

# An attempt to clarify the neutrino abundance and mass impact on CMB and $P(k)$

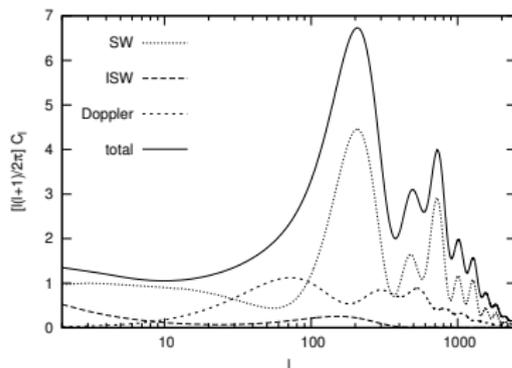
Julien Lesgourgues

EPFL & CERN

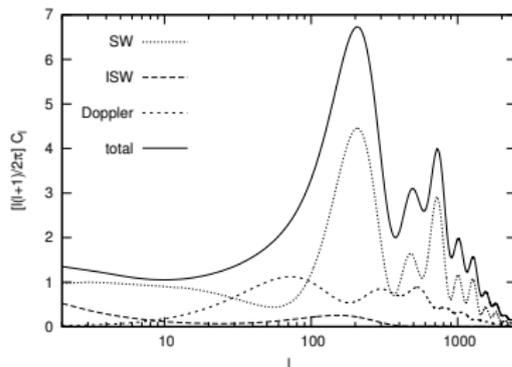
Benasque, 16.08.2012

- 1 CMB
  - CMB parameter dependence
  - Impact of  $N_{\text{eff}}$
  - Impact of  $M_\nu$
- 2 Matter power spectrum
  - $P(k)$  parameter dependence
  - Impact of  $N_{\text{eff}}$
  - Impact of  $M_\nu$
  - Impact of mass splitting

In neutrinoless  $\Lambda$ CDM model,  
 $C_l^{TT}$  controlled by 8  
effects/quantities:

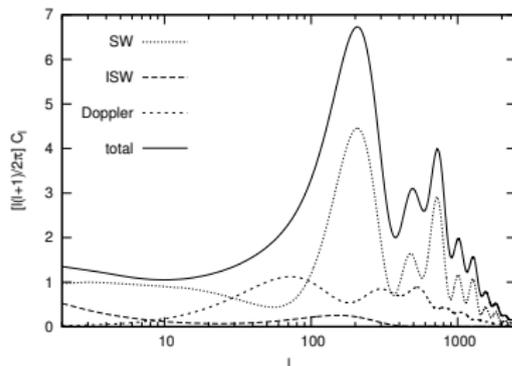


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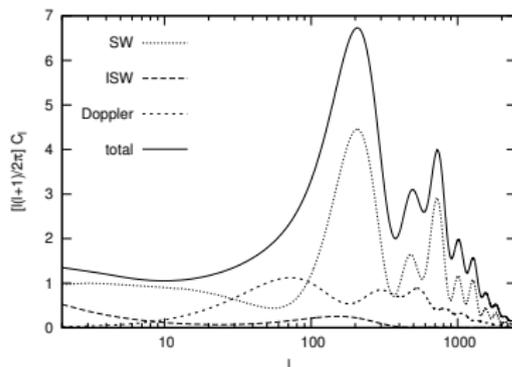
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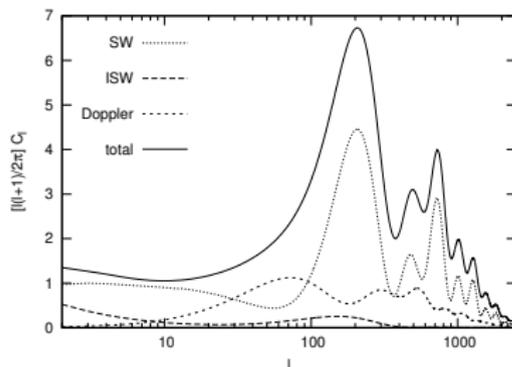
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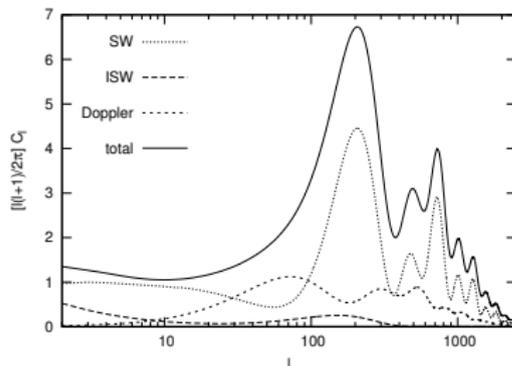
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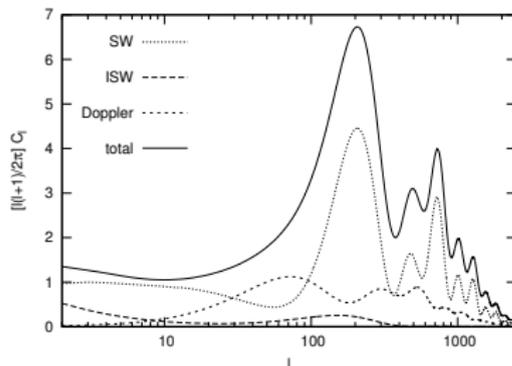
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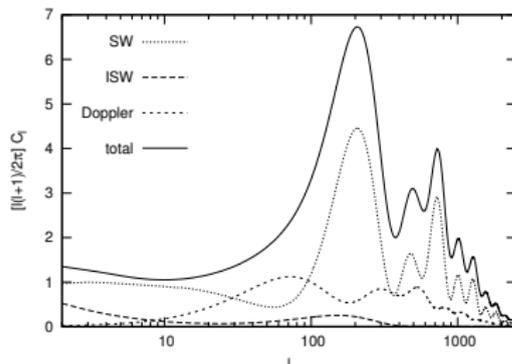
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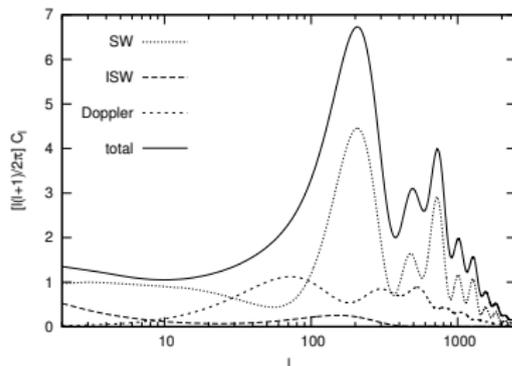
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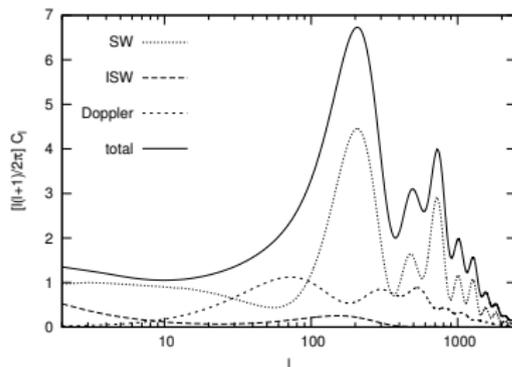


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In terms of parameters

$$\{\omega_m, \omega_b, \Omega_\Lambda, A_s, n_s, \tau_{\text{reio}}\}:$$

(with  $h = \sqrt{\omega_m/(1 - \Omega_\Lambda)}$  and  
 $\omega_m = \omega_b + \omega_c$ )



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## Effective neutrino number

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This method allows to see if the parameter is really detectable, and to “isolate” the direct perturbation effect. Applicable to other physical ingredients...

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- (C8) Relative amplitude for  $l \gg 40$  w.r.t  $l \ll 40$ : optical depth FIXED

We vary  $N_{\text{eff}}$  with fixed  $z_{\text{eq}}$ ,  $z_\Lambda$ . One can easily show that:

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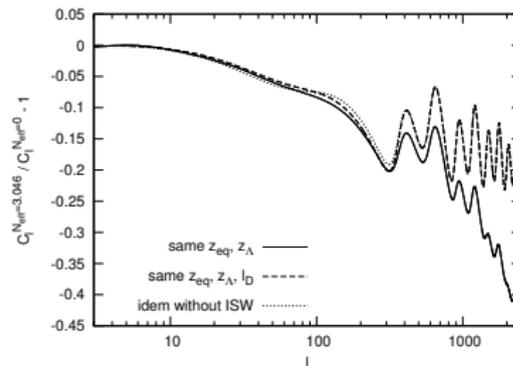
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**(C4)** cannot be compensated in minimal  $\Lambda$ CDM (.. but it can be kept fixed if  $Y_{\text{He}}$  is decreased [Bashinsky & Seljak 2004](#))

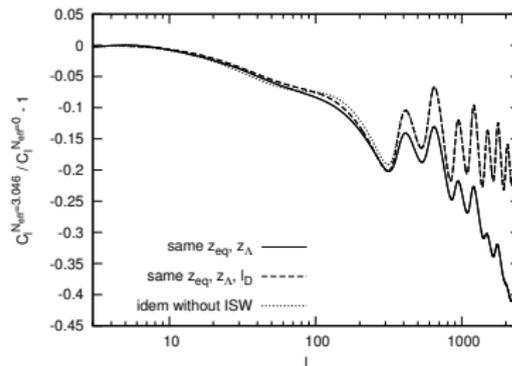
JL, Mangano, Miele, Pastor, in press

Example:  $N_{\text{eff}}$  increased from 0 to 3.046 with constant  $Y_{\text{He}}$  (solid) or (unrealistically) small  $Y_{\text{He}}$  (dashed)



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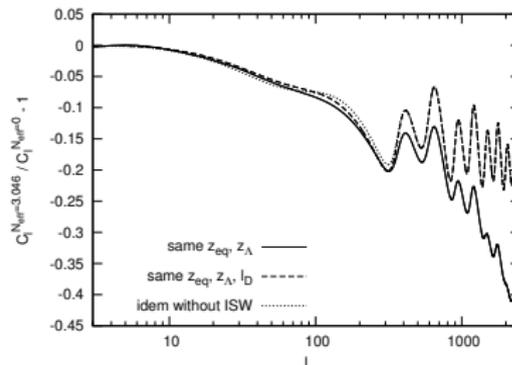
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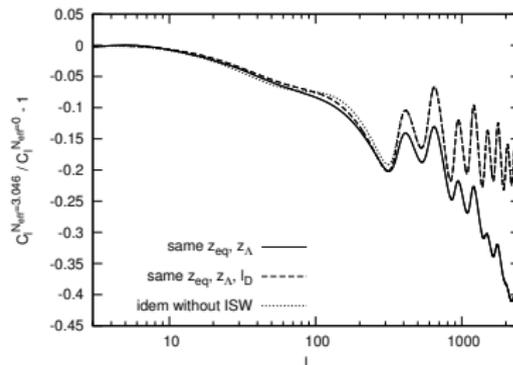
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- peak scale shifted because neutrinos propagate at  $c > c_s$ : “effective sound speed” enhanced (neutrino drag effect):

$$\Delta l \sim -3 \Delta N_{\text{eff}} \quad \text{Bashinsky \& Seljak 2004}$$

### Conclusions:

- $N_{\text{eff}}$  **clearly detectable with CMB** due to background and perturbation effects
- true for minimal  $\Lambda$ CDM and beyond (perturbation effects)
- accurate data at high- $l$  helps
- BBN prior on  $Y_{\text{He}}$  helps
- $H_0$  prior helps (if  $h$  fixed, cannot keep  $z_{\text{eq}}$ ,  $z_\Lambda$  fixed)

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Even without lensing information, WMAP gives  $M_\nu < 1.3 \text{ eV}$  (95% C.L.) and Planck expected to give  $M_\nu < 0.4 \text{ eV}$  (95% C.L.)

