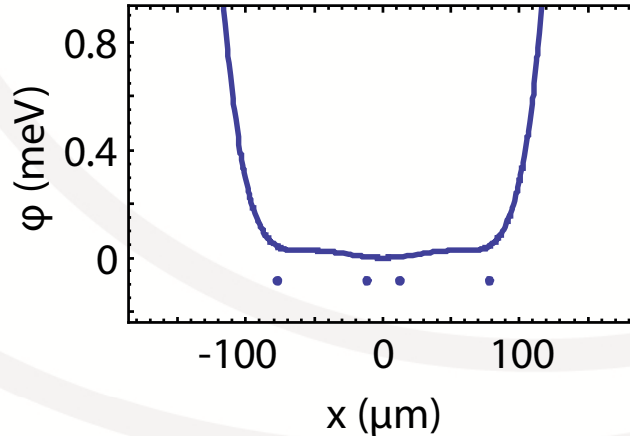
1. Tailoring Coupling Constants



- triple well:
 - localized modes within wells
 - ions within the same well are strongly coupled
 - inter-well coupling strongly reduced
 - local scaling by curvature and gradient



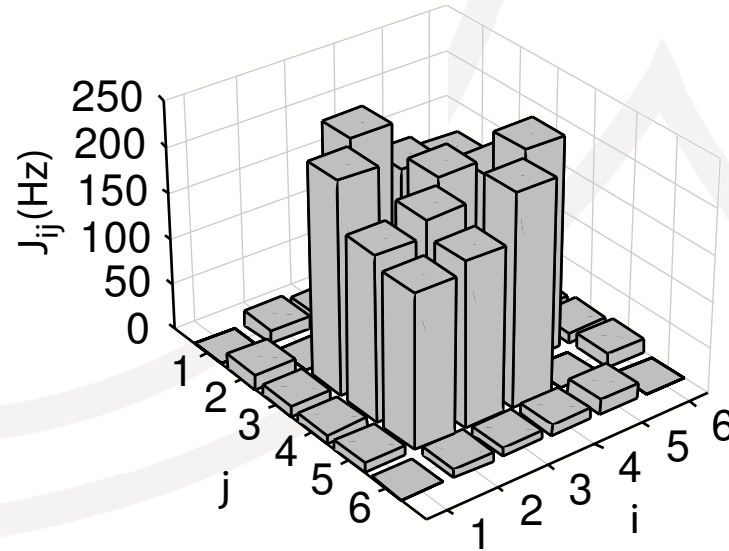
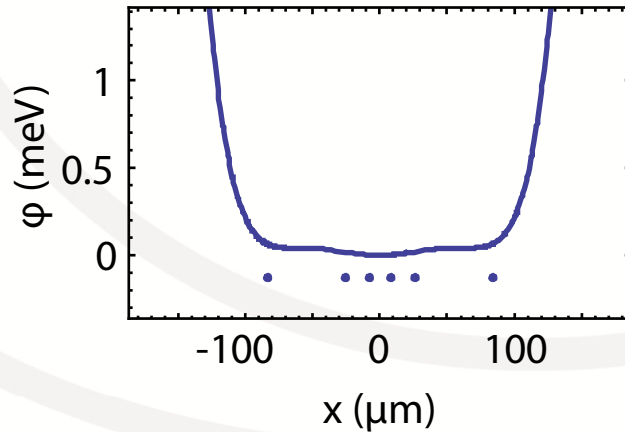
1. Tailoring Coupling Constants



- anharmonic trap:
 - some modes localized, localization scalable
 - scalable inter-well coupling



1. Tailoring Coupling Constants



- anharmonic trap
 - some modes localized, localization scalable
 - scalable inter-well coupling
 - similar conditions for more ions in centre well



2. Mimic Other Hamiltonians

- engineer Hamiltonian dynamics using microwave fields which sequentially drive the atomic spins

$$H = \sum_{j=1}^N \omega_j \sigma_j^z + \sum_{j,k} J_{j,k} \sigma_j^z \sigma_k^z - i \sum_m \Omega_m(t) \left[\sigma_{j_m}^+ e^{i\nu_{j_m} t} - \sigma_{j_m}^- e^{-i\nu_{j_m} t} \right]$$

