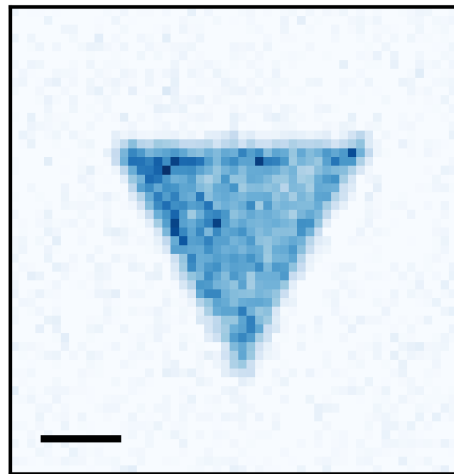


# Scale-invariant dynamics and breathers of an interacting 2D Bose gas

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# Introduction

- Scale invariance: When a physical system and its dilation behave similarly.
- Useful in high energy physics, in statistical physics, etc.
- How do we know a physical system is scale invariant?

→ Look at its action  $S = \int dt \mathcal{L}(\mathbf{q}, \dot{\mathbf{q}}, t)$

→ rescale space (and time), e.g:  $\mathbf{q} \rightarrow \mathbf{q}/\lambda$ ,  $\dot{\mathbf{q}} \rightarrow \lambda\dot{\mathbf{q}}$ ,  $t \rightarrow t/\lambda^2$ .

Is  $S$  invariant?

## Examples:

1. The ideal gas:  $\mathcal{L} = \frac{1}{2} \sum_i m_i \dot{\mathbf{q}}_i^2$

2.  $N$  particles interacting with  $U(\mathbf{q}_{ij}) \propto 1/\mathbf{q}_{ij}^2$

3. In cold atoms: Unitary Fermi gas

## Consequences:

$N = 1$ : Kepler problem.

Laplace-Runge-Lenz vector constant,  
closed orbits ( $E < 0$ )

Spectrum of the hydrogen atom

No bulk viscosity

# Introduction: the 2D Bose gas

- Weakly interacting bosons: description with a classical field  $\psi(\mathbf{r}, t)$ ,  
normalization  $\int d^2\mathbf{r} |\psi(\mathbf{r}, t)|^2 = 1$ , density  $n(\mathbf{r}, t) = N|\psi(\mathbf{r}, t)|^2$

- 2D gas: Interaction energy of  $N$  particles of mass  $m$ :

$$E_{\text{int}} = \frac{N^2 \hbar^2}{2m} \tilde{g} \int d^2\mathbf{r} |\psi(\mathbf{r}, t)|^4,$$

$\tilde{g}$ : dimensionless parameter. Description valid if  $\tilde{g} \ll 1$

- Full Lagrangian:

$$\mathcal{L}[\psi] = \int d^2\mathbf{r} \left[ i\hbar\psi^* \frac{\partial\psi}{\partial t} - \frac{N\hbar^2}{2m} |\nabla\psi|^2 - \frac{N^2\hbar^2}{2m} \tilde{g} |\psi|^4 \right]$$

- Transformation  $\mathbf{r} \rightarrow \mathbf{r}/\lambda$ ,  $t \rightarrow t/\lambda^2$ . Then  $\psi(\mathbf{r}) \rightarrow \lambda\psi(\lambda\mathbf{r})$  (normalization).

We have  $\mathcal{L} \rightarrow \lambda^2\mathcal{L}$ , and  $S$  is unchanged.

- Universal thermodynamics at equilibrium: Hung et al. Nature **140**, 236 (2011)  
Yefsah et al. PRL **107**, 130401 (2011)

# Introduction: the 2D Bose gas in a harmonic trap

- Add a harmonic trap of frequency  $\omega$ :

$$E_p = \frac{1}{2} \int d^2r m\omega^2 r^2 |\psi|^2$$

- Scaling transformation:  $E_p \rightarrow E_p/\lambda^2$

No scale invariance anymore!

However:

- Some transformations keep the action invariant.
- Interesting consequences on the dynamics of the gas: breathing mode at  $2\omega$

Pitaevskii and Rosch, PRA **55**, R853 (1997)

Chevy et al., PRL **88**, 250402 (2001)

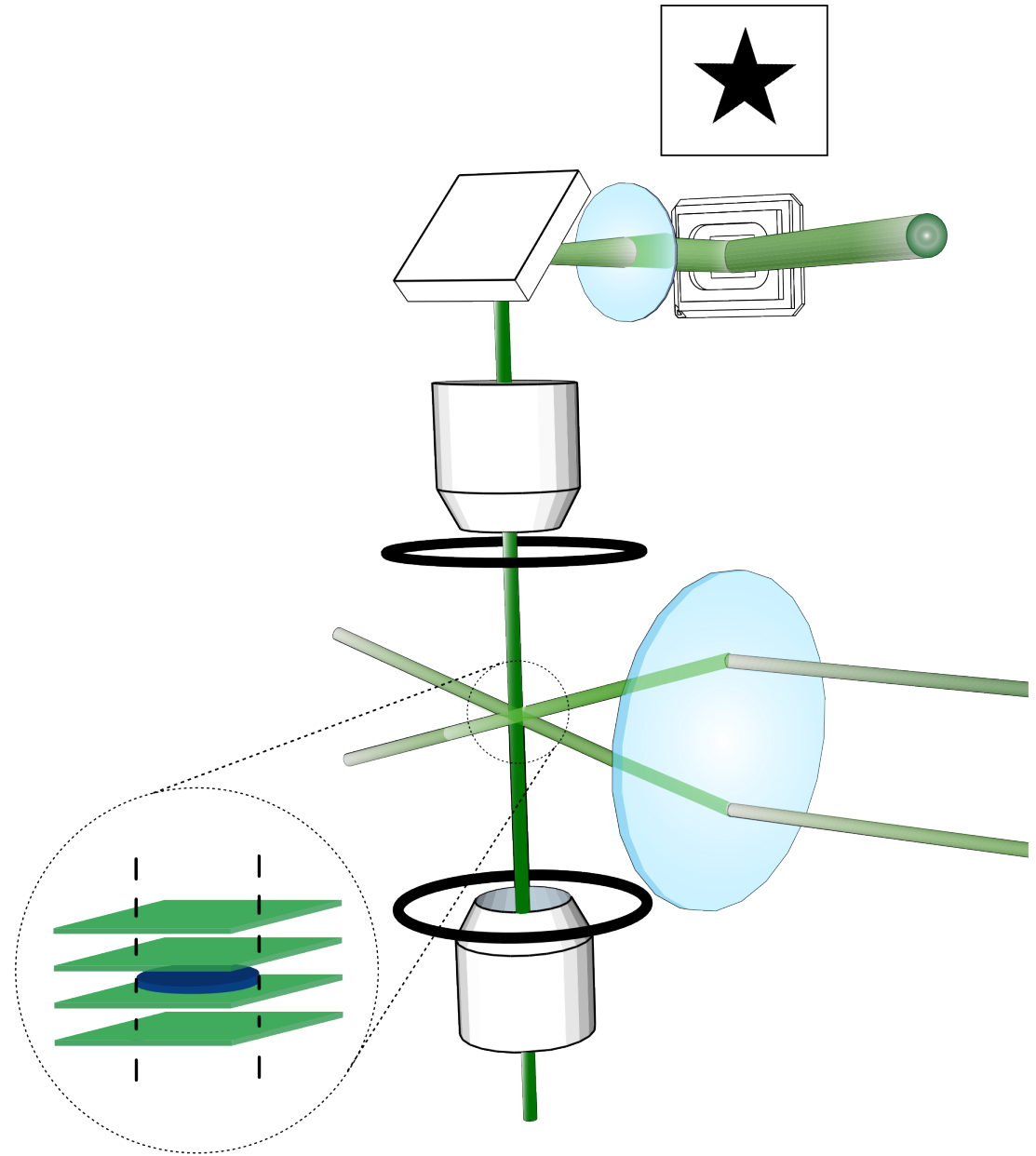
Vogt et al., PRL **108**, 070404 (2012)

# Experimental setup

- 2D confinement:  $\omega_z = 2\pi \cdot 4 \text{ kHz}$
- Repulsive walls:
  - uniform gas
  - arbitrary shape



- Internal state:  $F=1, m=0$ .
- Harmonic trap: with a magnetic field
- At  $t = 0$ , transfer to  $F=1, m=-1$ 
  - $\omega \approx 2\pi \cdot 20 \text{ Hz}$



**Goal of this talk:** Investigate the symmetries of this system with clouds far out-of-equilibrium in the harmonic trap.

# Outline of this talk

Introduction

- ① Symmetry group of the 2D Bose gas
- ② Linking evolutions of different clouds
- ③ Breathers of the Gross–Pitaevskii equation

Conclusion and outlook

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Conclusion and outlook

# SO(2,1) symmetry

Gross–Pitaevskii equation (GPE) of a free cloud:

$$i\hbar\frac{\partial\psi}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2\psi + \frac{\hbar^2}{m}\tilde{g}N|\psi|^2\psi.$$

Three types of interesting transformations:

- A time translation:  $\mathbf{r}, t \rightarrow \mathbf{r}, t + \beta$
- A scaling transformation:  $\mathbf{r}, t \rightarrow \mathbf{r}/\lambda, t/\lambda^2$
- An "expansion":  $\mathbf{r}, t \rightarrow \mathbf{r}/(\gamma t + 1), t/(\gamma t + 1)$

Combine them: 
$$\begin{cases} \mathbf{r} \\ t \end{cases} \rightarrow \begin{cases} \frac{\mathbf{r}}{\gamma t + \delta} \\ \frac{\alpha t + \beta}{\gamma t + \delta} \end{cases}, \quad \alpha\delta - \beta\gamma = 1$$

Matrix  $\begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \in \text{SL}(2, \mathbb{R})$ .

