

Equations with nonlinear boundary conditions in domains with rapidly varying boundaries

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General problem

Let us consider the following evolution problem in space X_0

$$\begin{cases} x' + A_0x = F_0(x), & t > 0 \\ x(0) = x_0 \in X_0 \end{cases}$$

Let us assume that appropriate conditions are satisfied so that we have global existence of solutions and continuous dependence with respect to initial data. Hence, the equation generates a **dynamical system** in X_0 (phase space):

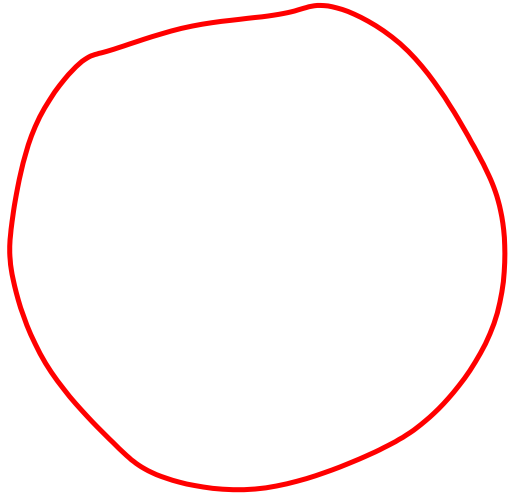
$$\begin{aligned} T_0(t) : X_0 &\longrightarrow X_0 \\ x_0 &\longrightarrow x(t, x_0) \end{aligned}$$

Under certain conditions on the equation we guarantee that T_0 is **dissipative** and **asymptotically compact**

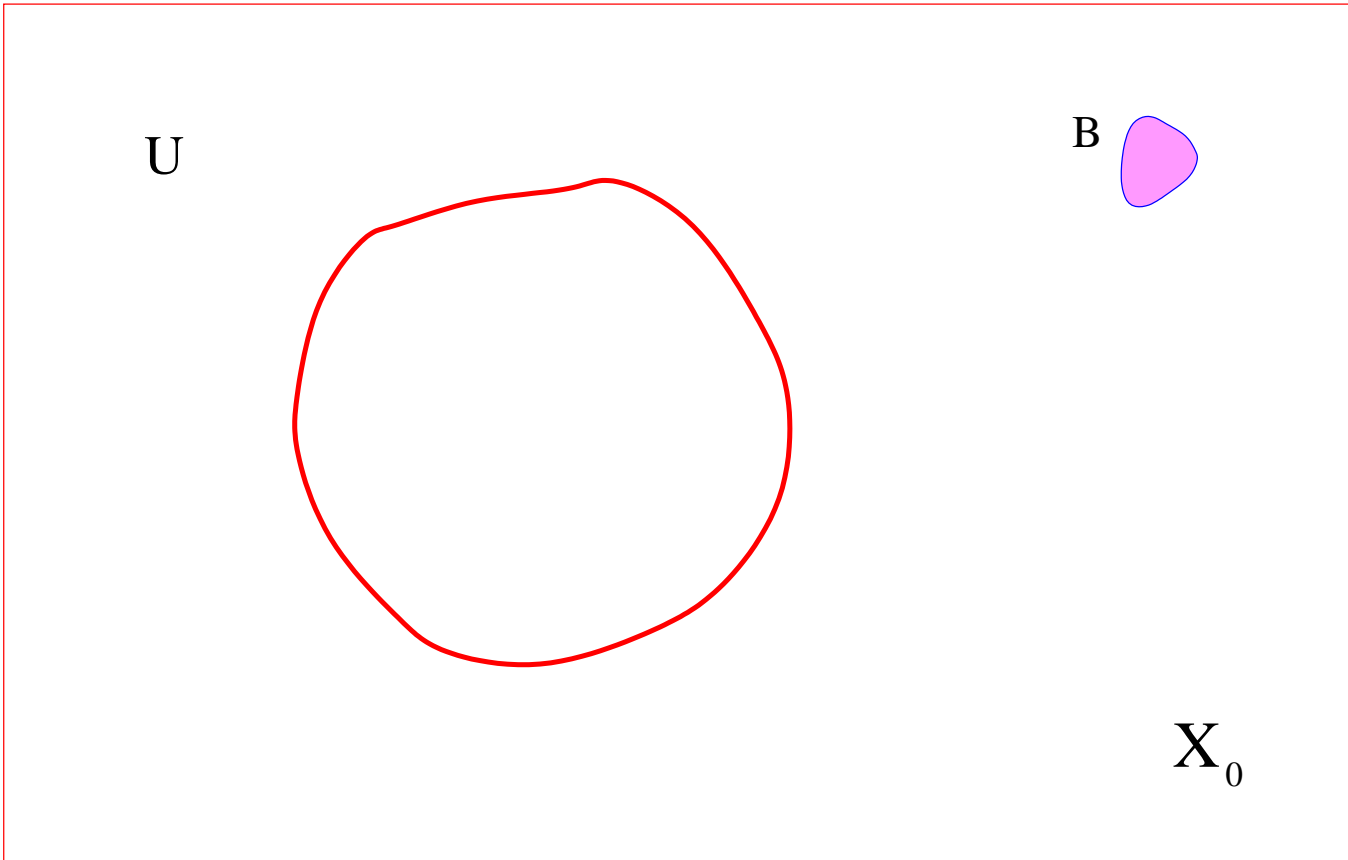
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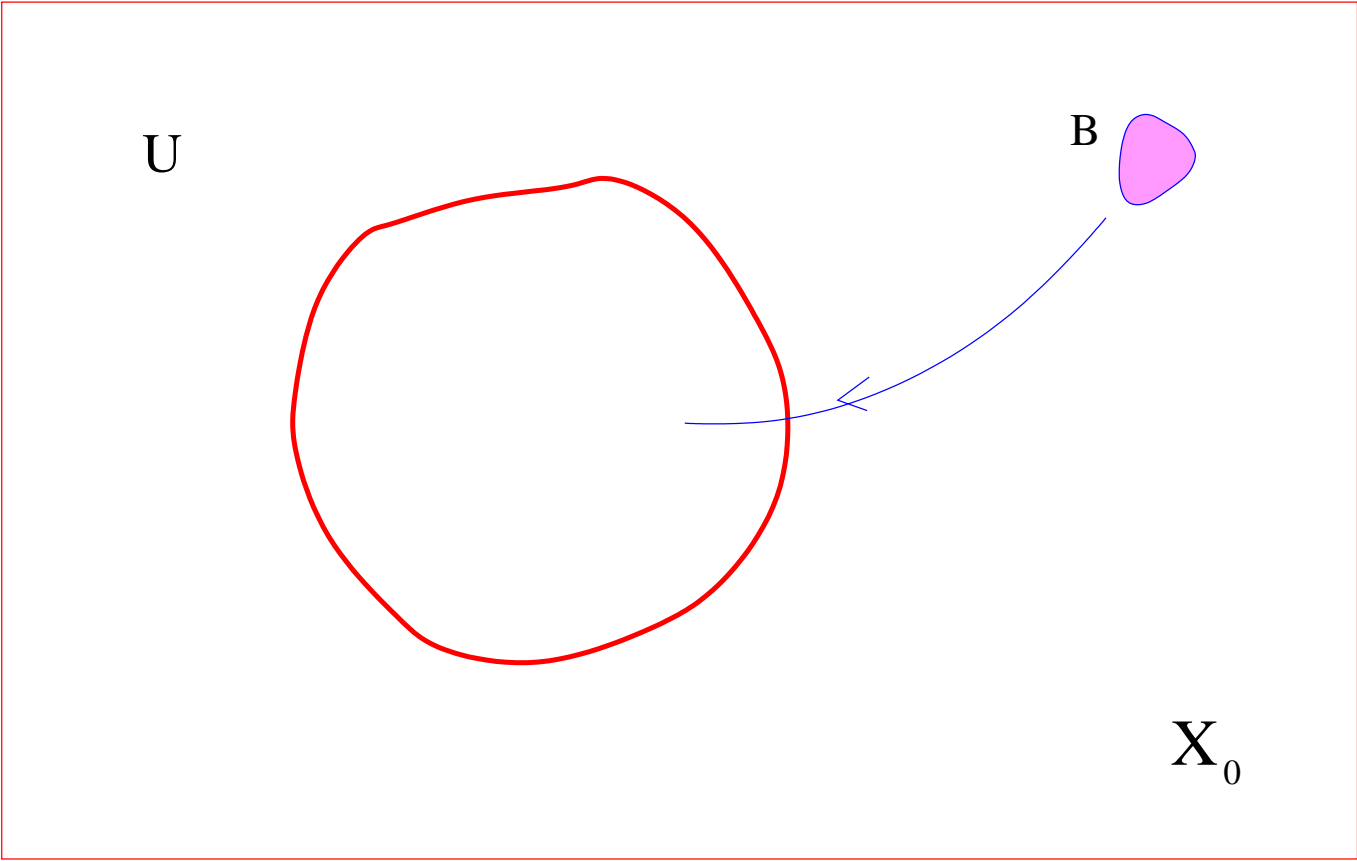
These two conditions guarantee the existence of the **attractor of the equation**, $\mathcal{A}_0 \subset X_0$.

