Decay of semilinear damped wave equations: cases without geometric control condition

Romain JOLY

Université Grenoble Alpes

Joint work with Camille Laurent, CNRS-Paris VI

Benasque 2019

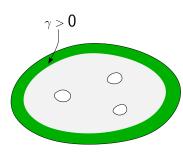
Benasque: a fruitful and friendly atmosphere

Benasque 2015



The semilinear damped wave equation

$$\partial_{tt}^2 u + \gamma(x) \partial_t u = \Delta u - f(x, u)$$



- Ω is a smooth compact manifold of dimension d = 2 with Dirichlet boundary conditions.
- the damping γ is in $\mathbb{L}^{\infty}(\Omega)$, $\gamma(x) \geq 0$
- f is smooth and of degree p $|f(x,u)| + |f'_{x}(x,u)| \le C(1+|u|)^{p}$ $|f'_{u}(x,u)| \le C(1+|u|)^{p-1}$
- f is of the sign of u:

$$f(x,u)u \geq 0$$

Notations

$$\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x, u)$$

Set $X = H_0^1(\Omega) \times L^2(\Omega)$ and

$$U = \begin{pmatrix} u \\ \partial_t u \end{pmatrix} \qquad A = \begin{pmatrix} 0 & Id \\ \Delta & -\gamma(x) \end{pmatrix} \qquad F(U) = \begin{pmatrix} 0 \\ -f(x, u) \end{pmatrix}$$

- $\Rightarrow e^{At}$ is a dissipative semigroup on X.
- \Rightarrow Since f is of degree $p < \infty$ and Ω is of dimension d = 2, $F: X \longrightarrow X$ is defined and Lipschitz on the bounded sets.

We consider in X the equation

$$\partial_t U = AU + F(U)$$
 $U(t=0) = U_0 \in X$

The gradient dynamics

Set $V(x, u) = \int_0^u f(x, \xi) d\xi$. The energy

$$\mathcal{E}(U) = \int_{\Omega} \frac{1}{2} (|\nabla u|^2 + |\partial_t u|^2) + V(x, u) dx$$

is non-increasing along the trajectories since

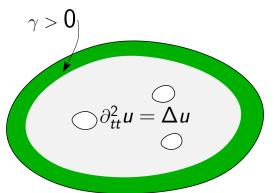
$$\partial_t \mathcal{E}(U(t)) = -\int_{\Omega} \gamma(x) |\partial_t u|^2 dx$$

⇒ Global existence of solutions

Motivations

The linear equation is dissipative and any solution goes to zero

$$\|e^{At}U\|_X \xrightarrow[t \to +\infty]{} 0.$$

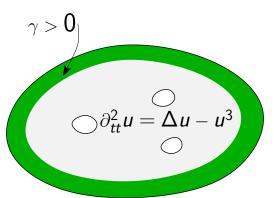


Motivations

The linear equation is **dissipative** and any solution goes to zero

$$\|e^{At}U\|_X \xrightarrow[t \to +\infty]{} 0.$$

Do we still have stabilization of the nonlinear problem? At which rate?

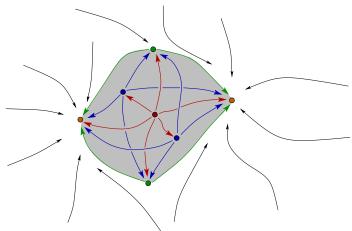


Motivations

Also possible to consider convergence to less trivial dynamics by assuming that f is only asymptotically of the sign of u

$$\forall |u| \geq R$$
, $f(x,u)u \geq 0$.

Existence of global attractor with gradient structure?

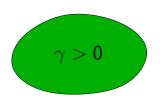


- A historic result
- 2 A standard extension
- The disk with two holes
- The disk with three holes
- Conclusion

$$\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x, u)$$

Assume:

- $\gamma(x) \ge \alpha > 0$ in Ω
- f is of degree $p < \infty$
- $f(x, u)u \geq 0$



Theorem – J.K. Hale (1985) and A. Haraux (1985)

With the above assumptions, any solution u(t) of the damped wave equation converges to 0 in $X = H_0^1(\Omega) \times L^2(\Omega)$. Moreover, the convergence is uniform in bounded sets of X.

Step 1: the trajectories are bounded.