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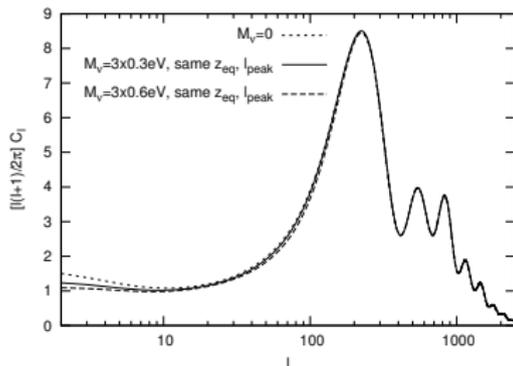
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- **(C8)** Relative amplitude for  $l \gg 40$  w.r.t  $l \ll 40$ : optical depth FIXED

In the parameter set  $\{M_\nu, \omega_c, \omega_b, \Omega_\Lambda, A_s, n_s, \tau_{\text{reio}}\}$ , still have possibility to vary  $\Omega_\Lambda$  in order to fix either

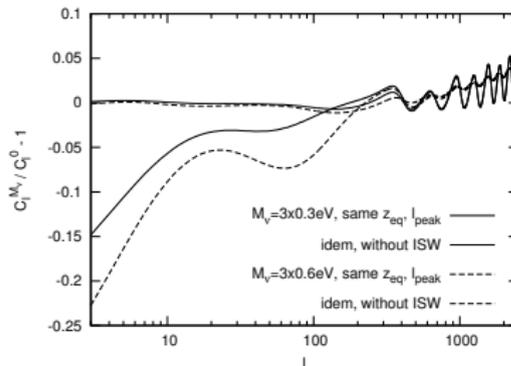
- $d_A(\eta_{LS})$  and **(C1)+(C4)**,
- or  $z_\Lambda$  and **(C7)**.

First option better motivated (cosmic variance).

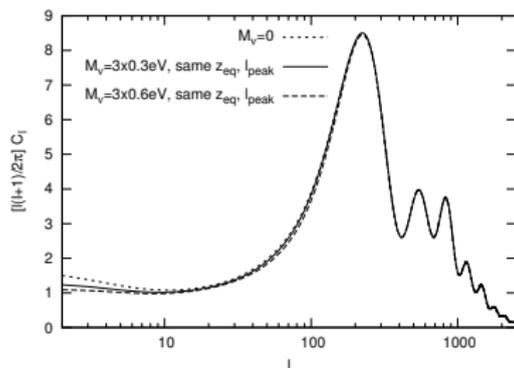
We vary  $M_\nu$  with fixed  $\omega_b$ ,  $\omega_c$ ,  $d_A(\eta_{LS})$ :



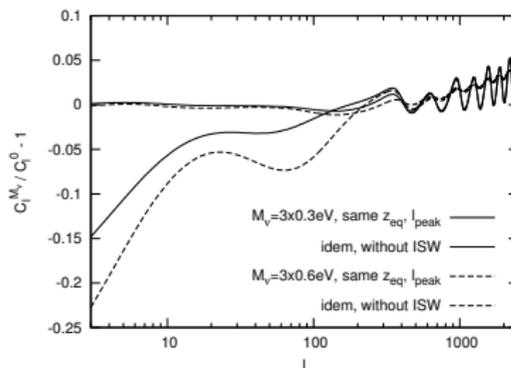
JL, Mangano, Miele, Pastor, in press



We vary  $M_\nu$  with fixed  $\omega_b$ ,  $\omega_c$ ,  $d_A(\eta_{\text{LS}})$ :



JL, Mangano, Miele, Pastor, in press



Only modified **late ISW effect (C7)** plus direct **perturbation effects of extra d.o.f.**  
Later mainly consists in extra early ISW ( $20 < l < 200$ ) due to metric variations when neutrinos become non-relativistic after decoupling. Amplitude:

$$\Delta C_l / C_l \sim [m_\nu / 10 \text{ eV}].$$

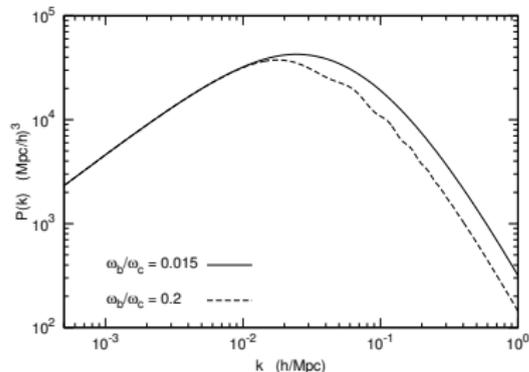
Also effects at  $l > 200$  due to the fact that neutrino not fully relativistic prior to recombination.

