



Overview of spectroscopies III

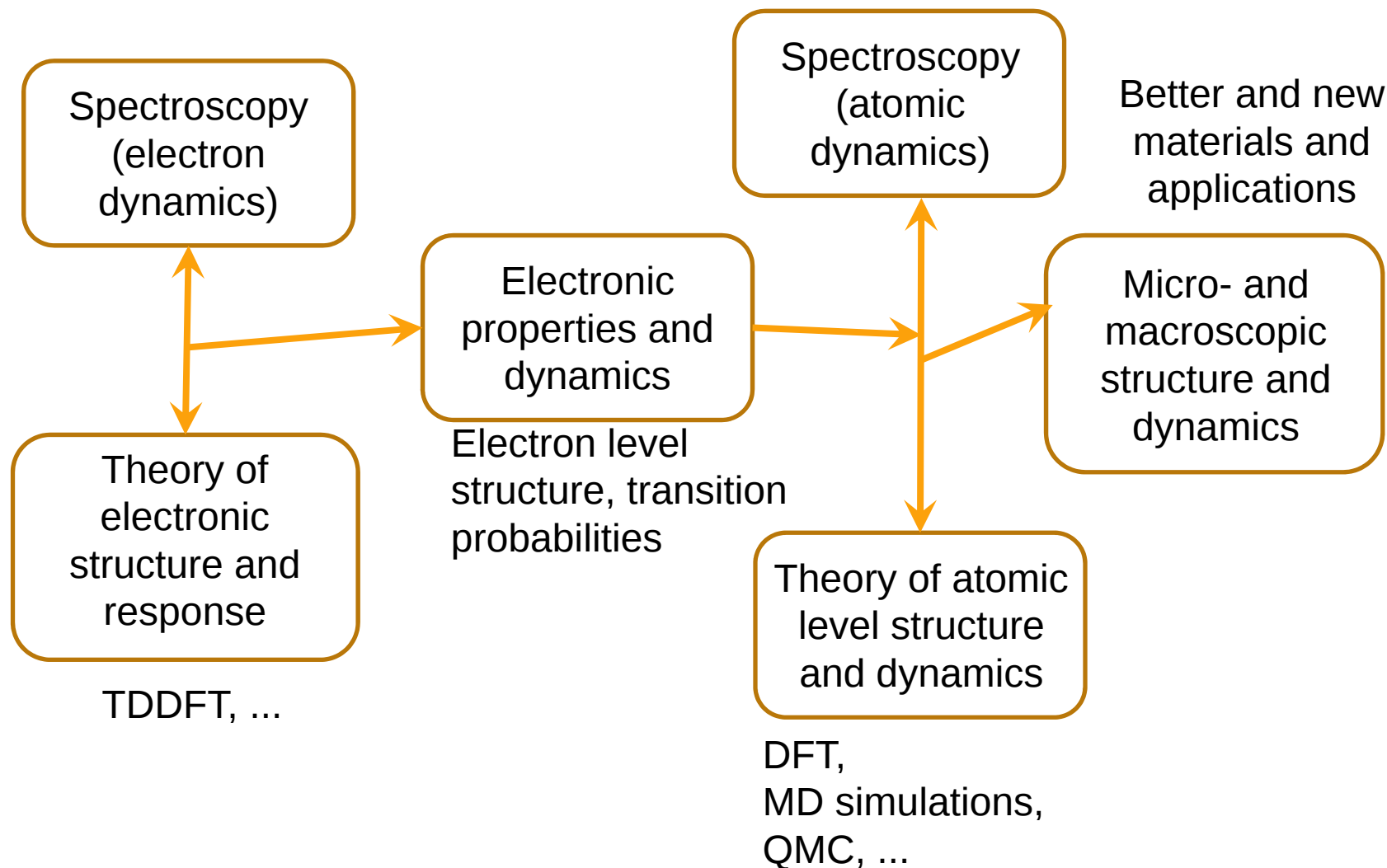
Simo Huotari

University of Helsinki, Finland

TDDFT school, Benasque, Spain, January 2012



Motivation: why we need theory



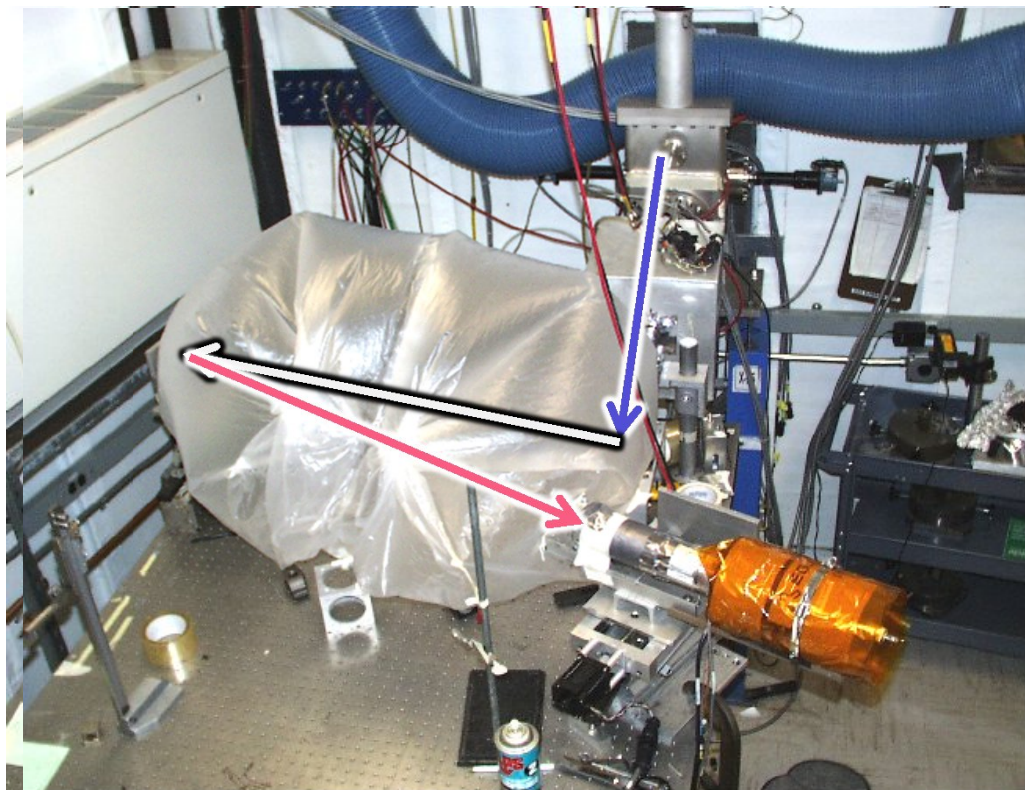


Measurement vs. experiment

Equipment for measurements

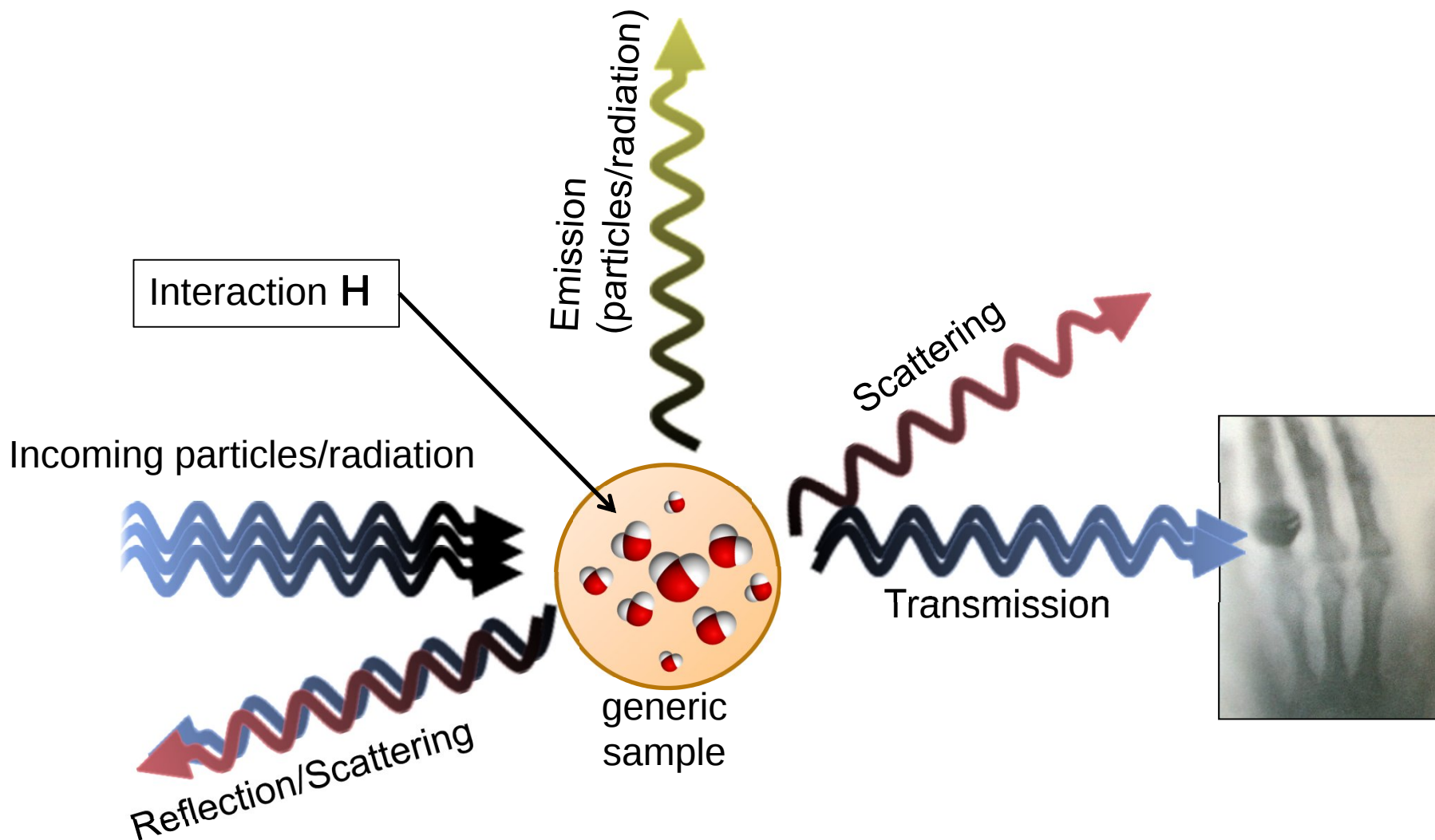


Equipment for experiments





Probes and messengers





Outline

Part 1 – Excitations and properties

Part 2 – Techniques (general)

Part 3 – Sample environments

Tomorrow – Techniques (examples)



Classification of excitations

(non-exhaustive)

	Collective	"Single particle"
Vibrational	Sound waves (phonons)	Molecular bending, stretching...
