

Chiral Projected Entangled Pair States

Thorsten B. Wahl



Rudolf Peierls Centre for Theoretical Physics, University of Oxford

19 February 2016



MPQ

Max-Planck Institute of Quantum
Optics, Garching, Germany

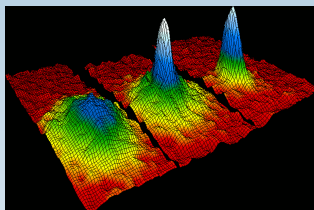
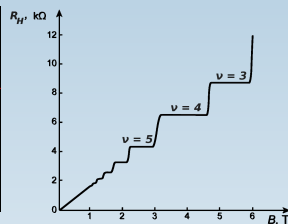
- QCCC Elitenetzwerk Bayern
- Alexander von Humboldt Foundation
- EU Integrated Project SIQS

H.-H. Tu, N. Schuch, S. Haßler, S. Yang, J. I. Cirac

source: Wikipedia



Superconductivity

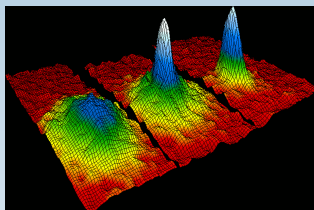
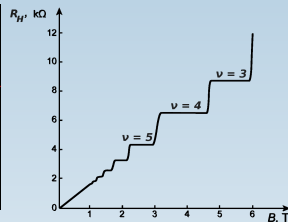
Bose-Einstein
condensate in Rb-87

Quantum Hall Effect

source: Wikipedia



Superconductivity

Bose-Einstein
condensate in Rb-87

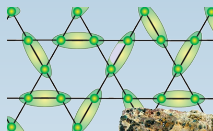
Quantum Hall Effect

Approxiation Schemes

- Mean Field Theory / Hartree-Fock



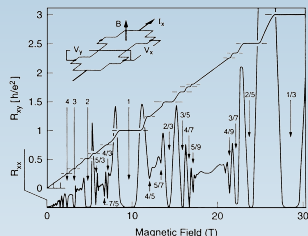
High temperature
Superconductivity



M. Schirber,
PRL Syn. 2013



Quantum Spin Liquids



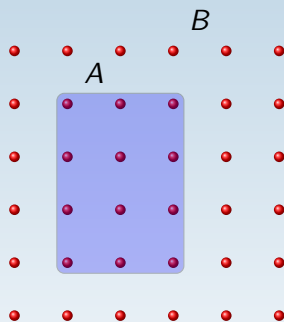
fractional QHE
fract. Chern insulators

Approximation Schemes

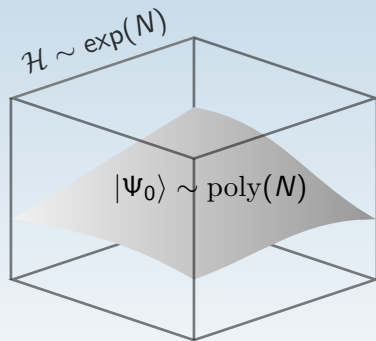
- Mean Field Theory / Hartree-Fock
- Quantum Monte Carlo
- Tensor Network States
- ...



The Area Law of Entanglement



$$S_{AB} := S(\rho_A) \propto \partial A$$

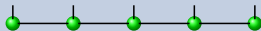


Tensor Network States

1D: Matrix Product States (MPS)

- efficient approximation of local Hamiltonians in 1D

F. Verstraete, and J. I. Cirac, Phys. Rev. B 73, 094423 (2006)



- classification of all phases in 1D

F. Pollmann, A. M. Turner, E. Berg, and M. Oshikawa, Phys. Rev. B 81, 064439 (2010)

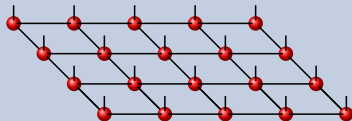
X. Chen, Z.-C. Gu, and X.-G. Wen, Phys. Rev. B 83, 035107 (2011)

N. Schuch, D. Pérez-García, and J. I. Cirac, Phys. Rev. B 84, 165139 (2011)

2D: Projected Entangled-Pair States (PEPS)

