Keller-Segel, Fast Diffusion and Functional Inequalities

J. A. Carrillo

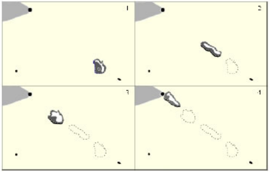
ICREA - Universitat Autònoma de Barcelona

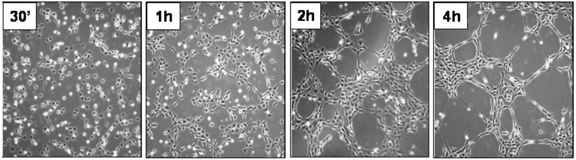
Benasque, Spain, 2011

Outline

Outline

Chemotaxis





Cell movement and aggregation by chemical interaction.

PKS System

$$
\begin{cases}
\dfrac{\partial \rho}{\partial t}(x,t) = \Delta \rho(x,t) - \nabla \cdot (\rho(x,t) \nabla c(x,t)) & x \in \mathbb{R}^2\,,\ t > 0\,, \\[2ex]
c(x,t) = -\dfrac{1}{2\pi} \displaystyle\int_{\mathbb{R}^2} \log|x-y|\, \rho(y,t)\, dy & x \in \mathbb{R}^2\,,\ t > 0\,, \\[2ex]
\rho(x, t=0) = \rho_0 \geq 0 & x \in \mathbb{R}^2\,.
\end{cases}
$$

Huge Literature: Horstmann reviews (2003& 2004), Perthame review (2004).

Conservations:

- Conservation of mass:

$$M := \int_{\mathbb{R}^2} \rho_0(x)\, dx = \int_{\mathbb{R}^2} \rho(x,t)\, dx$$

- Conservation of center of mass:

$$M_1 := \int_{\mathbb{R}^2} x\, \rho_0(x)\, dx = \int_{\mathbb{R}^2} x\, \rho(x,t)\, dx\,.$$

PKS System

$$\begin{cases} \dfrac{\partial \rho}{\partial t}(x,t) = \Delta \rho(x,t) - \nabla \cdot (\rho(x,t) \nabla c(x,t)) & x \in \mathbb{R}^2 \,,\ t > 0 \,, \\ c(x,t) = -\dfrac{1}{2\pi} \displaystyle\int_{\mathbb{R}^2} \log|x-y| \, \rho(y,t) \, dy & x \in \mathbb{R}^2 \,,\ t > 0 \,, \\ \rho(x, t=0) = \rho_0 \geq 0 & x \in \mathbb{R}^2 \,. \end{cases}$$

Huge Literature: Horstmann reviews (2003& 2004), Perthame review (2004).

Conservations:

- Conservation of mass:

$$M := \int_{\mathbb{R}^2} \rho_0(x) \, dx = \int_{\mathbb{R}^2} \rho(x,t) \, dx$$

- Conservation of center of mass:

$$M_1 := \int_{\mathbb{R}^2} x \, \rho_0(x) \, dx = \int_{\mathbb{R}^2} x \, \rho(x,t) \, dx \,.$$

PKS System

$$\begin{cases} \dfrac{\partial \rho}{\partial t}(x,t) = \Delta \rho(x,t) - \nabla \cdot (\rho(x,t) \nabla c(x,t)) & x \in \mathbb{R}^2 \,,\; t > 0 \,, \\ c(x,t) = -\dfrac{1}{2\pi} \displaystyle\int_{\mathbb{R}^2} \log|x-y| \, \rho(y,t) \, dy & x \in \mathbb{R}^2 \,,\; t > 0 \,, \\ \rho(x, t=0) = \rho_0 \geq 0 & x \in \mathbb{R}^2 \,. \end{cases}$$

Huge Literature: Horstmann reviews (2003& 2004), Perthame review (2004).

Conservations:

- Conservation of mass:

$$M := \int_{\mathbb{R}^2} \rho_0(x) \, dx = \int_{\mathbb{R}^2} \rho(x,t) \, dx$$

- Conservation of center of mass:

$$M_1 := \int_{\mathbb{R}^2} x \, \rho_0(x) \, dx = \int_{\mathbb{R}^2} x \, \rho(x,t) \, dx \,.$$

Second Moment

Distributional Solution:

We shall say that $\rho \in C^0\big([0,T); L^1_{\text{weak}}(\mathbb{R}^2)\big)$ is a weak solution to the PKS system if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$,

$$\frac{d}{dt}\int_{\mathbb{R}^2} \psi(x)\,\rho(x,t)\,dx =$$

$$\int_{\mathbb{R}^2} \Delta\psi(x)\,\rho(x,t)\,dx - \frac{1}{4\pi}\int_{\mathbb{R}^2}\int_{\mathbb{R}^2} [\nabla\psi(x) - \nabla\psi(y)]\cdot\frac{x-y}{|x-y|^2}\,\rho(x,t)\,\rho(y,t)\,dx\,dy$$

holds in the distributional sense in $(0,T)$ and $\rho(0) = \rho_0$.

Evolution of second moment:

$$\frac{d}{dt}\int_{\mathbb{R}^2} |x|^2\,\rho(x,t)\,dx = 4M - \frac{1}{2\pi}M^2,$$

Struggle between diffusion and aggregation. Balance between these two mechanisms happens precisely at the critical mass $M = 8\,\pi$.

Second Moment

Distributional Solution:

We shall say that $\rho \in C^0\left([0,T); L^1_{\text{weak}}(\mathbb{R}^2)\right)$ is a weak solution to the PKS system if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$,

$$\frac{d}{dt}\int_{\mathbb{R}^2} \psi(x)\,\rho(x,t)\,dx =$$

$$\int_{\mathbb{R}^2} \Delta\psi(x)\,\rho(x,t)\,dx - \frac{1}{4\pi}\int_{\mathbb{R}^2}\int_{\mathbb{R}^2} [\nabla\psi(x) - \nabla\psi(y)]\cdot\frac{x-y}{|x-y|^2}\,\rho(x,t)\,\rho(y,t)\,dx\,dy$$

holds in the distributional sense in $(0,T)$ and $\rho(0) = \rho_0$.

Evolution of second moment:

$$\frac{d}{dt}\int_{\mathbb{R}^2} |x|^2\,\rho(x,t)\,dx = 4M - \frac{1}{2\pi}M^2,$$

Struggle between diffusion and aggregation. Balance between these two mechanisms happens precisely at the critical mass $M = 8\,\pi$.

Second Moment

Distributional Solution:

We shall say that $\rho \in C^0\left([0,T); L^1_{\text{weak}}(\mathbb{R}^2)\right)$ is a weak solution to the PKS system if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$,

$$\frac{d}{dt}\int_{\mathbb{R}^2} \psi(x)\,\rho(x,t)\,dx =$$

$$\int_{\mathbb{R}^2} \Delta\psi(x)\,\rho(x,t)\,dx - \frac{1}{4\pi}\int_{\mathbb{R}^2}\int_{\mathbb{R}^2} [\nabla\psi(x) - \nabla\psi(y)]\cdot\frac{x-y}{|x-y|^2}\,\rho(x,t)\,\rho(y,t)\,dx\,dy$$

holds in the distributional sense in $(0,T)$ and $\rho(0) = \rho_0$.

Evolution of second moment:

$$\frac{d}{dt}\int_{\mathbb{R}^2} |x|^2\,\rho(x,t)\,dx = 4M - \frac{1}{2\pi}M^2,$$

Struggle between diffusion and aggregation. Balance between these two mechanisms happens precisely at the critical mass $M = 8\,\pi$.

Second Moment

Distributional Solution:

We shall say that $\rho \in C^0\left([0,T); L^1_{\text{weak}}(\mathbb{R}^2)\right)$ is a weak solution to the PKS system if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$,

$$\frac{d}{dt}\int_{\mathbb{R}^2} \psi(x)\,\rho(x,t)\,dx =$$

$$\int_{\mathbb{R}^2} \Delta\psi(x)\,\rho(x,t)\,dx - \frac{1}{4\pi}\int_{\mathbb{R}^2}\int_{\mathbb{R}^2} [\nabla\psi(x) - \nabla\psi(y)]\cdot\frac{x-y}{|x-y|^2}\,\rho(x,t)\,\rho(y,t)\,dx\,dy$$

holds in the distributional sense in $(0,T)$ and $\rho(0) = \rho_0$.

Evolution of second moment:

$$\frac{d}{dt}\int_{\mathbb{R}^2} |x|^2\,\rho(x,t)\,dx = 4M - \frac{1}{2\pi}M^2,$$

Struggle between diffusion and aggregation. Balance between these two mechanisms happens precisely at the critical mass $M = 8\,\pi$.

Second Moment

Distributional Solution:

We shall say that $\rho \in C^0\left([0,T); L^1_{\text{weak}}(\mathbb{R}^2)\right)$ is a weak solution to the PKS system if for all test functions $\psi \in \mathcal{D}(\mathbb{R}^2)$,

$$\frac{d}{dt}\int_{\mathbb{R}^2} \psi(x)\,\rho(x,t)\,dx =$$

$$\int_{\mathbb{R}^2} \Delta\psi(x)\,\rho(x,t)\,dx - \frac{1}{4\pi}\int_{\mathbb{R}^2}\int_{\mathbb{R}^2} [\nabla\psi(x) - \nabla\psi(y)]\cdot\frac{x-y}{|x-y|^2}\,\rho(x,t)\,\rho(y,t)\,dx\,dy$$

holds in the distributional sense in $(0,T)$ and $\rho(0) = \rho_0$.

Evolution of second moment:

$$\frac{d}{dt}\int_{\mathbb{R}^2} |x|^2\,\rho(x,t)\,dx = 4M - \frac{1}{2\pi}M^2,$$

Struggle between diffusion and aggregation. Balance between these two mechanisms happens precisely at the critical mass $M = 8\,\pi$.

