#### Controllability of the beam equation

## Cristina Urbani

#### Joint work with P. Cannarsa

#### GSSI - Gran Sasso Science Institute

#### VII Partial differential equations, optimal design and numerics

20/08-01/09 Centro de Ciencias de Benasque Pedro Pascual

• Bilinear control problems

- Bilinear control problems
- Model

- Bilinear control problems
- Model
- Previous results

Ξ.

- Bilinear control problems
- Model
- Previous results
- Controllability of the beam equation

Dynamical system:



Dynamical system:

$$\dot{y} = f(y, u)$$



Dynamical system:

$$\dot{y}=f(y,u)$$

- $y \in Y$  state of the system
- $\bullet \ u \in \mathcal{U} \text{ control}$

$$\dot{y} = Ay + uBy$$



Dynamical system:

$$\dot{y}=f(y,u)$$

- $y \in Y$  state of the system
- $u \in \mathcal{U}$  control

$$\dot{y} = Ay + uBy$$

• Problem: local controllability along a trajectory Let  $(\bar{y}, \bar{u})$  be a trajectory of the control system  $\dot{y} = f(y, u)$ . The control system is *locally controllable along the trajectory*  $(\bar{y}, \bar{u})$  if, for every  $\varepsilon > 0$ , there exists  $\nu > 0$  such that, for every  $(a, b) \in Y \times Y$  with  $|a - \bar{y}(0)| < \nu$  and  $|b - \bar{y}(T)| < \nu$ , there exists a trajectory (y, u) such that

$$y(0) = a, \quad y(T) = b,$$
  
$$|u(t) - \bar{u}(t)| \le \varepsilon, \quad t \in [0, T]$$



 ${\scriptstyle \bullet}$  trajectory  $\rightarrow$  linearized system



- ${\ \bullet\ }$  trajectory  $\rightarrow$  linearized system
  - $\bullet\,$  controllability of linearized system  $\rightarrow$  inversion theorem

K. Beauchard, 2008  $\Rightarrow$  Nash - Moser Theorem



- ${\ \bullet\ }$  trajectory  $\rightarrow$  linearized system
  - $\bullet\,$  controllability of linearized system  $\rightarrow\,$  inversion theorem

K. Beauchard, 2008  $\Rightarrow$  Nash - Moser Theorem

 $\bullet$  linearized system  $\rightarrow$  moment theory ( Ingham inequality )

Euler-Bernoulli model, 1750.



Euler-Bernoulli model, 1750. Assumptions:



Euler-Bernoulli model, 1750.

Assumptions:

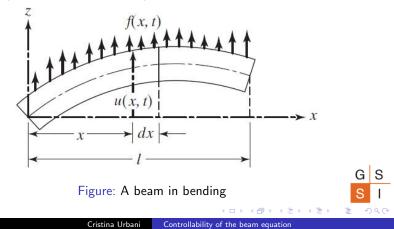
• linear elasticity of the material, Hooke's law



Euler-Bernoulli model, 1750.

Assumptions:

- linear elasticity of the material, Hooke's law
- plane sections remain plane and perpendicular to the neutral axis ( transverse vibrations ).



If u(t, x) is the beam deflection, the equation for the bending is

$$\rho A(x) \frac{\partial^2 u(t,x)}{\partial t^2} + \frac{\partial^2}{\partial x^2} \left[ EI(x) \frac{\partial^2 u(t,x)}{\partial x^2} \right] = f(t,x)$$

where  $\rho$  density per unit length, A(x) cross-sectional area, EYoung elastic modulus, I(x) cross-sectional area moment of inertia (about "z axis"), f(t, x) total external force.



