

Lüscher corrections in integrable and string sigma models

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- Introduction
- Lüscher correction for the mass gap
- Lüscher's F-term and the Bajnok-Janik formula
- Example: XXX model ($O(2)$ model)
- Example: Konishi operator

AdS/CFT conjecture

String σ -model on $AdS_5 \times S_5 \iff \mathcal{N} = 4 N_c \rightarrow \infty$ SYM 3+1 dim

Energy of string states \iff Anomalous dims of local operators

String model coupling $g \iff$ 't Hooft coupling $\lambda = 4\pi^2 g^2$

Integrability (both sides)!

- asymptotic BA: all power corrections, neglecting “wrapping”

Beisert, Staudacher '05

- generalized Lüscher approach: 4-loop Konishi (agree with PT)

Janik, Lukowski '07 Bajnok, Janik '09

$$\mathcal{K} = \text{Tr}\{D^2 Z^2 - (DZ)^2\}$$

- higher twist-two operators in sl_2 sector: agree with BFKL

Bajnok, Janik, Lukowski '09

- 5-loop Konishi

Bajnok, Hegedus, Janik, Lukowski '10

$$\Delta^{(10)} = \Delta_{\text{asympt}}^{(10)} + g^{10} \left\{ -\frac{81\zeta(3)^2}{16} + \frac{81\zeta(3)}{32} - \frac{45\zeta(5)}{4} + \frac{945\zeta(7)}{32} - \frac{2835}{256} \right\}$$

- 5-loop twist-two

Lukowski, Rej, Velizhanin '10

- TBA for $AdS_5 \times S_5$ string σ -model

Arutyunov, Frolov '09 Bombardelli, Fioravanti, Tateo '09

Gromov, Kazakov, Kozak, Vieira '09 Arutyunov, Frolov, Suzuki '10

- Linearized TBA \longrightarrow Lüscher numerically

Arutyunov, Frolov, Suzuki '10

Large λ expansion

- classical string:

Gubser, Plebanov, Polyakov '98

Arutyunov, Frolov, Staudacher '04

$$E_K(\lambda) = 2\sqrt[4]{\lambda} + \dots$$

- numerical fit from TBA:

Gromov, Kazakov, Vieira '09

$$E_K(\lambda) = \sqrt[4]{\lambda} \left\{ 2.0004 + \frac{1.988}{\sqrt{\lambda}} - \frac{2.60}{\lambda} + \frac{6.2}{\sqrt{\lambda}^3} + \dots \right\}$$

$$\lambda < 664 < \lambda_{\text{crit}}^{(1)}$$

- semiclassical expansion:

Roiban, Tseytlin '09

$$E_K(\lambda) = \sqrt[4]{\lambda} \left\{ 2 + \frac{1}{\sqrt{\lambda}} + \dots \right\}$$

- numerical fit from TBA:

Frolov '10

$$E_K(\lambda) = \sqrt[4]{\lambda} \left\{ 2.00045 + \frac{1.98}{\sqrt{\lambda}} - \frac{2.55}{\lambda} + \frac{6.7}{\sqrt{\lambda}^3} + \dots \right\}$$

$$\lambda < 2046 < \lambda_{\text{crit}}^{(1)}$$

- no $\lambda_{\text{crit}}^{(1)}$? ($\lambda_{\text{crit}}^{(1)} \approx 774$ from ABA)

Lüscher corrections to the mass gap

Lüscher '83 Klassen, Melzer '91

All order diagrammatic expansion: Lüscher's F-term

$$\Delta m_a^{(F)}(L)$$

$$= -\frac{dL}{4\pi m_a} \sum_b \int_{-\infty}^{\infty} d\theta \left(\frac{m_b \cosh \theta}{2\pi L} \right)^{d/2} K_{d/2-1}(m_b L \cosh \theta) F_{ab} \left(\theta + \frac{i\pi}{2} \right)$$

1+1 dimensional case: $(m_a = m_b = m)$

$$\Delta m_a^{(F)}(L) = -\frac{m}{2\pi} \int_{-\infty}^{\infty} d\theta \cosh \theta e^{-mL \cosh \theta} S_a \left(\theta + \frac{i\pi}{2} \right)$$

$$S_a(\theta) = -n + \sum_{b=1}^n S_{ab}^{ab}(\theta)$$

μ -term (1+1 dim)