### Conclusions:

- $M_\nu$  **difficult to measure**: background effect (= LISW) masked by cosmic variance, and perturbation effect very small.
- extra priors ( $H_0$ , BAO...) help: not possible to keep  $z_{\text{eq}} + d_A(\eta_{LS})$  fixed... then, significant background effect... that could still be compensated in more general cosmology (spatial curvature)
- detecting mass splitting in CMB is hopeless
- lensing extraction helps a lot

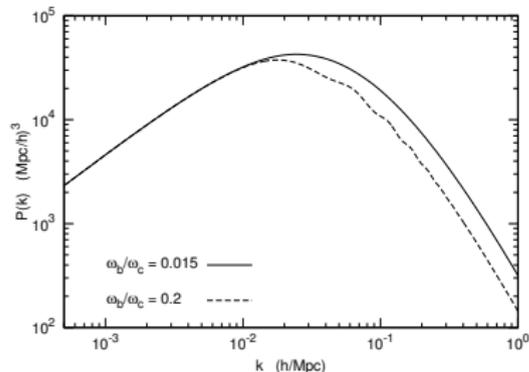
## Matter power spectrum

In neutrinoless  $\Lambda$ CDM model, linear  $P(k)$  (in  $[\text{Mpc}/h]^3$  vs  $[h/\text{Mpc}]$ ) controlled by 5 effects/quantities:



## Matter power spectrum

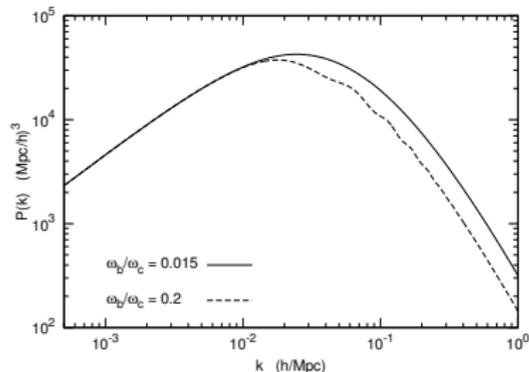
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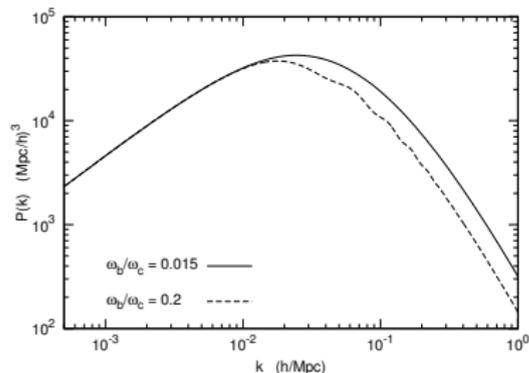
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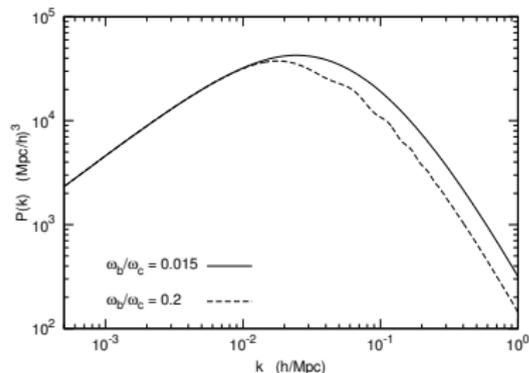
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**BAO amplitude** on Silk damping scale  $\lambda_d(\eta_d)$

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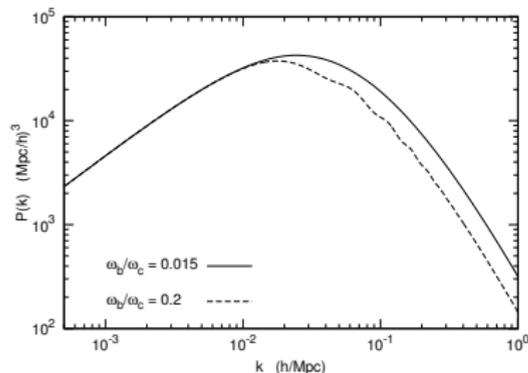
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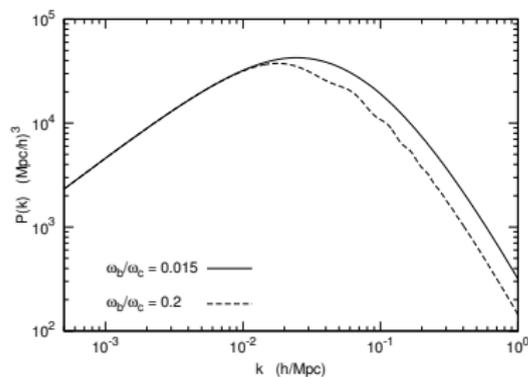
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- (P4) **Overall amplitude:** depends on  $\Omega_m$  and  $A_s$
- (P5) **Global tilt:**  $n_s$