 J. M. Ball, J. E. Marsdent, M, Slemrod, "Controllability for distributed bilinear systems", 1982 → the beam equation is not controllable in H<sup>2</sup><sub>0</sub>((0,1), ℝ) × L<sup>2</sup>((0,1), ℝ), with control p in L<sup>2</sup><sub>loc</sub>([0,+∞), ℝ)

- J. M. Ball, J. E. Marsdent, M, Slemrod, "Controllability for distributed bilinear systems", 1982 → the beam equation is not controllable in H<sup>2</sup><sub>0</sub>((0,1), ℝ) × L<sup>2</sup>((0,1), ℝ), with control p in L<sup>2</sup><sub>loc</sub>([0,+∞), ℝ)
- K. Beauchard, "Local controllability of one-dimentional beam equation", 2008  $\rightarrow$  local controllability in  $H^{5+\varepsilon}_{(0)}((0,1),\mathbb{R}) \times H^{3+\varepsilon}_{(0)}((0,1),\mathbb{R})$ , with  $\varepsilon > 0$  and control p in  $H^1_{loc}(\mathbb{R}_+,\mathbb{R})$ .



#### Control system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0. & (1) \end{cases}$$

 $\mu(x)$  is the effect of an axial force.



G S

#### Control system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0. \end{cases}$$
(1)

 $\mu(x)$  is the effect of an axial force. We introduce the operator A defined by

$$D(A) := H^4 \cap H^2_0((0,1), \mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

4

#### Control system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0. \end{cases}$$
(1)

 $\mu(x)$  is the effect of an axial force. We introduce the operator A defined by

$$D(A) := H^4 \cap H^2_0((0,1),\mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

whose eigenvalues and eigenvectors are

$$\lambda_k^2 = \mu_k = (k\pi)^4, \qquad \varphi_k = \sqrt{2}\sin(k\pi x), \qquad \forall k \in \mathbb{N}^*.$$

4

#### Control system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0. \end{cases}$$
(1)

 $\mu(x)$  is the effect of an axial force. We introduce the operator A defined by

$$D(A) := H^4 \cap H^2_0((0,1), \mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

whose eigenvalues and eigenvectors are

$$\lambda_k^2 = \mu_k = (k\pi)^4, \qquad \varphi_k = \sqrt{2}\sin(k\pi x), \qquad \forall k \in \mathbb{N}^*$$

 $\psi_k(t,x) = \varphi_k(x)e^{-i\lambda_k t}$  are solutions of (1) with control  $p \equiv 0$ .

#### The system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} + f(t) = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0, \\ u(0,x) = u_0(x), u_t(0,x) = u_1(x) \end{cases}$$
(2)

can be transformed into the Cauchy problem

$$\begin{cases} \frac{d\psi}{dt} = -\mathcal{A}\psi - p(t)\mu(x)\mathcal{B}\psi + F(t),\\ \psi(0) = \psi_0 \end{cases}$$
(3)

G S

#### The system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} + f(t) = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u(t,0) = u(t,1) = u_{xx}(t,0) = u_{xx}(t,1) = 0, \\ u(0,x) = u_0(x), u_t(0,x) = u_1(x) \end{cases}$$
(2)

can be transformed into the Cauchy problem

$$\begin{cases} \frac{d\psi}{dt} = -\mathcal{A}\psi - p(t)\mu(x)\mathcal{B}\psi + F(t),\\ \psi(0) = \psi_0 \end{cases}$$
(3)

where the linear operators  $\mathcal A$  e  $\mathcal B$  are defined as follows

$$D(\mathcal{A}) := H_{(0)}^4 \times H_0^2((0,1),\mathbb{R}), \qquad D(\mathcal{B}) := H_0^2 \times L^2((0,1),\mathbb{R})$$

$$\mathcal{A}\begin{pmatrix}\psi^1\\\psi^2\end{pmatrix} := \begin{pmatrix}-\psi^2\\\psi^1_{xxxx}\end{pmatrix}, \qquad \mathcal{B}\begin{pmatrix}\psi^1\\\psi^2\end{pmatrix} := \begin{pmatrix}0\\\psi^1_{xx}\end{pmatrix} \qquad \text{GS}$$
and  $F:(0,T) \to H_0^2 \times L^2((0,1),\mathbb{R}).$ 
Controllability of the heat equation

By choosing  $\psi = (u, u_t)$ ,  $\psi_0 = (u_0, u_1)$ , and F = (0, -f) the problems (2) e (3) are equivalent.