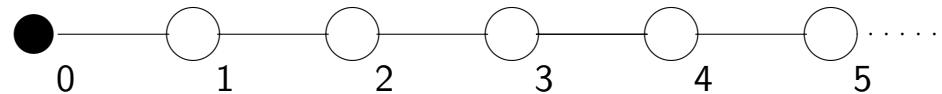
$$\Delta m_a^{(\mu)}(L) = -\frac{1}{8m_a^2} \sum_{bc} \theta(m_a^2 - |m_b^2 - m_c^2|) \frac{\lambda_{abc}^2}{\mu_{abc}} e^{-\mu_{abc}L}$$

λ_{abc} : 3-point coupling

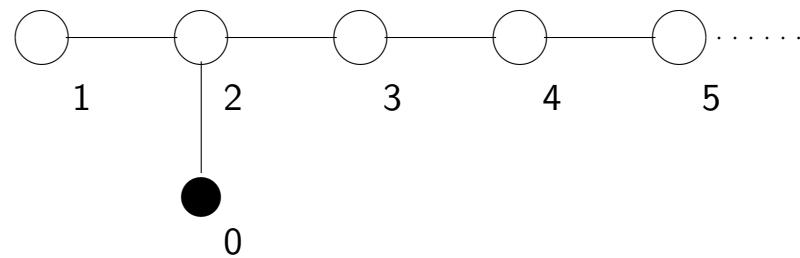
μ_{abc} : height of mass triangle (m_a, m_b, m_c)

- LWW coupling and “step scaling” function Lüscher, Weisz, Wolff '91
- Lattice MC measurements
- Sigma model perturbation theory

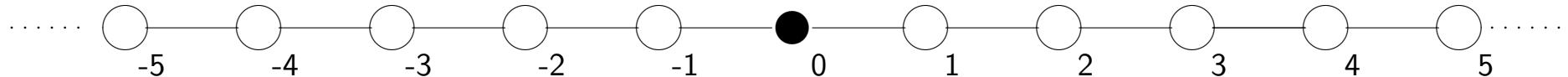
Generalized Lüscher formulae



O(2) model TBA diagram



O(3) model TBA diagram



O(4) model TBA diagram

SIGMA MODELS ($n = 2, 3, 4$)

Bethe-Yang limit energy:

$$E^{(0)} = \sum_{j=1}^r \mu \cosh \theta_j$$

Bethe-Yang equations (QC)

$$QC_k^{(0)}(\theta_1, \dots, \theta_r) = e^{i\mu L \sinh \theta_k} e^{i\mathcal{R}_k} = -1, \quad k = 1, \dots, r,$$

$$e^{i\mathcal{R}_k} = \sigma(\theta_k | \theta_1, \dots, \theta_r)$$

$\sigma(\theta | \theta_1, \dots, \theta_r)$: eigenvalue of the transfer matrix
(unitary and crossing symmetric S-matrix)

Lüscher's F-term:

$$E^{(1)} = E^{(0)} + \delta E =$$

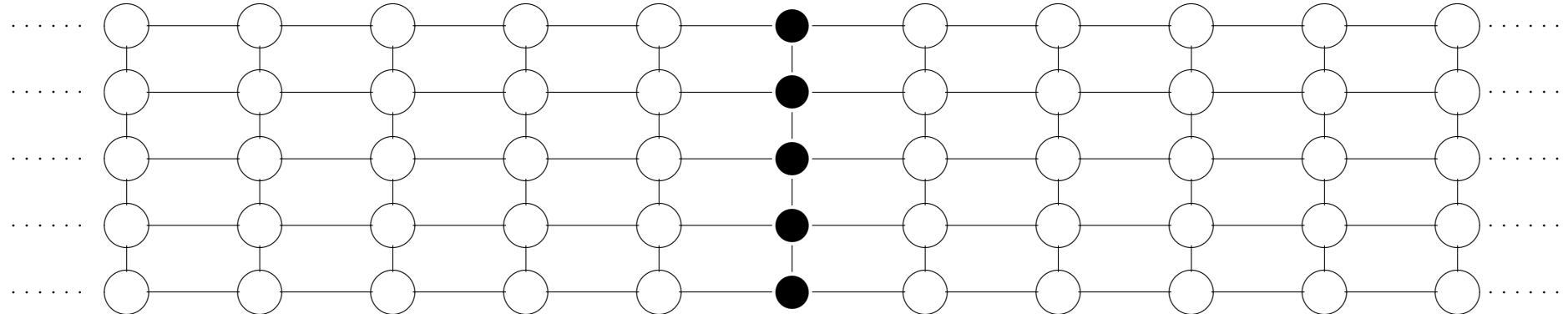
$$\delta E = -\frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta \mu \cosh \theta e^{-\mu L \cosh \theta} \sigma \left(\theta + \frac{i\pi}{2} \middle| \theta_1, \dots, \theta_r \right).$$

