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# Chowla-Selberg and other formulas useful for zeta regularization

EMILIO ELIZALDE

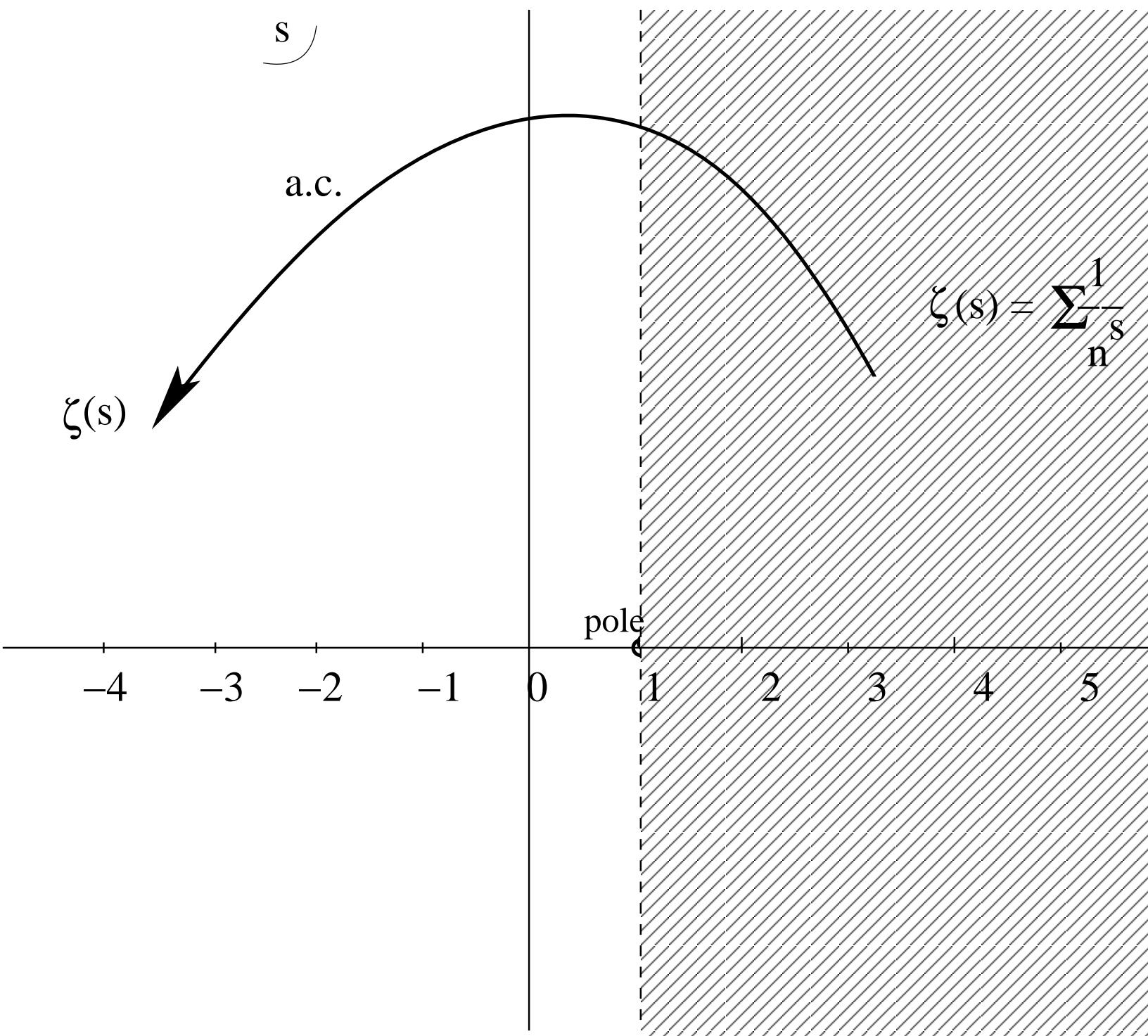
*ICE/CSIC & IIEC, UAB, Barcelona*

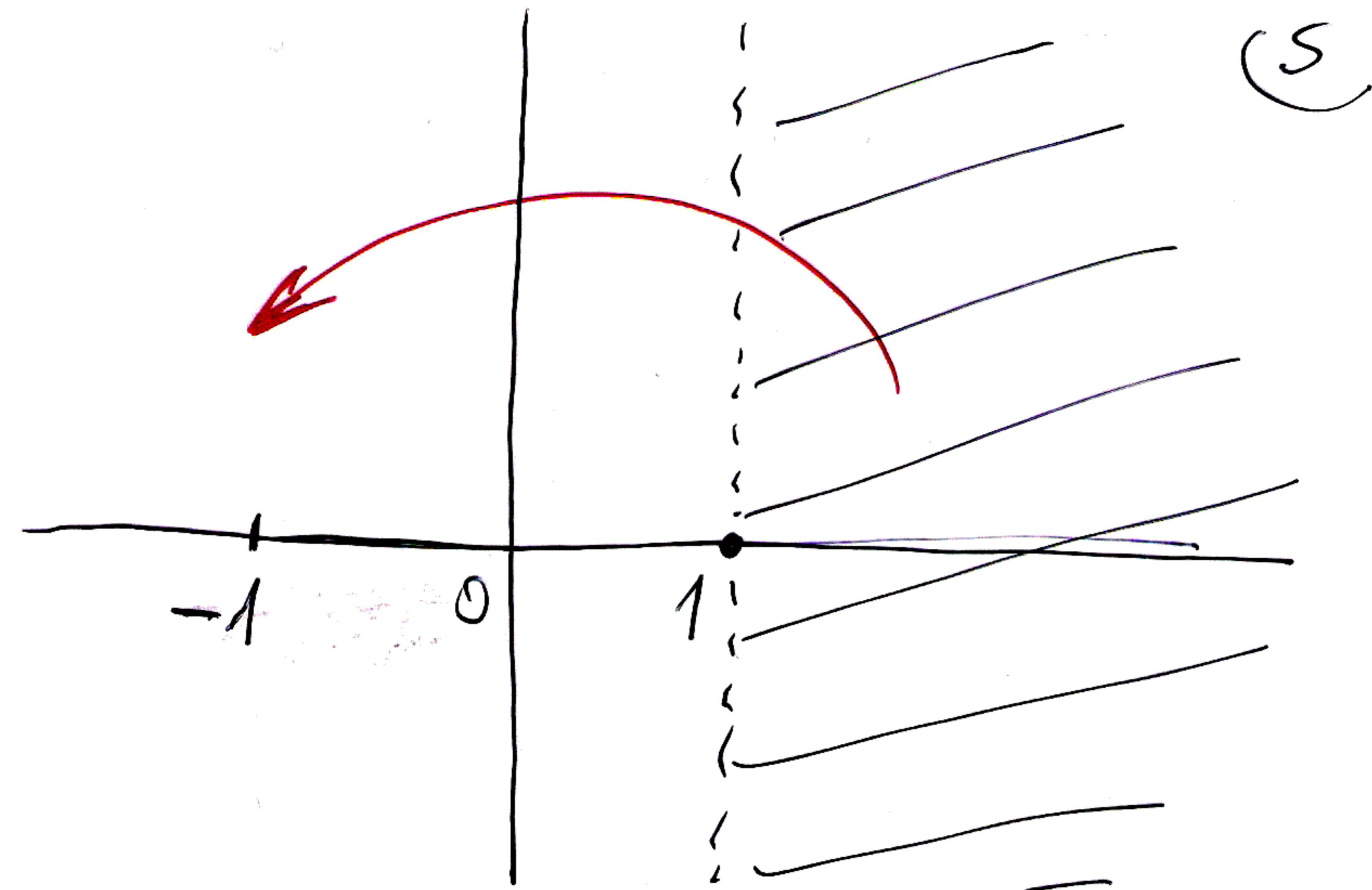
Mathematical Structures in Quantum Systems  
and applications