By choosing  $\psi = (u, u_t)$ ,  $\psi_0 = (u_0, u_1)$ , and F = (0, -f) the problems (2) e (3) are equivalent.

Let us consider as solution of the homogeneus problem

$$\hat{\psi}_1(t,x) = (\psi_1(t,x), (\psi_1(t,x))_t)$$

and initial datum

٠

$$\psi(0,x) = \hat{\varphi}_1(x) = (\varphi_1(x), \lambda_1 \varphi_1(x))$$

#### Theorem

Let T>0 and  $\mu\in H^3((0,1),\mathbb{R})$  such that

$$\exists c > 0 \quad \text{such that} \quad |\langle \mu(\varphi_1)_{xx}, \varphi_k \rangle| \ge \frac{c}{k^3}, \quad \forall k \in \mathbb{N}^*.$$
 (4)

There exists  $\delta > 0$  and a  $C^1$  map

$$\Gamma: \mathcal{V}_T \to L^2((0,T),\mathbb{R}),$$

$$\mathcal{V}_T := \{\Psi_f \in \mathcal{S}^2 \cap H^5_{(0)} \times H^3_{(0)}((0,1),\mathbb{C}); ||\Psi_f - \hat{\psi}_1(T)||_{H^5_{(0)} \times H^3_{(0)}} < \delta$$

such that,  $\Gamma(\hat{\psi}_1(T)) = 0$  and for all  $\Psi_f \in \mathcal{V}_T$  the solution of (3), with  $\psi = (u, u_t)$ ,  $\psi_0 = (u_0, u_1)$ , initial condition

U

$$\psi_0 = \hat{\varphi}_1$$

(5)

and control  $p = \Gamma(\Psi_f)$ , satisfies  $\psi(T) = \Psi_f$ .

Sketch of the proof:

• proof of existence, uniqueness and  $C^0([0,T], H_0^2 \times L^2((0,1), \mathbb{R}))$  regularity of the weak solution of

$$\begin{cases} \frac{d\psi}{dt} = -\mathcal{A}\psi - p(t)\mu(x)\mathcal{B}\psi + F(t),\\ \psi(0) = \psi_0, \end{cases}$$
(6)

Sketch of the proof:

• proof of existence, uniqueness and  $C^{0}([0,T], H_{0}^{2} \times L^{2}((0,1), \mathbb{R}))$  regularity of the weak solution of  $\int d\psi = A_{0}(x_{0}) R_{0}(x_{0}) R_{0}(x_{0}) = E(t)$ 

$$\begin{cases} \frac{a\psi}{dt} = -\mathcal{A}\psi - p(t)\mu(x)\mathcal{B}\psi + F(t),\\ \psi(0) = \psi_0, \end{cases}$$
(6)

definition of

$$V_T := \{ \xi = (\xi^1, \xi^2) \in L^2 \times L^2(0, 1); \Im \langle i\lambda_1 \xi^1 + \xi^2, \psi_1(T) \rangle = 0 \}$$

and the othogonal projection onto  $V_T$ 

$$P_T: L^2 \times L^2((0,T),\mathbb{R}) \to V_T.$$

• definition of the end point map

$$\Theta_T : L^2((0,T),\mathbb{R}) \to V_T \cap H^2_0 \times L^2(0,1)$$
$$u \mapsto P_T[\psi(T)],$$



• definition of the end point map

$$\Theta_T : L^2((0,T), \mathbb{R}) \to V_T \cap H^2_0 \times L^2(0,1)$$
$$u \mapsto P_T[\psi(T)],$$

• proof of  $C^1$  regularity of  $\Theta_T$ ,  $d\Theta_T(p)\cdot q=P_T(\Psi(T)) \text{ where } \Psi \text{ is solution of the linearized system,}$ 

$$\begin{cases} \Psi_t = -\mathcal{A}\Psi - p(t)\mu(x)\mathcal{B}\Psi - q(t)\mu(x)\mathcal{B}\psi, \\ \Psi(t,0) = \Psi(t,1) = 0, \Psi(0,x) = 0 \end{cases}$$