U

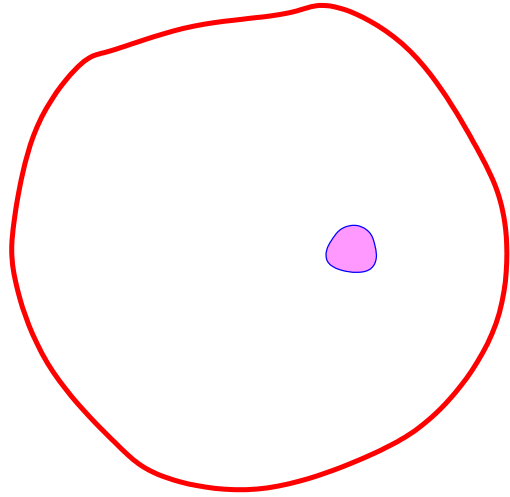


X_0



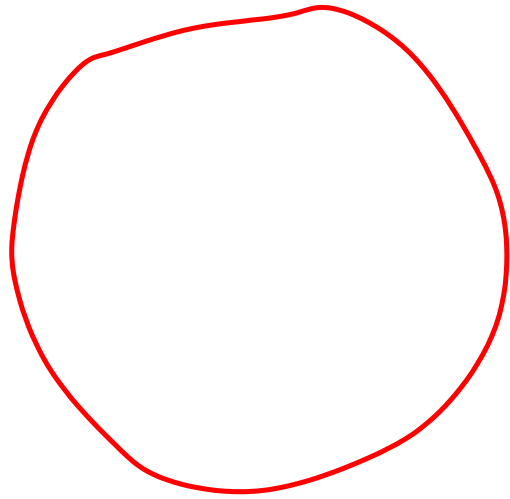


U



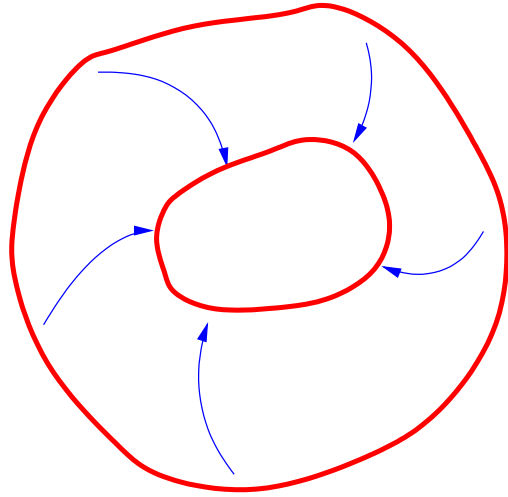
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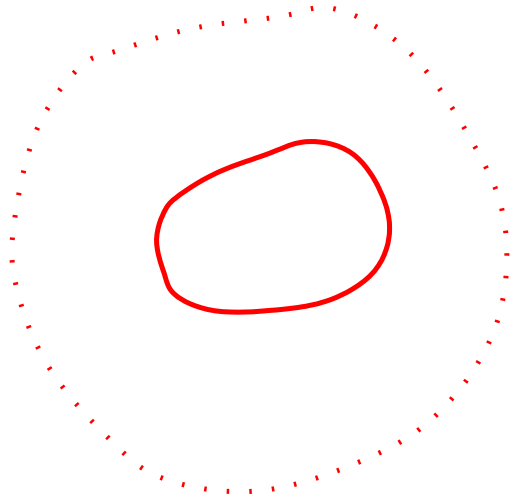
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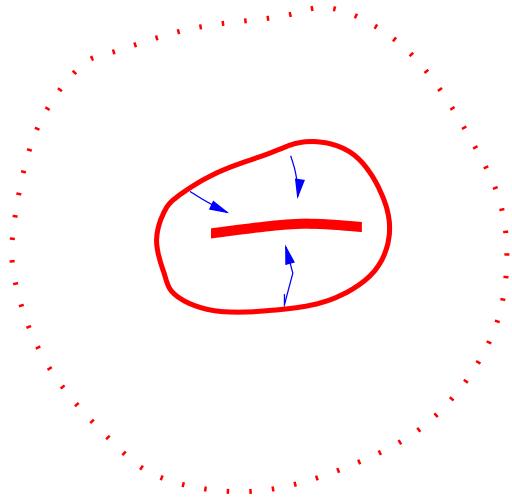
X_0

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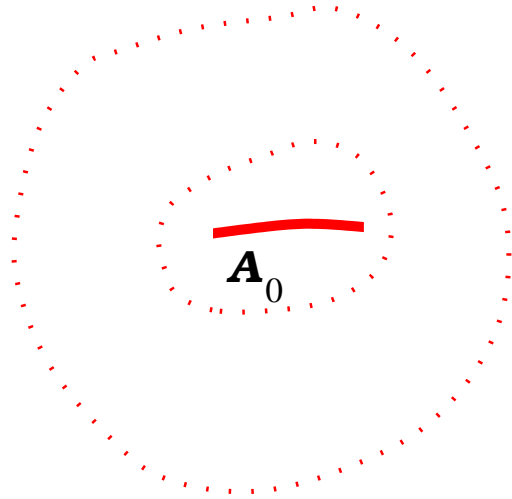
X_0

U



X_0

U



X_0

Attractor: largest compact, invariant set which attracts every bounded set of the phase space.

- It contains all global and bounded orbits: equilibria, periodic orbits, connecting orbits, etc ..
- The dynamics in the attractor contains all the asymptotic dynamics.

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J.K. Hale “Asymptotic behavior of dissipative systems” Mathematical Surveys and Monographs **25** American Mathematical Society, Providence 1988.

A. Babin, M.I. Vishik “Attractors of evolution equations” North Holland, 1992

R. Temam “Infinite dimensional dynamical systems in mechanics and physics”, Springer 1988

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Since we are comparing elements of X_0 with elements of X_ϵ , we need a concept of “closeness” or “convergence” for elements living in different spaces.

If for instance there exists an space Y so that $X_\epsilon \hookrightarrow Y$, $0 \leq \epsilon \leq \epsilon_0$, then we can talk of convergence in Y .

In each case we need to define this concept in a very precise way.

The attractor is a global entity of the dynamical system. Therefore, understanding its structure and its behavior under perturbations is a global problem, which is far away from being resolved in this generality.

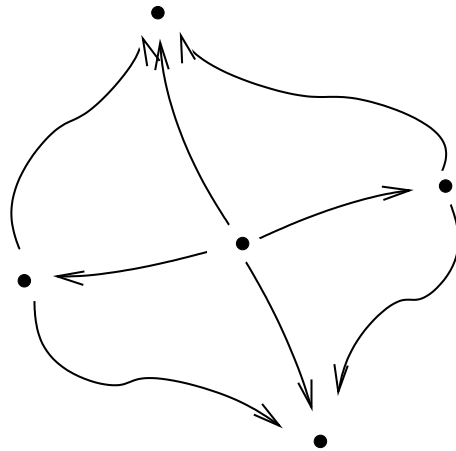
The attractor may have a very complicated structure and it is not easy to analyze its behavior under perturbations.

Nevertheless, if the dynamical system is **gradient**, the attractor structure is simpler. It is made of

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- Connections among equilibria.

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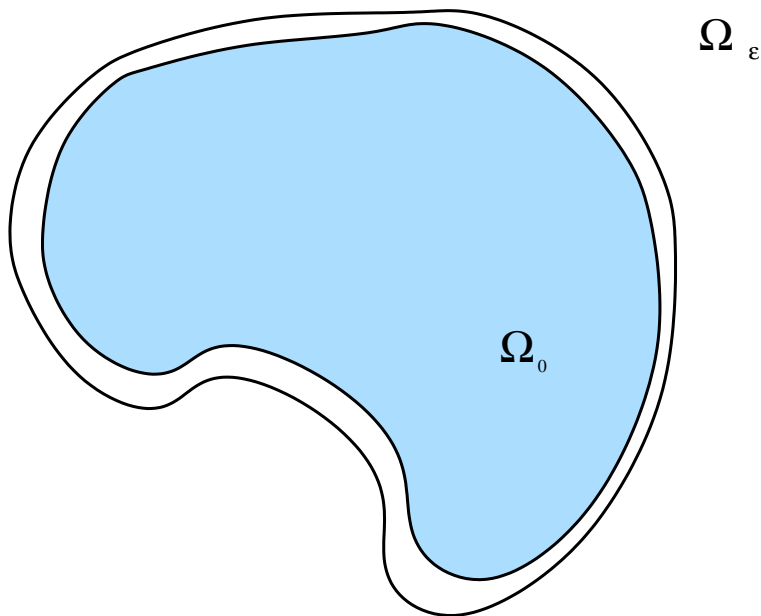
Domain Perturbation

Case 1. General domain perturbation and Neumann boundary conditions.

$$\begin{cases} u_t - \Delta u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

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J.A., A.N. Carvalho “ Spectral Convergence and Nonlinear Dynamics of Reaction-Diffusion Equations under Perturbations of the Domain ” Journal of Differential Equations 199 (2004) pp 143-178

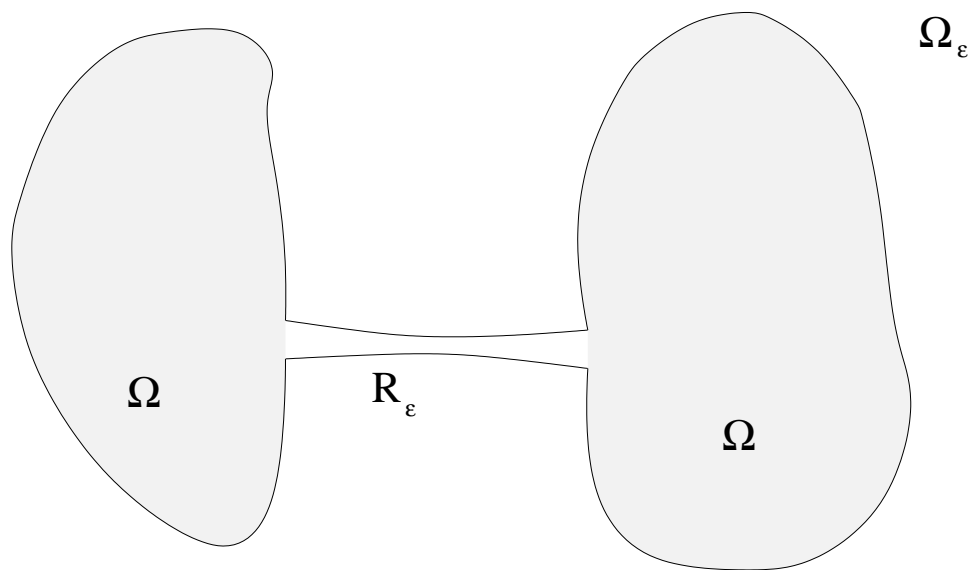
Case 2. Dumbbell type domain

$$\begin{cases} u_t - \Delta u = f(u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

