Entropy Conditions for some Flux Limited Diffusion Equations Benasque, August 29-September 9, 2011

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Introduction: RHE

$$u_t = \nu \operatorname{div} \left(\frac{u D u}{\sqrt{u^2 + \frac{\nu^2}{c^2} |Du|^2}} \right).$$
(1)

• $\nu > 0$ represents a mean free path, c the speed of light

This type of equations was introduced by J.R. Wilson (circa 1960) as a phenomenological model to control the speed of diffusion:

• when $c \to \infty$ (mean free path small with respect to *c*): the solution goes to

$$u_t = \nu \,\Delta u. \tag{2}$$

 \bullet when $\nu \to \infty$ (free streaming in transparent regions): the solution goes to

$$u_t = c \operatorname{div} \left(u \frac{Du}{|Du|} \right) \tag{3}$$

Introduction: General case

We are interested in equations of the type

• $0 \le u_0 \in L^1(\mathbb{R}^N)$

• $\mathbf{a}(z,\xi) = \nabla_{\xi} f(z,\xi)$ is continuous where $f(z,\xi)$ is convex and differentiable in ξ .

• f is coercive and satisfies the linear growth condition

$$C_0(z)\|\xi\| - D_0(z) \le f(z,\xi) \le M_0(z)(\|\xi\| + 1)$$
(5)

 $\forall (z,\xi) \in \mathbb{R} \times \mathbb{R}^N$, and some positive and continuous functions C_0 , D_0 , M_0 , such that $C_0(z) > 0 \ \forall z \neq 0$.

For the RHE: $f(z,\xi) = \frac{c^2}{\nu} |z| \sqrt{z^2 + \frac{\nu^2}{c^2} |\xi|^2}$.

Brenier's derivation of RHE

The RHE can be derived using a 'gradient descent' with

$$\rho_n^h := \arg \min_{\rho \in \mathcal{P}(\mathbb{R}^N)} \left\{ h W_k^h \left(\rho_{n-1}^h, \rho \right) + E(\rho) \right\}.$$

$$E(\rho) := \int_{\mathbb{R}^N} F(\rho(x)) \, dx, \quad F(\rho) := \nu \rho(\ln \rho - 1)$$
$$V_k^h(\rho_0, \rho_1) := \inf_{\gamma} \left\{ \int_{\mathbb{R}^N \times \mathbb{R}^N} k\left(\frac{x - y}{h}\right) \, d\gamma(x, y) \right\},$$

where γ is a probability measure in $\mathbb{R}^N \times \mathbb{R}^N$ with marginals ρ_0, ρ_1 , and

$$k(z) := \begin{cases} c^2 \left(1 - \sqrt{1 - \frac{|z|^2}{c^2}} \right) & \text{if } |z| \le c \\ +\infty & \text{if } |z| > c. \end{cases}$$
(6)

Formally, we have

$$\frac{\rho_n^h - \rho_{n-1}^h}{h} = \operatorname{div}\left\{\rho_n^h \nabla k^* \left[\nabla \left(F'(\rho_n^h)\right)\right]\right\} + A_n(h)$$

where $A_n(h) \to 0$ as $h \to 0+$ (in the dual of $W^{2,\infty}$).

Recently transformed into a rigorous approach by R. McCann and M. Puel.

Other phenomenological derivations by Ph. Rosenau (1990).

Plan of the talk

• Recall the notion of entropy solution to have existence and uniqueness result.

- \bullet Recall that for the RHE the support of solutions moves with speed c>0
- There are discontinuity fronts propagating at speed *c*
- Interpretation of the notion of entropy solution in terms of the evolution of the discontinuity fronts.
- Final comments

Entropy solutions: motivation

Start with the RHHE model (with c = 1)

$$u_t = \operatorname{div}\left(u\frac{Du}{|Du|}\right) = u\operatorname{div}\left(\frac{Du}{|Du|}\right) + |Du|.$$

Let $u_0(x) = \chi_C(x)$, *C* convex, for instance C = B(0, R). Look for solutions of the form

$$u(t,x) = \alpha(t)\chi_{C(t)}, \quad C(t) = C \oplus B(0,t).$$

Then

$$u_{t} = \alpha'(t)\chi_{C(t)} + \alpha(t)\mathcal{H}^{N-1} \bigsqcup \partial C(t)$$

$$\operatorname{div}\left(u\frac{Du}{|Du|}\right) = \alpha(t)\chi_{C(t)}\operatorname{div}\left(\frac{Du}{|Du|}\right) + \alpha(t)\mathcal{H}^{N-1} \bigsqcup \partial C(t).$$

$$\implies \alpha'(t) = \alpha(t)\operatorname{div}\left(\frac{Du}{|Du|}\right) \quad \text{in } C(t)$$

