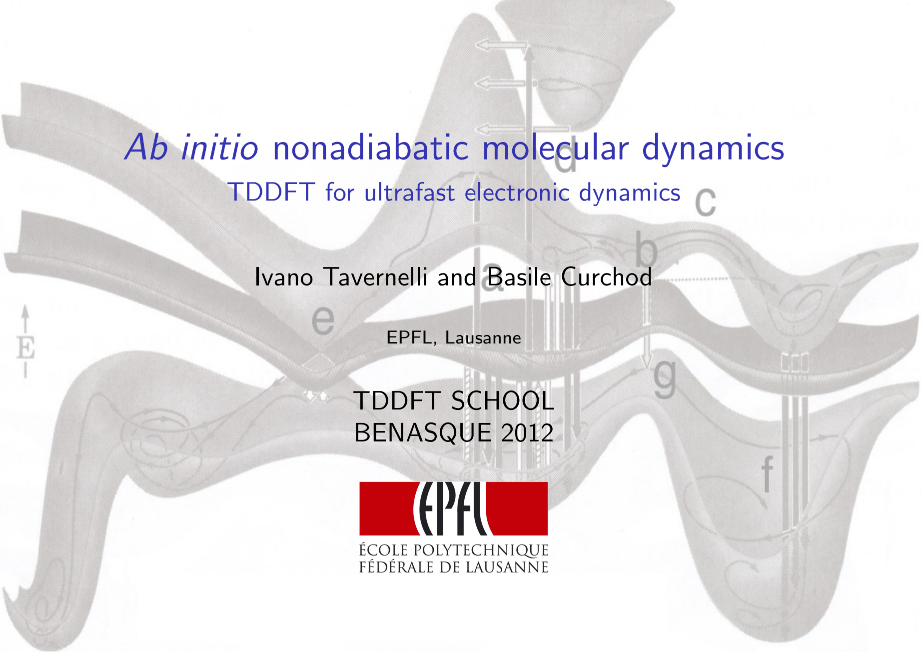

# *Ab initio* nonadiabatic molecular dynamics

## TDDFT for ultrafast electronic dynamics

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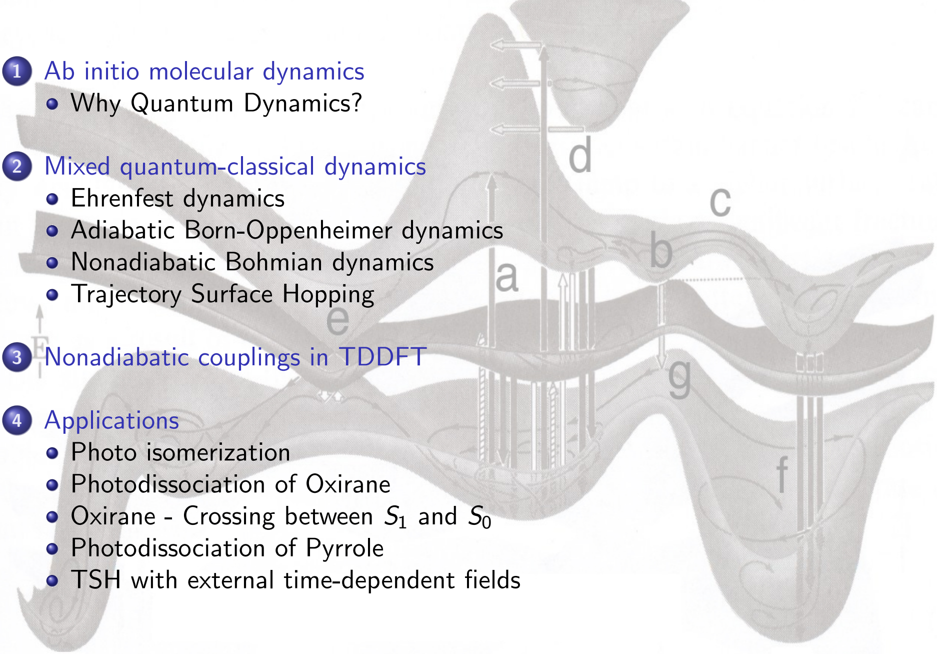
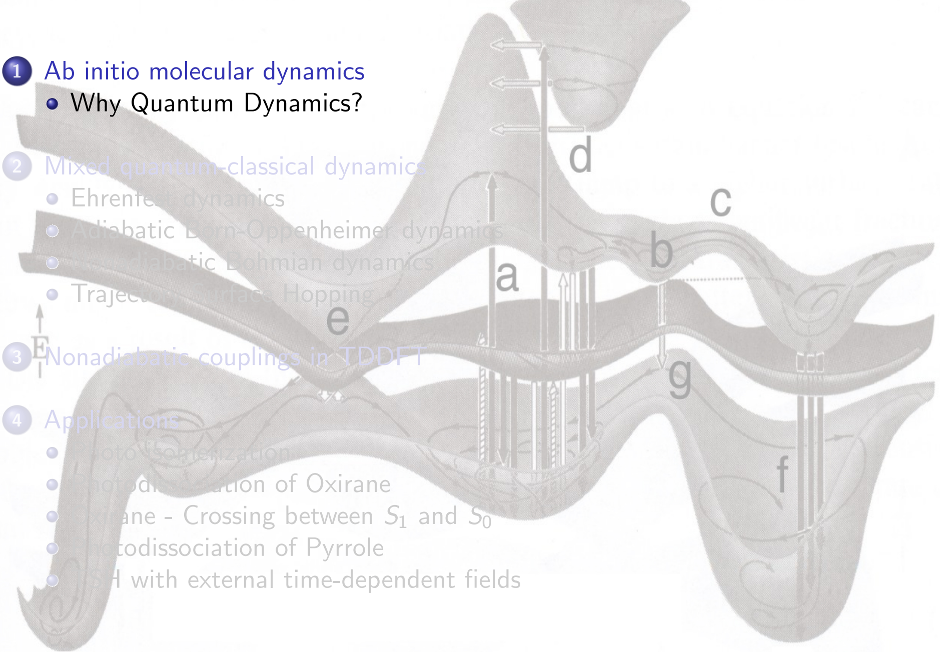
EPFL, Lausanne

TDDFT SCHOOL
BENASQUE 2012

ÉCOLE POLYTECHNIQUE
FÉDÉRALE DE LAUSANNE

1. Ab initio molecular dynamics
    - Why Quantum Dynamics?
2. Mixed quantum-classical dynamics
    - Ehrenfest dynamics
    - Adiabatic Born-Oppenheimer dynamics
    - Nonadiabatic Bohmian dynamics
    - Trajectory Surface Hopping
3. Nonadiabatic couplings in TDDFT
4. Applications
    - Photo isomerization
    - Photodissociation of Oxirane
    - Oxirane - Crossing between $S_1$ and $S_0$
    - Photodissociation of Pyrrole
    - TSH with external time-dependent fields





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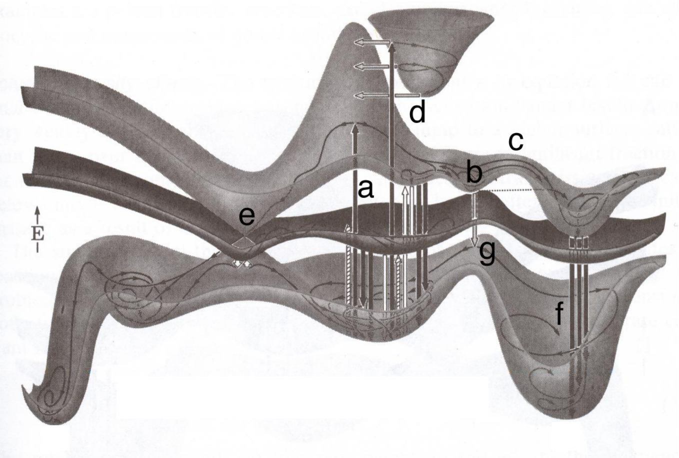
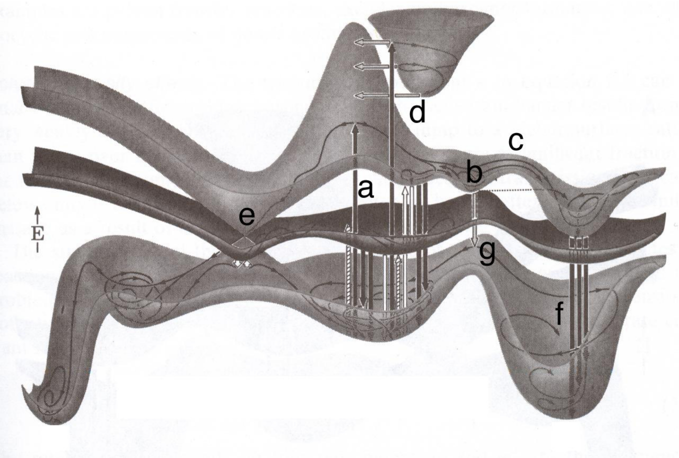
# Reminder from last lecture: potential energy surfaces



We have electronic structure methods for electronic ground and excited states...

Now, we need to propagate the nuclei...

# Reminder from last lecture: potential energy surfaces



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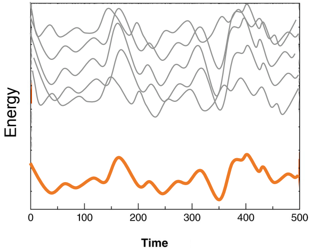
Now, we need to propagate the nuclei...

# Why Quantum dynamics?

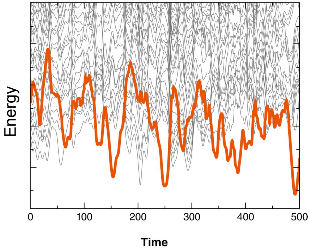


GS adiabatic dynamics (BO vs. CP)

BO $M_I \ddot{\mathbf{R}}_I(t) = -\nabla \min_\rho E_{KS}(\{\phi_i[\rho]\})$

CP $\mu_i |\ddot{\phi}_i(t)\rangle = -\frac{\delta}{\delta\langle\phi_i|} E_{KS}(\{\phi_i(\mathbf{r})\}) + \frac{\delta}{\delta\langle\phi_i|}\{\text{constr.}\}$

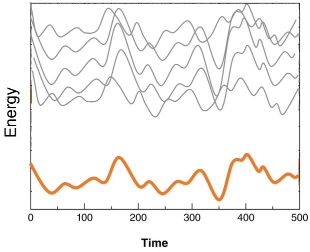
$M_I \ddot{\mathbf{R}}_I(t) = -\nabla E_{KS}(\{\phi_i(t)\})$



ES nonadiabatic quantum dynamics

- Wavepacket dynamics (MCTDH)
- Trajectory-based approaches
  - Tully's trajectory surface hopping (TSH)
  - Bohmian dynamics (quantum hydrodyn.)
  - Semiclassical (WKB, DR)
  - Path integrals (Pechukas)
  - Mean-field solution (Ehrenfest dynamics)
- Density matrix, Liouvillian approaches, ...

# Why Quantum dynamics?

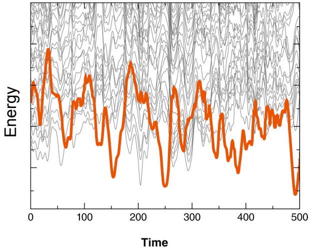


## GS adiabatic dynamics

First principles Heaven

Ab initio MD with WF methods
Ab initio MD with DFT & TDDFT [CP]
classical MD
Coarse-grained MD
...

No principles World (Mike Klein, Sept. 2011)



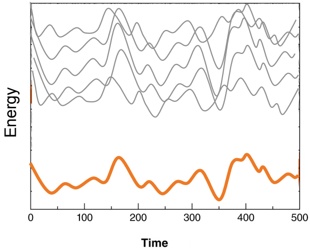
## ES nonadiabatic quantum dynamics

First principles Heaven

Ab initio MD with WF methods
Ab initio MD with DFT & TDDFT [CP]
↓
Models
↓?