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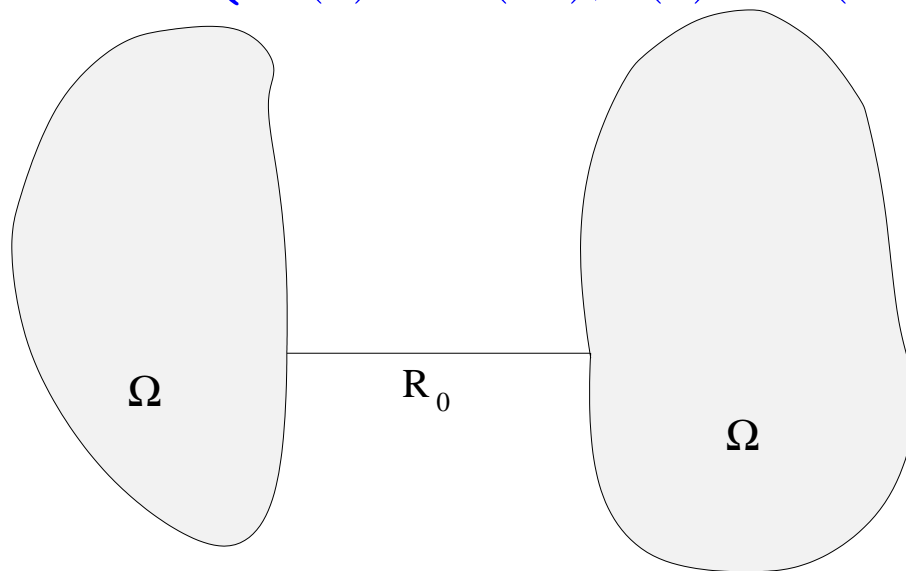
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$$\Omega_\epsilon = \Omega \cup R_\epsilon$$



The limit problem and limit “domain” are

$$\left\{ \begin{array}{l} w_t - \Delta w = f(w) \quad \text{in } \Omega \\ \frac{\partial w}{\partial n} = 0 \quad \text{on } \partial\Omega \\ \\ v_t - \frac{1}{g}(gv_x)_x = f(v), \quad x \in (0, 1) \\ v(0) = w(P_0), v(1) = w(P_1) \end{array} \right.$$

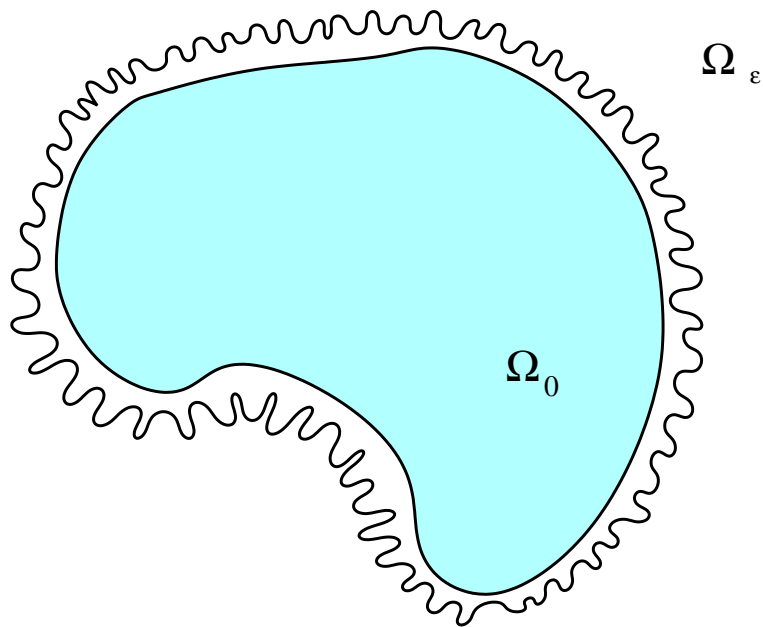


J.A., A.N. Carvalho, G. Lozada-Cruz “Dynamics in Dumbbell Domains I. Continuity of the set of equilibria”, *Journal of Differential Equations*, 231, Issue 2, pp. 551-597, (2006),

J.A., A.N. Carvalho, G. Lozada-Cruz “Dynamics in Dumbbell Domains II. Continuity of attractors”, In preparation

Case 3. Nonlinear boundary conditions and boundary oscillations

$$\begin{cases} u_t - \Delta u + u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$



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J.A., S.M. Bruschi “Boundary oscillations and nonlinear boundary conditions”, *C. R. Acad. Sci. Paris*, t. 343, Series I, pp. 99-104 (2006)

J.A., S.M. Bruschi “Rapidly varying boundaries in equations with nonlinear boundary conditions. The case of a Lipschitz deformation”, *Math. Methods and Models in Applied Science* (2007). **To appear.**

J.A., S.M. Bruschi “Very rapidly varying boundaries in equations with nonlinear boundary conditions.”, *In preparation*

Boundary oscillations

Joint work with [Simone M. Bruschi](#), UNESP, Brazil

$$\begin{cases} u_t - \Delta u + u = f(x, u) & \text{in } \Omega \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{en } \partial\Omega. \end{cases}$$

i) $\Omega \subset \mathbb{R}^N$ bounded smooth domain

ii) $f, g : \mathbb{R}^N \times \mathbb{R} \rightarrow \mathbb{R}$, regular enough and satisfying the dissipativeness conditions: $\exists M > 0$, such that

$$f(x, u)u \leq 0, \quad -g(x, u)u \leq 0, \quad \forall |u| \geq M, \quad x \in \mathbb{R}^N.$$

To simplify, let us assume that both $f(x, \cdot)$ and $g(x, \cdot)$ are globally Lipschitz functions, uniformly in $x \in \mathbb{R}^N$, that is:

$$|f(x, s) - f(x, r)| \leq L|s - r|$$

$$|g(x, s) - g(x, r)| \leq L|s - r|$$

This problem is well posed in spaces $L^2(\Omega)$ and $H^1(\Omega)$, it generates a dynamical system (semiflow, nonlinear semigroup) in both spaces.

If we denote $X(\Omega) = L^2(\Omega)$ or $H^1(\Omega)$:

$$\begin{aligned} T_{\Omega}(t) : X(\Omega) &\rightarrow X(\Omega) \\ \xi &\rightarrow u(t, \cdot, \xi) \end{aligned}$$

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$$\begin{aligned} T_\Omega(t) : X(\Omega) &\rightarrow X(\Omega) \\ \xi &\rightarrow u(t, \cdot, \xi) \end{aligned}$$

where $u(t, \cdot, \xi)$ is the solution at time t , with initial data $\xi \in X(\Omega)$, that is

$$\begin{cases} u_t - \Delta u + u = f(x, u) & \text{en } \Omega \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{en } \partial\Omega, \\ u(0, \cdot) = \xi \end{cases}$$

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► it has bounds in $L^\infty(\Omega)$, independent of Ω . Actually, by the maximum principle, we have

$$\|\varphi\|_{L^\infty(\Omega)} \leq M, \quad \forall \varphi \in \mathcal{A}_\Omega,$$

where M is such that $f(x, u)u \leq 0$ and $-g(x, u)u \leq 0$, for $|u| \geq M$.

The dynamical system is C^1 and gradient.

$$V(u) = \frac{1}{2} \int_{\Omega} (|\nabla u(x)|^2 + |u|^2) dx - \int_{\Omega} \int_0^{u(x)} f(x, s) ds dx + \\ + \int_{\partial\Omega} \int_0^{u(x)} g(x, s) ds dx$$

is a Lyapunov function.

As a matter of fact, we have

$$\frac{d}{dt}(V(u(t))) = - \int_{\Omega} |u_t(t, x)|^2 dx$$

Hence, if \mathcal{E}_Ω , is **the set of equilibria**, that is, all solutions of

$$\begin{cases} -\Delta u + u = f(x, u) & \text{en } \Omega \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{en } \partial\Omega. \end{cases}$$

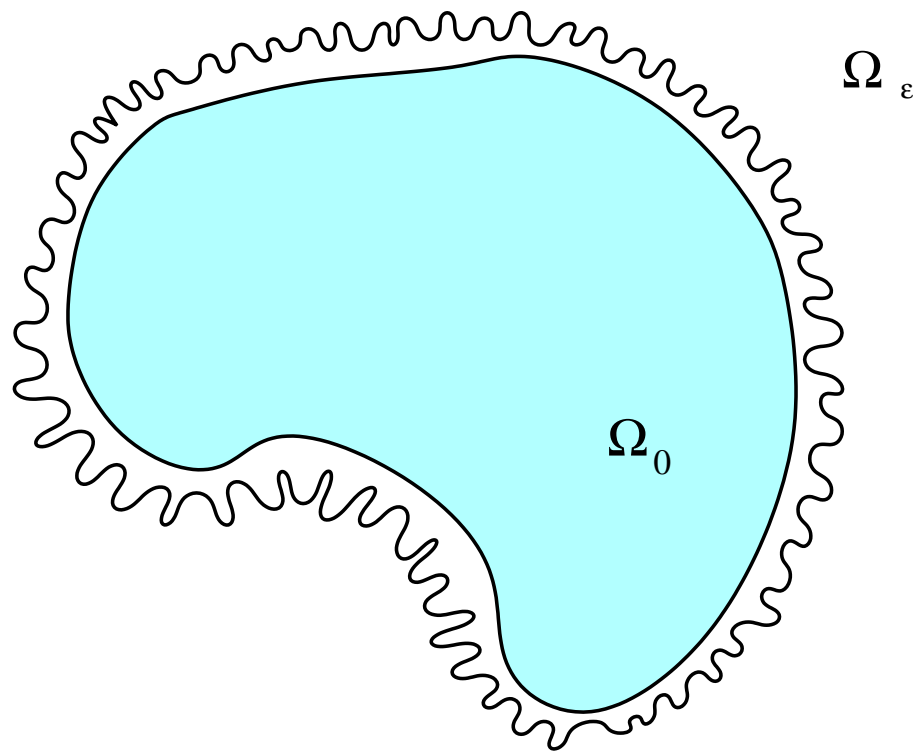
then \mathcal{A}_Ω is formed by

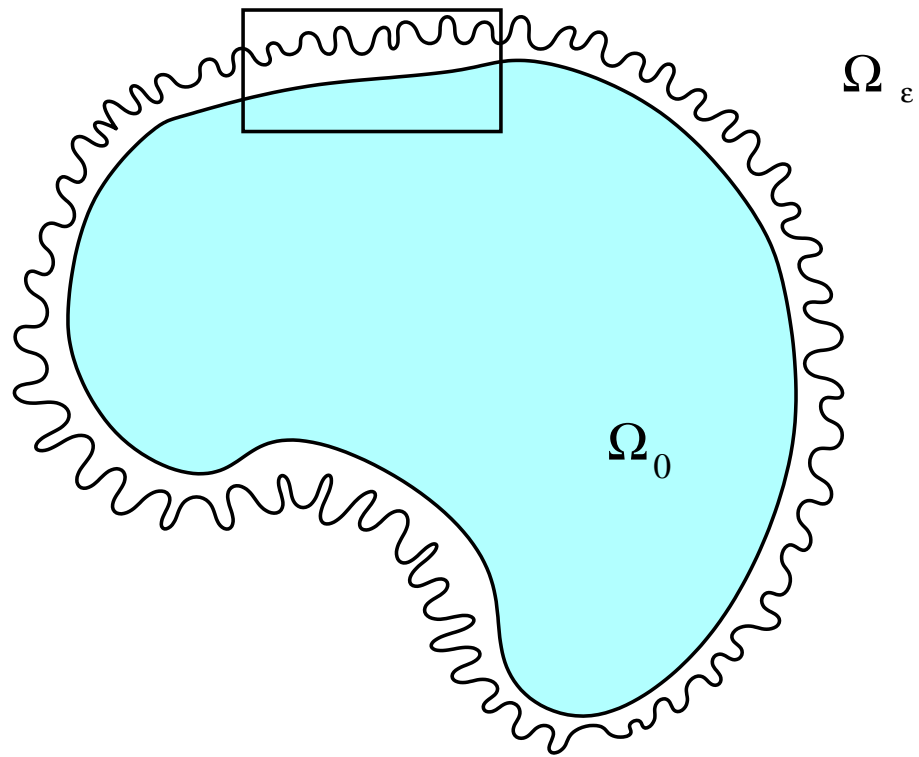
- \mathcal{E}_Ω
- Connections among elements of \mathcal{E}_Ω .

