



Probing electron excitations with inelastic x-ray scattering spectroscopies

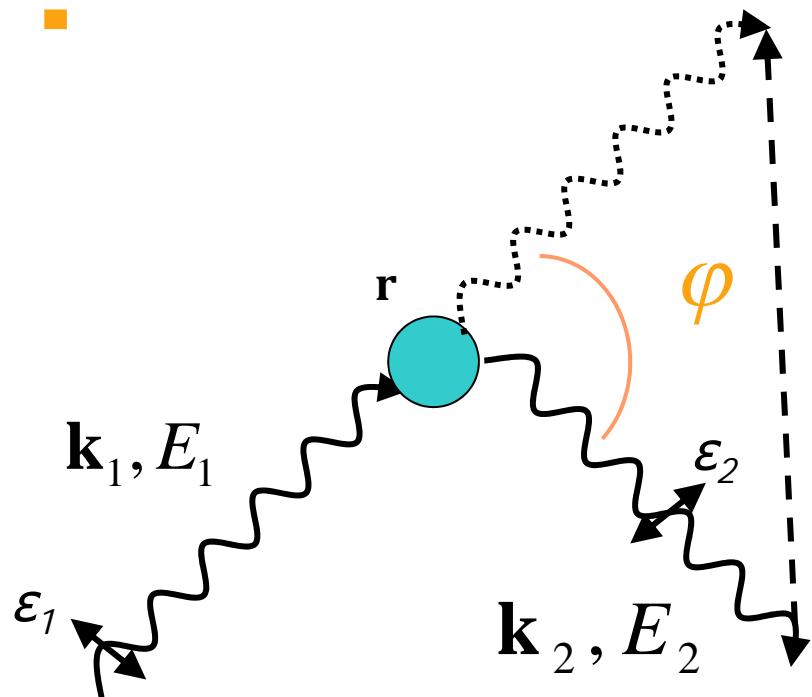
Simo Huotari

University of Helsinki, Finland

Benasque TDDFT workshop January 2012



Inelastic x-ray scattering



Dynamic structure factor $S(Q,E)$

$$S(\mathbf{Q}, E) \propto \sum_f |\langle \Psi_f | e^{i \mathbf{Q} \cdot \mathbf{r}} | \Psi_i \rangle|^2 \delta(E_i - E_f + E)$$

$$\mathbf{Q} = \mathbf{k}_1 - \mathbf{k}_2$$

$$S(\mathbf{Q}, E) = \frac{Q^2}{4\pi^2 n} \text{Im} \left[\frac{-1}{\varepsilon(\mathbf{Q}, E)} \right]$$

