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- Outline

Flavour Physics and New Physics:

High Precision calculations: Lattice QCD

Opportunity of new physics and Lattice QCD

Charged Current modes:  $K \rightarrow \mu v$ ,  $B \rightarrow \tau v$ ,  $B \rightarrow D \tau v$ :  $H^+$ , right-handed current ...

⇒ Flavour Changing Neutral Current observables:  $\Delta M_s$ ,  $\varepsilon_{K,}$ B→K\*µµ : H<sup>0</sup>, gluino ...

Super B, Benasque, January 18-21, 2011

#### Flavour Physics and New Physics

Next Hints for the Quark Flavour sector mostly driven by the modes



➡ For these modes, hadronic uncertainties are not a big issue!!:

However, to discriminate among NP models we will need much larger set of flavour-changing processes =>

 $|\mathbf{B} \rightarrow \tau v, \mathbf{B} \rightarrow \mathbf{D} \tau v, \mathbf{B} \rightarrow \pi(\rho) \mathbf{I} v, \Delta \mathbf{M}_{s}, \varepsilon_{K}, \mathbf{B} \rightarrow \mathbf{K}^{*} \mu \mu$ 

Super B

**EXEX POINT**: to keep under very good control hadronic uncertainties



#### Flavour Physics: New Physics vs QCD effects

Changes of quark flavour inside a hadron are weak interaction processes:

 $B \longrightarrow B^{-meson}$   $B \longrightarrow$ 

 $->L_{eff}(M_w)->OPE$ 

 Due to confinement, QCD effects are significant and non-perturbative QCD effects absorbed into hadronic matrix elements

A task for Lattice QCD -> tremendous progress in the last years!

## Chiral Regime of QCD

#### A problem with (old) lattice simulations: Heavy pion mass!



**Incompatibility with ChPT?** 

# Chiral Regime of QCD

#### No longer problems with (new) lattice simulations Good compatibility with ChPT



ETMC arXiV:0803.0224

2006-08, unquenched, Wilson-like fermions ©

#### Similar plots by other collaborations (using different approaches)





Lattice study from first principle action still need some time but progress already visible



#### **Lattice QCD and Flavour Physics**



## $f_{\rm K}/f_{\pi}$ : an example of Modern lattice measure



Hadronic uncertainties from  $\langle 0 | \bar{s} \gamma^{\mu} \gamma_{5} u | K \rangle = p^{\mu} f_{K}$   $\langle 0 | \bar{d} \gamma^{\mu} \gamma_{5} u | \pi \rangle = p^{\mu} f_{\pi}$ 

## $f_{\rm K}/f_{\pi}$ : lattice enters high precision era

\* No Competition from non-Lattice approaches:

 $f_K/f_{\pi} = 1 + (\mathbf{m}_K^2 - \mathbf{m}_{\pi}^2) \times \text{unknown-coef.}$ 







• BMW result: N<sub>F</sub>=2+1, *Clover*, a=0.06->0,  $m_{\pi} \ge 190$  MeV, L  $m_{\pi} \sim 4$  $f_{K}/f_{\pi}=1.192(7)$  (6) =>  $\sigma_{rel} \sim 0.75\%$  ( $\delta_{SU(3)} \sim 5\%$ )

Durr, Fodor, Hoebling, Katz, Krieg, Kurth, Lellouch, Lippert, Ramos, Szabo, hep-lat 1001.4692

<u>Constrain on scalar densities:</u>  $SM + (\overline{S}_R u_L) (\overline{\ell} v_L)$ 

$$\frac{B(K \to \mu v)}{B(\pi \to \mu v)} \times \frac{B(n \to p l v)}{B(K \to \pi l v)} = \left(\frac{f_K}{f_\pi} \frac{1}{f_+(0)}\right)^2 \times \left(1 + g^H \frac{m_K^2}{m_{H^+}^2}\right)^2 \qquad Br(B \to \tau v) \propto \left|V_{ub}\right|^2 f_B^2 m_B m_\tau^2 \times \left(1 + g^H \frac{m_B^2}{m_{H^+}^2}\right)^2$$



$$g^{H} = -\tan\beta^{2}/(1+\varepsilon_{0}\tan\beta)$$

Kaon: *NP highly suppressed* => but favourable *exp.* and *th.* scenario

B-physics: *large NP effects* => but *exp.* and *th.* accuracy under progress



Present picture very reassuring and compatible between B and K physics.