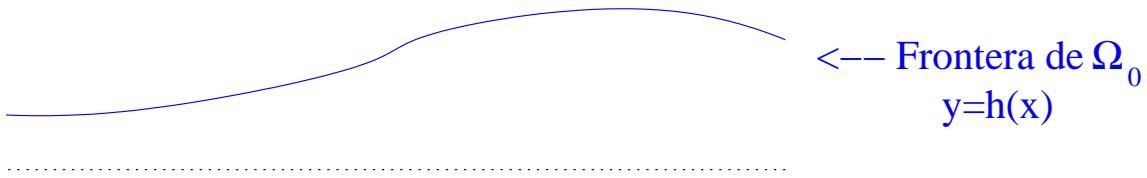
Hence, if we want to analyze the dependence of attractor $\mathcal{A}_\Omega \subset H^1(\Omega)$ as a function of the domain Ω , we better start understanding the dependence of the set of equilibria, \mathcal{E}_Ω , with respect to the domain.

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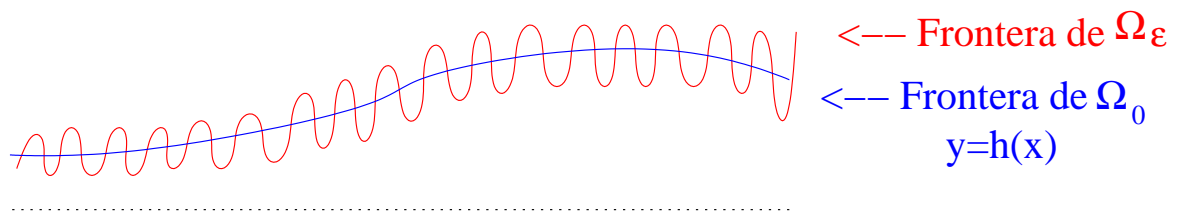
Hence, we consider a family of domains Ω_ϵ **converging in certain sense** to Ω_0 as $\epsilon \rightarrow 0$. We want to understand the behavior of $\mathcal{E}_{\Omega_\epsilon}$ as $\epsilon \rightarrow 0$.







$$y=h(x)+ \varepsilon \sin(x/ \varepsilon^\alpha) j(x)$$



What is the **limit problem**?

What is the **limit problem**?. Let us consider a simpler case:

$$\begin{cases} -\Delta u_\epsilon + u_\epsilon = 0 & \text{en } \Omega_\epsilon \\ \frac{\partial u_\epsilon}{\partial n} + g(x) = 0 & \text{en } \partial\Omega_\epsilon. \end{cases}$$

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($g \in C^0(\mathbb{R}^N)$).

It is equivalent to,

$$\int_{\Omega_\epsilon} (\nabla u_\epsilon \nabla \varphi + u_\epsilon \varphi) = - \int_{\partial\Omega_\epsilon} g \varphi, \quad \forall \varphi \in C^\infty(\mathbb{R}^N)$$

With a priori estimates, weak limits, etc... we have that there exist $u_0 \in H^1(\Omega_0)$ such that

$$\int_{\Omega_\epsilon} (\nabla u_\epsilon \nabla \varphi + u_\epsilon \varphi) \rightarrow \int_{\Omega_0} (\nabla u_0 \nabla \varphi + u_0 \varphi), \quad \forall \varphi \in C^\infty(\mathbb{R}^N)$$

Moreover,

$$\int_{\partial\Omega_\epsilon} g(x)\varphi(x) \rightarrow \int_{\partial\Omega_0} \gamma(x)g(x)\varphi(x)$$

where the function γ satisfies $1 \leq \gamma(x) \leq +\infty$.

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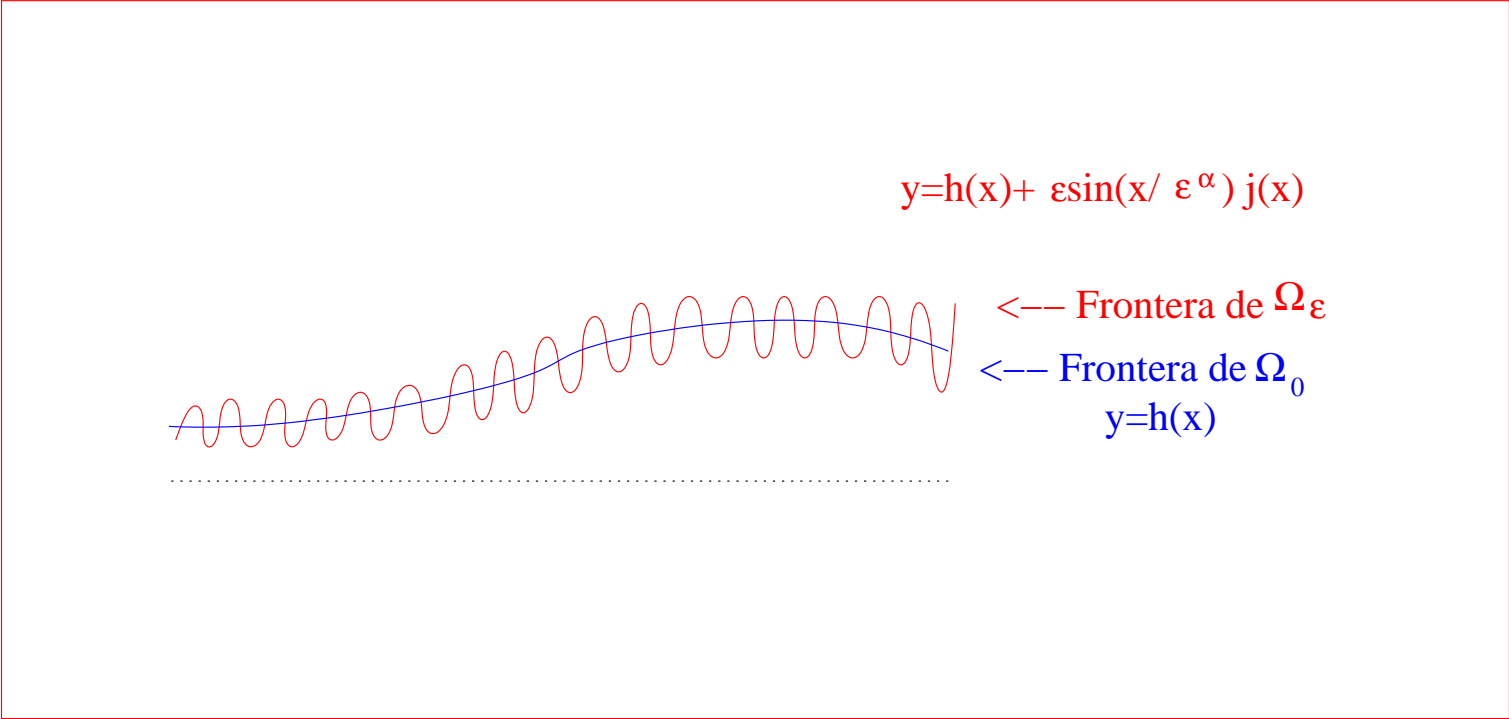
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For $x_0 \in \partial\Omega_0$, the value $\gamma(x_0)$ represents the relative measure of $\partial\Omega_\epsilon$ with respect to $\partial\Omega_0$ locally around x_0 . That is,

$$\gamma(x_0) \approx \frac{|\partial\Omega_\epsilon \cap B(x_0, r)|}{|\partial\Omega_0 \cap B(x_0, r)|}, \quad \text{as } \epsilon, r \rightarrow 0$$

For the case



- If $0 \leq \alpha < 1$, then $\gamma(x) \equiv 1$.
- If $\alpha > 1$, then $\gamma = +\infty$
- If $\alpha = 1$, then $1 \leq \gamma(x) \leq C$. For instance, if $h \equiv 0$, then

$$\gamma(x) = \frac{1}{2\pi} \int_0^{2\pi} \sqrt{1 + (j(x) \cos(z))^2} dz$$

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We will restrict our exposition to the case $\alpha \leq 1$.

In particular, domains Ω_ϵ satisfy

- $\bigcup_{0 \leq \epsilon \leq \epsilon_0} \bar{\Omega}_\epsilon \subset U$, where U is a bounded domain.
- $\Omega_\epsilon \rightarrow \Omega_0$, $\partial\Omega_\epsilon \rightarrow \partial\Omega_0$, in the sense of Hausdorff
- Ω_ϵ are smooth domains and **uniformly Lipschitz** in ϵ .
 - \exists extension operators $P_\epsilon : H^1(\Omega_\epsilon) \rightarrow H^1(U)$, with **norms uniformly bounded in ϵ** . Same for $W^{1,p}(\Omega_\epsilon)$, $C^\beta(\Omega_\epsilon)$.
 - The norms of Sobolev embeddings $W^{1,p}(\Omega_\epsilon) \hookrightarrow L^q(\Omega_\epsilon)$ and trace operators $W^{1,p}(\Omega_\epsilon) \rightarrow L^r(\partial\Omega_\epsilon)$ are **uniformly bounded in ϵ** .
- There exists a function γ defined on $\partial\Omega_0$, with $1 \leq \gamma \leq M$, such that, for all $f \in W^{1,1}(U)$,

$$\int_{\partial\Omega_\epsilon} f \rightarrow \int_{\partial\Omega_0} \gamma f$$

Hence, for the nonlinear problem

$$(P)_\epsilon \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

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E.N. Dancer & D. Daners, "Domain Perturbations for Elliptic Equations subject to Robin Boundary Conditions", *J. of Diff. Equations* 138, (1997) 86-132.