Infinitesimal generators:  $L_1 = \frac{i}{2}\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ ,  $L_2 = \frac{i}{2}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$  and  $L_3 = \frac{i}{2}\begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$

$$[L_1, L_2] = -iL_3, \quad [L_2, L_3] = iL_1, \quad [L_3, L_1] = iL_2 \quad \longrightarrow \text{SO}(2, 1)$$



# SO(2,1) symmetry: with a harmonic trap

- GPE:

$$i\hbar\frac{\partial\psi}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2\psi + \frac{\hbar^2}{m}\tilde{g}N|\psi|^2\psi + \frac{1}{2}m\omega^2 r^2\psi.$$

- Transformations:

$$\begin{cases} r \\ \eta = \tan(\omega t) \end{cases} \rightarrow \begin{cases} r/\lambda(t) \\ \eta' = (\alpha\eta + \beta)/(\gamma\eta + \delta) \end{cases}, \quad \alpha\delta - \beta\gamma = 1$$

$$\lambda(t) = [(\alpha \sin(\omega t) + \beta \cos(\omega t))^2 + (\gamma \sin(\omega t) + \delta \cos(\omega t))^2]^{1/2}$$

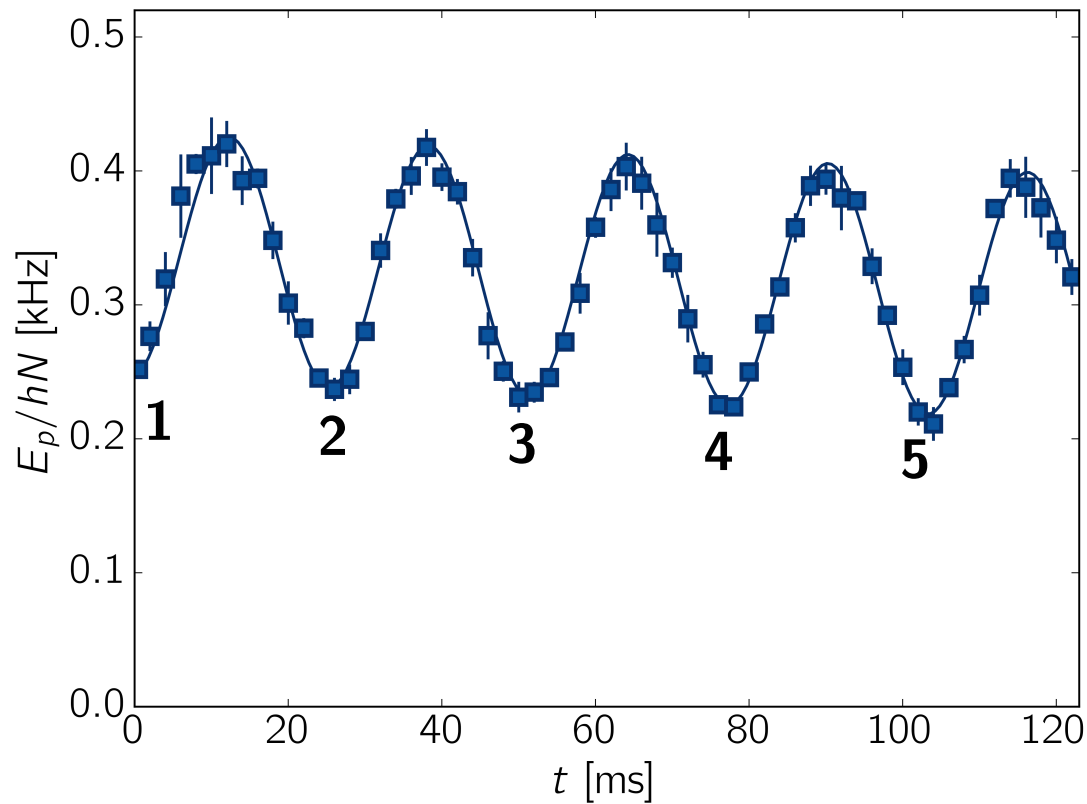
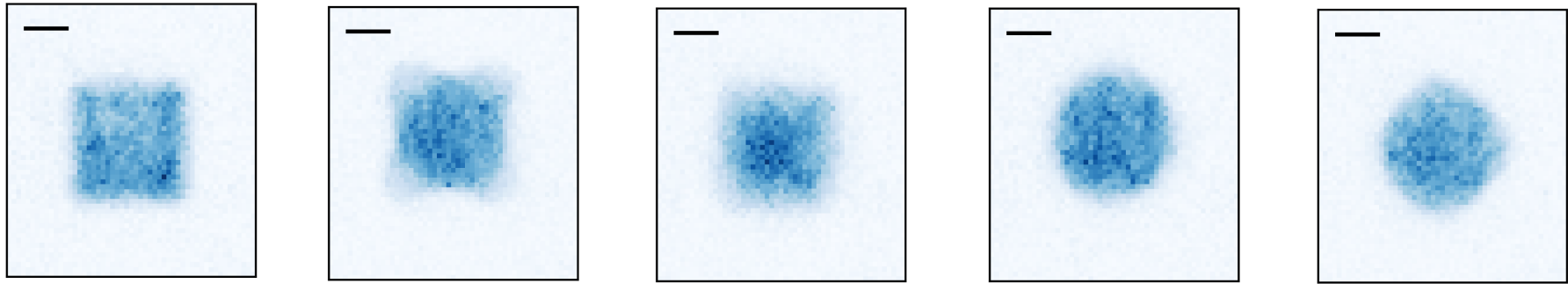
The symmetry group is still SO(2,1).

- A strong consequence:  $E_p \propto \sin(2\omega t)$

→ Breathing mode at  $2\omega$

We can test this on the experiment.

# Periodic potential energy



$$\omega_{\text{fit}} = 2\pi \cdot 38.5(1) \text{ Hz}$$
$$\omega = 2\pi \cdot 19.3(1) \text{ Hz}$$

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- ③ Breathers of the Gross–Pitaevskii equation

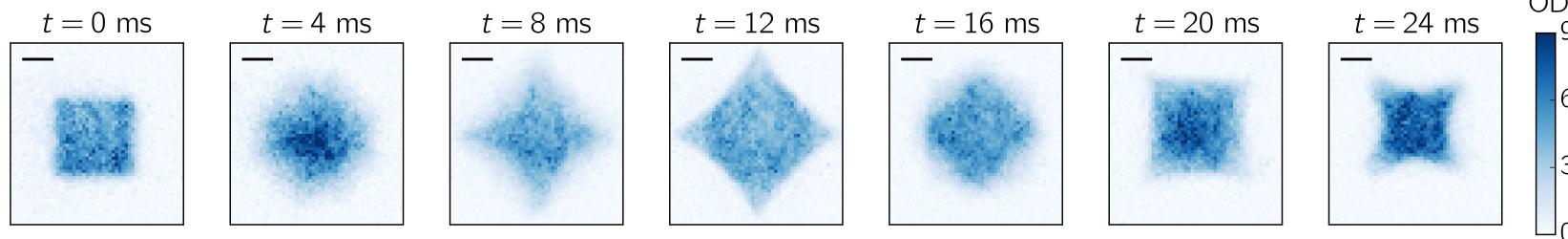
Conclusion and outlook

# Goal of this part

$$\omega_1 = \omega_2 = 2\pi \cdot 20 \text{ Hz}$$

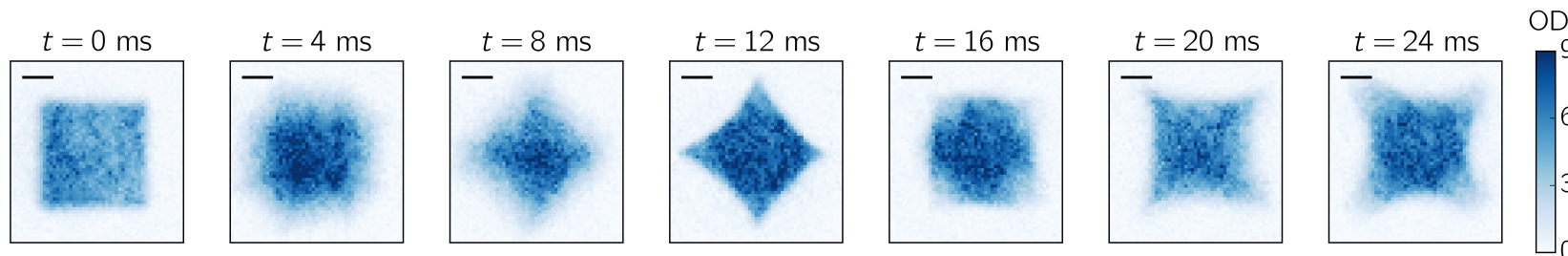
$$L_1 = 27 \mu\text{m}$$

$$N_1 = 3.7 \cdot 10^4$$



$$L_2 = 37 \mu\text{m}$$

$$N_2 = 5.4 \cdot 10^4$$



Evolution of  $\psi_1$

$$\left\{ \begin{array}{l} L_1 \\ N_1 \\ \omega_1 \\ \tilde{g}_1 \end{array} \right.$$

same initial  
shape



Symmetries

Evolution of  $\psi_2$

$$\left\{ \begin{array}{l} L_2 \\ N_2 \\ \omega_2 \\ \tilde{g}_2 \end{array} \right.$$