G S

• definition of the end point map

$$\Theta_T : L^2((0,T), \mathbb{R}) \to V_T \cap H^2_0 \times L^2(0,1)$$
$$u \mapsto P_T[\psi(T)],$$

• proof of  $C^1$  regularity of  $\Theta_T$ ,  $d\Theta_T(p)\cdot q=P_T(\Psi(T)) \text{ where } \Psi \text{ is solution of the linearized system,}$ 

$$\begin{cases} \Psi_t = -\mathcal{A}\Psi - p(t)\mu(x)\mathcal{B}\Psi - q(t)\mu(x)\mathcal{B}\psi, \\ \Psi(t,0) = \Psi(t,1) = 0, \Psi(0,x) = 0 \end{cases}$$

• proof of the existence of the map

$$d\Theta_T(0)^{-1}: V_T \cap H^2_0 \times L^2(0,1) \to L^2((0,T),\mathbb{R})$$

and its  $C^0$  regularity (controllability of the linearized system) GS

• definition of the end point map

$$\Theta_T : L^2((0,T), \mathbb{R}) \to V_T \cap H^2_0 \times L^2(0,1)$$
$$u \mapsto P_T[\psi(T)],$$

• proof of  $C^1$  regularity of  $\Theta_T$ ,  $d\Theta_T(p)\cdot q=P_T(\Psi(T)) \text{ where } \Psi \text{ is solution of the linearized system,}$ 

$$\begin{cases} \Psi_t = -\mathcal{A}\Psi - p(t)\mu(x)\mathcal{B}\Psi - q(t)\mu(x)\mathcal{B}\psi, \\ \Psi(t,0) = \Psi(t,1) = 0, \Psi(0,x) = 0 \end{cases}$$

• proof of the existence of the map

$$d\Theta_T(0)^{-1}: V_T \cap H^2_0 \times L^2(0,1) \to L^2((0,T),\mathbb{R})$$

and its  $C^0$  regularity (controllability of the linearized system) GS

 $\bullet\,$  moment problem  $\rightarrow$  (Corollary of) Ingham inequality

Cristina Urbani

Moment problem to solve

$$-\int_{0}^{T} q(s)e^{-i(\lambda_{k}-\lambda_{1})s}ds = d_{k-1}(\Psi_{f}) :=$$
$$= \left(i\frac{\langle\Psi_{f}^{1}(x),\varphi_{k}(x)\rangle}{\langle\mu(x)(\varphi_{1}(x))_{xx},\varphi_{k}(x)\rangle}\lambda_{k} + \frac{\langle\Psi_{f}^{2}(x),\varphi_{k}(x)\rangle}{\langle\mu(x)(\varphi_{1}(x))_{xx},\varphi_{k}(x)\rangle}\right)e^{-i\lambda_{k}T}$$

G S

Moment problem to solve

$$-\int_{0}^{T} q(s)e^{-i(\lambda_{k}-\lambda_{1})s}ds = d_{k-1}(\Psi_{f}) :=$$
$$= \left(i\frac{\langle\Psi_{f}^{1}(x),\varphi_{k}(x)\rangle}{\langle\mu(x)(\varphi_{1}(x))_{xx},\varphi_{k}(x)\rangle}\lambda_{k} + \frac{\langle\Psi_{f}^{2}(x),\varphi_{k}(x)\rangle}{\langle\mu(x)(\varphi_{1}(x))_{xx},\varphi_{k}(x)\rangle}\right)e^{-i\lambda_{k}T}$$

We need to show that given

$$\Psi_f = (\Psi_f^1, \Psi_f^2) \in V_T \cap H_0^2 \times L^2(0, 1),$$

there exists

$$q \in L^2((0,T),\mathbb{R})$$

that satisfies the moment problem.

GS

#### Corollary (of Ingham Theorem)

Let T>0 and  $(\omega_k)_{k\in\mathbb{N}}$  an increasing sequence in  $[0,+\infty)$  such that  $\omega_0=0,$  and

$$\omega_{k+1} - \omega_k \to +\infty$$
 when  $k \to +\infty$ .