Bajnok-Janik formula:

$$QC_k^{(1)}(\theta_1, \dots, \theta_r) = QC_k^{(0)}(\theta_1, \dots, \theta_r) \{1 + i \delta \mathcal{R}_k\} = -1,$$

$$\begin{aligned} k &= 1, \dots, r, \\ \delta \mathcal{R}_k(\theta_1, \dots, \theta_r) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta e^{-\mu L \cosh \theta} \partial_k \sigma \left(\theta + \frac{i\pi}{2} \middle| \theta_1, \dots, \theta_r \right) \\ \partial_k &= \partial / \partial \theta_k. \end{aligned}$$

PRINCIPAL MODELS ($SU(n)$)



$SU(6)$ model TBA diagram

several types of particles: masses μ_a

transfer matrix: $\sigma^a(\theta|\theta_1, \dots, \theta_r)$

(all particles in $a = n - 1$ (anti-vector) representation)

Lüscher corrections to the energy:

$$E^{(1)} = E^{(0)} + \delta E$$

$$E^{(0)} = \sum_{j=1}^r \mu_{n-1} \cosh \theta_j$$

$$\delta E = i \sum_{a=1}^{n-1} \mu_a \int_{-\infty}^{\infty} \frac{d\theta}{2\pi} \sinh \left(\theta + \frac{i\pi}{n} \right) e^{i\mu_a L \sinh \left(\theta + \frac{i\pi}{n} \right)} \sigma^a \left(\theta + \frac{i\pi}{n} \middle| \theta_1, \dots, \theta_r \right)$$

quantization conditions: $(k = 1, \dots, r)$

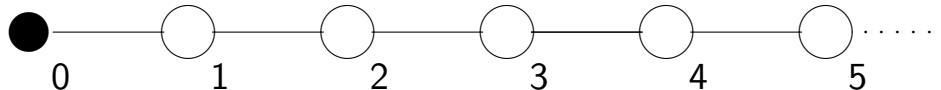
$$QC_k^{(1)} = QC_k^{(0)} \{1 + i \delta \mathcal{R}_k\}$$

$$= e^{i\mu_{n-1} L \sinh \theta_k} \sigma^{n-1}(\theta_k | \theta_1, \dots, \theta_r) \{1 + i \delta \mathcal{R}_k\} = -1,$$

Bajnok-Janik formula:

$$\delta \mathcal{R}_k(\theta_1, \dots, \theta_r) = \sum_{a=1}^{n-1} \int_{-\infty}^{\infty} \frac{d\theta}{2\pi} e^{i\mu_a L \sinh \left(\theta + \frac{i\pi}{n} \right)} \partial_k \sigma^a \left(\theta + \frac{i\pi}{n} \middle| \theta_1, \dots, \theta_r \right)$$

Example: $O(2)$ model



$O(2)$ model TBA diagram

BETHE-YANG LIMIT SOLUTION (BETHE ANSATZ):

$$T_m(\theta + \frac{i\pi}{2})T_m(\theta - \frac{i\pi}{2}) = T_{m-1}(\theta)T_{m+1}(\theta) + T_0(\theta + i\frac{m+1}{2}\pi)T_0(\theta - i\frac{m+1}{2}\pi)$$

$$m = 1, 2, \dots \quad (T_{-1}(\theta) \equiv 0 \text{ boundary cond.})$$

$$T_0(\theta) = \prod_k (\theta - \theta_k) \quad T_m(\theta) : \text{BA sol'n}$$

Bethe Ansatz

O(2) S-matrix: $S_{cd}^{ab}(\theta) = U(\theta) \{ \theta \delta^{ab} \delta_{cd} + (i\pi - \theta) \delta_d^a \delta_c^b \}$

$U(\theta)$: unitarity, crossing

Transfer matrix (3-particles)

$$T_{a_1 a_2 a_3}^{b_1 b_2 b_3}(\theta | \theta_1, \theta_2, \theta_3) = S_{u_3 a_1}^{u_1 b_1}(\theta - \theta_1) S_{u_1 a_2}^{u_2 b_2}(\theta - \theta_2) S_{u_2 a_3}^{u_3 b_3}(\theta - \theta_3)$$