No principles World

# Why Quantum dynamics?



## GS adiabatic dynamics

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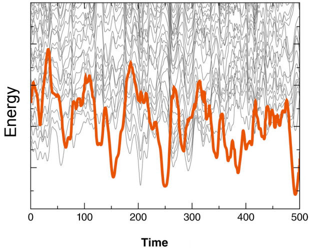
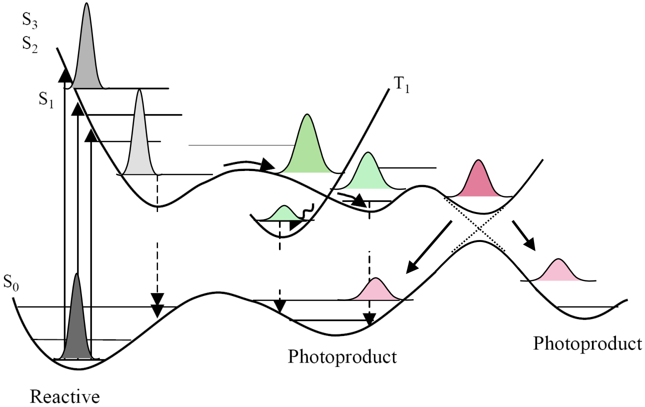


## ES nonadiabatic quantum dynamics

- (-) We cannot get read of electrons
- (-) Nuclei keep some QM flavor
- (-) Accuracy is an issue
- (-) Size can be large (diffuse excitons)
- (+) Time scales are usually short ($<$ *ps*)

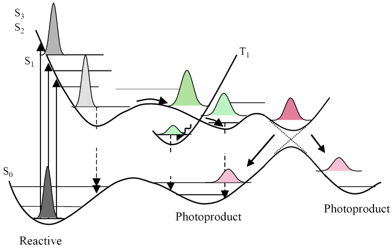
# Nonadiabatic effects requires quantum nuclear dynamics

The nuclear dynamics cannot be described by a single classical trajectory (like in the ground state -adiabatically separated- case)



# Why trajectory-based approaches?

- W1 In "conventional" nuclear wavepacket propagation potential energy surfaces are needed.
- W2 Difficulty to obtain and fit potential energy surfaces for large molecules.
- W3 Nuclear wavepacket dynamics is very expensive for large systems (6 degrees of freedom, 30 for MCTDH). Bad scaling.
- T1 Trajectory based approaches can be run *on-the-fly* (no need to parametrize potential energy surfaces).
- T2 Can handle large molecules in the full (unconstraint) configuration space.
- T3 They offer a good **compromise** between accuracy and computational effort.



# Starting point

The starting point is the molecular time-dependent Schrödinger equation :

$$\hat{H}\Psi(\mathbf{r},\mathbf{R},t) = i\hbar\frac{\partial}{\partial t}\Psi(\mathbf{r},\mathbf{R},t)$$

where $\hat{H}$ is the molecular time-independent Hamiltonian and $\Psi(\mathbf{r},\mathbf{R},t)$ the total wavefunction (nuclear + electronic) of our system.

In mixed quantum-classical dynamics the nuclear dynamics is described by a swarm of classical trajectories (taking a "partial" limit $\hbar \to 0$ for the nuclear wf).

In this lecture we will discuss two main approximate solutions based on the following Ansätze for the total wavefucntion

$$\Psi(\mathbf{r},\mathbf{R},t) \xrightarrow[\text{Huang}]{\text{Born-}} \sum_j^{\infty} \Phi_j(\mathbf{r};\mathbf{R})\Omega_j(\mathbf{R},t)$$

$$\Psi(\mathbf{r},\mathbf{R},t) \xrightarrow{\text{Ehrenfest}} \Phi(\mathbf{r},t)\Omega(\mathbf{R},t)\exp\left[\frac{i}{\hbar}\int_{t_0}^{t} E_{el}(t')dt'\right]$$

# Tarjectory-based quantum and mixed QM-CL solutions

We can “derive” the following trajectory-based solutions:

- Nonadiabatic Ehrenfest dynamics dynamics

  I. Tavernelli et al., *Mol. Phys.*, **103**, 963981 (2005).

- Adiabatic Born-Oppenheimer MD equations

- Nonadiabatic Bohmian Dynamics (NABDY)

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- Nonadiabatic Trajectory Surface Hopping (TSH) dynamics
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- Time dependent potential energy surface approach based on the exact decomposition: $\Psi(\mathbf{r}, \mathbf{R}, t) = \Omega(\mathbf{R}, t)\Phi(\mathbf{r}, t)$.

  A. Abedi, N. T. Maitra, E. K. U. Gross, *PRL*, **105**, 123002 (2010)

# Ehrenfest dynamics

$$\Psi(\mathbf{r},\mathbf{R},t) \xrightarrow{\text{Ehrenfest}} \Phi(\mathbf{r},t)\Omega(\mathbf{R},t)\exp\left[\frac{i}{\hbar}\int_{t_0}^{t} E_{el}(t')dt'\right]$$

Inserting this representation of the total wavefunction into the molecular td Schrödinger equation and multiplying from the left-hand side by $\Omega^*(\mathbf{R},t)$ and integrating over $\mathbf{R}$ we get

$$i\hbar\frac{\partial\Phi(\mathbf{r},t)}{\partial t} = -\frac{\hbar^2}{2m_e}\sum_i \nabla_i^2\Phi(\mathbf{r},t) + \left[\int d\mathbf{R}\ \Omega^*(\mathbf{R},t)\hat{V}(\mathbf{r},\mathbf{R})\Omega(\mathbf{R},t)\right]\Phi(\mathbf{r},t)$$

where $\hat{V}(\mathbf{r},\mathbf{R}) = \sum_{i<j}\frac{e^2}{|\mathbf{r}_i-\mathbf{r}_j|} - \sum_{\gamma,i}\frac{e^2 Z_\gamma}{|\mathbf{R}_\gamma-\mathbf{r}_i|}$ .
In a similar way, multiplying by $\Phi^*(\mathbf{r},t)$ and integrating over $\mathbf{r}$ we obtain

$$i\hbar\frac{\partial\Omega(\mathbf{R},t)}{\partial t} = -\frac{\hbar^2}{2}\sum_\gamma M_\gamma^{-1}\nabla_\gamma^2\Omega(\mathbf{R},t) + \left[\int d\mathbf{r}\ \Phi^*(\mathbf{r},t)\hat{\mathcal{H}}_{el}\Phi(\mathbf{r},t)\right]\Omega(\mathbf{R},t)$$

Conservation of energy has also to be imposed through the condition that $d\langle\hat{H}\rangle/dt \equiv 0$.

Note that both the electronic and nuclear parts evolve according to an average potential generated by the other component (in square brakets). These average potentials are time-dependent and are responsible for the feedback interaction between the electronic and nuclear components.

# Ehrenfest dynamics - the nuclear equation

We start from the polar representation of the nuclear wavefunction

$$\Omega(\mathbf{R}, t) = A(\mathbf{R}, t) \exp\left[\frac{i}{\hbar} S(\mathbf{R}, t)\right]$$

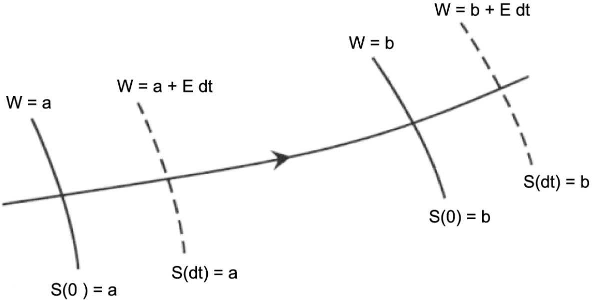
where the amplitude $A(\mathbf{R}, t)$ and the phase $S(\mathbf{R}, t)/\hbar$ are real functions. Inserting this representation for $\Omega(\mathbf{R}, t)$ and separating the real and the imaginary parts one gets for the phase $S$ in the classical limit $\hbar \to 0$

$$\frac{\partial S}{\partial t} = -\frac{1}{2} \sum_{\gamma} M_{\gamma}^{-1} \left(\nabla_{\gamma} S\right)^2 - \left[\int d\mathbf{r}\ \Phi^*(\mathbf{r}, t) \hat{\mathcal{H}}_{el}(\mathbf{r}, \mathbf{R}) \Phi(\mathbf{r}, t)\right]$$

This has the form of the "Hamilton-Jacobi" (HJ) equation of classical mechanics, which establishes a relation between the partial differential equation for $S(\mathbf{R}, t)$ in configuration space and the trajectories of the corresponding (quantum) mechanical systems.

# Ehrenfest dynamics - the nuclear equation

$$\frac{\partial S}{\partial t} = -\frac{1}{2}\sum_{\gamma} M_{\gamma}^{-1}\left(\nabla_{\gamma} S\right)^{2} - \left[\int d\mathbf{r}\ \Phi^{*}(\mathbf{r},t)\hat{\mathcal{H}}_{el}(\mathbf{r},\mathbf{R})\Phi(\mathbf{r},t)\right]$$



Instead of solving the field equation for $S(\mathbf{R},t)$, find the equation of motion for the corresponding trajectories (*characteristics*).

# Ehrenfest dynamics - the nuclear equation

The identification of $S(\mathbf{R}, t)$ with the "classical" action, defines a point-particle dynamics with Hamiltonian, $H_{cl}$ and momenta

$$\mathbf{P} = \nabla_{\mathbf{R}} S(\mathbf{R}).$$

The solutions of this Hamiltonian system iare curves (*characteristics*) in the $(\mathbf{R}, t)$-space, which are extrema of the action $S(\mathbf{R}, t)$ for given initial conditions $\mathbf{R}(t_0)$ and $\mathbf{P}(t_0) = \nabla_{\mathbf{R}} S(\mathbf{R})|_{\mathbf{R}(t_0)}$.
Newton-like equation for the nuclear trajectories corresponding to the HJ equation

$$\frac{d\mathbf{P}_\gamma}{dt} = -\nabla_\gamma \left[ \int d\mathbf{r}\ \Phi^*(\mathbf{r}, t) \hat{\mathcal{H}}_{el}(\mathbf{r}, \mathbf{R}) \Phi(\mathbf{r}, t) \right]$$

## Ehrenfest dynamics

$$i\hbar \frac{\partial \Phi(\mathbf{r}; \mathbf{R}, t)}{\partial t} = \hat{\mathcal{H}}_{el}(\mathbf{r}; \mathbf{R}) \Phi(\mathbf{r}; \mathbf{R}, t)$$

$$M_I \ddot{\mathbf{R}}_I = -\nabla_I \langle \hat{\mathcal{H}}_{el}(\mathbf{r}; \mathbf{R}) \rangle$$

# Ehrenfest dynamics - the nuclear equation

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$$\frac{d\mathbf{P}_\gamma}{dt} = -\nabla_\gamma \left[ \int d\mathbf{r}\ \Phi^*(\mathbf{r},t) \hat{\mathcal{H}}_{el}(\mathbf{r},\mathbf{R}) \Phi(\mathbf{r},t) \right]$$

**Ehrenfest dynamics - Densityfunctionalization ($\phi_k$: KS orbitals)**

$$i\hbar \frac{\partial}{\partial t} \phi_k(\mathbf{r},t) = -\frac{1}{2m_e} \nabla^2_{\mathbf{r}} \phi_k(\mathbf{r},t) + v_{\text{eff}}[\rho,\Phi_0](\mathbf{r},t)\, \phi_k(\mathbf{r},t)$$

$$M_I \ddot{\mathbf{R}}_I = -\nabla_I E[\rho(\mathbf{r},t)]$$

# Ehrenfest dynamics - Example

## Ehrenfest dynamics

$$i\hbar \frac{\partial \Phi(\mathbf{r};\mathbf{R},t)}{\partial t} = \hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi(\mathbf{r};\mathbf{R},t)$$

$$M_I \ddot{\mathbf{R}}_I = -\nabla_I \langle \hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R}) \rangle$$

# Ehrenfest dynamics and mixing of electronic states

## Ehrenfest dynamics

$$i\hbar\frac{\partial\Phi(\mathbf{r};\mathbf{R},t)}{\partial t} = \hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi(\mathbf{r};\mathbf{R},t)$$

$$M_I\ddot{\mathbf{R}}_I = -\nabla_I\langle\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\rangle$$

Consider the following expansion of $\Phi(\mathbf{r};\mathbf{R},t)$ in the *static* basis of electronic wavefucntions $\{\Phi_k(\mathbf{r};\mathbf{R})\}$

$$\Phi(\mathbf{r};\mathbf{R},t) = \sum_{k=0}^{\infty} c_k(t)\Phi_k(\mathbf{r};\mathbf{R})$$

The time-dependency is now on the set of coefficients $\{c_k(t)\}$ ($|c_k(t)|^2$ is the population of state $k$). Inserting in the Ehrenfest's equations...

# Ehrenfest dynamics and mixing of electronic states

## Ehrenfest dynamics

$$i\hbar\dot{c}_k(t) = c_k(t)E_k^{el} - i\hbar\sum_j c_j(t)D_{kj}$$

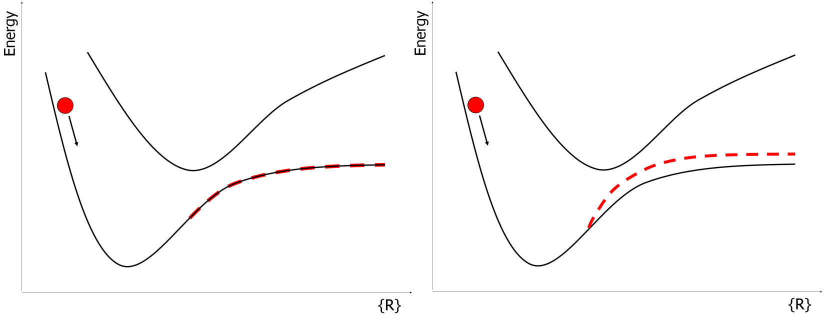
$$M_I\ddot{\mathbf{R}}_I = -\nabla_I\sum_{k=0}^{\infty}|c_k(t)|^2E_k^{el}$$

where

$$D_{kj} = \langle\Phi_k|\frac{\partial}{\partial t}|\Phi_j\rangle = \langle\Phi_k|\frac{\partial\mathbf{R}}{\partial t}\frac{\partial}{\partial\mathbf{R}}|\Phi_j\rangle = \dot{\mathbf{R}}\langle\Phi_k|\nabla|\Phi_j\rangle = \dot{\mathbf{R}}\cdot\mathbf{d}_{kj}$$

Thus we incorporate directly nonadiabatic effects.