Spin	Magnons, spin waves	Spin flip
Electronic	Plasmons, orbitons, polarons	Crystal fields, excitons, Compton, electron removal

+ multiples (bimagnons, double plasmons, ...)



Properties

Macroscopic **dielectric function** $\varepsilon_M(Q, E) = \varepsilon_1(Q, E) + i \varepsilon_2(Q, E)$

Complex **refractive index** $n + i \kappa$

Reflectivity R

Absorption coefficient $\alpha = 4\pi\kappa/\lambda$

Loss function $-\text{Im}[\varepsilon_M^{-1}(\mathbf{Q}, E)]$

Dynamic structure factor $S(\mathbf{Q}, E) = -\left(\frac{Q^2}{4\pi^2 e^2 n}\right) \text{Im}[\varepsilon_M^{-1}(\mathbf{Q}, E)]$

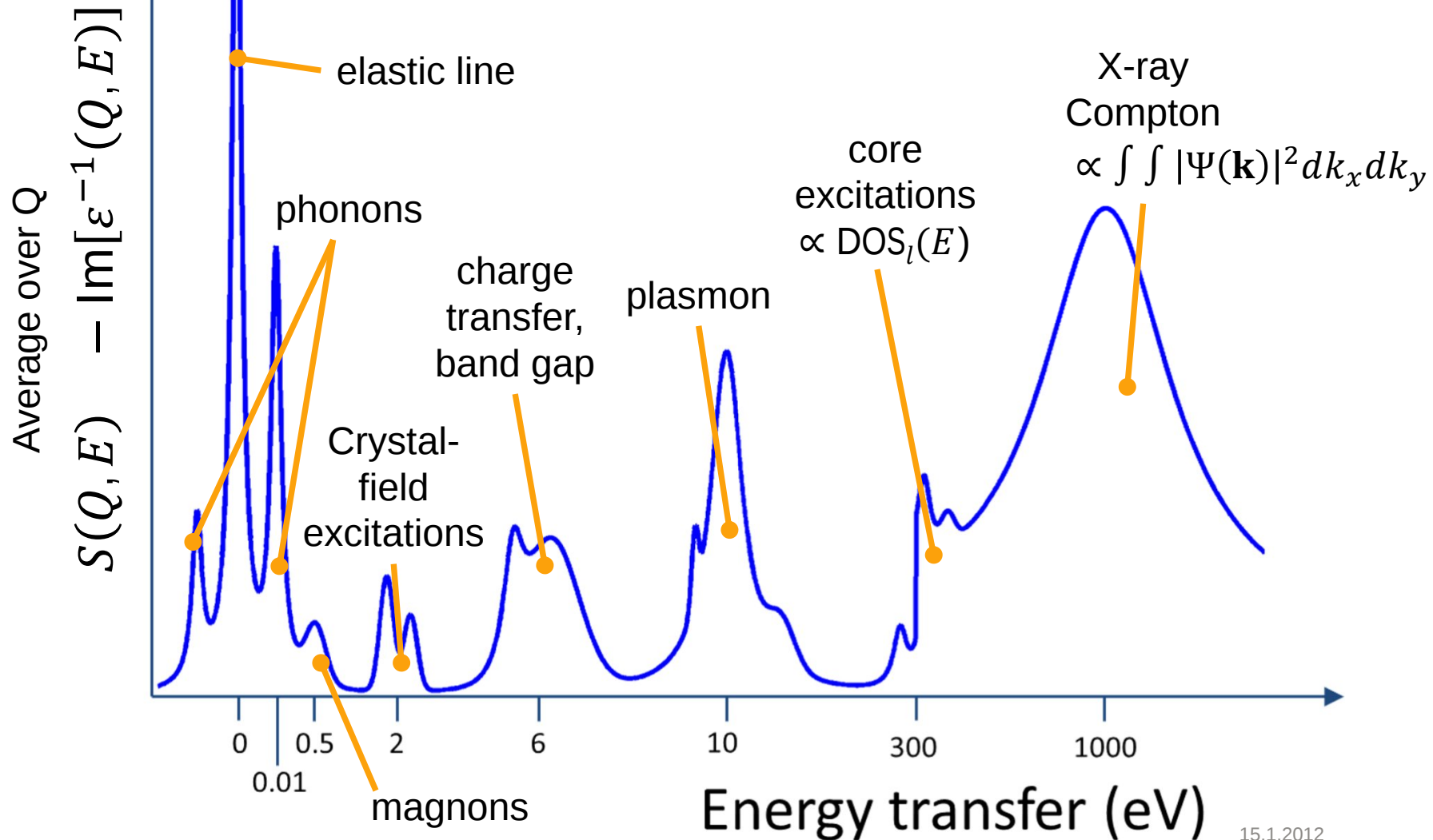
Spectral density function $A(k, \omega)$