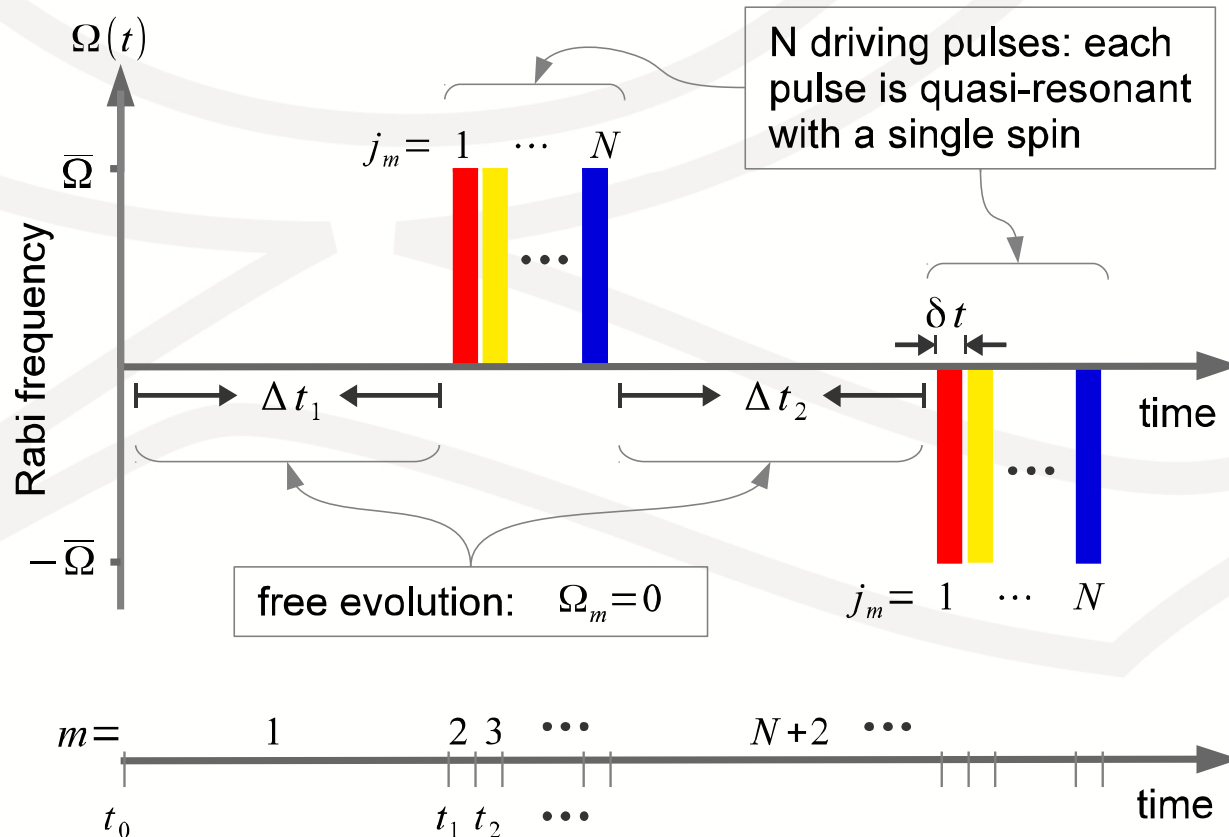
- every time-step: driving field is quasi-resonant with a single spin resonance; effect on the other off-resonant spins is negligible
- bracketing a free evolution between μ wave pulses mimics coupling along a different axis:

$$e^{-i\pi\sigma_j^y/4} \sigma_j^z e^{i\pi\sigma_j^y/4} = \sigma_j^x$$



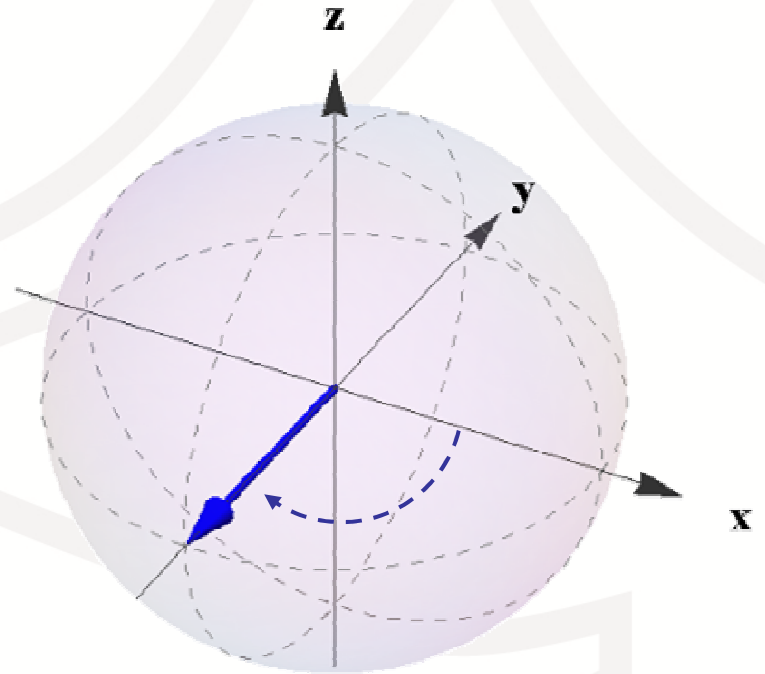
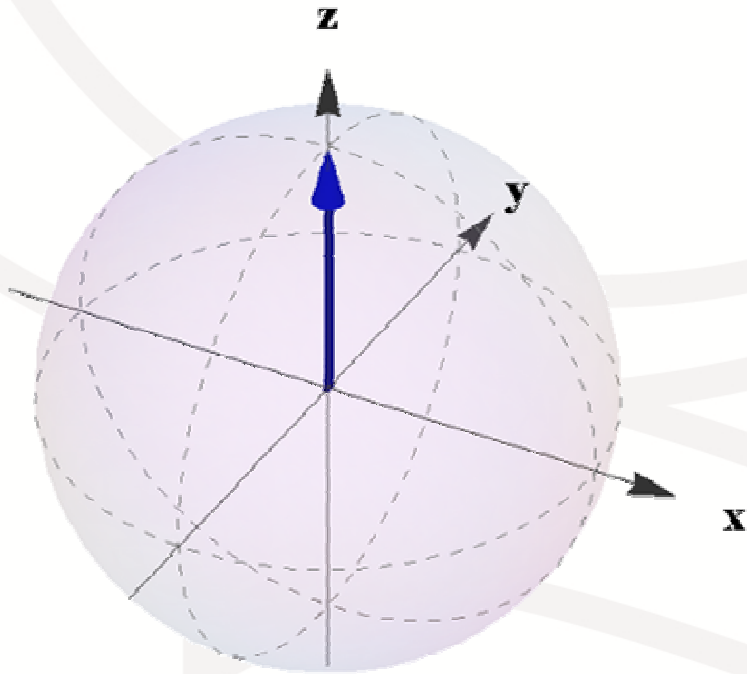
Sequential Forth and Back