In terms of parameters

$\{\omega_m, \omega_b, \Omega_\Lambda, A_s, n_s, \tau_{\text{reio}}\}$ :

(with  $h = \sqrt{\omega_m / (1 - \Omega_\Lambda)}$  and

$$\omega_m = \omega_b + \omega_c)$$



- **(P1)** Peak location: depends on  $k_{\text{eq}} \sim [\Omega_m(1 + z_{\text{eq}})]^{1/2}$
- **(P2)** Slope/amplitude for  $k \geq k_{\text{eq}}$ : baryon-to-cdm ratio
- **(P3)** BAO phase ( $d_s(\eta_d)$ ), BAO amplitude ( $\lambda_d(\eta_d)$ )
- **(P4)** Overall amplitude
- **(P5)** Global tilt

$\omega_m, \Omega_\Lambda$

$\omega_m, \omega_b$

$\omega_b$

$\Omega_\Lambda, A_s$

$n_s$

# Impact of $N_{\text{eff}}$

Like for CMB, we vary  $N_{\text{eff}}$  with fixed  $\{z_{\text{eq}}, z_\Lambda, \omega_b\}$ , i.e. same  $\Omega_\Lambda, \omega_b$  and varying  $\omega_m$ .

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- **(P3)** BAO phase and amplitude  $\lambda_d(\eta_d)$  **FIXED**
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We could have kept **(P2)** fixed by increasing  $\omega_b$  proportionally to  $\omega_c$ ... then **(P3)** modified (BAO).

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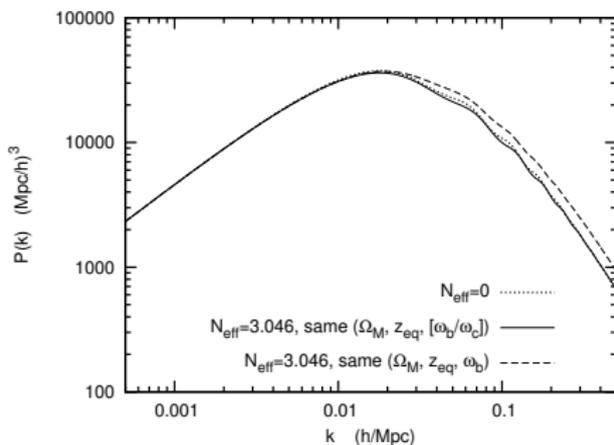
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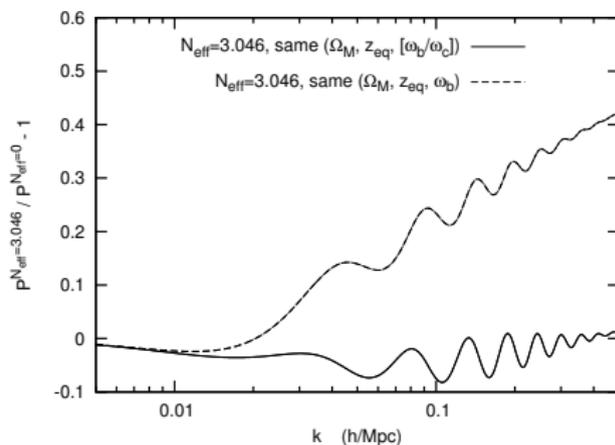
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In both cases, background effect (**(P2)** or **(P3)**) adds up with direct perturbation effect: we expect the amplitude/phase shift observed in CMB to show up in BAOs.