- presumably efficient approximation in 2D, proven for finite temperatures

M. B. Hastings, Phys. Rev. B 76, 035114 (2007)

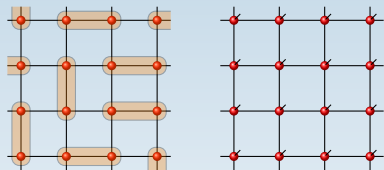


- no sign problem

Previous topological PEPS

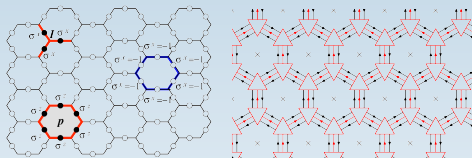
Resonating valence bond states

P. W. Anderson, Mater. Res. Bull. 8, 153 (1973)



String net models

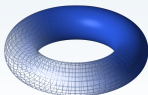
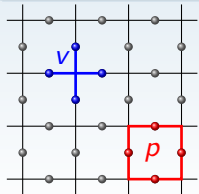
M. A. Levin, and X.-G. Wen, Phys. Rev. B 71, 045110 (2005)



O. Buerschaper, M. Aguado, and G. Vidal,
Phys. Rev. B 79, 085119 (2009)

Toric code

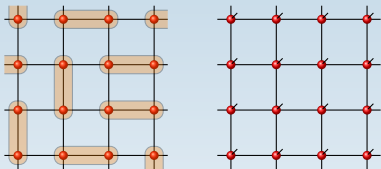
A. Kitaev, Ann. Phys. 303, 2 (2003)



Previous topological PEPS

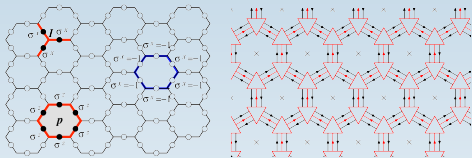
Resonating valence bond states

P. W. Anderson, Mater. Res. Bull. 8, 153 (1973)



String net models

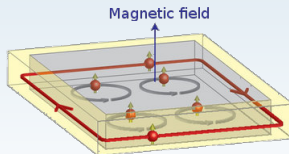
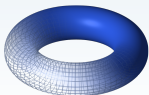
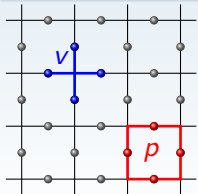
M. A. Levin, and X.-G. Wen, Phys. Rev. B 71, 045110 (2005)



O. Buerschaper, M. Aguado, and G. Vidal,
Phys. Rev. B 79, 085119 (2009)

Toric code

A. Kitaev, Ann. Phys. 303, 2 (2003)



D. Kong, and Y. Cui, Nature Chemistry 3, 845 (2011)

Table of content

1 Motivation

2 Construction of Tensor Network States

3 Free fermionic chiral PEPS

- T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 111, 236805 (2013)
- T. B. Wahl, S. T. Haßler, H.-H. Tu, J. I. Cirac, and N. Schuch, Phys. Rev. B 90, 115133 (2014)

4 Topologically ordered chiral PEPS

- S. Yang, T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 114, 106803 (2015)

Table of content

1 Motivation

2 Construction of Tensor Network States

3 Free fermionic chiral PEPS

- T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 111, 236805 (2013)
- T. B. Wahl, S. T. Haßler, H.-H. Tu, J. I. Cirac, and N. Schuch, Phys. Rev. B 90, 115133 (2014)

4 Topologically ordered chiral PEPS

- S. Yang, T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 114, 106803 (2015)

Tensor Network States - Overview

$$|\Psi\rangle = |\Psi(A_{v_1 v_2 \dots v_z}^i)\rangle$$

- in 1D: Matrix Product State (MPS): A_{lr}^i ($l, r = 1, \dots, D$)

$$|\Psi\rangle = \sum_{i_1, \dots, i_N} \underbrace{\text{tr} \left(A^{i_1} \dots A^{i_N} \right)}_{c_{i_1 \dots i_N}} |i_1 \dots i_N\rangle$$



