



Cold atoms meet flux quanta & microwave cavities

József Fortágh



Cold atoms on the superconducting chip

Surface at 4K temperature

↑ separation: down to microns ↓

atomic cloud $N \sim 10^4 - 10^6$
 $T \sim$ down to 100nK

$p \sim 10^{-12}$ mbar

↓

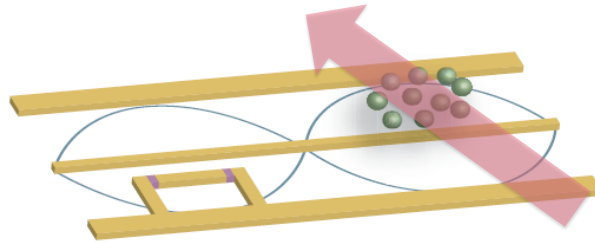
Measure and control interactions → Couple atoms to quantum electronic circuits

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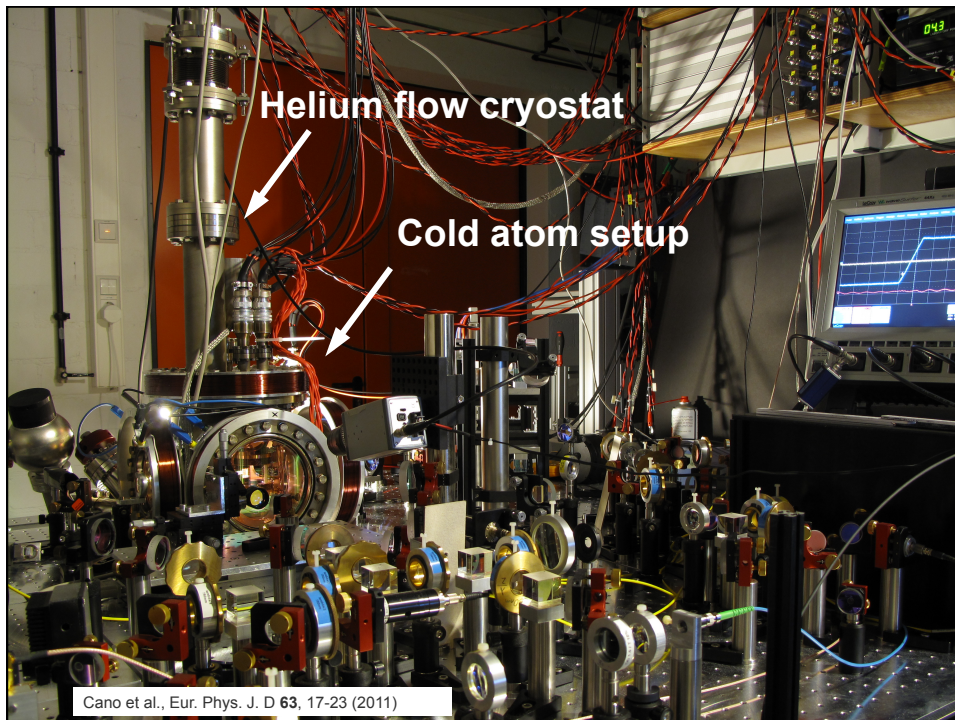
Vision: Quantum processor on a chip

Possible architecture: superconductors – microwaves – atoms – light

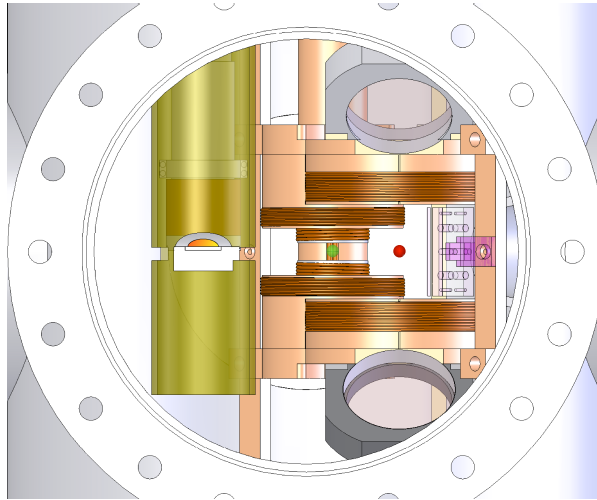


$$\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

Image: K. Tordrup and K. Molmer, PRA 77, 020301(R) (2008)

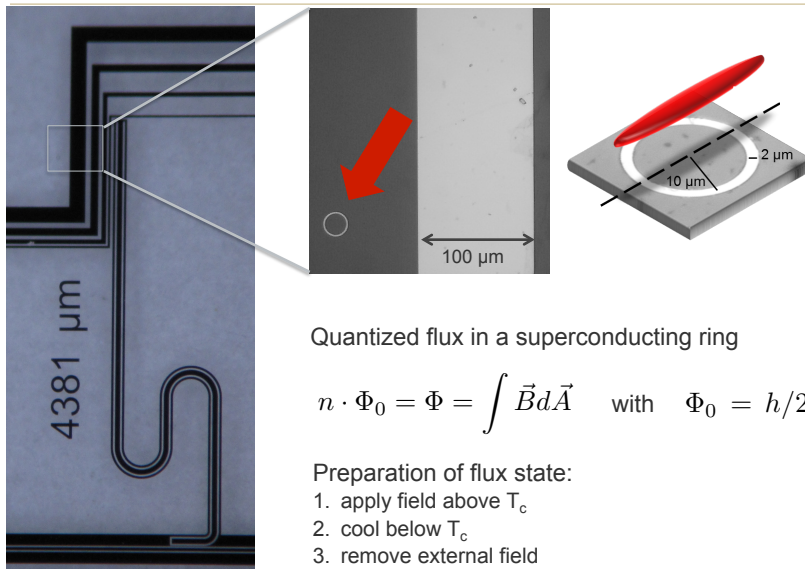


Magneto-optical trap, magnetic trap, optical tweezers



Cano et al., Eur. Phys. J. D **63**, 17-23 (2011)