# Ehrenfest dynamics: the mean-field potential



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# Born-Oppenheimer approximation

$$\Psi(\mathbf{r},\mathbf{R},t) \xrightarrow[\text{Huang}]{\text{Born-}} \sum_{j}^{\infty} \Phi_j(\mathbf{r};\mathbf{R})\Omega_j(\mathbf{R},t)$$

In this equation, $\{\Phi_j(\mathbf{r};\mathbf{R})\}$ describes a complete basis of electronic states solution of the time-independent Schrödinger equation:

$$\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi_j(\mathbf{r};\mathbf{R}) = E_{el,j}(\mathbf{R})\Phi_j(\mathbf{r};\mathbf{R})$$

$\mathbf{R}$ is taken as a parameter.
Eigenfunctions of $\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})$ are considered to be orthonormal, i.e. $\langle\Phi_j|\Phi_i\rangle = \delta_{ij}$.

# Born-Oppenheimer approximation

$$\Psi(\mathbf{r},\mathbf{R},t) \xrightarrow[\text{Huang}]{\text{Born-}} \sum_j^\infty \Phi_j(\mathbf{r};\mathbf{R})\Omega_j(\mathbf{R},t)$$

Electrons are *static*. Use your favorite el. str. method.

For the nuclei, insert this *Ansatz* into the molecular time-dependent Schrödinger equation

$$\hat{H}\Psi(\mathbf{r},\mathbf{R},t) = i\hbar\frac{\partial}{\partial t}\Psi(\mathbf{r},\mathbf{R},t)$$

After left multiplication by $\Phi_k^*(\mathbf{r};\mathbf{R})$ and integration over $\mathbf{r}$, we obtain the following equation (we used $\langle\Phi_j|\Phi_i\rangle = \delta_{ij}$) :

$$\left[-\sum_I \frac{\hbar^2}{2M_I}\nabla_I^2 + E_{el,k}(\mathbf{R})\right]\Omega_k(\mathbf{R},t) + \sum_j^\infty D_{kj}\Omega_j(\mathbf{R},t) = i\hbar\frac{\partial}{\partial t}\Omega_k(\mathbf{R},t)$$

# Born-Oppenheimer approximation

$$\left[-\sum_I \frac{\hbar^2}{2M_I}\nabla_I^2 + E_{el,k}(\mathbf{R})\right]\Omega_k(\mathbf{R},t) + \sum_j D_{kj}\Omega_j(\mathbf{R},t) = i\hbar\frac{\partial}{\partial t}\Omega_k(\mathbf{R},t)$$

- Equation for the nuclear “wavepacket”, $\Omega(\mathbf{R},t)$, dynamics.
- $E_{el,k}(\mathbf{R})$ represents a potential energy surface for the nuclei.

Important additional term : $D_{kj}$ ! NONADIABATIC COUPLING TERMS

$$D_{kj} = \int \Phi_k^*(\mathbf{r};\mathbf{R})\left[\sum_I \frac{\hbar^2}{2M_I}\nabla_I^2\right]\Phi_j(\mathbf{r};\mathbf{R})d\mathbf{r}$$

$$+\sum_I \frac{1}{M_I}\left\{\int \Phi_k^*(\mathbf{r};\mathbf{R})\left[-i\hbar\nabla_I\right]\Phi_j(\mathbf{r};\mathbf{R})d\mathbf{r}\right\}\left[-i\hbar\nabla_I\right]$$

# Born-Oppenheimer approximation

$$D_{kj} = \int \Phi_k^*(\mathbf{r};\mathbf{R}) \left[ \sum_I \frac{\hbar^2}{2M_I} \nabla_I^2 \right] \Phi_j(\mathbf{r};\mathbf{R}) d\mathbf{r}$$
$$+ \sum_I \frac{1}{M_I} \left\{ \int \Phi_k^*(\mathbf{r};\mathbf{R}) \left[ -i\hbar\nabla_I \right] \Phi_j(\mathbf{r};\mathbf{R}) d\mathbf{r} \right\} \left[ -i\hbar\nabla_I \right]$$

If we neglect all the $D_{kj}$ terms (diagonal and off-diagonal), we have the Born-Oppenheimer approximation.

$$\left[ -\sum_I \frac{\hbar^2}{2M_I} \nabla_I^2 + E_{el,k}(\mathbf{R}) \right] \Omega_k(\mathbf{R},t) = i\hbar \frac{\partial}{\partial t} \Omega_k(\mathbf{R},t)$$

Mainly for ground state dynamics or for dynamics on states that do not couple with others.
(Back to nonadiabatic dynamics later).

# Born-Oppenheimer approximation: the nuclear trajectories

$$\left[-\sum_I \frac{\hbar^2}{2M_I}\nabla_I^2 + E_{el,k}(\mathbf{R})\right]\Omega_k(\mathbf{R},t) = i\hbar\frac{\partial}{\partial t}\Omega_k(\mathbf{R},t)$$

Using a polar expansion for $\Omega_k(\mathbf{R},t)$, we may find a way to obtain classical equation of motions for the nuclei.

$$\Omega_k(\mathbf{R},t) = A_k(\mathbf{R},t)\exp\left[\frac{i}{\hbar}S_k(\mathbf{R},t)\right].$$

$A_k(\mathbf{R},t)$ represents an amplitude and $S_k(\mathbf{R},t)/\hbar$ a phase.

Further: insert the polar representation into the equation above, do some algebra, and separate real and imaginary part, we obtain an interesting set of equations:

# Born-Oppenheimer approximation: the nuclear trajectories

$$\frac{\partial S_k}{\partial t} = \frac{\hbar^2}{2} \sum_I M_I^{-1} \frac{\nabla_I^2 A_k}{A_k} - \frac{1}{2} \sum_I M_I^{-1} \left(\nabla_I S_k\right)^2 - E_k$$

$$\frac{\partial A_k}{\partial t} = -\sum_I M_I^{-1} \nabla_I A_k \nabla_I S_k - \frac{1}{2} \sum_I M_I^{-1} A_k \nabla_I^2 S_k$$

Dependences of the functions $S$ and $A$ are omitted for clarity ($k$ is a index for the electronic state; in principle there is only one state in the adiabatic case).

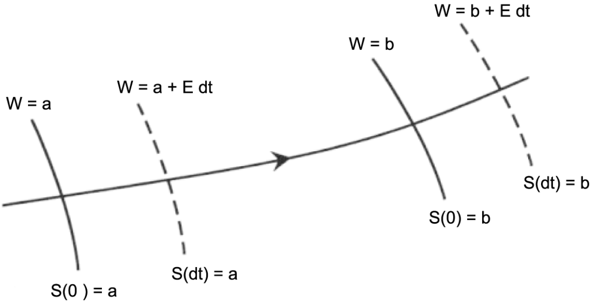
We have now a time-dependent equation for both the amplitude and the phase. Since we are in the adiabatic case there is only one PES and the second equation becomes trivially a diffusion continuity equation.

The nuclear dynamics is derived from the real part ($\frac{\partial S_k}{\partial t}$). This equation has again the form of a *classical* Hamilton-Jacobi equation.

# Born-Oppenheimer approximation: the nuclear trajectories

$$\frac{\partial S_k}{\partial t} = \frac{\hbar^2}{2}\sum_I M_I^{-1}\frac{\nabla_I^2 A_k}{A_k} - \frac{1}{2}\sum_I M_I^{-1}\left(\nabla_I S_k\right)^2 - E_k$$

$$\frac{\partial A_k}{\partial t} = -\sum_I M_I^{-1}\nabla_I A_k \nabla_I S_k - \frac{1}{2}\sum_I M_I^{-1} A_k \nabla_I^2 S_k$$



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# Born-Oppenheimer approximation: the nuclear trajectories

$$\frac{\partial S_k}{\partial t} = \frac{\hbar^2}{2}\sum_I M_I^{-1}\frac{\nabla_I^2 A_k}{A_k} - \frac{1}{2}\sum_I M_I^{-1}\left(\nabla_I S_k\right)^2 - E_k$$

The classical limit is obtained by taking[1]: $\hbar \longrightarrow 0$

$$\frac{\partial S_k}{\partial t} = -\frac{1}{2}\sum_I M_I^{-1}\left(\nabla_I S_k\right)^2 - E_k$$

These are the classical Hamilton-Jacobi equation and $S$ is the classical action related to a particle.

$$S(t) = \int_{t_0}^{t} L(t')dt' = \int_{t_0}^{t}\left[E_{kin}(t') - E_{pot}(t')\right]dt'$$

The momentum of a particle $I$ is related to

$$\nabla_I S = \mathbf{p}_I = \frac{\mathbf{v}_I}{M_I}$$

---

[1] Caution! This classical limit is subject to controversy...

# Born-Oppenheimer approximation: the nuclear trajectories

Therefore, taking the gradient,

$$-\nabla_J \frac{\partial S_k}{\partial t} = \frac{1}{2} \nabla_J \sum_I M_I^{-1} \left(\nabla_I S_k\right)^2 + \nabla_J E_k$$

and rearranging this equation using $\nabla_J S_k / M_J = \mathbf{v}_J^k$, we obtain the (familiar) Newton equation:

$$M_J \frac{d}{dt} \mathbf{v}_J^k = -\nabla_J E_k$$

In Summary:

**Adiabatic BO MD**

$$\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi_k(\mathbf{r};\mathbf{R}) = E_k^{el}(\mathbf{R})\Phi_k(\mathbf{r};\mathbf{R})$$

$$M_I \ddot{\mathbf{R}}_I = -\nabla_I E_k^{el}(\mathbf{R}) = -\underset{min \Phi_k}{\nabla_I} \langle \Phi_k | \hat{\mathcal{H}}_{el} | \Phi_k \rangle$$

# Mean-field vs. BO MD (adiabatic case)

## Ehrenfest dynamics

$$i\hbar\frac{\partial\Phi(\mathbf{r};\mathbf{R},t)}{\partial t}=\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi(\mathbf{r};\mathbf{R},t)$$

$$M_I\ddot{\mathbf{R}}_I=-\nabla_I\langle\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\rangle$$

Explicit time dependence of the electronic wavefunction.

## Born-Oppenheimer dynamics

$$\hat{\mathcal{H}}_{el}(\mathbf{r};\mathbf{R})\Phi_k(\mathbf{r};\mathbf{R})=E_k^{el}(\mathbf{R})\Phi_k(\mathbf{r};\mathbf{R})$$

$$M_I\ddot{\mathbf{R}}_I=-\nabla_I E_k^{el}(\mathbf{R})=-\underset{min\Phi_k}{\nabla_I}\langle\Phi_k|\hat{\mathcal{H}}_{el}|\Phi_k\rangle$$

The electronic wavefunction are *static* (only implicit time-dependence.

# Mean-field vs. BO MD (adiabatic case)

| Method | Born-Oppenheimer MD | Ehrenfest MD |
|---|---|---|
| | adiabatic MD (one PES) | nonadiabatic MD (mean-field) |
| | Efficient propagation of the nuclei | Get the "real" dynamics of the electrons |
| | Adiabatic nuclear propagation | Propagation of nuclei & electrons |
| | $\delta t$ ∼10-20 a.u. (0.25-0.5 fs) | $\delta t$ ∼0.01 a.u. (0.25 as) |
| | Simple algorithm | Common propagation of the nuclei and the electrons implies more sofisticated algorithm |

## Exact quantum dynamics?

Can we derive "exact" quantum equations of motion for the nuclei?
(without taking the classical limit $\hbar \to 0$?)

# Tarjectory-based quantum and mixed QM-CL solutions

We can "derive" the following trajectory-based solutions:

- Nonadiabatic Ehrenfest dynamics dynamics

  I. Tavernelli et al., *Mol. Phys.*, **103**, 963981 (2005).

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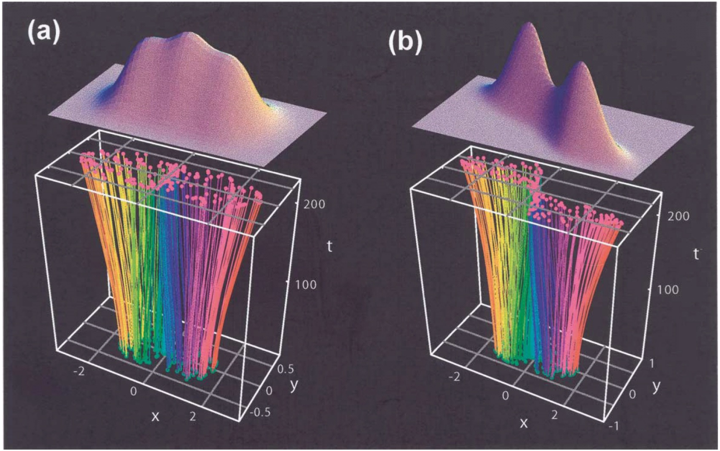
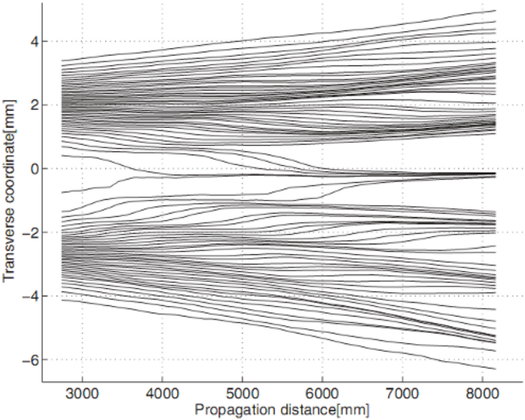
  A. Abedi, N. T. Maitra, E. K. U. Gross, *PRL*, **105**, 123002 (2010)

# Nonadiabatic dynamics: Multi-trajectory solutions



# Nonadiabatic Bohmian dynamics



Pioneers in quantum hydrodynamics: D. Bohm, P. R. Holland, R. E. Wyatt, and many others.