- in 2D: Projected Entangled Pair States (PEPS): A_{lrud}^i

MPS construction

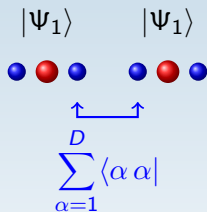
$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$

$$|\Psi_1\rangle$$



MPS construction

$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$



MPS construction

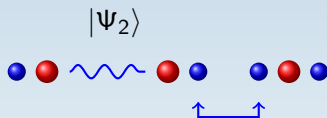
$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$

$$|\Psi_2\rangle$$



MPS construction

$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$



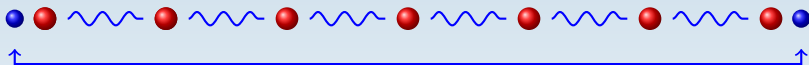
MPS construction

$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$



MPS construction

$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$



MPS construction

$$|\Psi_1\rangle = \sum_{i=1}^d \sum_{l,r=1}^D A^i_{lr} |i\rangle_{lr}$$



$$|\Psi_{\text{MPS}}\rangle = \sum_{i_1 \dots i_N} \text{tr} \left(A^{i_1} A^{i_2} \dots A^{i_N} \right) |i_1 i_2 \dots i_N\rangle$$

PEPS construction

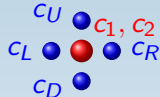
$$|\Psi_1\rangle =$$


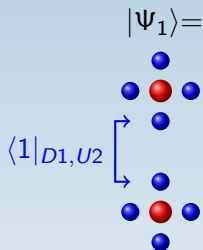
Diagram illustrating the construction of a PEPS state $|\Psi_1\rangle$. The state is represented by a central red circle (labeled c_1, c_2) connected to four blue circles (labeled c_U, c_L, c_D, c_R).

