GRAVITATIONAL WAVES PROBE OF THE EARLY UNIVERSE



School on Gravitational Waves for Cosmology and Astrophysics, Benasque, May 28 - June 10, 2017





Einstein 1916 ... aLIGO 2015/16

* O(10) Solar mass Black Holes (BH) exist

* We can test the BH's paradigm (or Neutron Star physics)

* We can further test General Relativity (GR) [so far no deviation]

* We can observe the Universe through GWs



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(binaries)

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Extremely interesting ! but ... (binaries)

... We will focus on something else !



*







School on Gravitational Waves for Cosmology and Astrophysics 2017, May 28 - June 10

D. Blas (CERN), C. Caprini (APC), V. Cardoso (CENTRA, Lisbon), G. Nardini (U. Bern)



- * Late Universe: Hubble diagram from Binaries
- * Early Universe: High Energy Particle Physics



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* Late Universe: Hul Lectures by BONVIN

* Early Universe: High Energy Particle Physics



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PS

* Late Universe: Hul Lectures by BONVIN

* Early Universe: High Energy Particle Physics



Inflationary Period



(Image: Google Search)

Preheating



(Fig. credit: Phys.Rev. D67 103501)

Phase Transitions



(Image: PRL 112 (2014) 041301)

Cosmic Defects



(Image: Daverio et al, 2013)

Inflationary Period



(Image: Google Search)



Preheating



(Fig. credit: Phys.Rev. D67 103501)

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Inflationary Period



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(Fig. credit: Phys.Rev. D67 103501)



Cosmic Defects



(Image: Daverio et al. 2013)

OUTLINE

1) GWs from Inflation

Early Universe

- 2) GWs from Preheating
- 3) GWs from Phase Transitions

4) GWs from Cosmic Defects

OUTLINE

0) Cosmology

1) GWs from Inflation

Early Universe

- 2) GWs from Preheating
- 3) GWs from Phase Transitions

4) GWs from Cosmic Defects













A COSMOLOGY PRIMER



(evolution of the Universe)

Cosmology

Inflation =
$$\begin{pmatrix} initial \\ cond. \end{pmatrix}$$







BASICS of COSMOLOGY General Relativity theoretical pillars **Cosmological Pple** hot Big Bang (hBB) (evolution of the Universe) **Expansion** observational pillars ▲ CMB **BBN** Cosmology **Cosmological Pple** (initial cond.) 'cures' hBB Inflation = **CMB/LSS**

Gravitational Framework

General Relativity (GR)



 $G_{\mu
u} = \frac{1}{m_p^2} T_{\mu
u}$ geometry matter

$$m_p = (8\pi G)^{-1/2} = 2.44 \cdot 10^{18} \,\mathrm{GeV}$$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$

DIFF:
$$x^{\mu} \to x'^{\mu}(x)$$

symmetry

Gravitational Framework





Gravitational Framework






$$G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$$
 geometry of matter within the Universe the Universe



 $\begin{array}{ll} G_{\mu\nu}=\frac{1}{m_p^2}T_{\mu\nu} \\ \text{geometry of} & \text{matter within} \\ \text{the Universe} & \text{the Universe} \end{array}$

Principle of Symmetry:

The Universe is Homogeneous & Isotropic



 $G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$ geometry of matter within the Universe the Universe

Principle of Symmetry:

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Principle of Symmetry:

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$$g^{[U]}_{\mu\nu} \equiv \operatorname{diag}\left(-1, \frac{a^2(t)}{1-kr^2}, a^2(t)r^2, a^2(t)r^2\sin^2\theta\right)$$

 $G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$

geometry of matter within

the Universe the Universe

FLRW Friedmann-Lemaître -Robertson-Walker

Principle of Symmetry:

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 $G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$

the Universe the Universe

matter within

geometry of

FLRW Friedmann-Lemaître -Robertson-Walker

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = -dt^{2} + a^{2}(t) \left\{ \frac{dr^{2}}{1 - (kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta d\varphi^{2}) \right\}$$
2 dof: Scale Factor (dynamical) Curvature
$$\begin{cases} k < 0, \text{Open} \\ k = 0, \text{Flat} \\ k > 0, \text{Close} \end{cases}$$