Superconducting ring



Quantized flux in a superconducting ring

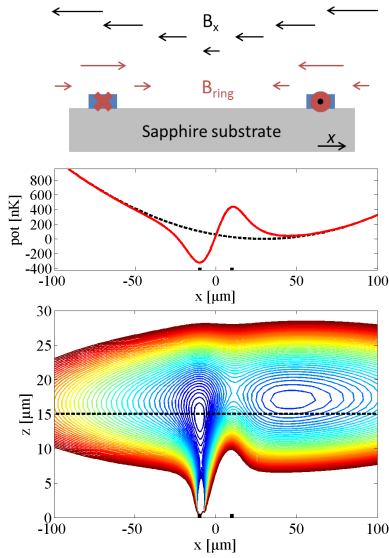
$$n \cdot \Phi_0 = \Phi = \int \vec{B} d\vec{A} \quad \text{with} \quad \Phi_0 = h/2e$$

Preparation of flux state:

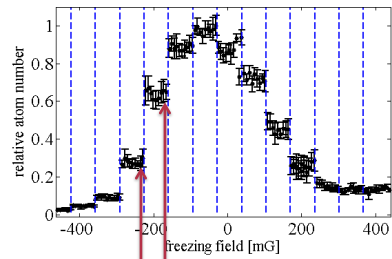
1. apply field above T_c
2. cool below T_c
3. remove external field

Weiss et al., PRL **114**, 113003 (2015)

Mapping the flux state of the ring to atomic clouds



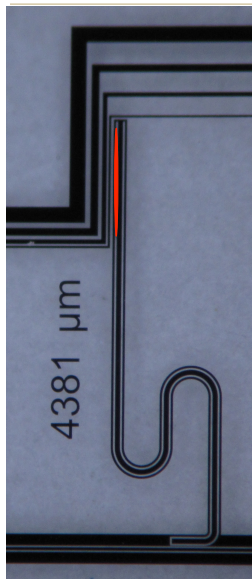
Total atom number after 1s storage



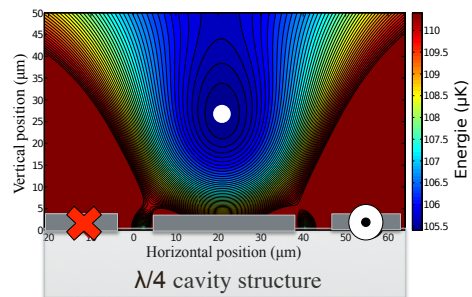
Trap depth varies in steps with the number of flux quanta

Weiss et al., PRL 114, 113003 (2015)

Trapping atoms on the SC chip



Persistent current trap

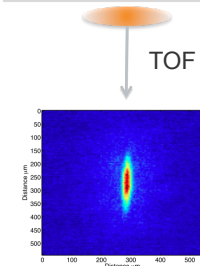
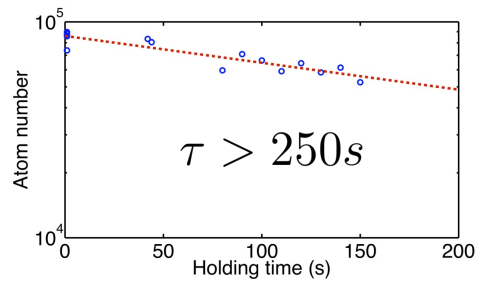
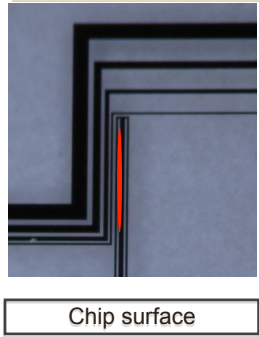


$\lambda/4$ cavity structure

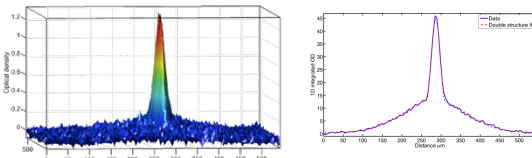
Bias Field

Bernon et al., Nat. Commun. 4, 2380 (2013)

Trapping atoms on the SC chip

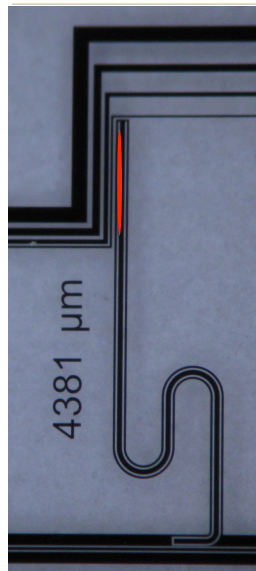


$N_C \approx 10^6$ in the F=1 or F=2 hyperfine state



Bernon et al., Nat. Commun. 4, 2380 (2013)

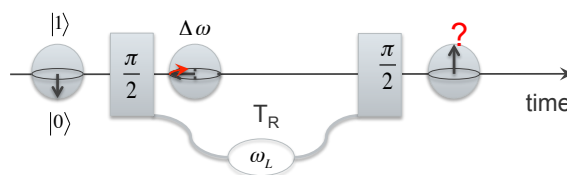
Atomic coherence at the superconducting coplanar cavity structure



$$|0\rangle \rightarrow \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

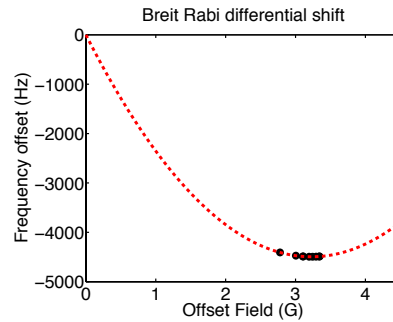
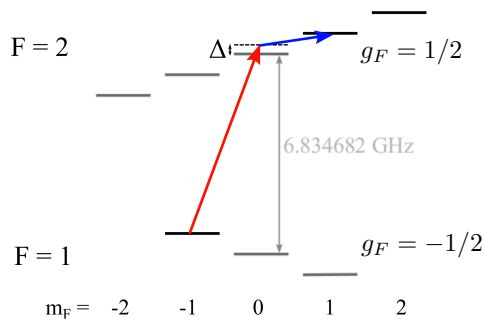
Time evolution

Coherence ?