They treat the case $g(x, u) = b(x)u$, with $b(x) \geq b_0 > 0$

How do we compare functions in Ω_0 with functions in Ω_ϵ ?

In what sense can we say that a sequence of functions u_ϵ , each of them defined in Ω_ϵ converges to a function u_0 defined in Ω_0 ?

We have an extension operator

$$E_\epsilon : H^1(\Omega_0) \rightarrow H^1(\Omega_\epsilon)$$

which is obtained as $E_\epsilon = R_\epsilon \circ E$, where $E : H^1(\Omega_0) \rightarrow H^1(\mathbb{R}^N)$ and R_ϵ is the restriction operator of functions defined in \mathbb{R}^N to functions defined in Ω_ϵ .

(E_ϵ is also an extension operator for functions $L^p; W^{1,p}, C^\beta$).

We define,

Definition. A sequence of elements $u_\epsilon \in H^1(\Omega_\epsilon)$ is said to be *E-convergent* to $u \in H^1(\Omega_0)$ if $\|u_\epsilon - E_\epsilon u\|_{H^1(\Omega_\epsilon)} \rightarrow 0$ as $\epsilon \rightarrow 0$. We write this as $u_\epsilon \xrightarrow{E} u$.

Theorem. (Continuity of the set of solutions)

i) If u_ϵ^ , $0 < \epsilon \leq \epsilon_0$, is a family of solutions of problem $(P)_\epsilon$ then there exists a subsequence, still denoted by u_ϵ^* , and a solution of problem $(P)_0$, $u_0^* \in H^1(\Omega_0)$, with the property that, for some $0 < \beta < 1$*

$$\|u_\epsilon^* - E_\epsilon u_0^*\|_{H^1(\Omega_\epsilon)} + \|u_\epsilon^* - E_\epsilon u_0^*\|_{C^\beta(\bar{\Omega}_\epsilon)} \rightarrow 0, \text{ as } \epsilon \rightarrow 0.$$

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ii) If $u_0^ \in H^1(\Omega_0)$ is a solution of $(P)_0$, which is hyperbolic ($\lambda = 0$ is not an eigenvalue of the linearized problem of $(P)_0$ around u_0^*), then, there exists $\delta > 0$ small such that problem $(P)_\epsilon$ has one and only one solution u_ϵ^* , satisfying $\|u_\epsilon^* - E_\epsilon u_0^*\|_{H^1(\Omega_\epsilon)} \leq \delta$ for ϵ small enough.*

u_0^* hyperbolic

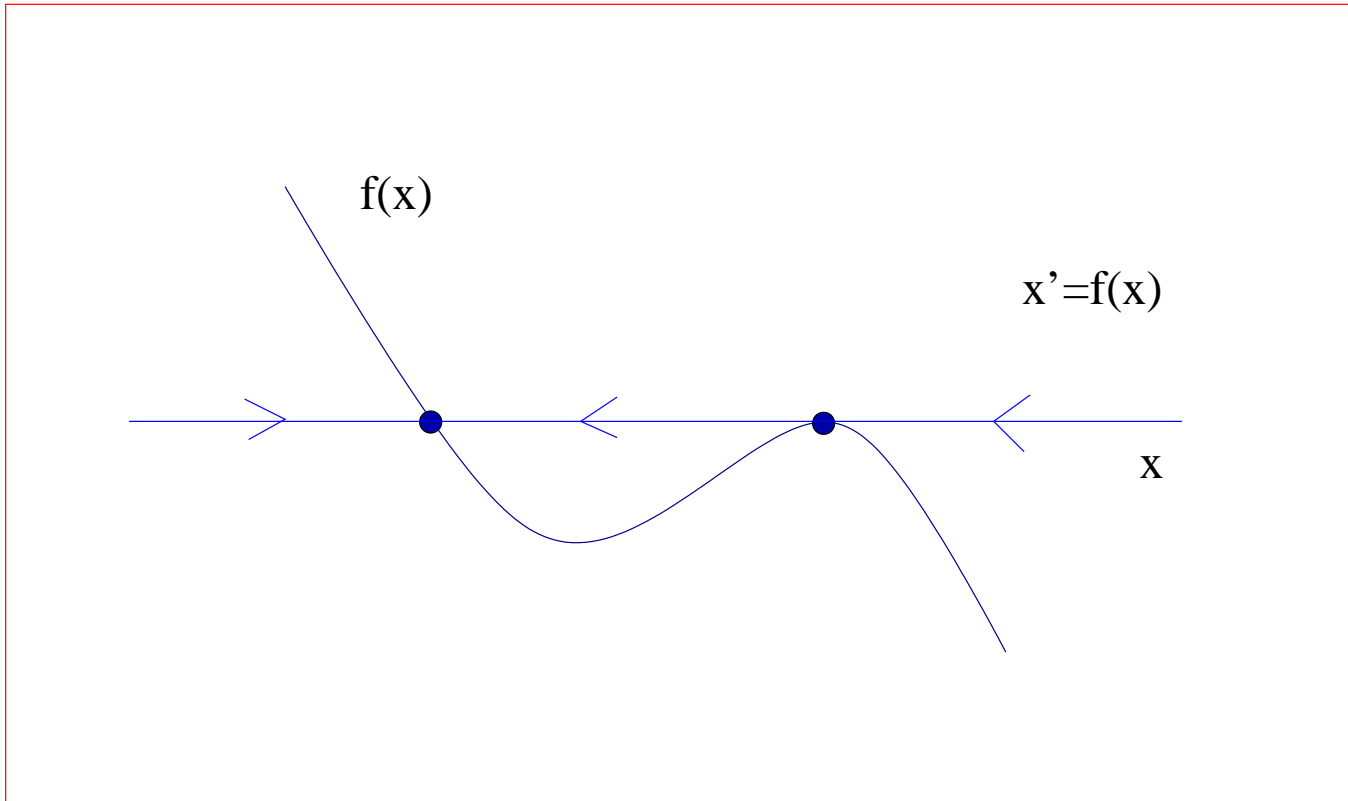
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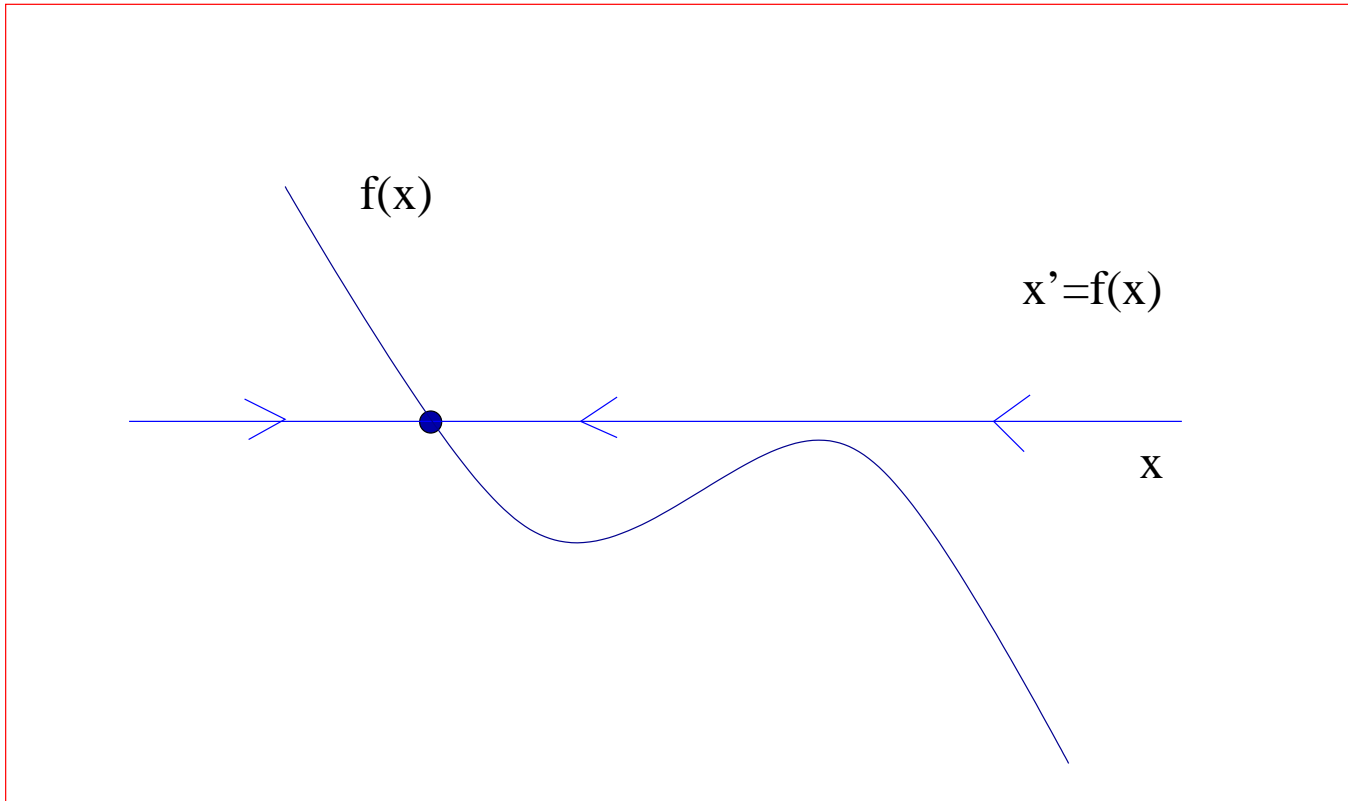
$$\begin{cases} -\Delta w + w - \partial_u f(x, u_0^*)w = \lambda w & \text{in } \Omega_0, \\ \frac{\partial w}{\partial n} + \gamma \partial_u g(x, u_0^*)w = 0 & \text{on } \partial\Omega_0. \end{cases} \quad (1)$$