Majorana modes: $c^\dagger = c$

$$\frac{1}{2}(c_1 - ic_2) = a$$

$$\frac{1}{2}(c_1 + ic_2) = a^\dagger$$

PEPS construction



Majorana modes: $c^\dagger = c$

$$\frac{1}{2}(c_1 - ic_2) = a$$

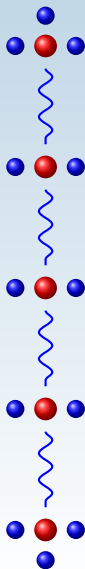
$$\frac{1}{2}(c_1 + ic_2) = a^\dagger$$

PEPS construction

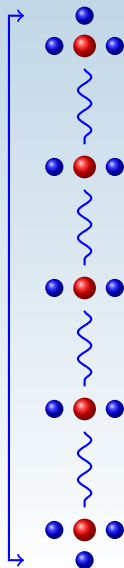
$$|\Psi_2\rangle =$$



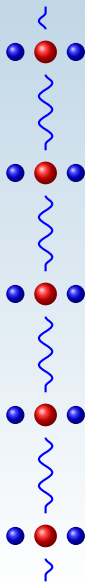
PEPS construction



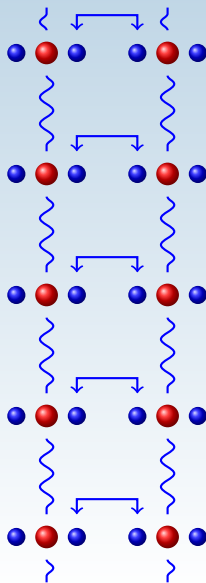
PEPS construction



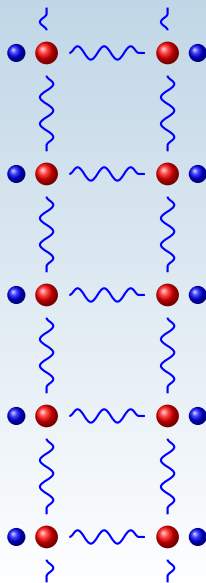
PEPS construction



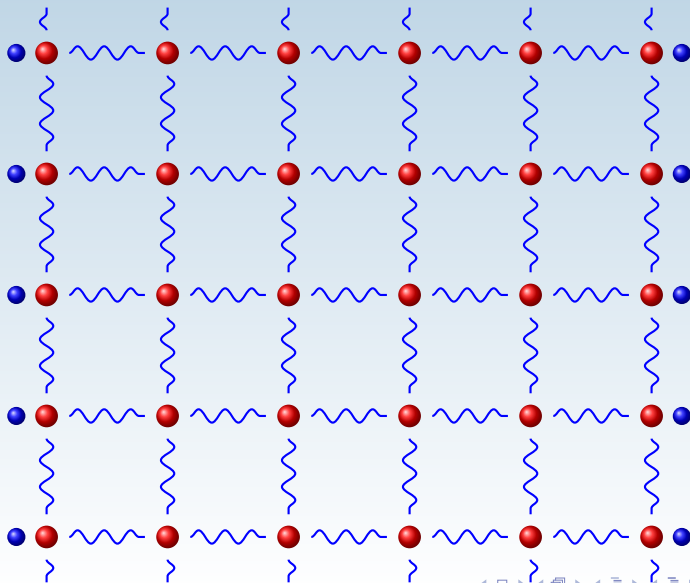
PEPS construction



PEPS construction



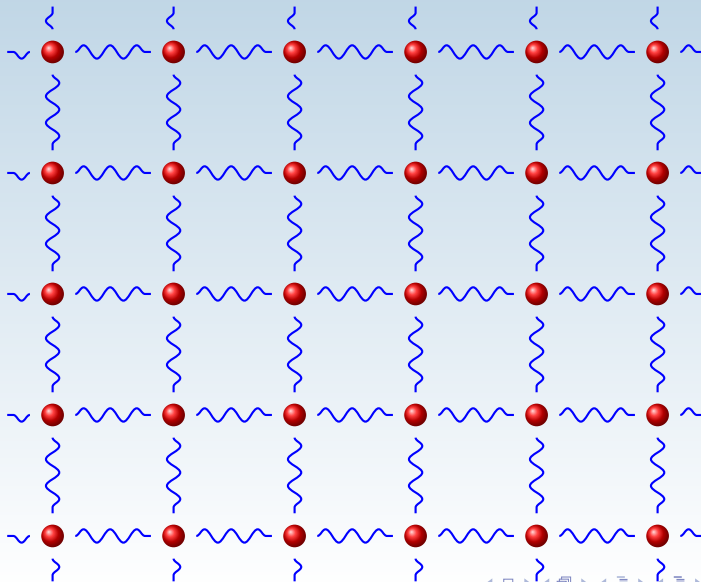
PEPS construction



PEPS construction

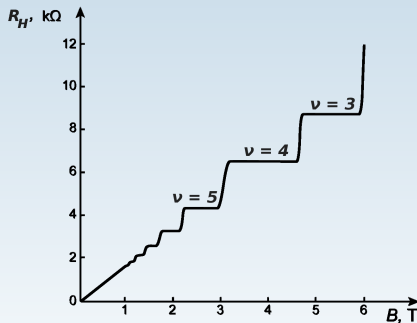
$$|\Psi_{\text{PEPS}}\rangle =$$

F. Verstraete and J. I. Cirac, Phys. Rev. A 70, 060302(R) (2004)



PEPS vs. Quantum Hall Effect

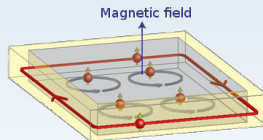
$$A_{lrud}^p \longrightarrow |\Psi_1\rangle = \begin{array}{c} c_U \text{ (blue)} \\ c_L \text{ (blue)} \text{ } c_1, c_2 \text{ (red)} \\ c_D \text{ (blue)} \end{array} \text{ } c_R \text{ (blue)}$$



$$R_H = \frac{h}{\nu e^2}$$

Chern number

$$\nu = C = 0, \pm 1, \pm 2, \dots$$



D. Kong, and Y. Cui, Nature Chemistry 3, 845 (2011)

Table of content

1 Motivation

2 Construction of Tensor Network States

3 Free fermionic chiral PEPS

- T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 111, 236805 (2013)
- T. B. Wahl, S. T. Haßler, H.-H. Tu, J. I. Cirac, and N. Schuch, Phys. Rev. B 90, 115133 (2014)

4 Topologically ordered chiral PEPS

- S. Yang, T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 114, 106803 (2015)

$$H = \sum_{k,j} T_{kj} a_k^\dagger a_j + \Delta_{kj} a_k^\dagger a_j^\dagger + \bar{\Delta}_{kj} a_j a_k$$

Free fermionic chiral PEPS

- all chiral PEPS with one Majorana mode
- general properties for more Majorana modes

The simplest chiral PEPS

$$|\Psi_1\rangle = \begin{array}{c} c_U \text{ (blue)} \\ c_L \text{ (blue)} \quad c_1, c_2 \text{ (red)} \quad c_R \text{ (blue)} \\ c_D \text{ (blue)} \end{array}$$