Principle of Symmetry:

The Universe is Homogeneous & Isotropic

$$g_{\mu\nu}^{[U]} \equiv \text{diag}\left(-1, \frac{a^2(t)}{1 - kr^2}, a^2(t)r^2, a^2(t)r^2 \sin^2\theta\right)$$
 Frie -Ro

 $G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$

geometry of matter within

the Universe the Universe

FLRW Friedmann-Lemaître -Robertson-Walker

$$ds^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = -dt^{2} + a^{2}(t)\left\{\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta d\varphi^{2})\right\}$$

invariant:
$$\begin{cases} k \to k/c^2 \\ r \to c \cdot r \\ a \to a/c \end{cases} \longrightarrow \begin{cases} a, r, k \text{ unphysical} \\ \frac{k}{a^2}, a \cdot r, kr^2 \text{ physical} \end{cases}$$



 $G_{\mu\nu} = \frac{1}{m_p^2} T_{\mu\nu}$ geometry of matter within the Universe the Universe

Principle of Symmetry:

The Universe is Homogeneous & Isotropic

Redshift

$$z_1 \equiv \frac{a_o - a_1}{a_1}$$

$$1 + z \equiv \frac{a(t_o)}{a(t)}$$







$$\begin{split} \mathbf{H} \, \mathbf{\&} \, \mathbf{I} \\ T_{\nu}^{\mu} &\equiv \operatorname{diag}(-\rho, p, p, p) \\ \mathbf{\bigvee} \\ \mathbf{\bigvee} \\ m_{p}^{2} G_{\nu}^{\mu} \left[g_{**}^{(FRW)} \right] = T_{\nu}^{\mu} \end{split}$$



(I)+(II)
$$\longrightarrow \frac{1}{\rho} \frac{d\rho}{dt} = -\frac{3}{a} \frac{da}{dt} (1+w)$$
 (III)

1) GR +

2) H & I

UNIVERSE:



$$\frac{1}{a}\frac{d^{2}a}{dt^{2}} = -\frac{\rho}{6m_{p}^{2}}(1+3w) \quad \text{(I)}$$

$$H^{2} \equiv \left(\frac{1}{a}\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}} \quad \text{(II)}$$

$$\frac{1}{\rho}\frac{d\rho}{dt} = -\frac{3}{a}\frac{da}{dt}(1+w) \quad \text{(III)}$$

$$\left(w \equiv \frac{p}{\rho}\right) \quad \text{Equation of}$$

$$\text{State (EoS)}$$

(II)
$$H^2 \equiv \left(\frac{da}{dt}\right)^2 = \frac{\rho}{3m_p^2} - \frac{K}{a^2} \longrightarrow \begin{bmatrix} \rho_c \equiv 3m_p^2 H^2 \end{bmatrix}$$

Critical density $(\rho = \rho_c \Leftrightarrow K = 0)$

(II)
$$H^{2} \equiv \left(\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}} \longrightarrow \begin{bmatrix} \rho_{c} \equiv 3m_{p}^{2}H^{2} \\ \text{Critical density} & (\rho = \rho_{c} \Leftrightarrow K = 0) \end{bmatrix}$$

 $\rho = \sum_{i} \rho_{i} \; ; \; \Omega_{i} \equiv \frac{\rho_{i}}{\rho_{c}} \implies \Omega \equiv \frac{\rho}{\rho_{c}} = \sum_{i} \Omega_{i} \implies \boxed{\Omega - 1 \equiv \frac{k}{a^{2}H^{2}}}$
Cosmic Sum

(II)
$$H^{2} \equiv \left(\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}} \longrightarrow \left[\begin{array}{c} \rho_{c} \equiv 3m_{p}^{2}H^{2} \\ \text{Critical density} \end{array}\right] (\rho = \rho_{c} \Leftrightarrow K = 0)$$