Cases

PKS Cases:

- **Subcritical Case, $M < 8\,\pi$:** Global existence: Jägger-Luckhaus (1992), Dolbeault-Perthame (2004), Blanchet-Dolbeault-Perthame (2006), Blanchet-Calvez-C. (2008).
- **Supercritical Case, $M > 8\,\pi$:** Blow-up: Herrero-Velazquez (1996), Velazquez (2002-2004), Dolbeault-Schmeiser (2009).
- **Critical Case, $M = 8\,\pi$:** Infinite-time aggregation, infinitely many stationary states: Biler-Karch-Laurençot-Nadzieja (2006), Blanchet-C.-Masmoudi (2008), Blanchet-Carlen-C. (2011), Carlen-C.-Loss (2010), Carlen-Figalli (preprint).

Cases

PKS Cases:

- **Subcritical Case, $M < 8\pi$:** Global existence: Jägger-Luckhaus (1992), Dolbeault-Perthame (2004), Blanchet-Dolbeault-Perthame (2006), Blanchet-Calvez-C. (2008).
- **Supercritical Case, $M > 8\pi$:** Blow-up: Herrero-Velazquez (1996), Velazquez (2002-2004), Dolbeault-Schmeiser (2009).
- **Critical Case, $M = 8\pi$:** Infinite-time aggregation, infinitely many stationary states: Biler-Karch-Laurençot-Nadzieja (2006), Blanchet-C.-Masmoudi (2008), Blanchet-Carlen-C. (2011), Carlen-C.-Loss (2010), Carlen-Figalli (preprint).

Cases

PKS Cases:

- **Subcritical Case, $M < 8\,\pi$:** Global existence: Jägger-Luckhaus (1992), Dolbeault-Perthame (2004), Blanchet-Dolbeault-Perthame (2006), Blanchet-Calvez-C. (2008).
- **Supercritical Case, $M > 8\,\pi$:** Blow-up: Herrero-Velazquez (1996), Velazquez (2002-2004), Dolbeault-Schmeiser (2009).
- **Critical Case, $M = 8\,\pi$:** Infinite-time aggregation, infinitely many stationary states: Biler-Karch-Laurençot-Nadzieja (2006), Blanchet-C.-Masmoudi (2008), Blanchet-Carlen-C. (2011), Carlen-C.-Loss (2010), Carlen-Figalli (preprint).

Outline

Log HLS Inequality by Carlen& Loss

Let f be a non-negative function in $L^1(\mathbb{R}^2)$ with mass M such that $f \log f$ and $f \log(e+|x|^2)$ belong to $L^1(\mathbb{R}^2)$. Then

$$\int_{\mathbb{R}^2} f(x) \log f(x)\, \mathrm{d}x + \frac{2}{M} \iint_{\mathbb{R}^2\times\mathbb{R}^2} f(x) f(y) \log|x-y|\, \mathrm{d}x\, \mathrm{d}y \geq -\, C(M)$$

with $C(M) := M(1+\log\pi - \log M)$.

Equality cases:

There is equality if and only if $f(x) = \bar\rho_\lambda(x-x_0)$ for some $\lambda > 0$ and some $x_0 \in \mathbb{R}^2$, where

$$\bar\rho_\lambda(x) := \frac{M}{\pi} \frac{\lambda}{(\lambda+|x|^2)^2}$$

Natural Liapunov Functional:

Free energy:

$$\mathcal{F}_{\mathrm{PKS}}[\rho] := \int_{\mathbb{R}^2} \rho(x) \log \rho(x)\, \mathrm{d}x + \frac{1}{4\pi} \iint_{\mathbb{R}^2\times\mathbb{R}^2} \rho(x)\, \rho(y)\, \log|x-y|\, \mathrm{d}x\, \mathrm{d}y$$

Log HLS Inequality by Carlen& Loss

Let f be a non-negative function in $L^1(\mathbb{R}^2)$ with mass M such that $f \log f$ and $f \log(e + |x|^2)$ belong to $L^1(\mathbb{R}^2)$. Then

$$\int_{\mathbb{R}^2} f(x) \log f(x)\, \mathrm{d}x + \frac{2}{M} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) f(y) \log|x-y|\, \mathrm{d}x\, \mathrm{d}y \geq -\, C(M)$$

with $C(M) := M(1 + \log \pi - \log M)$.

Equality cases:

There is equality if and only if $f(x) = \bar{\rho}_\lambda(x - x_0)$ for some $\lambda > 0$ and some $x_0 \in \mathbb{R}^2$, where

$$\bar{\rho}_\lambda(x) := \frac{M}{\pi} \frac{\lambda}{(\lambda + |x|^2)^2}$$

Natural Liapunov Functional:

Free energy:

$$\mathcal{F}_{\mathrm{PKS}}[\rho] := \int_{\mathbb{R}^2} \rho(x) \log \rho(x)\, \mathrm{d}x + \frac{1}{4\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} \rho(x)\, \rho(y)\, \log|x-y|\, \mathrm{d}x\, \mathrm{d}y$$

Log HLS Inequality by Carlen& Loss

Let f be a non-negative function in $L^1(\mathbb{R}^2)$ with mass M such that $f \log f$ and $f \log(e + |x|^2)$ belong to $L^1(\mathbb{R}^2)$. Then

$$\int_{\mathbb{R}^2} f(x) \log f(x)\, \mathrm{d}x + \frac{2}{M} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) f(y) \log|x-y|\, \mathrm{d}x\, \mathrm{d}y \geq -\, C(M)$$

with $C(M) := M(1 + \log \pi - \log M)$.

Equality cases:

There is equality if and only if $f(x) = \bar{\rho}_\lambda(x - x_0)$ for some $\lambda > 0$ and some $x_0 \in \mathbb{R}^2$, where

$$\bar{\rho}_\lambda(x) := \frac{M}{\pi} \frac{\lambda}{(\lambda + |x|^2)^2}$$

Natural Liapunov Functional:

Free energy:

$$\mathcal{F}_{\mathrm{PKS}}[\rho] := \int_{\mathbb{R}^2} \rho(x) \log \rho(x)\, \mathrm{d}x + \frac{1}{4\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} \rho(x)\, \rho(y)\, \log|x-y|\, \mathrm{d}x\, \mathrm{d}y$$

Formal Gradient Flow

A formal calculation shows that for all $u \in C_c^\infty(\mathbb{R}^2)$ with zero mean,

$$\lim_{\epsilon \to 0} \frac{1}{\epsilon} \left(\mathcal{F}_{\text{PKS}}[\rho + \epsilon u] - \mathcal{F}_{\text{PKS}}[\rho] \right) = \int_{\mathbb{R}^2} \frac{\delta \mathcal{F}_{\text{PKS}}(\rho)}{\delta \rho}(x)\, u(x)\, \mathrm{d}x$$

where

$$\frac{\delta \mathcal{F}_{\text{PKS}}(\rho)}{\delta \rho}(x) := \log \rho(x) + \frac{1}{2\pi} \int_{\mathbb{R}^2} \log|x-y| \rho(y)\, \mathrm{d}y = \log \rho(x) - G * \rho(x)\,.$$

The PKS equation can be rewritten as

$$\frac{\partial \rho}{\partial t}(t,x) = \operatorname{div}\left(\rho(t,x) \nabla \left[\frac{\delta \mathcal{F}_{\text{PKS}}(\rho(t))}{\delta \rho}(t,x) \right] \right)\,.$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{F}_{\text{PKS}}[\rho(t)] = -\int_{\mathbb{R}^2} \rho(t,x) \left| \nabla \frac{\delta \mathcal{F}_{\text{PKS}}(\rho(t))}{\delta \rho}(t,x) \right|^2 \mathrm{d}x\,.$$

Formal Gradient Flow

A formal calculation shows that for all $u \in C_c^\infty(\mathbb{R}^2)$ with zero mean,

$$\lim_{\epsilon \to 0} \frac{1}{\epsilon} \left(\mathcal{F}_{\text{PKS}}[\rho + \epsilon u] - \mathcal{F}_{\text{PKS}}[\rho] \right) = \int_{\mathbb{R}^2} \frac{\delta \mathcal{F}_{\text{PKS}}(\rho)}{\delta \rho}(x)\, u(x)\, \mathrm{d}x$$

where

$$\frac{\delta \mathcal{F}_{\text{PKS}}(\rho)}{\delta \rho}(x) := \log \rho(x) + \frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x-y| \rho(y)\, \mathrm{d}y = \log \rho(x) - G * \rho(x)\ .$$

The PKS equation can be rewritten as

$$\frac{\partial \rho}{\partial t}(t,x) = \operatorname{div} \left(\rho(t,x) \nabla \left[\frac{\delta \mathcal{F}_{\text{PKS}}(\rho(t))}{\delta \rho}(t,x) \right] \right)\ .$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{F}_{\text{PKS}}[\rho(t)] = -\int_{\mathbb{R}^2} \rho(t,x) \left| \nabla \frac{\delta \mathcal{F}_{\text{PKS}}(\rho(t))}{\delta \rho}(t,x) \right|^2 \mathrm{d}x\ .$$

Formal Gradient Flow

A formal calculation shows that for all $u \in C_c^\infty(\mathbb{R}^2)$ with zero mean,

$$\lim_{\epsilon\to 0} \frac{1}{\epsilon} \left(\mathcal{F}_{\rm PKS}[\rho+\epsilon u] - \mathcal{F}_{\rm PKS}[\rho]\right) = \int_{\mathbb{R}^2} \frac{\delta \mathcal{F}_{\rm PKS}(\rho)}{\delta \rho}(x)\, u(x)\, \mathrm{d}x$$

where

$$\frac{\delta \mathcal{F}_{\rm PKS}(\rho)}{\delta \rho}(x) := \log \rho(x) + \frac{1}{2\pi}\int_{\mathbb{R}^2} \log|x-y|\rho(y)\, \mathrm{d}y = \log\rho(x) - G*\rho(x)\,.$$

The PKS equation can be rewritten as

$$\frac{\partial \rho}{\partial t}(t,x) = \mathrm{div}\left(\rho(t,x)\nabla\left[\frac{\delta \mathcal{F}_{\rm PKS}(\rho(t))}{\delta\rho}(t,x)\right]\right)\,.$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{F}_{\rm PKS}[\rho(t)] = -\int_{\mathbb{R}^2} \rho(t,x)\left|\nabla \frac{\delta \mathcal{F}_{\rm PKS}(\rho(t))}{\delta\rho}(t,x)\right|^2\, \mathrm{d}x\,.$$

Critical Case: Stationary States & Main Result

The critical case $M = 8\pi$ has a family of explicit stationary solutions of the form

$$\bar{\rho}_\lambda(x) = \frac{8\lambda}{(\lambda + |x|^2)^2}$$

with $\lambda > 0$.