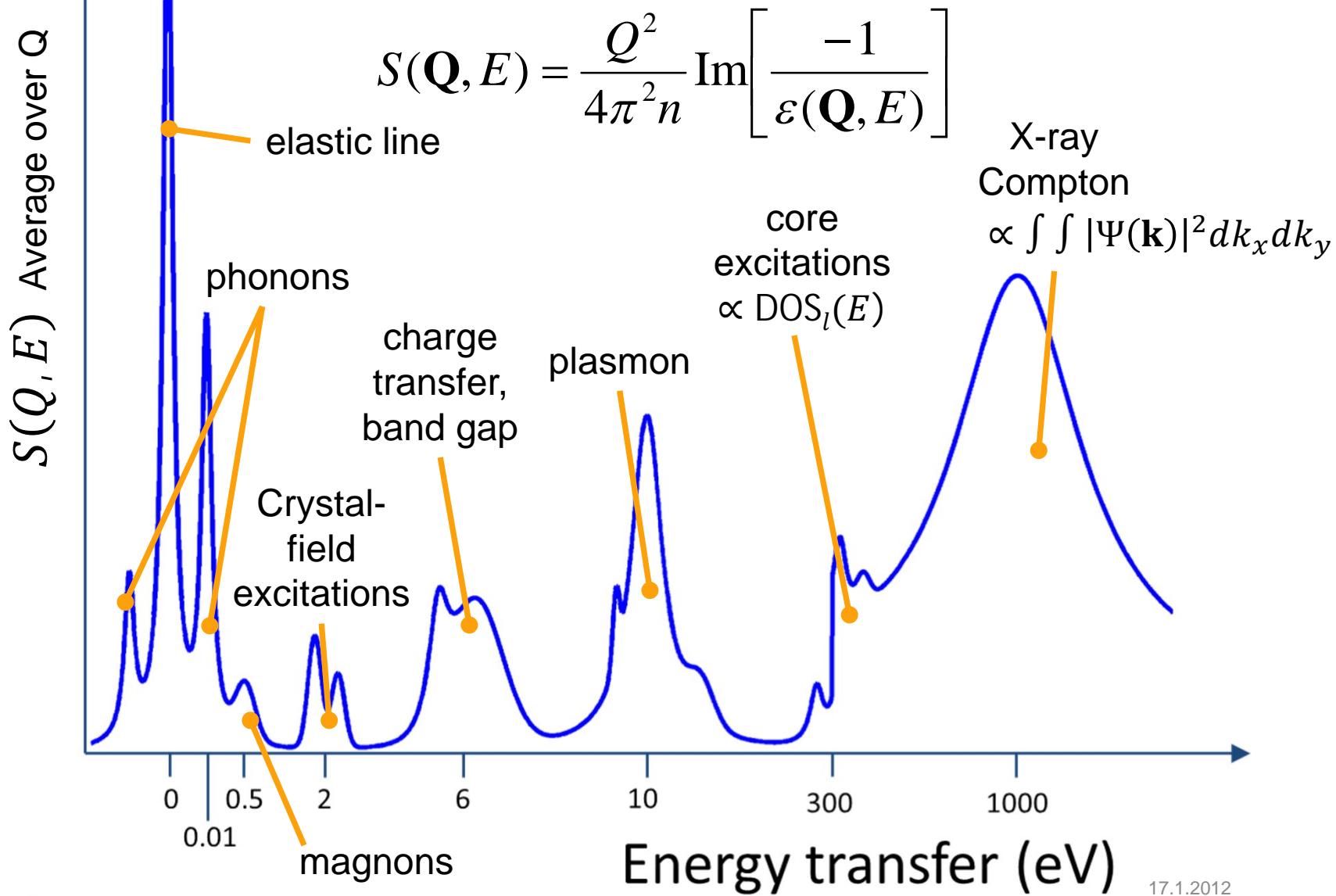
Momentum transfer

$$\mathbf{Q} = \mathbf{k}_1 - \mathbf{k}_2$$

Energy transfer $E = E_1 - E_2$



Dynamic structure factor





Outline

Part 1 Metal-to-insulator transition

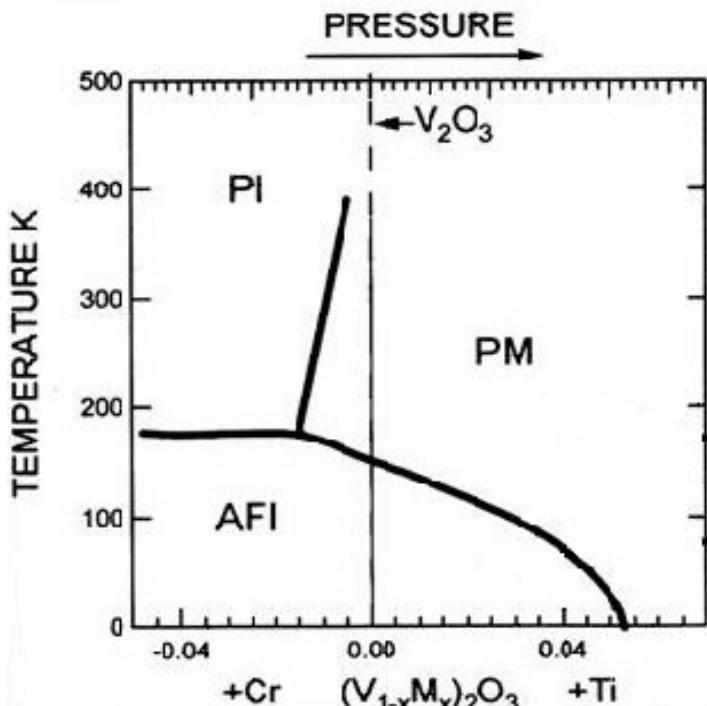
Part 2 Plasmons and e-h continuum

Part 3 Double plasmons

Part 4 Quasiparticle renormalisation

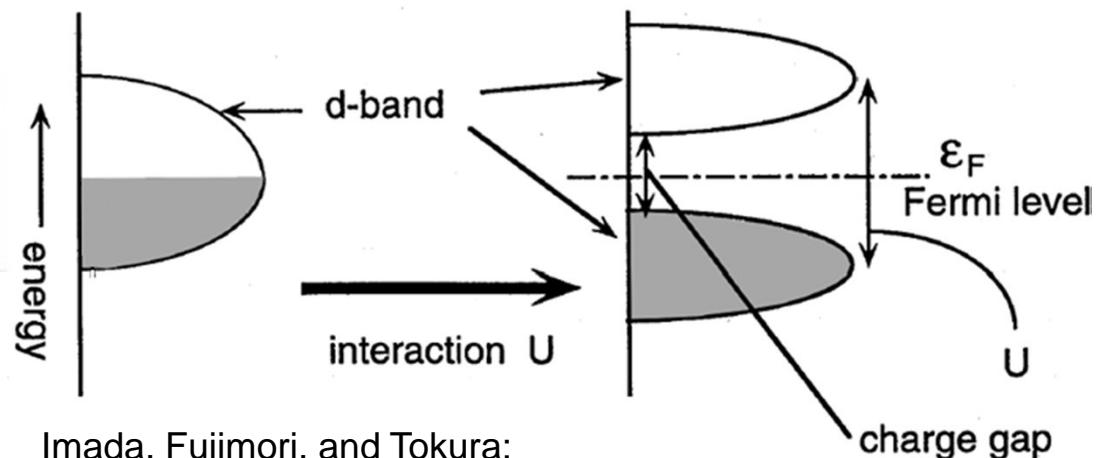


V_2O_3 and the Mott transition



McWhan et al., PRL 27 (1971)

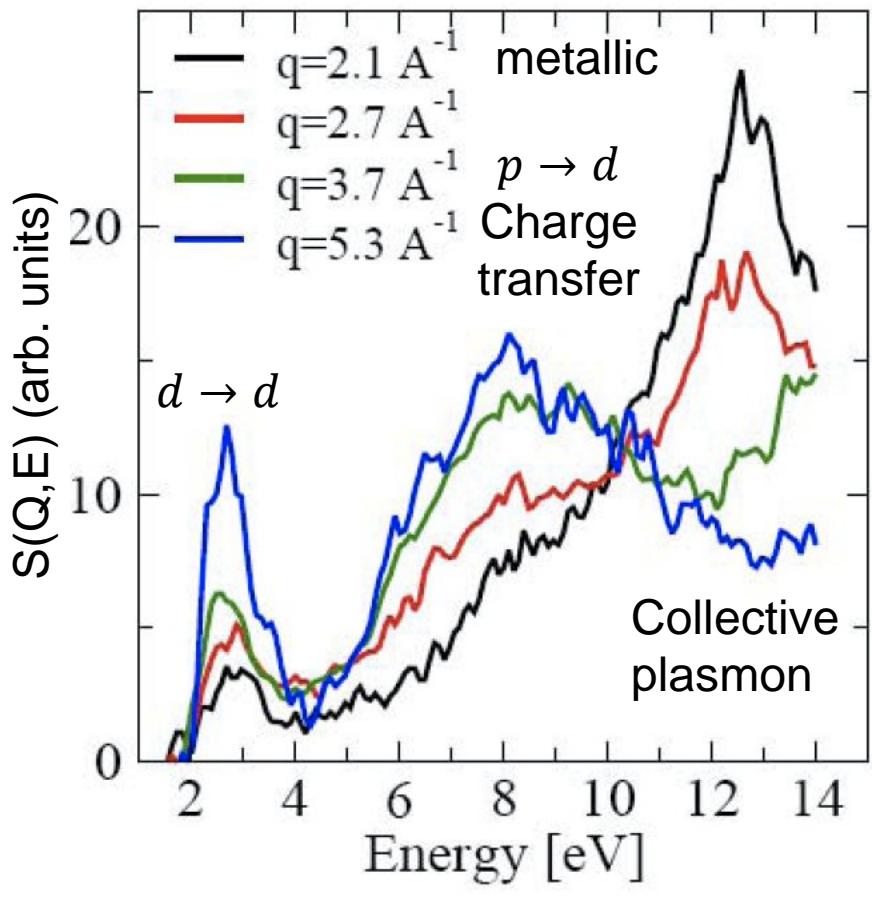
- V_2O_3 is one of the canonical model systems for a Mott-Hubbard insulator
- The Mott transition is usually explained within the Hubbard model with U = on-site Coulomb repulsion



Imada, Fujimori, and Tokura:
Metal-insulator transitions, RMP 70 1039 (1998)



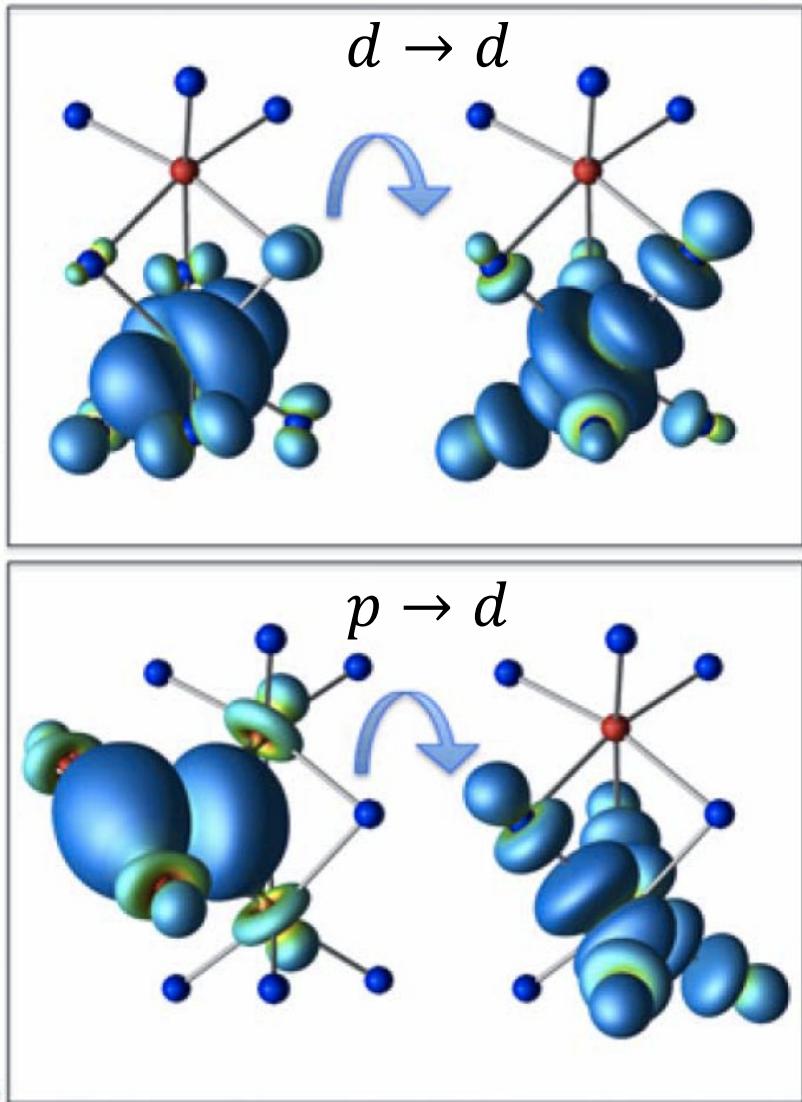
V₂O₃



$$e^{i\mathbf{Q}\cdot\mathbf{r}} = 1 + i\mathbf{Q} \cdot \mathbf{r} - (\mathbf{Q} \cdot \mathbf{r})^2/2 + \dots$$

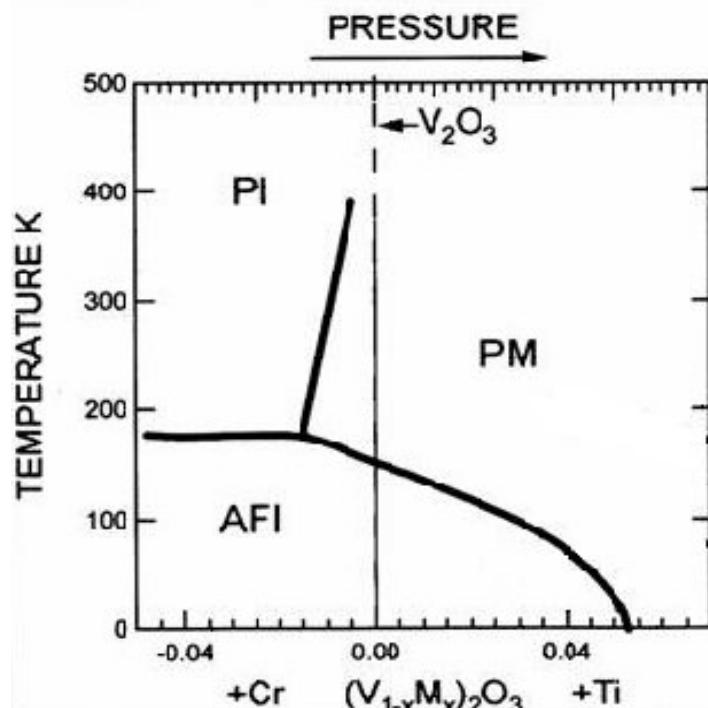
A.V.Kozhevnikov et al., in preparation

Simo Huotari, Benasque TDDFT workshop 2012



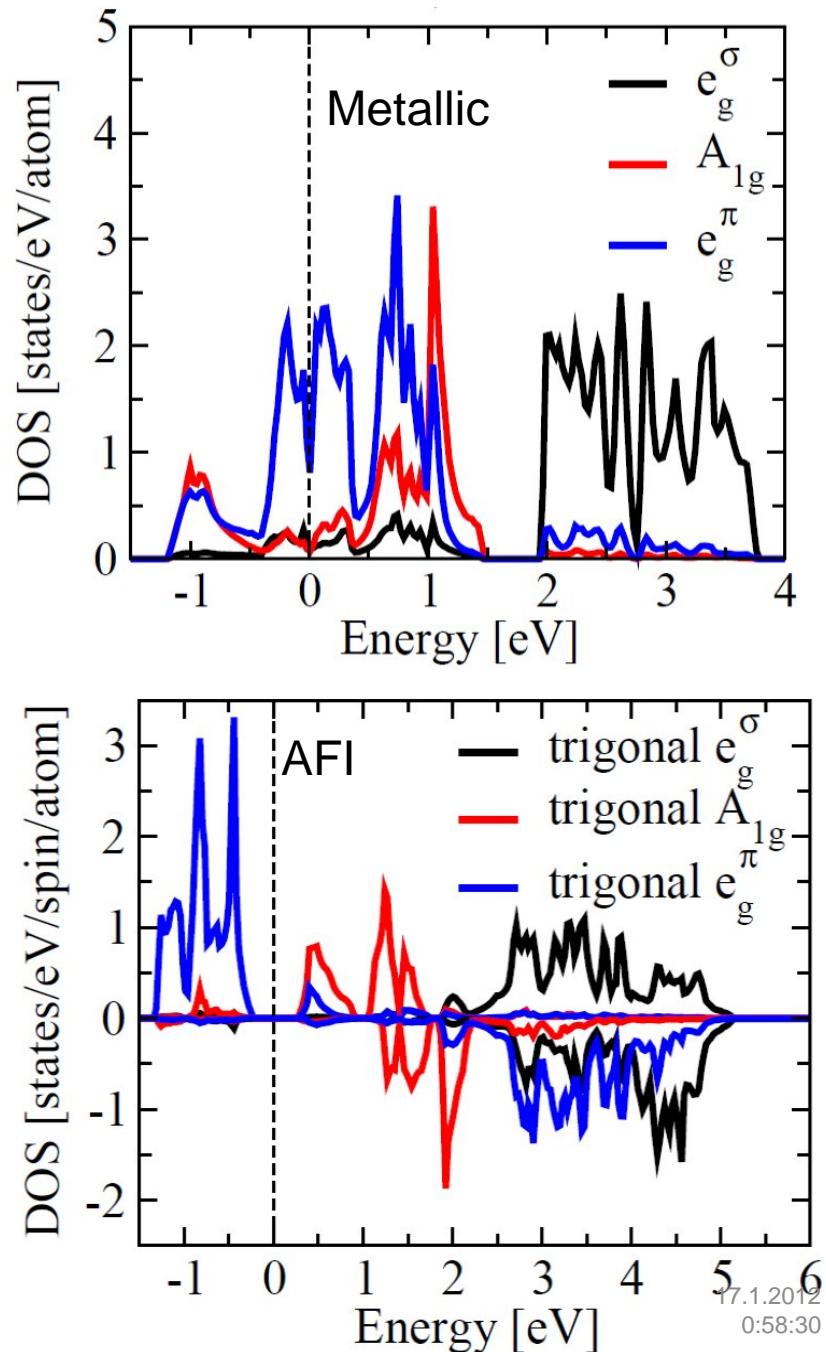


V_2O_3



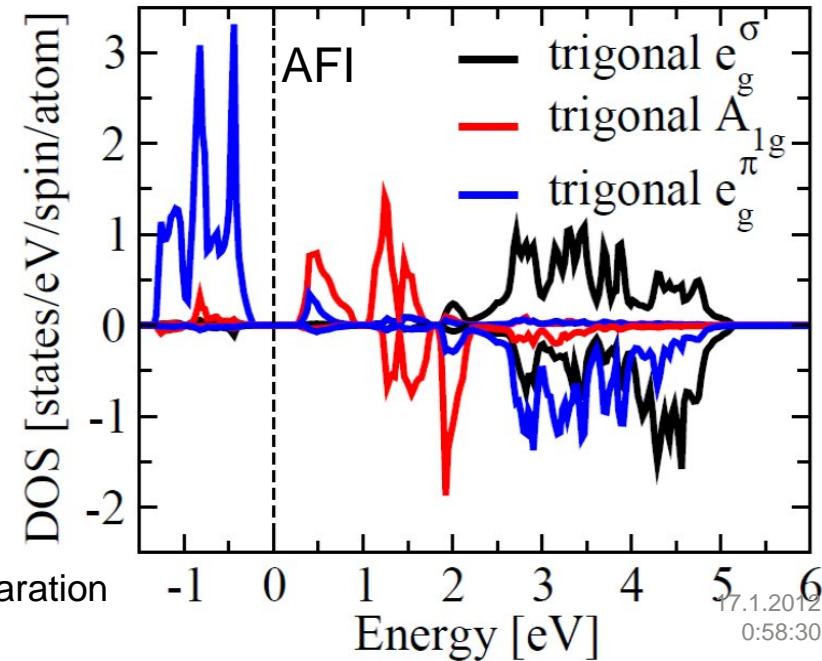
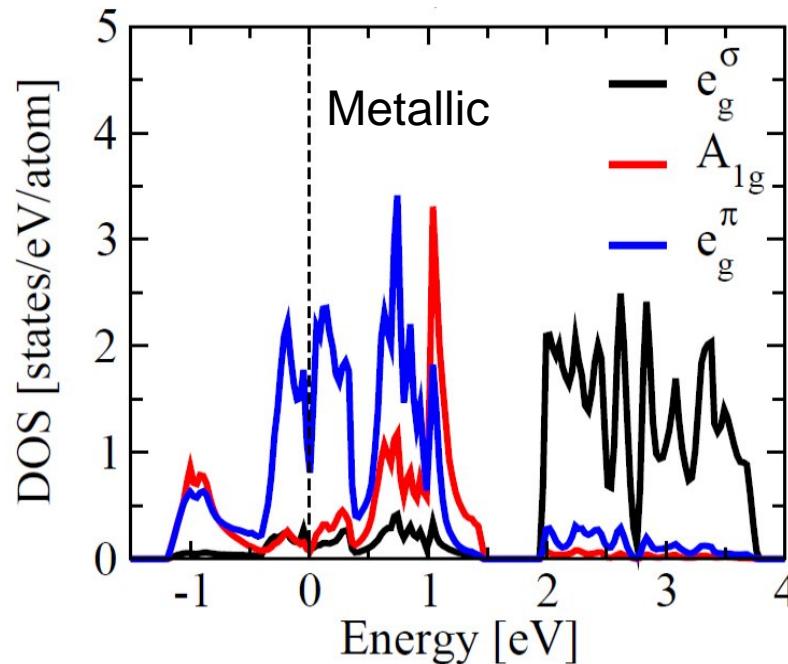
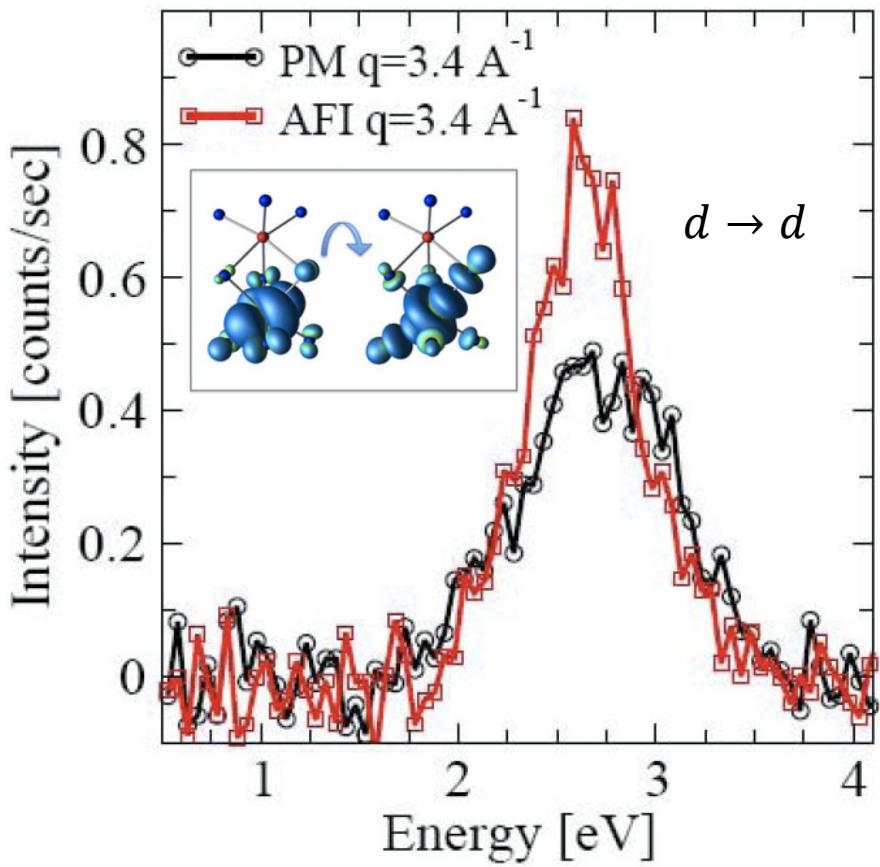
McWhan *et al.*, PRL 27 (1971)

A.V.Kozhevnikov *et al.*, in preparation





V_2O_3

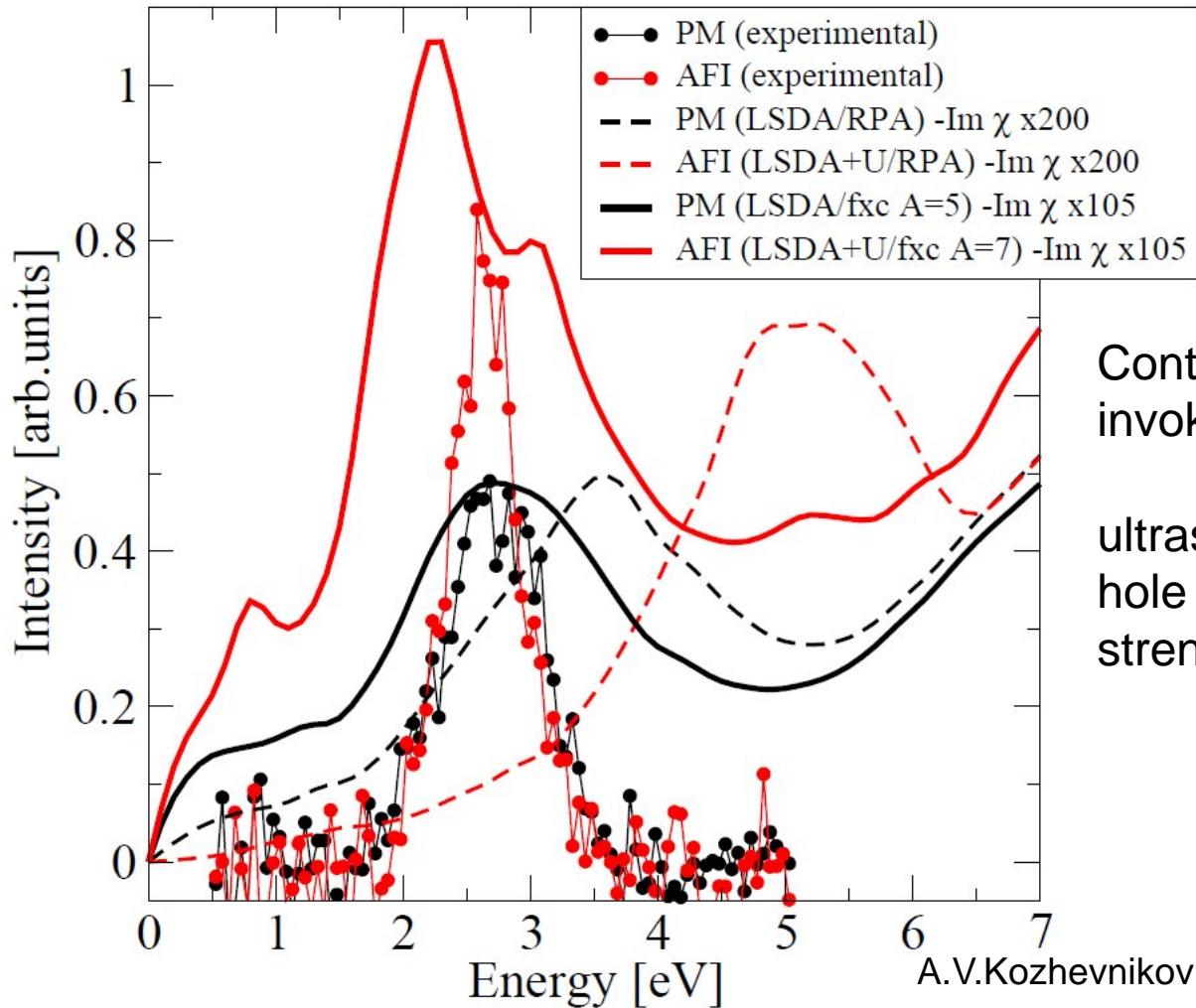


A.V.Kozhevnikov et al., in preparation

Simo Huotari, Benasque TDDFT workshop 2012



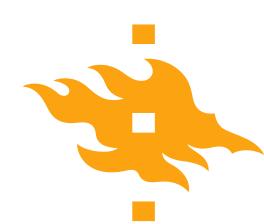
Metal-to-insulator transition in V_2O_3



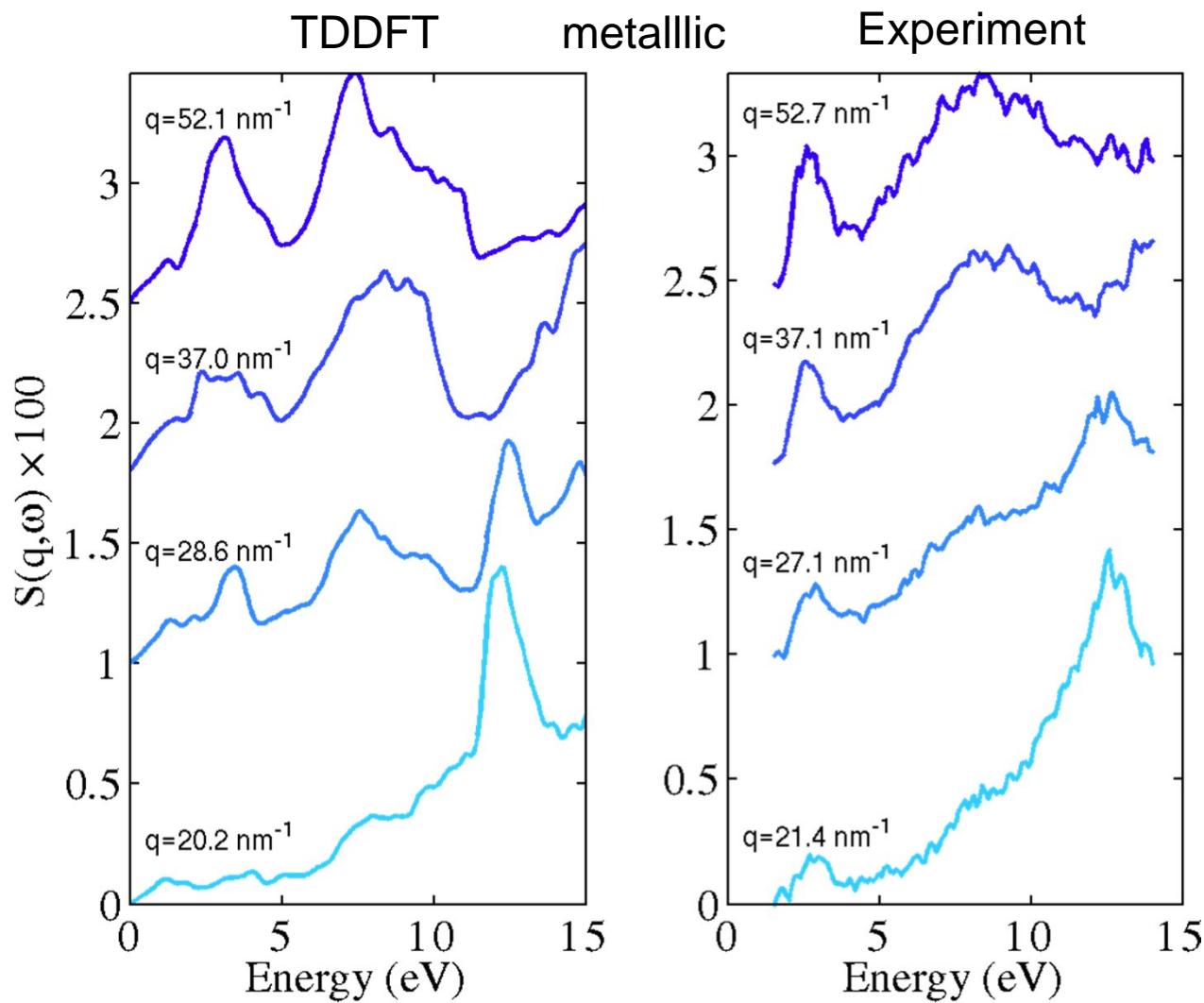
Contact exciton model invoked:

ultrashort-range electron-hole attraction, tunable strength with parameter A

A.V.Kozhevnikov et al., in preparation



V_2O_3



A.V.Kozhevnikov et al., in preparation

Simo Huotari, Benasque TDDFT workshop 2012

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Metal-to-insulator transition in V₂O₃

Conclusions for metal-insulator transition of V₂O₃:

- ❑ d-d excitation changes surprisingly little upon going to insulating phase
- ❑ RPA would predict much larger change
- ❑ Must imply a switching-on of a excitonic interaction upon the MIT?

A.V. Kozhevnikov, M.C. Troparevsky, T. C. Schulthess, A.G. Eguiluz,
T. Pylkkänen, L. Simonelli, G. Monaco, and S. Huotari, under preparation



Outline

Part 1 Metal-to-insulator transition

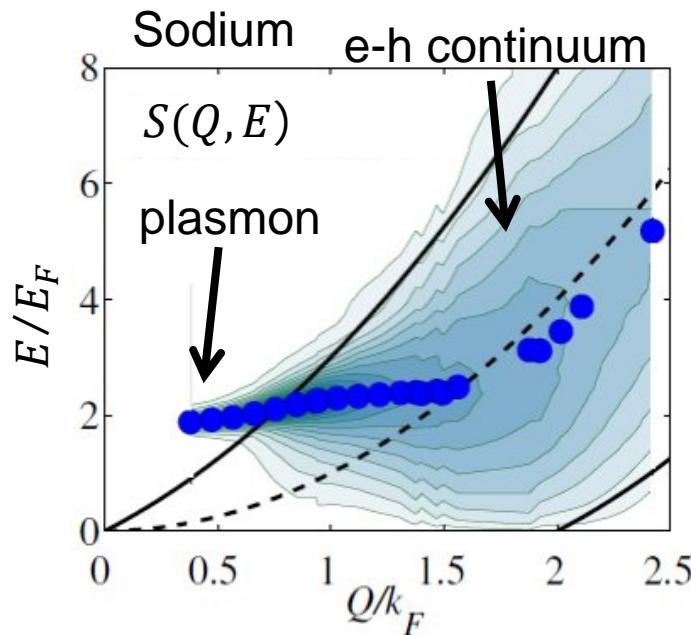
Part 2 Plasmons

Part 3 Double plasmons

Part 4 Quasiparticle renormalisation

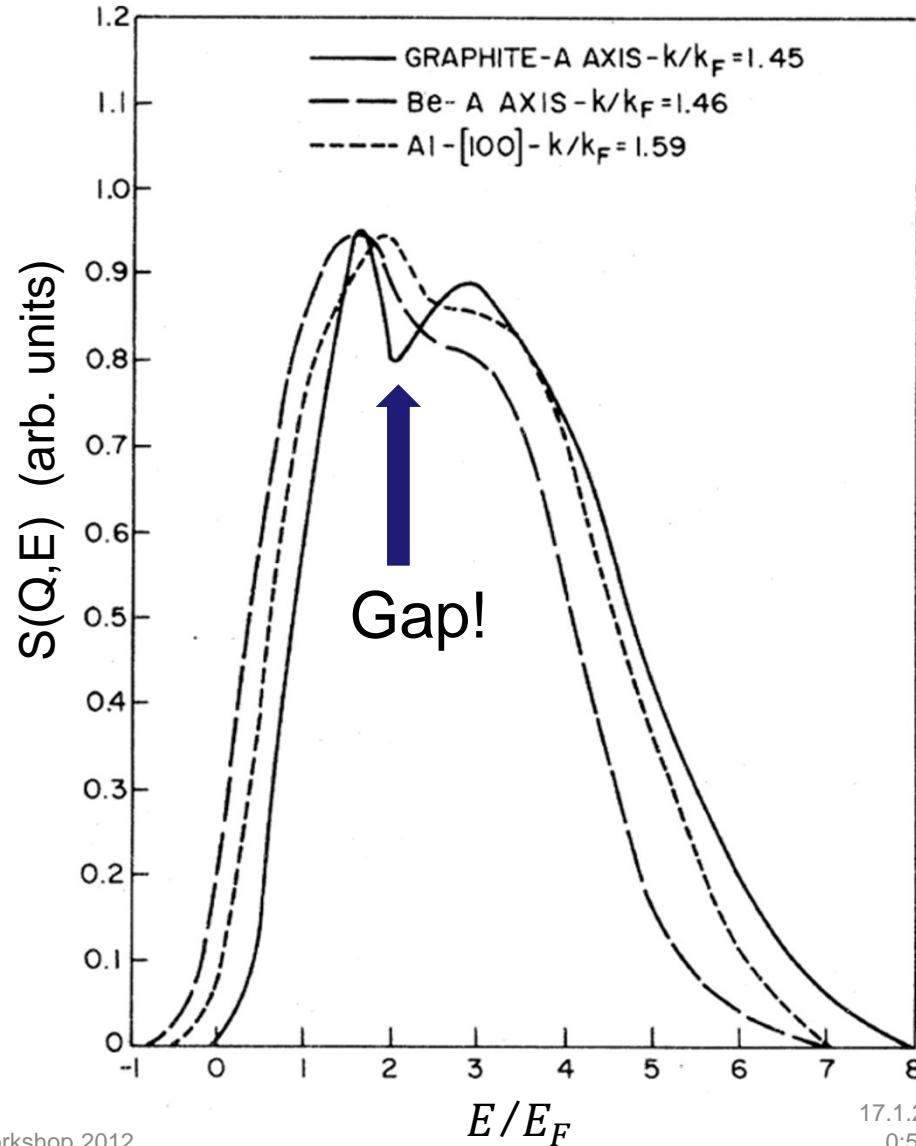


Electron gas response



In 1974, $S(Q, E)$ for many materials seemed to suggest a "universal shape" of the electron gas response function