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$$\begin{cases} \Omega > 1 \Rightarrow \operatorname{Close}(k > 0) \\ \Omega = 1 \Rightarrow \operatorname{Flat}(k = 0) \\ \Omega < 1 \Rightarrow \operatorname{Open}(k < 0) \end{cases} \quad \begin{array}{l} \text{one-to-one} \\ \text{correlation} \end{cases}$$

$$(II) \quad H^{2} \equiv \left(\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}} \longrightarrow \left[\begin{array}{c} \rho_{c} \equiv 3m_{p}^{2}H^{2} \\ \text{Critical density} \end{array}\right] \quad (\rho = \rho_{c} \Leftrightarrow K = 0)$$

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$$(III) \quad \frac{1}{\rho} \frac{d\rho}{dt} = -\frac{3}{a} \frac{da}{dt} (1+w) \implies \rho \propto e^{-3\int \frac{da}{a} (1+w)} = \begin{cases} 1/a^{3} & \text{, Mat.}(w=0) \\ 1/a^{4} & \text{, Rad.}(w=1/3) \\ \text{const.} & \text{, C.C.}(w=-1) \end{cases}$$

(II)
$$H^{2} \equiv \left(\frac{da}{dt}\right)^{2} = \frac{\rho}{3m_{p}^{2}} - \frac{K}{a^{2}} \longrightarrow \left[\begin{array}{c} \rho_{c} \equiv 3m_{p}^{2}H^{2} \\ \text{Critical density} \end{array}\right] (\rho = \rho_{c} \Leftrightarrow K = 0)$$

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(III) + (II) :
 $H^{2}(z) = H^{2} \left[\begin{array}{c} \Omega^{(p)} \left(\frac{a_{p}}{a_{p}}\right)^{4} + \Omega^{(p)} \left(\frac{a_{p}}{a_{p}}\right)^{3} + \Omega^{(p)} \left(\frac{a_{p}}{a_{p}}\right)^{2} + \Omega^{(p)} - 3\int \frac{da}{a}(1+w) \end{array}\right]$

$$H^{2}(a) = H_{o}^{2} \left\{ \Omega_{\mathrm{R}}^{(o)} \left(\frac{a_{o}}{a} \right)^{2} + \Omega_{\mathrm{M}}^{(o)} \left(\frac{a_{o}}{a} \right)^{2} + \Omega_{\mathrm{k}}^{(o)} \left(\frac{a_{o}}{a} \right)^{2} + \Omega_{\mathrm{DE}}^{(o)} e^{-3 \int \frac{da}{a} (1+w)} \right\}$$
$$\equiv H_{o}^{2} E^{2}(a)$$

$$E(a) \equiv \sqrt{\Omega_{\rm R}^{(o)} \left(\frac{a_o}{a}\right)^4 + \Omega_{\rm M}^{(o)} \left(\frac{a_o}{a}\right)^3 + \Omega_{\rm k}^{(o)} \left(\frac{a_o}{a}\right)^2 + \Omega_{\rm DE}^{(o)} e^{-3\int \frac{da}{a}(1+w)}} \qquad \Omega_{\rm k}^{(o)} \equiv -\frac{k}{a_o^2 H_o^2}$$



Past: particle ensemble

Statistical Mechanics

(III)
$$\frac{d\rho}{dt} + 3H(\rho + p) = 0 \longrightarrow \frac{dU}{dt} + p\frac{dV}{dt} = 0, \qquad \begin{cases} U = a^3\rho, \\ V = a^3 \end{cases}$$



Past: particle ensemble

Statistical Mechanics

$$\implies \begin{cases} \frac{dU}{dt} + p \frac{dV}{dt} = T \frac{dS}{dt}, \longrightarrow \text{Thermal Eq.} \\ \frac{dS}{dt} = 0, \longrightarrow \text{Adiabatic Exp.} \end{cases}$$



Thermal Eq.	(densities)
$f n = g_* \int d\vec{p} f(\vec{p}) ,$	number
$\rho = g_* \int d\vec{p} E(\vec{p}) f(\vec{p})$, energy
$p = g_* \int d\vec{p} \frac{ \vec{p} ^2}{3E(\vec{p})} f(\vec{p})$, pressure
dof Dispersion S relation D	Statistical istribution