- All of these stationary solutions have critical mass and infinite second moment.
- They are the cases of equality functions for the Log HLS inequality.

Critical Case: Stationary States & Main Result

The critical case $M = 8\pi$ has a family of explicit stationary solutions of the form

$$\bar{\rho}_\lambda(x) = \frac{8\lambda}{(\lambda + |x|^2)^2}$$

with $\lambda > 0$.

- All of these stationary solutions have critical mass and infinite second moment.
- They are the cases of equality functions for the Log HLS inequality.

Critical Case: Stationary States & Main Result

The critical case $M = 8\pi$ has a family of explicit stationary solutions of the form

$$\bar{\rho}_\lambda(x) = \frac{8\lambda}{(\lambda + |x|^2)^2}$$

with $\lambda > 0$.

- All of these stationary solutions have critical mass and infinite second moment.
- They are the cases of equality functions for the Log HLS inequality.

Outline

Critical Fast Diffusion

The nonlinear Fokker-Planck equation in $\mathbb{R}^2$ with exponent $1/2$:

$$\begin{cases} \dfrac{\partial v}{\partial t}(t,x) = \Delta\left(\sqrt{v(t,x)}\right) + \dfrac{1}{\sqrt{2\lambda}}\,\mathrm{div}(x\,v(t,x)) & t > 0\,,\ x \in \mathbb{R}^2\,, \\ v(0,x) = v_0(x) \geq 0 & x \in \mathbb{R}^2\,, \end{cases}$$

corresponding to the *fast diffusion equation* $\frac{\partial u}{\partial t} = \Delta\left(\sqrt{u}\right)$ by a self-similar change of variable.

For $\lambda > 0$, define the functional $\mathcal{H}_\lambda$ on the non-negative functions in $L^1(\mathbb{R}^2)$ by

$$\mathcal{H}_\lambda[v] := \int_{\mathbb{R}^2} \left(\sqrt{v(x)} - \sqrt{\bar\rho_\lambda(x)}\right)^2 \bar\rho_\lambda^{-1/2}(x)\ \mathrm{d}x$$

This functional is the relative entropy of the fast diffusion equation with respect to the stationary solution $\bar\rho_\lambda$. The unique minimizer of $\mathcal{H}_\lambda$ is $\bar\rho_\lambda$.

Critical Fast Diffusion

The nonlinear Fokker-Planck equation in $\mathbb{R}^2$ with exponent $1/2$:

$$\begin{cases} \dfrac{\partial v}{\partial t}(t,x) = \Delta\left(\sqrt{v(t,x)}\right) + \dfrac{1}{\sqrt{2\lambda}}\,\mathrm{div}(x\,v(t,x)) & t>0\,,\ x\in\mathbb{R}^2\,, \\ v(0,x) = v_0(x) \geq 0 & x\in\mathbb{R}^2\,, \end{cases}$$

corresponding to the *fast diffusion equation* $\frac{\partial u}{\partial t} = \Delta\left(\sqrt{u}\right)$ by a self-similar change of variable.

For $\lambda > 0$, define the functional $\mathcal{H}_\lambda$ on the non-negative functions in $L^1(\mathbb{R}^2)$ by

$$\mathcal{H}_\lambda[v] := \int_{\mathbb{R}^2} \left(\sqrt{v(x)} - \sqrt{\bar\rho_\lambda(x)}\right)^2 \bar\rho_\lambda^{-1/2}(x)\;\mathrm{d}x$$

This functional is the relative entropy of the fast diffusion equation with respect to the stationary solution $\bar\rho_\lambda$. The unique minimizer of $\mathcal{H}_\lambda$ is $\bar\rho_\lambda$.

Formal Gradient Flow

A formal calculation as before shows,

$$\frac{\delta \mathcal{H}_\lambda[v]}{\delta v} = \frac{1}{\sqrt{\bar\rho_\lambda}} - \frac{1}{\sqrt{v}}\ ,$$

and the critical fast diffusion equation can be rewritten as

$$\frac{\partial v}{\partial t}(t,x) = \mathrm{div}\left(v(t,x) \left[\nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right] \right)\ ,$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{H}_\lambda[v(t)] = -\int_{\mathbb{R}^2} v(t,x) \left| \nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right|^2 \,\mathrm{d}x.$$

Let us point out that the functional $\mathcal{H}_\lambda[v]$ can be written as

$$\mathcal{H}_\lambda[u] := \int_{\mathbb{R}^2} \left[\Phi(v(x)) - \Phi(\bar\rho_\lambda(x)) - \Phi'(\bar\rho_\lambda)(v(x) - \bar\rho_\lambda(x)) \right] \,\mathrm{d}x$$

with $\Phi(s) = -2\sqrt{s}$ a convex function.

Formal Gradient Flow

A formal calculation as before shows,

$$\frac{\delta \mathcal{H}_\lambda[v]}{\delta v} = \frac{1}{\sqrt{\bar\rho_\lambda}} - \frac{1}{\sqrt{v}}\,,$$

and the critical fast diffusion equation can be rewritten as

$$\frac{\partial v}{\partial t}(t,x) = \operatorname{div}\left(v(t,x) \left[\nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right] \right)\,,$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{H}_\lambda[v(t)] = -\int_{\mathbb{R}^2} v(t,x) \left| \nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right|^2 \,\mathrm{d}x.$$

Let us point out that the functional $\mathcal{H}_\lambda[v]$ can be written as

$$\mathcal{H}_\lambda[u] := \int_{\mathbb{R}^2} \left[\Phi(v(x)) - \Phi(\bar\rho_\lambda(x)) - \Phi'(\bar\rho_\lambda)(v(x) - \bar\rho_\lambda(x)) \right] \,\mathrm{d}x$$

with $\Phi(s) = -2\sqrt{s}$ a convex function.

Formal Gradient Flow

A formal calculation as before shows,

$$\frac{\delta \mathcal{H}_\lambda[v]}{\delta v} = \frac{1}{\sqrt{\bar\rho_\lambda}} - \frac{1}{\sqrt{v}}\,,$$

and the critical fast diffusion equation can be rewritten as

$$\frac{\partial v}{\partial t}(t,x) = \mathrm{div}\left(v(t,x)\left[\nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v}\right]\right)\,,$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}_\lambda[v(t)] = -\int_{\mathbb{R}^2} v(t,x)\left|\nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v}\right|^2\,\mathrm{d}x.$$

Let us point out that the functional $\mathcal{H}_\lambda[v]$ can be written as

$$\mathcal{H}_\lambda[u] := \int_{\mathbb{R}^2}\left[\Phi(v(x)) - \Phi(\bar\rho_\lambda(x)) - \Phi'(\bar\rho_\lambda)(v(x)-\bar\rho_\lambda(x))\right]\,\mathrm{d}x$$

with $\Phi(s) = -2\sqrt{s}$ a convex function.

Formal Gradient Flow

A formal calculation as before shows,

$$\frac{\delta \mathcal{H}_\lambda[v]}{\delta v} = \frac{1}{\sqrt{\bar\rho_\lambda}} - \frac{1}{\sqrt{v}} \,,$$

and the critical fast diffusion equation can be rewritten as

$$\frac{\partial v}{\partial t}(t,x) = \mathrm{div}\left(v(t,x) \left[\nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right] \right) \,,$$

with entropy dissipation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{H}_\lambda[v(t)] = - \int_{\mathbb{R}^2} v(t,x) \left| \nabla \frac{\delta \mathcal{H}_\lambda[v]}{\delta v} \right|^2 \, \mathrm{d}x.$$

Let us point out that the functional $\mathcal{H}_\lambda[v]$ can be written as

$$\mathcal{H}_\lambda[u] := \int_{\mathbb{R}^2} \left[\Phi(v(x)) - \Phi(\bar\rho_\lambda(x)) - \Phi'(\bar\rho_\lambda)(v(x) - \bar\rho_\lambda(x)) \right] \, \mathrm{d}x$$

with $\Phi(s) = -2\sqrt{s}$ a convex function.

Outline

1. Macroscopic Models: PKS system
   - Modelling Chemotaxis: First Properties
   - PKS as Gradient Flow
   - Critical Fast Diffusion as Gradient Flow
   - New Liapunov Functionals
2. Ideas of the Rigorous Proof
   - Concentration-Control Inequalities
3. A byproduct: A new proof of HLS inequalities
   - Log HLS via Fast Diffusion Flows
   - A New Proof of the HLS inequality with Equality Cases
4. Conclusions

New Liapunov Functional

Claim: The critical fast diffusion functional is also a Liapunov functional for the critical mass PKS. Formal Computation:

$$\begin{aligned}\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}_\lambda[\rho(t)] &= \int_{\mathbb{R}^2} \frac{\delta\mathcal{H}_\lambda[\rho]}{\delta\rho}\operatorname{div}\left(\rho(t,x)\nabla\left[\frac{\delta\mathcal{F}_{\mathrm{PKS}}[\rho]}{\delta\rho}\right]\right)\,\mathrm{d}x\\
&= -\int_{\mathbb{R}^2}\rho\nabla\left[\frac{\delta\mathcal{H}_\lambda[\rho]}{\delta\rho}\right]\cdot\nabla\left[\frac{\delta\mathcal{F}_{\mathrm{PKS}}[\rho]}{\delta\rho}\right]\,\mathrm{d}x\\
&= -\int_{\mathbb{R}^2}\rho\nabla\left[\frac{1}{\sqrt{\rho_\lambda}}-\frac{1}{\sqrt{\rho}}\right]\cdot\nabla\left[\log\rho - G*\rho\right]\,\mathrm{d}x\\
&= -\int_{\mathbb{R}^2}\left[2\sqrt{\frac{\pi}{\lambda M}}\,x\,\rho+\nabla\sqrt{\rho}\right]\cdot\nabla\left[\log\rho - G*\rho\right]\,\mathrm{d}x\end{aligned}$$