Three parameters:  $\delta = L_2/L_1$ ,

$$\mu = (\tilde{g}_2 N_2 / \tilde{g}_1 N_1)^{1/2},$$

$$\zeta = \omega_2 / \omega_1$$

# Evolution of clouds with the same $\tilde{g}N$

- GPE:

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + \frac{\hbar^2}{m} \tilde{g}N |\psi|^2 \psi + \frac{1}{2} m \omega^2 r^2 \psi.$$

- **Evolution of a dilation:** Apply a SO(2,1) transformation to a solution  $\psi(\mathbf{r}, t)$ :

$$\psi(\mathbf{r}, t) \rightarrow \lambda(t) \psi(\lambda(t)\mathbf{r}, \tau)$$

$$\lambda(t) = \left[ \delta^2 \cos^2(\omega t) + \frac{1}{\delta^2} \sin^2(\omega t) \right]^{-1/2}$$

$$\delta^2 \tan(\omega \tau) = \tan(\omega t)$$

$$L_1 \qquad L_2 = \delta L_1$$

$$\omega_1 \qquad \omega_2 = \omega_1$$

Choice  $\beta = \gamma = 0$ : cloud at rest at  $t = 0$ .

- **Evolution with a different  $\omega$ :** Apply a simple scaling transformation to  $\psi(\mathbf{r}, t)$ :

$$\psi(\mathbf{r}, t) \rightarrow \lambda \psi(\lambda \mathbf{r}, \lambda^2 t)$$

$$L_1 \qquad L_2 = L_1 / \lambda$$

$$\omega_1 \qquad \omega_2 = \omega_1 / \lambda^2$$

- Know  $\psi(\mathbf{r}, t) \rightarrow$  know the evolution of any dilation in a trap of any frequency

# Evolution of clouds with different $\tilde{g}N$ 's

- Only in the hydrodynamic regime: Translate GPE in terms of

$$n(\mathbf{r}) = N|\psi(\mathbf{r})|^2 \quad \mathbf{v}(\mathbf{r}) = \frac{\hbar \operatorname{Im}(\psi^*(\mathbf{r})\nabla\psi(\mathbf{r}))}{m n(\mathbf{r})} :$$

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{v}) = 0$$

$$\frac{\partial \mathbf{v}}{\partial t} + \nabla \cdot \left( \frac{1}{2}m\mathbf{v}^2 + \frac{\hbar^2}{m}\tilde{g}n + \frac{1}{2}m\omega^2\mathbf{r}^2 - \frac{\hbar^2 \nabla^2 \sqrt{n}}{2m \sqrt{n}} \right) = 0 \quad \text{Healing length } \xi \ll L$$

- Invariant under  $\mathbf{r} \rightarrow \mu\mathbf{r}$ ,  $t \rightarrow t/\mu$ ,  $\tilde{g}n \rightarrow \mu^2\tilde{g}n$ ,  $\mathbf{v} \rightarrow \mu\mathbf{v}$ ,  $\omega \rightarrow \mu\omega$ .

**Summary:** Link evolution  $\psi(\mathbf{r}, t)$  (size  $L_1$ ,  $N_1$ ,  $\omega_1$ ,  $\tilde{g}_1$ ) with evolution of:

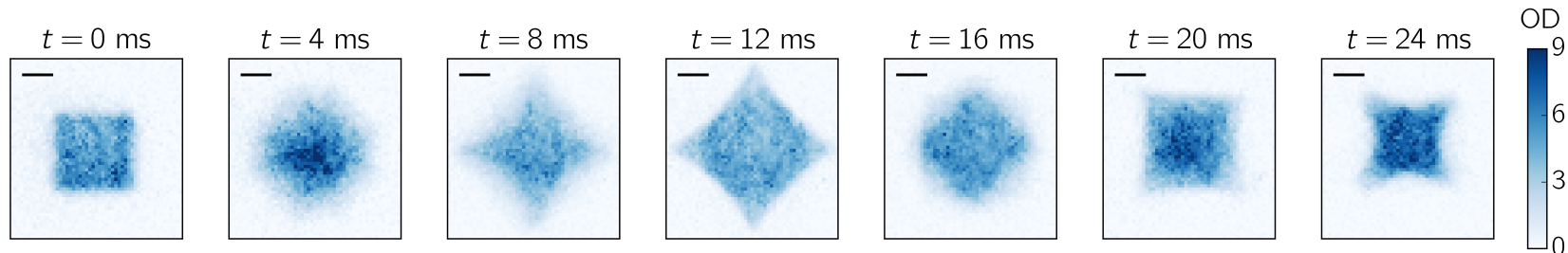
- (1) any dilation of  $\psi$  with size  $L_2 = \delta L_1$
- (2) in any trap frequency  $\omega_2 = \zeta\omega_1$
- (3) with any atom number  $N_2$  and parameter  $\tilde{g}_2$ , with  $\tilde{g}_2 N_2 = \mu^2 \tilde{g}_1 N_1$

$$\lambda(t) = \left[ \delta^2 \cos^2(\omega_2 t) + \left( \frac{\mu}{\delta\zeta} \right)^2 \sin^2(\omega_2 t) \right]^{-1/2} \quad \tan(\omega_1 \tau) = \frac{\mu}{\zeta\delta^2} \tan(\omega_2 t)$$

# On the experiment

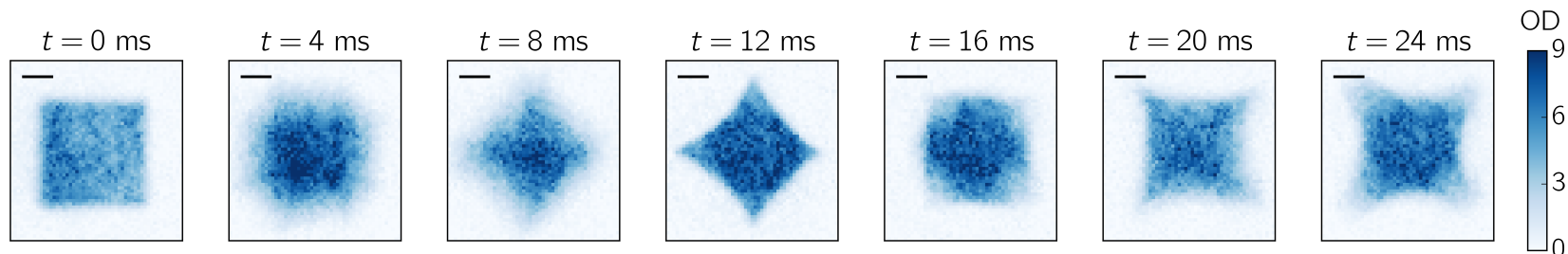
$$L_1 = 27 \mu\text{m}$$

$$N_1 = 3.7 \cdot 10^4$$



$$L_2 = 37 \mu\text{m}$$

$$N_2 = 5.4 \cdot 10^4$$



Overlap:

$$\mathcal{O}(n_1, n_2) = \max_{\lambda} \left( \lambda \int d^2r \frac{n_1(\lambda r) n_2(r)}{\|n_1\| \|n_2\|} \right)$$

Theory, no free parameter:

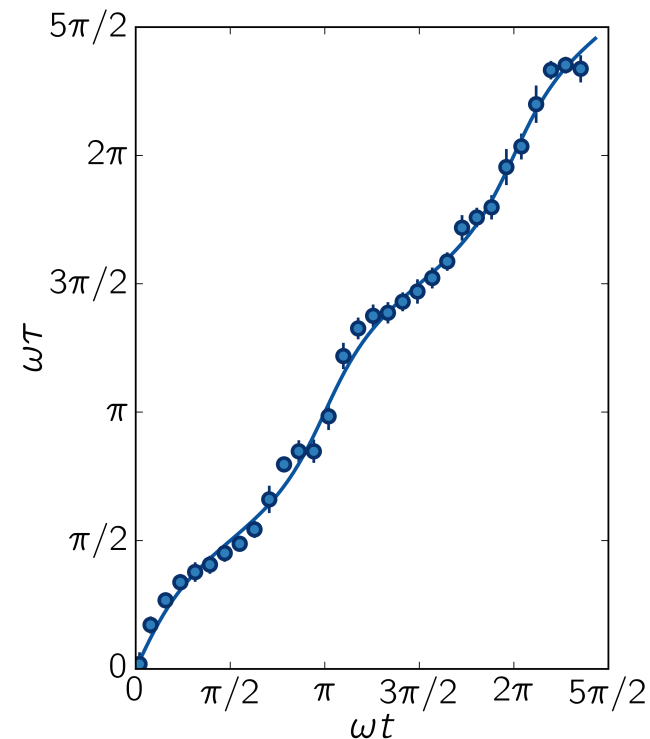
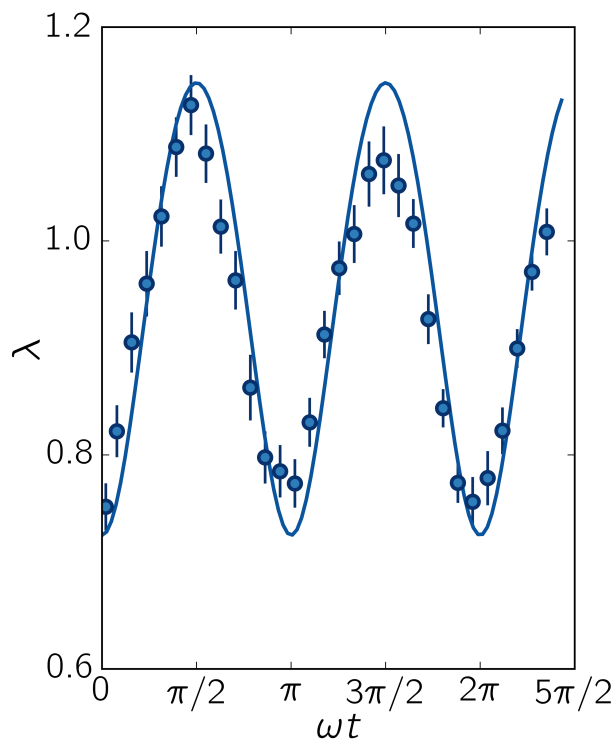
$$\lambda(t) = \left[ \delta^2 \cos^2(\omega t) + \left( \frac{\mu}{\delta} \right)^2 \sin^2(\omega t) \right]^{-1/2}$$