Past: Radiation Domination (RD)

$$\rho_R^{(i)} = f_i g_*^{(i)} \frac{\pi^2}{30} T_i^4 , \quad f_i = \begin{cases} 1, \ B\\ \frac{7}{8}, \ F \end{cases}$$







Adiabatic Exp:



$$a^3T^3g_*^{(s)}(T) = const.$$

$$g_{*}^{(s)}(T) \equiv \sum_{i} g_{*,i}^{(B)} \left(\frac{T_{i}}{T}\right)^{3} + \frac{7}{8} \sum_{i} g_{*,i}^{(F)} \left(\frac{T_{i}}{T}\right)^{3}$$



Adiabatic Exp:



$$g_*^{(s)}(T) \equiv \sum_i g_{*,i}^{(B)} \left(\frac{T_i}{T}\right)^3 + \frac{7}{8} \sum_i g_{*,i}^{(F)} \left(\frac{T_i}{T}\right)^3$$

When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \end{cases}$ Species Decoupling, $T \to T_i$, 2) Mass threshold, $T < 2m_i$,



BiGGER size, **SMALLER Temp**



TODAY [Galaxies, Clusters, ...] (13.700 Million years)

FIRST GALAXIES (500 Millions years)

ATOMS CREATION (300.000-400.000 years)

ATOMIC NUCLEI CREATION (3 minutes !)



hot early Universe $(\sim 1s)$



hot early Universe $(\sim 1s)$






When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \underbrace{\frac{\pi^{2}}{30}g_{*}(T)T^{4}}_{\rho_{R}} \Rightarrow t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$

When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

1

Decoupling:

$$\begin{split} \Gamma_{\rm int} &= \sigma \times \langle nv \rangle \longrightarrow \qquad N_{\rm int} = \int_t^{t+\Delta t} dt' \Gamma_{\rm int}(t') \sim \frac{\Gamma_{\rm int}(t)}{H(t)} \begin{cases} \ll 1 \Rightarrow \text{ decoupling} \\ \gg 1 \Rightarrow \text{ Thermal Eq.} \end{cases} \\ \text{section density } (\mathsf{c} = 1) \end{split}$$

$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \underbrace{\frac{\pi^{2}}{30}g_{*}(T)T^{4}}_{\rho_{R}} \Rightarrow t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$

When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

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$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \underbrace{\frac{\pi^{2}}{30}g_{*}(T)T^{4}}_{\rho_{R}} \Rightarrow t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$

When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

Decoupling:

$$\Gamma_{\rm int} = \sigma \times \langle nv \rangle \longrightarrow N_{\rm int} = \int_t^{t+\Delta t} dt' \Gamma_{\rm int}(t') \sim \frac{\Gamma_{\rm int}(t)}{H(t)} \begin{cases} \ll 1 \Rightarrow \text{decoupling} \\ \gg 1 \Rightarrow \text{Thermal Eq.} \end{cases}$$

Neutrino Decoupling:

$$\Gamma_{\nu} = \sigma_{\rm EW} \times \langle n \rangle \sim T^5 / M_W^4 \lesssim H(t) \quad \Leftrightarrow \quad T \lesssim T_{\nu-\rm dec} = 0.8 \,\mathrm{MeV}$$
$$\sim G_F^2 T^2 \sim T^3$$

$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \frac{\pi^{2}}{30} g_{*}(T)T^{4} \quad \Rightarrow \quad t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$

When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

Annihilation (mass threshold):

 $T < 2m_i \implies \begin{array}{l} {\rm Can't\ produce}\\ {\rm it\ anymore}\end{array}$ [Boltzman Supression $\ \sim e^{-m/T}$]

$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \frac{\pi^{2}}{30} g_{*}(T)T^{4} \quad \Rightarrow \quad t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$

When do
$$g_*(T), g_*^{(s)}(T)$$
 change ? $\begin{pmatrix} 1 \end{pmatrix}$ Species Decoupling, $T \to T_i$, 2) Mass threshold, $T < 2m_i$,

Annihilation (mass threshold):

$$T < 2m_i \Rightarrow {{\rm Can't\ produce}\atop{{\rm it\ anymore}}}$$

[Boltzman Supression $\sim e^{-m/T}$]