Benasque, July 8-14, 2012

# Outline

- The Riemann zeta function as a regularization tool.
- General scheme for Linear and Quadratic cases.  
Truncations.
- Spectrum only known Implicitly.
- The Chowla-Selberg formula in Number Theory.
- The Chowla-Selberg series formula (CS). Nontrivial Extensions (ECS).
- Operator Zeta Functions:  $\zeta_A$  for  $A$  a  $\Psi$ DO, Det's.
- Dixmier trace, Wodzicki Residue.
- Multiplicative Anomaly or Defect.





$$\zeta(s) = \sum_{n=1}^{\infty} n^{-s}$$

$$1 + \frac{1}{2} + \frac{1}{3} + \frac{1}{4} + \dots = \infty$$

$$\zeta(0) = -\frac{1}{2} \quad \text{or} \quad 1 + 1 + 1 + \dots = -\frac{1}{2}$$

$$\zeta(-1) = -\frac{1}{12} \quad \text{or} \quad 1 + 2 + 3 + \dots = -\frac{1}{12}$$

⋮

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{\text{prime}} \frac{1}{1 - \frac{1}{p^s}}$$

The prime number theorem

[Hadamard  
de la Vallée Poussin]

$$\pi(x) = \#\{\text{primes } p \leq x\} \sim \frac{x}{\log x}$$

$$\tilde{\zeta}(s) \equiv \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s) = \tilde{\zeta}(1-s)$$

(Gelbart + Miller, BAMS '03) (completed Z 1.)

(E) Entirety:  $\tilde{\zeta}(s)$  meromorphic c.,  $s=0, 1$  poles

$$(\text{FE}) \quad \tilde{\zeta}(s) = \tilde{\zeta}(1-s)$$

(BV) Bounded in vertical strips:

$$\tilde{\zeta}(s) + \frac{1}{s} + \frac{1}{1-s} \text{ bounded } -\infty < a < \text{Res} < b < \frac{x}{2}$$

Riemann (1859)

Poisson 5. for.

$$\sum_{n \in \mathbb{Z}} f(n) = \sum_{n \in \mathbb{Z}} \hat{f}(n)$$

Fouriért.

$$\hat{f}(r) = \int_{-\infty}^{+\infty} f(x) e^{-2\pi i rx} dx, \quad f \text{ Schwartz}$$

1) FE  $\zeta$

$$2) \quad \hat{f}(r) = \frac{1}{\sqrt{t}} e^{-\pi r^2/t} \rightarrow \text{Jacobi id.}$$

$$\theta(it) = \frac{1}{\sqrt{t}} \theta\left(\frac{i}{t}\right), \quad \theta(z) = \sum_{n \in \mathbb{Z}} e^{\pi i n^2 z}$$

$$\text{Dirichlet} \quad L(s, \chi) = \sum_{n=1}^{\infty} \chi(n) n^{-s}$$

Hamburger

# Basic strategies

- Jacobi's identity for the  $\theta$ -function

$$\theta_3(z, \tau) := 1 + 2 \sum_{n=1}^{\infty} q^{n^2} \cos(2nz), \quad q := e^{i\pi\tau}, \quad \tau \in \mathbb{C}$$

$$\theta_3(z, \tau) = \frac{1}{\sqrt{-i\tau}} e^{z^2/i\pi\tau} \theta_3\left(\frac{z}{\tau} \mid \frac{-1}{\tau}\right) \quad \text{equivalently:}$$

$$\sum_{n=-\infty}^{\infty} e^{-(n+z)^2 t} = \sqrt{\frac{\pi}{t}} \sum_{n=0}^{\infty} e^{-\frac{\pi^2 n^2}{t}} \cos(2\pi n z), \quad z, t \in \mathbb{C}, \quad \operatorname{Re} t > 0$$

- Higher dimensions: Poisson summ formula (Riemann)

$$\sum_{\vec{n} \in \mathbb{Z}^p} f(\vec{n}) = \sum_{\vec{m} \in \mathbb{Z}^p} \tilde{f}(\vec{m})$$

$\tilde{f}$  Fourier transform

[Gelbart + Miller, BAMS '03, Iwaniec, Morgan, ICM '06]