Now, integrating by parts once more on the term involving the Green's function,

$$\begin{aligned}\int_{\mathbb{R}^2}\nabla\sqrt{\rho}\cdot\nabla\left[\log\rho - G*\rho\right]\,\mathrm{d}x &= \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}+\int_{\mathbb{R}^2}\sqrt{\rho}\,\Delta G*\rho\\
&= \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}-\int_{\mathbb{R}^2}\rho^{3/2}\,\mathrm{d}x\,.\end{aligned}$$

New Liapunov Functional

Claim: The critical fast diffusion functional is also a Liapunov functional for the critical mass PKS. Formal Computation:

$$\frac{d}{dt}\mathcal{H}_\lambda[\rho(t)] = \int_{\mathbb{R}^2} \frac{\delta \mathcal{H}_\lambda[\rho]}{\delta\rho} \operatorname{div}\left(\rho(t,x)\nabla\left[\frac{\delta \mathcal{F}_{\rm PKS}[\rho]}{\delta\rho}\right]\right)\, dx$$

$$= -\int_{\mathbb{R}^2} \rho\nabla\left[\frac{\delta \mathcal{H}_\lambda[\rho]}{\delta\rho}\right]\cdot\nabla\left[\frac{\delta \mathcal{F}_{\rm PKS}[\rho]}{\delta\rho}\right]\, dx$$

$$= -\int_{\mathbb{R}^2} \rho\nabla\left[\frac{1}{\sqrt{\bar\rho_\lambda}} - \frac{1}{\sqrt{\rho}}\right]\cdot\nabla\left[\log\rho - G*\rho\right]\, dx$$

$$= -\int_{\mathbb{R}^2}\left[2\sqrt{\frac{\pi}{\lambda M}}\,x\,\rho + \nabla\sqrt{\rho}\right]\cdot\nabla\left[\log\rho - G*\rho\right]\, dx$$

Now, integrating by parts once more on the term involving the Green's function,

$$\int_{\mathbb{R}^2}\nabla\sqrt{\rho}\cdot\nabla\left[\log\rho - G*\rho\right]\, dx = \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}} + \int_{\mathbb{R}^2}\sqrt{\rho}\,\Delta G*\rho$$

$$= \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}} - \int_{\mathbb{R}^2}\rho^{3/2}\, dx\,.$$

New Liapunov Functional

Claim: The critical fast diffusion functional is also a Liapunov functional for the critical mass PKS. Formal Computation:

$$\begin{aligned}\frac{d}{dt}\mathcal{H}_\lambda[\rho(t)] &= \int_{\mathbb{R}^2} \frac{\delta \mathcal{H}_\lambda[\rho]}{\delta\rho} \operatorname{div}\left(\rho(t,x)\nabla\left[\frac{\delta\mathcal{F}_{\mathrm{PKS}}[\rho]}{\delta\rho}\right]\right)\,dx \\ &= -\int_{\mathbb{R}^2} \rho\nabla\left[\frac{\delta \mathcal{H}_\lambda[\rho]}{\delta\rho}\right]\cdot\nabla\left[\frac{\delta\mathcal{F}_{\mathrm{PKS}}[\rho]}{\delta\rho}\right]\,dx \\ &= -\int_{\mathbb{R}^2} \rho\nabla\left[\frac{1}{\sqrt{\bar\rho_\lambda}}-\frac{1}{\sqrt{\rho}}\right]\cdot\nabla\left[\log\rho - G*\rho\right]\,dx \\ &= -\int_{\mathbb{R}^2}\left[2\sqrt{\frac{\pi}{\lambda M}}\,x\,\rho+\nabla\sqrt{\rho}\right]\cdot\nabla\left[\log\rho-G*\rho\right]\,dx\end{aligned}$$

Now, integrating by parts once more on the term involving the Green's function,

$$\begin{aligned}\int_{\mathbb{R}^2}\nabla\sqrt{\rho}\cdot\nabla\left[\log\rho-G*\rho\right]\,dx &= \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}+\int_{\mathbb{R}^2}\sqrt{\rho}\,\Delta G*\rho \\ &= \frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}-\int_{\mathbb{R}^2}\rho^{3/2}\,dx\,.\end{aligned}$$

New Liapunov Functional 2

Also, $\int_{\mathbb{R}^2} x\cdot\nabla\rho \,\mathrm{d}x = -2M$ and, making the same symmetrization that led to the evolution of the second moment,

$$\int_{\mathbb{R}^2} \rho(x)\, x\cdot\nabla G*\rho(x)\,\mathrm{d}x = \frac{1}{4\pi}\int_{\mathbb{R}^2\times\mathbb{R}^2} \rho(t,x)\,(x-y)\cdot\frac{x-y}{|x-y|^2}\,\rho(t,y)\,\mathrm{d}x\,\mathrm{d}y = -\frac{M^2}{4\pi}\,.$$

Collecting the above computations:

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}_\lambda[\rho(t)] = -\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}\,\mathrm{d}x + \int_{\mathbb{R}^2}\rho^{3/2}\,\mathrm{d}x + 4\sqrt{\frac{M\,\pi}{\lambda}}\left(1-\frac{M}{8\pi}\right)\,.$$

Notice that the constant term vanishes in critical mass case $M = 8\pi$. Thus, in the critical mass case, formal calculation yields that for all $T > 0$,

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T\left[\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}(t,x)\,\mathrm{d}x - \int_{\mathbb{R}^2}\rho^{3/2}(t,x)\,\mathrm{d}x\right]\mathrm{d}t \le \mathcal{H}_\lambda[\rho_0]\,.$$

New Liapunov Functional 2

Also, $\int_{\mathbb{R}^2} x \cdot \nabla\rho \, \mathrm{d}x = -2M$ and, making the same symmetrization that led to the evolution of the second moment,

$$\int_{\mathbb{R}^2} \rho(x)\, x \cdot \nabla G * \rho(x)\, \mathrm{d}x = \frac{1}{4\pi} \int_{\mathbb{R}^2\times\mathbb{R}^2} \rho(t,x)\,(x-y)\cdot\frac{x-y}{|x-y|^2}\,\rho(t,y)\,\mathrm{d}x\,\mathrm{d}y = -\frac{M^2}{4\pi}\,.$$

Collecting the above computations:

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}_\lambda[\rho(t)] = -\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}\,\mathrm{d}x + \int_{\mathbb{R}^2}\rho^{3/2}\,\mathrm{d}x + 4\sqrt{\frac{M\,\pi}{\lambda}}\left(1-\frac{M}{8\pi}\right).$$

Notice that the constant term vanishes in critical mass case $M = 8\pi$. Thus, in the critical mass case, formal calculation yields that for all $T > 0$,

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \left[\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}(t,x)\,\mathrm{d}x - \int_{\mathbb{R}^2}\rho^{3/2}(t,x)\,\mathrm{d}x\right]\mathrm{d}t \le \mathcal{H}_\lambda[\rho_0]\,.$$

New Liapunov Functional 2

Also, $\int_{\mathbb{R}^2} x \cdot \nabla \rho \, \mathrm{d}x = -2M$ and, making the same symmetrization that led to the evolution of the second moment,

$$\int_{\mathbb{R}^2} \rho(x)\, x \cdot \nabla G * \rho(x)\, \mathrm{d}x = \frac{1}{4\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} \rho(t,x)\,(x-y) \cdot \frac{x-y}{|x-y|^2}\, \rho(t,y)\, \mathrm{d}x\, \mathrm{d}y = -\frac{M^2}{4\pi} .$$

Collecting the above computations:

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{H}_\lambda[\rho(t)] = -\frac{1}{2} \int_{\mathbb{R}^2} \frac{|\nabla \rho|^2}{\rho^{3/2}}\, \mathrm{d}x + \int_{\mathbb{R}^2} \rho^{3/2}\, \mathrm{d}x + 4\sqrt{\frac{M\pi}{\lambda}} \left(1 - \frac{M}{8\pi}\right) .$$

Notice that the constant term vanishes in critical mass case $M = 8\pi$. Thus, in the critical mass case, formal calculation yields that for all $T > 0$,

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \left[\frac{1}{2} \int_{\mathbb{R}^2} \frac{|\nabla \rho|^2}{\rho^{3/2}}(t,x)\, \mathrm{d}x - \int_{\mathbb{R}^2} \rho^{3/2}(t,x)\, \mathrm{d}x \right] \mathrm{d}t \le \mathcal{H}_\lambda[\rho_0] .$$

New Liapunov Functional 2

Also, $\int_{\mathbb{R}^2} x\cdot\nabla\rho\,\mathrm{d}x = -2M$ and, making the same symmetrization that led to the evolution of the second moment,

$$\int_{\mathbb{R}^2} \rho(x)\, x\cdot\nabla G*\rho(x)\,\mathrm{d}x = \frac{1}{4\pi}\int_{\mathbb{R}^2\times\mathbb{R}^2} \rho(t,x)\,(x-y)\cdot\frac{x-y}{|x-y|^2}\,\rho(t,y)\,\mathrm{d}x\,\mathrm{d}y = -\frac{M^2}{4\pi}\,.$$

Collecting the above computations:

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}_\lambda[\rho(t)] = -\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}\,\mathrm{d}x + \int_{\mathbb{R}^2}\rho^{3/2}\,\mathrm{d}x + 4\sqrt{\frac{M\pi}{\lambda}}\left(1-\frac{M}{8\pi}\right)\,.$$

Notice that the constant term vanishes in critical mass case $M = 8\pi$. Thus, in the critical mass case, formal calculation yields that for all $T>0$,

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T\left[\frac{1}{2}\int_{\mathbb{R}^2}\frac{|\nabla\rho|^2}{\rho^{3/2}}(t,x)\,\mathrm{d}x - \int_{\mathbb{R}^2}\rho^{3/2}(t,x)\,\mathrm{d}x\right]\mathrm{d}t \le \mathcal{H}_\lambda[\rho_0]\,.$$

New Liapunov Functional 3

The missing link to exploit the previous relation is a particular case of the Gagliardo-Nirenberg-Sobolev inequalities[1].