There exist a linear and continuous map

$$L: l_r^2(\mathbb{N}, \mathbb{C}) \to L^2((0, T), \mathbb{R}),$$
$$d \mapsto L(d)$$

such that, for all  $d = (d_k)_{k \in \mathbb{N}} \in l_r^2(\mathbb{N}, \mathbb{C})$ , the function v := L(d) solves

$$\int_0^1 v(t)e^{i\omega_k t}dt = d_k, \qquad \forall k \in \mathbb{N}.$$

We should ensure that

•  $d_0 \in \mathbb{R}$ :

$$d_0 = \frac{\langle i\lambda_1 \Psi_f^1 + \Psi_f^2, \varphi_1 \rangle}{\langle \mu(x)(\varphi_1)_{xx}, \varphi_1 \rangle} e^{-i\lambda_1 T} \in \mathbb{R}$$



э

We should ensure that

•  $d_0 \in \mathbb{R}$ :

 $d_0 = \frac{\langle i\lambda_1 \Psi_f^1 + \Psi_f^2, \varphi_1 \rangle}{\langle \mu(x)(\varphi_1)_{xx}, \varphi_1 \rangle} e^{-i\lambda_1 T} \in \mathbb{R}$ 

•  $(d_k)_{k\in\mathbb{N}}\in l^2(\mathbb{N},\mathbb{C})$ :



We should ensure that

•  $d_0 \in \mathbb{R}$ :

 $d_0 = \frac{\langle i\lambda_1 \Psi_f^1 + \Psi_f^2, \varphi_1 \rangle}{\langle \mu(x)(\varphi_1)_{xx}, \varphi_1 \rangle} e^{-i\lambda_1 T} \in \mathbb{R}$ 

•  $(d_k)_{k\in\mathbb{N}}\in l^2(\mathbb{N},\mathbb{C})$ :

$$\begin{split} \sum_{k=1}^{\infty} |d_k(\Psi_f)|^2 &= \\ &= \sum_{k=1}^{\infty} \left| \left( i \frac{\langle \Psi_f^1(x), \varphi_k(x) \rangle}{\langle \langle (\mu(x)\varphi_1(x))_{xx}, \varphi_k(x) \rangle} \lambda_k + \frac{\langle \Psi_f^2(x), \varphi_k(x) \rangle}{\langle (\mu(x)\varphi_1(x))_{xx}, \varphi_k(x) \rangle} \right) e^{-i\lambda_k T} \right|^2 \\ &\leq 2 \sum_{k=1}^{\infty} \left| \frac{\langle \Psi_f^1(x), \varphi_k(x) \rangle}{\langle (\mu(x)\varphi_1(x))_{xx}, \varphi_k(x) \rangle} k^2 \pi^2 \right|^2 + \left| \frac{\langle \Psi_f^2(x), \varphi_k(x) \rangle}{\langle (\mu(x)\varphi_1(x))_{xx}, \varphi_k(x) \rangle} \right|^2 \\ &\leq 2 \sum_{k=1}^{\infty} \left| \frac{|\langle \Psi_f^1(x), \varphi_k(x) \rangle|}{c} k^5 \pi^2 \right|^2 + \left| \frac{|\langle \Psi_f^2(x), \varphi_k(x) \rangle|}{c} k^3 \right|^2 \\ &= 2 \sum_{k=1}^{\infty} \frac{|\langle k^5 \Psi_f^1(x), \varphi_k(x) \rangle|^2}{c} \pi^4 + \frac{|\langle k^3 \Psi_f^2(x), \varphi_k(x) \rangle|^2}{c} \\ &\leq 0 \\ \leq 0$$

Cristina Urbani Controllability of the beam equation

Recall: the space

$$H^s_{(0)}(I,\mathbb{C})$$

is equipped with the norm

$$||\phi||_{H^s_{(0)}} = \sum_{k=1}^{\infty} |k^s \langle \phi, \varphi_k \rangle|^2.$$



Recall: the space

$$H^s_{(0)}(I,\mathbb{C})$$

is equipped with the norm

$$||\phi||_{H^s_{(0)}} = \sum_{k=1}^{\infty} |k^s \langle \phi, \varphi_k \rangle|^2.$$