- Truncated sums  $\longrightarrow$  asymptotic series

## $\zeta$ : EXPLICIT CALCULATIONS

Epstein zeta functions (quadratic)

$$\zeta_E = \sum_{\vec{n} \in \mathbb{Z}^d} Q(\vec{n})^{-s}$$

$Q$  quadratic form

Barnes zeta functions (linear)

$$\zeta_B = \sum_{\vec{n} \in N^d} L(\vec{n})^{-s}$$

$L$  affine form  
(coeff's  $\in \mathbb{Q}^+$ )

Extensions:

$$\zeta_E \xrightarrow{\text{Q} + L \text{ affine}}$$

$$\sum_{\vec{n} \in N^d} \quad \text{(truncation)}$$

$$\zeta_B \xrightarrow{\quad} \zeta'_B(0) \quad \text{(new formulas)}$$

$$\sum'_{\vec{n} \in \mathbb{Z}^d} \quad \text{(by analyt. cont.)}$$

## $\zeta$ -REGULARIZ: SPECTRUM KNOWN IMPLICITLY

- Example of the ball:

- Operator

$$(-\Delta + m^2)$$

on the  $D$ -dim ball  $B^D = \{x \in R^D; |x| \leq R\}$

with Dirichlet, Neumann or Robin BC

- The zeta function

$$\zeta(s) = \sum_k \lambda_k^{-s}$$

- Eigenvalues implicitly obtained from

$$(-\Delta + m^2)\phi_k(x) = \lambda_k \phi_k(x) \quad + \quad BC$$

- In spherical coordinates:

$$\phi_{l,m,n}(r, \Omega) = r^{1-\frac{D}{2}} J_{l+\frac{D-2}{2}}(w_{l,n}r) Y_{l+\frac{D}{2}}(\Omega)$$

$J_{l+(D-2)/2}$  Bessel functions

$Y_{l+D/2}$  hyperspherical harmonics

- Eigenvalues  $w_{l,n} (> 0)$  determined through BC

$$J_{l+\frac{D-2}{2}}(w_{l,n}R) = 0,$$

for Dirichlet BC

$$\frac{u}{R} J_{l+\frac{D-2}{2}}(w_{l,n} R) + w_{l,n} J'_{l+\frac{D-2}{2}}(w_{l,n} r) \Big|_{r=R} = 0, \text{ for Robin BC}$$

- Now,  $\lambda_{l,n} = w_{l,n}^2 + m^2$

$$\zeta(s) = \sum_{n=0}^{\infty} \sum_{l=0}^{\infty} d_l(D) (w_{l,n}^2 + m^2)^{-s}$$

$w_{l,n} (> 0)$  is defined as the n-th root of the l-th equation,  $d_l(D) = (2l + D - 2) \frac{(l+D-3)!}{l! (D-2)!}$

### • Procedure:

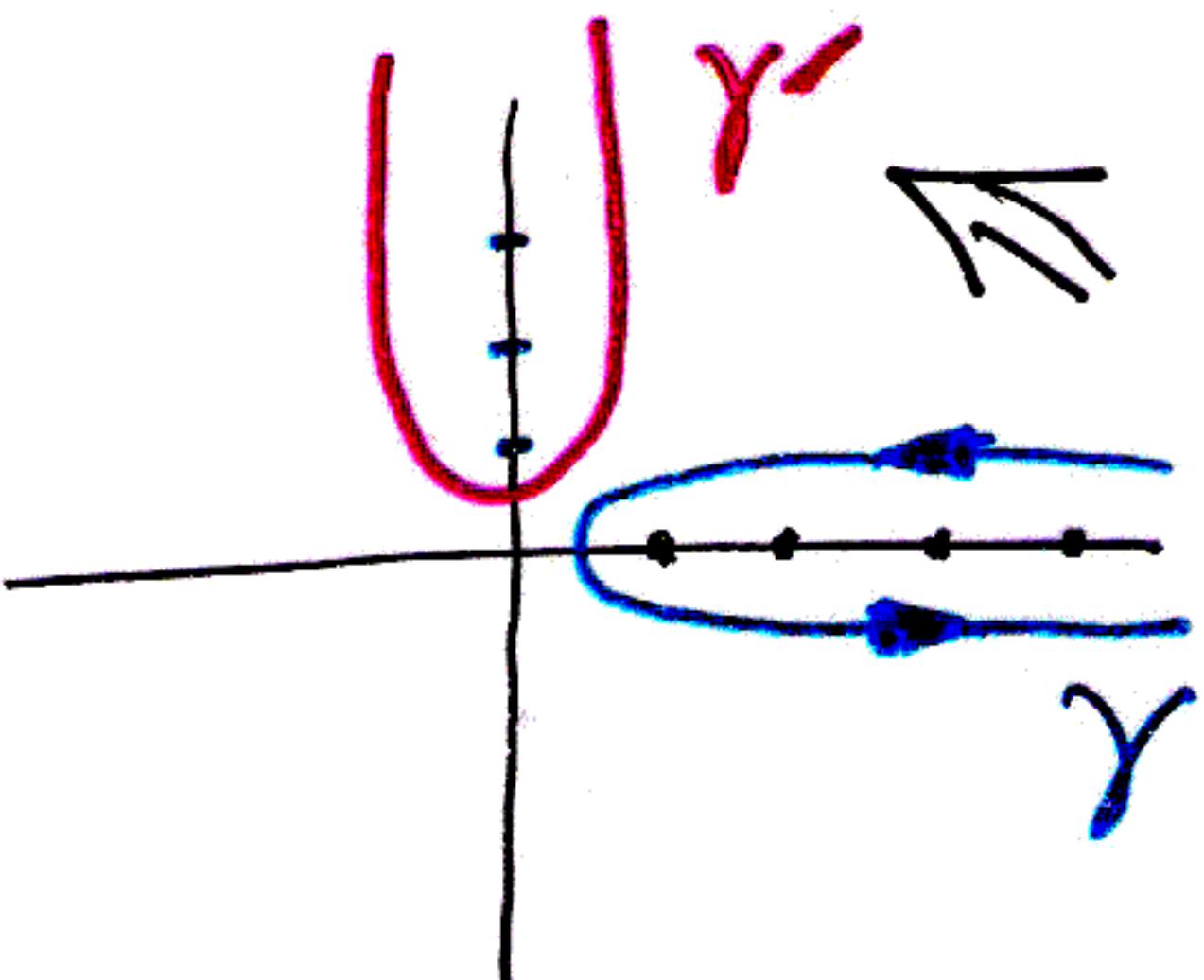
- Contour integral on the complex plane

$$\zeta(s) = \sum_{l=0}^{\infty} d_l(D) \int_{\gamma} \frac{dk}{2\pi i} (k^2 + m^2)^{-s} \frac{\partial}{\partial k} \ln \Phi_{l+\frac{D-2}{2}}(kR)$$

$\gamma$  runs counterclockwise and must enclose all the solutions [Ginzburg, Van Kampen, EE + I. Brevik]

• Obtained: [with Bordag, Kirsten, Leseduarte, Vassilievich,...]