Bethe roots: $\{u_j\}_{j=1}^M$

Baxter's Q operator: $Q(\theta) = \prod_{j=1}^M (\theta - u_j)$

Bethe equations:

$$\frac{Q(u_j+i\pi)}{Q(u_j-i\pi)} = -\frac{T_0(u_j+\frac{i\pi}{2})}{T_0(u_j-\frac{i\pi}{2})} \quad j = 1, \dots, M$$

$T_k(\theta)$ polynomials:

$$\xi(\theta) = \frac{T_0(\theta)}{Q(\theta+\frac{i\pi}{2}) Q(\theta-\frac{i\pi}{2})}$$

$$T_k(\theta) = Q(\theta + i \frac{k+1}{2} \pi) Q(\theta - i \frac{k+1}{2} \pi) \sum_{j=0}^k \xi(\theta + i \frac{k-2j}{2} \pi)$$

Y – system \iff T – system (gauge)

$$y_m(\theta) = \frac{T_{m-1}(\theta) T_{m+1}(\theta)}{T_0(\theta + i\frac{m+1}{2}\pi) T_0(\theta - i\frac{m+1}{2}\pi)}$$

$$Y_m(\theta) = 1 + y_m(\theta) = \frac{T_m(\theta + \frac{i\pi}{2}) T_m(\theta - \frac{i\pi}{2})}{T_0(\theta + i\frac{m+1}{2}\pi) T_0(\theta - i\frac{m+1}{2}\pi)}$$

Y-system equations:

$$y_m(\theta + \frac{i\pi}{2}) y_m(\theta - \frac{i\pi}{2}) = Y_{m-1}(\theta) Y_{m+1}(\theta) \quad m = 1, 2 \dots$$

y_m roots: T_{m+1} roots + T_{m-1} roots

only T_0 has roots: \implies only y_1 has roots: $\{\theta_k\}$

T_m : transfer matrix in spin $\frac{m}{2}$ representation (up to scale factor)

T_1 : transfer matrix in defining representation

$$\frac{T_1(\theta + \frac{i\pi}{2}) T_1(\theta - \frac{i\pi}{2})}{T_0(\theta + i\pi) T_0(\theta - i\pi)} = Y_1(\theta)$$

physical transfer matrix (unitary + crossing symmetric):

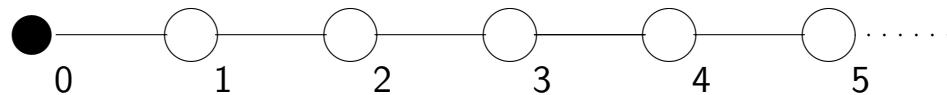
$$\Lambda(\theta) = \sigma(\theta + \frac{i\pi}{2})$$

$$\Lambda(\theta + \frac{i\pi}{2}) \Lambda(\theta - \frac{i\pi}{2}) = Y_1(\theta)$$

Bethe-Yang equations:

$$e^{i\ell \sinh \theta_k} \sigma(\theta_k) = -1 = e^{i\ell \sinh \theta_k} \Lambda(\theta_k - \frac{i\pi}{2})$$

EXACT TBA EQUATIONS:



$$y_m^e(\theta); \quad Y_m^e(\theta) = 1 + y_m^e(\theta); \quad L_m^e(\theta) = \ln Y_m^e(\theta); \quad m = 0, 1, \dots$$

exact Y-system equations:

$$y_m^e\left(\theta + \frac{i\pi}{2}\right) y_m^e\left(\theta - \frac{i\pi}{2}\right) = Y_{m-1}^e(\theta) Y_{m+1}^e(\theta) \quad m = 0, 1, \dots$$

$(Y_{-1}^e(\theta) \equiv 1$ by convention)

Asymptotic condition: $y_0^e(\theta) = e^{-\ell \cosh \theta} \Lambda^e(\theta)$

$$\Lambda^e\left(\theta + \frac{i\pi}{2}\right) \Lambda^e\left(\theta - \frac{i\pi}{2}\right) = Y_1^e(\theta)$$

Quantization conditions (QC): $y_1^e(\theta_k) = 0 \Rightarrow Y_0^e(\theta_k \pm \frac{i\pi}{2}) = 0$

TBA Lemma

$$f(\theta + \frac{i\pi}{2}) f(\theta - \frac{i\pi}{2}) = F(\theta)$$

$\{t_k\}$ roots in physical strip: $f(t_k) = 0$

$$f(\theta) = \left\{ \prod_k \tau(\theta - t_k) \right\} \exp \{(K \star \ln F)(\theta)\}$$

$$\tau(\theta) = \tanh \frac{\theta}{2} \quad \frac{\tau'(\theta)}{\tau(\theta)} = \frac{1}{\sinh \theta}$$

$$K(\theta) = \frac{1}{\cosh \theta} \quad (A \star B)(x) = \int_{-\infty}^{\infty} dy A(x-y)B(y)$$