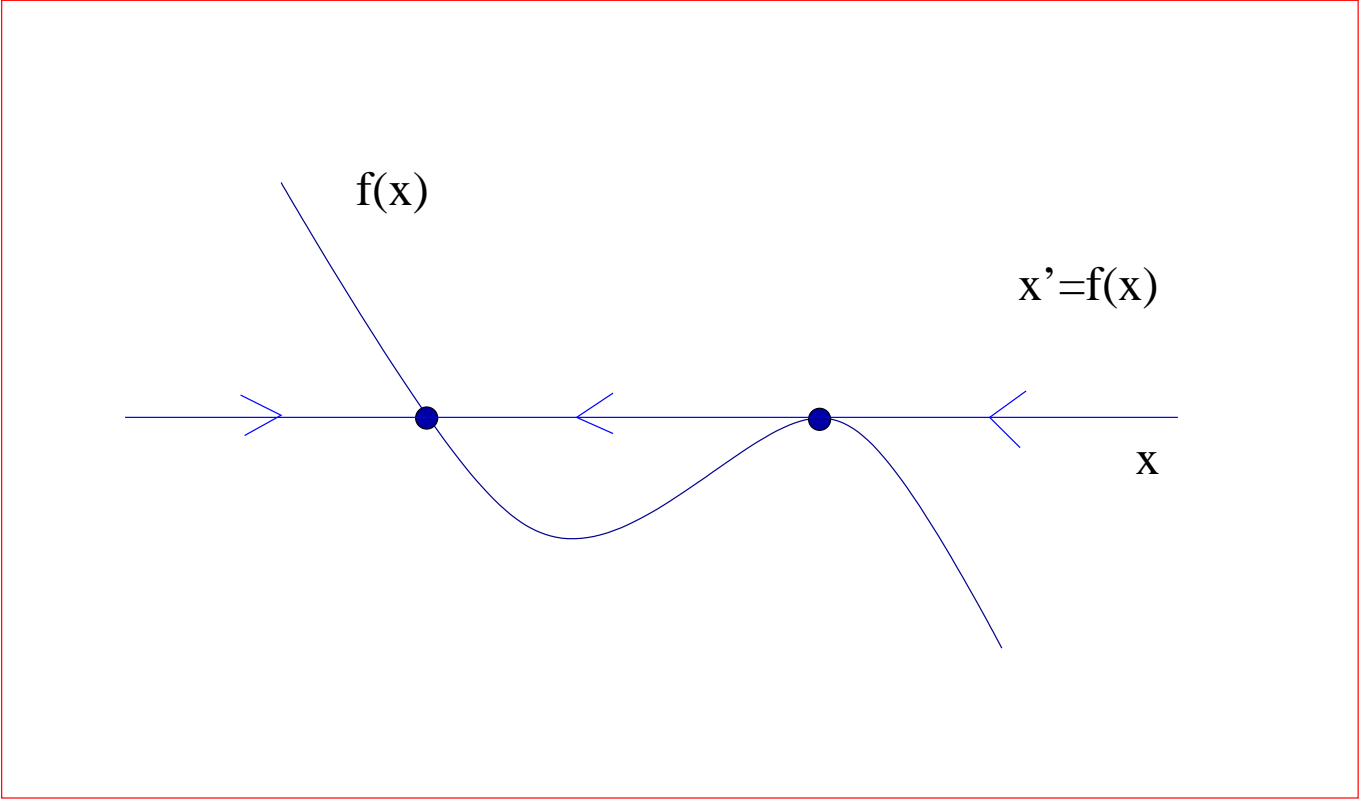
u_0^* hyperbolic $\iff \lambda = 0$ is not an eigenvalue of

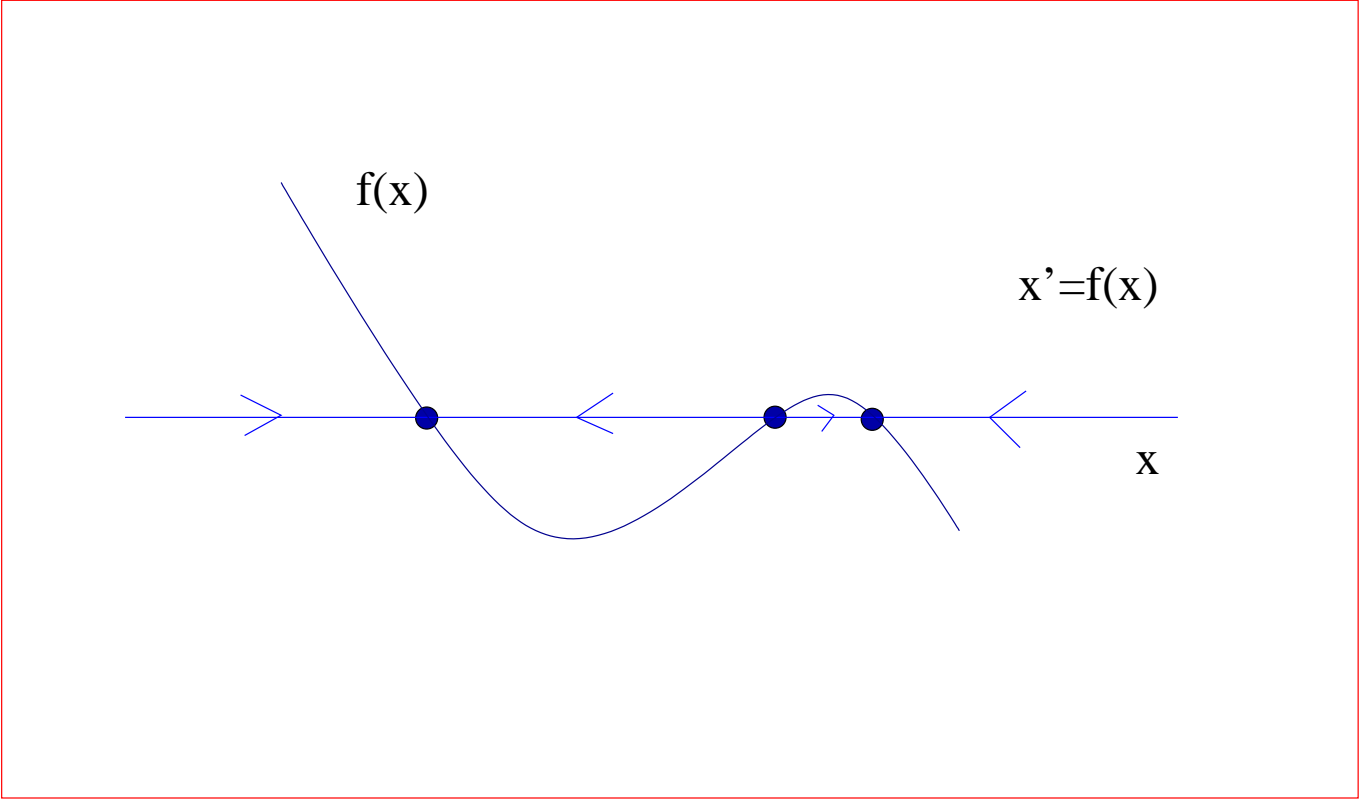
$$\begin{cases} -\Delta w + w - \partial_u f(x, u_0^*)w = \lambda w & \text{in } \Omega_0, \\ \frac{\partial w}{\partial n} + \gamma \partial_u g(x, u_0^*)w = 0 & \text{on } \partial\Omega_0. \end{cases} \quad (2)$$

This condition is not technical.









Idea of the proof.

- We use a functional analysis for operators defined in different spaces.

Let H_ϵ , $0 < \epsilon \leq \epsilon_0$ and H be Hilbert spaces. Let also

$$E_\epsilon : H \rightarrow H_\epsilon$$

be a bounded linear operator, such that $\|E_\epsilon u\|_{H_\epsilon} \rightarrow \|u\|_H$.

For instance, $H_\epsilon = H^1(\Omega_\epsilon)$, $H = H^1(\Omega_0)$ and $E_\epsilon : H \rightarrow H_\epsilon$ is the constructed extension operator.

Definition. A sequence of elements $u_\epsilon \in H_\epsilon$ is said to be *E-convergent* to $u \in H$ if $\|u_\epsilon - E_\epsilon u\|_{H_\epsilon} \rightarrow 0$ as $\epsilon \rightarrow 0$. We write this as $u_\epsilon \xrightarrow{E} u$.

Definition. A sequence of elements $u_\epsilon \in H_\epsilon$ is said to be *E-precompact* if for any subsequence $\{u_{\epsilon_n}\}$ there exist a subsequence $\{u_{\epsilon_{n'}}\}$ and $u \in H$ such that $u_{\epsilon_{n'}} \xrightarrow{E} u$, as $n' \rightarrow \infty$.

Definition. A family of operators $T_\epsilon : H_\epsilon \rightarrow H_\epsilon$, $\epsilon \in (0, 1]$, *E-converges* to $T : H \rightarrow H$, if $T_\epsilon u_\epsilon \xrightarrow{E} T u$, whenever $u_\epsilon \xrightarrow{E} u$. We denote this by $T_\epsilon \xrightarrow{EE} T$.

Definition. A family of compact operators $T_\epsilon : H_\epsilon \rightarrow H_\epsilon$, $\epsilon \in (0, 1]$ *converges compactly* to a compact $T : H \rightarrow H$ if for any family u_ϵ with $\|u_\epsilon\|_{H_\epsilon}$ bounded, the family $\{T_\epsilon u_\epsilon\}$ is *E-precompact* and $T_\epsilon \xrightarrow{EE} T$. We write $T_\epsilon \xrightarrow{CC} T$.

Theorem. Let $T_\epsilon : H_\epsilon \rightarrow H_\epsilon$ be a family of compact operators such that $T_\epsilon \xrightarrow{CC} T$. Let u_ϵ be a fixed point of T_ϵ such that $\|u_\epsilon\|_{H_\epsilon}$ is uniformly bounded. Then, there exists a subsequence u_{ϵ_k} and $u \in H$ a fixed point of T , such that $u_{\epsilon_k} \xrightarrow{E} u$.

Proof. If $u_\epsilon = T_\epsilon u_\epsilon$ and $\|u_\epsilon\|_{H_\epsilon} \leq M$, by compact convergence $T_\epsilon \xrightarrow{CC} T$, there exists a subsequence $\epsilon_k \rightarrow 0$, and a $u \in H$ such that $T_{\epsilon_k} u_{\epsilon_k} \xrightarrow{E} u$. Hence $u_{\epsilon_k} = T_{\epsilon_k} u_{\epsilon_k} \xrightarrow{E} u$ and therefore, $T_{\epsilon_k} u_{\epsilon_k} \xrightarrow{E} Tu$. This implies $u = Tu$.