Atomic coherence on a SC chip

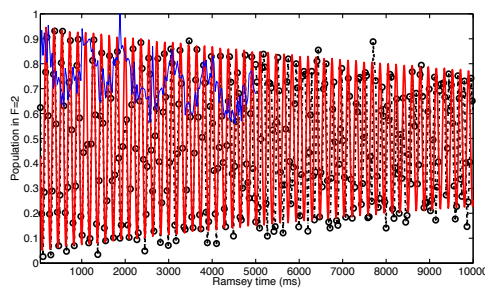


$$\Delta\omega(B) \cong \omega_f + \beta \cdot (B - B_0)^2$$

D.M. Harber *et al* Phys. Rev. A **66**, 053616 (2002)
P. Treutlein *et al* Phys. Rev. Lett. **92**, 203005 (2004)



Coherence and stability of operation



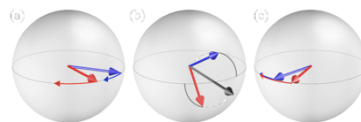
Coherence time $T_2 > 10$ s

Expected due to residual frequency inhomogeneity in the trap $T_2 \approx 6.5$ s

Identical spin rotation effect synchronizes the clock

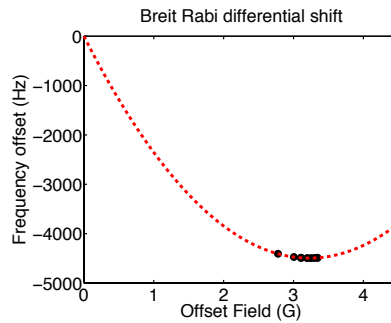
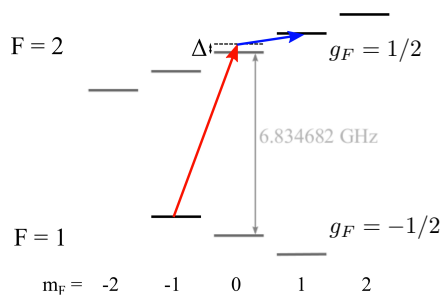
C. Deutsch *et al* Phys. Rev. Lett. **105**, 020401 (2010)

Bernon *et al.*, Nat. Commun. **4**, 2380 (2013)





Differential Zeeman shift: dephasing and phase shift

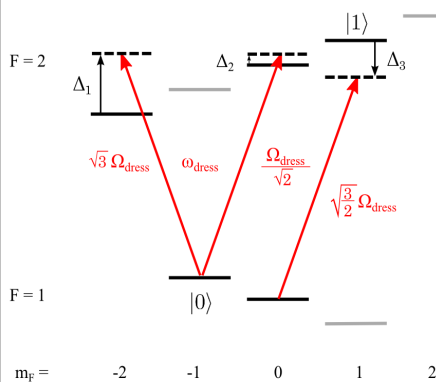


$$\Delta\omega(B) \cong \omega_f + \beta \cdot (B - B_0)^2$$

Suppression of the differential shift?



Dressing the clock transition



Dressing field with perpendicular polarisation to the magnetic field at the trap centre (quantisation axis).

$$\Delta\omega(B) \cong \omega_f + \beta \cdot (B - B_0)^2$$

$$+ \Omega_{dress}^2 \cdot \left(\frac{3}{\Delta_1(B)} + \frac{1/2}{\Delta_2(B)} - \frac{3/2}{\Delta_3(B)} \right)$$

$$\frac{d\Delta\omega}{dB}, \frac{d^2\Delta\omega}{dB^2} = 0 \Rightarrow \Omega_{dress}, \omega_{dress}$$

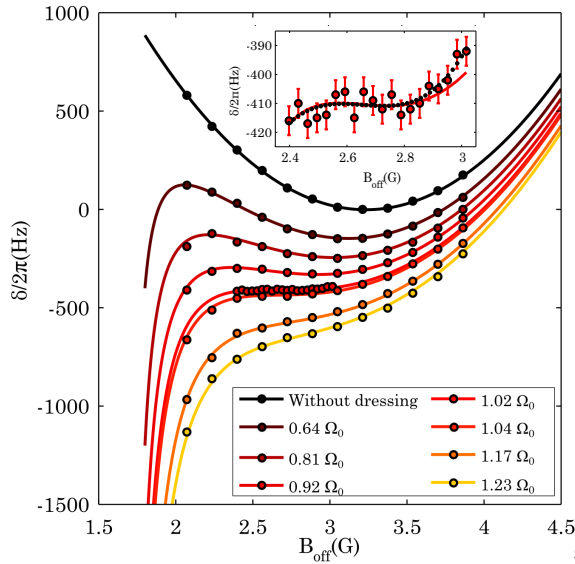
$$\Omega_{dress} = \Omega_{dress}(B_{center}) !$$

$$\omega_{dress} = \omega_{dress}(B_{center}) !$$

For any offset field the **differential Zeeman shift can be suppressed up to 2nd order.**



MW-control of the differential shift



$$\Delta_{\text{dress}} = -2\pi \times 1.19 \text{ MHz}$$

$$\Omega_0 = 2\pi \times 20.1 \text{ kHz}$$

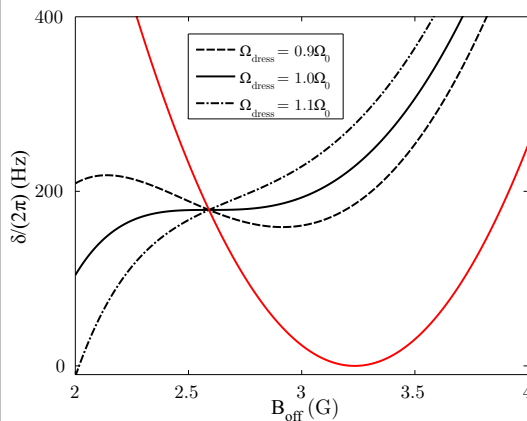
For any offset field the differential shift can be suppressed up to 2nd order.

Sárkány et al., PRA **90**, 053416 (2014)



Doubly protected clock states

For certain offset fields the differential shift becomes independent also from the microwave power.

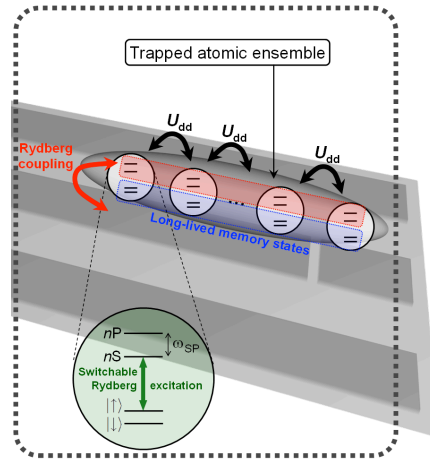


These “double magic points” are the preferable working points.

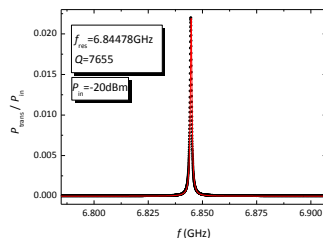
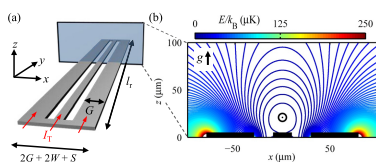
Sárkány et al., PRA **90**, 053416 (2014)

Intermediate status

- Rubidium atomic clouds (BEC) at a superconducting chip
- Coherence of hyperfine superposition states: seconds
- Noise-protected memory states
- Atom-cavity coupling?



Fabricated superconducting cavity



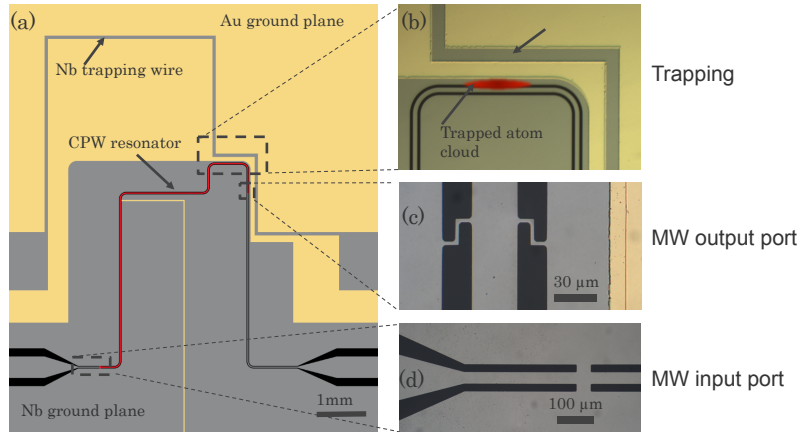
Inductively coupled co-planar superconducting cavity [1]

- Niobium on sapphire
- Fundamental frequency 6.8 GHz
- Measured $Q \sim 10^4$ for the fundamental mode at 4K
- Higher modes $Q \sim 10^3$

With integrated traps

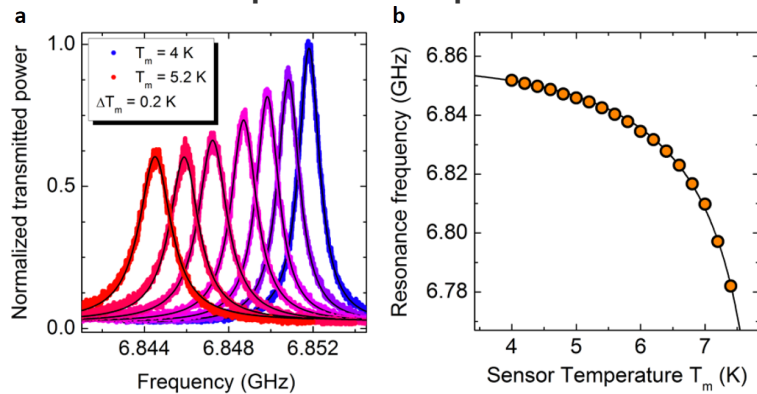
[1] Bothner et al., New J. Phys. 15, 093024 (2013)
In-house collaboration with the group of Reinhold Kleiner & Dieter Koelle.

Fabricated atom chip



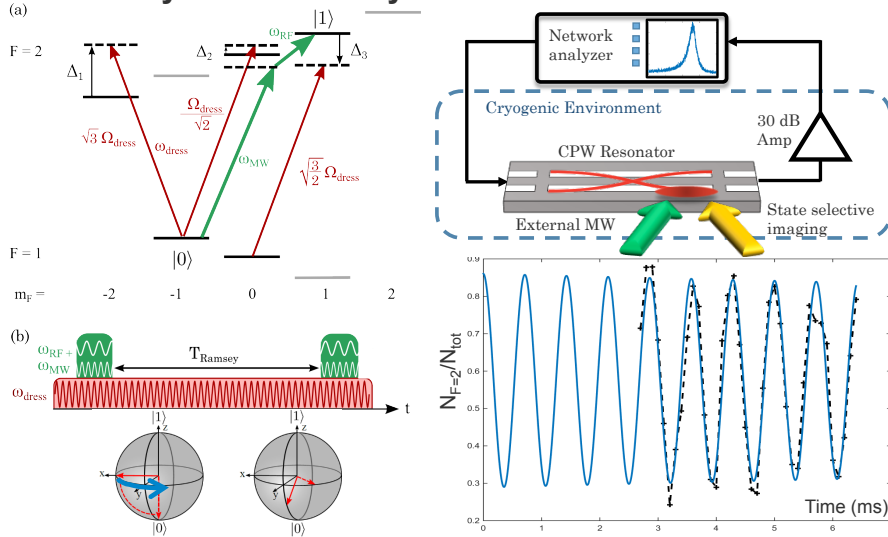
Ground planes are partially made of gold in order to avoid the Meissner effect

Resonance frequency vs. temperature



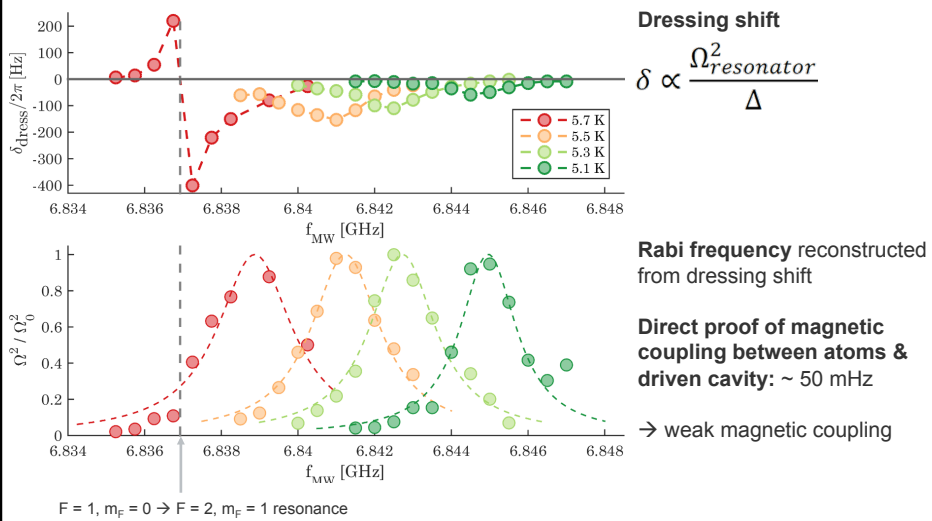
- Fundamental frequency tunable with temperature over 30 MHz
- Quality factor $Q \sim 10^4$
- Cavity resonance matches HF-Splitting at $T \sim 6\text{K}$

Ramsey interferometry in the resonator

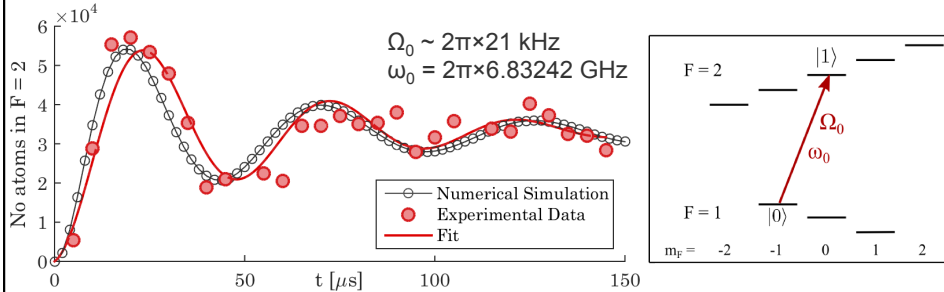


The Ramsey frequency varies with the dressing power and frequency.

Off-resonant dressing by the cavity field



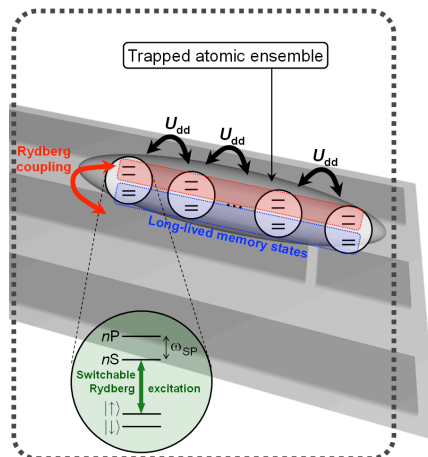
Resonant Rabi oscillation in a driven cavity



- Dephasing due to inhomogeneity of cavity field at the position of the atomic cloud
- Coupling is extremely weak (calculated $g \sim 0.05 \text{ Hz}$ for single atom and single photon) \rightarrow Strong far of reach

Intermediate status

- Rubidium atomic clouds (BEC) at a superconducting chip
- Coherence of hyperfine superposition states: seconds
- Noise-protected memory states
- Hyperfine ground states couple to the magnetic field of the cavity: weak (50 mHz)



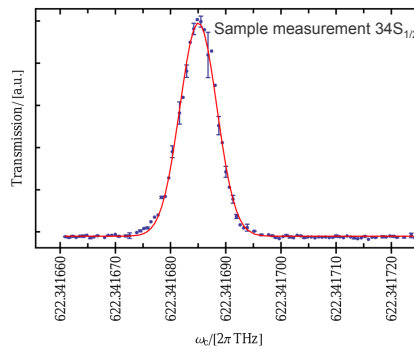
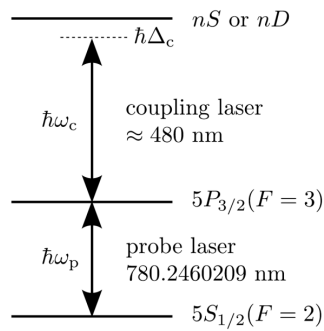
Coupling Rydberg states?

Expected coupling: strong (5 MHz)



Spectroscopy of ^{87}Rb Rydberg states

Electromagnetically induced transparency (EIT)

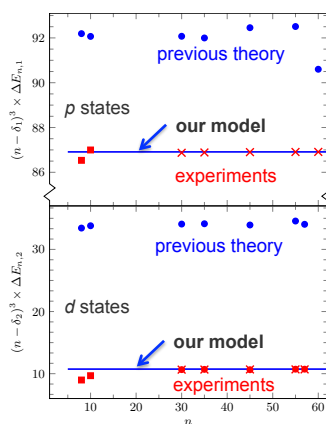


Lasers locked to frequency comb & wavemeter.

Measurement of quantum defects and ground state ionization energy of ^{87}Rb
(≤ 1 MHz abs. accuracy)
Mack et al., Phys. Rev. A **83**, 052515 (2011)



Core potential



Previous models:
Marinescu et al., PRA **49**, 982 (1994)
Pawlak et al., PRA **89**, 042506 (2014)

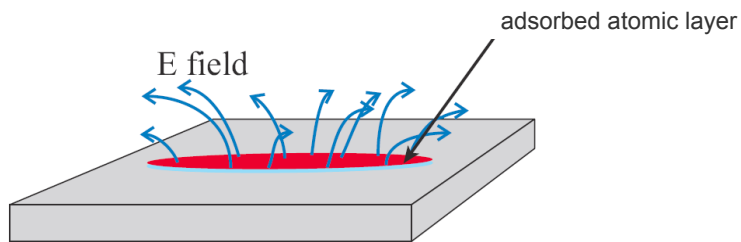
Our model:
Sanayei et al., PRA **91**, 032509 (2015)

Fine splitting of ^{87}Rb Rydberg
 p - and d -states



State of research

Polarized adatoms at the surface produce inhomogeneous electric fields



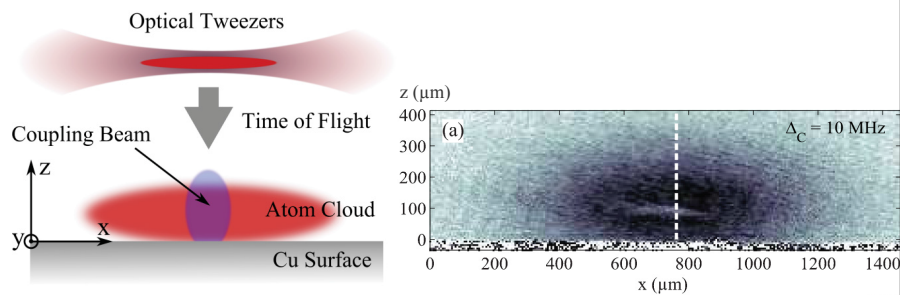
McGuirk et al., Phys. Rev. A **69**, 62905 (2004)
 Tauschinsky Phys. Rev. A **81**, 063411 (2010)
 Hattermann et al, Phys. Rev. A **86**, 022511 (2012)
 Carter et al., Phys. Rev. A **86**, 053401 (2012)
 Chan et al., Phys. Rev. Lett. **112**, 026101 (2014)
 Sedlacek et al., Phys. Rev. Lett. **116**, 133201 (2016)
 ...

Solution:

Selection of Rydberg state pairs
 - to match the cavity resonance
 - to suppress differential shift



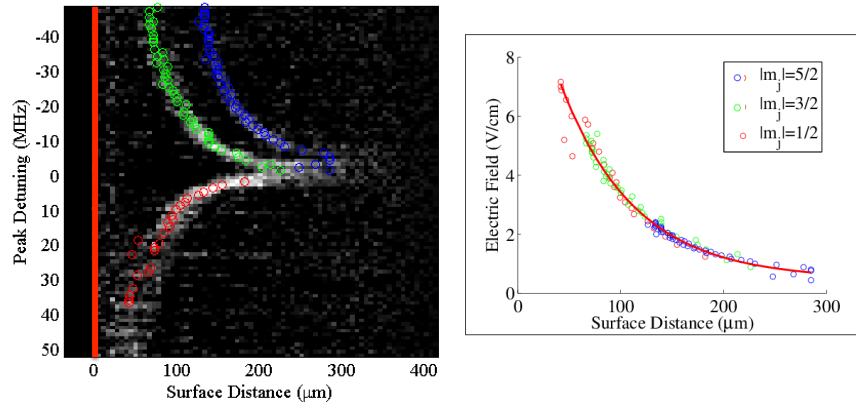
Stark shift near a dipole layer (copper surface with Rb adsorbates)



Energy shift measured with electromagnetically induced transparency (EIT)

Amsterdam: Tauschinsky Phys. Rev. A **81**, 063411 (2010)
Tübingen: Hattermann et al, Phys. Rev. A **86**, 022511 (2012)

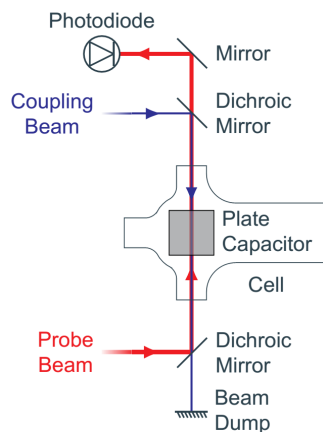
Rubidium Rydberg level shifts near a copper surface



Coupling laser detuning to the $5P_{3/2} \rightarrow 35D_{5/2}$ transition of ^{87}Rb
 We observe the stark shift of the $|m_j|=1/2$, $|m_j|=3/2$, and $|m_j|=5/2$ states

Hattermann et al, Phys. Rev. A **86**, 022511 (2012)

Stark-map in a cell



Setup

- Room-temp. vapor cell with electrodes
- EIT spectroscopy on ^{87}Rb , $n=35$ to 70
- Lasers locked to frequency comb & wavemeter

Measured Stark maps

- $E = 0\text{-}500$ V/cm, 3mV/cm field steps
- 2 MHz optical resolution

Improved numerical calculation

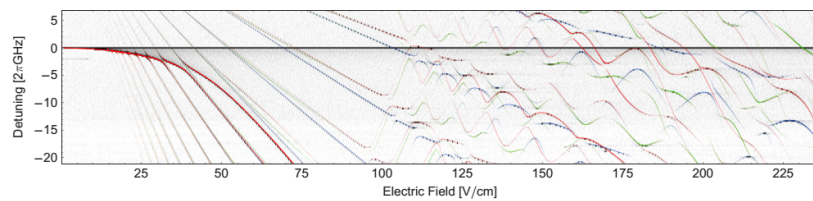
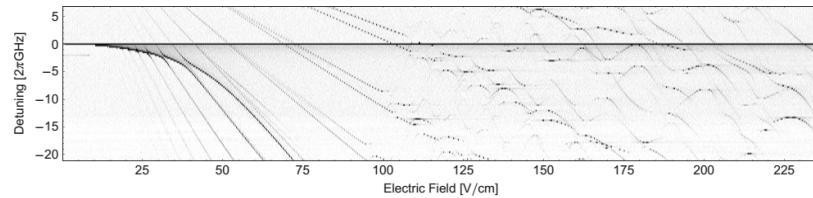
based on Zimmerman PRA **20**, 2251 (1979)

incl. transition strengths:

- Grimm et al., NJP **17**, 053005 (2015)



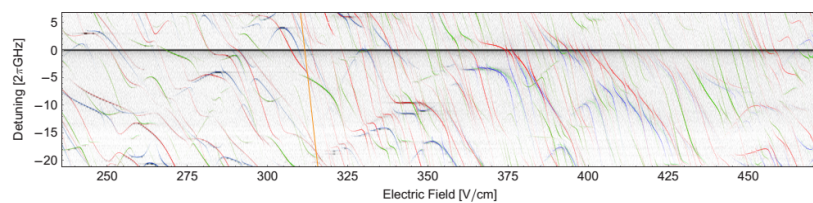
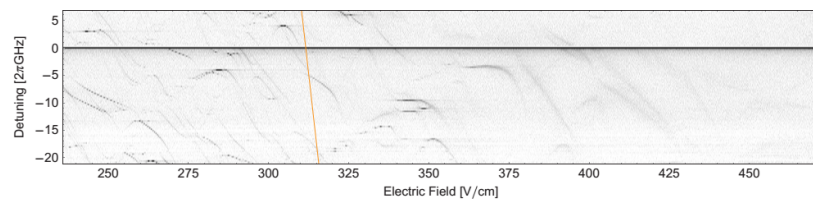
$35\text{S}_{1/2}$ Stark Map – EIT measurement (greyscale) and results from numerical calculations (colors)



Grimmel et al., NJP 17, 053005 (2015)



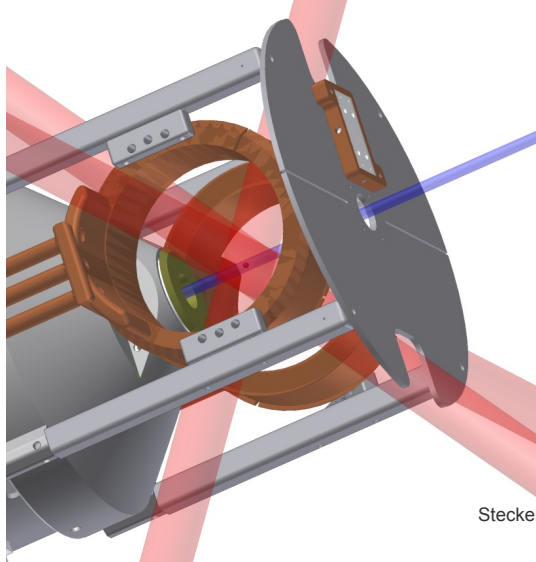
$35\text{S}_{1/2}$ Stark Map – EIT measurement (greyscale) and results from numerical calculations (colors)



Grimmel et al., NJP 17, 053005 (2015)



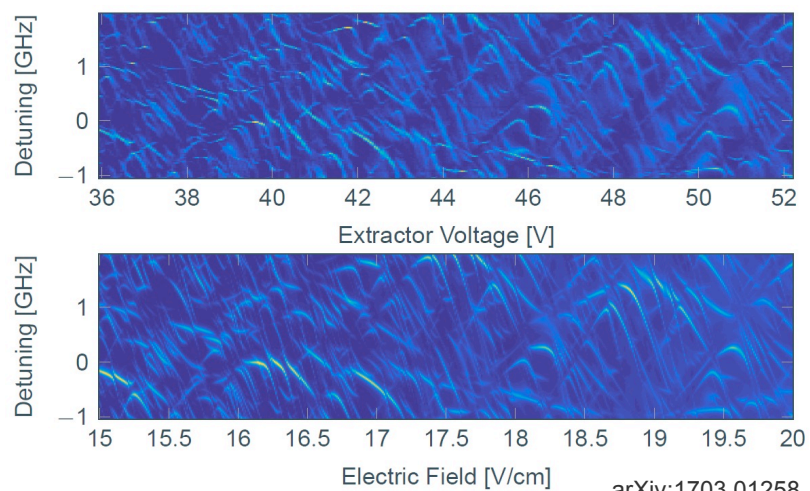
Ionization Stark spectra from a MOT



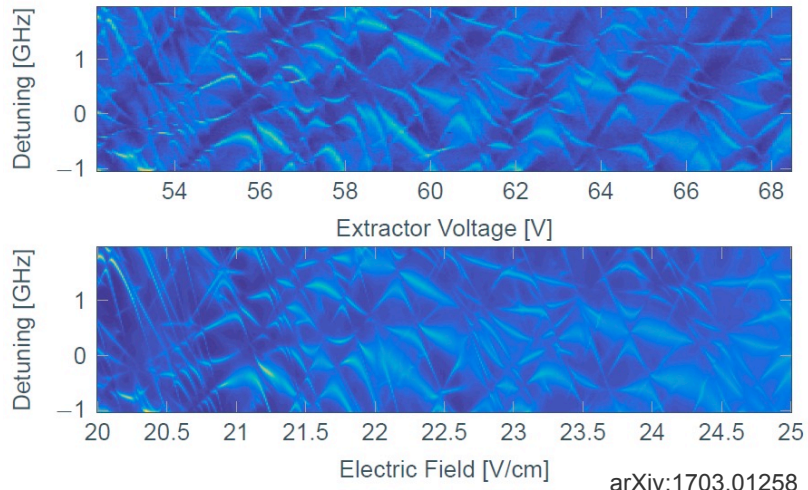
Stecker et al., NJP 19, 043020 (2017)



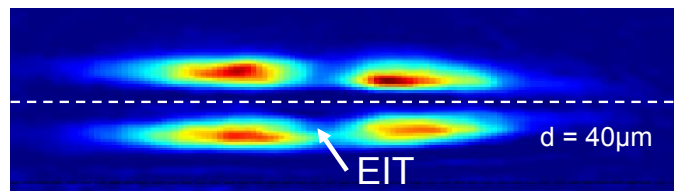
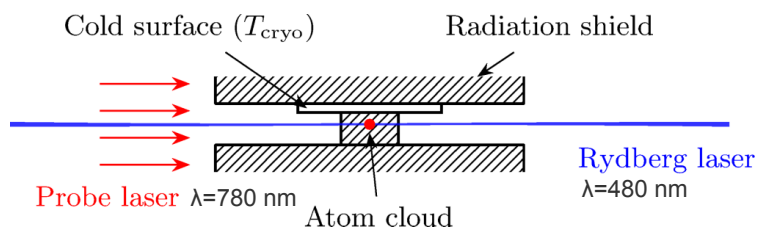
Ionization spectra near $70S_{1/2}$ from a MOT



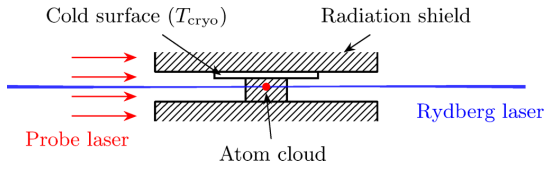
Ionization spectra near $70S_{1/2}$ from a MOT



Rydberg EIT on the SC chip

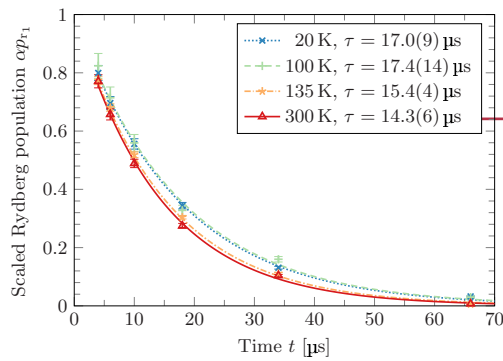


Lifetime of the $30S_{1/2}$ Rydberg state



All optical detection of Rydberg populations & coherences

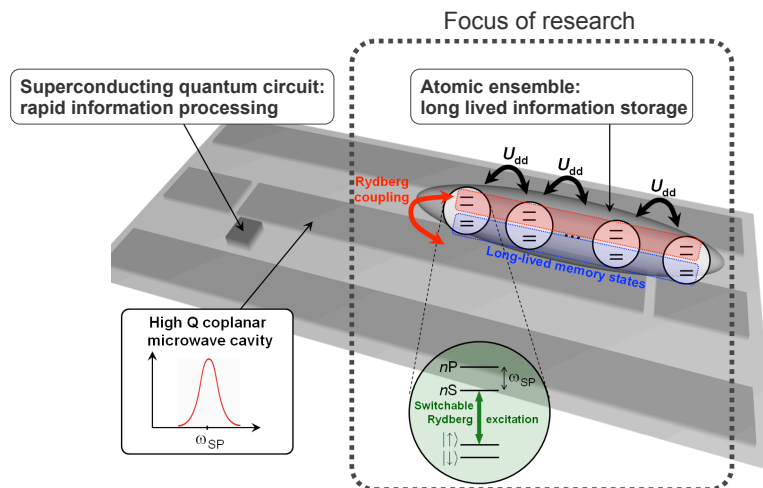
PRA 91, 0434422 (2015)



Previous literature value for $T = 300$ K
 $\tau = (14.5 \pm 1.2) \mu\text{s}$
Nascimento et al.,
PRA 74, 054501 (2006)

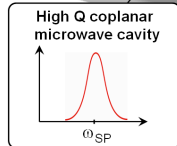
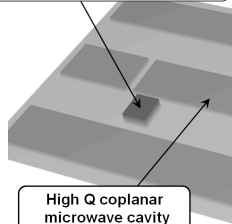
This lifetime data:
room T & cryogenic environment
PRA 92, 012517 (2015)

Outlook



Outlook

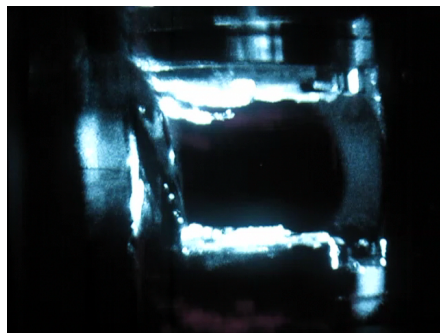
Superconducting quantum circuit:
rapid information processing



**Coherence of solid-state-qubits
is short at 4K (<math><1\mu\text{s}</math>)**

→ **50mK cryostat** ($\sim 100\mu\text{s}$)
MOT & magnetic traps achieved
in the 4K stage of the cryostat

**MOT & magnetic traps achieved
in the 4K stage of dilution fridge**



PT1 plate
($\approx 50\text{--}65\text{K}$)

PT2 plate
($\approx 4\text{--}6\text{K}$)

Still plate
($\approx 1.5\text{K}$)

100mK plate
($\approx 130\text{mK}$)

MC plate
($\approx 25\text{--}40\text{mK}$)

Zeeman
slower

Add. 4K plate
($\approx 4\text{--}6\text{K}$)

$^3\text{He}/^4\text{He}$
dilution fridge

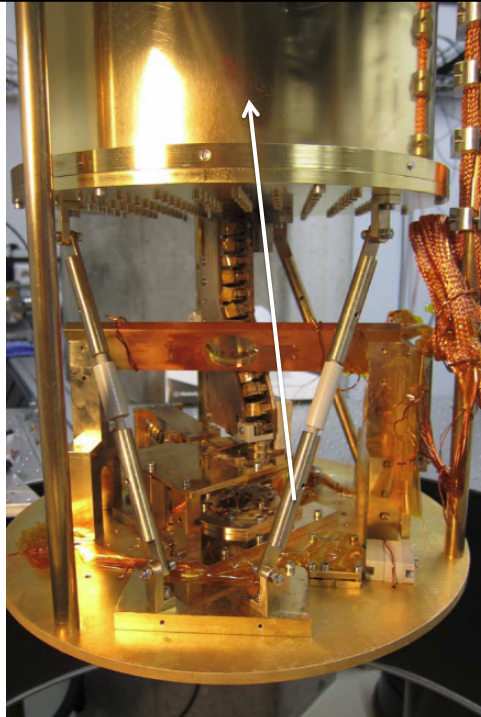
MOT

25 mK stage



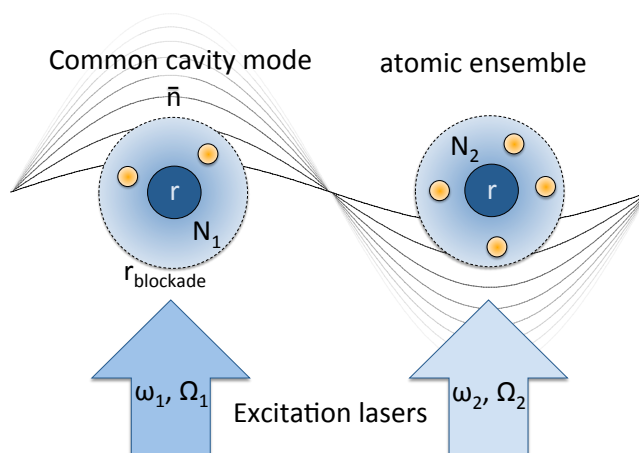
magnetic
transport

4K stage



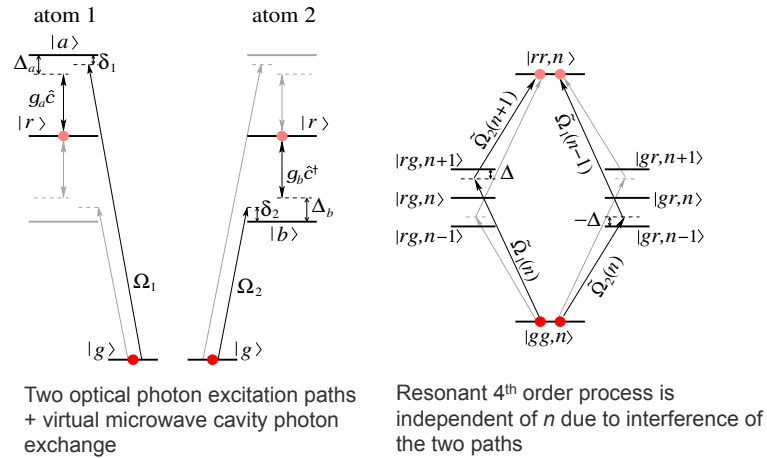
Jessen et al., Appl. Phys. B 116, 665-671 (2014)

Long distance quantum gate in a thermal cavity ?



Sárkány, Fortágh, Petrosyan, PRA 90 (R), 030303 (2015)

Sørensen-Mølmer scheme: Conditional excitation of a pair of Rydberg atoms in a cavity with n photons



Sárkány, Fortágh, Petrosyan, PRA **90** (R), 030303 (2015)

Contributors to this work

Cold-atoms & superconductors

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Helge Hattermann
Lőrinc Sárkány
Florian Jessen

Patrizia Weiß, Simon Bell, Simon Bernon,
Solid state research group of
Reinhold Kleiner & Dieter Koelle

Rydberg atoms

Jens Grimmel
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Markus Stecker
Andreas Günther
David Petrosyan

Markus Mack, Florian Karlewski
Nóra Sándor, Group of Nils Schopohl