Assuming that C(t) is calibrable

$$\operatorname{div}\left(\frac{Du}{|Du|}\right) = -\frac{P(C(t))}{|C(t)|} \,\chi_{C(t)}.$$

Hence

$$\alpha'(t) = -\alpha(t) \frac{P(C(t))}{|C(t)|}$$

$$\alpha(t) = \alpha \frac{|C|}{|C(t)|}, \qquad u(t,x) = \alpha \frac{|C|}{|C(t)|} \chi_{C(t)}.$$

Observe that

 $u \in C([0,T], L^1(\mathbb{R}^N))$ $u \in L^1_{\text{loc}}((0,T), BV(\mathbb{R}^N))$ (weakly measurable) $u_t \in \mathcal{M}((0,T) \times \mathbb{R}^N)$ (in this case).

Basic formal estimates:

• Mass preservation:

$$\frac{d}{dt}\int u = \int \operatorname{div} \mathbf{a}(u, Du) = 0$$

• *L*^{*p*}-estimates:

$$\frac{d}{dt}\int u^{p+1} = -(p+1)\int \mathbf{a}(u, Du) \cdot Du^p \le 0$$

• BV-estimate (case RHE)

$$\frac{d}{dt}\int u^2 + 2\int \frac{uDu}{\sqrt{u^2 + |Du|^2}} \cdot Du \le 0$$

Use that $\frac{u|\xi|^2}{\sqrt{u^2+|\xi|^2}} \ge u|\xi| - u^2$ to obtain $\frac{d}{dt} \int u^2 + \int |Du^2| \le C(||u_0||_2)$ $\implies T_{a,b}(u(t)) - a \in BV(\mathbb{R}^N) \quad t > 0, \ 0 < a < b.$

Entropy solutions: motivation

Recall the usual FORMAL uniqueness proof of solutions of:

$$\frac{\partial u}{\partial t} = \operatorname{div} \mathbf{a}(u, Du) \tag{7}$$

If u, v are two solutions with values $u(t), v(t) \in BV(\mathbb{R}^N)$, then

$$\frac{1}{2}\frac{d}{dt}\int |u-v|dx = \frac{1}{2}\int (u-v)_t \operatorname{sign}(u-v) =$$

$$-\int (\mathbf{a}(u, Du) - \mathbf{a}(v, Dv)) \cdot (Du - Dv)\delta_0(u - v) =$$

$$-\int (\mathbf{a}(u, Du) - \mathbf{a}(u, Dv)) \cdot (Du - Dv) \delta_0(u - v) \le 0.$$

Thinking in a Kruzkov's type of proof:

- we need test functions of the form $T(u-l) \in L^1_{loc}((0,T), BV(\mathbb{R}^N))$.
- we have to give sense to the above integrals.

Kruzkov's definition of solution:

Theorem. (Andreu-C-Mazón) $u \in C([0,T]; L^1(\mathbb{R}^N))$ $T_{a,b}(u(\cdot)) \in L^1_{loc,w}(0,T,BV(\mathbb{R}^N)) \ \forall 0 < a < b$, where

 $T_{a,b}(r) = \max\left(\min(r,b),a\right)$

- (i) $u_t = \operatorname{div} \mathbf{a}(u, \nabla u)$ in $\mathcal{D}'(Q_T)$
- (ii) Kruzkov inequalities hold.

 $(i) + (ii) \Longrightarrow$ uniqueness

Let us explain how to write Kruzkov's inequalities.