Parameterization of simplest chiral PEPS

$$|\Psi_1\rangle = (1 + a^\dagger b^\dagger) |\text{vac}\rangle$$

$$b = \alpha(c_L \pm ic_R) + \beta(c_U \pm ic_D)$$

Virtual Symmetry:

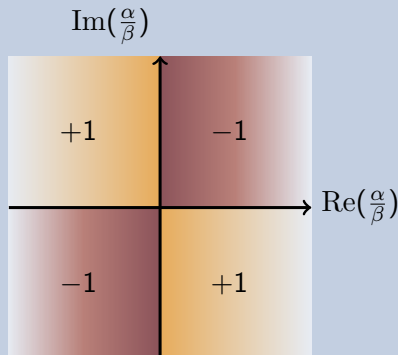
$$d_1 |\Psi_1\rangle = 0$$

$$d_1 = -\bar{\beta}(c_L \pm ic_R) + \bar{\alpha}(c_U \pm ic_D)$$

T. B. Wahl, H.-H. Tu, N. Schuch, J. I. Cirac, PRL '13

T. B. Wahl, S. T. Haßler, H.-H. Tu, J. I. Cirac, N. Schuch, PRB '14

see also: J. Dubail and N. Read, PRB '15



Local vs. global Symmetries

a)

$$\mathbf{d}_1 \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \end{array} = 0$$

b)

$$\mathbf{d}_2 \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \end{array} = 0$$

c)

$$\mathbf{d}_0 \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \end{array} = 0$$

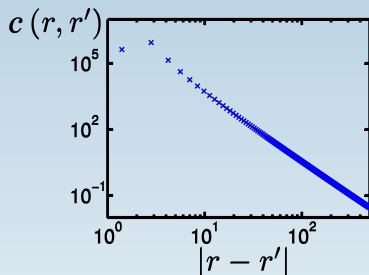
d)

$$\mathbf{d} \begin{array}{c} \bullet \\ \bullet \\ \bullet \\ \bullet \end{array} = 0$$

Free fermionic chiral PEPS

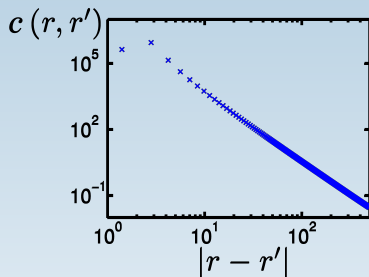
- general properties for more Majorana modes

General properties of chiral free fermionic PEPS



Rigorously shown in:
J. Dubail and N. Read,
PRB 92, 205307 (2015)

General properties of chiral free fermionic PEPS



Rigorously shown in:
J. Dubail and N. Read,
PRB 92, 205307 (2015)

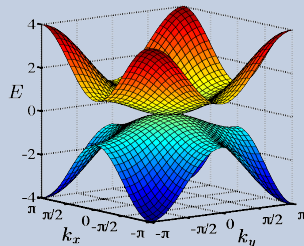
Flat band Hamiltonian

- long-range
- stable to perturbations

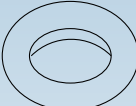
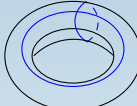
Frustration free Hamiltonian

$$\mathcal{H} = \sum_{\mathbf{r}} h_{\mathbf{r}}$$

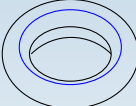
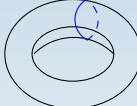
- local
- gapless



General properties of chiral free fermionic PEPS


 $= 0,$

 $= |\Phi_0\rangle$

insert string operators $\sum_x c_x$:


 $\propto$

 $= |\Phi_1\rangle$

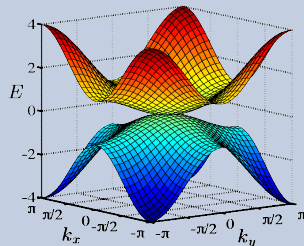
Flat band Hamiltonian

- long-range
- stable to perturbations

Frustration free Hamiltonian

$$\mathcal{H} = \sum_{\mathbf{r}} h_{\mathbf{r}}$$

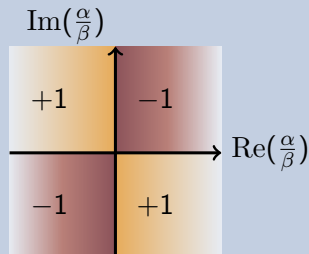
- local
- gapless



Summary: chiral free fermionic PEPS

PEPS with one Majorana mode

- chiral cases: $|\Psi_1(\alpha, \beta)\rangle$
- virtual symmetry $d_1|\Psi_1\rangle = 0$



any chiral free fermionic PEPS

- polynomially decaying correlations
- critical local Hamiltonians
- long range gapped topological Hamiltonians