G. Vainikko, Approximative methods for nonlinear equations (two approaches to the convergence problem), *Nonlinear Analysis, Theory, Methods & Applications*, vol2, 6 (1978) 647-687.

A. Carvalho, S. Piskarev, A general approximation scheme for attractors of abstract parabolic problems, *Numer. Funct. Anal. Optim*, 27 (2006), no. 7-8, 785–829.

For our case, we have $H_\epsilon = H^1(\Omega_\epsilon)$, $H = H^1(\Omega_0)$, E_ϵ the extension operator.

$T_\epsilon : H^1(\Omega_\epsilon) \rightarrow H^1(\Omega_\epsilon)$ is given by $T_\epsilon(z_\epsilon) = u_\epsilon$, where u_ϵ is the unique solution of

$$\begin{cases} -\Delta u_\epsilon + u_\epsilon = f(x, z_\epsilon) & \text{in } \Omega_\epsilon \\ \frac{\partial u_\epsilon}{\partial n} + g(x, z_\epsilon) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

$T : H^1(\Omega_0) \rightarrow H^1(\Omega_0)$ is given by $T(z) = u$, where u is the unique solution of

$$\begin{cases} -\Delta u + u = f(x, z) & \text{in } \Omega_0 \\ \frac{\partial u}{\partial n} + \gamma(x)g(x, z) = 0 & \text{on } \partial\Omega_0. \end{cases}$$

That is,

$$T_\epsilon = A_\epsilon^{-1} \circ h_\epsilon$$

where $A_\epsilon = -\Delta + I$ with homogeneous Neumann boundary conditions and $h_\epsilon : H^1(\Omega_\epsilon) \rightarrow H^{-\alpha}(\Omega_\epsilon)$ is defined as

$$\langle h_\epsilon(z_\epsilon), \varphi_\epsilon \rangle = \int_{\Omega_\epsilon} f(x, z_\epsilon) \varphi_\epsilon(x) dx - \int_{\partial\Omega_\epsilon} g(x, z_\epsilon) \varphi_\epsilon(x) d\sigma_\epsilon$$

Similarly,

$$T = A^{-1} \circ h$$

where $A = -\Delta + I$ with homogeneous Neumann boundary conditions and $h : H^1(\Omega_0) \rightarrow H^{-\alpha}(\Omega_0)$ is defined as

$$\langle h(z), \varphi \rangle = \int_{\Omega_0} f(x, z) \varphi(x) dx - \int_{\partial\Omega_0} \gamma(x) g(x, z) \varphi(x) d\sigma$$

To show that $T_\epsilon \xrightarrow{CC} T$, we get a sequence of $z_\epsilon \in H^1(\Omega_\epsilon)$, $\|z_\epsilon\|_{H^1(\Omega_\epsilon)} \leq M$ and let $u_\epsilon = T_\epsilon(z_\epsilon)$, that is,

$$\begin{cases} -\Delta u_\epsilon + u_\epsilon = f(x, z_\epsilon) & \text{in } \Omega_\epsilon \\ \frac{\partial u_\epsilon}{\partial n} + g(x, z_\epsilon) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

We have $\|u_\epsilon\|_{H^1(\Omega_\epsilon)} \leq M'$. Using the extension operators P_ϵ and getting subsequences, we obtain z and u , such that

$$P_{\epsilon_k} z_{\epsilon_k} \rightharpoonup z \in H^1(U), \quad P_{\epsilon_k} u_{\epsilon_k} \rightharpoonup u \in H^1(U)$$

Using the weak formulation of the equation,

$$\int_{\Omega_\epsilon} \nabla u_\epsilon \nabla \varphi + u_\epsilon \varphi = \int_{\Omega_\epsilon} f(x, z_\epsilon) \varphi - \int_{\partial\Omega_\epsilon} g(x, z_\epsilon) \varphi$$

We can easily pass to the limit in the first three terms. For the boundary term, we have

$$\int_{\partial\Omega_\epsilon} g(x, z_\epsilon) \varphi = \int_{\partial\Omega_\epsilon} (g(x, z_\epsilon) - g(x, z)) \varphi + \int_{\partial\Omega_\epsilon} g(x, z) \varphi$$

and

$$\left| \int_{\partial\Omega_\epsilon} (g(x, z_\epsilon) - g(x, z)) \varphi \right| \leq C \int_{\partial\Omega_\epsilon} |z_\epsilon - z| \rightarrow 0$$

Moreover,

$$\int_{\partial\Omega_\epsilon} g(x, z) \varphi(x) \rightarrow \int_{\partial\Omega_0} \gamma(x) g(x, z) \varphi(x)$$

Hence, passing to the limit

$$\int_{\Omega_0} \nabla u \nabla \varphi + u \varphi = \int_{\Omega_0} f(x, z) \varphi - \int_{\partial\Omega_0} \gamma(x) g(x, z) \varphi$$

which means that u is solution of

$$\begin{cases} -\Delta u + u = f(x, z) & \text{in } \Omega_0 \\ \frac{\partial u}{\partial n} + \gamma(x) g(x, z) = 0 & \text{on } \partial\Omega_0. \end{cases}$$

To show that actually

$$\|u_\epsilon - E_\epsilon u\|_{H^1(\Omega_\epsilon)} \rightarrow 0,$$

we show the convergence of the norms, that is $\|u_\epsilon\|_{H^1(\Omega_\epsilon)} \rightarrow \|u\|_{H^1(\Omega_0)}$.

Convergence in C^β -norms.

We are able to obtain Hölder estimates of u_ϵ^* independent of ϵ .

Uniqueness of solutions near a hyperbolic solution.

Let u_0^* be a hyperbolic solution of the limit problem. That is, $\lambda = 0$ is not an eigenvalue of

$$\begin{cases} -\Delta\phi + \phi - \partial_u f(x, u_0^*)\phi = \lambda\phi & \text{in } \Omega_0 \\ \frac{\partial\phi}{\partial n} + \gamma(x)\partial_u g(x, u_0^*)\phi = 0 & \text{on } \partial\Omega_0. \end{cases}$$

In particular,

- u_0^* is an isolated solution of $(P)_0$.
- With index-degree theory it is possible to show that for all $0 < \epsilon \leq \epsilon_0$ there exists at least one solution u_ϵ^* of $(P)_\epsilon$ near u_0^* . In particular, there exists a sequence $u_\epsilon^* \xrightarrow{E} u_0^*$.

Let us see that there exists a $\delta > 0$, such that no other function $w_\epsilon \in H^1(\Omega_\epsilon)$ with $\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)} \leq \delta$ is a fixed point of $T_\epsilon = A_\epsilon^{-1} \circ h_\epsilon$.

$$\begin{aligned}
& \|w_\epsilon - T_\epsilon(w_\epsilon)\|_{H^1(\Omega_\epsilon)} = \|w_\epsilon - u_\epsilon^* - (T_\epsilon(w_\epsilon) - T_\epsilon(u_\epsilon^*))\|_{H^1(\Omega_\epsilon)} \\
& \geq \|w_\epsilon - u_\epsilon^* - B_\epsilon(w_\epsilon - u_\epsilon^*) + (T_\epsilon(w_\epsilon) - T_\epsilon(u_\epsilon^*) - B_\epsilon(w_\epsilon - u_\epsilon^*))\|_{H^1(\Omega_\epsilon)} \\
& \geq \|(I - B_\epsilon)(w_\epsilon - u_\epsilon^*)\|_{H^1(\Omega_\epsilon)} - \\
& \quad - \|T_\epsilon(w_\epsilon) - T_\epsilon(u_\epsilon^*) - B_\epsilon(w_\epsilon - u_\epsilon^*)\|_{H^1(\Omega_\epsilon)}
\end{aligned}$$

The linear operator B_ϵ , is defined by $v_\epsilon = B_\epsilon(z_\epsilon)$ where

$$\begin{cases} -\Delta v_\epsilon + v_\epsilon = \partial_u f(x, u_\epsilon^*) z_\epsilon & \text{in } \Omega_\epsilon \\ \frac{\partial v_\epsilon}{\partial n} + \partial_u g(x, u_\epsilon^*) z_\epsilon = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

We can show $B_\epsilon \xrightarrow{CC} B_0$, where $v = B_0(z)$ where

$$\begin{cases} -\Delta v + v = \partial_u f(x, u_0^*)z & \text{in } \Omega_0 \\ \frac{\partial v}{\partial n} + \gamma(x)\partial_u g(x, u_0^*)z = 0 & \text{on } \partial\Omega_0. \end{cases}$$

The fact that u_0^* is a hyperbolic solution, is equivalent to say that $1 \notin \sigma(B_0)$ and therefore $I - B_0$ is invertible, that is $\|(I - B_0)^{-1}\| \leq C$.