The Kruzkov's inequalities

We need truncatures S(u - l) to approximate the sign(u - l), $l \in \mathbb{R}$. Write them as S(u). We need truncatures T(u) so that $T(u(t)) \in BV(\mathbb{R}^N)$ ($S = T_{a,b} - a$)

Then formally multiply $u_t = \operatorname{div} \mathbf{a}(u, Du)$ by $S(u)T(u)\phi(t, x)$

$$-\int_0^T \int_{\mathbb{R}^N} J_{TS}(u(t))\phi_t(t)\,dxdt$$

$$\begin{split} \int_{0}^{T} \int_{\mathbb{R}^{N}} \phi \underbrace{S(u) \mathbf{a}(u, Du) \cdot DT(u)}_{S(u)h(u, DT(u))} dt + \int_{0}^{T} \int_{\mathbb{R}^{N}} \phi \underbrace{T(u) \mathbf{a}(u, Du) \cdot DS(u)}_{T(u)h(u, DS(u))} dt + \\ \int_{0}^{T} \int_{\mathbb{R}^{N}} \mathbf{a}(u, Du) \cdot \nabla \phi \ T(u)S(u) \, dx dt \leq 0 \end{split}$$

 $\forall \phi$ test function, J_{TS} the primitive of T(r)S(r).

The term $S(u)\mathbf{a}(u,Du)\cdot DT(u)$

• Let $F(z,\xi)$ be one of the functions

$$f(z,\xi)$$
 or $h(z,\xi)$

Then if $w \in BV(\mathbb{R}^N)$ we know how to define the measure S(w)F(w,Dw):

$$\begin{aligned} \langle S(w)F(w,Dw),\phi\rangle &:= \int \phi(x)S(w)F(w,\nabla w) \, dx \\ &+ \int \phi(x)(SF)^{\circ}(\tilde{w}(x),\frac{D^{c}w}{|D^{c}w|}(x)) \, d|D^{c}w| \\ &+ \int \int_{w^{-}(x)}^{w^{+}(x)} \phi(x)(SF)^{\circ}(s,\nu_{w}(x)) \, d\mathcal{H}^{N-1} \end{aligned}$$

where $(SF)^{\circ}(z,\xi) := \lim_{t\to 0+} tS(u)F(u,\frac{\xi}{t})$ is the recession function. $\tilde{w}(x)$ is the approximate limit of w at x

Functional calculus:

If $T = T_{a,b} + c$, $c \in \mathbb{R}$,

 $\langle S(u)F(u,DT(u)),\phi\rangle := \langle S(T_{a,b}(u))F(T_{a,b}(u),DT_{a,b}(u)),\phi\rangle +$

$$\int_{[u < a]} \phi(x) [S(u)F(u, 0) - S(a)F(a, 0)] + \int_{[u > b]} \phi(x) [S(u)F(u, 0) - S(b)F(b, 0)]$$

When need the lower semicontinuity results for F(w, Dw) when $F(z, \xi)$ is convex in ξ (G.Buttazzo, G. Dal Masso, V. De Cicco- N. Fusco- A. Verde).

Entropy inequalities

Entropy inequalities:

 $S(u)h(u, DT(u)) + T(u)h(u, DS(u)) \le -(J_{TS}(u(t)))_t + \operatorname{div}(\mathbf{z}T(u)S(u))$

where $\mathbf{z} = \mathbf{a}(u, \nabla u)$

 $\forall S, T \in TSUB \cup TSUPER$ (positive truncatures)

 $\forall \phi$ test funtion.

 $(S,T) \in \mathcal{TSUB}$ truncatures: $S \ge 0, S' \ge 0$ and $T \ge 0, T' \ge 0$.

 $(S,T) \in \mathcal{TSUPER}$ truncatures: $S \leq 0, S' \geq 0$ and $T \geq 0, T' \leq 0$.

Analysis of the entropy conditions requires test functions in TSUB and TSUPER.

Existence and Uniqueness

Theorem: Existence and Uniqueness (Andreu-C-Mazón)

Under assumptions on $f(z,\xi)$ and $\mathbf{a}(z,\xi)$:

For any $0 \le u_0 \in L^{\infty}(\mathbb{R}^N) \cap L^1(\mathbb{R}^N)$ there exists a unique entropy solution u of

$$u_t = \operatorname{div} \mathbf{a}(u, Du) \quad \text{in } Q_T = (0, T) \times \mathbb{R}^N$$

 $u(0) = u_0$

If u(t), $\overline{u}(t)$ are the entropy solutions corresponding to initial data u_0 , $\overline{u}_0 \in (L^{\infty}(\mathbb{R}^N) \cap L^1(\mathbb{R}^N))^+$, then

$$|(u(t) - \overline{u}(t))^+||_1 \le ||(u_0 - \overline{u}_0)^+||_1$$
 for all $t \ge 0$. (8)

Use Crandall-Liggett's scheme and Kruzkov's technique.

Evolution of the support of RHE

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Theorem (Andreu-C-Mazón-Moll)
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Let C be an open bounded set in \mathbb{R}^N .

Let $u_0 \in (L^1(\mathbb{R}^N) \cap L^\infty(\mathbb{R}^N))^+$ with support equal to \overline{C} .

Let u(t) be the entropy solution of RHE with $u(0) = u_0$. If we assume that

(*) $u_0 >> 0$ inside its support,

then

 $\operatorname{supp}(u(t)) = \overline{\operatorname{supp}(u_0) \oplus B(0, ct)}$ for all $t \ge 0$.

The Rankine-Hugoniot condition

Let $u \in BV_{loc}((0,T) \times \mathbb{R}^N)$ and let $\mathbf{z} \in L^{\infty}([0,T] \times \mathbb{R}^N, \mathbb{R}^N)$ be such that $u_t = \operatorname{div} \mathbf{z}$. Let (ν_t, ν_x) be the normal to J_u .

We define the speed of the discontinuity set of u as $v(t, x) = \frac{\nu_t(t, x)}{|\nu_x(t, x)|}$ \mathcal{H}^N -a.e. on J_u .

Observe that $H^N(\{(t, x) \in J_u : \nu_x(t, x) = 0\}) = 0.$

Proposition For \mathcal{L}^1 almost any t > 0 we have

 $[u(t)](x)v(t,x) = [[\mathbf{z} \cdot \nu^{J_{u(t)}}]]_{+-} \qquad \mathcal{H}^{N-1}\text{-a.e. in } J_{u(t)}, \tag{9}$

where $[[\mathbf{z} \cdot \nu^{J_{u(t)}}]]_{+-}$ denotes the difference of traces from both sides of $J_{u(t)}$.

The Rankine-Hugoniot condition

Let *B* be a Borel set of J_u contained in the boundary of an open Lipschitz set. Then by integration by parts:

$$[u]\nu_t \mathcal{H}^N|_B = [[\mathbf{z} \cdot \nu^{J_{u(t)}}]]_{+-} \mathcal{H}^{N-1}|_{J_{u(t)} \cap B} dt.$$

We have by disintegration:

$$\nu_x \mathcal{H}^N|_{J_u} = \nu^{J_{u(t)}} \mathcal{H}^{N-1}|_{J_{u(t)}} dt$$

Hence (by the definition of v):

$$\nu_t \mathcal{H}^N|_{J_u} = \frac{\nu_t}{|\nu_x|} |\nu_x| \mathcal{H}^N|_{J_u} = v \mathcal{H}^{N-1}|_{J_{u(t)}} dt.$$

We obtain

$$[u]v\mathcal{H}^{N-1}|_{J_{u(t)}\cap B} dt = [[\mathbf{z} \cdot \nu^{J_{u(t)}}]]_{+-}\mathcal{H}^{N-1}|_{J_{u(t)}\cap B} dt.$$

This implies the conclusion.

Is u_t a Radon measure ?

Proposition For the RHE: $u_0 \in L^1(\mathbb{R}^N) \cap L^\infty(\mathbb{R}^N)$, $u_0(x) \ge 0$. Let Γ_i , $i = 0, \ldots, \ell$, be the boundaries of bounded open sets of class $C^{1,1}$. Assume that

- (i) dist $(\Gamma_i, \Gamma_j) > 0$ for any $i \neq j$,
- (ii) $u_0 \in W^{2,1}(\mathbb{R}^N \setminus \cup_{i=0}^{\ell} \Gamma_i)$ and $\nabla u_0 \in L^{\infty}(\mathbb{R}^N \setminus \cup_{i=0}^{\ell} \Gamma_i)$,
- (iii) u_0 is discontinuous in Γ_i .
- (iv) u_0 is either 0 or is bounded away from zero in any connected component of $\mathbb{R}^N \setminus \bigcup_{i=0}^{\ell} \Gamma_i$.
- (v) When the trace $u_0^-|_{\Gamma_i}$ is bounded away from zero, and we are on the side corresponding to the upper (resp. lower) trace of u_0 , the direction of the gradient of u_0 and the normal to Γ_i are not aligned near the points where u_0 is increasing (resp. decreasing) towards Γ_i .