 $2008 f_{\text{Ds}}$  puzzle: 3-4 $\sigma$  deviation with experiments?

## $2008 f_{Ds}$ puzzle: 3-4 $\sigma$ deviation with experiments?





"Puzzle" seems to disappear:

No conclusive evidence for New Physics in the charm quark sector yet, but the  $D_{(s)}$  leptonic decays will continue to help constraining SM extensions

 $\Delta F=2$  observables:  $\mathcal{E}_{K}(B_{K}), \Delta M_{d}$  and  $\Delta M_{s}(fB and Bb)$ 

Important role for the SM Unitarity Triangle analysis [UTfit and CKMfitter] :
strong constraints for non-minimal flavour models:

- SMSSM -> generic down squark insertion [Gabbiani et al. 1994, ...]
- CLittle Higgs [A. Buras et al. 2004]

Lattice role and New Physics Opportunities!

# **CP-Violation in K – K Mixing:** $\varepsilon_K$ and $B_K$



$$\varepsilon_{K} \sim \langle \bar{K}^{0} | \mathcal{H}_{\text{eff}}^{\Delta S=2} | K^{0} \rangle = C(\mu) \cdot \langle \bar{K}^{0} | \overbrace{\bar{s}\gamma_{\mu}(1-\gamma_{5})d}^{Q(\mu)} \overline{s\gamma_{\mu}(1-\gamma_{5})d} | K^{0} \rangle$$

$$\langle \bar{K}^{0} | \boldsymbol{Q}(\boldsymbol{\mu}) | K^{0} \rangle = \frac{8}{3} f_{K}^{2} m_{K}^{2} B_{K}(\boldsymbol{\mu})$$

Dramatic progress in  $B_{\kappa}$  evaluation from recent simulations



Average by FLAG WG (1011.4408 [hep-lat]):  $B_{K} = 0.724(30)$  (4%)

> current value lower than old  $N_F=0$ :  $B_K \sim 0.83(6)$ 

**key observation (2000):** to exploit χ–improved action -> among others Carlos Pena (UAM), Pilar Hernandez (IFIC)

increase the tension in the UTA

## Challenge of B-physics: the multi scale-problem of QCD



hierarchy of disparate physical scales to be covered:

$$\begin{split} \Lambda_{\text{IR}} &= L^{-1} \ll \mathfrak{m}_{\pi} \,, \, \dots \,, \, \mathfrak{m}_{\text{D}} \,, \, \mathfrak{m}_{\text{B}} \, \ll \, \mathfrak{a}^{-1} \,= \Lambda_{\text{UV}} \\ \\ O(e^{-L\mathfrak{m}_{\pi}}) \Rightarrow & L \gtrsim \frac{4}{\mathfrak{m}_{\pi}} \sim 6 \, \text{fm} \, \Big\} \ \curvearrowright \ L/\mathfrak{a} \gtrsim 120 \ \curvearrowleft \ \Big\{ \mathfrak{a}\mathfrak{m}_{\text{D}} \lesssim \frac{1}{2} \Rightarrow \mathfrak{a} \approx 0.05 \, \text{fm} \Big\} \end{split}$$

Currently, **b** quarks cannot be directly simulate at their physical mass due to large discretization errors (a m<sub>b</sub> « 1)

Generative theory: like NRQCD action

Simulate heavy quark in the charm region and extrapolate to the B:

## $f_B$ - $f_{Bs}$ : unquenched results



## **Comments:**

- Effective theory approach: discretized NRQCD action
  - > Quite Sophisticated procedure!
    - ⇒ larger set of  $1/m_Q$  corrections on the lattice w.r.t the continuum >Improved through O[ $1/m_Q^2$ ,  $a^2$ ] and leading terms O[ $1/m_Q^3$ ] >O[ $\alpha_s^n/(am_Q)$ ] divergences to be subtracted to get the continuum limit
  - On the other hand, large experience by HPQCD (Eduardo Follana, Elvira Gamiz (Granada)).
  - Pratical approach (a la San Tommaso) welcome

- Pratical approach (a la San Tommaso): ETMC



➤ HQET-guided extrapolation
➤ Combined fit with the static point, m<sub>b</sub>=∞
➤ Larger safe windows with smaller a~6 GeV