- Zeta functions
- Determinants
- Seeley [heat-kernel] coefficients



# The Chowla-Selberg Formula (CS)

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## ON EPSTEIN'S ZETA FUNCTION (I)

BY S. CHOWLA AND A. SELBERG

INSTITUTE FOR ADVANCED STUDY, PRINCETON, N. J.

Communicated by H. Weyl, May 18, 1949

1. This paper contains a short account of results whose detailed proofs  
will be published later.

We define the function  $Z(s)$  by

$$Z(s) = \sum' (am^2 + bmn + cn^2)^{-s} \quad (1)$$

where  $s = \sigma + it$  ( $\sigma$  and  $t$ , real),  $\sigma > 1$ , and the summation is for all integers  $m, n$  (each going from  $-\infty$  to  $+\infty$ ), while the dash indicates that  $m = n = 0$  is excluded from the summation; further  $a$  and  $c$  are positive numbers while  $b$  is real and subject to  $4ac - b^2 = \Delta > 0$ .

It is well known that the function  $Z(s)$ , defined for  $\sigma > 1$  by (1), can be continued analytically over the whole  $s$ -plane, and satisfies a functional equation similar to the one satisfied by the Riemann Zeta Function. The function  $Z(s)$ , thus defined, is a meromorphic function with a simple pole at  $s = 1$ .

Deuring (*Math. Ztschr.*, 37, 403–413 (1933)) obtained an important formula for  $Z(s)$ . Deuring's work led Heilbronn (*Quart. J. Maths., Oxford*, 5, 150 (1934)) to the proof of the following famous conjecture of Gauss on the class-number of binary quadratic forms with a negative fundamental discriminant: let  $h(-\Delta)$  denote the number of classes of binary quadratic forms of negative fundamental discriminant  $-\Delta = b^2 - 4ac$ , then

$$h(-\Delta) \rightarrow \infty \quad \text{as} \quad \Delta \rightarrow \infty \quad (2)$$

Again using the ideas of Heilbronn and Deuring, Siegel proved that

$$h(-\Delta) > \Delta^{1/2 - \epsilon} \quad [\Delta > \Delta_0(\epsilon)] \quad (3)$$

which is a great advance on (2).

Our starting point is the formula:

$$Z(s) = 2\zeta(2s)a^{-s} + \frac{2^{2s}a^{s-1}\sqrt{\pi}}{\Gamma(s)\Delta^{s-1/2}}\zeta(2s-1)\Gamma(s-1/2) + Q(s) \quad (4)$$

where

$$Q(s) = \frac{\pi^s \cdot 2^s + \frac{s}{2}}{a^{1/2}\Gamma(s)\Delta^{s/2-1/4}} \sum_{n=1}^{\infty} n^{s-1/2} \sigma_{1-2s}(n) \cos\left(\frac{n\pi b}{a}\right) \int_0^{\infty} \phi^{s-1/2} \exp\left\{-\frac{\pi n \Delta^{1/2}}{2a} (\phi + \phi^{-1})\right\} d\phi \quad (4)$$

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# On Epstein's Zeta-function

By Ale Selberg at Princeton (N. J.), and S. Chowla at State College (Pa.)

## Introduction

This paper was written in the Spring of 1949, and a resumé appeared in the note: On Epstein's zeta Function (I), Proceedings of the National Academy of Sciences (U. S. A.), 35 (1949), 371--374.

Meanwhile, the following papers which have reference to the Proceedings paper, came to our attention:

1. *J. B. Rosser*, Real roots of real Dirichlet  $L$ -series, Jour. Research National Bureau of Standards, 45 (1950), 505--514.

2. *E. A. Anferteva*, On an identity of Chowla and Selberg (Russian), Izvestija Vyssik Učebnyh Zavadenii Mathematika (Kazan), No. 3 (10) (1959), 13--21.

3. *P. T. Bateman* and *E. Grosswald*, On Epstein's zeta Function, Acta Arithmetica, 9 (1964), 365--373.

4. *K. Ramachandra*, Some applications of Kronecker's limit formulas, Annals of Mathematics 80 (1964), 104--148.

## § 1.

We define the function  $Z(s)$  by

$$(1) \quad Z(s) = \Sigma' (am^2 + bmn + cn^2)^{-s}$$

where  $s = \sigma + it$  ( $\sigma$  and  $t$ , real),  $\sigma > 1$ , and the summation is for all integers  $m, n$  (each going from  $-\infty$  to  $+\infty$ ), while the dash indicates that  $m = n = 0$  is excluded from the summation; further  $a$  and  $c$  are positive numbers while  $b$  is real and subject to  $4ac - b^2 = \Delta > 0$ .

It is well-known that the function  $Z(s)$ , defined for  $\sigma > 1$  by (1), can be continued analytically over the whole  $s$ -plane. The function  $Z(s)$ , thus defined, is a meromorphic function with a simple pole at  $s = 1$ .