$$\tan(\omega \tau) = \frac{\mu}{\delta^2} \tan(\omega t)$$

Three experiments to vary

$\delta$ ,  $\mu$ ,  $\zeta$  independently:

arXiv:1903.04528



# Outline of this talk

Introduction

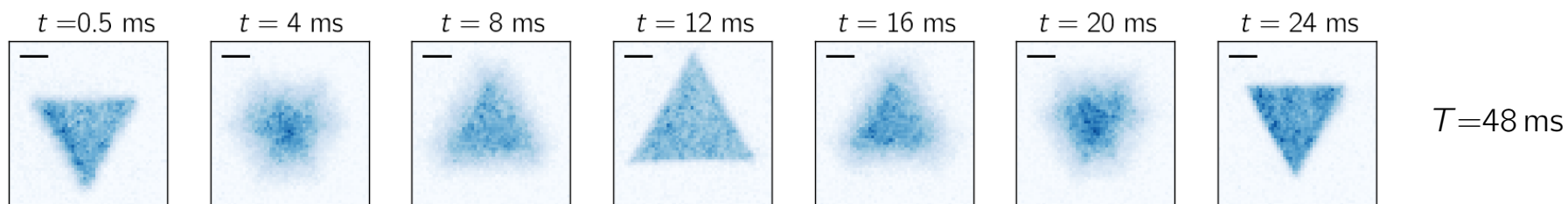
- ① Symmetry group of the 2D Bose gas
- ② Linking evolutions of different clouds
- ③ **Breathers of the Gross–Pitaevskii equation**

Conclusion and outlook

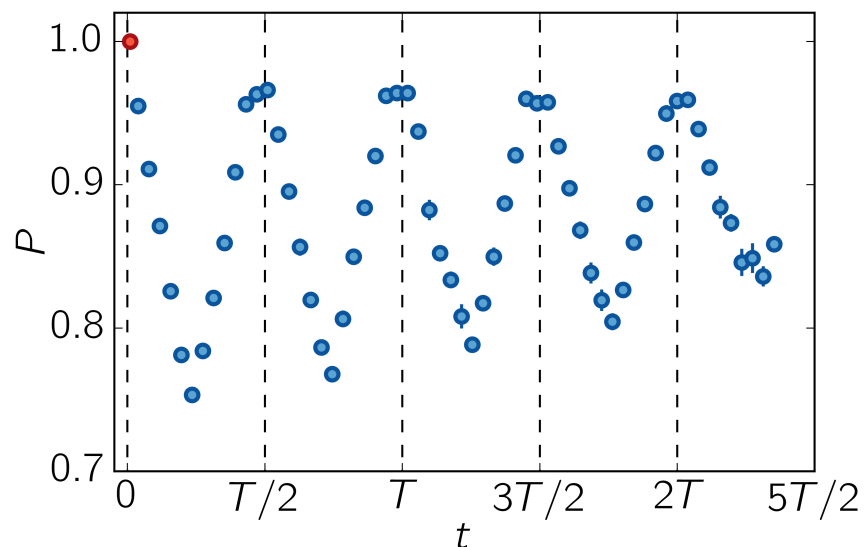


# Periodic evolution of a wave function

- Wave function  $\psi$  in a harmonic trap: not necessarily periodic.
- A triangle seems to be periodic, period  $T/2 \rightarrow$  breather of the GPE

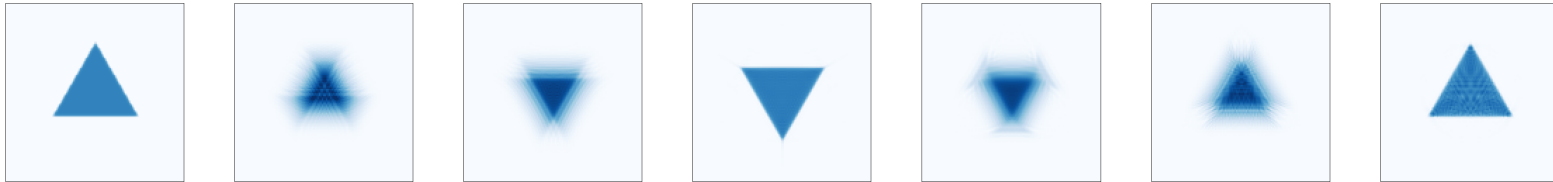


- Overlap of the images: 
$$P = \frac{\int d^2r n(\mathbf{r},0) n(\mathbf{r},t)}{\|n(\mathbf{r},0)\| \|n(\mathbf{r},t)\|}$$

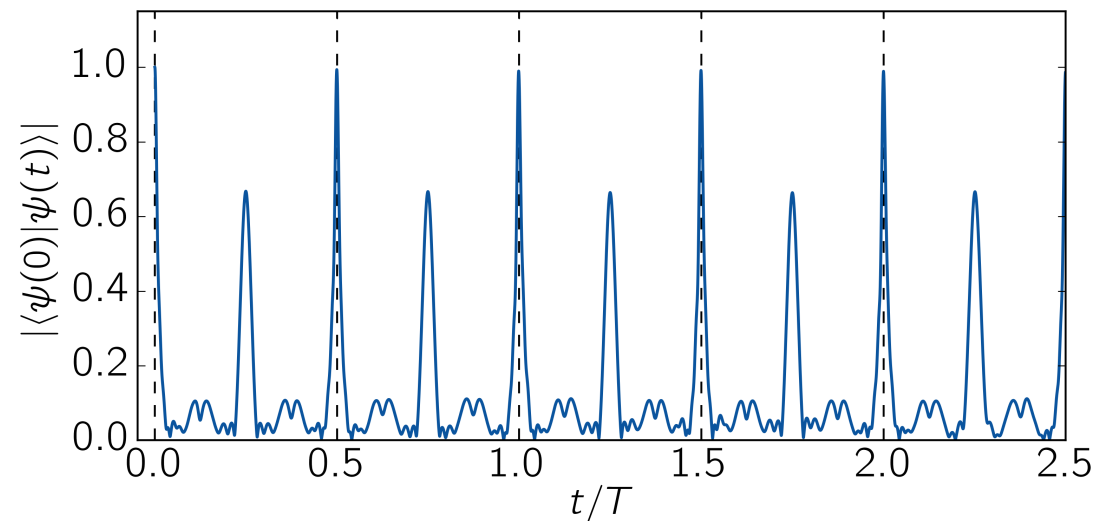


# Numerical simulations

- Grid size  $N_s \times N_s$  with  $N_s = 128$ , healing length  $\xi = 0.5\ell$



- Overlap between the wave functions:  $|\langle \psi(0) | \psi(t) \rangle|$

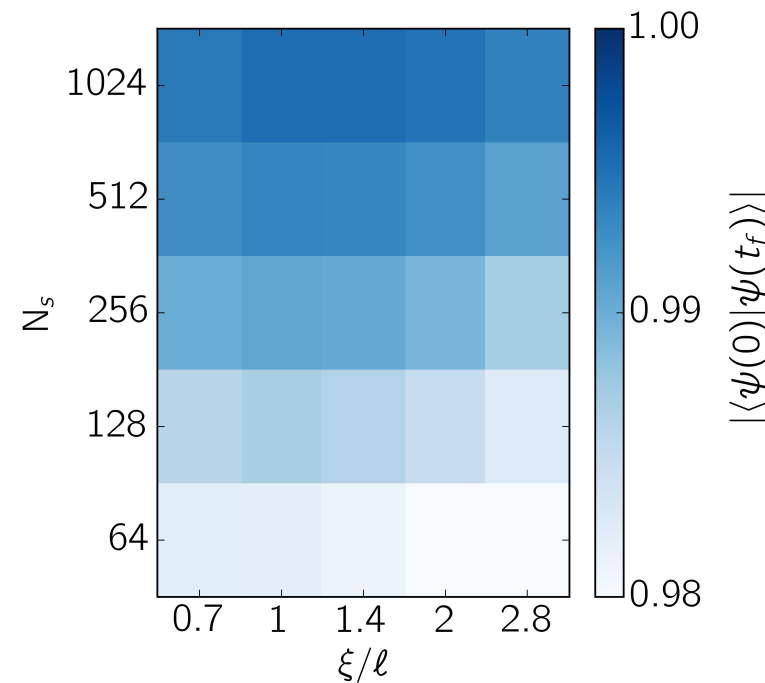


- No analytical proof of periodicity
- Finite-size scaling on the numerical simulations