$$e^{-iHt} \simeq \left(\prod_j e^{-\frac{i\pi}{4} \sigma_j^y} \right) e^{-iH_I^{(z)} \Delta t} \left(\prod_j e^{\frac{i\pi}{4} \sigma_j^y} \right) = e^{-iH_I^{(x)} \Delta t}$$





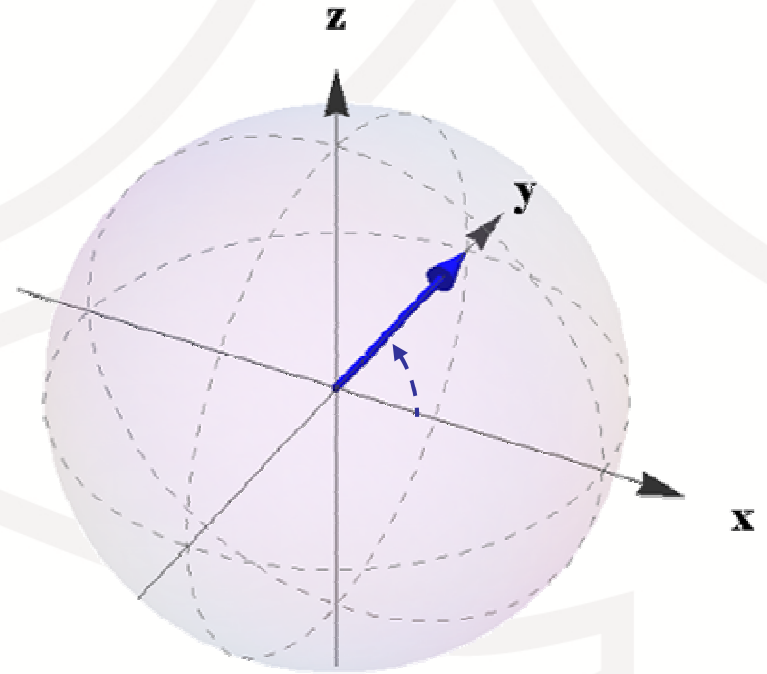
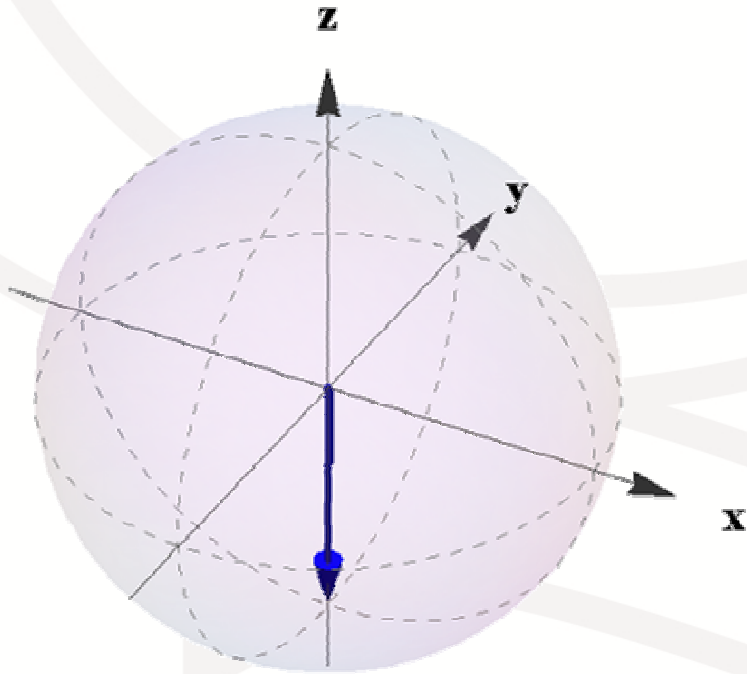
$\sigma^z \sigma^z$ -Coupling