Resonant Raman spectra

Dynamic structure factor is also the Fourier transform of the **density correlation function** - which is the probability to find two different particles separated by $\mathbf{r}$ and t .



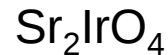
Energy-loss spectrum



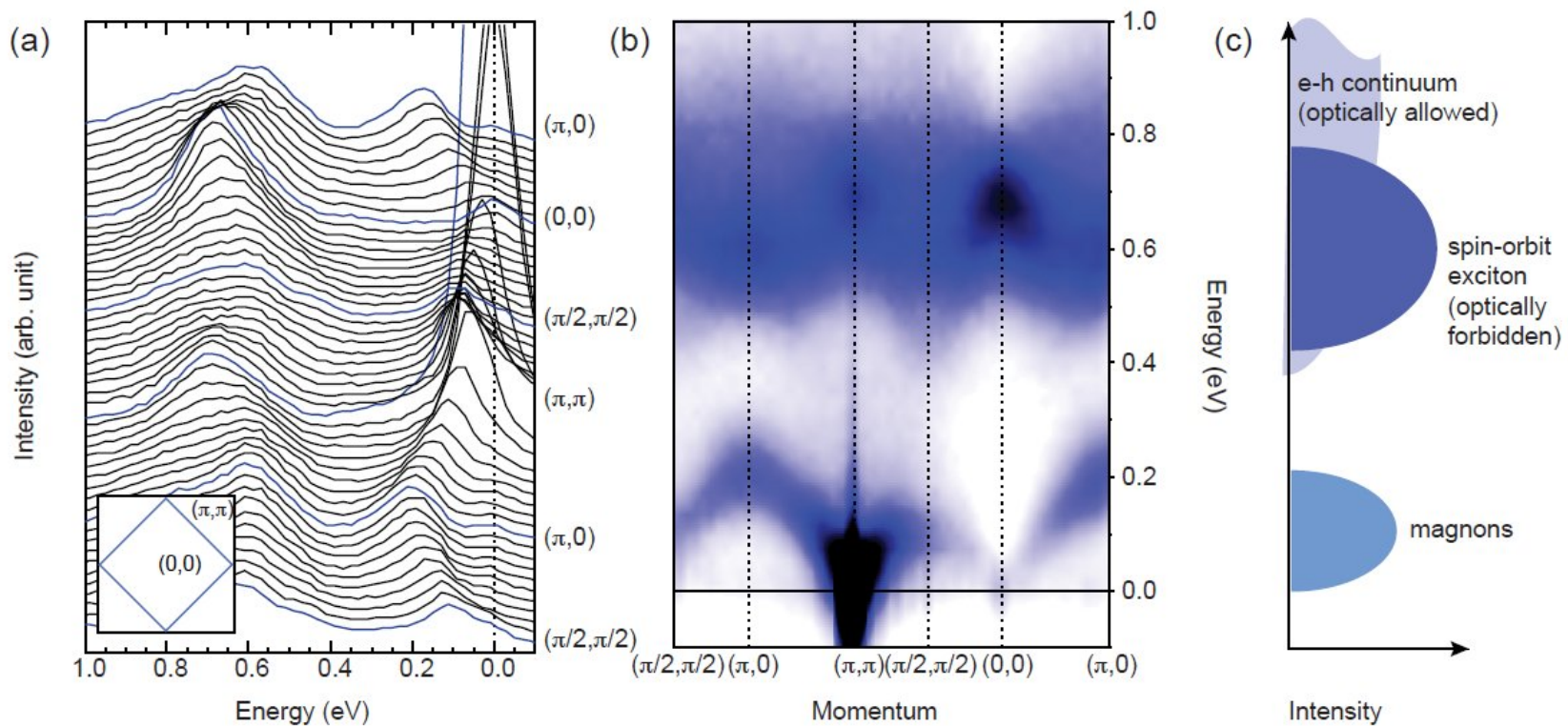


Magnons: spin waves

Magnetic excitations: orbital physics in transition metal oxides

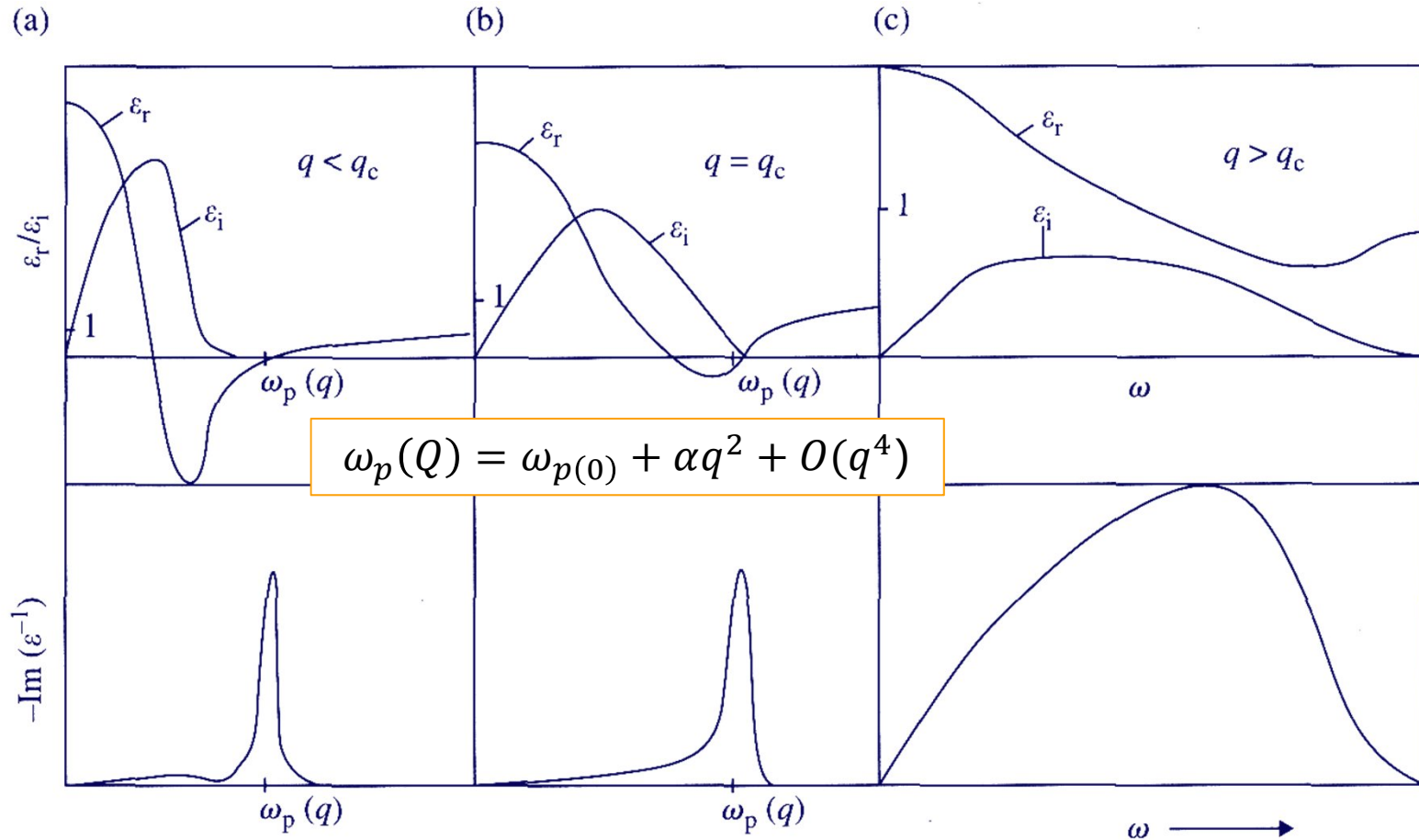


J. Kim et al. 2011, arXiv:1110.0759v1 [cond-mat.str-el]





Plasmons



W. Schülke: Electron dynamics by inelastic x-ray scattering (Oxford Univ. Press)

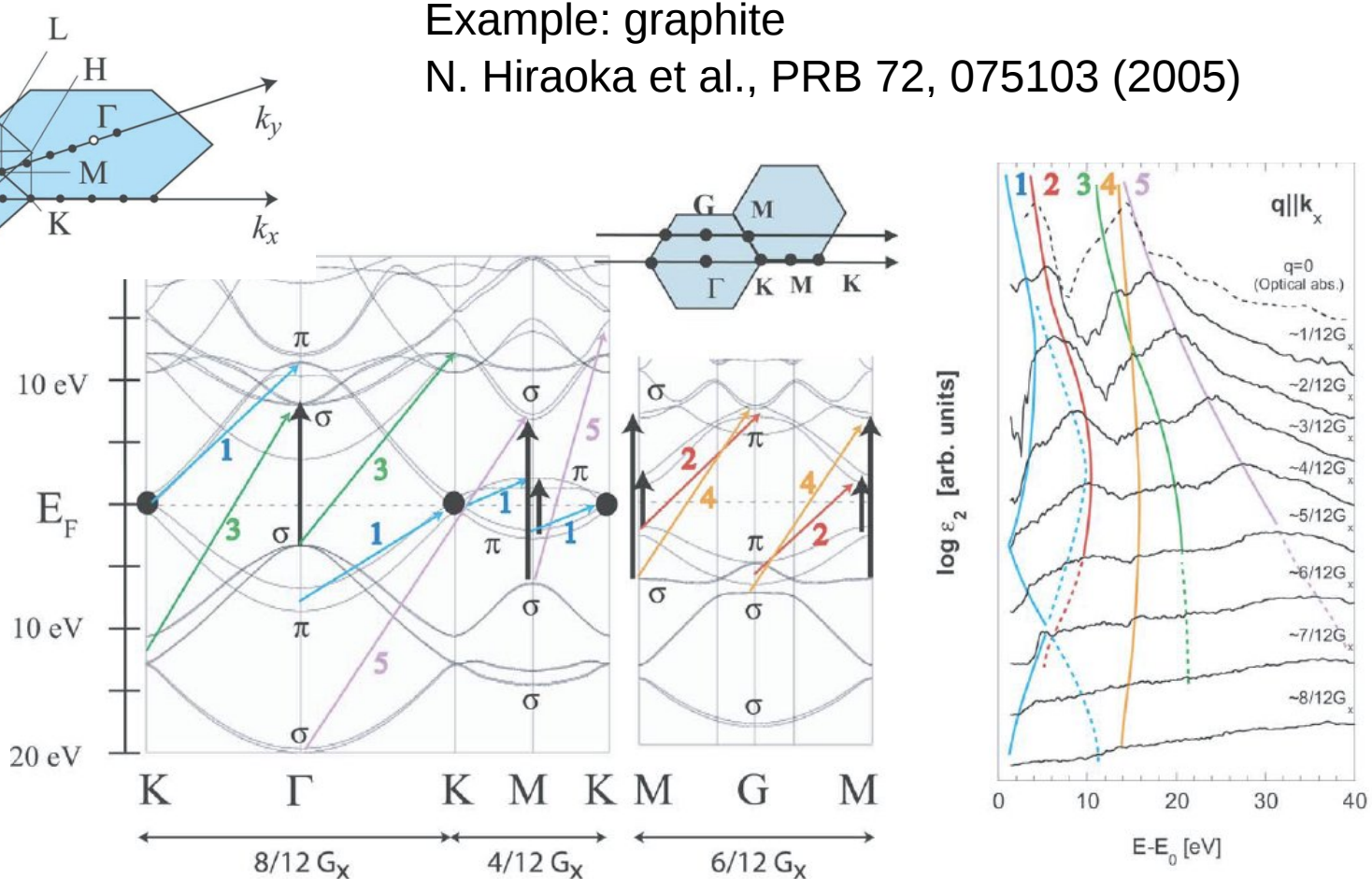


Band structure picture

Both **energy** and **momentum** can be controlled

Example: graphite

N. Hiraoka et al., PRB 72, 075103 (2005)





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Outgoing particle

Incoming particle

	Nothing	Photon	Electron
Nothing	Skiing at the Pyrenees	Light emission, photoluminescence	Electron emission
Photon	Absorption (IR, UV/vis, vacuum UV, x-ray..) CD; XMCD	Ellipsometry, reflectometry, ATR, Raman, Brillouin, x-ray scattering, Compton scattering	Photoemission Auger spectroscopy
Electron	Electron absorption	Inverse photoemission Cathodo-luminescence	Electron energy loss Tunneling