# NABDY: “exact” trajectory-based nonadiabatic dynamics

Using

- $\Psi(\mathbf{r}, \mathbf{R}, t) = \sum_j^\infty \Phi_j(\mathbf{r}; \mathbf{R}) \Omega_j(\mathbf{R}, t)$
- $\Omega_j(\mathbf{R}, t) = A_j(\mathbf{R}, t) \exp\left[\frac{i}{\hbar} S_j(\mathbf{R}, t)\right]$

in the exact time-dependent Schrödinger equation for the nuclear wavefucntion we get

$$\begin{aligned}
-\frac{\partial S_j(\mathbf{R}, t)}{\partial t} = & \sum_\gamma \frac{1}{2M_\gamma} \left(\nabla_\gamma S_j(\mathbf{R}, t)\right)^2 + E_j^{el}(\mathbf{R}) - \sum_\gamma \frac{\hbar^2}{2M_\gamma} \frac{\nabla_\gamma^2 A_j(\mathbf{R}, t)}{A_j(\mathbf{R}, t)} \\
& + \sum_{\gamma i} \frac{\hbar^2}{2M_\gamma} D_{ji}^\gamma(\mathbf{R}) \frac{A_i(\mathbf{R}, t)}{A_j(\mathbf{R}, t)} \Re\left[e^{i\phi}\right] - \sum_{\gamma, i \neq j} \frac{\hbar^2}{M_\gamma} \mathbf{d}_{ji}^\gamma(\mathbf{R}) \frac{\nabla_\gamma A_i(\mathbf{R}, t)}{A_j(\mathbf{R}, t)} \Re\left[e^{i\phi}\right] \\
& + \sum_{\gamma, i \neq j} \frac{\hbar}{M_\gamma} \mathbf{d}_{ji}^\gamma(\mathbf{R}) \frac{A_i(\mathbf{R}, t)}{A_j(\mathbf{R}, t)} \nabla_\gamma S_i(\mathbf{R}, t) \Im\left[e^{i\phi}\right]
\end{aligned}$$

and

$$\begin{aligned}
\frac{\partial A_j(\mathbf{R}, t)}{\partial t} = & -\sum_\gamma \frac{1}{M_\gamma} \nabla_\gamma A_j(\mathbf{R}, t) \nabla_\gamma S_j(\mathbf{R}, t) - \sum_\gamma \frac{1}{2M_\gamma} A_j(\mathbf{R}, t) \nabla_\gamma^2 S_j(\mathbf{R}, t) \\
& + \sum_{\gamma i} \frac{\hbar}{2M_\gamma} D_{ji}^\gamma(\mathbf{R}) A_i(\mathbf{R}, t) \Im\left[e^{i\phi}\right] - \sum_{\gamma, i \neq j} \frac{\hbar}{M_\gamma} \mathbf{d}_{ji}^\gamma(\mathbf{R}) \nabla_\gamma A_i(\mathbf{R}, t) \Im\left[e^{i\phi}\right] \\
& - \sum_{\gamma, i \neq j} \frac{1}{M_\gamma} \mathbf{d}_{ji}^\gamma(\mathbf{R}) A_i(\mathbf{R}, t) \nabla_\gamma S_i(\mathbf{R}, t) \Re\left[e^{i\phi}\right],
\end{aligned}$$

where both $S_j(\mathbf{R}, t)$ and $A_j(\mathbf{R}, t)$ are real fields and $\phi = \frac{1}{\hbar}(S_i(\mathbf{R}, t) - S_j(\mathbf{R}, t))$.

# NABDY: “exact” trajectory-based nonadiabatic dynamics

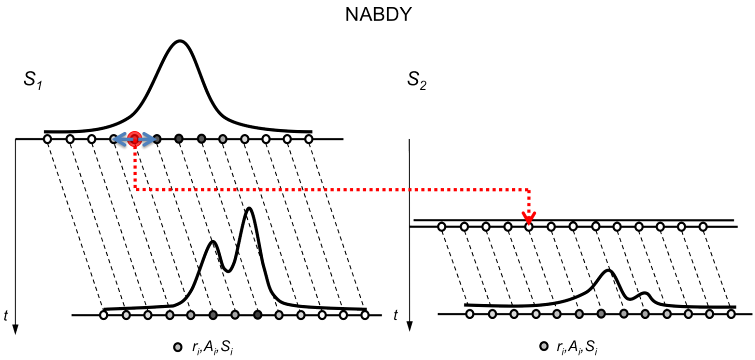


Figure: NABDY: sketch of the dynamics

# NABDY: "exact" trajectory-based nonadiabatic dynamics

From the NABDY equations we can obtain a Newton-like equation of motion (using the HJ definition of the momenta $\nabla_\beta S_j(\mathbf{R}, t) = \mathbf{P}^j_\beta$)

$$M_\beta \frac{d^2 \mathbf{R}_\beta}{(dt^j)^2} = -\nabla_\beta \left[ E^j_{el}(\mathbf{R}) + \mathcal{Q}_j(\mathbf{R}, t) + \sum_i D_{ij}(\mathbf{R}, t) \right]$$

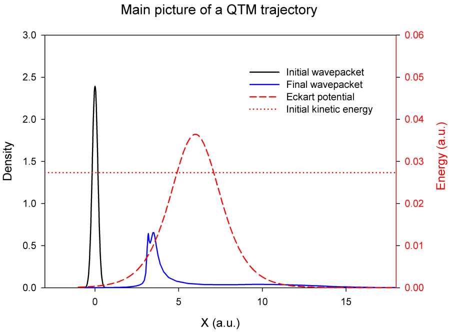
where $\mathcal{Q}_j(\mathbf{R}, t)$ is the quantum potential responsible for all coherence/decoherence "intrasurface" QM effects, and $\mathcal{D}_j(\mathbf{R}, t)$ is the *nonadiabatic* potential responsible for the amlpitude transfer among the different PESs.

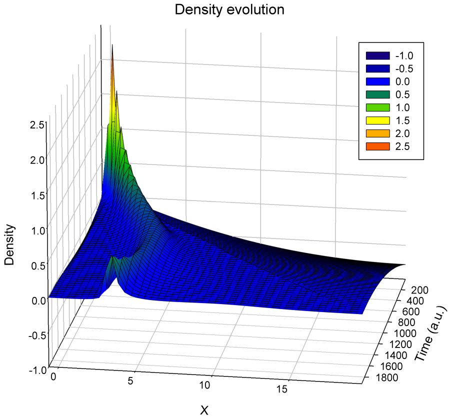
For more informations see:
B. Curchod, IT, U. Rothlisberger, *PCCP*, **13**, 3231 – 3236 (2011)

## NABDY limitations

- Mainly numerical challenges
- Instabilities induced by the quantum potential
- Compute derivatives in the $3N$ dimensional($\mathbb{R}^{3N}$) configuration space

# Gaussian wavepacket on an Eckart potential ($E_k = 3/4V$)





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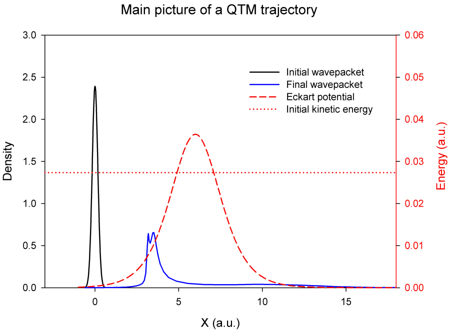
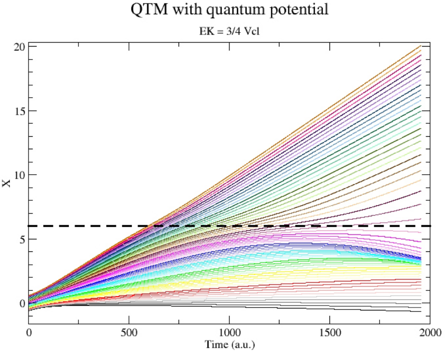
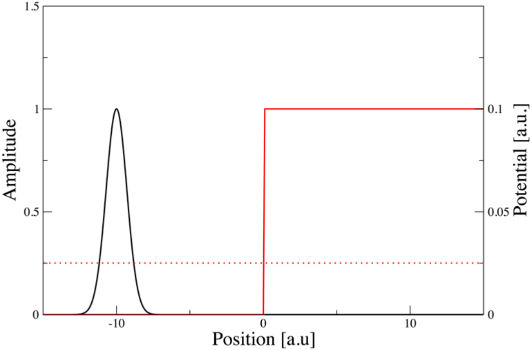




# NABDY: The strange world of Bohmian dynamics

**Bohmian trajectories cannot cross in the configuration space!**

This is the reason why Bohmian dynamics becomes difficult in low dimensions (contrary to all other quantum dynamics approaches)



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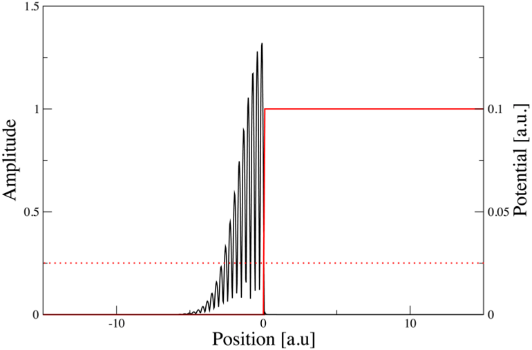

Nodes do not arise because of negative interference, BUT because of the avoiding crossing rule!

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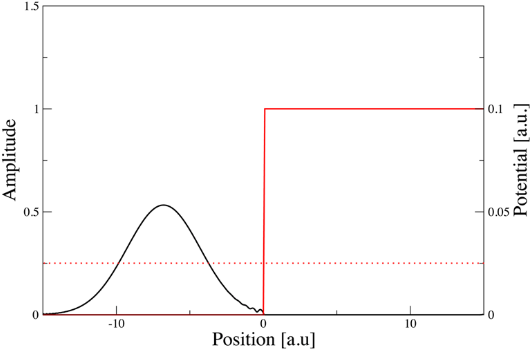

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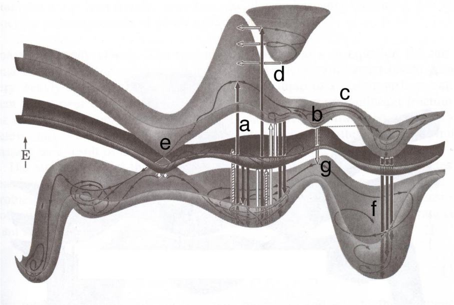
# Applications in Photochemistry and Photophysics

Trajectory-based solutions of the “exact” nonadiabatic equations are still impractical.



Approximate solutions are available. Among the most popular is

**Trajectory Surface Hopping (TSH)**

# Tarjectory-based mixed quantum-classical solutions

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# TSH nonadiabatic MD

There is no derivation of TSH dynamics. The fundamental hypothesis beyond TSH is that it is possible to design a dynamics that consists of:

- classical (*adiabatic*) time evolution of the nuclear trajectories on adiabatic states solution of the Schrödinger equation for the electronic sub-system
- propagation of a “quantum” amplitude, $C_k^\alpha(t)$, associated to each PES, $I$

  $$\Psi^\alpha(\mathbf{r}, \mathbf{R}, t) = \sum_k^\infty C_k^\alpha(t)\Phi_k(\mathbf{r}; \mathbf{R})$$

  (the label $\alpha$ is to recall that we have a different contribution from each different trajectory.)
- transitions (hops) of the trajectories between electronic states according to a stochastic algorithm, which depends on the nonadiabatic couplings and the amplitudes $C_k^\alpha(t)$

See also: J. Tully, *Faraday discussion*, **110**, 407 (1998) and B. Curchod, U. Rothlisberger, IT, in preparation.