 $\Delta M_d$  and  $\Delta M_s$ 

$$\Delta Ms :< B \mid \overline{b}_L \gamma^{\mu} s_L \overline{b}_L \gamma^{\mu} s_L \mid \overline{B} > \propto \widehat{B}_{Bs} f_{Bs}^2.$$

$$\Delta Md :< B \mid \overline{b}_L \gamma^{\mu} d_L \overline{b}_L \gamma^{\mu} d_L \mid \overline{B} > \propto \widehat{B}_{Bd} f_{Bd}^2$$

Solve Solve State St

ℬ Continuum Limit – still coarse lattice spacing ℬ

$$\hat{B}_{Bd} = 1.26(11) - > (9\%)$$
  $\hat{B}_{Bs} = 1.33(6) - > (5\%)$ 

reassuring scenario but we need further calculation from several groups!

Average (Laiho, Lunghi, Van de Water '10)

 $f_{Bs}\sqrt{B_{Bs}} = 275(13) \rightarrow 5\%!!? \qquad \xi = 1.243(28) \rightarrow 2\% \qquad \xi \equiv \frac{f_{Bs}\sqrt{B_{Bs}}}{f_{Bd}\sqrt{B_{Bd}}}$ 

 $\Delta F=1 b \rightarrow u, b \rightarrow s transitions:$  $B \rightarrow \pi(\rho) l V$ ,  $B \rightarrow K^* \mu^+ \mu^-,$  $B \rightarrow K^* (\phi) \gamma$ 

very preliminary unquenched activities
-> pratically as bad as quenched era

 $\ensuremath{\varnothing}$  further complication -> form factors = f( $M_B$ ,  $M_{\pi}$ ,  $q^2$ )

$$A(B \to K^* \mu \mu) = \left\langle K^* \mid Heff \mid B \right\rangle = \begin{cases} C_7 \times \overline{b} \sigma^{\mu\nu} sF_{\mu\nu} \\ C_9 \times \overline{b} \gamma_L^{\mu} s \overline{l} \gamma^{\mu} l + C_{10} \times \overline{b} \gamma_L^{\mu} s \overline{l} \gamma^{\mu} \gamma_5 l \\ C_S \times \overline{b}_L s_R \overline{l} l \end{cases}$$

➢Matrix elements ⇔ 7 form factors (Lattice QCD)

$$\begin{split} \langle \mathcal{K}^{*}(p_{\mathcal{K}^{*}}) | \bar{s} \gamma_{\mu} P_{L,R} b | \mathcal{B}(p) \rangle &= i \epsilon_{\mu\nu\alpha\beta} \epsilon^{\nu*} p^{\alpha} q^{\beta} \frac{\mathcal{V}(q^{2})}{m_{B} + m_{K^{*}}} \mp \\ & \mp \frac{1}{2} \bigg\{ \epsilon^{*}_{\mu}(m_{B} + m_{K^{*}}) \mathcal{A}_{1}(q^{2}) - (\epsilon^{*} \cdot q) (2p - q)_{\mu} \frac{\mathcal{A}_{2}(q^{2})}{m_{B} + m_{K^{*}}} - \\ & - \frac{2m_{K^{*}}}{q^{2}} (\epsilon^{*} \cdot q) [\mathcal{A}_{3}(q^{2}) - \mathcal{A}_{0}(q^{2})] q_{\mu} \bigg\}, \\ \langle \mathcal{K}^{*}(p_{K^{*}}) | \bar{s} i \sigma_{\mu\nu} q^{\nu} P_{R,L} b | \mathcal{B}(p) \rangle &= -i \epsilon_{\mu\nu\alpha\beta} \epsilon^{\nu*} p^{\alpha} q^{\beta} T_{1}(q^{2}) \pm \\ & \pm \frac{1}{2} \bigg\{ [\epsilon^{*}_{\mu}(m_{B}^{2} - m_{K^{*}}^{2}) - (\epsilon^{*} \cdot q) (2p - q)_{\mu}] T_{2}(q^{2}) + \\ & + (\epsilon^{*} \cdot q) \bigg[ q_{\mu} - \frac{q^{2}}{m_{B}^{2} - m_{K^{*}}^{2}} (2p - q)_{\mu} \bigg] T_{3}(q^{2}) \bigg\}. \end{split}$$

 $\odot$  Up to now, only  $T_1(0)$  on the Lattice but quenched (Becirevic, Lubicz, F.M. '06)



### $A_{FB}(B_d \rightarrow K^*l^+l^-)$ , already plays a special role: $\Rightarrow$ very large experimental error; $\bigcirc$ Necessary to improve lattice form factors!