We can conclude the proof defining

$$\Gamma(\Psi_f) =: \Theta_T^{-1}[P_T \Psi_f].$$

G S

Recall: the space

$$H^s_{(0)}(I,\mathbb{C})$$

is equipped with the norm

$$||\phi||_{H^s_{(0)}} = \sum_{k=1}^{\infty} |k^s \langle \phi, \varphi_k \rangle|^2.$$

We can conclude the proof defining

$$\Gamma(\Psi_f) =: \Theta_T^{-1}[P_T \Psi_f].$$

Therefore the solution of the control system with  $p=\Gamma(\Psi_f)$  satisfies

$$\begin{split} \psi(T) &= P_T(\psi(T)) + \sqrt{1 - ||P_T\psi(T)||_{L^2 \times L^2}^2} \tilde{\psi}(T) = \\ &= P_T(\Psi_f) + \sqrt{1 - ||P_T\Psi_f||_{L^2 \times L^2}^2} \tilde{\psi}(T) = \Psi_f \end{split} \qquad \begin{array}{c} \mathbf{G} \in \mathbf{G} \\ \mathbf{S} \end{array}$$

Let us consider the system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u_x(t,0) = u_x(t,1) = u_{xxx}(t,0) = u_{xxx}(t,1) = 0. \end{cases}$$
(7)



Let us consider the system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u_x(t,0) = u_x(t,1) = u_{xxx}(t,0) = u_{xxx}(t,1) = 0. \end{cases}$$
(7)

We introduce the operator A defined by

$$D(A) := H^4 \cap H^2_0((0,1), \mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

.

GS

Let us consider the system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u_x(t,0) = u_x(t,1) = u_{xxx}(t,0) = u_{xxx}(t,1) = 0. \end{cases}$$
(7)

4

We introduce the operator  $\boldsymbol{A}$  defined by

$$D(A) := H^4 \cap H^2_0((0,1),\mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

whose eigenvalues and eigenvectors are

$$\mu_k = (k\pi)^4, \qquad \varphi_k = \sqrt{2}\cos(k\pi x), \qquad \forall k \in \mathbb{N}^*.$$

Let us consider the system

$$\begin{cases} u_{tt} + u_{xxxx} + p(t)\mu(x)u_{xx} = 0, & (t,x) \in \mathbb{R}_+ \times (0,1), \\ u_x(t,0) = u_x(t,1) = u_{xxx}(t,0) = u_{xxx}(t,1) = 0. \end{cases}$$
(7)

We introduce the operator  $\boldsymbol{A}$  defined by

$$D(A) := H^4 \cap H^2_0((0,1),\mathbb{R}), \qquad Av := \frac{d^4v}{dx^4}.$$

whose eigenvalues and eigenvectors are

$$\mu_k = (k\pi)^4, \qquad \varphi_k = \sqrt{2}\cos(k\pi x), \qquad \forall k \in \mathbb{N}^*.$$

Again the system can be transformed into

$$\left\{ \begin{array}{ll} \frac{d\psi}{dt} = -\mathcal{A}\psi - p(t)\mu(x)\mathcal{B}\psi + F(t), & \mbox{G S} \\ \psi(0) = \psi_0 & \mbox{S I} \end{array} \right.$$

#### Theorem

Let T > 0 and  $\mu \in H^3((0,1),\mathbb{R})$  such that

$$\exists c>0 \quad \textit{tale che} \quad |\langle \mu(\varphi_1)_{xx},\varphi_k\rangle| \geq \frac{c}{k^2}, \quad \forall k\in \mathbb{N}^*.$$

There exists  $\Gamma > 0$  and a  $C^1$  map

$$\Gamma: \mathcal{V}_T \to L^2((0,T),\mathbb{R}),$$

$$\mathcal{V}_T := \{\Psi_f \in \mathcal{S}^2 \cap H^4_{(0)} \times H^2_{(0)}((0,1),\mathbb{C}); ||\Psi_f - \hat{\psi}_1(T)||_{H^4_{(0)} \times H^2_{(0)}} < \delta\}$$

such that,  $\Gamma(\hat{\psi}_1(T)) = 0$  and for every  $\Psi_f \in \mathcal{V}_T$  the solution of (7), with  $\psi = (u, u_t)$ ,  $\psi_0 = (u_0, u_1)$ , initial condition

$$\psi_0 = \hat{\varphi}_1$$

and control  $p = \Gamma(\Psi_f)$ , satisfies  $\psi(T) = \Psi_f$ .