If $f(x, u)u \ge 0$ and f is of degree p, then the energy is well defined, non-negative and bounded on bounded sets.

$$\frac{1}{2}\|U\|_X^2 + \min V \leq \int_{\Omega} \frac{1}{2} (|\nabla u|^2 + |\partial_t u|^2) + V(x, u) \, \mathrm{d}x \leq K(\|U\|_X) .$$

Since \mathcal{E} is non-increasing, the trajectories of bounded sets are bounded.

Step 2: the asymptotic compactness.

The linear semigroup is stabilized:

$$\forall t \geq 0 , \|e^{At}\|_{\mathcal{L}(X)} \leq Me^{-\lambda t}$$

Moreover, if f is of degree p, then

$$U(t) = e^{At} U_0 + \int_0^t e^{A(t-s)} F(U(s)) ds$$

 \Rightarrow the bounded sets admits compact $\omega-$ limit sets.

Step 3: a unique continuation property.

It is sufficient to show that the ω -limit sets consists of equilibrium points. By Lasalle's principle, the trajectories $U(t)=(u,\partial_t u)$ in the ω -limit sets have constant energy. So we have

$$\partial_t \mathcal{E}(U(t)) = -\int_{\Omega} \gamma(x) |\partial_t u|^2 dx = 0.$$

Since $\gamma(x) \ge \alpha > 0$, we have $\partial_t u \equiv 0$ and thus u is an equilibrium point. Due to the sign assumption, we finally obtain $u \equiv 0$.

- Asymptotic compactness ⇔ high frequencies are not really modified by the nonlinearity
- Unique continuation ⇔ classify low-frequency solutions

Key arguments where we use γ positive:

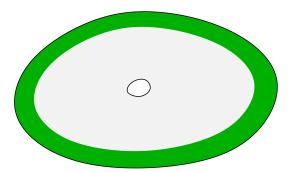
- ② if $\mathcal{E}(U(t))$ is constant, then $\int \gamma(x)|u_t|^2=0$ and u(t) is constant.

A standard extension

- A historic result
- 2 A standard extension
- The disk with two holes
- The disk with three holes
- Conclusion

A standard extension

What happens when $\gamma(x)$ may vanish?



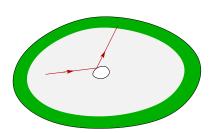
The decay of the linear semigroup

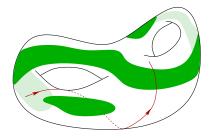
Theorem – J. Rauch and M. Taylor (1974) C. Bardos, G. Lebeau and J. Rauch (1992)

$$\|e^{At}\|_{\mathcal{L}(X)} \leq Me^{-\lambda t}$$

$$\iff$$

Any long enough geodesic meets the support of the damping γ





The unique continuation property

If $U_{\infty}(t)$ belongs to an $\omega-$ limit set,

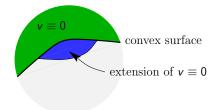
$$\partial_t \mathcal{E}(U_\infty(t)) = -\int_\Omega \gamma(x) |\partial_t u_\infty|^2 dx = 0.$$

So $v(t) = \partial_t u_\infty(t)$ vanishes in ω the support of γ . Thus, we have

$$v\equiv 0 \text{ in } \omega imes\mathbb{R} \quad \text{ and } \quad \partial^2_{tt}v=\Delta v-f_u'(x,u_\infty(x,t))v \ .$$

To conclude that $v \equiv 0$ everywhere, we need to use a unique continuation property.

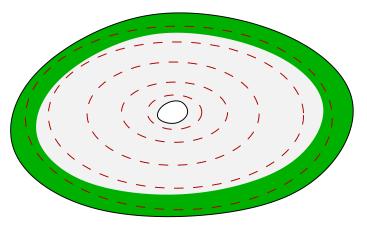
Basically, we may extend the zone where $v \equiv 0$ through **convex surfaces**.