cw precession



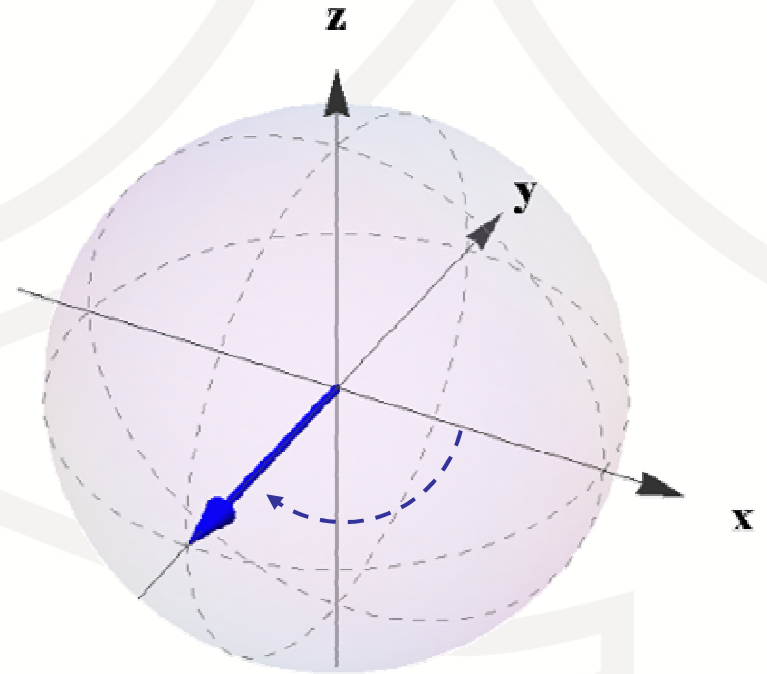
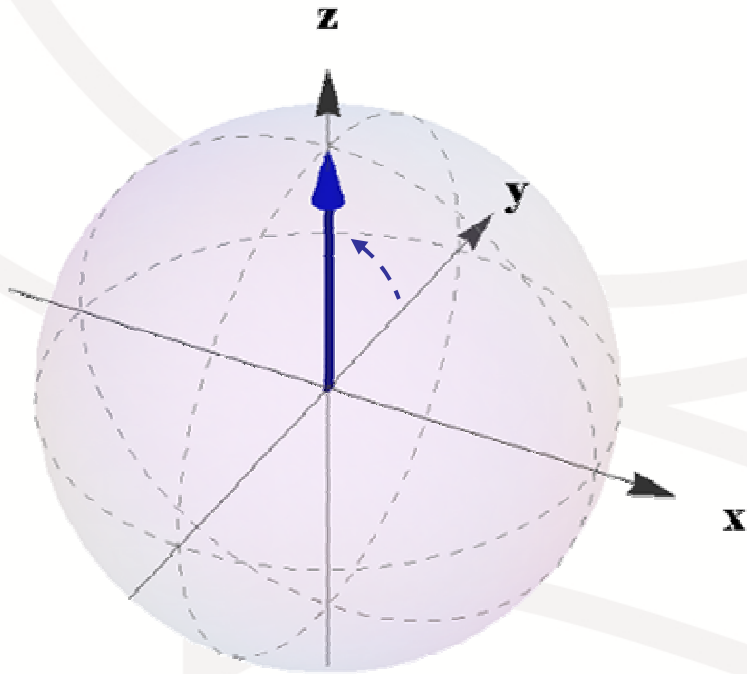
$\sigma^z \sigma^z$ -Coupling