In 1933, Deuring obtained an important formula for  $Z(s)$ . Deuring's work led Heilbronn to his proof of a famous conjecture of Gauss on the class number of binary quadratic forms with a negative fundamental discriminant. If  $h(-\Delta)$  is the number of classes of binary quadratic forms of negative fundamental discriminant  $-\Delta = b^2 - 4ac$ , Gauss conjectured that

$$(2) \quad h(-\Delta) \rightarrow \infty \text{ as } \Delta \rightarrow \infty.$$

Transforming this we get

$$\sum_{j=1}^h \log \Delta \left( \frac{b_j + i\sqrt{|d|}}{2a_j} \right) = 6 \left\{ h\gamma + \log \prod_{j=1}^h \frac{a_j}{|d|^{\frac{j}{h}}} \right\} - \frac{3w}{\pi} \sqrt{|d|} L'_d(1).$$

Inserting here the value (obtained like (58))

$$L'_d(1) = -\frac{\pi}{\sqrt{|d|}} \sum_{m=1}^{|d|} \left( \frac{d}{m} \right) \log \Gamma \left( \frac{m}{|d|} \right) + \frac{2h\pi(\gamma + \log 2\pi)}{w\sqrt{|d|}}$$

one gets, writing  $\tau_j = \frac{b_j + i\sqrt{|d|}}{2a_j}$ ,

$$(2) \quad \prod_{j=1}^h \Delta(\tau_j) = \frac{\prod_{j=1}^h a_j^6}{(2\pi |d|)^{6h}} \left\{ \prod_{m=1}^{|d|} \Gamma \left( \frac{m}{|d|} \right) \left( \frac{d}{m} \right)^{\frac{3w}{h}} \right\}$$

Now let  $\tau = i \frac{K'}{K}$  be a number from the field  $k(\sqrt{d})$ , then from Lemma 3 we get

$$\frac{\Delta(\tau_j)}{\Delta(\tau)} = \lambda_j,$$

where  $\lambda_j$  are algebraic numbers. Thus (2) gives

$$(3) \quad \Delta(\tau) = \frac{\lambda'}{\pi^6} \left\{ \prod_{m=1}^{|d|} \Gamma \left( \frac{m}{|d|} \right) \left( \frac{d}{m} \right)^{\frac{3w}{h}} \right\},$$

where  $\lambda'$  is an algebraic number. Finally we have from (48)

$$\Delta(\tau) = \left( \frac{2K}{\pi} \right)^{12} \cdot 2^{-s} (kk')^4 = \lambda'' \left( \frac{K}{\pi} \right)^{12},$$

where  $\lambda''$  is an algebraic number. This gives, when inserted in (3)

$$(4) \quad K = \lambda''' \sqrt{\pi} \left\{ \prod_{m=1}^{|d|} \Gamma \left( \frac{m}{|d|} \right) \left( \frac{d}{m} \right)^{\frac{w}{4h}} \right\},$$

which is the desired expression for  $K$  in finite terms.

### References (in the order of appearance in the text)

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- [2] S. Chowla, Acta Arithmetica 1 (1935), 113—114.
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- P. Deligne, *Valeurs de fonctions L et periodes d'integrales*, PSPM 33 (1979) 313-346

# History

- Lerch (1897):

$$\sum_{\lambda=1}^{|D|} \left( \frac{D}{\lambda} \right) \log \Gamma \left( \frac{\lambda}{D} \right) = h \log |D| - \frac{h}{3} \log(2\pi) - \sum_{(a,b,c)} \log a + \frac{2}{3} \sum_{(a,b,c)} \log [\theta'_1(0|\alpha)\theta'_1(0|\beta)]$$

$D$  discriminant,  $\theta'_1 \sim \eta^3$

$h$  class number of binary quadratic forms  $(a, b, c)$

# History

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$$\sum_{\lambda=1}^{|D|} \left( \frac{D}{\lambda} \right) \log \Gamma \left( \frac{\lambda}{D} \right) = h \log |D| - \frac{h}{3} \log(2\pi) - \sum_{(a,b,c)} \log a$$

$$+ \frac{2}{3} \sum_{(a,b,c)} \log [\theta'_1(0|\alpha)\theta'_1(0|\beta)]$$

$D$  discriminant,  $\theta'_1 \sim \eta^3$

$h$  class number of binary quadratic forms  $(a, b, c)$

- Eta evaluations Dedekind eta function for  $\operatorname{Im}(\tau) > 0$

$$\eta(\tau) = q^{1/24} \prod_{n=1}^{\infty} (1 - q^n), \quad q := e^{2\pi i \tau}$$

It is a 24-th root of the discriminant func  $\Delta(\tau)$  of an elliptic curve  $\mathbb{C}/L$  from a lattice  $L = \{a\tau + b \mid a, b \in \mathbb{Z}\}$

$$\Delta(\tau) = (2\pi)^{12} q \prod_{n=1}^{\infty} (1 - q^n)^{24}$$

# Properties & Recent Results

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  - T. Yang, The Chowla-Selberg formula and the Colmez conjecture, Canad. J. Math. 62 (2010), pp. 456-472

# Extended CS Series Formulas (ECS)

- Consider the zeta function ( $\operatorname{Re}s > p/2$ ,  $A > 0$ ,  $\operatorname{Re}q > 0$ )

$$\zeta_{A,\vec{c},q}(s) = \sum'_{\vec{n} \in \mathbb{Z}^p} \left[ \frac{1}{2} (\vec{n} + \vec{c})^T A (\vec{n} + \vec{c}) + q \right]^{-s} = \sum'_{\vec{n} \in \mathbb{Z}^p} [Q(\vec{n} + \vec{c}) + q]^{-s}$$

**prime:** point  $\vec{n} = \vec{0}$  to be excluded from the sum  
(inescapable condition when  $c_1 = \dots = c_p = q = 0$ )

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EE JPA34 (2001) 3025

[ECS1]

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EE JPA34 (2001) 3025 [ECS1]

- Pole:  $s = p/2$  Residue:

$$\operatorname{Res}_{s=p/2} \zeta_{A,\vec{c},q}(s) = \frac{(2\pi)^{p/2}}{\Gamma(p/2)} (\det A)^{-1/2}$$

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- **Case**  $c_1 = \dots = c_p = q = 0$  [true extens of CS, diag subcase]

$$\zeta_{A_p}(s) = \frac{2^{1+s}}{\Gamma(s)} \sum_{j=0}^{p-1} (\det A_j)^{-1/2} \left[ \pi^{j/2} a_{p-j}^{j/2-s} \Gamma\left(s - \frac{j}{2}\right) \zeta_R(2s-j) + \right.$$

$$\left. 4\pi^s a_{p-j}^{\frac{j}{4} - \frac{s}{2}} \sum_{n=1}^{\infty} \sum_{\vec{m}_j \in \mathbb{Z}^j} n^{j/2-s} (\vec{m}_j^t A_j^{-1} \vec{m}_j)^{s/2-j/4} K_{j/2-s} \left( 2\pi n \sqrt{a_{p-j} \vec{m}_j^t A_j^{-1} \vec{m}_j} \right) \right]$$