# Tully's surface hopping - How does it work?

The main claim of TSH is that, the collection of a large enough set of independent trajectories gives an accurate representation of the nuclear wave packet

$$\rho_k^{CL}(\mathbf{R}^\alpha, t^\alpha) = \frac{N_k^\alpha(\mathbf{R}^\alpha, dV, t^\alpha)}{N_{tot}} \frac{1}{dV} \sim |\Omega_k(\mathbf{R}^\alpha, t^\alpha)|^2 \sim |C^\alpha_{k,\mathbf{R}^\alpha,t^\alpha}|^2$$

Inserting

$$\Psi^\alpha(\mathbf{r}, \mathbf{R}, t) = \sum_k^\infty C_k^\alpha(t)\Phi_k(\mathbf{r}; \mathbf{R})$$

into the time-dependent Schrödinger equation for the electrons and after some rearrangement, we obtain:

$$i\hbar\dot{C}_k^\alpha(t) = \sum_j C_j^\alpha(t)(H_{kj} - i\hbar\dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{kj}^\alpha)$$

with $H_{kj} = \langle\Phi_k(\mathbf{r}; \mathbf{R})\rangle|\hat{\mathcal{H}}_{el}|\Phi_j(\mathbf{r}; \mathbf{R})\rangle$.
In the adiabatic representation, we have $H_{kk} = E_k^{el}$ and $H_{kj} = 0$ .

# Tully's surface hopping - How does it work?

Equation of motion for the amplitudes:

$$i\hbar\dot{C}_k^\alpha(t) = \sum_j C_j^\alpha(t)(H_{kj} - i\hbar\dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{kj}^\alpha)$$

Switching algorithm:

In the fewest switches algorithm, the transition probability from state j to state k in the time interval [t, t + dt] is given by:

$$g_{jk}^\alpha(t, t+dt) \approx 2\int_t^{t+dt} d\tau \frac{\Im[C_k^\alpha(\tau)C_j^{\alpha*}H_{kj}(\tau)] - \Re[C_k^\alpha(\tau)C_j^{\alpha*}(\tau)\Xi_{kj}^\alpha(\tau)]}{C_j^\alpha(\tau)C_j^{\alpha*}(\tau)}$$

where $\Xi_{kj}^\alpha(\tau) = \dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{kj}^\alpha(\tau)$. A hop occurs between j and k if

$$\sum_{l\leq k-1} g_{jl}^\alpha < \zeta < \sum_{l\leq k} g_{jl}^\alpha ,$$

where $\zeta$ is generated randomly in the interval $[0,1]$.

# Tully's surface hopping - Summary

## Tully's surface hopping

$$i\hbar \dot{C}_k^\alpha(t) = \sum_j C_j^\alpha(t)(H_{kj} - i\hbar \dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{kj}^\alpha)$$

$$M_I \ddot{\mathbf{R}}_I = -\nabla_I E_k^{el}(\mathbf{R})$$

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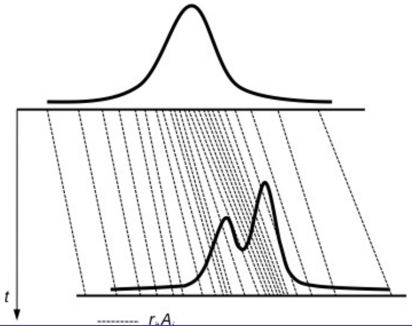

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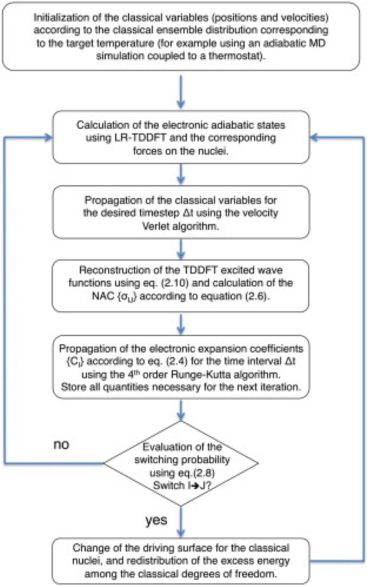
$$\sum_{l\leq k-1} g_{jl}^\alpha < \zeta < \sum_{l\leq k} g_{jl}^\alpha ,$$

**Some warnings:**

1. Evolution of classical trajectories (no QM effects – such as tunneling – are possible).
2. Rescaling of the nuclei velocities after a surface hop (to ensure energy conservation) is still a matter of debate.
3. Depending on the system studied, many trajectories could be needed to obtain a complete statistical description of the non-radiative channels.

For more details (and warnings) about Tully's surface hopping, see G. Granucci and M. Persico, J Chem Phys 126, 134114 (2007).

# Tully's surface hopping - The algorithm



Initialization of the classical variables (positions and velocities) according to the classical ensemble distribution corresponding to the target temperature (for example using an adiabatic MD simulation coupled to a thermostat).

Calculation of the electronic adiabatic states using LR-TDDFT and the corresponding forces on the nuclei.

Propagation of the classical variables for the desired timestep $\Delta t$ using the velocity Verlet algorithm.

Reconstruction of the TDDFT excited wave functions using eq. (2.10) and calculation of the NAC $\{\sigma_{IJ}\}$ according to equation (2.6).

Propagation of the electronic expansion coefficients $\{C_I\}$ according to eq. (2.4) for the time interval $\Delta t$ using the 4th order Runge-Kutta algorithm. Store all quantities necessary for the next iteration.

Evaluation of the switching probability using eq.(2.8) Switch I→J?

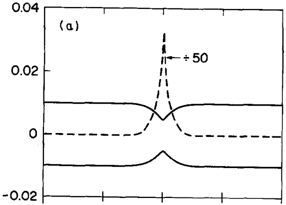
no

yes

Change of the driving surface for the classical nuclei, and redistribution of the excess energy among the classical degrees of freedom.

# Tully's surface hopping - Examples

1D systems



J.C. Tully, J. Chem. Phys. (1990), 93, 1061

http://wavepacket.sourceforge.net/

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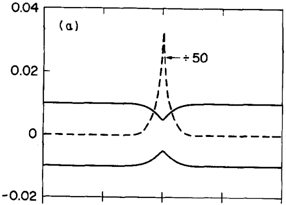

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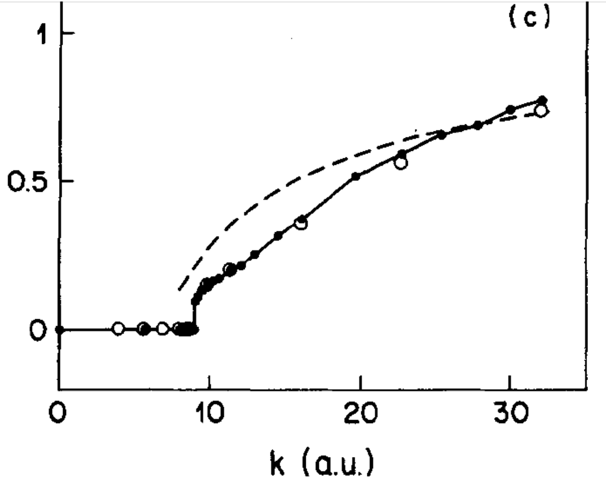
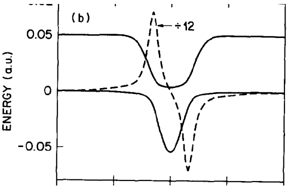


On the right: population of the upper state (k=mom)

- ○ exact
- ● TSH
- – Landau-Zener

# Tully's surface hopping - Examples

1D systems



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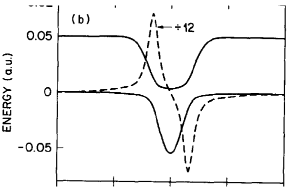

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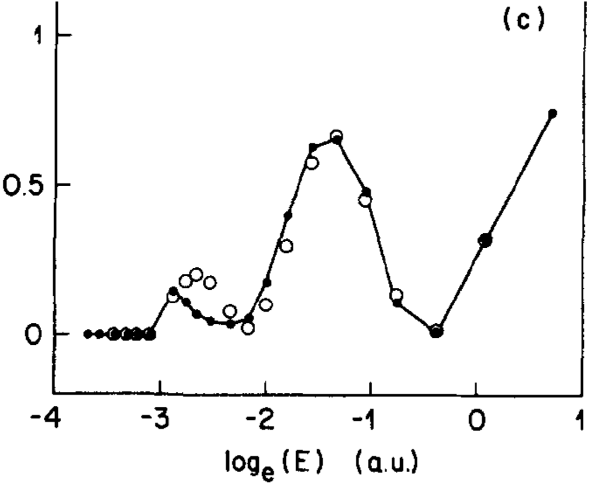

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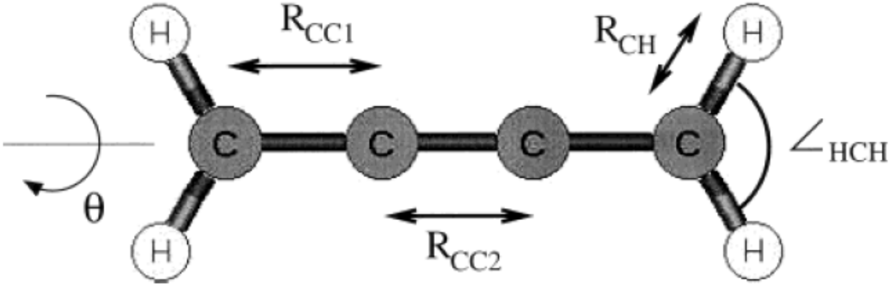
# Comparison with wavepacket dynamics

Butatriene molecule: dynamics of the radical cation in the first excited state.



JPCA,107,621 (2003)

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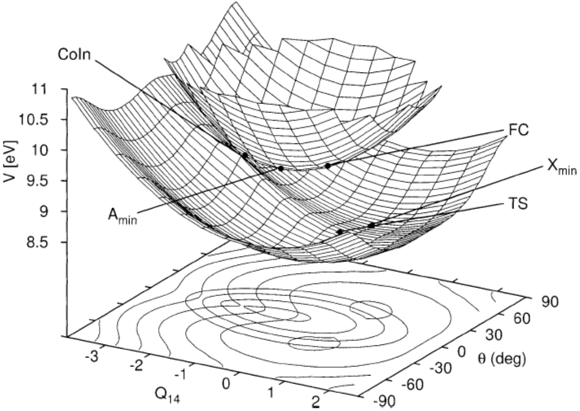
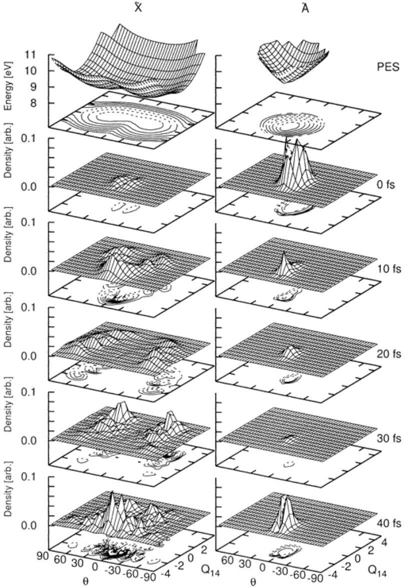


JPCA,107,621 (2003)

CASSCF PESs for the radical cation ($Q_{14}$: symmetric stretch, $\theta$: torsional angle).

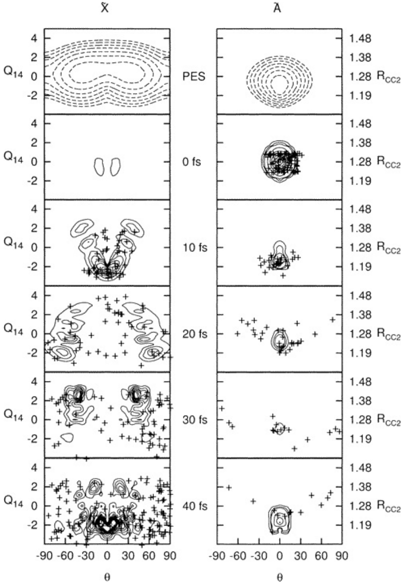
# Comparison with wavepacket dynamics



Nuclear wavepacket dynamics on fitted potential energy surfaces (using MCTDH with 5 modes) . Reappearing of the wavepacket in $S_1$ after $\sim$ 40fs.

JPCA,107,621 (2003)

# Comparison with wavepacket dynamics



On-the-fly dynamics with 80 trajectories (crosses).

Trajectories are not coming back close to the conical intersection.

What is the reason for this discrepancy? The independent trajectory approximation?, i.e. the fact that trajectories are not correlated?
(Or it has to do with differences in the PESs?)

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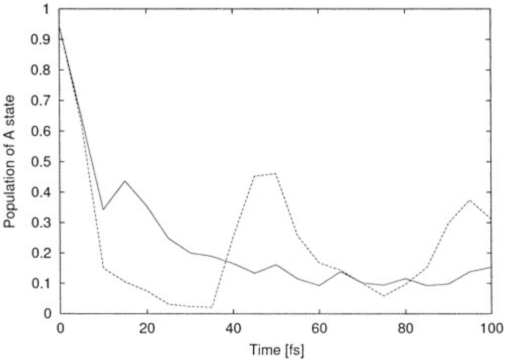

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(Or it has to do with differences in the PESs?)