Exact TBA equations:

$$\Lambda^e(\theta) = \exp \{(K \star L_1^e)(\theta)\}$$

$$y_1^e(\theta) = \left\{ \prod_k \tau(\theta - \theta_k) \right\} \exp \{(K \star L_0^e)(\theta) + (K \star L_2^e)(\theta)\}$$

$$y_m^e(\theta) = \exp \{(K \star L_{m+1}^e)(\theta) + (K \star L_{m-1}^e)(\theta)\} \quad m = 2, 3, \dots$$

Exact Bethe-Yang equations (QC):

$$e^{i\ell \sinh \theta_k} \Lambda^e(\theta_k - \frac{i\pi}{2}) = -1$$

$$\Lambda^e(\theta_k - \frac{i\pi}{2}) = e^{i\mathcal{B}^e(\theta_k)}$$

$$\mathcal{B}^e(x) = -\frac{\mathcal{P}}{2\pi} \int_{-\infty}^{\infty} dy \frac{L_1^e(y)}{\sinh(y-x)} \quad \mathcal{P} : \text{superfluous}$$

QC (exact Bethe-Yang equation):

$$\mathcal{B}^e(\theta_k) + \ell \sinh \theta_k = 2\pi I_k \quad (\text{half integer})$$

Energy:

$$\epsilon = \sum_k \cosh \theta_k - \frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta \cosh \theta L_0^e(\theta)$$

Back to B-Y limit:

$$\begin{aligned}\Lambda(\theta) &= \exp \{(K \star L_1)(\theta)\} \\ y_1(\theta) &= \left\{ \prod_k \tau(\theta - \theta_k) \right\} \exp \{(K \star L_2)(\theta)\} \\ y_m(\theta) &= \exp \{(K \star L_{m+1})(\theta) + (K \star L_{m-1})(\theta)\} \quad m = 2, 3, \dots\end{aligned}$$

$$\Lambda(\theta - \frac{i\pi}{2}) = \sigma(\theta) = e^{i\mathcal{B}(\theta)}$$

B-Y equations:

$$\mathcal{B}(\theta_k) + \ell \sinh \theta_k = 2\pi I_k$$

LINEARIZED TBA EQUATIONS:

$$f^e = f + \delta f$$

$$\frac{\delta y_1}{y_1} = K \star L_0^e + K \star \delta L_2$$

$$\frac{\delta y_m}{y_m} = K \star \delta L_{m+1} + K \star \delta L_{m-1} \quad m = 2, 3, \dots$$

$$L_0^e(\theta) = \ln [1 + e^{-\ell \cosh \theta} \Lambda^e(\theta)] \approx e^{-\ell \cosh \theta} \Lambda(\theta) = \lambda(\theta)$$

Lüscher's F-term:

$$-\frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta \cosh \theta L_0^e(\theta) \approx -\frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta \cosh \theta e^{-\ell \cosh \theta} \Lambda(\theta)$$

$$= -\frac{1}{2\pi} \int_{-\infty}^{\infty} d\theta \cosh \theta e^{-\ell \cosh \theta} \sigma(\theta + \frac{i\pi}{2})$$

Structure of linear system: $\mathbf{M}\xi = J$

$$\mathbf{M} = \begin{pmatrix} D_1 & -K\star & 0 & 0 & \dots \\ -K\star & D_2 & -K\star & 0 & \dots \\ 0 & -K\star & D_3 & -K\star & \dots \\ 0 & 0 & -K\star & D_4 & \dots \\ & & & \vdots & \end{pmatrix}$$

$$\xi = \begin{pmatrix} \delta L_1 \\ \delta L_2 \\ \delta L_3 \\ \vdots \end{pmatrix} \quad J = \begin{pmatrix} j(\theta) \\ 0 \\ 0 \\ \vdots \end{pmatrix}$$

$$j(\theta) = K \star \lambda$$

Bethe-Yang equation:

$$\mathcal{B}(\theta_k) + \delta\mathcal{B}(\theta_k) + \ell \sinh \theta_k = 2\pi I_k$$

Bajnok-Janik:

$$\delta\mathcal{B}_k = \delta\mathcal{B}(\theta_k) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} dy \frac{\delta L_1(y)}{\sinh(y-\theta_k)}$$

$$\delta\mathcal{B}_k = -\frac{1}{2\pi} \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dx \frac{R_{11}(y,x) j(x)}{\sinh(y-\theta_k)}$$