ccw precession



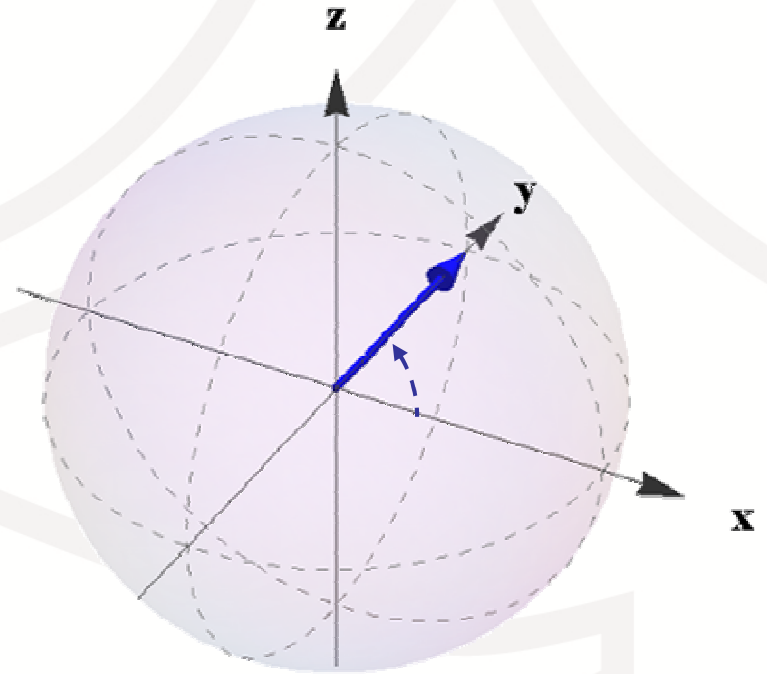
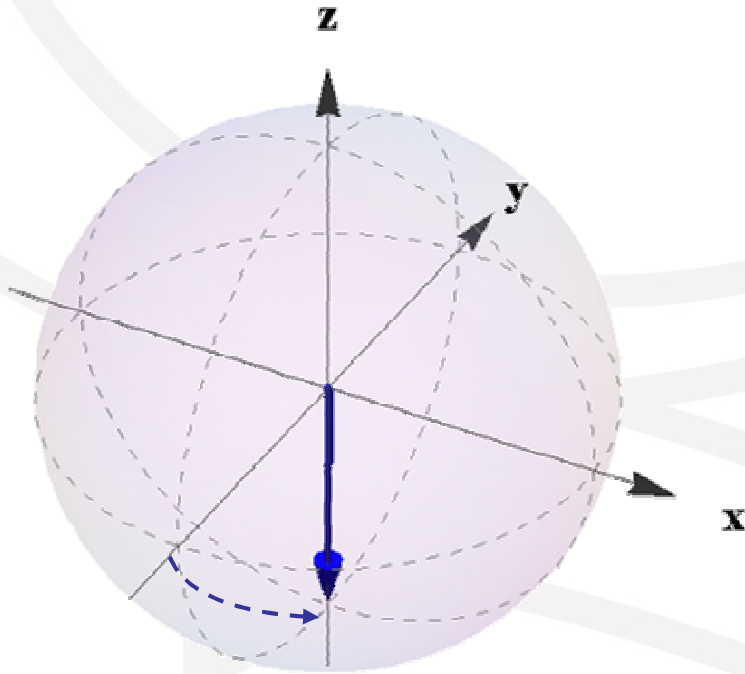
$\sigma^y \sigma^z$ -Coupling



cw precession



$\sigma^y \sigma^z$ -Coupling



ccw precession

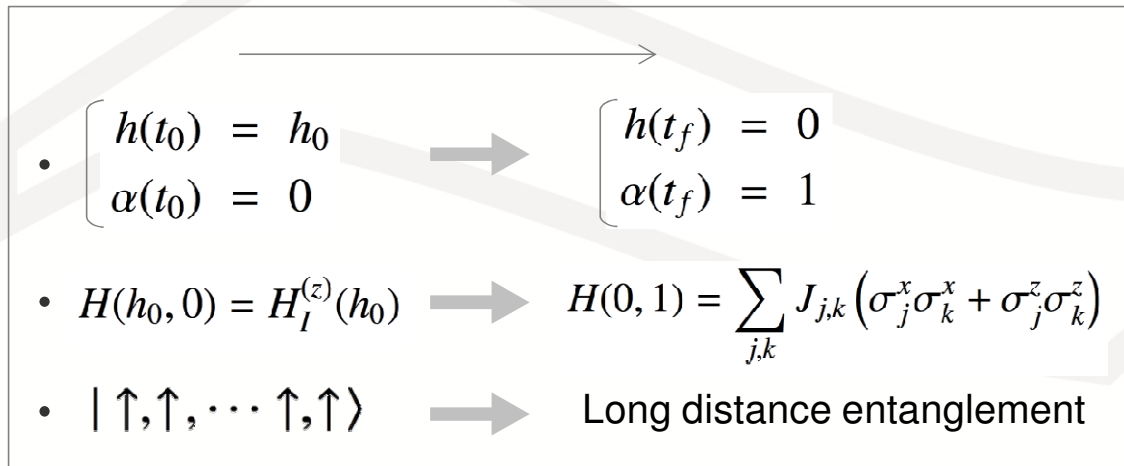
$$\text{or: } e^{-i\pi\sigma_j^y/4} \sigma_j^z e^{i\pi\sigma_j^y/4} = \sigma_j^x$$



3. Adiabatic Prep. of Ground State

Under a slow variation of some Hamiltonian parameter, a system initially in an eigenstate will follow the instantaneous eigenstate