[ECS3d]

# QFT in s-t with non-comm toroidal part

- $D$ -dim non-commut manifold:  $M = \mathbb{R}^{1,d} \otimes \mathbb{T}_\theta^p$ ,  $D = d + p + 1$   
 $\mathbb{T}_\theta^p$  a  $p$ -dim non-commutative torus:  $[x_j, x_k] = i\theta\sigma_{jk}$   
 $\sigma_{jk}$  a real, nonsingular, antisymmetric matrix of  $\pm 1$  entries  
 $\theta$  the non-commutative parameter. [Quresh's talk]

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- Interest recently, in connection with  $M$ -theory & string theory  
[Connes, Douglas, Seiberg, Cheung, Chu, Chomerus, Ardalan, ...]
- Unified treatment: only one zeta function, nature of field  
(bosonic, fermionic) as a parameter, together with # of  
compact, noncompact, and noncommutative dimensions

$$\zeta_\alpha(s) = \frac{V \Gamma(s - (d+1)/2)}{(4\pi)^{(d+1)/2} \Gamma(s)} \sum_{\vec{n} \in \mathbb{Z}^p} Q(\vec{n})^{(d+1)/2-s} [1 + \Lambda \theta^{2-2\alpha} Q(\vec{n})^{-\alpha}]^{(d+1)/2-s}$$

$$\alpha = 2 \text{ bos}, \quad \alpha = 3 \text{ ferm}, \quad V = \text{Vol}(\mathbb{R}^{d+1}) \text{ of non-compact part}$$

$$Q(\vec{n}) = \sum_{j=1}^p a_j n_j^2 \text{ a diag quadratic form, } R_j = a_j^{-1/2} \text{ compactific radii}$$

$$\begin{aligned}
\zeta_\alpha(s) = & \frac{2^{s-d} V}{(2\pi)^{(d+1)/2} \Gamma(s)} \sum_{l=0}^{\infty} \frac{\Gamma(s+l-(d+1)/2)}{l! \Gamma(s+\alpha l - (d+1)/2)} (-2^\alpha \Lambda \theta^{2-2\alpha})^l \sum_{j=0}^{p-1} (\det A_j)^{-\frac{1}{2}} \\
& \times \left[ \pi^{j/2} a_{p-j}^{-s-\alpha l + (d+j+1)/2} \Gamma(s+\alpha l - (d+j+1)/2) \zeta_R(2s+2\alpha l - d - j - 1) \right. \\
& + 4\pi^{s+\alpha l -(d+1)/2} a_{p-j}^{-(s+\alpha l)/2 - (d+j+1)/4} \sum_{n=1}^{\infty} \sum'_{\vec{m}_j \in \mathbb{Z}^j} n^{(d+j+1)/2 - s - \alpha l} \\
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\end{aligned}$$

$p \setminus D$	even	odd
odd	(1a) pole / finite ( $l \geq l_1$ )	(2a) pole / pole
even	(1b) double pole / pole ( $l \geq l_1, l_2$ )	(2b) pole / double pole ( $l \geq l_2$ )

- General pole structure of  $\zeta_\alpha(s)$ , for the possible values of  $D$  and  $p$  being odd or even. Magenta, type of behavior corresponding to lower values of  $l$ ; behavior in blue corresponds to larger values of  $l$

After some calculations,

$$\zeta_\alpha(s) = \frac{V}{(4\pi)^{(d+1)/2}} \sum_{l=0}^{\infty} \frac{\Gamma(s + l - \frac{d+1}{2})}{l! \Gamma(s)} (-\Lambda \theta^{2-2\alpha})^l \zeta_{Q,\vec{0},0}(s + \alpha l - \frac{d+1}{2})$$

for all radii equal to  $R$ , with  $I(\vec{n}) = \sum_{j=1}^p n_j^2$ ,

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where we use the notation  $\zeta_E(s) := \zeta_{I,\vec{0},0}(s)$

e.g., the Epstein zeta function for the standard quadratic form

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- Rich pole structure:** pole of Epstein zf at  
 $s = p/2 - \alpha k + (d+1)/2 = D/2 - \alpha k$ , combined with  
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- Classify** the different possible cases according to the values of  $d$  and  $D = d + p + 1$ . We obtain, at  $s = 0$ :

For  $d = 2k$

$$\begin{cases} \text{if } D \neq \overline{2\alpha} & \implies \zeta_\alpha(0) = 0 \\ \text{if } D = \overline{2\alpha} & \implies \zeta_\alpha(0) = \text{finite} \end{cases}$$

For  $d = 2k - 1$

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- Pole structure of the zeta function  $\zeta_\alpha(s)$ , at  $s = 0$ , according to the different possible values of  $d$  and  $D$  ( $\overline{2\alpha}$  means multiple of  $2\alpha$ )

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- Pole structure of the zeta function  $\zeta_\alpha(s)$ , at  $s = 0$ , according to the different possible values of  $d$  and  $D$  ( $\overline{2\alpha}$  means multiple of  $2\alpha$ )
- ⇒ Explicit analytic continuation of  $\zeta_\alpha(s)$ ,  $\alpha = 2, 3$ , & specific pole structure

# Pseudodifferential Operator ( $\Psi$ DO)

- $A$   $\Psi$ DO of order  $m$        $M_n$  manifold
- Symbol of  $A$ :  $a(x, \xi) \in S^m(\mathbb{R}^n \times \mathbb{R}^n) \subset C^\infty$  functions such that for any pair of multi-indices  $\alpha, \beta$  there exists a constant  $C_{\alpha, \beta}$  so that

$$\left| \partial_\xi^\alpha \partial_x^\beta a(x, \xi) \right| \leq C_{\alpha, \beta} (1 + |\xi|)^{m - |\alpha|}$$