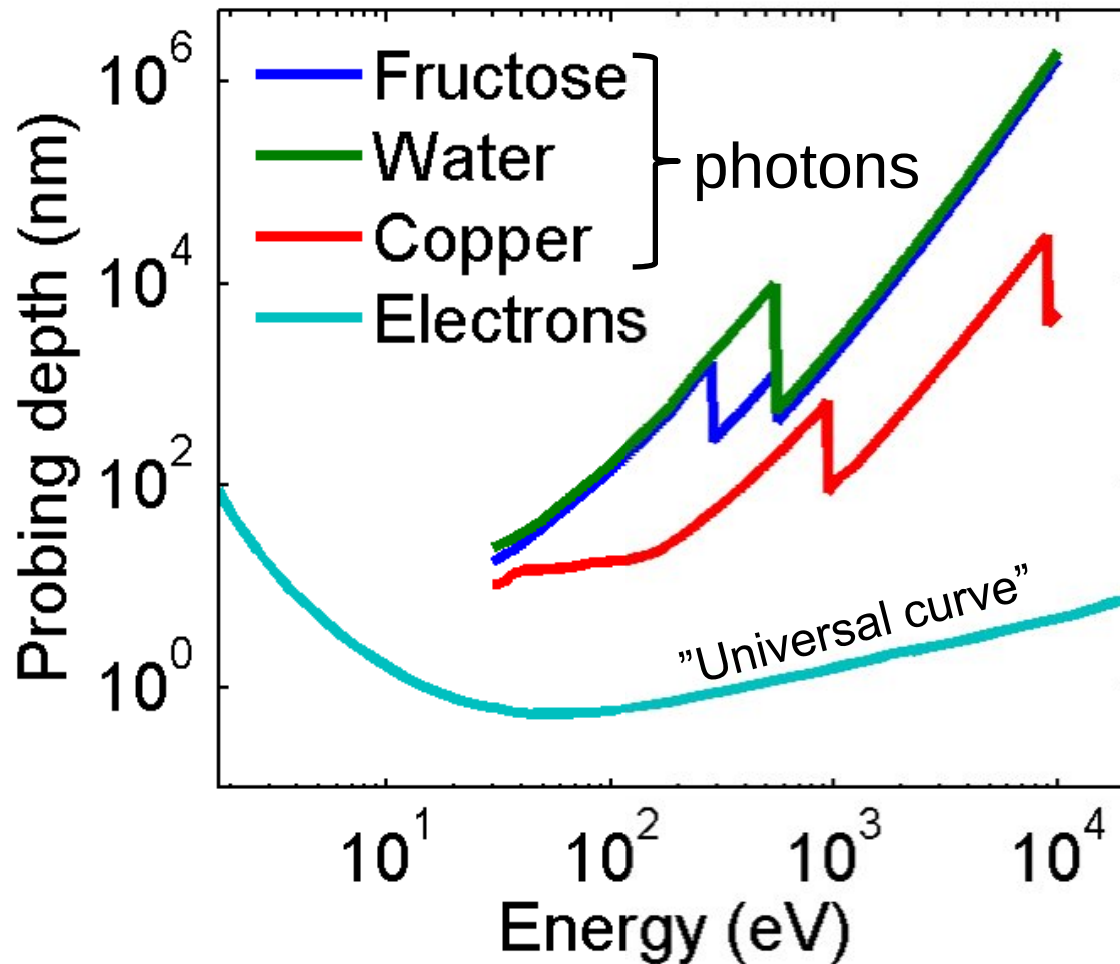
- + Time-of-flights
- + Non-linear phenomena

Differences in:

- Dynamic range of Q and E, coupling to spin, nuclei, ...
- Bulk or surface sensitivity
- Resolving power in energy, momentum, time, space...
- Element specificity (useful for complicated systems)