Table of content

1 Motivation

2 Construction of Tensor Network States

3 Free fermionic chiral PEPS

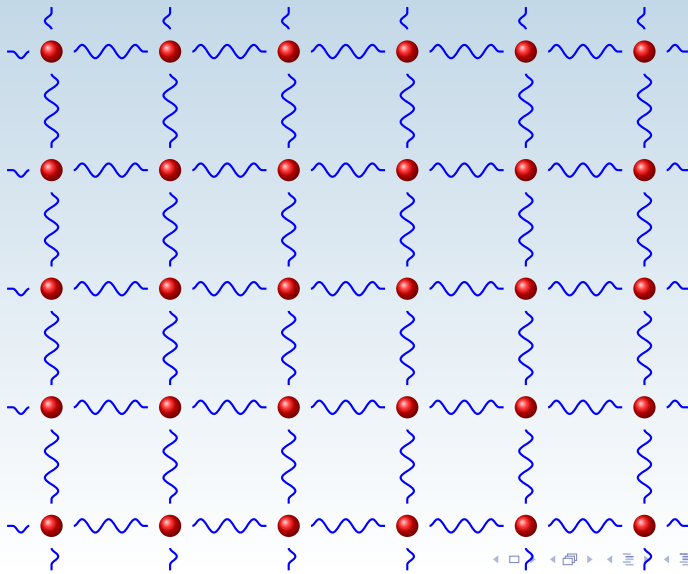
- T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 111, 236805 (2013)
- T. B. Wahl, S. T. Haßler, H.-H. Tu, J. I. Cirac, and N. Schuch, Phys. Rev. B 90, 115133 (2014)

4 Topologically ordered chiral PEPS

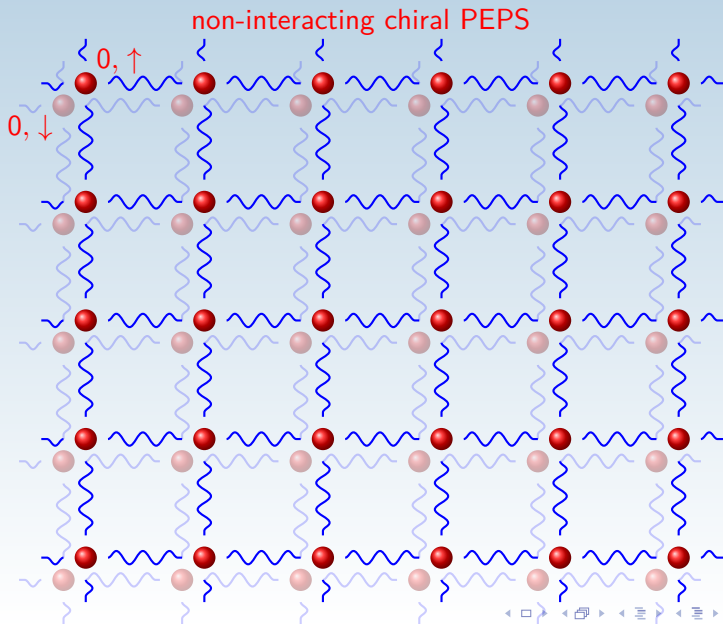
- S. Yang, T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, Phys. Rev. Lett. 114, 106803 (2015)