Model-independent Analysis in MFV!

FCNC constraints from  $Br(B_d \rightarrow X_s \gamma)$ ,  $Br(B_d \rightarrow X_s l^+ l^-)$ ,  $Br(B_d \rightarrow \mu^+ \mu^-)$ ,  $\Delta M_s$ 

Hurth, Isidori, Kamenik, F.M '08



Lattice QCD for D, NF=2 unquenched



**Figure 3:** Vector  $[f_+(q^2)]$  and scalar  $[f_0(q^2)]$  form factors for the  $D \to \pi \ell v_\ell$  decay versus the squared

#### ETMC et al., '09

# Lattice Flavour Physics in Spain

# ③ 3 international collaborations: ETMC, HPQCD, CLS <u>⑧ 1 Mare Nostrum!</u>



## **Conclusions:**

### LATTICE QCD -> touchable progress in recent years:

⇒ reliable unquenched simulations with pions close to the physical point =>  $m_{\pi}$ =156 MeV (PACS-CS),  $m_{\pi}$ =190 MeV (BMW)

- $\supset$   $f_{\kappa}/f_{\pi}$  paradigma of present lattice progress!
- many other calculations should be come soon!

In particular, for B-physics observables



#### Super B and Lattice QCD Alliance

Constraining  $V_{us}$ ,  $V_{ub}$ ,  $\rho$ ,  $\eta$  as precisely as possible from as many independent modes as possible in order to search for New Physics!







 $K_{I3}$ : Higgs effects =>  $f_0(q^2) \rightarrow f_0(q^2)(1 + g^H q^2 / m_H^2)$ 

• to test  $g^{H} =$  <u>calculate</u> slope and curvature of  $f_{0}(q^{2})!$ 

 $|V_{us}|^{KI3}$  is NP free as soon as we use  $f_0(q^2)$  from data

 $K_{I2}$ : Higgs effects =>  $|V_{us}|^{KI2}$ ->  $|V_{us}|^{KI2}$  (1 +  $g^H m^2_K / m^2_H$ )

• to test  $g^{H} = f_{K}$  from theory:  $f_{K}/f_{\pi}$  better determined

$$\frac{B(K \to \mu v)}{B(\pi \to \mu v)} \times \frac{B(n \to plv)}{B(K \to \pi lv)} = \left(\frac{f_K}{f_\pi} \frac{1}{f_+(0)}\right)^2 \times \left(1 + g^H \frac{m_K^2}{m_{H^+}^2}\right)^2 (\text{known f.s})$$

similarly search on H<sup>+</sup> from

 $B^{\pm} \rightarrow \tau^{\pm} \nu$ 

$$Br(B \to \tau \nu) \propto \left| V_{ub} \right|^2 f_B^2 m_B m_\tau^2 \times \left( 1 + g^H \frac{m_B^2}{m_{H^+}^2} \right)^2$$

hadronic uncertainty  $-> f_B$ : ~5% accuracy from Lattice (Lunghi et al. 2009)

Best for indirect H<sup>+</sup> searches but only feasible at •<u>SuperB</u>

**Only scalar component sensitivity** to H<sup>+</sup>: opportunity for <u>Lhcb?</u>

## but different theoretical and experimental prospects



Quenched simulation.

Staggered fermions

$$\frac{d\Gamma(B \to D\tau \nu)}{dq^2} \propto \left| V_{cb} \right|^2 \rho_{\nu}(q^2) \times \left( 1 - \frac{m_{\tau}^2}{m_B^2} \left| 1 + g^H \frac{q^2}{m_{H^+}^2} \right|^2 \rho_S(q^2) \right)$$

$$dq^{2} \times \left(1 - \frac{m_{\tau}^{2}}{m_{B}^{2}} \left|1 + g^{H} \frac{q^{2}}{m_{H^{+}}^{2}}\right|^{2} \rho_{S}(q^{2})\right)$$

 $B^{\pm} \rightarrow D \tau^{\pm} v$ 

2. hadronic uncertainty -> 
$$\rho_{v}$$
,  
->  $\rho_{s}$ 