# Finite-size scaling

- Two choices to improve a simulation:
  - (a) Increase  $N_s$  (reduce pixel size  $\ell$ ) and keep  $\xi$  the same
  - (b) Reduce  $\xi$ : more in the Thomas-Fermi regime

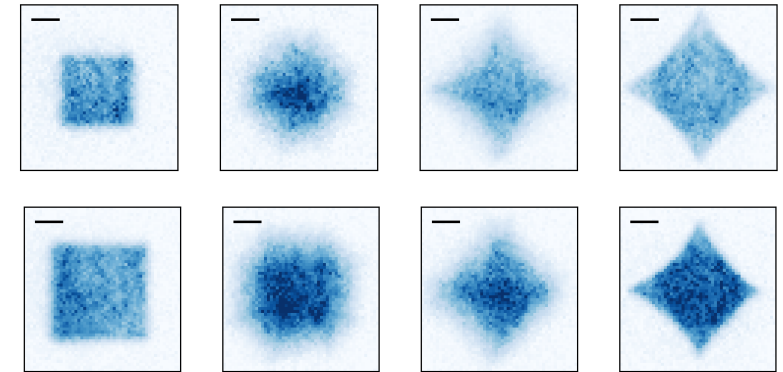
Overlap between  $\psi(0)$  and  $\psi(T/2)$ :



- Best overlap here:  $> 0.995$
- Other shapes? Square, pentagon, hexagon, star, etc. not periodic  
But a disk-shaped cloud also seems to evolve periodically.

# Conclusion

- Dynamical symmetries of the 2D Bose gas:  
Vary  $L$ ,  $N$ ,  $\tilde{g}$ ,  $\omega \rightarrow$  same universal dynamics



- Breathers of the GPE: this dynamics is periodic

- **Open questions:**

- $\rightarrow$  Periodic only with contact interactions?

- $\rightarrow$  Resilient to quantum anomaly (increase  $\tilde{g}$ )?

- Olshanii et al. PRL, **105**, 095302 (2010)

- Holten et al. PRL, **121**, 120401 (2018)

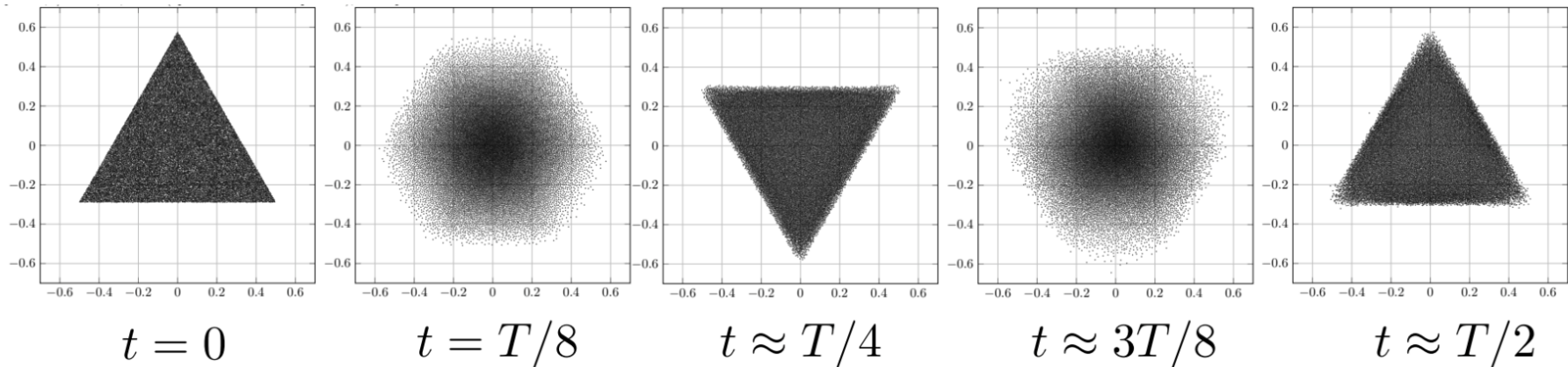
- $\rightarrow$  Effect of temperature?

- $\rightarrow$  Breathers for unitary Fermi gas in 3D?

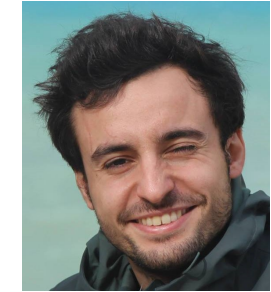
- $\rightarrow$  ...

# Conclusion

- Classical particles with interaction  $\propto 1/r^2$
- Numerical simulation with 20 000 particles
- Initial condition: uniform filling of a triangle,  $\mathbf{v}_i = 0$



# The team

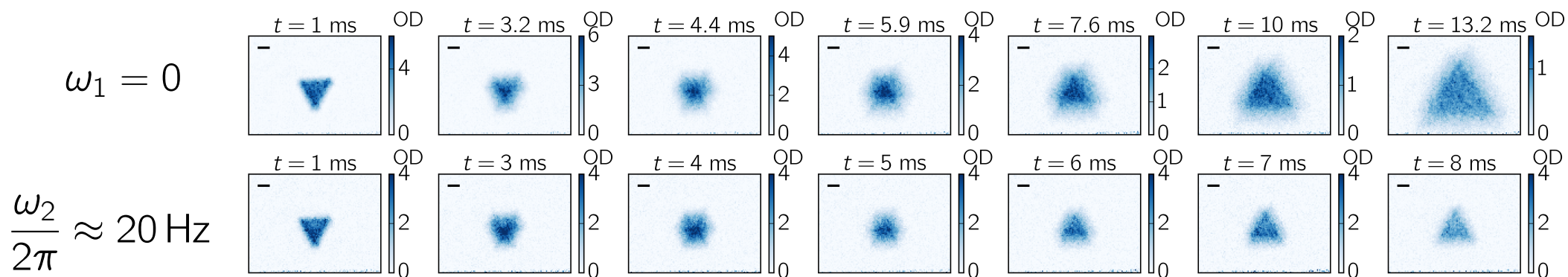


Brice

**Thank you for your attention!**

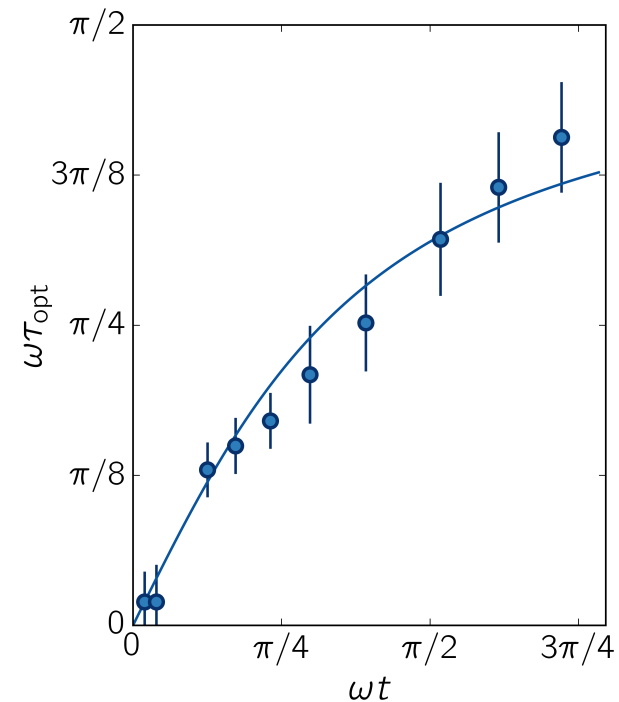
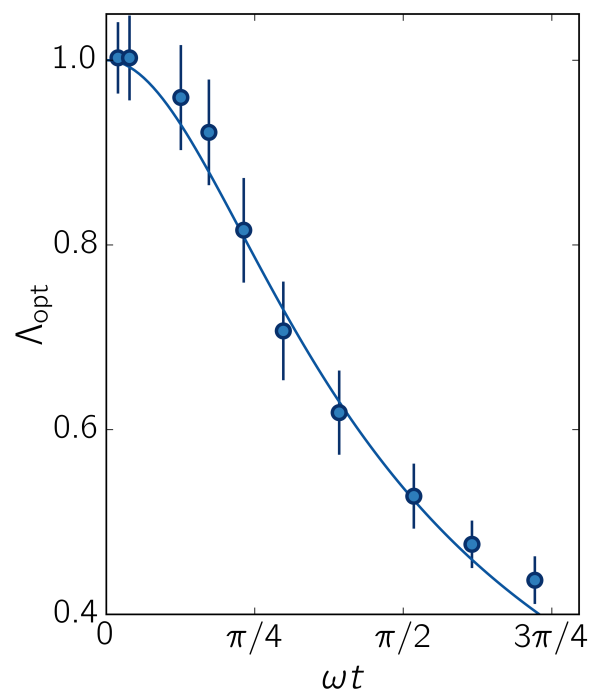