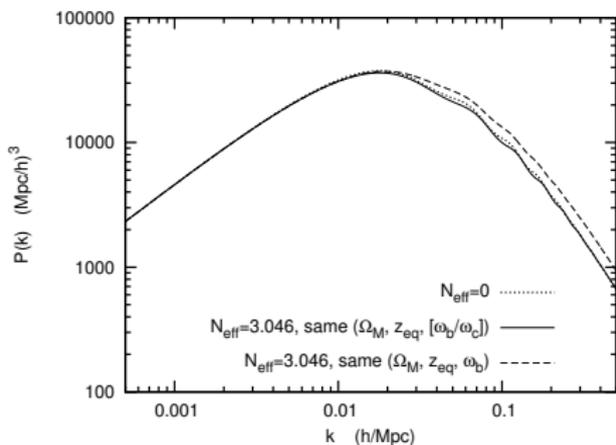
We increase  $N_{\text{eff}}$  with fixed  $z_{\text{eq}}$ ,  $\Omega_\Lambda$ , and either fixed  $\omega_b/\omega_c$  or  $\omega_b$ :



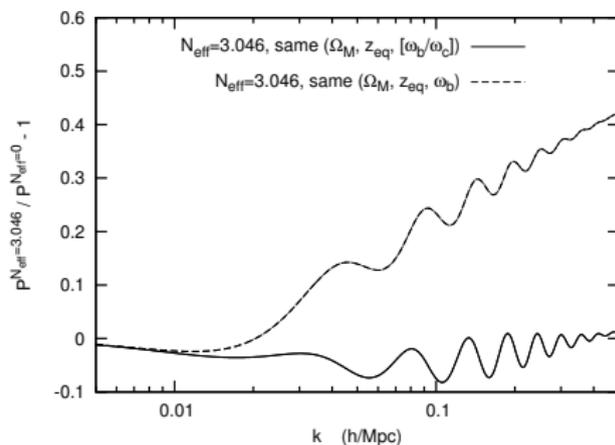
JL, Mangano, Miele, Pastor, in press



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JL, Mangano, Miele, Pastor, in press



Background effect = either change of slope for  $k \geq k_{\text{eq}}$  or in BAO phase.

Always an additional BAO phase shift from perturbation effect (neutrino drag).

### Conclusions:

- matter power spectrum = **complementary probe of  $N_{\text{eff}}$** .
- main signatures in slope (for fixed  $n_s$ ) and in BAO phase.
- currently:  $\sigma(N_{\text{eff}}) \sim 0.7$ ,  $1.0\sigma - 1.9\sigma$  excess.  
Planck with lensing extraction:  $\sigma(N_{\text{eff}}) \sim 0.3$ ,  
Planck + Euclid:  $\sigma(N_{\text{eff}}) \sim 0.1$ .  
No prospects to test with precision standard value 3.046

# Impact of $M_\nu$

Famous **neutrino free-streaming effect**: on small scales, not only neutrinos do not cluster, but also the growth of CDM perturbations is modified (scale-dependent growth factor):

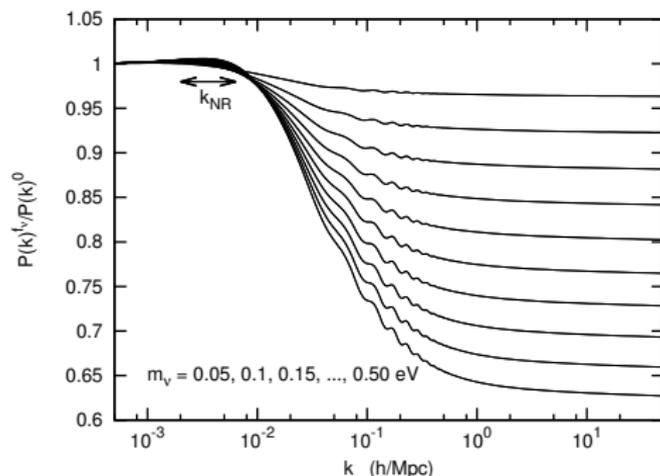
$$\delta_c \propto a^{1-\frac{3}{5}f_\nu}, \quad f_\nu \equiv \omega_\nu/\omega_m.$$