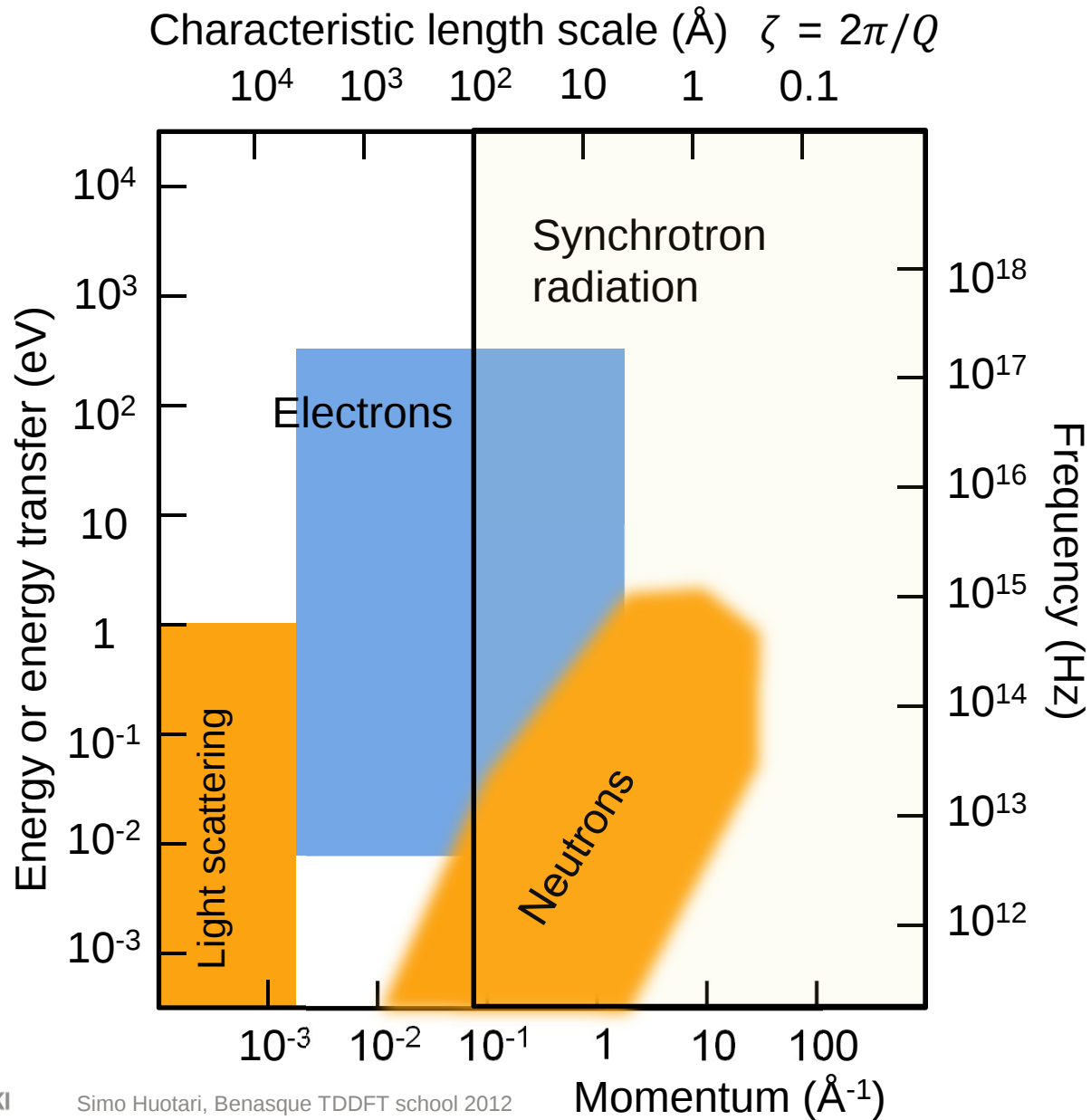


Particle probing depth





Dynamic ranges for scattering





Spectroscopy - interaction

Photon-electron

$$H = \frac{e^2}{2mc^2} \mathbf{A}^2 + \frac{e}{mc} \mathbf{A} \cdot \mathbf{p}$$



Scattering

Absorption and emission in first order, resonant scattering in second order

Electron-electron

$$V = e^2 / |\mathbf{r} - \mathbf{r}_i|$$

Transition rate from state $|A\rangle$ to state $|B\rangle$ is given by the Fermi's Golden Rule:

$$W_{BA} = \frac{2\pi}{\hbar} \left| \langle B | H | A \rangle + \sum_I \frac{\langle B | H | I \rangle \langle I | H | A \rangle}{E_A - E_I + i\Gamma/2} \right|^2 \delta(E_B - E_A) + O(H^3)$$



Kramers-Heisenberg formula

From $\mathbf{A} \cdot \mathbf{A}$ - Non-resonant scattering (Raman, inelastic x-ray,)

$$\begin{aligned} \frac{d^2\sigma}{d\Omega d\omega_2} = & r_0^2 \left(\frac{\omega_2}{\omega_1} \right) \left| \langle \Psi_f | e^{i\mathbf{Q} \cdot \mathbf{r}} | \Psi_i \rangle (\boldsymbol{\epsilon}_1 \cdot \boldsymbol{\epsilon}_2) \right. \\ & + \frac{1}{m} \sum_n \left[\frac{\langle \Psi_f | \boldsymbol{\epsilon}_2 \cdot \mathbf{p} e^{-i\mathbf{k}_2 \cdot \mathbf{r}} | \Psi_n \rangle \langle \Psi_n | \boldsymbol{\epsilon}_1 \cdot \mathbf{p} e^{i\mathbf{k}_1 \cdot \mathbf{r}} | \Psi_i \rangle}{E_i - E_n + \omega_1 + i\Gamma_n} \right. \\ & \left. \left. + \frac{\langle \Psi_f | \boldsymbol{\epsilon}_1 \cdot \mathbf{p} e^{i\mathbf{k}_1 \cdot \mathbf{r}} | \Psi_n \rangle \langle \Psi_n | \boldsymbol{\epsilon}_2 \cdot \mathbf{p} e^{-i\mathbf{k}_2 \cdot \mathbf{r}} | \Psi_i \rangle}{E_i - E_n - \omega_2} \right] \right|^2 \delta(E_i - E_f + \omega) \end{aligned}$$

From $\mathbf{p} \cdot \mathbf{A}$ - Resonant scattering (Resonant Raman, resonant inelastic x-ray, fluorescence....)



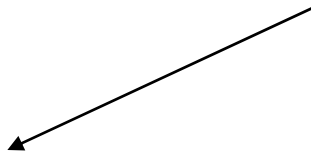
Dynamic structure factor

Average density-density correlation function

$$G(\mathbf{r}, t) = \frac{1}{N} \int \langle \rho(\mathbf{r}' - \mathbf{r}, t) \rho(\mathbf{r}', 0) \rangle d\mathbf{r}'$$

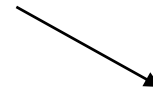
Dynamic structure factor

$$S(\mathbf{Q}, E) = \int \int G(\mathbf{r}, t) e^{i\mathbf{Q} \cdot \mathbf{r}} e^{-iEt}$$



Small Q

interference effects are important
(collective excitations)



Large Q

independent particle picture
(Compton scattering)



Absorption and scattering

Absorption	$P(\omega) \propto \sum_f \langle \Psi_f \mathbf{p} \cdot \mathbf{A} \Psi_i \rangle ^2 \delta(E_i - E_f + \omega)$	Dipole selection rule
Scattering	$S(\mathbf{Q}, \omega) \propto \sum_f \langle \Psi_f \sum_j e^{i \mathbf{Q} \cdot \mathbf{r}_j} \Psi_i \rangle ^2 \delta(E_i - E_f + \omega)$	

Use the momentum transfer dependence of the scattering matrix element: as $q \rightarrow 0$ then

$$e^{i \mathbf{Q} \cdot \mathbf{r}} = 1 + i \mathbf{Q} \cdot \mathbf{r} - (\mathbf{Q} \cdot \mathbf{r})^2 / 2 + \dots$$