If u(t) is the solution of the RHE with $u(0) = u_0$, then u_t is a Radon measure.

Characterization of entropy conditions

Let $u \in C([0,T]; L^1(\mathbb{R}^N)) \cap BV_{\text{loc}}((0,T) \times \mathbb{R}^N).$

Assume that u is the entropy solution of

$$(RHEm) u_t = \nu \operatorname{div} \left(\frac{u^m \nabla u}{\sqrt{u^2 + \frac{\nu^2}{c^2} |\nabla u|^2}} \right). (10)$$

$$u(0) = u_0 \in (L^1(\mathbb{R}^N) \cap L^\infty(\mathbb{R}^N))^+ \cap BV(\mathbb{R}^N),$$

where $m \ge 1$.

Characterization of entropy conditions

Recall $\mathbf{z}(t, x) = \mathbf{a}(u, Du)$.

Then u is an entropy solution of (RHEm) if and only if for any $(T, S) \in TSUB$ (for any $(T, S) \in TSUB \cup TSUPER$) we have

 $(S(u)h(u, DT(u))^{c} + (T(u)h(u, DS(u))^{c} \le (\mathbf{z}(t, x) \cdot D(T(u)S(u)))^{c}$

and for \mathcal{L}^1 -almost any t > 0 the inequality

 $[STu(t)^{m}]_{+-} - [J_{TS(u^{m})'}(u(t))]_{+-}$

 $\leq -v[J_{TS}(u(t))]_{+-} + [[\mathbf{z}(t) \cdot \nu^{J_{u(t)}}]T(u(t))S(u(t))]_{+-}$

holds \mathcal{H}^{N-1} a.e. on $J_{u(t)}$.

(Write the EC on the jump set.)

Vertical contact angle at the boundary of the evolving front

If m = 1, we assume that $u \in BV_{loc}((0,T) \times \mathbb{R}^N)$. If m > 1 and $u_0 \in BV(\mathbb{R}^N)$, then $u \in BV_{loc}((0,T) \times \mathbb{R}^N)$. Then the entropy conditions hold on the jump set if and only if

$$\left[\frac{u^{m}\nabla u}{\sqrt{u^{2} + \frac{\nu^{2}}{c^{2}}|\nabla u|^{2}}} \cdot \nu^{J_{u(t)}}\right]_{+} = (u^{+})^{m}$$

$$\left[\frac{u^m \nabla u}{\sqrt{u^2 + \frac{\nu^2}{c^2} |\nabla u|^2}} \cdot \nu^{J_{u(t)}}\right]_{-} = (u^-)^m.$$

Moreover the velocity of the discontinuity fronts is

$$v = \frac{(u^+)^m - (u^-)^m}{u^+ - u^-}.$$
(11)

In case m = 1, we have v = 1.

Asymptotic limit as $c \to \infty$

Theorem

Let u_c be the entropy solution of (RHE) with $u(0,x) = u_0(x) \in (L^{\infty}(\mathbb{R}^N) \cap L^1(\mathbb{R}^N))^+$. As $c \to \infty$, u_c converges in $C([0,T], L^1(\mathbb{R}^N))$ to the solution of U the heat equation

$$u_t = \nu \Delta u$$

with $u(0, x) = u_0(x)$.

Basic estimate:

Assume that $u_0 > 0$ satisfies $\left\| \frac{\nabla u_0}{u_0} \right\|_{\infty} < \infty$.

Let u be the entropy solution of RHE. Then for any t > 0, u(t) satisfies

$$\sup_{[0,T]\times\mathbb{R}^N} \frac{|\nabla u(t,x)|}{|u(t,x)|} \le \left\|\frac{\nabla u_0}{u_0}\right\|_{\infty}.$$
(12)

Asymptotic limit as $c \to \infty$

If u(x) > 0 for all $x \in \mathbb{R}^N$, define $u = e^v$. If u is an entropy solution of

$$u_t = \operatorname{div}\left(\frac{uDu}{\sqrt{u^2 + |Du|^2}}\right). \tag{13}$$

Then v(t, x) is a solution of

$$v_t = \operatorname{div}\left(\frac{\nabla v}{\sqrt{1+|\nabla v|^2}}\right) + \frac{|\nabla v|^2}{\sqrt{1+|\nabla v|^2}}.$$
(14)

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