If we compute explicitly

$$\begin{aligned} \langle \mu(x)(\varphi_1)_{xx},\varphi_k\rangle &= \mu'(1) \left[ \frac{(-1)^k 2(k^2+1)}{(k^2-1)^2} \right] + \mu'(0) \left[ \frac{2(k^2+1)}{(k^2-1)^2} \right] + \\ &- \int_0^1 \mu'''(x) \left( \frac{\sin((k+1)\pi x)}{(k+1)^3\pi} + \frac{\sin((k-1)\pi x)}{(k-1)^3\pi} \right) dx \end{aligned}$$



If we compute explicitly

$$\begin{aligned} \langle \mu(x)(\varphi_1)_{xx},\varphi_k \rangle &= \mu'(1) \left[ \frac{(-1)^k 2(k^2+1)}{(k^2-1)^2} \right] + \mu'(0) \left[ \frac{2(k^2+1)}{(k^2-1)^2} \right] + \\ &- \int_0^1 \mu'''(x) \left( \frac{\sin((k+1)\pi x)}{(k+1)^3\pi} + \frac{\sin((k-1)\pi x)}{(k-1)^3\pi} \right) dx \end{aligned}$$

we have

If we compute explicitly

$$\begin{aligned} \langle \mu(x)(\varphi_1)_{xx},\varphi_k \rangle &= \mu'(1) \left[ \frac{(-1)^k 2(k^2+1)}{(k^2-1)^2} \right] + \mu'(0) \left[ \frac{2(k^2+1)}{(k^2-1)^2} \right] + \\ &- \int_0^1 \mu'''(x) \left( \frac{\sin((k+1)\pi x)}{(k+1)^3\pi} + \frac{\sin((k-1)\pi x)}{(k-1)^3\pi} \right) dx \end{aligned}$$

we have

$$\begin{split} \sum_{k=1}^{\infty} |d_k(\Psi_f)|^2 &= \sum_{k=1}^{\infty} \left| \left( i \frac{\langle \Psi_f^1(x), \varphi_k(x) \rangle}{\langle \langle (\mu(x)\varphi_1(x) \rangle_{xx}, \varphi_k(x) \rangle} \lambda_k + \frac{\langle \Psi_f^2(x), \varphi_k(x) \rangle}{\langle \langle (\mu(x)\varphi_1(x) \rangle_{xx}, \varphi_k(x) \rangle} \right) e^{-i\lambda_k T} \right|^2 \\ &\leq 2 \sum_{k=1}^{\infty} \left| \frac{|\langle \Psi_f^1(x), \varphi_k(x) \rangle|}{|\langle (\mu(x)\varphi_1(x) \rangle_{xx}, \varphi_k(x) \rangle|} k^2 \pi^2 \right|^2 + \left| \frac{|\langle \Psi_f^2(x), \varphi_k(x) \rangle|}{|\langle (\mu(x)\varphi_1(x) \rangle_{xx}, \varphi_k(x) \rangle|} \right|^2 \\ &\leq 2 \sum_{k=1}^{\infty} \left| \frac{\langle \Psi_f^1(x), \varphi_k(x) \rangle}{c} k^4 \pi^2 \right|^2 + \left| \frac{|\langle \Psi_f^2(x), \varphi_k(x) \rangle|}{c} k^2 \right|^2 \\ &= 2 \sum_{k=1}^{\infty} \frac{|\langle k^4 \Psi_f^1(x), \varphi_k(x) \rangle|^2}{c} \pi^4 + \frac{|\langle k^2 \Psi_f^2(x), \varphi_k(x) \rangle|^2}{c} \end{split}$$

# Thanks for your attention!

Cristina Urbani Controllability of the beam equation