[N. Lerner and L. Robbiano, 1985], [L. Hörmander, 1985], [Tataru, 1996]

The unique continuation property

The stabilization holds for the domain with zero or one hole

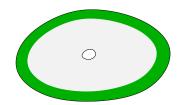


A classic result

$$\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x, u)$$

Assume:

- ullet Ω is a two dimensional convex compact domain with or without a convex hole
- $\gamma(x) \ge \alpha > 0$ in a neighborhood of the exterior boundary of Ω .
- $f(x, u)u \ge 0$ and f of degree p



Theorem

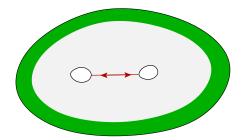
With the above assumptions, the semilinear damped wave equation is stabilized. More precisely, there exists $\lambda>0$ such that, for any R>0, there exists M_R such that

$$\begin{aligned} \|(u_0,u_1)\|_{H_0^1\times L^2} &\leq R \\ \Longrightarrow \|(u,\partial_t u)(t)\|_{H_0^1\times L^2} &\leq M_R e^{-\lambda t} \xrightarrow[t\longrightarrow +\infty]{} 0. \end{aligned}$$

- A historic result
- 2 A standard extension
- The disk with two holes
- 4 The disk with three holes
- Conclusion

Without the geometric control condition

In some cases, the geometric control condition does not hold, but very few geodesics miss the support of the damping.



$$\|e^{At}U_0\|_{H^1\times L^2} \leq Me^{-\lambda t^{1/3}}\|U_0\|_{H^2\times H^1}$$

[N. Burq, 1993]

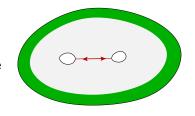
[N. Burq and M. Zworski, 2004]

[R.J. and C. Laurent, 2018]

$$\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x, u)$$

Assume:

- Ω is a convex compact domain of dimension 2 with two convex holes
- $\gamma(x) \ge \alpha > 0$ in a neighborhood of the exterior boundary of Ω .
- $f(x, u)u \ge 0$ and f of degree $p < \infty$



Theorem – R.J. and C. Laurent (2018)

With the above assumptions, the semilinear damped wave equation is semi-stabilized. More precisely, $\|(u,\partial_t u)(t)\|_{H_0^1 \times L^2} \xrightarrow[t \to +\infty]{} 0$.

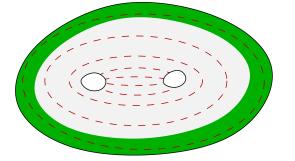
Moreover, there exists $\tilde{\lambda}$ such that, for any R and $\sigma \in (0,1]$, there exists $C_{R,\sigma}$ such that

$$\|(u_0, u_1)\|_{H^{1+\sigma} \times H^{\sigma}} \le R \implies \|(u, \partial_t u)(t)\|_{H^1_0 \times L^2} \le C_{R,\sigma} e^{-\sigma \tilde{\lambda} t^{1/3}}$$

Main arguments:

$$\|e^{At}U_0\|_{H^1 \times L^2} \le Me^{-\lambda t^{1/3}} \|U_0\|_{H^2 \times H^1}$$

$$U(t) = e^{At}U_0 + \int_0^t e^{A(t-s)} F(U(s)) ds$$



The basic idea to obtain the estimate is the following.

Assume $f(u) = u^3$, for u small in $H^{1+\sigma}(\Omega)$, we have

$$||f(u)||_{H^1} \le \delta ||u||_{H^{1+\sigma}}$$
 with δ small .

Thus,

$$e^{\sigma \lambda t^{1/3}} U(t) = e^{\sigma \lambda t^{1/3}} e^{At} U_0 + e^{\sigma \lambda t^{1/3}} \int_0^t e^{A(t-s)} F(U(s)) ds$$

and

$$\max_{t \in [0,T]} \|e^{\sigma \lambda t^{1/3}} U(t)\|_{X} \\
\leq C + \delta \max_{s \in [0,T]} \|e^{\sigma \lambda s^{1/3}} U(s)\|_{X} \int_{0}^{T} e^{\sigma \lambda (T^{1/3} - s^{1/3})} e^{-\lambda (T - s)^{1/3}} ds .$$

The disk with three holes

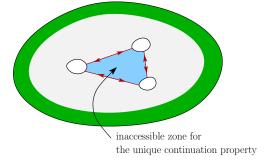
- A historic result
- 2 A standard extension
- The disk with two holes
- 4 The disk with three holes
- Conclusion

The disk with three holes

If there are three holes or more, with additional technical assumptions, we still have the decay

$$\|e^{At}U_0\|_{H^1\times L^2} \le Me^{-\lambda t^{1/3}}\|U_0\|_{H^2\times H^1}$$

But the unique continuation property of Lerner-Robbiano-Hörmander-Tataru cannot be used:



An analytic unique continuation property

Theorem – L. Robbiano and C. Zuily (1998) L. Hörmander (1997)

Assume $\omega \neq \emptyset$ and $v(t) = \partial_t u(t)$ solves

$$v \equiv 0$$
 in $\omega imes \mathbb{R}$ and $\partial^2_{tt} v = \Delta v - f_u'(x,u(x,t))v$.