Example e+/e- Annihilation

$$e^+ + e^- \leftrightarrow 2\gamma$$
, $T > 511 \text{ keV}$
 $e^+ + e^- \rightarrow 2\gamma$, $T < 511 \text{ keV}$

$$T_{\gamma}(t > t_{e^{\pm}}) = \left(\frac{g_s^{<}}{g_s^{>}}\right)^{1/3} T_{\nu} = \left(\frac{11}{4}\right)^{1/3} T_{\nu}$$



When do $g_*(T), g_*^{(s)}(T)$ change ? $\begin{cases} 1 \\ 2 \end{cases}$ Mass threshold, $T < 2m_i$,

$$T_{\gamma}(t > t_{e^{\pm}}) = \left(\frac{g_{s}^{<}}{g_{s}^{>}}\right)^{1/3} T_{\nu}$$

$$= (11/4)^{1/3} T_{\nu}$$

$$T_{\nu} \sim 0.8 \text{ MeV} \qquad (\nu) \qquad T = 0.5 \text{ MeV}$$

$$g_{s}^{<} = \frac{7}{8} \cdot 2 \cdot 2 + 2 = \frac{11}{2} \qquad g_{s}^{>} = 2$$





How can we probe the early Universe ? GWs !

$$H^{2} = \frac{1}{(2t)^{2}} = \frac{1}{3m_{p}^{2}} \frac{\pi^{2}}{30} g_{*}(T) T^{4} \quad \Rightarrow \quad t \sim \left(\frac{T}{\text{MeV}}\right)^{-2} \text{ s}$$



How can we probe the early Universe ? G





Big Bang NucleosynthesisFormation of atomic nuclei (1s - 3 mins)



Protons, Neutrons Interact strongly

Big Bang NucleosynthesisFormation of atomic nuclei (1s - 3 mins)



Protons, Neutrons Interact strongly

Universe cools down...

... protons and neutrons don't have sufficient energy anymore

Then they join together forming atomic nuclei: Nuclear Physics!

Big Bang NucleosynthesisFormation of atomic nuclei (1s - 3 mins)



Protons, Neutrons Interact strongly

Universe cools down...

... protons and neutrons don't have sufficient energy anymore

Then they join together forming atomic nuclei: Nuclear Physics!

Atomic Nuclei created !

Big Bang Nucleosynthesis

Formation of atomic nuclei (Is - 3 mins)



NUCLEAR PHYSICS (measured in the lab)

Leads to predict abundances of

 $H, {}^{4}\!He, D, {}^{3}\!He, {}^{7}\!Li, \dots$

Another definitive proof of hot Big Bang framework !



Protons, Electrons, Photons



Interact electromagnetically

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

Thomson Scattering

Protons, Electrons, Photons

$$e^-, p + \leftrightarrow \gamma$$

Interact electromagnetically

Thomson Scattering
$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

$$m_e \sim 10^{-3} m_p \Rightarrow \sigma_T^{(e)} \sim 10^6 \sigma_T^{(p)}$$

Only electro-photon scattering matters !

Protons, Electrons, Photons



Interact electromagnetically

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

Thomson Scattering



Protons, Electrons, Photons

$$e^-, p + \leftrightarrow \gamma$$

Interact electromagnetically

Thomson Scattering
$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

$$X_e \equiv \frac{n_e}{n_H}$$

Protons, Electrons, Photons

Interact electromagnetically

$$e^-, p + \leftrightarrow \gamma$$

Thomson Scattering

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$



Protons, Electrons, Photons



Interact electromagnetically

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

Thomson Scattering





Protons, **Electrons**, **Photons**

Past,

 $e^-, p + \leftrightarrow \gamma$

Interact electromagnetically

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

Thomson Scattering



Later,

colder







Protons, Electrons, **Photons**

Past,

 $e^-, p + \leftrightarrow \gamma$

Interact electromagnetically

$$\sigma_T = \frac{8\pi}{3} \frac{e^4}{(4\pi)^2 m^2}$$

Consequence of all this ?