JPCA,107,621 (2003)

# Tully's surface hopping - On-the-fly dynamics

## Tully's surface hopping

$$i\hbar\dot{C}_k^\alpha(t) = \sum_j C_j^\alpha(t)(H_{kj} - i\hbar\dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{kj}^\alpha)$$

$$M_I\ddot{\mathbf{R}}_I = -\nabla_I E_k^{el}(\mathbf{R})$$

$$\sum_{l\leq k-1} g_{jl}^\alpha < \zeta < \sum_{l\leq k} g_{jl}^\alpha \, ,$$

**What about the electronic structure method for on-the-fly dynamics? We need:**

- Potential energy surfaces → MR-CISD, LR-TDDFT, semiempirical, ...
- Forces on the nuclei → MR-CISD, LR-TDDFT, semiempirical methods, ... .
- Nonadiabatic coupling terms → MR-CISD, LR-TDDFT (?), semiempirical methods, ... .

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- Forces on the nuclei → MR-CISD, LR-TDDFT, semiempirical methods, ... .
- Nonadiabatic coupling terms → MR-CISD, LR-TDDFT (?), semiempirical methods, ... .

# Nonadiabatic couplings with LR-TDDFT?

Nonadiabatic coupling vectors are defined in terms of electronic wavefunctions:

$$\mathbf{d}_{kj} = \langle \Phi_k(\mathbf{R}) | \nabla_{\mathbf{R}} | \Phi_j(\mathbf{R}) \rangle = \frac{\langle \Phi_k(\mathbf{R}) | \nabla_{\mathbf{R}} \hat{\mathcal{H}}_{el} | \Phi_j(\mathbf{R}) \rangle}{E_j(\mathbf{R}) - E_k(\mathbf{R})}$$

The main challenge is to compute all these quantities as a functional of the ground state electronic density (or equivalently, of the occupied Kohn-Sham orbitals).

$$\mathbf{d}_{kj} \rightarrow \mathbf{d}_{kj}[\rho]$$

Different approaches for the calculation of $\mathbf{d}_{0j}[\rho]$ are available [2].
Here we will use the method based on the auxiliary many-electron wavefunctions.

---

[2] V. Chernyak and S. Mukamel, J. Chem. Phys. 112, 3572 (2000); R. Baer, Chem. Phys. Lett. 364, 75 (2002); E. Tapavicza, I. Tavernelli, and U. Rothlisberger, Phys. Rev. Lett. 98, 023001 (2007); C. P. Hu, H. Hirai, and O. Sugino, J. Chem. Phys. 127, 064103 (2007).

# The density response SOS formula

In TDDFT the density response $\chi(\omega)$ is

$$\chi(\omega) = \mathbb{S}^{-1/2}(\omega^2\mathbb{I} - \Omega(\omega))^{-1}\mathbb{S}^{-1/2}$$

where $\mathbf{Z}_n$ are the TDDFT eigenvectors of the pseudoeigenvalue equation, $\left(S_{ij\sigma,kl\tau} = \frac{\delta_{ik}\delta_{jl}\delta_{\sigma\tau}}{(f_{k\sigma}-f_{l\sigma})(\epsilon_{l\sigma}-\epsilon_{k\sigma})}\right)$

$$\mathbf{\Omega Z}_n = \omega_{0n}^2\mathbf{Z}_n \,,$$

with

$$\Omega_{ij\sigma,kl\tau} = \delta_{\sigma\tau}\delta_{ik}\delta_{jl}(\epsilon_{l\tau} - \epsilon_{k\sigma})^2 + 2\sqrt{(f_{i\sigma} - f_{j\sigma})(\epsilon_{j\sigma} - \epsilon_{i\sigma})}K_{ij\sigma,kl\tau}\sqrt{(f_{k\tau} - f_{l\tau})(\epsilon_{l\tau} - \epsilon_{k\tau})}$$

Using the spectral representation of the $(\omega^2\mathbb{I} - \Omega(\omega))^{-1}$, we can write

$$(\omega^2\mathbb{I} - \Omega(\omega))^{-1} = \sum_n \frac{\mathbf{Z}_n\mathbf{Z}_n^\dagger}{\omega_n^2 - \omega^2}$$

Therefore [3]

$$\chi(\omega) = \sum_n \frac{\mathbb{S}^{-1/2}\mathbf{Z}_n\mathbf{Z}_n^\dagger\mathbb{S}^{-1/2}}{\omega_n^2 - \omega^2}$$

and finally the perturbation of any observable $(\delta\mathcal{O}(\omega) = \sum_{ij\sigma} o_{ij\sigma}\delta P_{ij\sigma})$

$$\delta\mathcal{O}^{TDDFT}(\omega) = \sum_n \sum_{ij\sigma,kl\tau} o_{ij\sigma}\frac{(\mathbb{S}^{-1/2}\mathbf{Z}_n)_{ik\sigma}(\mathbf{Z}_n^\dagger\mathbb{S}^{-1/2})_{kl\tau}}{\omega_n^2 - \omega^2} v'_{kl\tau}E(\omega) \,.$$

---

[3] M. E. Casida, in Recent Advances in Density Functional Methods, edited by D. P. Chong, World Scientific, Singapore (1995), JCP, 130, 124107 (2007)

In MBPT the density response $\chi(\omega)$ is

$$\delta\mathcal{O}^{MBPT}(\omega) = \sum_n \frac{2\omega_{i0}\langle\Psi_0|\hat{\mathcal{O}}|\Psi_n\rangle\langle\Psi_n|\hat{v}'E(\omega)|\Psi_0\rangle}{\omega_{n0}^2 - \omega^2}$$

Equating $\delta\mathcal{O}^{TDDFT}(\omega)$ with $\delta\mathcal{O}^{MBPT}(\omega)$ residue-by-residue,

$$\langle\Psi_0|\hat{\mathcal{O}}|\Psi_n\rangle = \sum_{ij\sigma}^{(f_{i\sigma}-f_{j\sigma})>0} \frac{1}{\sqrt{\omega_n}} o_{ij\sigma}(\mathbb{S}^{-1/2}\mathbf{Z}_n)_{ij\sigma}$$

For any one-body operator, $\hat{\mathcal{O}}$, a mapping between MBPT and TDDFT quantities gives (for the moment, we only consider transitions from the ground state $\Psi_0$)

$$\mathbf{O}^\dagger\mathbb{S}^{-1/2}\mathbf{Z}_n = \omega_{0n}^{1/2}\langle\Psi_0|\hat{\mathcal{O}}|\Psi_n\rangle$$

where the operator $\hat{\mathcal{O}} = \sum_{ia\sigma} o_{ia\sigma}\hat{a}_{i\sigma}^\dagger\hat{a}_{a\sigma}$ has components $o_{ia\sigma} = \langle\phi_{i\sigma}|\hat{\mathcal{O}}|\psi_{a\sigma}\rangle$ [4] with $\omega_{0n} = E_n - E_0$. All matrices and vectors are given in the basis of KS orbitals $\{\phi_{i\sigma}\}$ with corresponding occupations $f_{i\sigma}$ and orbital energies $\epsilon_{i\sigma}$.

---

[4] $\sum_{ia\sigma}$ stands for $\sum_{i=1}^N \sum_{a=1}^\infty \sum_{\sigma\in\{\alpha,\beta\}}$.

# The auxiliary wavefunction

For practical purposes we introduce the auxiliary linear-response many-electron wavefunctions [5] as a linear combination of singly excited Slater determinants

$$\tilde{\Phi}_k[\{\phi.\}] = \sum_{ia\sigma} c^k_{ia\sigma}\, \hat{a}^\dagger_{a\sigma} \hat{a}_{i\sigma} \tilde{\Phi}_0[\{\phi.\}]\,,$$

with

$$c^k_{ia\sigma} \equiv \sqrt{\frac{S^{-1}_{ia\sigma}}{\omega_{0k}}} e^k_{ia\sigma}$$

where $\tilde{\Phi}_0[\{\phi.\}]$ is the Slater determinant of all occupied KS orbitals $\{\phi_{i\sigma}\}^N_{i=1}$, which, at a turn, are promoted into a virtual (unoccupied) orbitals, $\psi_{a\sigma}$.

We therefore have (in linear response!)

$$\langle \Psi_0 | \hat{\mathcal{O}} | \Psi_n \rangle = \langle \tilde{\Phi}_0 | \hat{\mathcal{O}} | \tilde{\Phi}_n \rangle$$

---

[5] JCP, 130, 124107 (2007), JCP, 131, 196101 (2009).

# Nonadiabatic couplings

- The *nonadiabatic coupling elements* at the mid step $t + \delta t/2$ of a LR-TDDFT AIMD can therefore be calculated as

$$\dot{\mathbf{R}} \cdot \mathbf{d}_{0k}|_{t+\delta t/2}[\{\phi_{\cdot}\}] = \left\langle \tilde{\Phi}_0(\mathbf{r};\mathbf{R}(t)) \,\middle|\, \nabla_{\mathbf{R}} \,\middle|\, \tilde{\Phi}_k(\mathbf{r};\mathbf{R}(t)) \right\rangle \cdot \dot{\mathbf{R}} = \left\langle \tilde{\Phi}_0(\mathbf{r};\mathbf{R}(t)) \,\middle|\, \frac{\partial}{\partial t} \,\middle|\, \tilde{\Phi}_k(\mathbf{r};\mathbf{R}(t)) \right\rangle$$

$$\simeq \frac{1}{2\delta t}\left[\langle \tilde{\Phi}_0(\mathbf{r};\mathbf{R}(t)) | \tilde{\Phi}_k(\mathbf{r};\mathbf{R}(t+\delta t))\rangle - \langle \tilde{\Phi}_0(\mathbf{r};\mathbf{R}(t+\delta t)) | \tilde{\Phi}_k(\mathbf{r};\mathbf{R}(t))\rangle\right]$$

- The *nonadiabatic coupling vectors* between **pairs of excites states** (second order response)

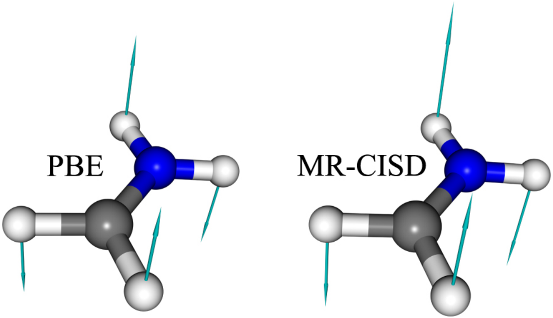
$$\mathbf{d}_{kj}[\{\phi_{\cdot}\}] = \frac{\langle \tilde{\Phi}_k(\mathbf{R}) | \nabla_{\mathbf{R}} \hat{\mathcal{H}}_{el} | \tilde{\Phi}_j(\mathbf{R}) \rangle}{E_j(\mathbf{R}) - E_k(\mathbf{R})}$$

  - Auxiliary many-electron wavefunctions give exact couplings between ground state and any (singly) excited state.
  - Auxiliary many-electron wavefunctions give *high quality* couplings between pairs of (singly) excited states ("exact" in the TDA and up to $\mathcal{O}(\delta\rho^3)$ in full response).

# Nonadiabatic couplings - examples

As always, the quality of the nonadiabatic couplings will depend on the *xc*-functional used...

$$\mathbf{d}_{kj}[\{\phi.\}] = \langle\tilde{\Phi}_k(\mathbf{R})|\nabla_{\mathbf{R}}|\tilde{\Phi}_j(\mathbf{R})\rangle$$



Protonated formaldimine: nonadiabatic coupling vectors $\mathbf{d}_{01}$ with LR-TDDFT/TDA.

# Nonadiabatic couplings - examples

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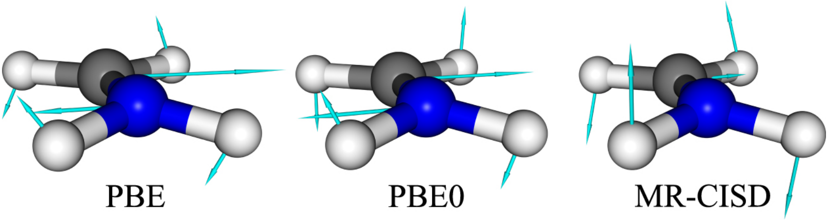

Protonated formaldimine: nonadiabatic coupling vectors $\mathbf{d}_{12}$ with LR-TDDFT/TDA.

# Protonated formaldimine

Protonated formaldimine ($\mathrm{CH_2NH_2^+}$) is the simplest model system of rhodopsin's chromophore retinal.

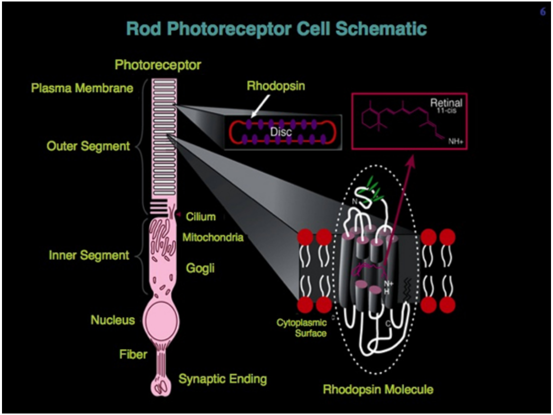

Figure: Rod photoreceptor cell responsible for dim-light vision. Credit: Rochester University.