Gagliardo-Nirenberg-Sobolev inequality

For all functions f in $\mathbb{R}^2$ with a square integrable distributional gradient ∇f,

$$\pi \int_{\mathbb{R}^2} |f|^6 \, dx \leq \int_{\mathbb{R}^2} |\nabla f|^2 \, dx \int_{\mathbb{R}^2} |f|^4 \, dx$$

and there is equality if and only if f is a multiple of a translate of $\bar\rho_\lambda^{1/4}$, $\lambda > 0$.

Dissipation of $\mathcal{H}_\lambda$

Applying the GNS to $f = \rho^{1/4}$: For all densities ρ of mass $M = 8\pi$,

$$\mathcal{D}[\rho] := \frac{1}{2} \int_{\mathbb{R}^2} \frac{|\nabla \rho(x)|^2}{\rho^{3/2}(x)} \, dx - \int_{\mathbb{R}^2} \rho^{3/2}(x) \, dx \geq 0 \,,$$

and moreover, there is equality if and only ρ is a translate of $\bar\rho_\lambda$ for some $\lambda > 0$.

[1] Dolbeault-DelPino, JMPA 2002

New Liapunov Functional 3

The missing link to exploit the previous relation is a particular case of the Gagliardo-Nirenberg-Sobolev inequalities[1].

Gagliardo-Nirenberg-Sobolev inequality

For all functions f in $\mathbb{R}^2$ with a square integrable distributional gradient ∇f,

$$\pi \int_{\mathbb{R}^2} |f|^6 \, dx \leq \int_{\mathbb{R}^2} |\nabla f|^2 \, dx \int_{\mathbb{R}^2} |f|^4 \, dx$$

and there is equality if and only if f is a multiple of a translate of $\bar{\rho}_\lambda^{1/4}$, $\lambda > 0$.

Dissipation of $\mathcal{H}_\lambda$

Applying the GNS to $f = \rho^{1/4}$: For all densities ρ of mass $M = 8\pi$,

$$\mathcal{D}[\rho] := \frac{1}{2} \int_{\mathbb{R}^2} \frac{|\nabla \rho(x)|^2}{\rho^{3/2}(x)} \, dx - \int_{\mathbb{R}^2} \rho^{3/2}(x) \, dx \geq 0 \, ,$$

and moreover, there is equality if and only ρ is a translate of $\bar{\rho}_\lambda$ for some $\lambda > 0$.

[1] Dolbeault-DelPino, JMPA 2002

New Liapunov Functional 3

The missing link to exploit the previous relation is a particular case of the Gagliardo-Nirenberg-Sobolev inequalities[1].

Gagliardo-Nirenberg-Sobolev inequality

For all functions f in $\mathbb{R}^2$ with a square integrable distributional gradient ∇f,

$$\pi \int_{\mathbb{R}^2} |f|^6 \, dx \leq \int_{\mathbb{R}^2} |\nabla f|^2 \, dx \int_{\mathbb{R}^2} |f|^4 \, dx$$

and there is equality if and only if f is a multiple of a translate of $\bar{\rho}_\lambda^{1/4}$, $\lambda > 0$.

Dissipation of $\mathcal{H}_\lambda$

Applying the GNS to $f = \rho^{1/4}$: For all densities ρ of mass $M = 8\pi$,

$$\mathcal{D}[\rho] := \frac{1}{2} \int_{\mathbb{R}^2} \frac{|\nabla \rho(x)|^2}{\rho^{3/2}(x)} \, dx - \int_{\mathbb{R}^2} \rho^{3/2}(x) \, dx \geq 0 \, ,$$

and moreover, there is equality if and only ρ is a translate of $\bar{\rho}_\lambda$ for some $\lambda > 0$.

[1] Dolbeault-DelPino, JMPA 2002

Main Result Blanchet-Carlen-C.

Bassin of attractions in the Critical Mass PKS

Given any density ρ_0 in $\mathbb{R}^2$ with total mass 8π such that there exists $\lambda > 0$ with

$$\mathcal{H}_\lambda[\rho_0] < \infty$$

Then there exists $\rho \in \mathcal{AC}^0([0,T],\mathcal{P}_2(\mathbb{R}^2))$, with $\rho(t) \in L^1(\mathbb{R}^2)$ for all $t \geq 0$ being a global-in-time weak solution of the critical mass PKS. Moreover, the solutions constructed satisfy that

$$\mathcal{H}_\lambda[\rho(t)] \leq \mathcal{H}_\lambda[\rho(s)] \qquad \text{and} \qquad \mathcal{F}_{\mathrm{PKS}}[\rho(t)] \leq \mathcal{F}_{\mathrm{PKS}}[\rho(s)]\,,$$

for all $0 \leq s \leq t$. Moreover, we can show that the constructed weak solutions are *dissipative solutions, i.e.*, they satisfy for all $T > 0$

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \mathcal{D}[\rho(t)]\,\mathrm{d}t \leq \mathcal{H}_\lambda[\rho_0]$$

and $\lim_{t\to\infty} \|\rho(t) - \bar\rho_\lambda\|_{L^1(\mathbb{R}^2)} = 0.$

Main Result Blanchet-Carlen-C.

Bassin of attractions in the Critical Mass PKS

Given any density ρ_0 in $\mathbb{R}^2$ with total mass 8π such that there exists $\lambda > 0$ with

$$\mathcal{H}_\lambda[\rho_0] < \infty$$

Then there exists $\rho \in \mathcal{AC}^0([0,T],\mathcal{P}_2(\mathbb{R}^2))$, with $\rho(t) \in L^1(\mathbb{R}^2)$ for all $t \geq 0$ being a global-in-time weak solution of the critical mass PKS. Moreover, the solutions constructed satisfy that

$$\mathcal{H}_\lambda[\rho(t)] \leq \mathcal{H}_\lambda[\rho(s)] \qquad \text{and} \qquad \mathcal{F}_{\text{PKS}}[\rho(t)] \leq \mathcal{F}_{\text{PKS}}[\rho(s)]\,,$$

for all $0 \leq s \leq t$. Moreover, we can show that the constructed weak solutions are *dissipative solutions, i.e.*, they satisfy for all $T > 0$

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \mathcal{D}[\rho(t)]\,\mathrm{d}t \leq \mathcal{H}_\lambda[\rho_0]$$

and $\lim_{t\to\infty} \|\rho(t) - \bar{\rho}_\lambda\|_{L^1(\mathbb{R}^2)} = 0.$

Main Result Blanchet-Carlen-C.

Bassin of attractions in the Critical Mass PKS

Given any density ρ_0 in $\mathbb{R}^2$ with total mass 8π such that there exists $\lambda > 0$ with

$$\mathcal{H}_\lambda[\rho_0] < \infty$$

Then there exists $\rho \in \mathcal{AC}^0([0,T],\mathcal{P}_2(\mathbb{R}^2))$, with $\rho(t) \in L^1(\mathbb{R}^2)$ for all $t \geq 0$ being a global-in-time weak solution of the critical mass PKS. Moreover, the solutions constructed satisfy that

$$\mathcal{H}_\lambda[\rho(t)] \leq \mathcal{H}_\lambda[\rho(s)] \qquad \text{and} \qquad \mathcal{F}_{\mathrm{PKS}}[\rho(t)] \leq \mathcal{F}_{\mathrm{PKS}}[\rho(s)]\,,$$

for all $0 \leq s \leq t$. Moreover, we can show that the constructed weak solutions are *dissipative solutions, i.e.*, they satisfy for all $T > 0$

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \mathcal{D}[\rho(t)]\,\mathrm{d}t \leq \mathcal{H}_\lambda[\rho_0]$$

and $\lim_{t\to\infty} \|\rho(t) - \bar\rho_\lambda\|_{L^1(\mathbb{R}^2)} = 0.$

Main Result Blanchet-Carlen-C.

Bassin of attractions in the Critical Mass PKS

Given any density ρ_0 in $\mathbb{R}^2$ with total mass 8π such that there exists $\lambda > 0$ with

$$\mathcal{H}_\lambda[\rho_0] < \infty$$

Then there exists $\rho \in \mathcal{AC}^0([0,T],\mathcal{P}_2(\mathbb{R}^2))$, with $\rho(t) \in L^1(\mathbb{R}^2)$ for all $t \geq 0$ being a global-in-time weak solution of the critical mass PKS. Moreover, the solutions constructed satisfy that

$$\mathcal{H}_\lambda[\rho(t)] \leq \mathcal{H}_\lambda[\rho(s)] \qquad \text{and} \qquad \mathcal{F}_{\mathrm{PKS}}[\rho(t)] \leq \mathcal{F}_{\mathrm{PKS}}[\rho(s)]\,,$$

for all $0 \leq s \leq t$. Moreover, we can show that the constructed weak solutions are *dissipative solutions, i.e.*, they satisfy for all $T > 0$

$$\mathcal{H}_\lambda[\rho(T)] + \int_0^T \mathcal{D}[\rho(t)]\,\mathrm{d}t \leq \mathcal{H}_\lambda[\rho_0]$$