Later,

colder

Atoms form (first time) ! Photon background freed !



Later, colder

Photon background set free ! → Origin of CMB !



Later, colder

Photon background set free ! → Origin of CMB !

$$\Gamma_{\gamma} = \sigma_T^{(e)} \langle n_e v \rangle = \sigma_T^{(e)} X_e \eta_b n_{\gamma} \le H(T) \quad \Leftrightarrow \quad T \le T_{\text{dec}}^{(\gamma)} = 0.26 \text{ eV}$$
$$[\text{at } z_{\text{dec}} = 1100]$$

Recombination & release of t Cosmic¹ Microwave Background (CMB)



Atom Formation: Free propagation of light ! (Recombination)



¿Where is that light? Everywhere!



¿Where is that light? Everywhere!







(almost-)ISOTROPIC

But

There are small 'Anisotropies' (variations of 1/100.000 only !)



$$T_{\rm dec} + \Delta T \iff \rho_{\gamma} (1 + \delta \rho_{\gamma})$$

$$\Delta T(\hat{n}) \equiv \frac{T(\hat{n}) - \bar{T}_{\rm CMB}}{\bar{T}_{\rm CMB}}$$



$$T_{\rm dec} + \Delta T \iff \rho_{\gamma} (1 + \delta \rho_{\gamma})$$

$$\Delta T(\hat{n}) \equiv \frac{T(\hat{n}) - \bar{T}_{\rm CMB}}{\bar{T}_{\rm CMB}}$$

$$\Delta T(\hat{n}) = \sum_{l=1}^{\infty} \sum_{m=-l}^{+l} a_{lm} Y_{lm}(\hat{n})$$

$$a_{lm} \equiv \int d\hat{\Omega} Y_{lm}^*(\hat{n}) \Delta T(\hat{n})$$

$$\langle \Delta T(\hat{n}) \rangle_{4\pi} = 0 \Rightarrow \langle a_{lm} \rangle_{4\pi} = 0$$



However...

$$\langle \Delta T(\hat{n})T(\hat{m}) \rangle_{4\pi} = F(\hat{n} \cdot \hat{m}) \neq 0$$

$$F(\hat{n}\cdot\hat{m}) = \sum_{l} \frac{(2l+1)}{4\pi} C_l P_l(\hat{n}\cdot\hat{m})$$

$$T_{\rm dec} + \Delta T \iff \rho_{\gamma} (1 + \delta \rho_{\gamma})$$

$$\Delta T(\hat{n}) \equiv \frac{T(\hat{n}) - \bar{T}_{\rm CMB}}{\bar{T}_{\rm CMB}}$$

$$\Delta T(\hat{n}) = \sum_{l=1}^{\infty} \sum_{m=-l}^{+l} a_{lm} Y_{lm}(\hat{n})$$

$$a_{lm} \equiv \int d\hat{\Omega} Y_{lm}^*(\hat{n}) \Delta T(\hat{n})$$

$$\langle \Delta T(\hat{n}) \rangle_{4\pi} = 0 \Rightarrow \langle a_{lm} \rangle_{4\pi} = 0$$

$$\left\langle a_{lm}a_{l'm}^{*}\right\rangle = C_{l}\delta_{ll'}\delta_{mm'}$$

$$C_l \equiv \int d\hat{\Omega} F(\hat{n} \cdot \hat{m}) P_l(\hat{n} \cdot \hat{m})$$


$$\left\langle |\Delta T(\hat{n})|^2 \right\rangle = \sum_l \frac{(2l+1)}{4\pi} C_l$$

$$Y_{lm}(\hat{n}) \longrightarrow 2l \ 'zeros' : \Delta \theta \sim \frac{\pi}{l}$$



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$$\frac{L_*(\Omega_M, \Omega_R, \Omega_k, \ldots)}{d_A(z_{\text{dec}})} \equiv \Delta \theta_* = \frac{\pi}{l_*}$$