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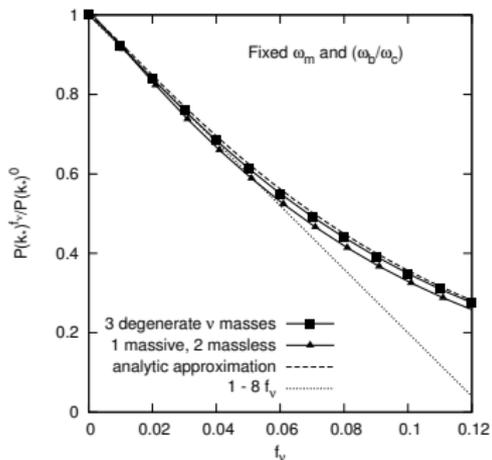
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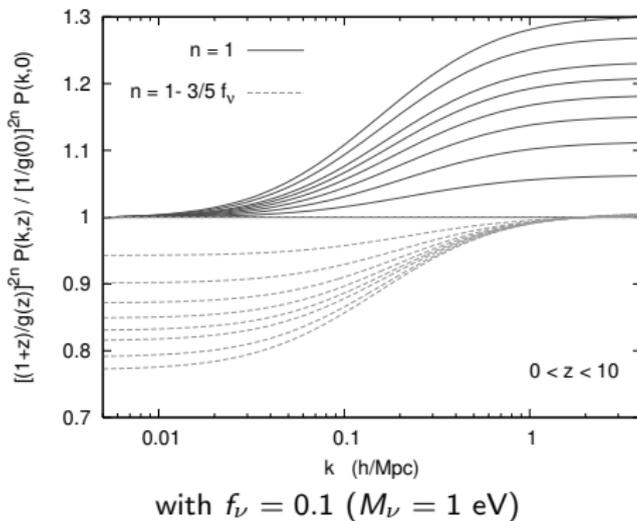
Best seen when most background effects are cancelled: below we fix  $\omega_m, \Omega_\Lambda, \omega_b/\omega_c$ :



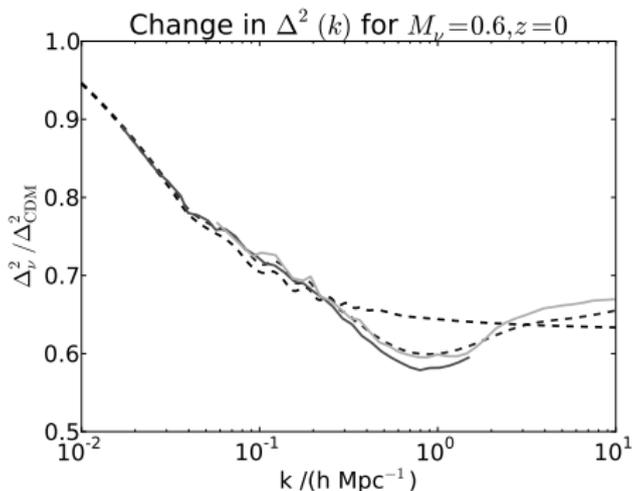
JL, Mangano, Miele, Pastor, in press



Effect is strongly redshift dependent:



Non-linear corrections well understood on mildly non-linear scales:



Bird, Haehnelt, Viel 2011; see also Brandbyge et al.2010

- $P(k)$  very sensitive to  $M_\nu$ ,  $\Delta P/P \sim -8(10)f_\nu \sim [M_\nu/1 \text{ eV}]$ : at least 5% for normal hierarchy scenario, 10% for inverted hierarchy scenario.
- scale-dependent growth factor  $g(z, k)$ : data at different redshift helps; not degenerate with usual extensions of  $\Lambda$ CDM (curvature, running...)
- $\sigma(M_\nu) \sim 0.1 \text{ eV}$  for Planck with lensing extraction or with BOSS,  
0.06 with Planck+DES,  
0.03 with Planck + Euclid CS,  
0.015 with Planck + Euclid  $P(k)$  ...

# Impact of mass splitting

- times of non-relativistic transitions depend on individual masses
- hence 3 free-streaming scales depending on each mass
- total small-scale suppression due to reduced cdm growth rate between non-relativistic transition and now: small dependance on individual mass.

Unlikely to be ever detected even with CMBPol/Core + SKA