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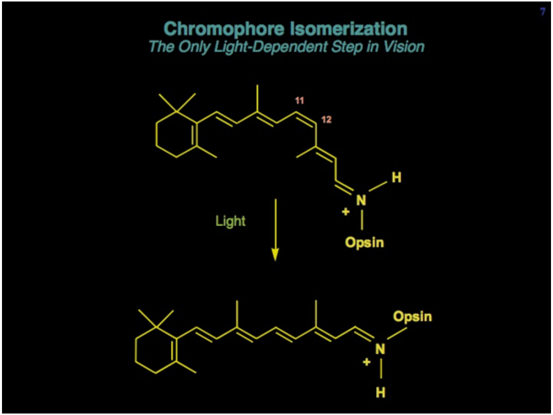

**Chromophore Isomerization**

*The Only Light-Dependent Step in Vision*

Figure: The visual chromophore, 11-cis-retinal, derived from vitamin A. Its isomerization after light absorption is the first step of vision. Credit: Rochester University.

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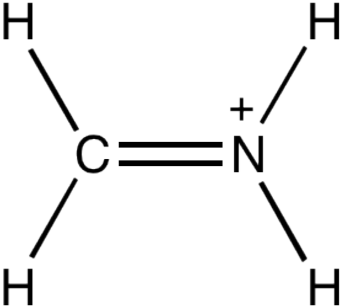
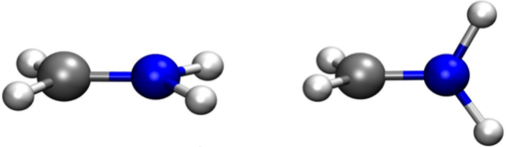


Figure: Protonated formaldimine

# Protonated formaldimine

The protonated formaldimine is a model compound for the study of isomerization in rhodopsin chromophore retinal.
In addition to the ground state (GS), two excited electronic states are of interest:

1. $S_1 : \sigma \rightarrow \pi^*$ (low oscillator strength)
2. $S_2 : \pi \rightarrow \pi^*$ (high oscillator strength)



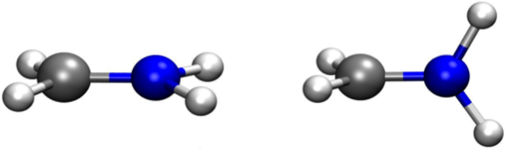
$CH_2NH_2^+$

# Protonated formaldimine

## Computational details

- Isolated system
- LR-TDDFT/PBE/TDA
- SH-AIMD
- 50 trajectories (NVT) each of ∼100 fs.

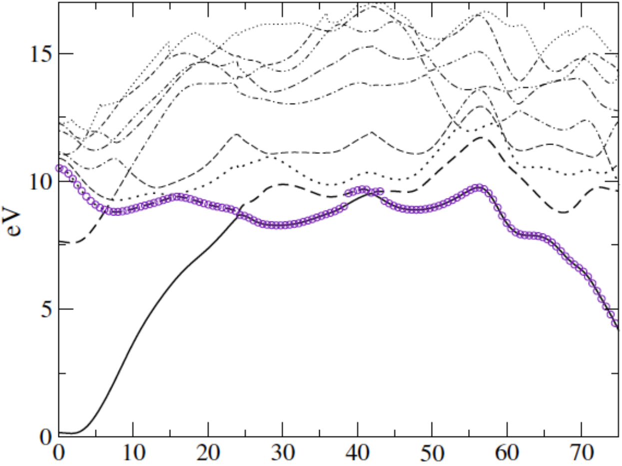
PRL, 98, 023001 (2007); THEOCHEM, 914, 22 (2009).



$CH_2NH_2^+$

# Protonated formaldimine

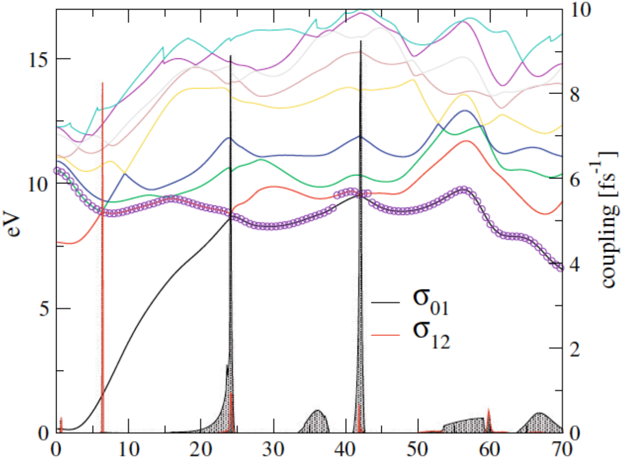
## Typical trajectory



PRL, 98, 023001 (2007); THEOCHEM, 914, 22 (2009).

# Protonated formaldimine

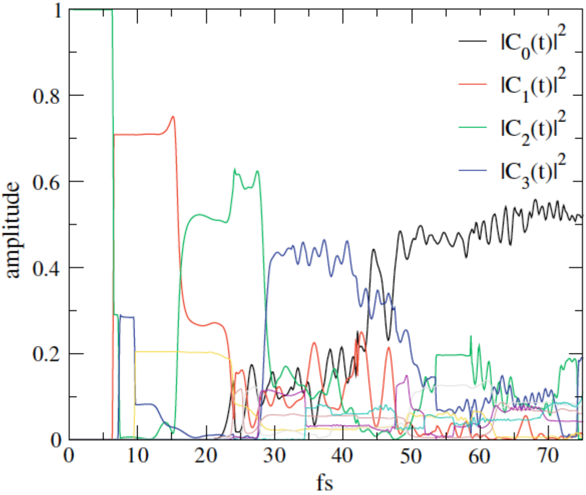
**Nonadiabatic couplings** $\sigma_{kj} = \dot{\mathbf{R}}^{\alpha} \cdot \mathbf{d}_{kj}^{\alpha}$



PRL, 98, 023001 (2007); THEOCHEM, 914, 22 (2009).

# Protonated formaldimine

## States population

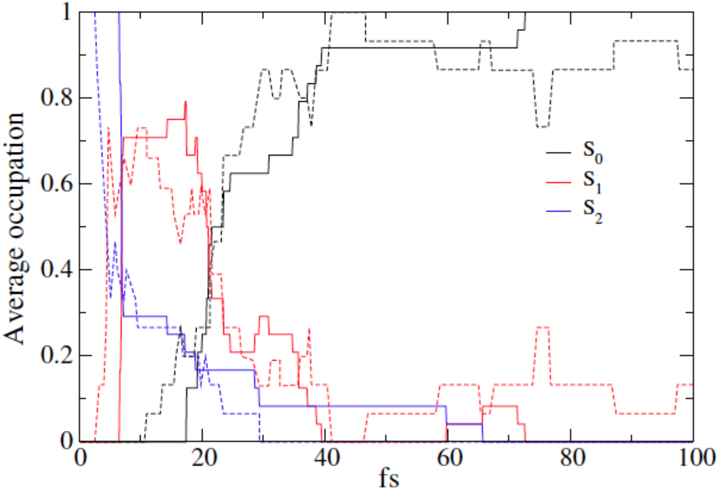


PRL, 98, 023001 (2007); THEOCHEM, 914, 22 (2009).

# Protonated formaldimine

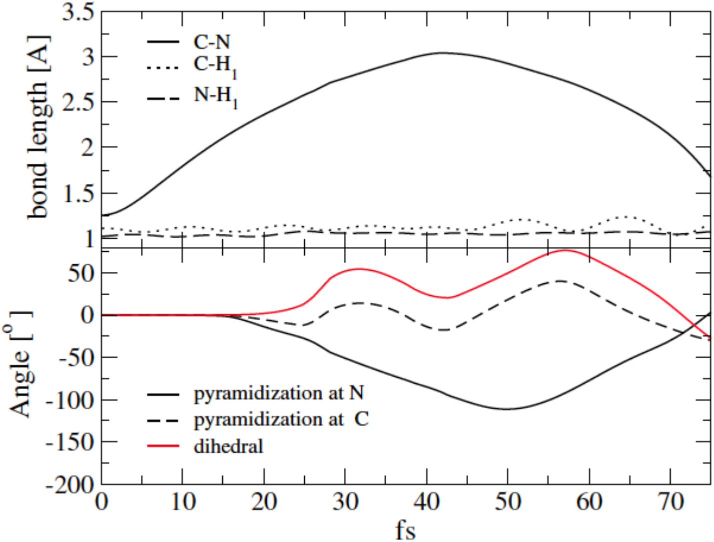
**States population** - **Average over many trajectories**.

Dashed line = CASSCF result.



# Protonated formaldimine

## Geometrical modifications



# Protonated formaldimine

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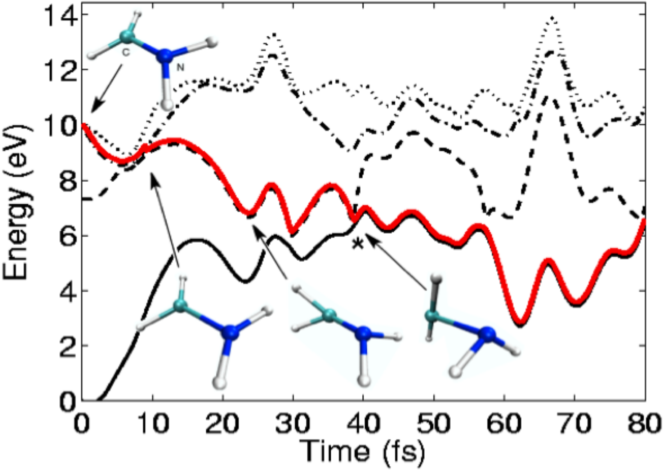
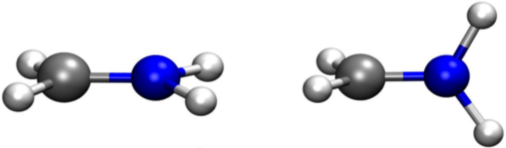


# Protonated formaldimine

## Comparison with experiment and model calculations

- In addition to the isomerization channel, intra-molecular proton transfer reactions was observed (formation of $\mathrm{CH_3NH^+}$).
- $\mathrm{H_2}$ abstraction is also observed in some cases.
- Structures and life times are in good agreement with reference calculations performed using high level wavefunction based methods.



$CH_2NH_2^+$

# Oxirane

Oxirane has interesting non-radiative decay channels, during which ring opening and dissociation can occur.

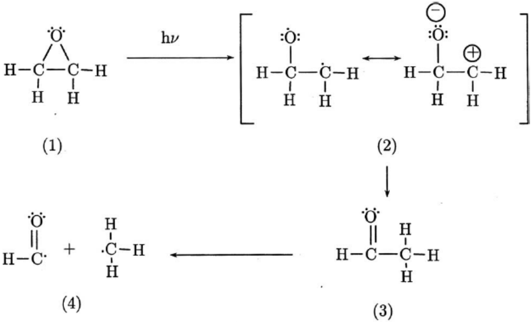


Figure: Mechanism proposed by Gomer and Noyes

# Oxirane

Oxirane has interesting non-radiative decay channels, during which ring opening and dissociation can occur.

## Computational details

- Isolated system
- LR-TDDFT/PBE/TDA
- SH-AIMD
- 30 trajectories (NVT) each of ∼100 fs.

JCP, 129, 124108 (2009).

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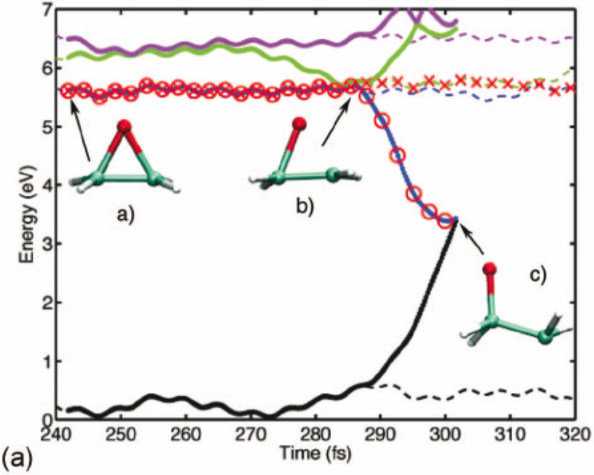
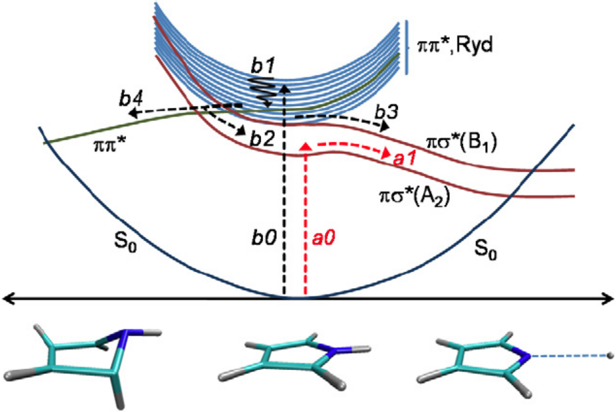


JCP, 129, 124108 (2009).