Photon-Decoupling: Sound horizon (and multiples)

$$l_1 \leftrightarrow c_s H_{dec}^{-1}(\Omega_M, \Omega_R, ...)$$



 $l_* = \frac{\pi d_A(z_{\text{dec}})}{L_*(\Omega_M, \Omega_R, \Omega_k, \ldots)}$

prediction

measurement

$$\left\langle |\Delta T(\hat{n})|^2 \right\rangle = \sum_l \frac{(2l+1)}{4\pi} C_l$$

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$$d_A(z_{\rm dec}) = \frac{H_o^{-1}}{(1+z_{\rm dec})|\Omega_k|^{1/2}} \sin\left\{ |\Omega_k|^{1/2} \int_0^{z_{\rm dec}} \frac{dz'}{E(z')} \right\}$$

$$E(a) \equiv \sqrt{\Omega_{\rm R}^{(o)} \left(\frac{a_o}{a}\right)^4 + \Omega_{\rm M}^{(o)} \left(\frac{a_o}{a}\right)^3 + \Omega_{\rm k}^{(o)} \left(\frac{a_o}{a}\right)^2 + \Omega_{\rm DE}^{(o)} e^{-3\int \frac{da}{a}(1+w)}}$$

$$l_{*} = \frac{\pi d_{A}(z_{\text{dec}})}{L_{*}(\Omega_{M}, \Omega_{R}, \Omega_{k}, ...)} \qquad \qquad d_{A}(z_{\text{dec}}) = \frac{H_{o}^{-1}}{(1 + z_{\text{dec}})|\Omega_{k}|^{1/2}} \sin\left\{|\Omega_{k}|^{1/2} \int_{0}^{z_{\text{dec}}} \frac{dz'}{E(z')}\right\}$$

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Properties of the Anisotropies (position of 1st acoustic peak) then Geometry of the Universe !

The UNIVERSE has FLAT GEOMETRY (k = 0)!



hot Big Bang (hBB)

BiGGER size, **SMALLER Temp**



BASICS of COSMOLOGY



BASICS of COSMOLOGY



Shortcomings of the hBB framework

Friedmann Equations



$$w \equiv \frac{p}{\rho}$$









1) Horizon Problem — Causality Violation ! hBB







Today : $|\Omega - 1|_o \lesssim 0.1$

DBB shortcompase
(motivation for inflation)
1) Horizon Problem Causality Violation !
1) Horizon Problem
$$|\Omega - 1| = \frac{|k|}{a^2 H^2} \begin{cases} \propto a^2, \text{ RD} \\ \propto a, \text{ MD} \end{cases}$$

(If curvature $\neq 0$, it grows unstable!) $|\Omega - 1| = \frac{|k|}{a^2 H^2} \begin{cases} \propto a^2, \text{ RD} \\ \propto a, \text{ MD} \end{cases}$
Today : $|\Omega - 1|_o \lesssim 0.1 \Rightarrow |\Omega - 1|_{\text{BBN}} = \left(\frac{a_{\text{eq}}}{a_o}\right) \left(\frac{a_{\text{BBN}}}{a_{\text{eq}}}\right)^2 |\Omega - 1|_o$
 $\approx \frac{1}{(1 + z_{\text{eq}})} \left(\frac{T_{\text{eq}}}{T_{\text{BBN}}}\right)^2 |\Omega - 1|_o$
 $\approx 10^{-3} 10^{-12} |\Omega - 1|_o \lesssim 10^{-18}$





Today : $|\Omega - 1|_o \lesssim 0.1 \Rightarrow |\Omega - 1|_{BBN} \lesssim 10^{-18}$



Today : $|\Omega - 1|_o \lesssim 0.1 \Rightarrow |\Omega - 1|_{BBN} \lesssim 10^{-18}$

$$\begin{aligned} |\Omega - 1|_{\rm GUT} &= \left(\frac{a_{\rm GUT}}{a_{\rm BBN}}\right)^2 |\Omega - 1|_{\rm BBN} \sim (T_{\rm BBN}/T_{\rm GUT})^2 |\Omega - 1|_{\rm BBN} \\ &\simeq 10^{-38} |\Omega - 1|_{\rm BBN} \lesssim 10^{-56} \end{aligned}$$



It might well be that $k = 0 \dots$





Need extra 'Ingredient'! ----> INFLATION !