# Evolution with different $\omega$ 's



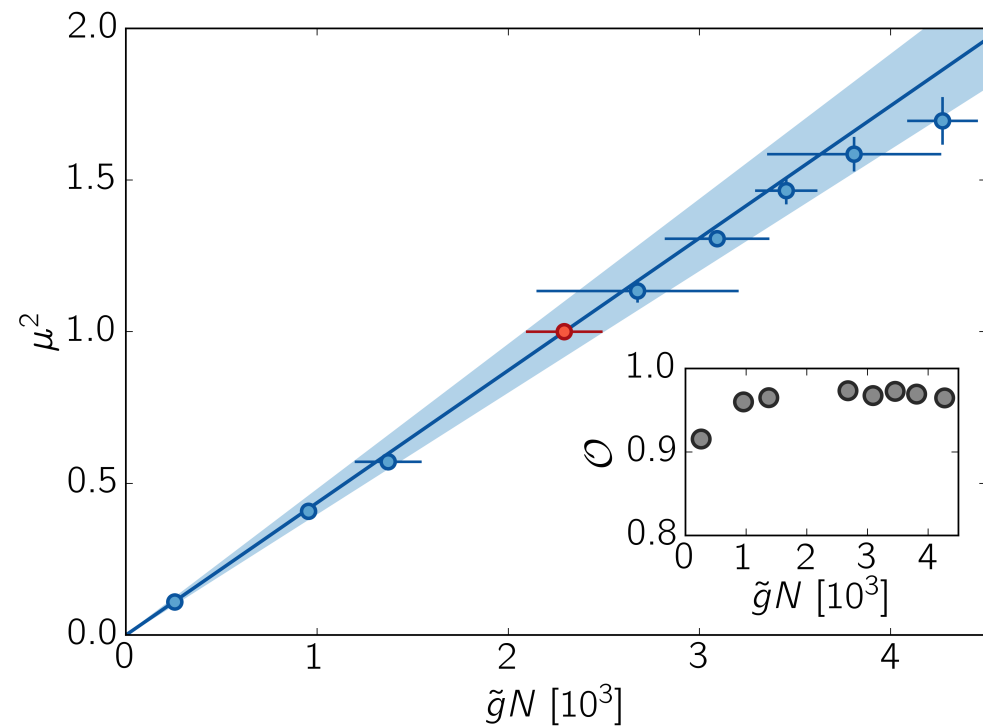
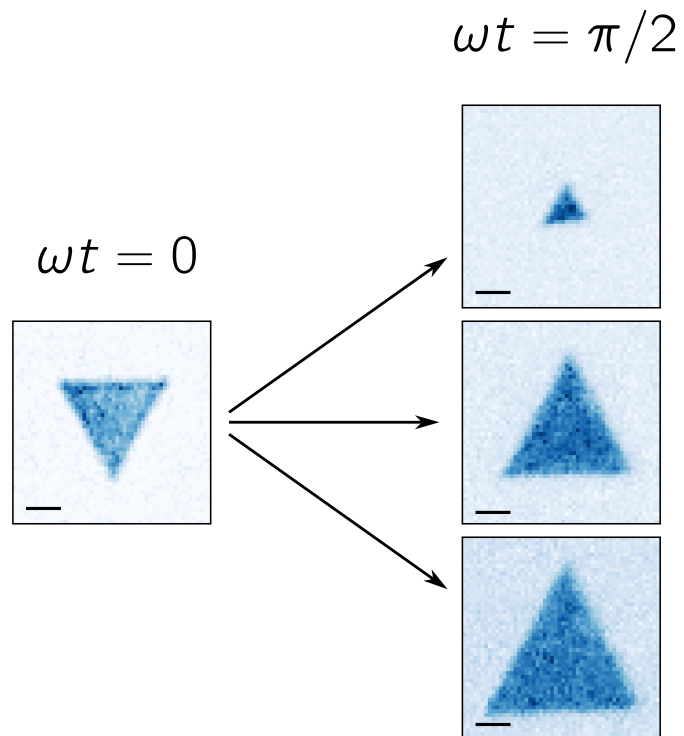
$$\Lambda(t) = (1 + \omega_2^2 t^2)^{-1/2}$$

$$\omega_2 \tau(t) = \arctan(\omega_2 t)$$





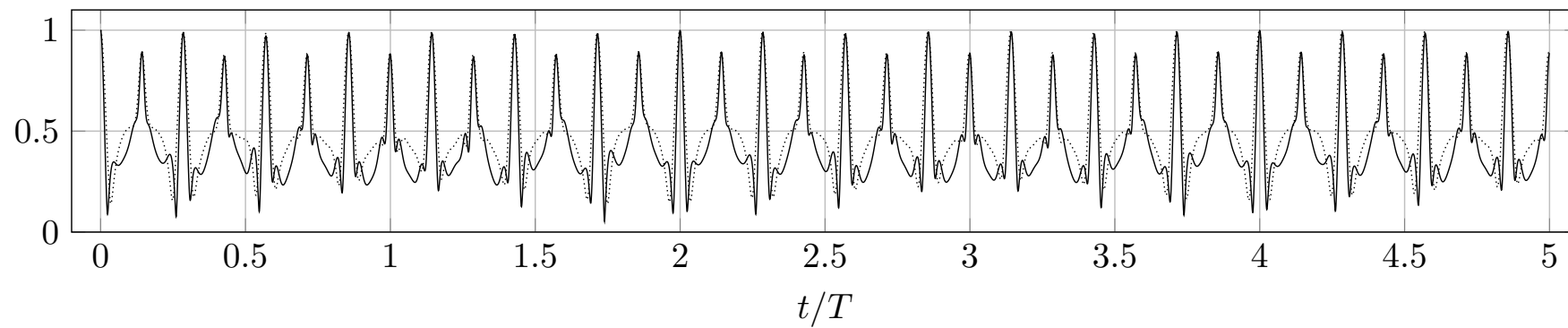
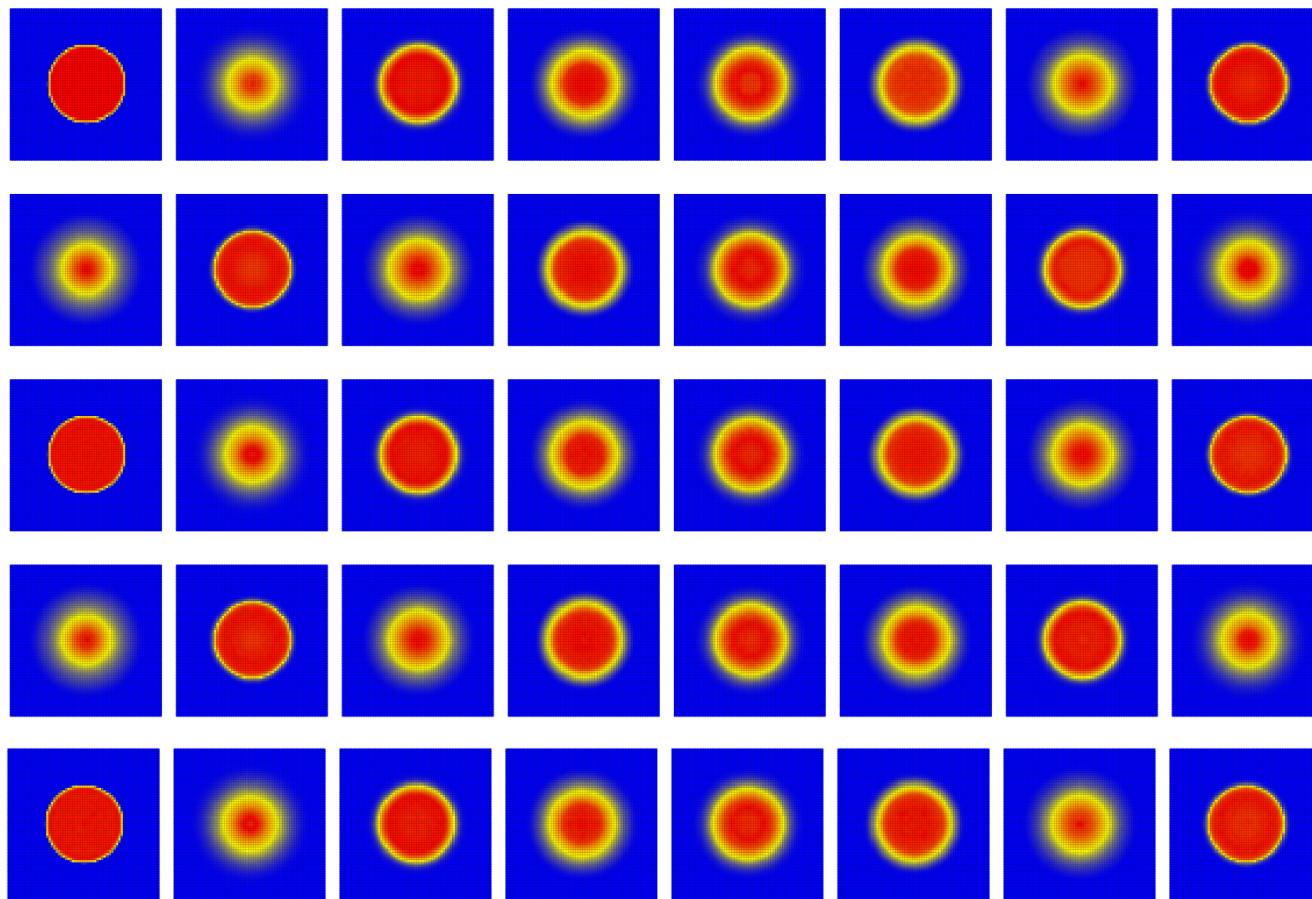
# Evolution with different $\tilde{g}N$ 's



$$\lambda(t) = \left[ \delta^2 \cos^2(\omega_2 t) + \left( \frac{\mu}{\delta \zeta} \right)^2 \sin^2(\omega_2 t) \right]^{-1/2}$$