Gutzwiller projection

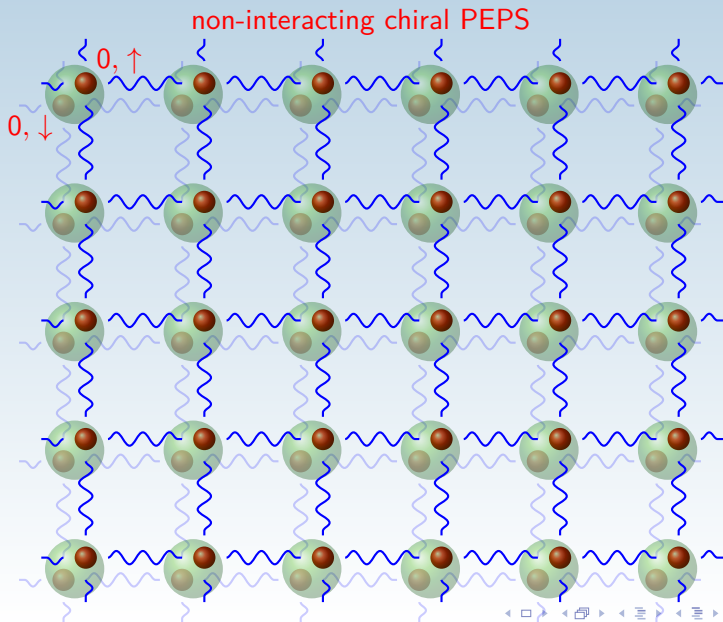
non-interacting chiral PEPS



Gutzwiller projection

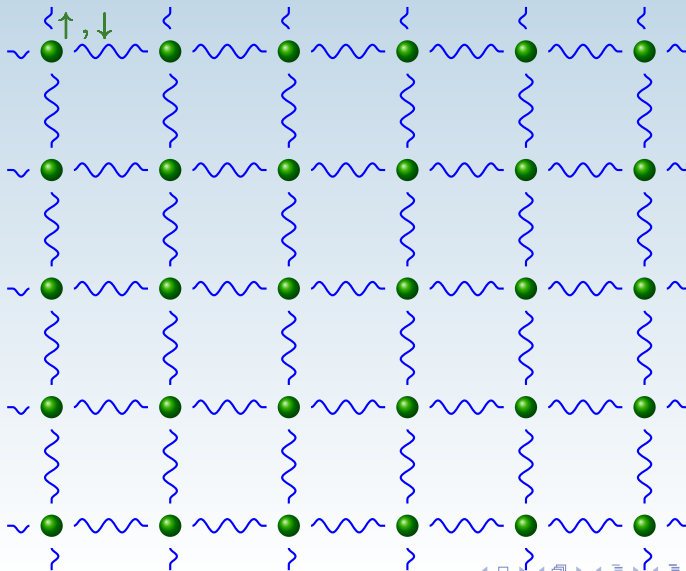


Gutzwiller projection



Gutzwiller projection

interacting chiral PEPS

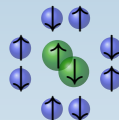


Topologically ordered chiral PEPS

$$|\Phi_1\rangle = \left(a_{\uparrow}^{\dagger} b_{\uparrow}^{\dagger} + a_{\downarrow}^{\dagger} b_{\downarrow}^{\dagger} \right) |\text{vac}\rangle$$

$$d_{\uparrow,\downarrow} |\Phi_1\rangle = 0$$

$$\sigma_L^z \sigma_R^z \sigma_U^z \sigma_D^z |\Phi_1\rangle = -|\Phi_1\rangle$$

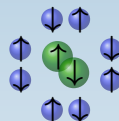


Topologically ordered chiral PEPS

$$|\Phi_1\rangle = \left(a_{\uparrow}^{\dagger} b_{\uparrow}^{\dagger} + a_{\downarrow}^{\dagger} b_{\downarrow}^{\dagger} \right) |\text{vac}\rangle$$

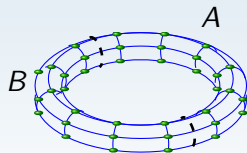
$$d_{\uparrow,\downarrow} |\Phi_1\rangle = 0$$

$$\sigma_L^Z \sigma_R^Z \sigma_U^Z \sigma_D^Z |\Phi_1\rangle = -|\Phi_1\rangle$$

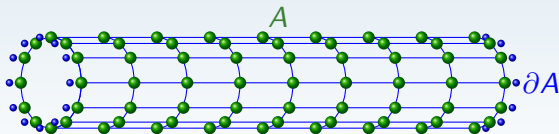


Properties:

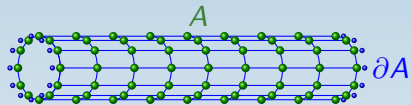
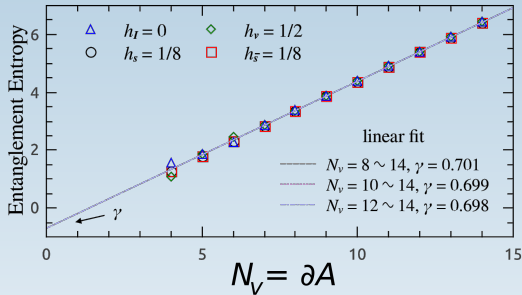
- algebraically decaying correlations
- $\mathcal{H} = \sum_r h_r$ has 5 ground states $\rightarrow$ **Topological degeneracy?**



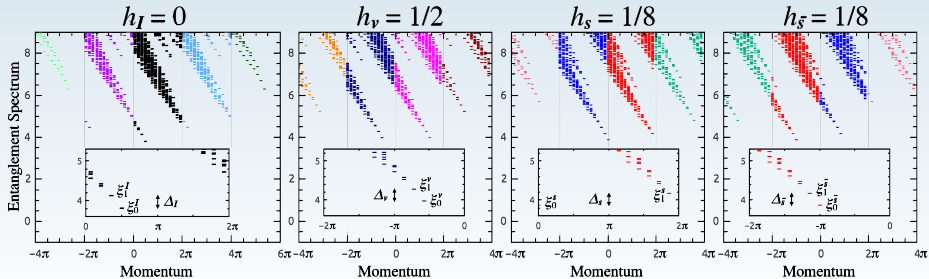
$$\rho_A = \text{tr}_B(|\Phi\rangle\langle\Phi|) \propto e^{-H_{\text{ent}}}$$



Area Law and Entanglement Spectrum



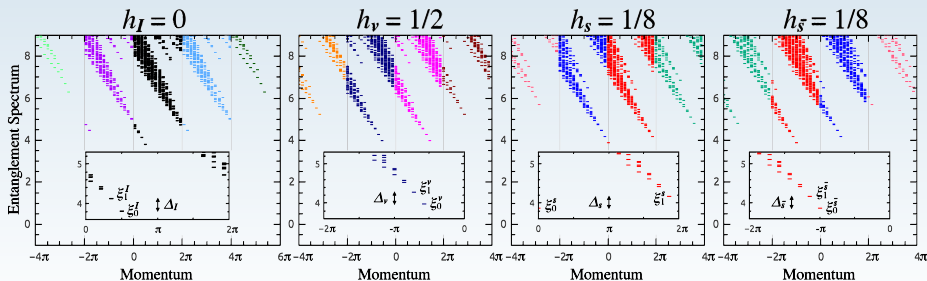
$$S(A) = c \partial A - \ln(2)$$



Area Law and Entanglement Spectrum

Properties:

- algebraically decaying correlations
- $\mathcal{H} = \sum_{\mathbf{r}} h_{\mathbf{r}}$ has 5 ground states $\rightarrow$ (4 topological, 1 gapless)
- $\text{SO}(2)_1$ CFT with $c = 1$
 topological correction to area law $S(A) = c \partial A - \gamma$; $\gamma = \ln(2)$
 entanglement spectra with the primary fields:



S. Yang, T. B. Wahl, H.-H. Tu, N. Schuch, and J. I. Cirac, PRL '15

Conclusions

Conclusions:

- PEPS can represent Quantum Hall systems
- non-interacting chiral PEPS have long range correlations
(gapless local Hamiltonians, long range flat band Hamiltonians)
- also true for example of interacting chiral PEPS

Conclusions

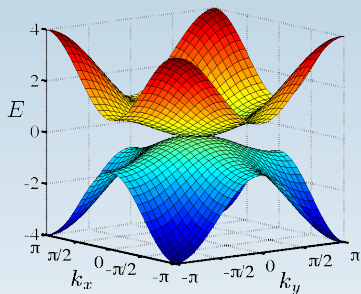
Conclusions:

- PEPS can represent Quantum Hall systems
- non-interacting chiral PEPS have long range correlations (gapless local Hamiltonians, long range flat band Hamiltonians)
- also true for example of interacting chiral PEPS

Outlook:

- disprove existence of short range chiral PEPS
- numerical simulations of chiral systems

Frustration free Hamiltonian



$$\mathcal{H} = \mathcal{H}_{\text{ff}} + \mu_0 \mathcal{H}_{\text{on-site}} + \nu_0 \mathcal{H}_{\text{hopping}}$$

insert string operators $\sum_x c_x$:

