# Data Analysis IV Data Analysis for PTA

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#### Structure of the lectures

- > DA I: Statistic basis for DA: Likelihood, frequentist/Bayesian
- **GW** Obs I: History, response to GW
- **GW Obs II: LISA: LISAPathfinder, noises, ...**
- DA II: 3 main classes of signal, parameter estimations,
   Fisher Matrix
- **DA III: LISA DA: Global analaysis, MBHB, stochastic, ...**
- GW Obs III: LIGO
- GW Obs IV: PTA
- **DA IV: PTA data analysis**



## Overview

- PTA data
- Fitting the model of the pulsar
- Continuous Gravitational Wave :
  - Frequentist
  - Bayesian
  - Upper limit
- Stochastic background
- Global analysis



## Pulsar Timing Array Data



#### Pulsar model

Model of the observed arrival time of the pulsar and radio wave propagation:





## PTA data



#### Residuals

- ► Time Of Arrivals (TOAs) are used to "fit" a model
- ► TOAs "best model" = residuals => input to GW analysis



#### Residuals

#### Error in position: annual effect



MJD

#### Residuals



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## Model imperfection

- The fitting of the model is not perfect in particular because we did not consider GWs as the same time as other "pulsars" parameters
- We need to integrate these imprecisions in GWs data analysis.
  - Global analysis ... many parameters
  - Analytical marginalization based on a [van Haasteren et al. 2009, 2012] (see later)



#### Data analysis

• The general likelihood is:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^n det(C)}} \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T C^{-1}\left(\vec{\delta t} - \vec{h}\right)\right)$$

- n : number of data points
- *C* : correlation matrix
- $\vec{\delta t}$  : residual = data points
- $\vec{\theta}$  : parameters
- Unequally sample data



#### Including pulsar model errors [van Haasteren et al. 2009, 2012]

- During the fitting of the pulsar model some residual errors
- Assumption:
  - random Gaussian process  $\vec{\delta t_i}^G$
  - + some contamination by several systematic signals with known functional forms  $f_p(t_i)$  but a-priori unknown amplitudes  $\xi_p$ :

$$\delta \vec{t}_i = \delta \vec{t}_i^G + \sum_p \xi_p f_p(t_i)$$
$$= \delta \vec{t}_i^G + M \vec{\xi_p}$$

M : "design matrix" : n x m : m number of pulsar fitting parameters



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#### Data analysis [van Haasteren et al. 2009, 2012]

#### Then the likelihood can be rewritten as

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m}det(G^T C G)}} \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(G^T C G)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

- n : number of data points
- *C* : correlation matrix
- $\vec{\delta t}$  : residual = data points
- $\vec{\theta}$  : parameters
- G derived from M:
  - $M = U \Sigma V$ :  $(n \times n) (n \times m) (m \times m)$
  - $U = (F G) : ((n \times m) (n \times (n-m)))$



$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m}det(G^T C G)}}$$
$$\times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(G^T C G)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$



$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(G^T C G)}}$$
$$\times \exp\left(-\frac{1}{2} \left(\vec{\delta t} - \vec{h}\right)^T G(G^T C G)^{-1} G^T \left(\vec{\delta t} - \vec{h}\right)\right)$$

GW in



$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m}det(G^T C G)}}$$
$$\times \exp\left(-\frac{1}{2}\left(\vec{\delta t} + \vec{h}\right)^T G(G^T C G)^{-1} G^T\left(\vec{\delta t} + \vec{h}\right)\right)$$

#### GW in

• h : deterministic GW like binaries



$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(G^TCG)}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t} + \vec{h}\right)^T G(G^TCG)^{-1}G^T\left(\vec{\delta t} + \vec{h}\right)\right)$$

#### GW in

- h : deterministic GW like binaries
- C : stochastic background as correlation between pulsars



$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} \det(G(C\theta))}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t}+\vec{h}\right)^T G(G(C\theta))^{-1} G^T\left(\vec{\delta t}+\vec{h}\right)\right)$$

#### GW in

- h : deterministic GW like binaries
- C : stochastic background as correlation between pulsars
- Noises in C



#### Gravitational wave signal

#### • GW signal in the pulsar residual:



- n : direction of the pulsar
- *L* : distance Earth pulsar
- k : direction of the GW propagation
- $h_{ij}$  : GW strain



#### Continous GWs



#### Sources

- SuperMassive Black Hole Binaries (SMBHB)
  - $10^7$  to  $10^9$  solar masses
  - far from the merger =>
    - quasi-monochromatic sources
    - spins can be neglected

Continuous waves: sources that can be resolved individually



#### **Distribution of sources**



► GW waveform:  $h_+(t) = \mathcal{A} \left(1 + \cos^2 i\right) \cos \left(\phi(t) + \phi_0\right)$  $h_\times(t) = -2\mathcal{A} \cos i \sin \left(\phi(t) + \phi_0\right)$ 

• with 
$$\mathcal{A} = 2 \frac{\mathcal{M}_c^{5/3}}{D_L} (\pi f)^{2/3}$$

- Parameters:
  - $\mathcal{M}_c$  : chirp mass
  - $D_L$  : luminosity distance
  - $f=2\ \pi\ \omega$  : frequency of GW
  - *i* : inclination
  - $\phi(t)$  : phase
  - $\phi_0$  : initial phase



GW signal in the residual: 2 terms pulsar term & Earth term

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e(I)

$$r_{a}(t) = r_{a}^{2}(t) - r_{a}(t)$$

$$e_{a}^{e}(t) = \frac{\mathcal{A}}{\omega} \left\{ (1 + \cos^{2} \iota) F_{a}^{+} \left[ \sin(\omega t + \Phi_{0}) - \sin \Phi_{0} \right] + 2\cos \iota F_{a}^{\times} \left[ \cos(\omega t + \Phi_{0}) - \cos \Phi_{0} \right] \right\},$$

$$r_{a}^{p}(t) = \frac{\mathcal{A}_{a}}{\omega_{a}} \left\{ (1 + \cos^{2} \iota) F_{a}^{+} \left[ \sin(\omega_{a} t + \Phi_{a} + \Phi_{0}) - \sin(\Phi_{a} + \Phi_{0}) \right] + 2\cos \iota F_{a}^{\times} \left[ \cos(\omega_{a} t + \Phi_{a} + \Phi_{0}) - \cos(\Phi_{a} + \Phi_{0}) \right] \right\}.$$
Parameters

$$F_{a}^{+} = \frac{1}{2} \frac{(\hat{n}^{a} \cdot \hat{p})^{2} - (\hat{n}^{a} \cdot \hat{q})^{2}}{1 + \hat{n}^{a} \cdot \hat{k}}$$
$$F_{a}^{\times} = \frac{(\hat{n}^{a} \cdot \hat{p})(\hat{n}^{a} \cdot \hat{q})}{1 + \hat{n}^{a} \cdot \hat{k}}$$

- $\phi_0$ : initial phase at Earth
- $\phi_a$ : initial phase at pulsar
- $F_a^+$ ,  $F_a^{\times}$ : beam patterns depending on direction & polarisation



- ► Beam patterns  $F_{a}^{+}$ ,  $F_{a}^{\times}$  are:  $F_{a}^{+} = \frac{1}{2} \frac{(\hat{n}^{a} \cdot \hat{p})^{2} (\hat{n}^{a} \cdot \hat{q})^{2}}{1 + \hat{n}^{a} \cdot \hat{k}}$  $F_{a}^{\times} = \frac{(\hat{n}^{a} \cdot \hat{p})(\hat{n}^{a} \cdot \hat{q})}{1 + \hat{n}^{a} \cdot \hat{k}}$
- ▶ with  $\hat{k} = -\{\sin\theta_S \cos\phi_S, \sin\theta_S \sin\phi_S, \cos\theta_S\}$   $\hat{p} = \hat{u}\cos\psi + \hat{v}\sin\psi$   $\hat{q} = \hat{u}\cos\psi + \hat{v}\sin\psi$   $\hat{u} = \{\cos\theta_S \cos\phi_S, \cos\theta_S \sin\phi_S, -\sin\theta_S\}$   $\hat{v} = \{\sin\phi_S, -\cos\phi_S, 0\}$ ▶ Parameters:
- Parameters.
  - $\theta_S, \phi_S$  : equatorial sky position of GW source
  - $\psi$  : polarisation
  - $p^a$  : direction of the pulsar



## DA pipeline (EPTA)





# Continous GWs Frequentist analysis



#### F-statistic Earth term only

- We consider only the Earth term which is the coherent term between all pulsars
- To apply F-statistic we want to rewrite the signal in the form

$$r_a^E(t) = \sum_j a_{(j)} h^a_{(j)}$$



• Beam patterns  $F_a^+$ ,  $F_a^{\times}$  can be rewritten in the form:

$$F_a^+ = F_c^a \cos(2\psi) + F_s^a \sin(2\psi)$$
$$F_a^{\times} = -F_s^a \cos(2\psi) + F_c^a \sin(2\psi)$$

$$\begin{aligned} \bullet \text{ with } \qquad F_c^a &= \left\{ \frac{1}{4} (\sin^2(\theta_a) - 2\cos^2(\theta_a)) \sin^2(\theta_S) - \\ &\quad \frac{1}{2} \cos(\theta_a) \sin(\theta_a) \sin(2\theta_S) \cos(\phi_S - \phi_a) + \\ &\quad \frac{1}{4} (1 + \cos^2(\theta_S)) \sin^2(\theta_a) \cos(2\phi_S - 2\phi_a) \right\} \frac{1}{1 + \hat{n}^a.\hat{k}} \\ F_s^a &= \left\{ \cos(\theta_a) \sin(\theta_a) \sin(\theta_S) \sin(\phi_S - \phi_a) + \\ &\quad \frac{1}{2} \sin^2(\theta_a) \cos(\theta_S) \sin(2\phi_a - 2\phi_S) \right\} \frac{1}{1 + \hat{n}^a.\hat{k}}. \end{aligned}$$

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#### F-statistic Earth term only Thus, we have $r_a^E(t) = \sum a_{(j)}h_{(j)}^a$

▶ with

$$h_{(1)} = F_c^a \sin(\Phi(t)), \quad h_{(2)} = F_s^a \sin(\Phi(t)),$$
  
$$h_{(3)} = F_c^a \cos(\Phi(t)), \quad h_{(4)} = F_s^a \cos(\Phi(t)),$$

) and

$$a_{(1)} = \frac{\mathcal{A}}{2\pi f} [(1 + \cos^2 \iota) \cos(2\psi) \cos(\Phi_0) - 2\cos\iota\sin(2\psi)\sin(\Phi_0)],$$
  

$$a_{(2)} = \frac{\mathcal{A}}{2\pi f} [(1 + \cos^2 \iota) \sin(2\psi) \cos(\Phi_0) + 2\cos\iota\cos(2\psi)\sin(\Phi_0)],$$
  

$$a_{(3)} = \frac{\mathcal{A}}{2\pi f} [(1 + \cos^2 \iota) \cos(2\psi)\sin(\Phi_0) + 2\cos\iota\sin(2\psi)\cos(\Phi_0)],$$
  

$$a_{(4)} = \frac{\mathcal{A}}{2\pi f} [(1 + \cos^2 \iota) \sin(2\psi)\sin(\Phi_0) - 2\cos\iota\cos(2\psi)\cos(\Phi_0)].$$

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#### F-statistic Earth term only

• We parametrize only the deterministic signal h, not the noises and the stochastic background

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m}det(G^T C G)}} \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{r}\right)^T G(G^T C G)^{-1} G^T\left(\vec{\delta t} - \vec{r}\right)\right)$$

=> the log likelihood is:  $\log \mathcal{L} = <\delta \vec{t}_a | \vec{r_a} > -\frac{1}{2} < \vec{r_a} | \vec{r_a} >$ 

with inner product:  $\langle \vec{x} | \vec{y} \rangle = \vec{x} G(G^T C G)^{-1} G^T \vec{y}$ 



## F-statistic Earth term only

- Analytic maximisation of the likelihood over  $a_{(j)}$ :  $\frac{\partial \log(\mathcal{L})}{\partial a_{(j)}} = 0$
- Solution:  $a_{(j)} = M_{kj}^{-1} X_j$ 
  - with  $X_j \equiv \sum_{\alpha=1}^P \langle x_\alpha | h_{(j)}^{\alpha} \rangle$  and  $M_{jk} \equiv \sum_{\alpha=1}^P \langle h_{(j)}^{\alpha} | h_{(k)}^{\alpha} \rangle$
- New statistic:  $\mathcal{F}_e = \{\log(\mathcal{L})\}_{max} = \frac{1}{2}X_k M_{jk}^{-1} X_j$
- To be maximized over the 3 remaining parameters (intrinsic):
  - Equatorial latitude:  $\theta_S$ ,
  - Equatorial latitude:  $\phi_S$
  - Frequency: f







# Search algorithm MS-GA: Multi Search

# Framework to run in parallel several dedicated search methods :

- → "Global searches" looks for new good candidates avoiding the ones already found.
- "Local searches" explores in details the best candidates found at the previous step.
- "Modes separation" : the results are combined to find a new set of best candidates using some criterions (high SNR and not to close to the others).
- Each search is done by a GA with a special tuning or a parallel MCMC (EMCEE)



# Individual sources and data

Template :

- Several individual sources : Earth term only, non-eccentric, fixed frequencies.
- Fstatistic : analytic maximisation over 4 parameters (Petiteau & al. 2012, Babak & Sesana 2012, Ellis & al. 2012) ⇒ search for Nsrc x 3 parameters (sky positions & frequencies).
  - GW background (GWB) and red-noise on pulsars (RN) can be taken into account in likelihood computation (use of design matrix based on Van Hasterren & Levin 2012 ) ...
     MS-GA can search for GWB + RN + individual source parameters (ex. in Stas Babak's talk).
     Search :
- MS-GA coupled with several technics for removing «ghost detections» (checking of correlation between pulsars, high-pass filtering, cyclic removing of pulsars).
  - Characterisation of errors and likelihood distribution via MCMC (Antoine Lassus's talk).

[Data : simplified dataset or par/tim dataset ( $\Rightarrow$  MS-GA can work with real dataset).


# Results on simulated data

Npulsars	noise	dataset	Nsrc	ind. src SNR	signal	Results
30-50	noiseless or white 50-100 ns	simplified	1-5	> 10	earth term, same frequency	Pilot study Babak & Sesana 2012
30-50	white 30-200 ns	simplified	3-8	> 10	earth term, <b>≠ freq.</b>	OK Petiteau & al. 2012

- MS-GA successfully identified all the injected sources in all datasets
- MS-GA found all source parameters : sky position offset by less than few degrees and frequencies found with precision better than 0.1 nHz (errors characterisation : Antoine Lassus' talk)





▶ 2 Fe is distributed as X<sup>2</sup> with n degrees of freedom [Janaroski & Krolak LRR, 15 (2012)]

$$p'_{0}(\mathcal{F}) = p_{\chi^{2}}(2\mathcal{F}) = \frac{(2\mathcal{F})^{n/2-1}e^{-(2\mathcal{F})/2}}{2^{n/2} \Gamma(n/2)} = \frac{(\mathcal{F})^{n/2-1}e^{-\mathcal{F}}}{2 \Gamma(n/2)}$$

• After normalization and approximations:

$$p_0(\mathcal{F}) = e^{-\mathcal{F}} \frac{\mathcal{F}^{n/2-1}}{(n/2-1)!} \left( \text{exact} : e^{-\mathcal{F}} \frac{\mathcal{F}^{n/2-1}}{\Gamma(n/2)} \right)$$



► The false alarm probability,  $P_F$  that  $\mathcal{F}_e$  exceed  $\mathcal{F}_{th}$  when there is no signal is  $P_F(\mathcal{F}_{th}) = \int_{\mathcal{F}_{th}}^{\infty} p_0(\mathcal{F}) d\mathcal{F}$ 

$$= e^{-\mathcal{F}_{th}} \sum_{k=0}^{n/2-1} \frac{\mathcal{F}_{th}^k}{k!} \left( \text{exact} : \frac{\Gamma(n/2, \mathcal{F}_{th})}{\Gamma(n/2)} \right)$$

- Probability that  $\mathcal{F} < \mathcal{F}_{th}$  for one template is:  $1 P_F(\mathcal{F}_{th})$
- Probability that  $\mathcal{F} < \mathcal{F}_{th}$  for one template is:  $[1 P_F(\mathcal{F}_{th})]^{N_{cell}}$
- ▶ Finally, probability that 𝓕>𝓕<sub>th</sub> for at least one template is the total false alarm probability:

$$P_F^T(\mathcal{F}_{th}) = 1 - [1 - P_F(\mathcal{F}_{th})]^{N_{cell}}$$



In the case of *Fe*, we marginalize over 4 parameters per source so  $N = 4 N_{src}$  then :  $\sum_{2N_{src}-1} \mathcal{F}_{a,th}^{k}$ 

$$P_F(\mathcal{F}_{e,th}) = e^{-\mathcal{F}_{e,th}} \sum_{k=0} \frac{\mathcal{F}_{e,th}}{k!}$$

• The total false alarm probability is than

$$P_{F}^{T} = 1 - \left[1 - e^{-\mathcal{F}_{e,th}} \sum_{k=0}^{2N_{src}-1} \frac{\mathcal{F}_{e,th}^{k}}{k!}\right]^{N_{cell}}$$

• We need  $N_{cell}$  the total number of templates ?



- We need  $N_{cell}$  the total number of templates:
  - criterion for considering that 2 templates defined by parameters  $\xi$  and  $\xi'$  are independent is that the autocovariance function is

$$C(\xi,\xi') \le \rho C(\xi,\xi) = \rho \frac{n}{2}$$

- with ho=0.5
- Stochastic template bank
- ▶ 1 source: False Alarm Probabilty

$$P_F^T = 1 - [1 - e^{-\mathcal{F}_{e,th}} (1 + \mathcal{F}_{e,th})]^{N_{cell,1}}$$

By inverting, we got the threshold corresponding to the false alarm probability



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### Apply on real-data

- Data: EPTA data release 2
  - 10- of 41 millisecond pulsars
- Noises analysis:
  - Hardest part: estimate C
  - Estimation of the noise contribution for each pulsar using various technics:
    - Red noise
    - Dispersion measurements
    - White noise component per back-end per pulsar
      - EFAC
      - EQUAD



### Apply on real-data

#### Detection results of Fp on real data



![](_page_42_Picture_3.jpeg)

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### **Upper** limit

![](_page_43_Figure_1.jpeg)

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### Results

#### Babak et al. EPTA MNRAS 455.2 (2016)

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![](_page_44_Figure_2.jpeg)

### **Fp** statistic

#### Ellis et al. (2012)

- Excess power in which we basically search for extra power at a given frequency in each pulsar data
- Maximisation over all parameters except frequency f

$$r_a(t) = \sum_{j=1}^2 b_{(j,a)}(\mathcal{A}, \theta_S, \phi_S, \Psi, \iota, \Phi_0, \Phi_a) \kappa_{(j)}(\omega, t)$$

- Distribution of Fp statistic:
  - Gaussian noise: central :  $p_0(\mathcal{F}_p) = \frac{\mathcal{F}_p^{n/2-1}}{(n/2-1)!} \exp(-\mathcal{F}_p)$
  - Signal: non-central with optimal SNR  $\rho$  :

$$p_1(\mathcal{F}_p, \rho) = \frac{(2\mathcal{F}_p)^{(n/2-1)/2}}{\rho^{n/2-1}} I_{n/2-1}(\rho \sqrt{2\mathcal{F}_p}) e^{-\mathcal{F}_p - \frac{1}{2}\rho}$$

![](_page_45_Picture_9.jpeg)

### Fe detection results Babak et al. EPTA MNRAS 455.2 (2016)

- Fp evaluated at 99 independent frequencies
- ▶ p-value = 0.93
- Take into account the uncertainty in the noise parameters by sampling from their posterior distribution derived from the single pulsar analysis

![](_page_46_Figure_4.jpeg)

### Results

#### Babak et al. EPTA MNRAS 455.2 (2016)

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![](_page_47_Figure_2.jpeg)

## Continous GWs Bayesian analysis

![](_page_48_Picture_1.jpeg)

### Bayesian analysis 1

- No evolving sources: frequency at the pulsar is the same as the earth frequency.
  - we should sample on  $7+N_{pulsar}$  parameters

$$(\mathcal{A}, \theta_S, \phi_S, \Psi, \iota, \omega, \Phi_0, \Phi_a)$$

- numerical marginalization over the pulsar phase φ<sub>a</sub> [Taylor et al., 2014]
- MultiNest
- Analysis:
  - 41 pulsars with fixed noise
  - 6 pulsars with varying noise

![](_page_49_Picture_9.jpeg)

### Bayesian analysis 2

#### Full response:

- $7 + 2 N_{pulsar}$  parameter space
- Parallel tempering MCMC
- Analysis:
  - 41 pulsars with Earth term only
  - 6 pulsars with pulsar and Earth terms

![](_page_50_Picture_7.jpeg)

### Results

- Bayes factor  $\mathcal{B} = \frac{\int \mathcal{L}(\vec{\theta}, \vec{\lambda} | \vec{\delta t}) \pi(\vec{\theta}, \vec{\lambda}) d\vec{\theta} d\vec{\lambda}}{\int \mathcal{L}(\vec{\theta} | \vec{\delta t}) \pi(\vec{\theta}) d\vec{\theta}}.$ 
  - Non evolving:  $\log(\mathcal{B}) = -0.27$
  - Earth term only:  $\log(\mathcal{B}) = -0.31$ 
    - => no detection
      - => upper limit

![](_page_51_Picture_6.jpeg)

![](_page_52_Figure_0.jpeg)

### **Results EPTA**

![](_page_53_Figure_1.jpeg)

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### **Upper limit**

![](_page_54_Figure_1.jpeg)

### Horizon

#### Invert amplitude

![](_page_55_Figure_2.jpeg)

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### Results from NANOGrav Arzoumanian et al. NANOGrav (2014)

![](_page_56_Figure_1.jpeg)

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### Stochastic background

![](_page_57_Picture_1.jpeg)

#### Likelihood:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m}det(G^T C G)}}$$
$$\times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(G^T C G)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

- Stochastic background shape: power-law
- Parametrization of the correlation matrix:

$$C_{GWB} = \zeta_{\alpha\beta} A^2 \left(\frac{1yr^{-1}}{f_L}\right)^{\gamma-1} \left[ \Gamma(1-\gamma) \sin\frac{\pi\gamma}{2} (f_{L\tau_{ij}})^{\gamma-1} - \sum_{n=0}^{\infty} \frac{(f_{L\tau_{ij}})^{2n}}{(2n)!(2n+1-\gamma)} \right]$$
$$\zeta_{\alpha\beta} = \frac{3}{2} y \ln y - \frac{1}{4} y + \frac{1}{2} + \frac{1}{2} \delta_{\alpha\beta} , \quad y = \frac{1-\cos\theta_{\alpha\beta}}{2} , \quad \tau_{ij} = 2\pi |t_i - t_j|$$

![](_page_58_Picture_6.jpeg)

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#### Likelihood:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(GTCG)}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(GTCG)^{-1}G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

Stochastic background shape: power-law

Parametrization of the correlation matrix:

$$C_{GWB} = \zeta_{\alpha\beta} A^2 \left(\frac{1yr^{-1}}{f_L}\right)^{\gamma-1} \left[ \Gamma(1-\gamma) \sin\frac{\pi\gamma}{2} (f_{L\tau_{ij}})^{\gamma-1} - \sum_{n=0}^{\infty} \frac{(f_{L\tau_{ij}})^{2n}}{(2n)!(2n+1-\gamma)} \right]$$
$$\zeta_{\alpha\beta} = \frac{3}{2} y \ln y - \frac{1}{4} y + \frac{1}{2} + \frac{1}{2} \delta_{\alpha\beta} , \quad y = \frac{1-\cos\theta_{\alpha\beta}}{2} , \quad \tau_{ij} = 2\pi |t_i - t_j|$$

![](_page_59_Picture_6.jpeg)

DEROJ

#### Likelihood:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(GTCG)}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(GTCG)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

- Stochastic background shape: power-law
- Parametrization of the correlation matrix:

$$C_{GWB} = \zeta_{\alpha\beta} A^2 \left( \frac{1yr^{-1}}{f_L} \right)^{\gamma-1} \left[ \Gamma(1-\gamma) \sin \frac{\pi\gamma}{2} (f_{L\tau_{ij}})^{\gamma-1} - \sum_{n=0}^{\infty} \frac{(f_{L\tau_{ij}})^{2n}}{(2n)!(2n+1-\gamma)} \right]$$
$$\zeta_{\alpha\beta} = \frac{3}{2} y \ln y - \frac{1}{4} y + \frac{1}{2} + \frac{1}{2} \delta_{\alpha\beta} , \quad y = \frac{1-\cos\theta_{\alpha\beta}}{2} , \quad \tau_{ij} = 2\pi |t_i - t_j|$$

![](_page_60_Picture_6.jpeg)

DEROJ

#### Likelihood:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(GTCG)}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(GTCG)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

- Stochastic background shape: power-law
- Parametrization of the correlation matrix:

$$C_{GWB} = \zeta_{\alpha\beta} A^2 \left( \frac{1yr^{-1}}{f_L} \right)^{\gamma - 1} \left[ \Gamma(1 - \gamma) \sin \frac{\pi(\gamma)}{2} (f_{L\tau_{ij}})^{\gamma - 1} - \sum_{n=0}^{\infty} \frac{(f_{L\tau_{ij}})^{2n}}{(2n)!(2n+1-\gamma)} \right]$$
$$\zeta_{\alpha\beta} = \frac{3}{2} y \ln y - \frac{1}{4} y + \frac{1}{2} + \frac{1}{2} \delta_{\alpha\beta} , \quad y = \frac{1 - \cos \theta_{\alpha\beta}}{2} , \quad \tau_{ij} = 2\pi |t_i - t_j|$$

![](_page_61_Picture_6.jpeg)

#### Likelihood:

$$p(\vec{\delta t}|\vec{\theta}) = \frac{1}{\sqrt{(2\pi)^{n-m} det(GTCG)}} \times \exp\left(-\frac{1}{2}\left(\vec{\delta t} - \vec{h}\right)^T G(GTCG)^{-1} G^T\left(\vec{\delta t} - \vec{h}\right)\right)$$

- Stochastic background shape: power-law
- Parametrization of the correlation matrix:

$$C_{GWB} = \zeta_{\alpha\beta} A^{2} \left( \frac{1yr^{-1}}{f_{L}} \right)^{\gamma - 1} \left[ \Gamma(1 - \gamma) \sin \frac{\pi(\gamma)}{2} (f_{L\tau_{ij}})^{\gamma - 1} - \sum_{n=0}^{\infty} \frac{(f_{L\tau_{ij}})^{2n}}{(2n)!(2n+1-\gamma)} \right]$$
$$\left( \zeta_{\alpha\beta} = \frac{3}{2}y \ln y - \frac{1}{4}y + \frac{1}{2} + \frac{1}{2}\delta_{\alpha\beta} \right), \quad y = \frac{1 - \cos\theta_{\alpha\beta}}{2} , \quad \tau_{ij} = 2\pi |t_{i} - t_{j}|$$

#### Hellings Downs curve

### Hellings Downs curve

 Overlap reduction function: expected correlation in pulsar timing residuals due to an isotropic stochastic GW background

![](_page_63_Figure_2.jpeg)

Data Analysis IV : PTA - A. Petiteau - GW School - Benasque - 5 to 9 June 2017

DIDEROT

### Search for stochastic background

- Parameterization of C and inverting C is time consuming because its a matrix n x n, with n the total number of measurements.
- Several methods:
  - Bayesian / Frequentist
  - Fixed noise / varying noise
  - Fixed slope / varying slope
  - Various samplers

![](_page_64_Picture_7.jpeg)

### Search for stochastic background

#### Bayesian:

- Use the likelihood previously define
- Sampling algorithm:
  - MultiNest [EPTA]
  - parallel tempering MCMC [Ellis]
- Inputs: priors
- Results: posterior distribution for:
  - amplitude and slope of the background
  - noises parameters for individual pulsars
  - common noises

![](_page_65_Picture_11.jpeg)

### Search for stochastic background Bayesian: Lentati et al. EPTA (2015)

• Prior on parameters

Parameter	Description	Prior range		
White noise				
α	Global EFAC	uniform in [0.5 , 1.5]	1 parameter per pulsar (total 6)	
Spin-noise				
A <sub>SN</sub>	Spin-noise power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	1 parameter per pulsar (total 6)	
γsn	Spin-noise power law spectral index	uniform in [0,7]	1 parameter per pulsar (total 6)	
DM variations				
A <sub>DM</sub>	DM variations power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	1 parameter per pulsar (total 6)	
ΫDM	DM variations power law spectral index	uniform in [0,7]	1 parameter per pulsar (total 6)	
Common noise				
A <sub>CN</sub>	Uncorrelated common noise power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	1 parameter for the array	
γcn	Uncorrelated common noise power law spectral index	uniform in [0,7]	1 parameter for the array	
Aclk	Clock error power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	1 parameter for the array	
Yclk	Clock error power law spectral index	uniform in [0,7]	1 parameter for the array	
$A_{\rm eph}$	Solar System ephemeris error power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	3 parameters for the array (x, y, z)	
Yeph	Solar System ephemeris error power law spectral index	uniform in [0,7]	3 parameters for the array (x, y, z)	
Stochastic GWB				
Α	GWB power law amplitude	uniform in $[10^{-20}, 10^{-10}]$	1 parameter for the array	
γ	GWB power law spectral index	uniform in [0,7]	1 parameter for the array	
$\rho_i$	GWB power spectrum coefficient at frequency $i/T$	uniform in $[10^{-20}, 10^{0}]$	1 parameter for the array per frequency in	
			unparameterised GWB power spectrum model (total 20)	
Stochastic background angular correlation function				
c <sub>14</sub>	Chebyshev polynomial coefficient	uniform in [-1, 1]	see Eq. (36)	
$\Gamma_{IJ}$	Correlation coefficient between pulsars (I,J)	uniform in [-1, 1]	1 parameter for the array per unique pulsar pair (total 15)	

![](_page_66_Picture_3.jpeg)

ARI

Francois Ara

### Search for stochastic background

#### • Example: EPTA

#### Lentati et al. EPTA (2015)

• data

![](_page_67_Figure_4.jpeg)

Pulsar	J0613-0200	J1012+5307	J1600-3053	J1713+0747	J1744-1134	J1909-3744
Dataspan (yr)	16.05	16.83	7.66	17.66	17.25	9.38
N <sub>sys</sub> <sup>a</sup>	14	15	4	14	9	3
$\sigma(\mu s)^{b}$	1.691	1.610	0.563	0.679	0.801	0.131
Log <sub>10</sub> A <sub>SN</sub>	-13.58 ± 0.40 (-13.41)	$-13.05 \pm 0.09 (-13.04)$	-13.71 ± 0.54 (-13.42)	-14.31 ± 0.46 (-14.20)	-13.63 ± 0.27 (-13.60)	$-14.22 \pm 0.42$ (-13.98)
γsn	$2.50 \pm 0.99$ (2.09)	$1.56 \pm 0.37 (1.56)$	1.91 ± 1.05 (1.38)	$3.50 \pm 1.16 (3.51)$	$2.21 \pm 0.82$ (2.16)	$2.23 \pm 0.89$ (2.17)
Log <sub>10</sub> A <sub>DM</sub>	$-11.61 \pm 0.12$ (-11.57)	$-12.25 \pm 0.47$ (-11.92)	-11.75 ± 0.39 (-11.67)	-11.97 ± 0.14 (-11.90)	$-12.19 \pm 0.38$ (-11.93)	$-12.76 \pm 0.53$ (-12.51)
γ <sub>DM</sub>	$1.36 \pm 0.48 (1.11)$	$1.26 \pm 0.97 (0.27)$	$1.64 \pm 0.80 (1.46)$	$2.03 \pm 0.55 (1.82)$	1.41 ± 1.09 (0.36)	2.23 ± 1.07 (2.16)
Global EFAC	$1.01 \pm 0.02 (1.01)$	$0.98 \pm 0.02 \ (0.98)$	$1.03 \pm 0.04 (1.03)$	$1.00 \pm 0.02 (1.00)$	$1.01 \pm 0.03 (1.00)$	$1.02 \pm 0.04 (1.01)$
95% upper limit <sup>c</sup>	$9.7 \times 10^{-15}$	$8.3 \times 10^{-15}$	$2.1 \times 10^{-14}$	$4.4 \times 10^{-15}$	$7.0 \times 10^{-15}$	$5.2 \times 10^{-15}$

# Search for stochastic background Example: EPTA Lentati et al. EPTA (2015)

![](_page_68_Figure_1.jpeg)

### Search for stochastic background

#### **Example: EPTA**

#### Lentati et al. EPTA (2015)

Model	95% upper limit (×10 <sup>-15</sup> )
Bayesian Analysis	
Fixed Noise - Fixed Spectral Index	1.7
Varying Noise - Fixed Spectral Index	3.0
Additional Common Signals - Fixed Spectral Index	3.0
Fixed Noise - Varying Spectral Index	8.0
Varying Noise - Varying Spectral Index	13
Additional Common Signals - Varying Spectral Index	13

![](_page_69_Figure_4.jpeg)

### Others bayesian methods

- Measure power in frequency bins [Lentati et al.]
- Unparameterised power spectrum analysis for a correlated Gravitational Wave Background

95% upper limits from an unparameterised power spectrum analysis for a correlated GWB (red points), and uncorrelated common red noise process (blue points) for the 6 pulsar

![](_page_70_Figure_4.jpeg)