Inverse operator:

$$\mathbf{M}^{-1} = \mathbf{R} \quad \xi_m(x) = \sum_s \int_{-\infty}^{\infty} dy R_{ms}(x,y) J_s(y)$$

$$\mathbf{M} = \mathbf{M}^T \quad \text{symmetric:} \quad \mathbf{R} = \mathbf{R}^T$$

$$R_{ms}(x,y) = R_{sm}(y,x)$$

DERIVATIVE OF B-Y LIMIT:

$$\partial_k = \frac{\partial}{\partial \theta_k}$$

$$\frac{\partial_k \Lambda}{\Lambda} = K \star \partial_k L_1$$

$$\frac{\partial_k y_1}{y_1} = -\frac{1}{\sinh(\theta-\theta_k)} + K \star \partial_k L_2$$

$$\frac{\partial_k y_m}{y_m} = K \star \partial_k L_{m+1} + K \star \partial_k L_{m-1} \quad m = 2, 3, \dots$$

same linear system!

$$\xi = \begin{pmatrix} \partial_k L_1 \\ \partial_k L_2 \\ \vdots \end{pmatrix} \quad J = \begin{pmatrix} j_k(\theta) \\ 0 \\ \vdots \end{pmatrix} \quad j_k(\theta) = -\frac{1}{\sinh(\theta-\theta_k)}$$

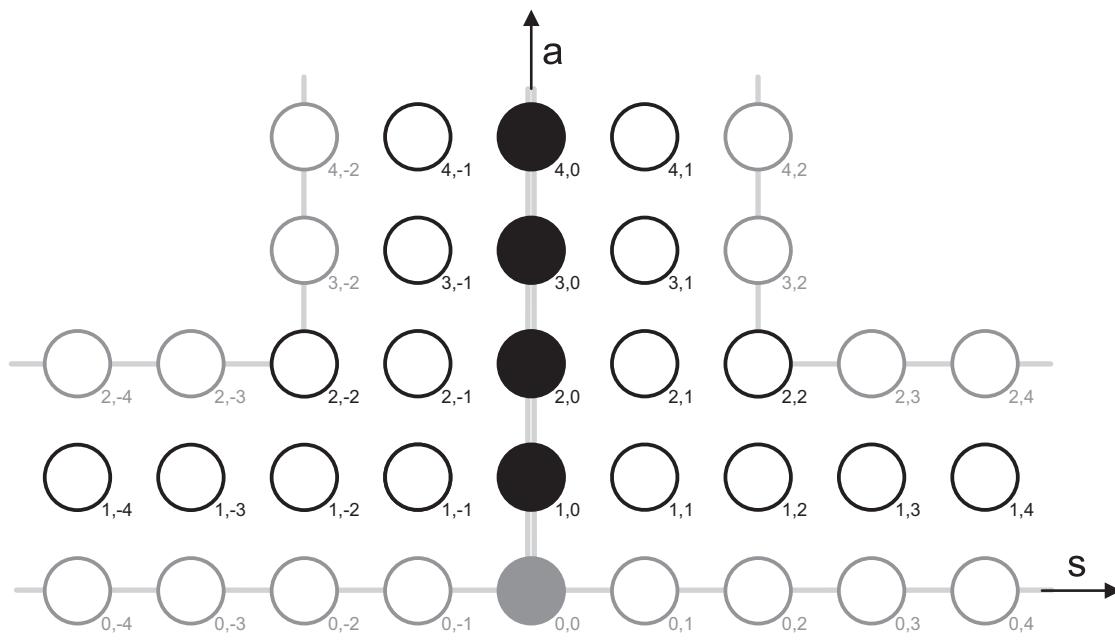
Relation from the derivative case:

$$\partial_k L_m(x) = - \int_{-\infty}^{\infty} dy R_{m1}(x, y) \frac{1}{\sinh(y - \theta_k)}$$

Huge piece of luck: just what we need!

$$\begin{aligned}\delta \mathcal{B}_k &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx \partial_k L_1(x) j(x) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx \partial_k L_1(x) (K \star \lambda)(x) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx \lambda(x) (K \star \partial_k L_1)(x) \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} dx \lambda(x) \frac{\partial_k \Lambda(x)}{\Lambda(x)} = \frac{1}{2\pi} \partial_k \int_{-\infty}^{\infty} dx e^{-\ell \cosh x} \Lambda(x)\end{aligned}$$