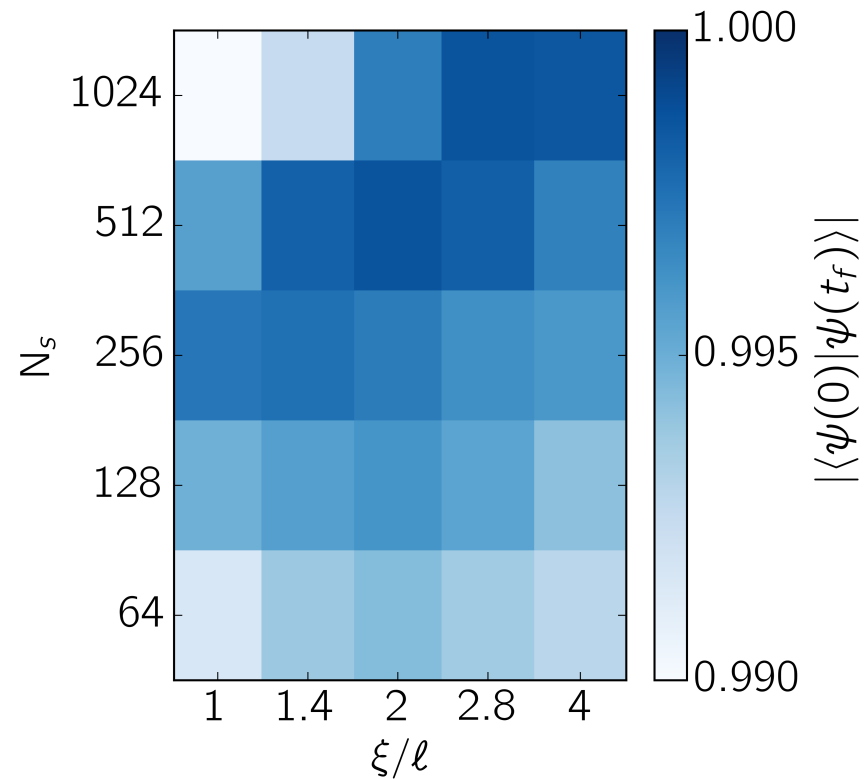
$$\lambda(t = \pi/2\omega) = \frac{\zeta \delta}{\mu}$$

# Breathers: disk



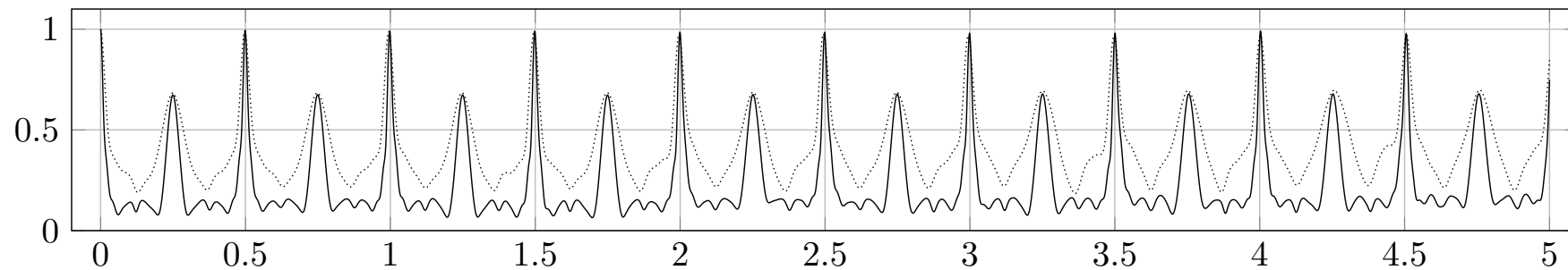
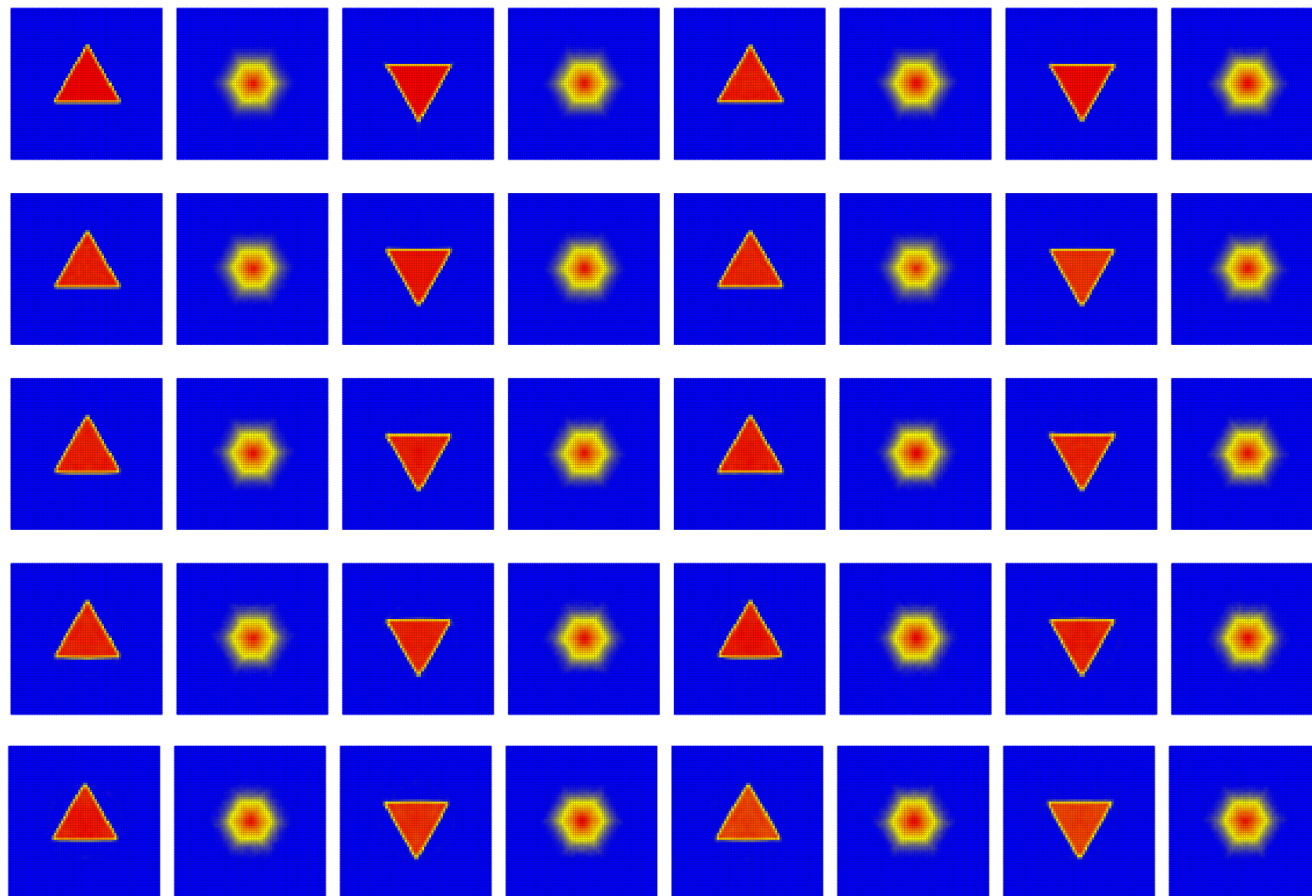
# Finite-size scaling: disk

Overlap between  $\psi(t = 0)$  and  $\psi(t = 4\pi/\omega)$ .