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Definition of  $A$  (in the distribution sense)

$$Af(x) = (2\pi)^{-n} \int e^{i \langle x, \xi \rangle} a(x, \xi) \hat{f}(\xi) d\xi$$

- $f$  is a smooth function  
 $f \in \mathcal{S} = \{f \in C^\infty(\mathbb{R}^n); \sup_x |x^\beta \partial^\alpha f(x)| < \infty, \forall \alpha, \beta \in \mathbb{N}^n\}$
- $\mathcal{S}'$  space of tempered distributions
- $\hat{f}$  is the Fourier transform of  $f$

# $\Psi$ DOs are useful tools

The symbol of a  $\Psi$ DO has the form:

$$a(x, \xi) = a_m(x, \xi) + a_{m-1}(x, \xi) + \cdots + a_{m-j}(x, \xi) + \cdots$$

being  $a_k(x, \xi) = b_k(x) \xi^k$

$a(x, \xi)$  is said to be elliptic if it is invertible for large  $|\xi|$  and if there exists a constant  $C$  such that  $|a(x, \xi)^{-1}| \leq C(1 + |\xi|)^{-m}$ , for  $|\xi| \geq C$

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— —  $\Psi$ DOs are basic tools both in Mathematics & in Physics — —

1. Proof of uniqueness of Cauchy problem [Calderón-Zygmund]
2. Proof of the Atiyah-Singer index formula
3. In QFT they appear in any analytical continuation process —as complex powers of differential operators, like the Laplacian [Seeley, Gilkey, ...]
4. Basic starting point of any rigorous formulation of QFT & gravitational interactions through  $\mu$ localization (the most important step towards the understanding of linear PDEs since the invention of distributions)

[K Fredenhagen, R Brunetti, ... R Wald '06, R Haag EPJH35 '10]

# Existence of $\zeta_A$ for $A$ a $\Psi$ DO

1.  $A$  a positive-definite elliptic  $\Psi$ DO of positive order  $m \in \mathbb{R}^+$
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(b)  $\zeta_A(s)$  has a meromorphic continuation to the whole complex plane  $\mathbb{C}$  (regular at  $s = 0$ ), provided the principal symbol of  $A$ ,  $a_m(x, \xi)$ , admits a spectral cut:  $L_\theta = \{\lambda \in \mathbb{C}; \operatorname{Arg} \lambda = \theta, \theta_1 < \theta < \theta_2\}$ ,  $\operatorname{Spec} A \cap L_\theta = \emptyset$  (the Agmon-Nirenberg condition)

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(c) The definition of  $\zeta_A(s)$  depends on the position of the cut  $L_\theta$

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(a) The zeta function is defined as:

$$\zeta_A(s) = \operatorname{tr} A^{-s} = \sum_j \lambda_j^{-s}, \quad \operatorname{Re} s > \frac{n}{m} := s_0$$

$\{\lambda_j\}$  ordered spect of  $A$ ,  $s_0 = \dim M/\operatorname{ord} A$  abscissa of converg of  $\zeta_A(s)$

(b)  $\zeta_A(s)$  has a meromorphic continuation to the whole complex plane  $\mathbb{C}$  (regular at  $s = 0$ ), provided the principal symbol of  $A$ ,  $a_m(x, \xi)$ , admits a spectral cut:  $L_\theta = \{\lambda \in \mathbb{C}; \operatorname{Arg} \lambda = \theta, \theta_1 < \theta < \theta_2\}$ ,  $\operatorname{Spec} A \cap L_\theta = \emptyset$  (the Agmon-Nirenberg condition)

(c) The definition of  $\zeta_A(s)$  depends on the position of the cut  $L_\theta$

(d) The only possible singularities of  $\zeta_A(s)$  are poles at

$$s_j = (n - j)/m, \quad j = 0, 1, 2, \dots, n - 1, n + 1, \dots$$

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C. Soulé et al, Lectures on Arakelov Geometry, CUP 1992; A. Voros, ...

# Properties

- The definition of the determinant  $\det_{\zeta} A$  only depends on the homotopy class of the cut
- A zeta function (and corresponding determinant) with the same meromorphic structure in the complex  $s$ -plane and extending the ordinary definition to operators of complex order  $m \in \mathbb{C} \setminus \mathbb{Z}$  (they do not admit spectral cuts), has been obtained [Kontsevich and Vishik]
- Asymptotic expansion for the heat kernel:

$$\text{tr } e^{-tA} = \sum'_{\lambda \in \text{Spec } A} e^{-t\lambda}$$

$$\sim \alpha_n(A) + \sum_{n \neq j \geq 0} \alpha_j(A) t^{-s_j} + \sum_{k \geq 1} \beta_k(A) t^k \ln t, \quad t \downarrow 0$$

$$\alpha_n(A) = \zeta_A(0), \quad \alpha_j(A) = \Gamma(s_j) \text{Res}_{s=s_j} \zeta_A(s), \quad s_j \notin -\mathbb{N}$$

$$\alpha_j(A) = \frac{(-1)^k}{k!} [\text{PP } \zeta_A(-k) + \psi(k+1) \text{Res}_{s=-k} \zeta_A(s)],$$

$$\beta_k(A) = \frac{(-1)^{k+1}}{k!} \text{Res}_{s=-k} \zeta_A(s), \quad k \in \mathbb{N} \setminus \{0\} \quad s_j = -k, \quad k \in \mathbb{N}$$

$$\text{PP } \phi := \lim_{s \rightarrow p} \left[ \phi(s) - \frac{\text{Res}_{s=p} \phi(s)}{s-p} \right]$$

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$$\text{Dtr } T = \lim_{N \rightarrow \infty} \frac{1}{\log N} \sigma_N(T)$$

provided that the Cesaro means  $M(\sigma)(N)$  of the sequence in  $N$  are convergent as  $N \rightarrow \infty$  [remember:  $M(f)(\lambda) = \frac{1}{\ln \lambda} \int_1^\lambda f(u) \frac{du}{u}$  ]

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- The **Hardy-Littlewood theorem** can be stated in a way that connects the Dixmier trace with the residue of the zeta function of the operator  $T^{-1}$  at  $s = 1$  [Connes] 
$$\text{Dtr } T = \lim_{s \rightarrow 1^+} (s - 1) \zeta_{T^{-1}}(s)$$