The compact convergence of B_ϵ to B_0 implies that for ϵ small $1 \notin \sigma(B_\epsilon)$ and that $\|(I - B_\epsilon)^{-1}\| \leq C'$. Hence, if $\eta = 1/C'$,

$$\|(I - B_\epsilon)\chi\|_{H^1(\Omega_\epsilon)} \geq \eta\|\chi\|_{H^1(\Omega_\epsilon)}$$

In particular,

$$\|(I - B_\epsilon)(w_\epsilon - u_\epsilon^*)\|_{H^1(\Omega_\epsilon)} \geq \eta\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}$$

On the other hand, if we denote by

$$v_\epsilon = T_\epsilon(w_\epsilon) - T_\epsilon(u_\epsilon^*) - B_\epsilon(w_\epsilon - u_\epsilon^*)$$

we have

$$\begin{cases} -\Delta v_\epsilon + v_\epsilon = f(x, w_\epsilon) - f(x, u_\epsilon^*) - \partial_u f(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*) & \text{in } \Omega_\epsilon \\ \frac{\partial v_\epsilon}{\partial n} + g(x, w_\epsilon) - g(x, u_\epsilon^*) - \partial_u g(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

$$\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2 = \int_{\Omega_\epsilon} (f(x, w_\epsilon) - f(x, u_\epsilon^*) - \partial_u f(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon$$

$$- \int_{\partial\Omega_\epsilon} (g(x, w_\epsilon) - g(x, u_\epsilon^*) - \partial_u g(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon$$

$$\begin{aligned}
\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2 &= \int_{\Omega_\epsilon} (f(x, w_\epsilon) - f(x, u_\epsilon^*) - \partial_u f(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\quad - \int_{\partial\Omega_\epsilon} (g(x, w_\epsilon) - g(x, u_\epsilon^*) - \partial_u g(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{2+\frac{2}{N-1}} + \frac{1}{2}\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2
\end{aligned}$$

$$\begin{aligned}
\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2 &= \int_{\Omega_\epsilon} (f(x, w_\epsilon) - f(x, u_\epsilon^*) - \partial_u f(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\quad - \int_{\partial\Omega_\epsilon} (g(x, w_\epsilon) - g(x, u_\epsilon^*) - \partial_u g(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{2+\frac{2}{N-1}} + \frac{1}{2}\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2
\end{aligned}$$

which implies

$$\|v_\epsilon\|_{H^1(\Omega_\epsilon)} \leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{1+\frac{1}{N-1}}$$

$$\begin{aligned}
\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2 &= \int_{\Omega_\epsilon} (f(x, w_\epsilon) - f(x, u_\epsilon^*) - \partial_u f(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\quad - \int_{\partial\Omega_\epsilon} (g(x, w_\epsilon) - g(x, u_\epsilon^*) - \partial_u g(x, u_\epsilon^*)(w_\epsilon - u_\epsilon^*))v_\epsilon \\
&\leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{2+\frac{2}{N-1}} + \frac{1}{2}\|v_\epsilon\|_{H^1(\Omega_\epsilon)}^2
\end{aligned}$$

which implies

$$\|v_\epsilon\|_{H^1(\Omega_\epsilon)} \leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{1+\frac{1}{N-1}}$$

Equivalently,

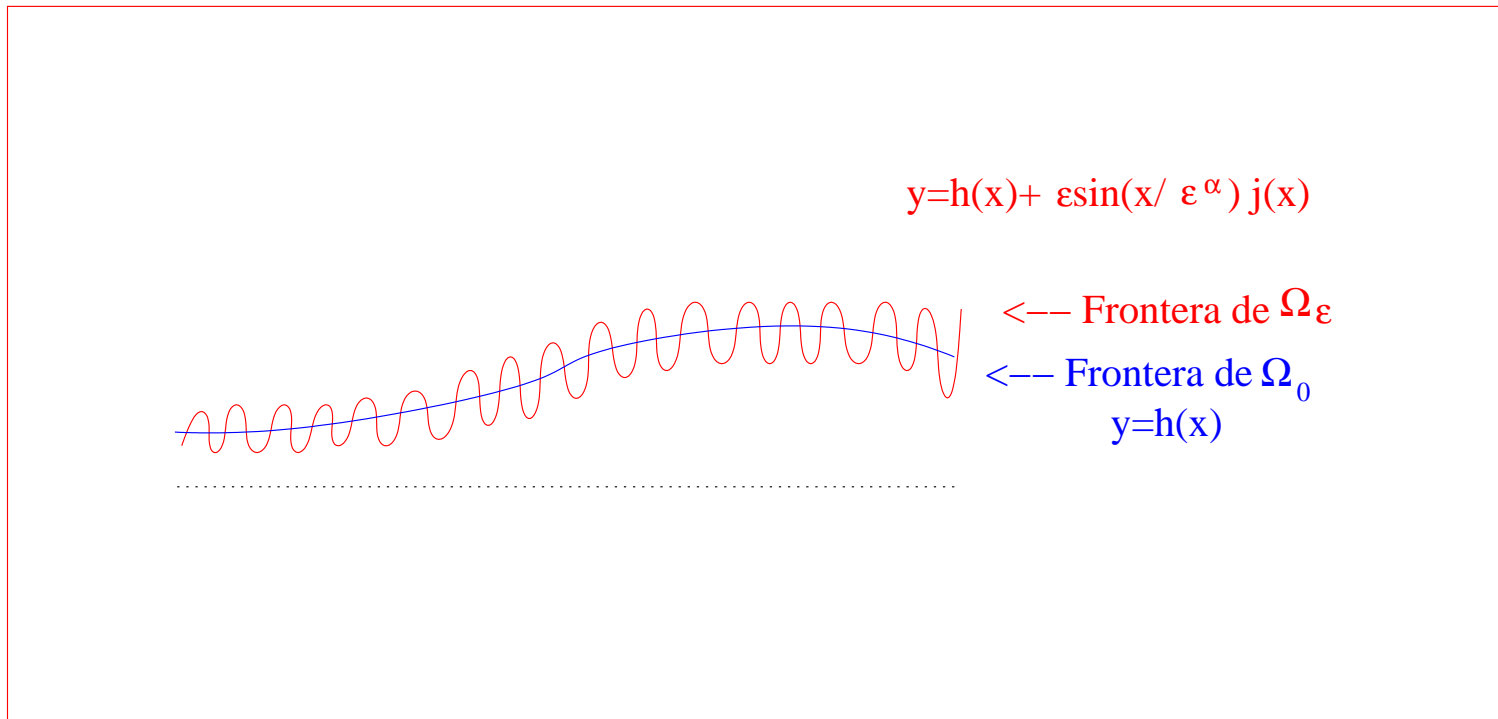
$$\|T_\epsilon(w_\epsilon) - T_\epsilon(u_\epsilon^*) - B_\epsilon(w_\epsilon - u_\epsilon^*)\|_{H^1(\Omega_\epsilon)} \leq C\|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{1+\frac{1}{N-1}}$$

Therefore,

$$\begin{aligned} \|w_\epsilon - T_\epsilon(w_\epsilon)\|_{H^1(\Omega_\epsilon)} &\geq \eta \|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)} - C \|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}^{1+\frac{1}{N-1}} \\ &\geq \frac{\eta}{2} \|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)}, \quad \text{if } \|w_\epsilon - u_\epsilon^*\|_{H^1(\Omega_\epsilon)} \leq |\eta/2C|^{N-1} \end{aligned}$$

Hence, the unique fixed point of T_ϵ in $B_{H^1(\Omega_\epsilon)}(u_\epsilon^*, |\eta/2C|^{N-1})$ is u_ϵ^* .

Case $\alpha > 1$. Very rapid oscillations.



Formally we have $\gamma = +\infty$.

For the nonlinear problem

$$(P)_\epsilon \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

When $\alpha \leq 1$, the limit problem was given by

$$(P)_0 \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_0 \\ \frac{\partial u}{\partial n} + \gamma(x)g(x, u) = 0 & \text{on } \partial\Omega_0. \end{cases}$$

For the nonlinear problem

$$(P)_\epsilon \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} + g(x, u) = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

When $\alpha \leq 1$, the limit problem was given by

$$(P)_0 \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_0 \\ \frac{\partial u}{\partial n} + \gamma(x)g(x, u) = 0 & \text{on } \partial\Omega_0. \end{cases}$$

The factor γ amplifies the dissipativity properties of the boundary condition.

Hence, if $\alpha > 1$ so that $\gamma = +\infty$ and the boundary condition is dissipative: for instance $g(x, u)u \geq \eta|u|^2$ for some $\eta > 0$, then the limit problem is:

$$(P)_0 \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_0 \\ u = 0 & \text{on } \partial\Omega_0. \end{cases}$$

J.A. & S. Bruschi, "Boundary oscillations and nonlinear boundary conditions",
C.R.A.S. 343 (2006)

J.A. & S. Bruschi, "Very rapidly varying boundaries in equations with nonlinear boundary conditions."(In preparation)

If $\alpha > 1$ and the boundary condition is not dissipative: say $g(x, u) = -u$

$$(P)_\epsilon \quad \begin{cases} -\Delta u + u = f(x, u) & \text{in } \Omega_\epsilon \\ \frac{\partial u}{\partial n} - u = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

The behavior as $\epsilon \rightarrow 0$ is not clear and it may be very complicated.

For instance, if $f(x, u) = f(u)$ and $f(0) = 0$ and if we denote by $f'(0) = a$, then the eigenvalue problem of the linearized equation around the trivial solution is given by:

$$(EP)_\epsilon \quad \begin{cases} -\Delta \phi + \phi - a\phi = \lambda\phi & \text{in } \Omega_\epsilon \\ \frac{\partial \phi}{\partial n} - \phi = 0 & \text{on } \partial\Omega_\epsilon. \end{cases}$$

The Raleigh quotient is:

$$J(\phi) = \frac{\int_{\Omega_\epsilon} |\nabla \phi|^2 + (1 - a)|\phi|^2 - \int_{\partial\Omega_\epsilon} |\phi|^2}{\int_{\Omega_\epsilon} |\phi|^2}$$

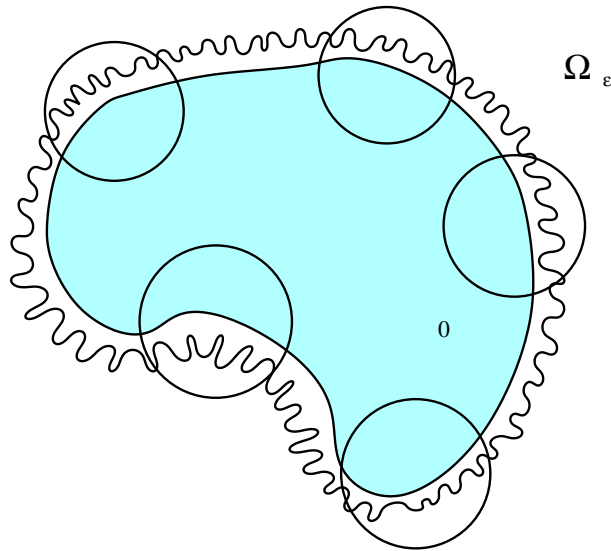
It is not difficult to prove that if λ_n^ϵ is the n-th eigenvalue, we have

$$\lambda_n^\epsilon \rightarrow -\infty \text{ as } \epsilon \rightarrow 0$$

for all $n = 1, 2, \dots$

The Raleigh quotient is:

$$J(\phi) = \frac{\int_{\Omega_\epsilon} |\nabla \phi|^2 + (1 - a)|\phi|^2 - \int_{\partial\Omega_\epsilon} |\phi|^2}{\int_{\Omega_\epsilon} |\phi|^2}$$



This is an indication that as $\epsilon \rightarrow 0$, the trivial solution becomes more and more unstable, it undergoes a sequence of bifurcations and the dynamics complicates tremendously.