Assume moreover that $t \mapsto f'_u(x, u(x, t))$ is analytic then $v \equiv 0$ everywhere.

[J.K. Hale and G. Raugel, 2003] let us hope that if f(x, u) is analytic in u, then a function u in the attractor should be analytic in time and thus $f'_u(x, u(x, t))$ is also analytic.

An analytic unique continuation property

In the proofs of [J.K. Hale and G. Raugel, 2003], a global solution u is split between the **low-frequencies** P_nu and the high-frequencies Q_nu . It is used that

$$\|e^{At}U\|_X \le Me^{-\lambda t}\|U\|_X \implies \|e^{Q_nAQ_nt}Q_nU\|_X \le Ne^{-\mu t}\|Q_nU\|_X$$
.

In our case, we would like to obtain

$$\|e^{At}U\|_X \leq Me^{-\lambda t^{1/3}}\|U\|_{D(A)} \implies \|e^{Q_nAQ_nt}Q_nU\|_X \leq h(t)\|Q_nU\|_{D(A)}.$$

 \implies we adapt the ideas of [J.K. Hale and G. Raugel, 2003] but several technical problems have to be overcome.

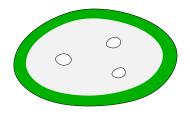
[C.J.K. Batty and Th. Duyckaerts, 2008], [A. Borichev and Y. Tomilov, 2010], [N. Anantharaman and M. Léautaud, 2014]

The disk with three holes

$$\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x, u)$$

Assume:

- $\bullet \ \Omega$ is as opposite and the holes are not aligned and small enough
- f(x, u) is analytic in u
- $f(x, u)u \ge 0$ and f of degree $p < \infty$



Theorem – R.J. and C. Laurent (2018)

With the above assumptions, the semilinear damped wave equation is semi-stabilized. More precisely, $\|(u, \partial_t u)(t)\|_{H_0^1 \times L^2} \xrightarrow{t \longrightarrow +\infty} 0$.

Moreover, there exists $\tilde{\lambda}$ such that, for any R and $\sigma \in (0,1]$, there exists $C_{R,\sigma}$ such that

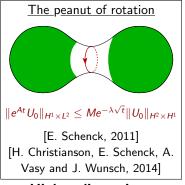
$$\|(u_0, u_1)\|_{H^{1+\sigma} \times H^{\sigma}} \le R \implies \|(u, \partial_t u)(t)\|_{H^1_0 \times L^2} \le C_{R,\sigma} e^{-\sigma \tilde{\lambda} t^{1/3}}$$

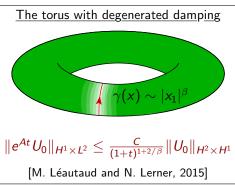
Conclusion

- A historic result
- 2 A standard extension
- The disk with two holes
- The disk with three holes
- Conclusion

Additional results

Other geometries are possible





Higher dimension

In dimension d=3, assume that f is Sobolev-subcritical, that is of degree p with p<3. It should also be possible to go to f energy-subcritical, that is of degree p with p<5 by using Strichartz estimates, see [B. Dehman, G. Lebeau and E. Zuazua, 2003], [R.J. and C. Laurent, 2013]

Open problem

Global control? At least approximate global controllability?

The problem is the local controllability close to zero.

Open problem

$$U(t) = e^{At} U_0 + \int_0^t e^{A(t-s)} F(U(s)) ds$$

Main open question:

How important is the integrability of the linear decay?

For example, if the linear decay is simply

$$\|e^{At}U_0\|_{H^1\times L^2} \leq \frac{C}{\ln(2+t)}\|U_0\|_{H^2\times H^1}$$

does the asymptotic compactness hold?

Thanks for your attention!

- R.J. and C. Laurent, Semi-uniform decay for some semilinear damped wave equations, preprint 2018.
- R.J. and C. Laurent, A note on the global controllability of the semilinear wave equation, SIAM Journal on Control and Optimization n°52 (2014), pp. 439–450.
- R.J. and C. Laurent, Stabilization for the semilinear wave equation with geometric control condition, Analysis and PDE n°6 (2013), pp. 1089–1119.