# Pyrrole

Study of the chemical products formed after light irradiation of pyrrole.



M. Barbatti, J. Pittner, M. Pederzoli, U. Werner, R. Mitric, V. Bonacic-Koutecky, H. Lischka, *Chem. Phys.*, **375**, 26, 2010; M. Vazdar, M. Eckert-Maksic, M. Barbatti, H. Lischka, *Mol. Phys.*, **107**, 845 (2009).

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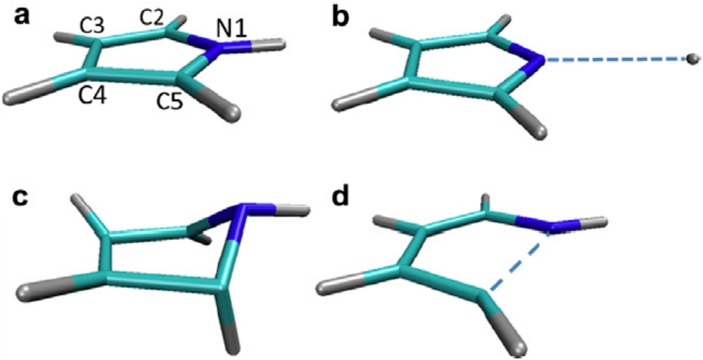

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# Pyrrole

Study of the chemical products formed after light irradiation of pyrrole.

## Computational details

- Isolated system
- LR-TDDFT/PBE0
- SH-AIMD
- $\sim$ 115 trajectories
- Comparison with AIMD-TSH using MR-CISD

Chem. Phys., 375, 26 (2010); Mol. Phys., 107, 845 (2009).

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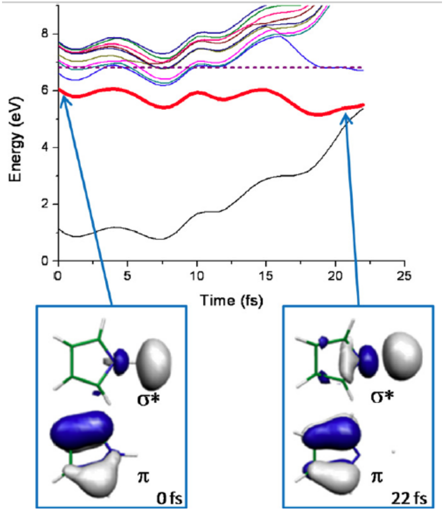

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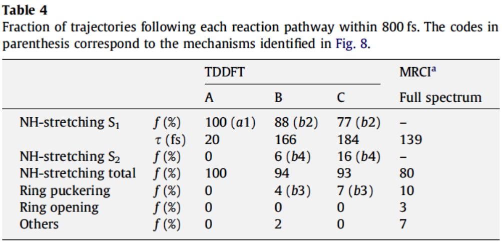
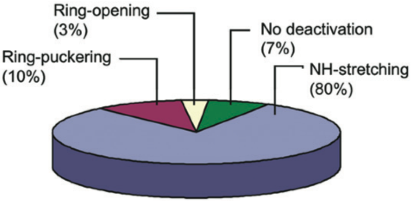
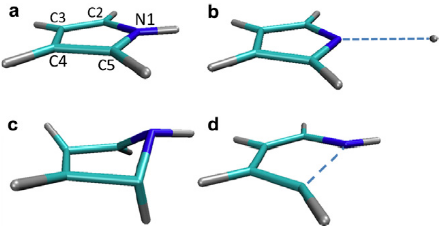


**Table 4**
Fraction of trajectories following each reaction pathway within 800 fs. The codes in parenthesis correspond to the mechanisms identified in Fig. 8.

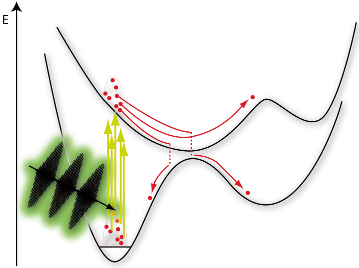
| | | TDDFT | | | MRCI[a] |
|---|---|---|---|---|---|
| | | A | B | C | Full spectrum |
| NH-stretching $S_1$ | $f$ (%) | 100 (*a*1) | 88 (*b*2) | 77 (*b*2) | – |
| | $\tau$ (fs) | 20 | 166 | 184 | 139 |
| NH-stretching $S_2$ | $f$ (%) | 0 | 6 (*b*4) | 16 (*b*4) | – |
| NH-stretching total | $f$ (%) | 100 | 94 | 93 | 80 |
| Ring puckering | $f$ (%) | 0 | 4 (*b*3) | 7 (*b*3) | 10 |
| Ring opening | $f$ (%) | 0 | 0 | 0 | 3 |
| Others | $f$ (%) | 0 | 2 | 0 | 7 |





# TSH with external time-dependent fields

**Addition of an external field within the equations of motion of TSH:**



### Startegy

The idea is to induce electronic excitations through the direct interaction with the time-dependent (td) electric field instead of “artificially” promote the system into one of its excited states.

**Method**: extended TSH nonadiabatic dynamics.

IT, B. Curchod, U. Rothlisberger, Phys. Rec. A **81**, 052508 (2010)

# TSH with external time-dependent fields

**Short summary of the theory**

The interaction Hamiltonian between the electrons and the td electric field is

$$\hat{H}_{int} = -\frac{e}{2m_e c} \sum_i \mathbf{A}(\mathbf{r}_i, t) \cdot \hat{\mathbf{p}}_i$$

where $\mathbf{A}(\mathbf{r}, t)$ is the (classical) vector potential of the electromagnetic field, $\hat{\mathbf{p}}_i$ is the momentum operator of electron $i$, $e$ is the electron charge, $m_e$ is the electron mass, and $c$ is the speed of light.

Remark

We are in the *dipole approximation* and therefore we do not need TDCDFT.

IT, B. Curchod, U. Rothlisberger, Phys. Rec. A **81**, 052508 (2010)

# External field within TSH

It can be shown (Phys. Rev. A **81** 052508 (2010)) that through the coupling with the td electric field, Tully's propagation equations acquire an additional term

$$i\hbar\dot{C}_J^\alpha(t) = \sum_I C_I^\alpha(t)(H_{JI} - i\hbar\dot{\mathbf{R}}^\alpha \cdot \mathbf{d}_{JI}^\alpha + i\omega_{JI}\frac{A_0}{c}\boldsymbol{\epsilon}^\lambda \cdot \boldsymbol{\mu}_{JI}^\alpha e^{-i\omega t})$$

with

$$i\omega_{JI}\frac{\mathbf{A}_0(t)}{c} \cdot \boldsymbol{\mu}_{JI} = \langle\Phi_J|\hat{H}_{int}|\Phi_I\rangle$$

and where $\mathbf{A}_0(t) = A_0\boldsymbol{\epsilon}^\lambda e^{-i\omega t}$ is the vector potential of the external td electric field,

$$\boldsymbol{\mu}_{JI} = -e\langle\Phi_J|\sum_i \hat{\mathbf{r}}_i|\Phi_I\rangle$$

is the the transition dipole vector, and $\omega_{JI} = (E_J - E_I)/\hbar$.

Note that Tully's hops probability should be modified accordingly.

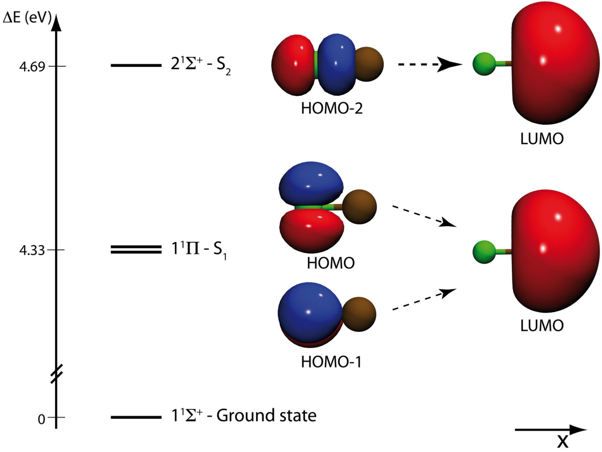
IT, B. Curchod, U. Rothlisberger, Phys. Rec. A **81**, 052508 (2010)

# Effect of an electromagnetic field - Lithium fluoride

Different excitations can be obtained, depending on the polarization vector of the laser pulse.

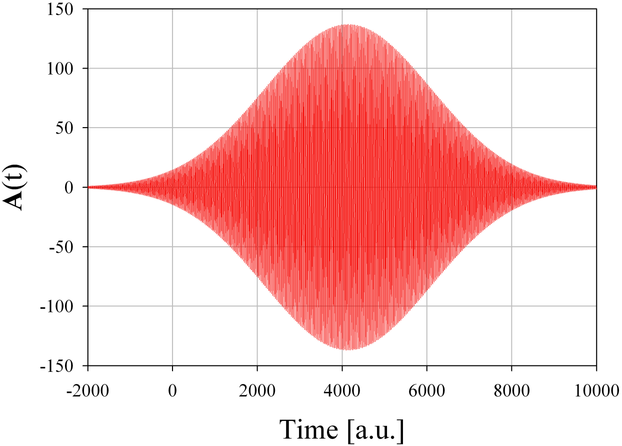
## Electronic structure of LiF

- Ground state - Σ symmetry (GS) .
- First excited state (doubly degenerate) - Π symmetry ($S_1$) .
- Second excited state - Σ symmetry ($S_2$) .
- Avoided crossing between GS and $S_2$



# Effect of an electromagnetic field - Lithium fluoride

Pulse: $\mathbf{A}(t) = -A_0 \epsilon^{\lambda} \exp\left(-\frac{(t-t_0)^2}{T^2}\right) \sin(\omega t)$



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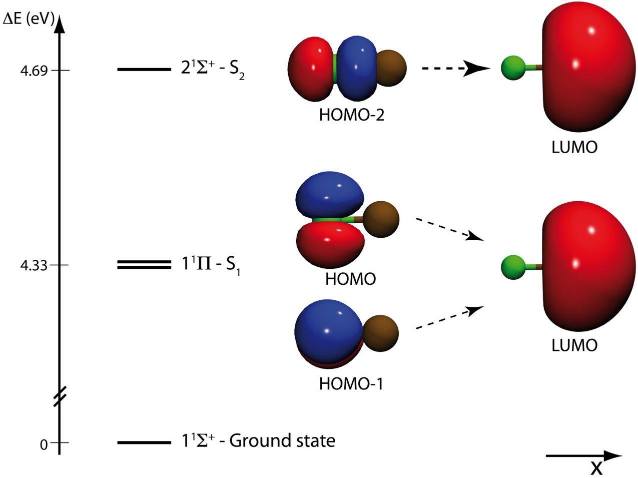
x-polarized pulse: $\epsilon^{\lambda} = (1, 0, 0)$



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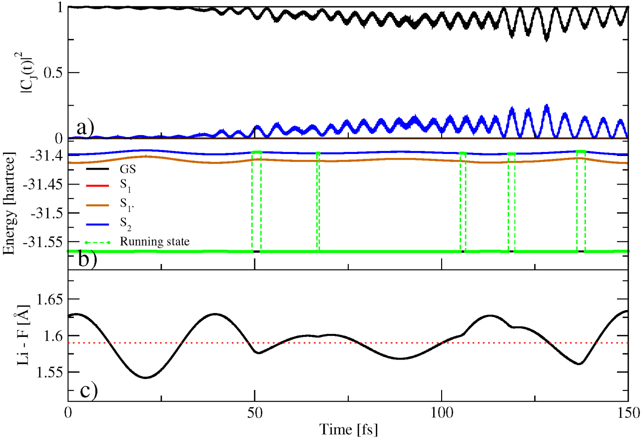
$\epsilon^{\lambda} = (1, 0, 0)$



IT, B. Curchod, U. Rothlisberger, Phys. Rec. A **81**, 052508 (2010)

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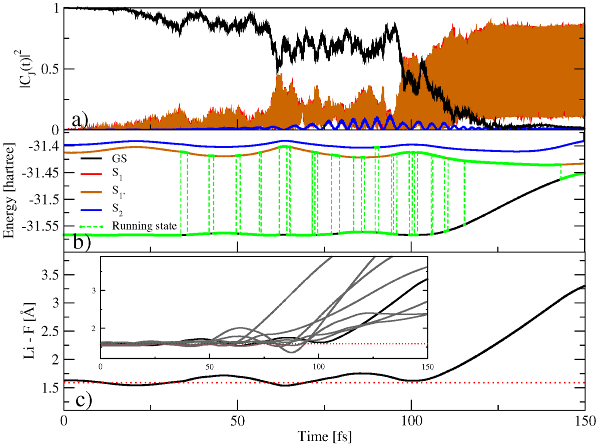
$$\epsilon^{\lambda} = \frac{1}{\sqrt{3}}(1, 1, 1)$$



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