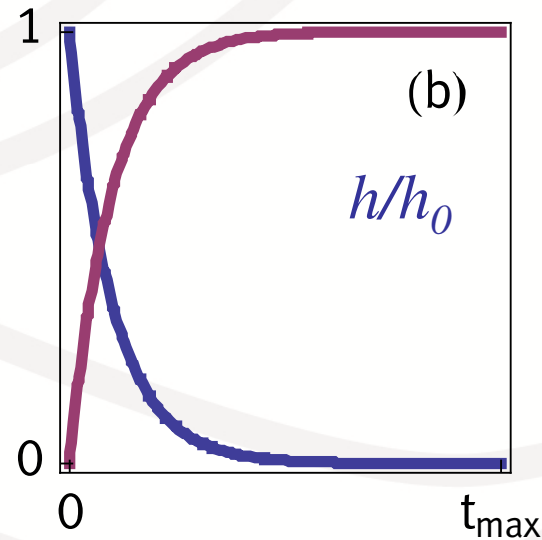
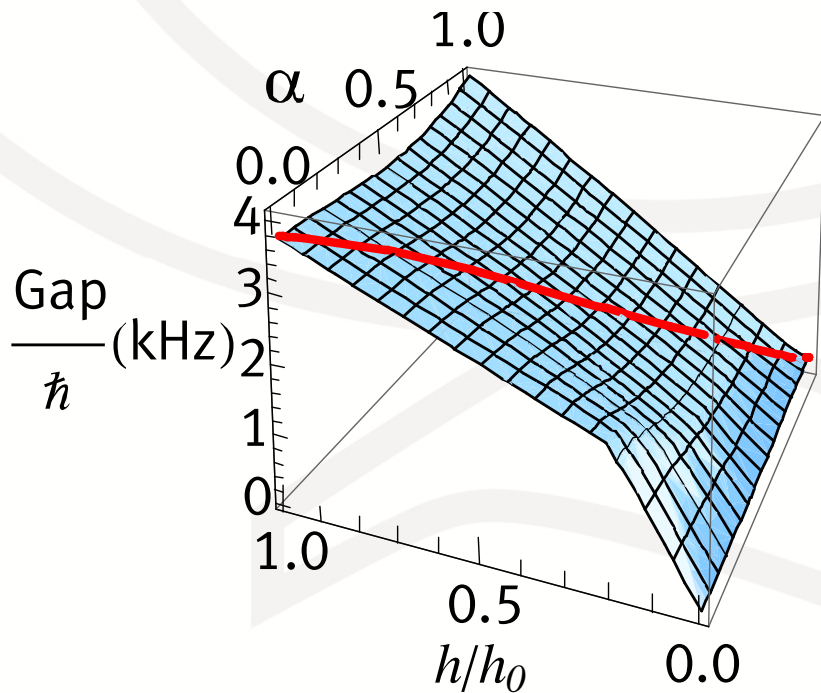
$$H_{\text{eff}}(h, \alpha) = H_I^{(z)}(h) + \alpha H_I^{(x)}(h)$$





Parameters: Time Dependence

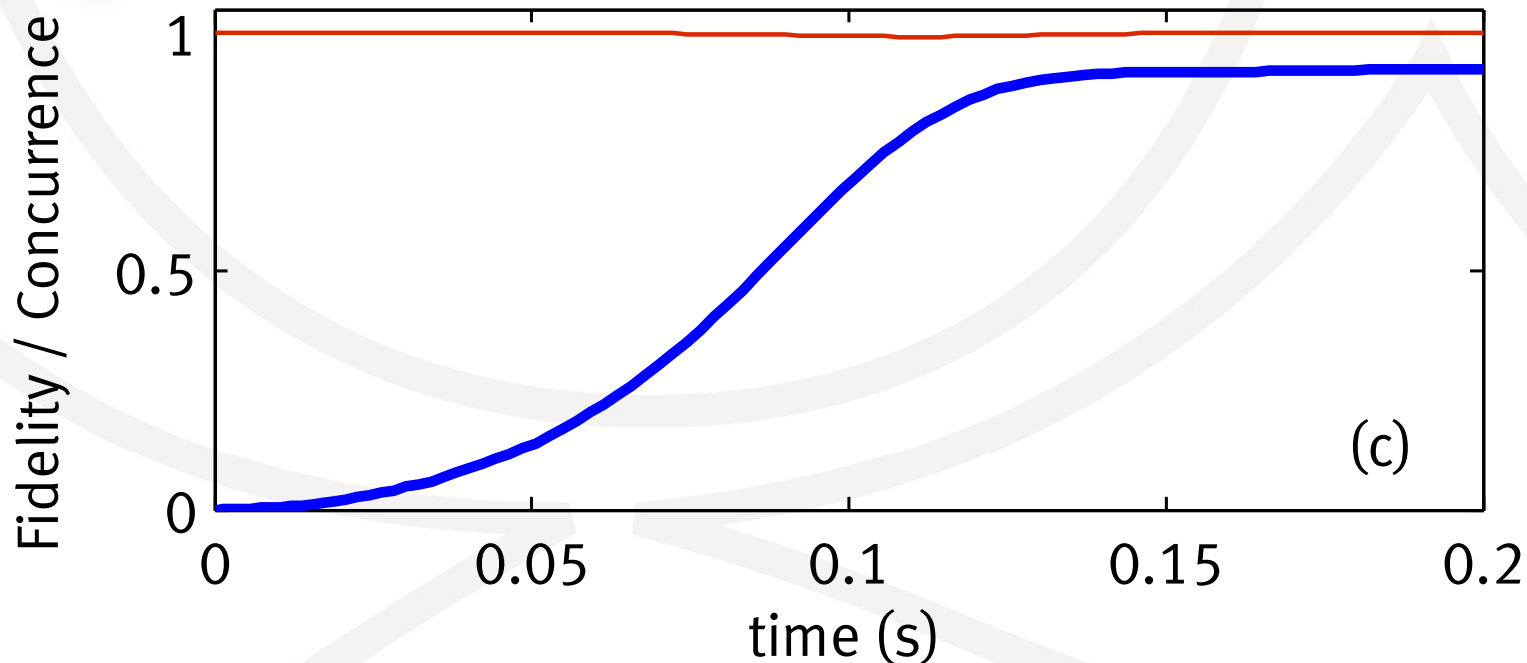
$$e^{-i(H_{I \sin g}^x + H_{I \sin g}^z)\tau} \approx \left(e^{-iH_{I \sin g}^x \frac{\tau}{n}} e^{-iH_{I \sin g}^z \frac{\tau}{n}} \right)^n \alpha$$