↑
dipole

↑
Higher order
multipoles



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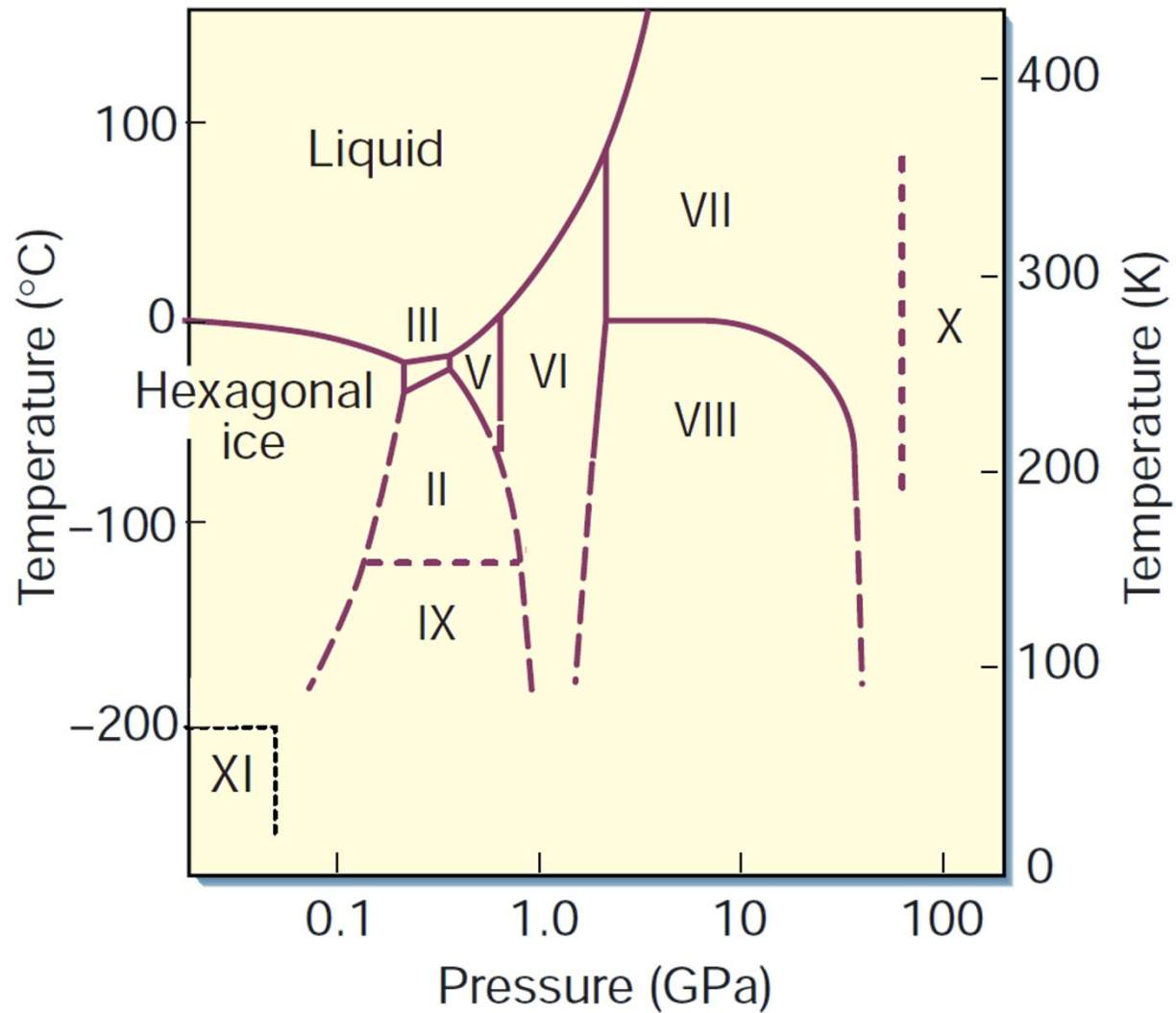
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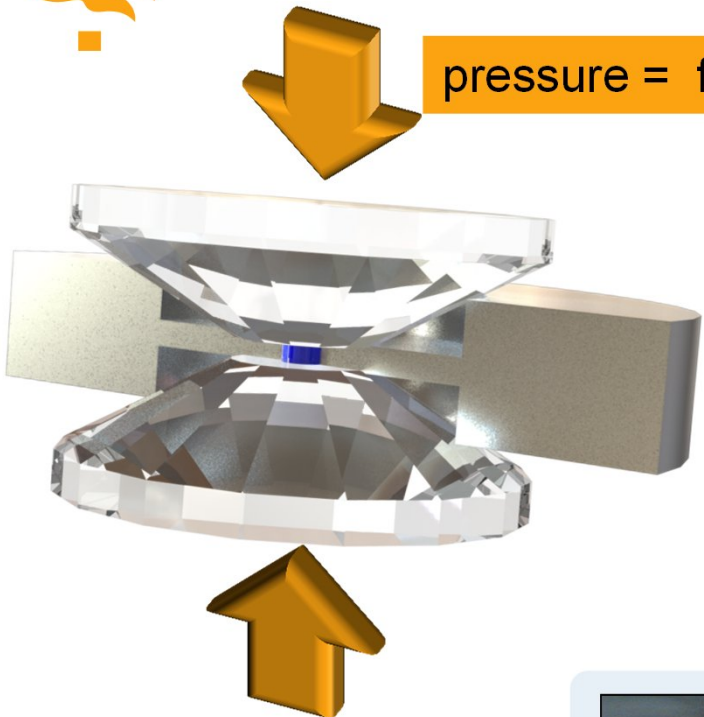


Phase diagram of water





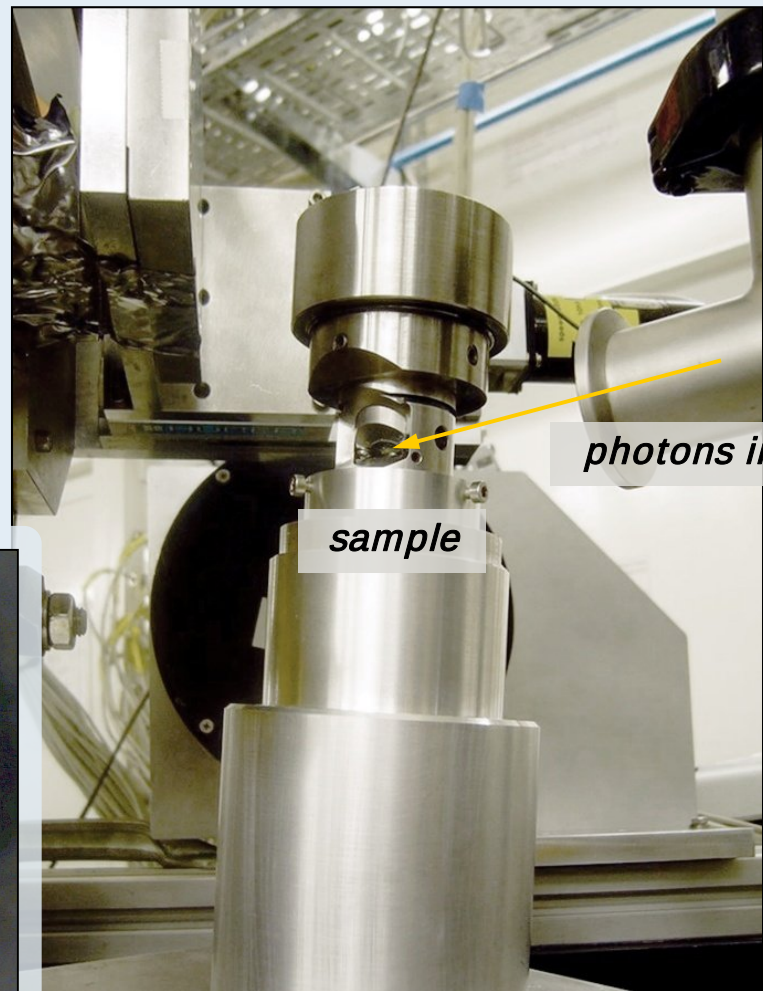
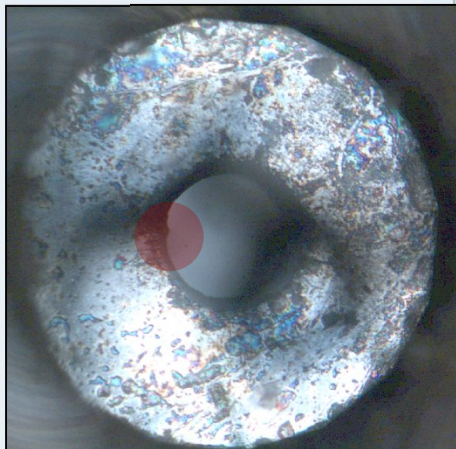
Extremely high pressures


$$\text{pressure} = \text{force} / \text{area}$$

V.M.Giordano, T. Pylkkänen et al. ESRF ID16

Experimental details

- 1 mm diamond tip (culet)
- $\varnothing$ 5 mm Be gasket
- $\varnothing$ 350 micron sample size
- X-ray beam $100 \times 50 \mu\text{m}^2$
- Ruby chip for P calibration