and $\lim_{t\to\infty} \|\rho(t) - \bar\rho_\lambda\|_{L^1(\mathbb{R}^2)} = 0.$

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\text{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\text{PKS}}$ takes its minimum value for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar{\rho}_\mu] \to \infty$ as $\mu \to 0$ since $\bar{\rho}_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\|\nabla(\bar{\rho}_\mu^{1/4})\|_2$ or $\|\bar{\rho}_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\text{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\text{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\text{PKS}}[\rho] = \mathcal{F}_{\text{PKS}}[\bar{\rho}_\lambda]$, then, up to translation, $\rho = \bar{\rho}_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\text{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\mathrm{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\mathrm{PKS}}$ takes its minimum value for $\rho = \bar\rho_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar\rho_\mu] \to \infty$ as $\mu \to 0$ since $\bar\rho_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar\rho_\mu$ for all $\mu > 0$, but $\|\nabla(\bar\rho_\mu^{1/4})\|_2$ or $\|\bar\rho_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\mathrm{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\mathrm{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\mathrm{PKS}}[\rho] = \mathcal{F}_{\mathrm{PKS}}[\bar\rho_\lambda]$, then, up to translation, $\rho = \bar\rho_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\mathrm{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\text{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\text{PKS}}$ takes its minimum value for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar{\rho}_\mu] \to \infty$ as $\mu \to 0$ since $\bar{\rho}_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\|\nabla(\bar{\rho}_\mu^{1/4})\|_2$ or $\|\bar{\rho}_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\text{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\text{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\text{PKS}}[\rho] = \mathcal{F}_{\text{PKS}}[\bar{\rho}_\lambda]$, then, up to translation, $\rho = \bar{\rho}_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\text{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\text{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\text{PKS}}$ takes its minimum value for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar{\rho}_\mu] \to \infty$ as $\mu \to 0$ since $\bar{\rho}_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\|\nabla(\bar{\rho}_\mu^{1/4})\|_2$ or $\|\bar{\rho}_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\text{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\text{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\text{PKS}}[\rho] = \mathcal{F}_{\text{PKS}}[\bar{\rho}_\lambda]$, then, up to translation, $\rho = \bar{\rho}_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\text{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\text{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\text{PKS}}$ takes its minimum value for $\rho = \bar\rho_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar\rho_\mu] \to \infty$ as $\mu \to 0$ since $\bar\rho_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar\rho_\mu$ for all $\mu > 0$, but $\|\nabla(\bar\rho_\mu^{1/4})\|_2$ or $\|\bar\rho_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\text{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\text{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\text{PKS}}[\rho] = \mathcal{F}_{\text{PKS}}[\bar\rho_\lambda]$, then, up to translation, $\rho = \bar\rho_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\text{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\text{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\text{PKS}}$ takes its minimum value for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar{\rho}_\mu] \to \infty$ as $\mu \to 0$ since $\bar{\rho}_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\|\nabla(\bar{\rho}_\mu^{1/4})\|_2$ or $\|\bar{\rho}_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\text{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\text{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\text{PKS}}[\rho] = \mathcal{F}_{\text{PKS}}[\bar{\rho}_\lambda]$, then, up to translation, $\rho = \bar{\rho}_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\text{PKS}}$, it mainly helps to quantify the decay rate.

Mathematical Difficulties

- For a density of mass 8π, an upper bound on $\mathcal{F}_{\mathrm{PKS}}[\rho]$ does not provide any upper bound on the entropy $\mathcal{E}[\rho]$.

  Indeed, $\mathcal{F}_{\mathrm{PKS}}$ takes its minimum value for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\mathcal{E}[\bar{\rho}_\mu] \to \infty$ as $\mu \to 0$ since $\bar{\rho}_\mu$ converges weakly-* as measures towards a Dirac delta at the origin with mass 8π.

- The dissipation functional $\mathcal{D}[\rho]$ is well defined as long as $\|\rho\|_{3/2} < \infty$, but an upper bound on $\mathcal{D}[\rho]$ does not give an upper bound on either $\|\nabla(\rho^{1/4})\|_2$ or $\|\rho\|_{3/2}$ since it is the difference of both terms.

  Indeed, $\mathcal{D}[\rho] = 0$ for $\rho = \bar{\rho}_\mu$ for all $\mu > 0$, but $\|\nabla(\bar{\rho}_\mu^{1/4})\|_2$ or $\|\bar{\rho}_\mu\|_{3/2}$ both diverge to infinity as μ tends to zero.

- Stability properties of $\mathcal{H}_\lambda$ over $\mathcal{F}_{\mathrm{PKS}}$. In fact, not much was known about stability for $\mathcal{F}_{\mathrm{PKS}}$ till very recently (E. Carlen, A. Figalli, preprint).

  We do know that if some density ρ with mass 8π satisfies $\mathcal{F}_{\mathrm{PKS}}[\rho] = \mathcal{F}_{\mathrm{PKS}}[\bar{\rho}_\lambda]$, then, up to translation, $\rho = \bar{\rho}_\mu$ for some $\mu > 0$. Even knowing a quantitative estimate for the error in $\mathcal{F}_{\mathrm{PKS}}$, it mainly helps to quantify the decay rate.

Outline

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\,\mathrm{d}x \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\,\mathrm{d}x = \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\,\mathrm{d}x \leq \frac{M}{2}.$$

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\, dx \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\, dx = \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\, dx \leq \frac{M}{2}\,.$$

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\, dx \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\, dx = \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\, dx \leq \frac{M}{2}.$$

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\, \mathrm{d}x \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\, \mathrm{d}x \;=\; \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\, \mathrm{d}x \leq \frac{M}{2}\,.$$

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\, \mathrm{d}x \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\, \mathrm{d}x \;=\; \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\, \mathrm{d}x \leq \frac{M}{2}\,.$$

Thick Tails & Localization by $\mathcal{H}_\lambda$

- Rigorous Proof of Displacement Convexity of $\mathcal{H}_\lambda$ by regularization.
- Proof of Talagrand inequality to control $W_2(\rho, \varrho_\lambda)$ in terms of $\mathcal{H}_\lambda[\rho]$.

Thick Tails

A bound on $\mathcal{H}_\lambda[\rho]$ implies a "matching tail" estimate.- Let ρ be any density of mass M such that $\mathcal{H}_\lambda[\rho] < \infty$. Then for $\eta_* := \frac{1}{5}e^{-1/5}$ and any $s > 1$

$$\int_{|x|^2 \geq \lambda s^2} \rho(x)\,\mathrm{d}x \geq \eta_*\, e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]} \int_{|x|^2 \geq \lambda s^2} \varrho_\lambda(x)\,\mathrm{d}x = \frac{M\eta_*}{1+s^2} e^{-\frac{4}{\sqrt{\pi M \lambda}}\mathcal{H}_\lambda[\rho]}.$$

Solid core

A bound on $\mathcal{H}_\lambda[\rho]$ implies localization.- Let ρ be a density of mass M such that for some $\lambda > 0$, $\mathcal{H}_\lambda[\rho] < \infty$. Then

$$\int_{\{|x| \geq 4\sqrt{\lambda} + 4(\lambda/M\pi)^{1/4}\sqrt{\mathcal{H}_\lambda[\rho]}\}} \rho\,\mathrm{d}x \leq \frac{M}{2}\,.$$

Concentration-Control for the Entropy I

In summary, we can show there exists $R > 1$ so that for some $0 < a < 8\pi$,

$$\int_{|x|>R-1} \rho \,\mathrm{d}x \le 8\pi - a \qquad \text{and} \qquad \int_{|x|<R+1} \rho \,\mathrm{d}x \le 8\pi - a \,.$$

The densities lying in sublevel sets of $\mathcal{F}_{\mathrm{PKS}}[\rho]$ and $\mathcal{H}_\lambda[\rho]$ are compact.

Entropy bound

Let ρ be any density with mass $M = 8\pi$, with $\mathcal{H}_\lambda[\rho] < \infty$ for some $\lambda > 0$. Then there exist positive computable constants γ_1 and C_{CCF} depending only on λ and $\mathcal{H}_\lambda[\rho]$ such that

$$\gamma_1 \int_{\mathbb{R}^2} \rho \log_+ \rho \,\mathrm{d}x \le \mathcal{F}_{\mathrm{PKS}}[\rho] + C_{\mathrm{CCF}} \,.$$

Concentration-Control for the Entropy I

In summary, we can show there exists $R > 1$ so that for some $0 < a < 8\pi$,

$$\int_{|x|>R-1} \rho \, \mathrm{d}x \le 8\pi - a \qquad \text{and} \qquad \int_{|x|<R+1} \rho \, \mathrm{d}x \le 8\pi - a \, .$$

The densities lying in sublevel sets of $\mathcal{F}_{\mathrm{PKS}}[\rho]$ and $\mathcal{H}_\lambda[\rho]$ are compact.

Entropy bound

Let ρ be any density with mass $M = 8\pi$, with $\mathcal{H}_\lambda[\rho] < \infty$ for some $\lambda > 0$. Then there exist positive computable constants γ_1 and C_{CCF} depending only on λ and $\mathcal{H}_\lambda[\rho]$ such that

$$\gamma_1 \int_{\mathbb{R}^2} \rho \log_+ \rho \, \mathrm{d}x \le \mathcal{F}_{\mathrm{PKS}}[\rho] + C_{\mathrm{CCF}} \, .$$

Concentration-Control for the Entropy I

In summary, we can show there exists $R > 1$ so that for some $0 < a < 8\pi$,

$$\int_{|x|>R-1} \rho \, \mathrm{d}x \le 8\pi - a \qquad \text{and} \qquad \int_{|x|<R+1} \rho \, \mathrm{d}x \le 8\pi - a \, .$$

The densities lying in sublevel sets of $\mathcal{F}_{\mathrm{PKS}}[\rho]$ and $\mathcal{H}_\lambda[\rho]$ are compact.

Entropy bound

Let ρ be any density with mass $M = 8\pi$, with $\mathcal{H}_\lambda[\rho] < \infty$ for some $\lambda > 0$. Then there exist positive computable constants γ_1 and C_{CCF} depending only on λ and $\mathcal{H}_\lambda[\rho]$ such that

$$\gamma_1 \int_{\mathbb{R}^2} \rho \log_+ \rho \, \mathrm{d}x \le \mathcal{F}_{\mathrm{PKS}}[\rho] + C_{\mathrm{CCF}} \, .$$

Concentration-Control for Entropy-Dissipation

The densities lying in sublevel sets of $\mathcal{F}_{\text{PKS}}[\rho]$, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{D}[\rho]$ are compact.

Concentration control for $\mathcal{D}$

Let ρ be any density with mass 8π, $\mathcal{F}_{\text{PKS}}[\rho]$ finite, and $\mathcal{H}_\lambda[\rho]$ finite for some $\lambda > 0$. Then there exist positive computable constants γ_2 and C_{CCD} depending only on λ, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{F}_{\text{PKS}}[\rho]$ such that

$$\gamma_2 \int_{\mathbb{R}^2} |\nabla \rho^{1/4}|^2 \, dx \leq \pi \mathcal{D}[\rho] + C_{\text{CCD}} \ .$$

This allows to control the $L^{3/2}$-norm of ρ too.