- initially the large gap allows for fast sweeping
- slow down towards smaller gaps to remain adiabatic



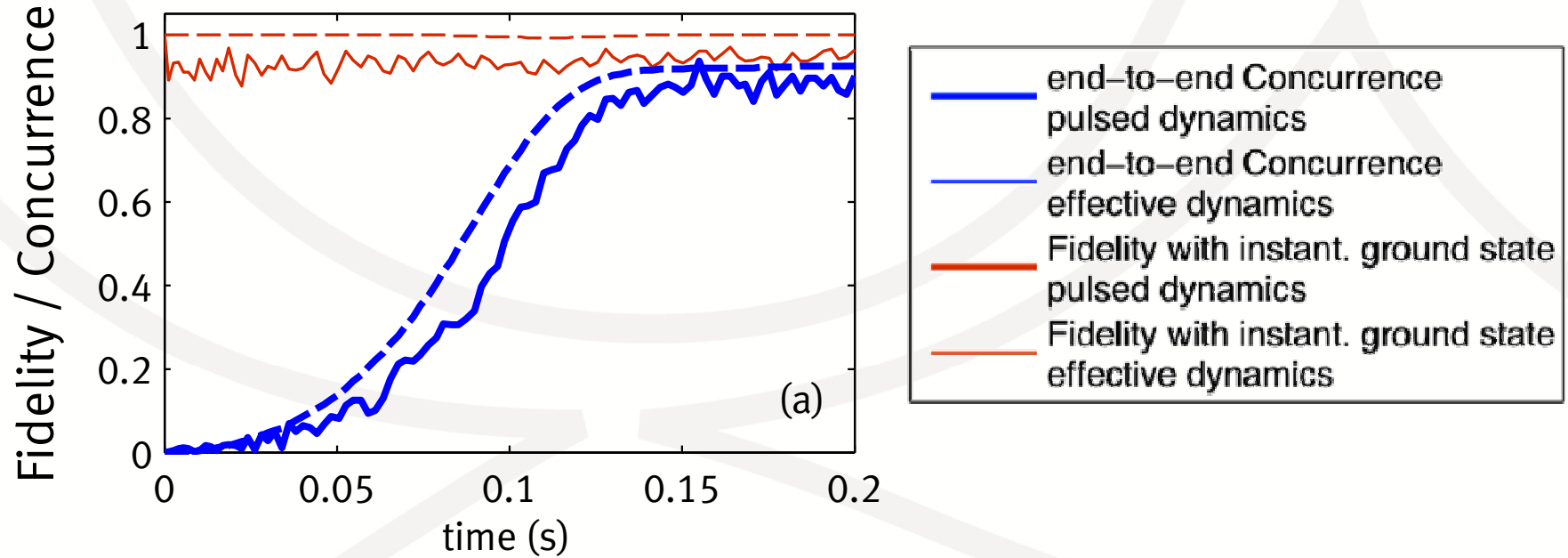
Adiabatic transformation



- adiabaticity well fulfilled
- final state: high fidelity with instantaneous ground state (red curve)
- large concurrence / entanglement (blue curve)



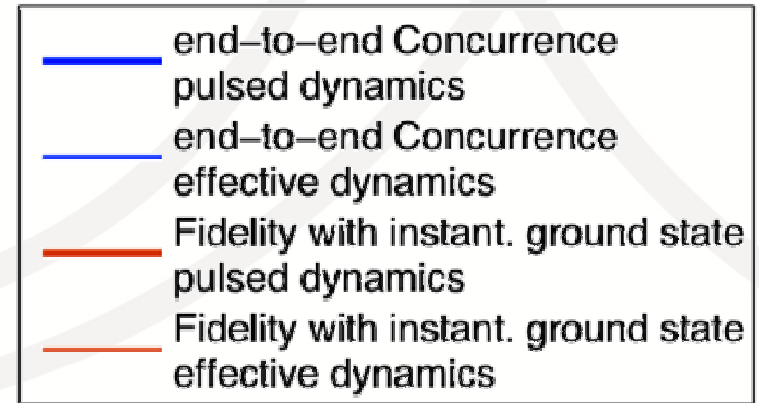
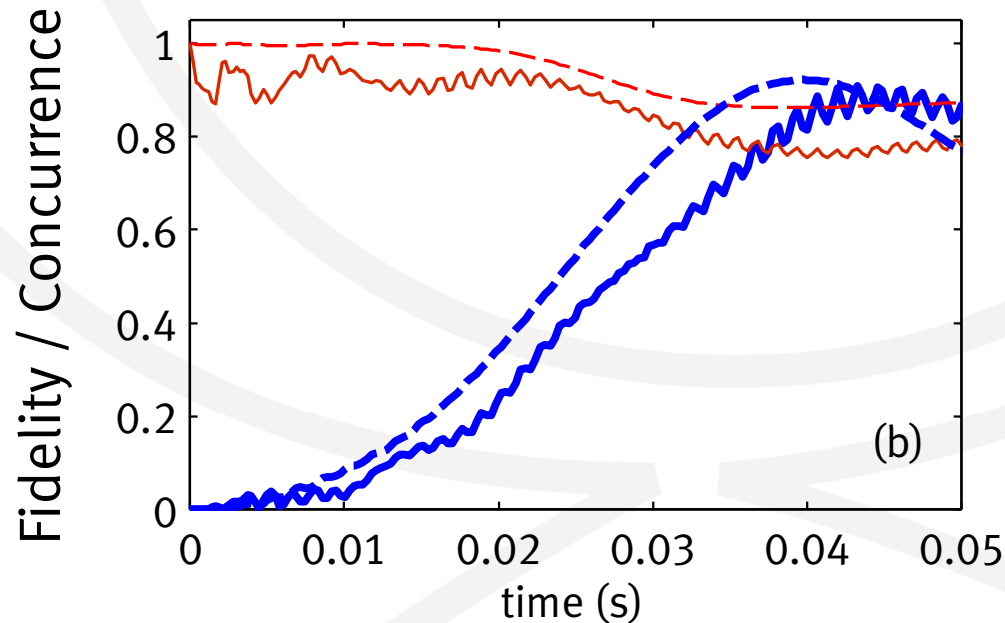
Adiabatic Transformation + Pulsed Dynamics