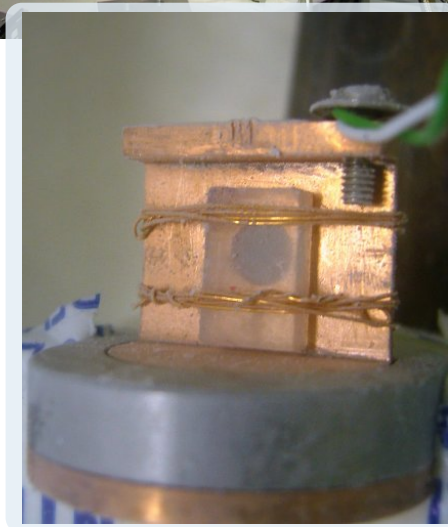


Low temperatures

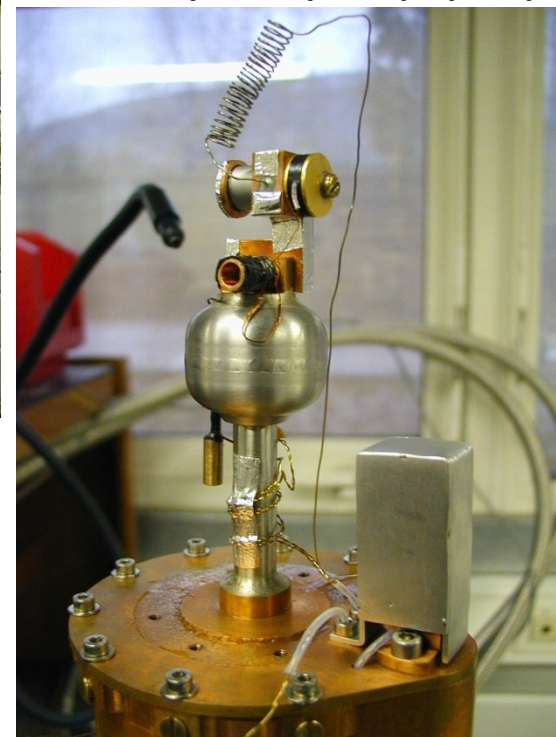
Liquid nitrogen cryo-jet (77 K)



Liquid He Cryostat (4 K)



He sorption pump (1 K)



F. Albergamo et al., ESRF

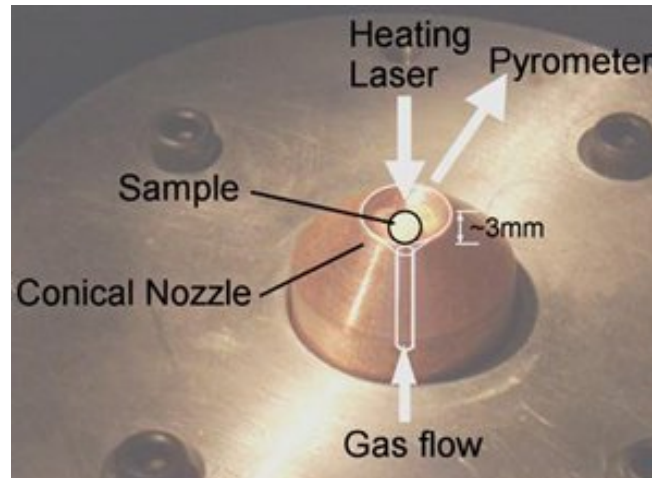


Extremely high temperatures (thousands of deg): Laser heating and aerodynamic levitation

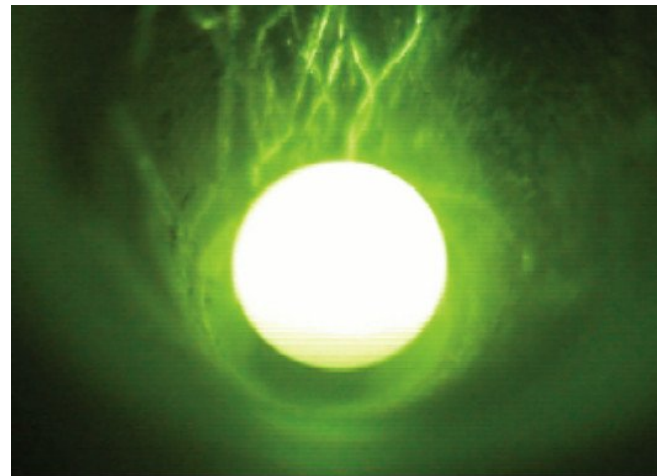
High temperatures
(up to 1000 deg C):
furnaces, hot air guns



ESRF, Sample group



Wikipedia:
Aerodynamic
levitation



Droplet of liquid basalt BCR-2 during levitation. The sample was heated from the top using a CO₂-laser. The diameter of the sphere was ~2 mm. A. Pack et al., Geochemical transactions 2010, 11:4



Modern 3rd generation synchrotrons

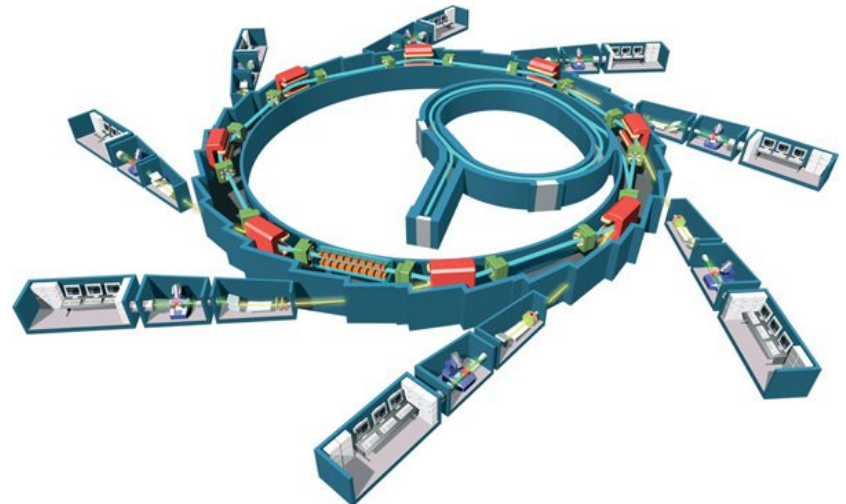
Advanced Photon Source, USA



Super Photon Ring 8 (Spring-8), Japan



European Synchrotron Radiation Facility, France





Moore's law