Concentration-Control for Entropy-Dissipation

The densities lying in sublevel sets of $\mathcal{F}_{\mathrm{PKS}}[\rho]$, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{D}[\rho]$ are compact.

Concentration control for $\mathcal{D}$

Let ρ be any density with mass 8π, $\mathcal{F}_{\mathrm{PKS}}[\rho]$ finite, and $\mathcal{H}_\lambda[\rho]$ finite for some $\lambda > 0$. Then there exist positive computable constants γ_2 and C_{CCD} depending only on λ, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{F}_{\mathrm{PKS}}[\rho]$ such that

$$\gamma_2 \int_{\mathbb{R}^2} |\nabla \rho^{1/4}|^2 \, dx \leq \pi \mathcal{D}[\rho] + C_{\mathrm{CCD}} \ .$$

This allows to control the $L^{3/2}$-norm of ρ too.

Concentration-Control for Entropy-Dissipation

The densities lying in sublevel sets of $\mathcal{F}_{\text{PKS}}[\rho]$, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{D}[\rho]$ are compact.

Concentration control for $\mathcal{D}$

Let ρ be any density with mass 8π, $\mathcal{F}_{\text{PKS}}[\rho]$ finite, and $\mathcal{H}_\lambda[\rho]$ finite for some $\lambda > 0$. Then there exist positive computable constants γ_2 and C_{CCD} depending only on λ, $\mathcal{H}_\lambda[\rho]$ and $\mathcal{F}_{\text{PKS}}[\rho]$ such that

$$\gamma_2 \int_{\mathbb{R}^2} |\nabla \rho^{1/4}|^2 \, \mathrm{d}x \le \pi \mathcal{D}[\rho] + C_{\text{CCD}} \ .$$

This allows to control the $L^{3/2}$-norm of ρ too.

“Duality” between CFD and PKS

Observation of Matthes, McCann & Savaré, CPDE 2009

Consider two smooth functions Φ and Ψ on $\mathbb{R}^d$, and consider the two ordinary differential equations describing gradient flow:

$$\dot{x}(t) = -\nabla\Phi[x(t)] \qquad \text{and} \qquad \dot{y}(t) = -\nabla\Psi[y(t)]\ .$$

Then of course $\Phi[x(t)]$ and $\Psi[y(t)]$ are monotone decreasing. Now differentiate each function along the others flow:

$$\frac{\mathrm{d}}{\mathrm{d}t}\Phi[y(t)] = -\langle\nabla\Phi[y(t)], \nabla\Psi[y(t)]\rangle$$
$$\frac{\mathrm{d}}{\mathrm{d}t}\Psi[x(t)] = -\langle\nabla\Psi[x(t)], \nabla\Phi[x(t)]\rangle\ .$$

Thus, Φ is decreasing along the gradient flow of Ψ for any initial data if and only if Ψ is decreasing along the gradient flow of Φ for any initial data.

"Duality" between CFD and PKS

Observation of Matthes, McCann & Savaré, CPDE 2009

Consider two smooth functions Φ and Ψ on $\mathbb{R}^d$, and consider the two ordinary differential equations describing gradient flow:

$$\dot{x}(t) = -\nabla\Phi[x(t)] \qquad \text{and} \qquad \dot{y}(t) = -\nabla\Psi[y(t)]\ .$$

Then of course $\Phi[x(t)]$ and $\Psi[y(t)]$ are monotone decreasing. Now differentiate each function along the others flow:

$$\frac{\mathrm{d}}{\mathrm{d}t}\Phi[y(t)] = -\langle\nabla\Phi[y(t)], \nabla\Psi[y(t)]\rangle$$
$$\frac{\mathrm{d}}{\mathrm{d}t}\Psi[x(t)] = -\langle\nabla\Psi[x(t)], \nabla\Phi[x(t)]\rangle\ .$$

Thus, Φ is decreasing along the gradient flow of Ψ for any initial data if and only if Ψ is decreasing along the gradient flow of Φ for any initial data.

"Duality" between CFD and PKS

- An analog of this holds for well-behaved gradient flows in the 2-Wasserstein sense as used in (Matthes, McCann & Savare, CPDE 2009).
- Our case: Apply it to the Log-HLS functional in $d = 2$.-

Since $\mathcal{H}_\lambda$ is decreasing along the 2-Wasserstein gradient flow for $\mathcal{F}_{\text{PKS}}$, i.e., the Patlak-Keller-Segel equation, one can expect that $\mathcal{F}_{\text{PKS}}$ of the Log-HLS functional is decreasing along the 2-Wasserstein gradient flow for $\mathcal{H}_\lambda$, i.e., the critical fast diffusion.

"Duality" between CFD and PKS

- An analog of this holds for well-behaved gradient flows in the 2-Wasserstein sense as used in (Matthes, McCann & Savare, CPDE 2009).
- Our case: Apply it to the Log-HLS functional in $d = 2$.-

Since $\mathcal{H}_\lambda$ is decreasing along the 2-Wasserstein gradient flow for $\mathcal{F}_{\mathrm{PKS}}$, i.e., the Patlak-Keller-Segel equation, one can expect that $\mathcal{F}_{\mathrm{PKS}}$ of the Log-HLS functional is decreasing along the 2-Wasserstein gradient flow for $\mathcal{H}_\lambda$, i.e., the critical fast diffusion.

Outline

The connection from FD perspective

The log HLS functional $\mathcal{F}$ is defined by

$$\mathcal{F}[f] := \int_{\mathbb{R}^2} f(x) \log f(x)\, \mathrm{d}x + \frac{2}{\int_{\mathbb{R}^2} f(x)\, \mathrm{d}x} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) \log|x-y| f(y)\, \mathrm{d}x\, \mathrm{d}y\,.$$

The critical fast diffusion is

$$\frac{\partial}{\partial t} u(x,t) = \Delta u^{1/2}(x,t)\,,$$

with associated Fokker-Planck equation

$$\frac{\partial}{\partial t} v(x,t) = \Delta v^{1/2}(x,t) + \nabla \cdot [x v(x,t)]\,.$$

which has as stationary solution

$$h(x) = \frac{M}{\pi}\left(\frac{1}{1+|x|^2}\right)^2$$

with mass M.

Scaling: $f_{(a)} := a^2 f(ax)$. Then, $\mathcal{F}[f_{(a)}] = \mathcal{F}[f]$ for all a. $v(x,t) := e^{2t} u(e^t x, e^t)$ and thus $\mathcal{F}[v(\cdot,t)] = \mathcal{F}[u(\cdot, e^t)]$.

The connection from FD perspective

The log HLS functional $\mathcal{F}$ is defined by

$$\mathcal{F}[f] := \int_{\mathbb{R}^2} f(x) \log f(x) \, \mathrm{d}x + \frac{2}{\int_{\mathbb{R}^2} f(x) \, \mathrm{d}x} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) \log|x-y| f(y) \, \mathrm{d}x \, \mathrm{d}y \, .$$

The critical fast diffusion is

$$\frac{\partial}{\partial t} u(x,t) = \Delta u^{1/2}(x,t) \, ,$$

with associated Fokker-Planck equation

$$\frac{\partial}{\partial t} v(x,t) = \Delta v^{1/2}(x,t) + \nabla \cdot [x v(x,t)] \, .$$

which has as stationary solution

$$h(x) = \frac{M}{\pi} \left(\frac{1}{1+|x|^2} \right)^2$$

with mass M.

Scaling: $f_{(a)} := a^2 f(ax)$. Then, $\mathcal{F}[f_{(a)}] = \mathcal{F}[f]$ for all a. $v(x,t) := e^{2t} u(e^t x, e^t)$ and thus $\mathcal{F}[v(\cdot,t)] = \mathcal{F}[u(\cdot,e^t)]$.

The connection from FD perspective

The log HLS functional $\mathcal{F}$ is defined by

$$\mathcal{F}[f] := \int_{\mathbb{R}^2} f(x) \log f(x) \, dx + \frac{2}{\int_{\mathbb{R}^2} f(x) \, dx} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) \log |x - y| f(y) \, dx \, dy \,.$$

The critical fast diffusion is

$$\frac{\partial}{\partial t} u(x,t) = \Delta u^{1/2}(x,t) \,,$$

with associated Fokker-Planck equation

$$\frac{\partial}{\partial t} v(x,t) = \Delta v^{1/2}(x,t) + \nabla \cdot [x v(x,t)] \,.$$

which has as stationary solution

$$h(x) = \frac{M}{\pi} \left(\frac{1}{1+|x|^2} \right)^2$$

with mass M.

Scaling: $f_{(a)} := a^2 f(ax)$. Then, $\mathcal{F}[f_{(a)}] = \mathcal{F}[f]$ for all a. $v(x,t) := e^{2t} u(e^t x, e^t)$ and thus $\mathcal{F}[v(\cdot,t)] = \mathcal{F}[u(\cdot, e^t)]$.

The connection from FD perspective

The log HLS functional $\mathcal{F}$ is defined by

$$\mathcal{F}[f] := \int_{\mathbb{R}^2} f(x) \log f(x) \, \mathrm{d}x + \frac{2}{\int_{\mathbb{R}^2} f(x) \, \mathrm{d}x} \iint_{\mathbb{R}^2\times\mathbb{R}^2} f(x) \log|x-y| f(y) \, \mathrm{d}x \, \mathrm{d}y \, .$$

The critical fast diffusion is

$$\frac{\partial}{\partial t} u(x,t) = \Delta u^{1/2}(x,t) \, ,$$

with associated Fokker-Planck equation

$$\frac{\partial}{\partial t} v(x,t) = \Delta v^{1/2}(x,t) + \nabla \cdot [x v(x,t)] \, .$$

which has as stationary solution

$$h(x) = \frac{M}{\pi} \left(\frac{1}{1+|x|^2} \right)^2$$

with mass M.

Scaling: $f_{(a)} := a^2 f(ax)$. Then, $\mathcal{F}[f_{(a)}] = \mathcal{F}[f]$ for all a. $v(x,t) := e^{2t} u(e^t x, e^t)$ and thus $\mathcal{F}[v(\cdot,t)] = \mathcal{F}[u(\cdot,e^t)]$.