- reasonable agreement with effective Hamiltonian
- quantum simulation
- ground state entanglement



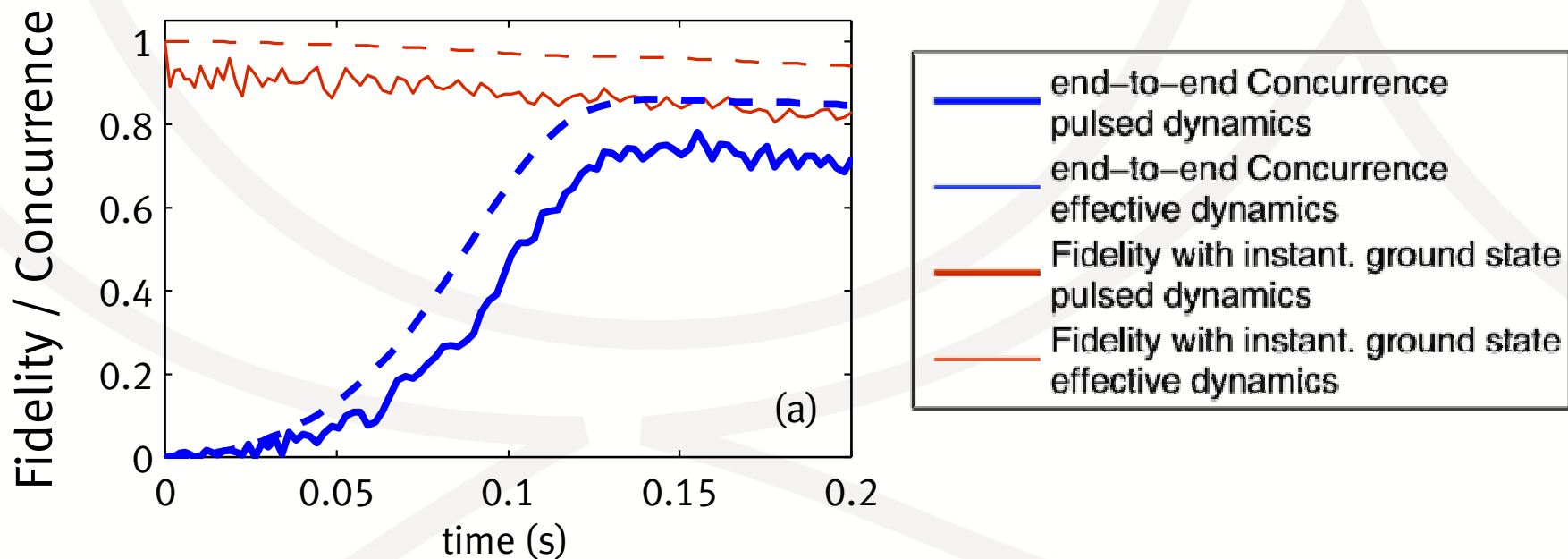
Increase Ramp Speed



- lower fidelity
- oscillatory behavior
- still good concurrence



Adiabatic Transformation + Pulsed Dynamics + Dephasing



- fidelity decreases over time
- still good concurrence



LDE Outlook

- tailoring the axial trapping potential can generate coupling patterns useful for LDE
- an effective XZ Hamiltonian by bracketing the evolution time with resonant microwave pulses
- changing the relative evolution times allows to sweep from Ising to XZ Hamiltonian
- simulations show ground state population, concurrence
- challenges: large gradients, coherence



Microtrap Witches And Wizards



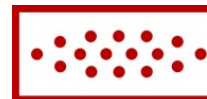
Thank you for your attention!

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