- Example of frequentist method used: Optimal statistic
  - weak signal maximum likelihood for GWB spectral amplitude (Anholm et al. 2009; Siemens et al. 2013; Chamberlin et al. 2014)
  - The statistic is:

$$\hat{A}^{2} = \frac{\sum_{IJ} \delta \mathbf{t}_{I}^{\mathrm{T}} \mathbf{P}_{I}^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_{J}^{-1} \delta \mathbf{t}_{J}}{\sum_{IJ} \mathrm{tr} \left[ \mathbf{P}_{I}^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_{J}^{-1} \tilde{\mathbf{S}}_{JI} \right]},$$

- with:
  - autocovariance of the post-fit residuals  $\mathbf{P}_I = \langle \delta \mathbf{t}_I \delta \mathbf{t}_I^{\mathrm{T}} \rangle$
  - signal term  $A^2 \tilde{\mathbf{S}}_{IJ} = \langle \delta \mathbf{t}_I \delta \mathbf{t}_J^{\mathrm{T}} \rangle = \mathbf{S}_{IJ}$

![](_page_71_Picture_8.jpeg)
# DA stochastic background

- Example of frequentist method used: Optimal statistic
  - weak signal maximum likelihood for GWB spectral amplitude (Anholm et al. 2009; Siemens et al. 2013; Chamberlin et al. 2014)
  - The statistic is:

$$\hat{A}^{2} = \frac{\sum_{IJ} \delta \mathbf{t}_{I}^{\mathrm{T}} \mathbf{P}_{I}^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_{J}^{-1} \delta \mathbf{t}_{J}}{\sum_{IJ} \mathrm{tr} \left[ \mathbf{P}_{I}^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_{J}^{-1} \tilde{\mathbf{S}}_{JI} \right]},$$

- with:
  - autocovariance of the post-fit residuals  $\mathbf{P}_I = \langle \delta \mathbf{t}_I \delta \mathbf{t}_I^{\mathrm{T}} \rangle$
  - signal term  $A^2 \tilde{\mathbf{S}}_{IJ} = \langle \delta \mathbf{t}_I \delta \mathbf{t}_J^{\mathrm{T}} \rangle = \mathbf{S}_{IJ}$



# DA stochastic background

- Example of frequentist method used: *Optimal statistic* 
  - SNR is:  $\rho = \frac{\hat{A}^2}{\sigma_0} = \frac{\sum_{IJ} \delta \mathbf{t}_I^{\mathrm{T}} \mathbf{P}_I^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_J^{-1} \delta \mathbf{t}_J}{\left(\sum_{IJ} \operatorname{tr} \left[\mathbf{P}_I^{-1} \tilde{\mathbf{S}}_{IJ} \mathbf{P}_J^{-1} \tilde{\mathbf{S}}_{JI}\right]\right)^{1/2}}$

measures how likely it is that we have found a crosscorrelated signal in our data rather than an uncorrelated signal

# • Measurement of cross-power values $\chi_{IJ} = \frac{\delta \mathbf{t}_I^T \mathbf{P}_I^{-1} \hat{\mathbf{S}}_{IJ} \mathbf{P}_J^{-1} \delta \mathbf{t}_J}{\operatorname{tr} \left[ \mathbf{P}_I^{-1} \hat{\mathbf{S}}_{IJ} \mathbf{P}_J^{-1} \hat{\mathbf{S}}_{JI} \right]}$

with  $\mathbf{S}_{IJ} = A^2 \zeta_{IJ} \hat{\mathbf{S}}_{IJ}$  and error  $\sigma_{0,IJ} = \left( \operatorname{tr} \left[ \mathbf{P}_I^{-1} \hat{\mathbf{S}}_{IJ} \mathbf{P}_J^{-1} \hat{\mathbf{S}}_{JI} \right] \right)^{-1/2}$ 

=> high SNR limit: cross-power values = Hellings Downs curve



## Results

#### Results of optimal-statistic on EPTA data



Lentati et al. EPTA (2015) <sup>66</sup> Data Analysis



# Upper limit result EPTA data Lentati et al. EPTA (2015)



# Upper limit on cosmic strings

- Background from cosmic string network.
- Parameters
  - string tension  $G\mu/c^2$  ,
  - $\alpha_{cs}$ : the birth-scale of loops relative to the horizon.



Lentati et al. EPTA (2015)

Separated approach ...



#### Separated approach including noises



- Global analysis including:
  - Pulsar + propagation parameters
  - Noises
  - GWs: continuous wave sources + backgrounds



- Global analysis including:
  - Pulsar + propagation parameters
  - Noises
  - GWs: continuous wave sources + backgrounds
- Could also include the pulse template matching
- Work in progress



# Future

#### More data:

- Continue to observe the pulsar
- Group all data in IPTA: EPTA, NANOGrav, PPTA
- Use more pulsars
- New instruments: SKA and it's precursor







### Future



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PARIS

# Future

#### More data:

- Continue to observe the pulsar
- Group all data in IPTA: EPTA, NANOGrav, PPTA
- Use more pulsars
- New instruments:
  SKA and
  it's precursor







# Thank you