Example: Konishi operator



$AdS_5 \times S_5$ string model TBA diagram

Gromov, Kazakov, Kozak, Vieira '09

Arutyunov, Frolov '09

$AdS_5 \times S_5$ STRING MODEL

String energies:

$$E = J + \sum_{i=1}^N \mathcal{E}(p_i) - \frac{1}{2\pi} \sum_{Q=1}^{\infty} \int_{-\infty}^{\infty} du \frac{d\tilde{p}^Q}{du} \log(1 + Y_Q)$$

Dispersion relation:

$$\mathcal{E}(p) = \sqrt{1 + 4g^2 \sin^2 \frac{p}{2}}$$

Leading order “wrapping” corrections:

$$\pi(2n_k + 1) = J p_k + i \sum_{j=1}^N \log S_{\mathfrak{sl}(2)}^{1*1*}(u_j, u_k) + \delta\mathcal{R}_k^{(\text{BJ})} + O(g^9)$$

Bajnok-Janik formula:

$$\delta\mathcal{R}_k^{(\text{BJ})} = \frac{1}{2\pi} \sum_{Q=1}^{\infty} \int_{-\infty}^{\infty} du \left. \frac{\partial}{\partial u_k} Y_Q^o(u) \right|_{\{u_j\} = \{u_j^o\}}$$

The exact Bethe equations

$$\begin{aligned}
& \pi i(2n_k + 1) = \log Y_{1*}(u_k) = iL p_k - \sum_{j=1}^N \log S_{\mathfrak{sl}(2)}^{1*1*}(u_j, u_k) \\
& + 2 \sum_{j=1}^N \log \text{Res}(S) \star K_{vwx}^{11*}(u_j^-, u_k) - 2 \sum_{j=1}^N \log \left(u_j - u_k - \frac{2i}{g} \right) \frac{x_j^- - \frac{1}{x_k^-}}{x_j^- - x_k^+} \\
& - 2 \sum_{j=1}^{n_1} \left(\log S \hat{\star} K_{y1*}(r_j^{(1)-}, u_k) - \log S(r_j^{(1)} - u_k) \right) \\
& + \log (1 + Y_Q) \star \left(K_{\mathfrak{sl}(2)}^{Q1*} + 2 s \star K_{vwx}^{Q-1,1*} \right) \\
& + 2 \log (1 + Y_{1|vw}) \star (s \hat{\star} K_{y1*} + \tilde{s}) - 2 \log \frac{1 - Y_-}{1 - Y_+} \hat{\star} s \star K_{vwx}^{11*} \\
& + \log \frac{1 - \frac{1}{Y_-}}{1 - \frac{1}{Y_+}} \hat{\star} K_1 + \log \left(1 - \frac{1}{Y_-} \right) \left(1 - \frac{1}{Y_+} \right) \hat{\star} K_{y1*}
\end{aligned}$$

LINEARIZED STRING TBA EQUATIONS

Arutyunov, Frolov, Suzuki '10

$$Y = Y^o(1 + \mathcal{Y}).$$

- **$m|w$ -strings:** $m \geq 1$, $\mathcal{Y}_{0|w} = 0$

$$\mathcal{Y}_{m|w} = (A_{m-1|w}\mathcal{Y}_{m-1|w} + A_{m+1|w}\mathcal{Y}_{m+1|w}) \star s + \delta_{m1} \left(\frac{\mathcal{Y}_+}{1-Y_+^o} - \frac{\mathcal{Y}_-}{1-Y_-^o} \right) \hat{s}$$

coefficients: $A_{m|w} = \frac{Y_{m|w}^o}{1+Y_{m|w}^o}$

- **$m|vw$ -strings:** $m \geq 1$, $\mathcal{Y}_{0|vw} = 0$

$$\mathcal{Y}_{m|vw} = (A_{m-1|vw}\mathcal{Y}_{m-1|vw} + A_{m+1|vw}\mathcal{Y}_{m+1|vw}) \star s - Y_{m+1}^o \star s$$

$$+ \delta_{m1} \left(\frac{\mathcal{Y}_-}{1-\frac{1}{Y_-^o}} - \frac{\mathcal{Y}_+}{1-\frac{1}{Y_+^o}} \right) \hat{s}$$

coefficients: $A_{m|vw} = \frac{Y_{m|vw}^o}{1+Y_{m|vw}^o}$

- y -particles:

$$\mathcal{Y}_+ - \mathcal{Y}_- = Y_Q^o \star K_{Qy}$$

Perturbative orders:

$$Y_Q^o(u) = O(g^8)$$

$$\mathcal{Y}_+ - \mathcal{Y}_- = O(g^9)$$

Y_+^o and Y_-^o coincide at leading order:

$$\frac{Y_+^o(u)}{Y_-^o(u)} = 1 + O(g^2).$$

Decoupling:

$$\mathcal{Y}_{m|w} = O(g^9) \quad \mathcal{Y}_{m|vw} = O(g^8)$$

Leading order (rescaled):

$$\frac{\delta y_m}{y_m} - K \star (\delta L_{m-1} + \delta L_{m+1}) = -K \star Y_{m+1}^o, \quad m = 1, 2, \dots$$

$Y_{m|vw}^o = y_m$: XXX model $Q(u) = u$ solution $N = 2$

More general states in the sl_2 sector: $N > 2$ (roots for all y_m)

wrapping corrected (leading order) Bethe equations

$$\pi(2n_k + 1) = J p_k + i \sum_{j=1}^N \log S_{\mathfrak{sl}(2)}^{1*1*} \left(\frac{u_j}{g}, \frac{u_k}{g} \right) + \delta \mathcal{R}_k + O(g^9)$$

$$\delta \mathcal{R}_k = \delta \mathcal{R}_k^{(1)} + \delta \mathcal{R}_k^{(2)} + \delta \mathcal{R}_k^{(3)}$$

$\delta\mathcal{R}_k^{(1)}$ & $\delta\mathcal{R}_k^{(3)}$: small g expansion of the Y_Q convolution terms

$\delta\mathcal{R}_k^{(2)}$: convolution terms containing $Y_{1|vw}$

$$\delta\mathcal{R}_k^{(1)} = \frac{1}{\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_m^o(u) \frac{u-u_k}{(m+1)^2 + (u-u_k)^2}$$

$$\delta\mathcal{R}_k^{(2)} = \int_{-\infty}^{\infty} du \frac{\delta L_1(u)}{2 \sinh \frac{\pi}{2}(u-u_k)}$$

$$\delta\mathcal{R}_k^{(3)} = \frac{1}{\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_{m+1}^o(u) \left\{ \mathcal{F}_m(u - u_k) - \frac{u-u_k}{m^2 + (u-u_k)^2} \right\}$$

$$\mathcal{F}_m(u) = \frac{-i}{4} \left\{ \psi\left(\frac{m+iu}{4}\right) - \psi\left(\frac{m-iu}{4}\right) - \psi\left(\frac{m+2+iu}{4}\right) + \psi\left(\frac{m+2-iu}{4}\right) \right\}$$

$$\psi(z) = \Gamma'(z)/\Gamma(z)$$

Nontrivial part: $\delta\mathcal{R}_k^{(2)}$

Same problem solved for O(2) model !

$$\begin{aligned}\delta\mathcal{R}_k^{(2)} &= \frac{1}{\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_{m+1}^o(u) \left\{ \partial_k \log \hat{t}_m(u) \right\} \\ &= \frac{1}{\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_{m+1}^o(u) \left\{ \partial_k \log t_m(u) - \frac{r'_m(u-u_k)}{r_m(u-u_k)} \right\}. \\ \frac{r'_m(x)}{r_m(x)} &= \mathcal{F}_m(x) - \frac{2x}{m^2+x^2}\end{aligned}$$

transcendental parts cancel

result: rational functions

$$\delta\mathcal{R}_k = \frac{1}{\pi} \int_{-\infty}^{\infty} du Y_1^o(u) \frac{u-u_k}{4+(u-u_k)^2}$$

$$+ \frac{1}{\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_{m+1}^o(u) \left\{ \partial_k \log t_m(u) + \frac{u-u_k}{m^2+(u-u_k)^2} + \frac{u-u_k}{(m+2)^2+(u-u_k)^2} \right\}$$

can be written compactly

$$\delta\mathcal{R}_k = \frac{1}{2\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du Y_m^o(u) \partial_k \log j_m(u)$$

$$j_m(u) = \frac{64g^8}{(u^2+m^2)^4} \frac{t_{m-1}^2(u)}{t_0(u-im-i) t_0(u-im+i) t_0(u+im-i) t_0(u+im+i)}$$

important observation:

$$Y_m^o(u) = j_m(u)$$

Bajnok-Janik formula reproduced:

$$\delta\mathcal{R}_k = \delta\mathcal{R}_k^{(\text{BJ})} = \frac{1}{2\pi} \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} du \partial_k Y_m^o(u)$$