The connection from FD perspective

The log HLS functional $\mathcal{F}$ is defined by

$$\mathcal{F}[f] := \int_{\mathbb{R}^2} f(x) \log f(x)\, \mathrm{d}x + \frac{2}{\int_{\mathbb{R}^2} f(x)\, \mathrm{d}x} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) \log|x-y| f(y)\, \mathrm{d}x\, \mathrm{d}y\,.$$

The critical fast diffusion is

$$\frac{\partial}{\partial t} u(x,t) = \Delta u^{1/2}(x,t)\,,$$

with associated Fokker-Planck equation

$$\frac{\partial}{\partial t} v(x,t) = \Delta v^{1/2}(x,t) + \nabla \cdot [x v(x,t)]\,.$$

which has as stationary solution

$$h(x) = \frac{M}{\pi}\left(\frac{1}{1+|x|^2}\right)^2$$

with mass M.

Scaling: $f_{(a)} := a^2 f(ax)$. Then, $\mathcal{F}[f_{(a)}] = \mathcal{F}[f]$ for all a. $v(x,t) := e^{2t} u(e^t x, e^t)$ and thus $\mathcal{F}[v(\cdot,t)] = \mathcal{F}[u(\cdot,e^t)]$.

The connection from FD perspective

Log HLS via fast diffusion flow

Let f be a non-negative measurable functions on $\mathbb{R}^2$ such that $f\log f$ and $f\log(e+|x|^2)$ belong to $L^1(\mathbb{R}^2)$. Suppose also that

$$\int_{\mathbb{R}^2} f(x)\ \mathrm{d}x = \int_{\mathbb{R}^2} h(x)\ \mathrm{d}x = M$$

and

$$\sup_{|x|>R} f(x)|x|^4 < \infty\ .$$

Let $u(x,t)$ be the solution of critical fast diffusion with $u(x,1)=f(x)$. Then

$$\mathcal{F}[f] = \mathcal{F}[h] + \int_1^\infty \mathcal{D}[u^{1/4}(\cdot,e^t)]\ \mathrm{d}t \ge \mathcal{F}[h]\ ,$$

where

$$\mathcal{D}[g] = \int_{\mathbb{R}^2} |\nabla g(x)|^2\ \mathrm{d}x \int_{\mathbb{R}^2} g^4(x)\ \mathrm{d}x - \pi \int_{\mathbb{R}^2} g^6(x)\ \mathrm{d}x\ ,$$

which is non-negative by the sharp GNS inequality.

The connection from FD perspective

Log HLS via fast diffusion flow

Let f be a non-negative measurable functions on $\mathbb{R}^2$ such that $f\log f$ and $f\log(e+|x|^2)$ belong to $L^1(\mathbb{R}^2)$. Suppose also that

$$\int_{\mathbb{R}^2} f(x)\ \mathrm{d}x = \int_{\mathbb{R}^2} h(x)\ \mathrm{d}x = M$$

and

$$\sup_{|x|>R} f(x)|x|^4 < \infty\ .$$

Let $u(x,t)$ be the solution of critical fast diffusion with $u(x,1)=f(x)$. Then

$$\mathcal{F}[f] = \mathcal{F}[h] + \int_1^\infty \mathcal{D}[u^{1/4}(\cdot, e^t)]\ \mathrm{d}t \geq \mathcal{F}[h]\ ,$$

where

$$\mathcal{D}[g] = \int_{\mathbb{R}^2} |\nabla g(x)|^2\ \mathrm{d}x \int_{\mathbb{R}^2} g^4(x)\ \mathrm{d}x - \pi \int_{\mathbb{R}^2} g^6(x)\ \mathrm{d}x\ ,$$

which is non-negative by the sharp GNS inequality.

Outline

Two Inequalities, One Equation

Sharp Hardy-Littlewood-Sobolev inequality

[a] It states that for all non-negative measurable functions f on $\mathbb{R}^d$, and all $0 < \lambda < d$,

$$\frac{\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} f(x)\frac{1}{|x-y|^\lambda} f(y)\, \mathrm{d}x\, \mathrm{d}y}{\|f\|_p^2} \le \frac{\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} h(x)\frac{1}{|x-y|^\lambda} h(y)\, \mathrm{d}x\, \mathrm{d}y}{\|h\|_p^2}$$

where

$$h(x) = \left(\frac{1}{1+|x|^2}\right)^{(2d-\lambda)/2} .$$

and $p = \frac{2d}{2d-\lambda}$.

Moreover, there is equality if and only if for some $x_0 \in \mathbb{R}^d$ and $s \in \mathbb{R}_+$, f is a non-zero multiple of $h(x/s - x_0)$.

[a] E. Lieb, Ann. Math 1983.

For $d \ge 3$, the $\lambda = d-2$ case of the sharp HLS inequality can be proved by using the fast diffusion flow for $m = d/(d+2)$.

Two Inequalities, One Equation

Sharp Hardy-Littlewood-Sobolev inequality

[a] It states that for all non-negative measurable functions f on $\mathbb{R}^d$, and all $0 < \lambda < d$,

$$\frac{\displaystyle\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} f(x)\frac{1}{|x-y|^\lambda}f(y)\,\mathrm{d}x\,\mathrm{d}y}{\|f\|_p^2} \le \frac{\displaystyle\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} h(x)\frac{1}{|x-y|^\lambda}h(y)\,\mathrm{d}x\,\mathrm{d}y}{\|h\|_p^2}$$

where

$$h(x) = \left(\frac{1}{1+|x|^2}\right)^{(2d-\lambda)/2} .$$

and $p = \frac{2d}{2d-\lambda}$.

Moreover, there is equality if and only if for some $x_0 \in \mathbb{R}^d$ and $s \in \mathbb{R}_+$, f is a non-zero multiple of $h(x/s - x_0)$.

[a] E. Lieb, Ann. Math 1983.

For $d \ge 3$, the $\lambda = d - 2$ case of the sharp HLS inequality can be proved by using the fast diffusion flow for $m = d/(d+2)$.

Two Inequalities, One Equation

Sharp Hardy-Littlewood-Sobolev inequality

[a] It states that for all non-negative measurable functions f on $\mathbb{R}^d$, and all $0 < \lambda < d$,

$$\frac{\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} f(x)\frac{1}{|x-y|^\lambda} f(y)\, \mathrm{d}x\, \mathrm{d}y}{\|f\|_p^2} \leq \frac{\int_{\mathbb{R}^d}\int_{\mathbb{R}^d} h(x)\frac{1}{|x-y|^\lambda} h(y)\, \mathrm{d}x\, \mathrm{d}y}{\|h\|_p^2}$$

where

$$h(x) = \left(\frac{1}{1+|x|^2}\right)^{(2d-\lambda)/2} .$$

and $p = \frac{2d}{2d-\lambda}$.

Moreover, there is equality if and only if for some $x_0 \in \mathbb{R}^d$ and $s \in \mathbb{R}_+$, f is a non-zero multiple of $h(x/s - x_0)$.

[a] E. Lieb, Ann. Math 1983.

For $d \geq 3$, the $\lambda = d - 2$ case of the sharp HLS inequality can be proved by using the fast diffusion flow for $m = d/(d+2)$.

Conclusions

- A new Liapunov functional unveiled for the critical mass PKS: a bassin of attraction determined for each stationary state.
- Optimal Transportation Tools are crucially used to construct the solutions.
- A technique of Concentration Controlled inequalities developed to cope with the lack of compactness.
- The duality between the PKS in the critical mass case and the critical nonlinear fast diffusion equation in 2D leads to the discovery of new proofs of log-HLS and HLS inequalities.
- References:

1. Blanchet-Carlen-C. (to appear in J. Funct. Analysis).
2. Carlen-C.-Loss (PNAS 2010).
3. Carlen-Figalli (preprint).

Conclusions

- A new Liapunov functional unveiled for the critical mass PKS: a bassin of attraction determined for each stationary state.
- Optimal Transportation Tools are crucially used to construct the solutions.
- A technique of Concentration Controlled inequalities developed to cope with the lack of compactness.
- The duality between the PKS in the critical mass case and the critical nonlinear fast diffusion equation in 2D leads to the discovery of new proofs of log-HLS and HLS inequalities.
- References:

1. Blanchet-Carlen-C. (to appear in J. Funct. Analysis).
2. Carlen-C.-Loss (PNAS 2010).
3. Carlen-Figalli (preprint).

Conclusions

- A new Liapunov functional unveiled for the critical mass PKS: a bassin of attraction determined for each stationary state.
- Optimal Transportation Tools are crucially used to construct the solutions.
- A technique of Concentration Controlled inequalities developed to cope with the lack of compactness.
- The duality between the PKS in the critical mass case and the critical nonlinear fast diffusion equation in 2D leads to the discovery of new proofs of log-HLS and HLS inequalities.
- References:

1. Blanchet-Carlen-C. (to appear in J. Funct. Analysis).
2. Carlen-C.-Loss (PNAS 2010).
3. Carlen-Figalli (preprint).

Conclusions

- A new Liapunov functional unveiled for the critical mass PKS: a bassin of attraction determined for each stationary state.
- Optimal Transportation Tools are crucially used to construct the solutions.
- A technique of Concentration Controlled inequalities developed to cope with the lack of compactness.
- The duality between the PKS in the critical mass case and the critical nonlinear fast diffusion equation in 2D leads to the discovery of new proofs of log-HLS and HLS inequalities.
- References:

1. Blanchet-Carlen-C. (to appear in J. Funct. Analysis).
2. Carlen-C.-Loss (PNAS 2010).
3. Carlen-Figalli (preprint).

Conclusions

- A new Liapunov functional unveiled for the critical mass PKS: a bassin of attraction determined for each stationary state.
- Optimal Transportation Tools are crucially used to construct the solutions.
- A technique of Concentration Controlled inequalities developed to cope with the lack of compactness.
- The duality between the PKS in the critical mass case and the critical nonlinear fast diffusion equation in 2D leads to the discovery of new proofs of log-HLS and HLS inequalities.
- References:

1. Blanchet-Carlen-C. (to appear in J. Funct. Analysis).
2. Carlen-C.-Loss (PNAS 2010).
3. Carlen-Figalli (preprint).