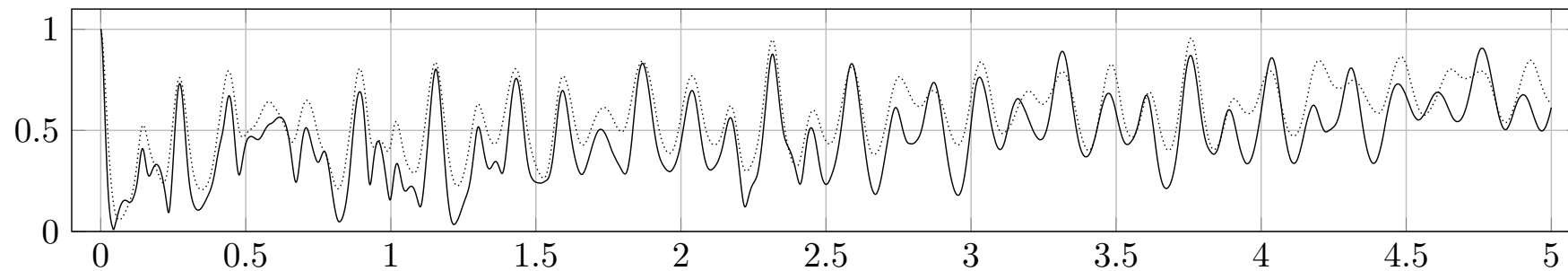
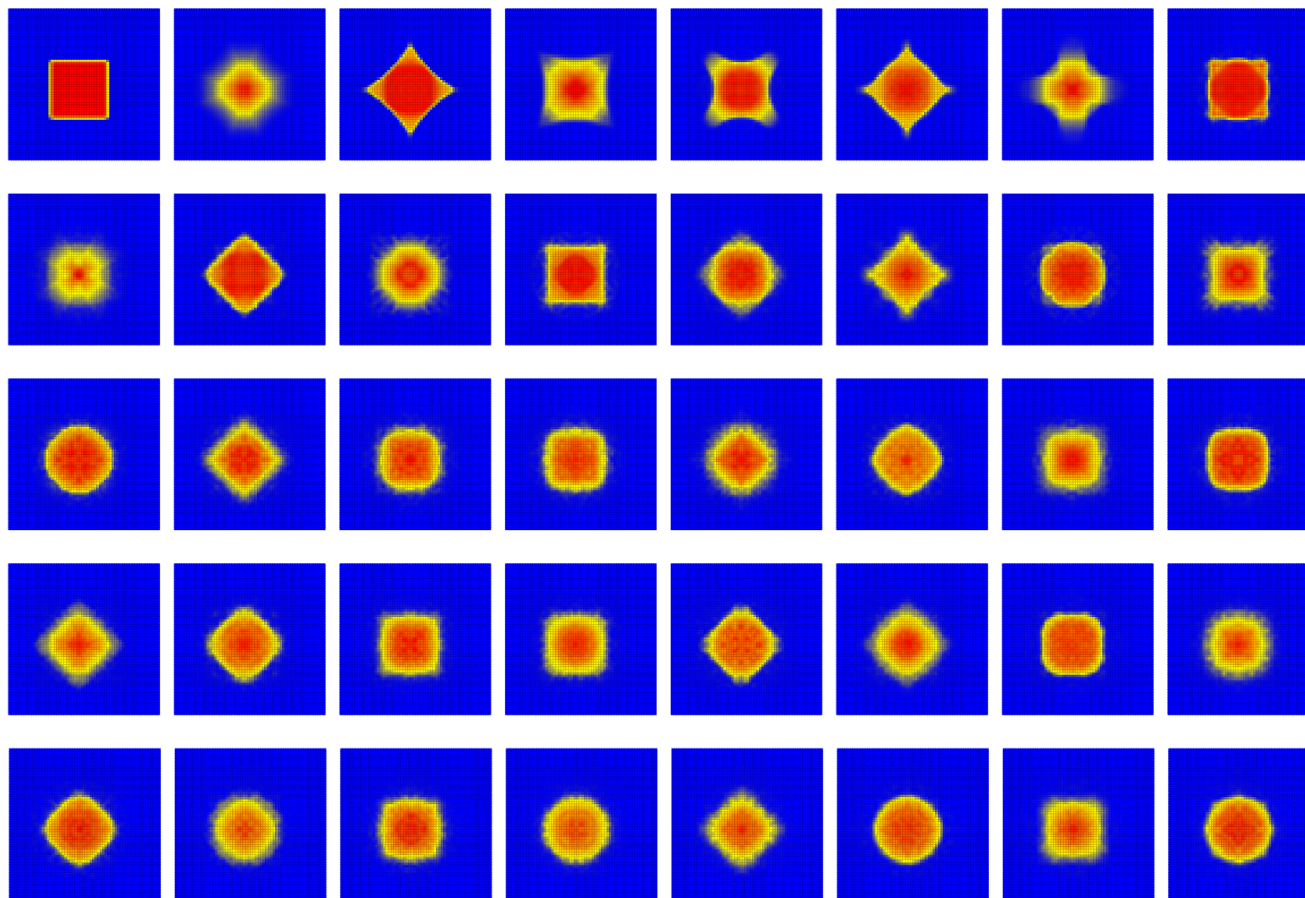


Best overlap  $> 0.998$ .

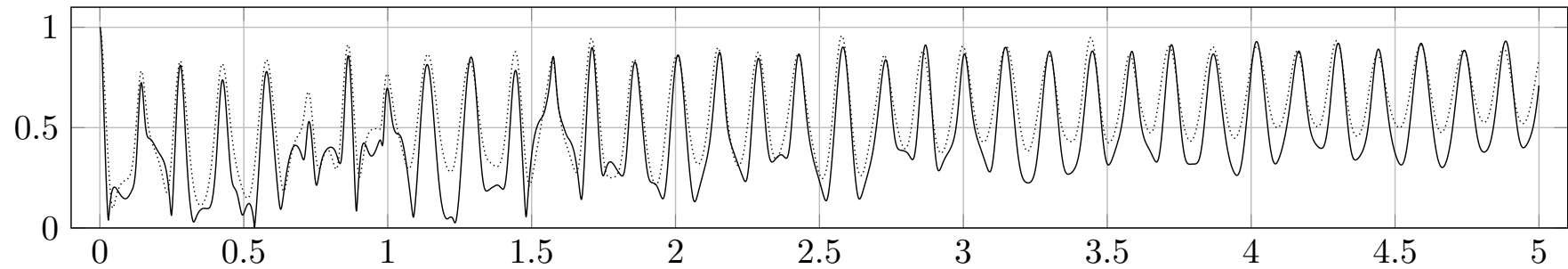
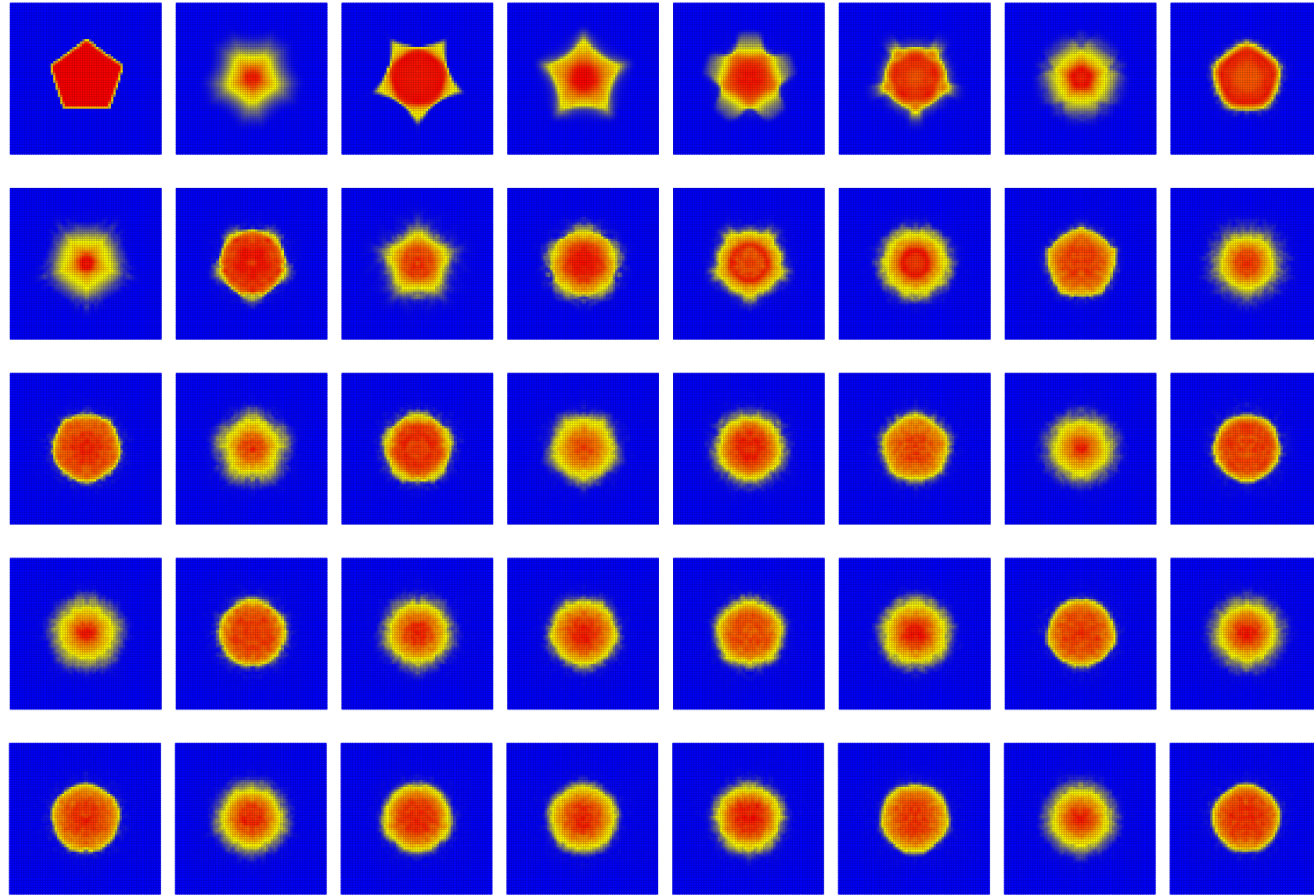
# Numerical simulations: Triangle



# Numerical simulations: Square



# Numerical simulations: Pentagon



# Numerical simulations: Hexagon