# The Wodzicki Residue

- The Wodzicki (or noncommutative) residue is the only extension of the Dixmier trace to  $\Psi$ DOs which are not in  $\mathcal{L}^{(1,\infty)}$
- Only trace one can define in the algebra of  $\Psi$ DOs (up to multipl const)
- Definition:  $\text{res } A = 2 \text{Res}_{s=0} \text{tr}(A\Delta^{-s})$ ,  $\Delta$  Laplacian
- Satisfies the trace condition:  $\text{res } (AB) = \text{res } (BA)$
- Important!: it can be expressed as an integral (local form)

$$\text{res } A = \int_{S^* M} \text{tr } a_{-n}(x, \xi) d\xi$$

with  $S^* M \subset T^* M$  the co-sphere bundle on  $M$  (some authors put a coefficient in front of the integral: Adler-Manin residue)

- If  $\dim M = n = -\text{ord } A$  ( $M$  compact Riemann,  $A$  elliptic,  $n \in \mathbb{N}$ ) it coincides with the Dixmier trace, and  $\text{Res}_{s=1} \zeta_A(s) = \frac{1}{n} \text{res } A^{-1}$
- The Wodzicki residue makes sense for  $\Psi$ DOs of arbitrary order. Even if the symbols  $a_j(x, \xi)$ ,  $j < m$ , are not coordinate invariant, the integral is, and defines a trace

# Singularities of $\zeta_A$

- A complete determination of the meromorphic structure of some zeta functions in the complex plane can be also obtained by means of the Dixmier trace and the Wodzicki residue
- Missing for full descript of the singularities: **residua** of all poles
- As for the regular part of the analytic continuation: specific methods have to be used (see later)
- **Proposition.** Under the conditions of existence of the zeta function of  $A$ , given above, and being the symbol  $a(x, \xi)$  of the operator  $A$  analytic in  $\xi^{-1}$  at  $\xi^{-1} = 0$ :

$$\text{Res}_{s=s_k} \zeta_A(s) = \frac{1}{m} \text{res } A^{-s_k} = \frac{1}{m} \int_{S^* M} \text{tr } a_{-n}^{-s_k}(x, \xi) d^{n-1}\xi$$

- **Proof.** The homog component of degree  $-n$  of the corresp power of the principal symbol of  $A$  is obtained by the appropriate derivative of a power of the symbol with respect to  $\xi^{-1}$  at  $\xi^{-1} = 0$  :

$$a_{-n}^{-s_k}(x, \xi) = \left( \frac{\partial}{\partial \xi^{-1}} \right)^k \left[ \xi^{n-k} a^{(k-n)/m}(x, \xi) \right] \Big|_{\xi^{-1}=0} \xi^{-n}$$

# Multipl or N-Comm Anomaly, or Defect

- Given  $A$ ,  $B$ , and  $AB$   $\psi$ DOs, even if  $\zeta_A$ ,  $\zeta_B$ , and  $\zeta_{AB}$  exist, it turns out that, in general,

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- The multiplicative (or noncommutative) anomaly (defect) is defined as

$$\delta(A, B) = \ln \left[ \frac{\det_\zeta(AB)}{\det_\zeta A \det_\zeta B} \right] = -\zeta'_{AB}(0) + \zeta'_A(0) + \zeta'_B(0)$$

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- Wodzicki formula

$$\delta(A, B) = \frac{\text{res} \{ [\ln \sigma(A, B)]^2 \}}{2 \text{ ord } A \text{ ord } B (\text{ord } A + \text{ord } B)}$$

where  $\sigma(A, B) = A^{\text{ord } B} B^{-\text{ord } A}$

# Consequences of the Multipl Anomaly

- In the path integral formulation

$$\int [d\Phi] \exp \left\{ - \int d^D x \left[ \Phi^\dagger(x) (\quad) \Phi(x) + \dots \right] \right\}$$

Gaussian integration:  $\longrightarrow \det(\quad)^\pm$

$$\begin{pmatrix} A_1 & A_2 \\ A_3 & A_4 \end{pmatrix} \longrightarrow \begin{pmatrix} A & \\ & B \end{pmatrix}$$

$$\det(AB) \quad \text{or} \quad \det A \cdot \det B \quad ?$$

- In a situation where a superselection rule exists,  $AB$  has no sense (much less its determinant):  $\Rightarrow \det A \cdot \det B$
- But if diagonal form obtained after change of basis (diag. process), the preserved quantity is:  $\Rightarrow \det(AB)$

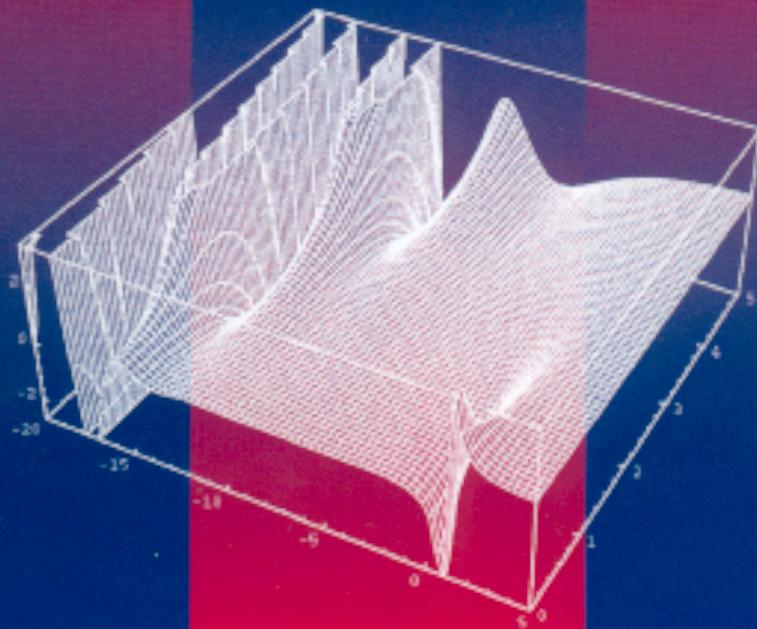
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