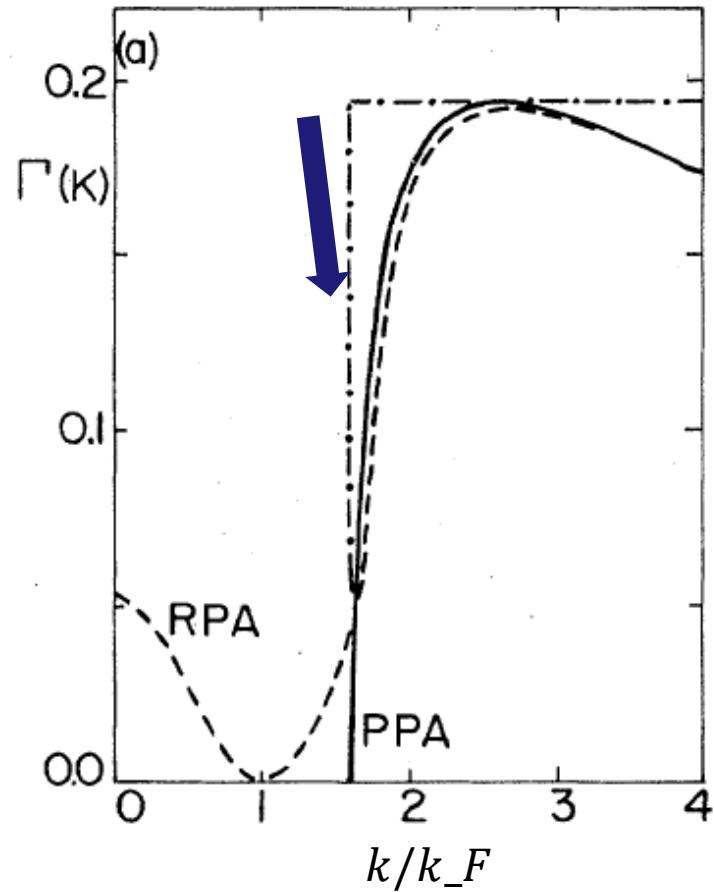
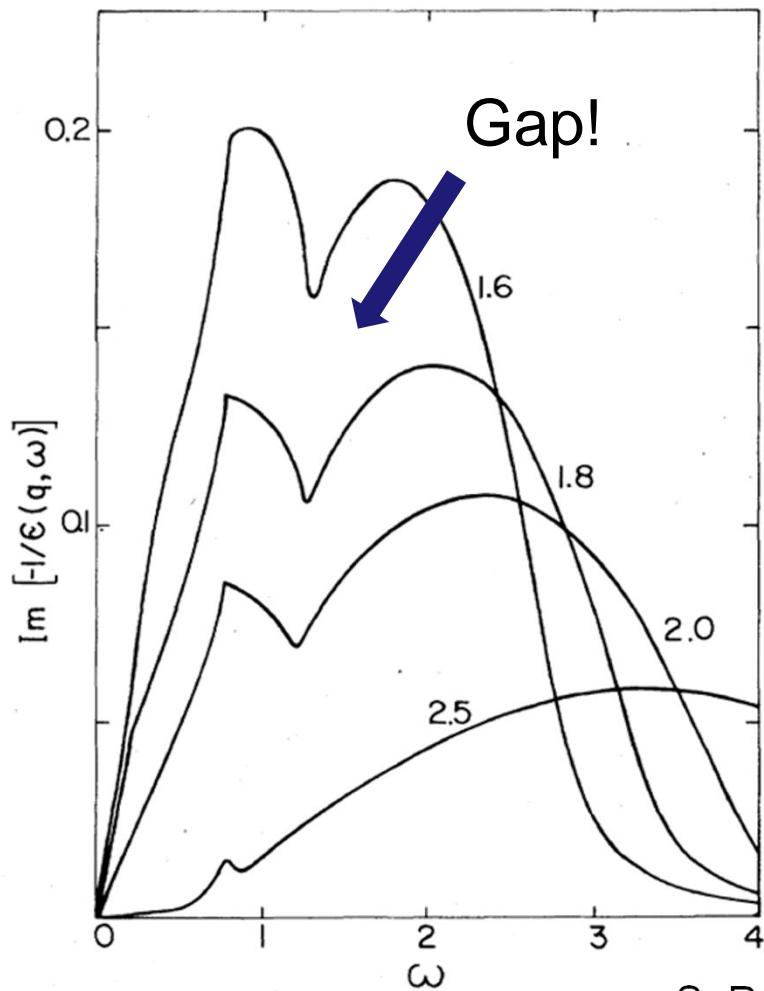
P.M. Platzman and P. Eisenberger,
PRL 33, 152 (1974)





Theory for electron gas

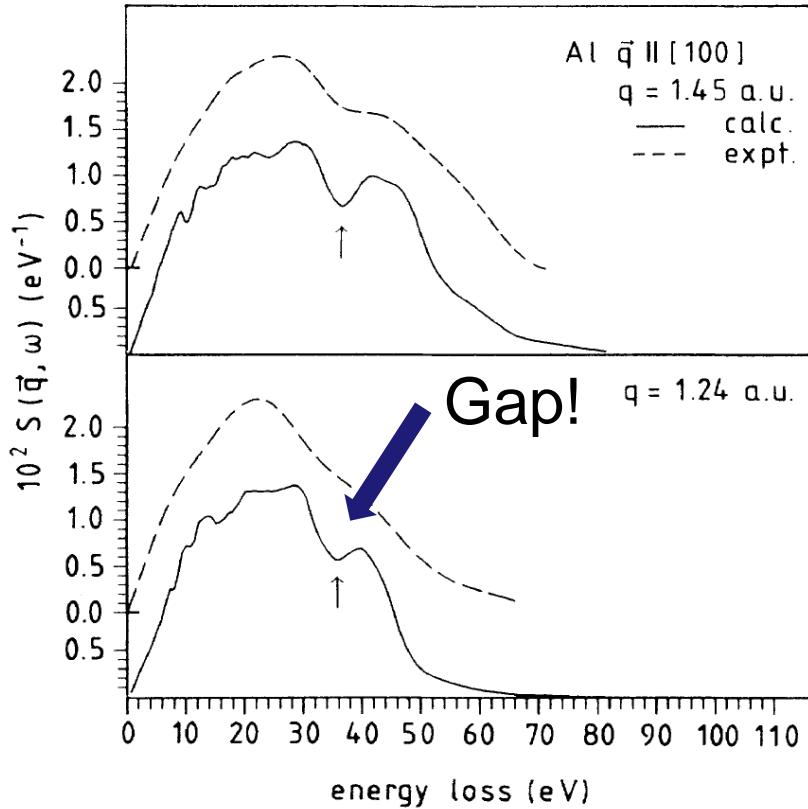
Gap structure appears when taking into account quasiparticle lifetimes



S. Rahman and G. Vignale, PRB 30, 6951 (1984)

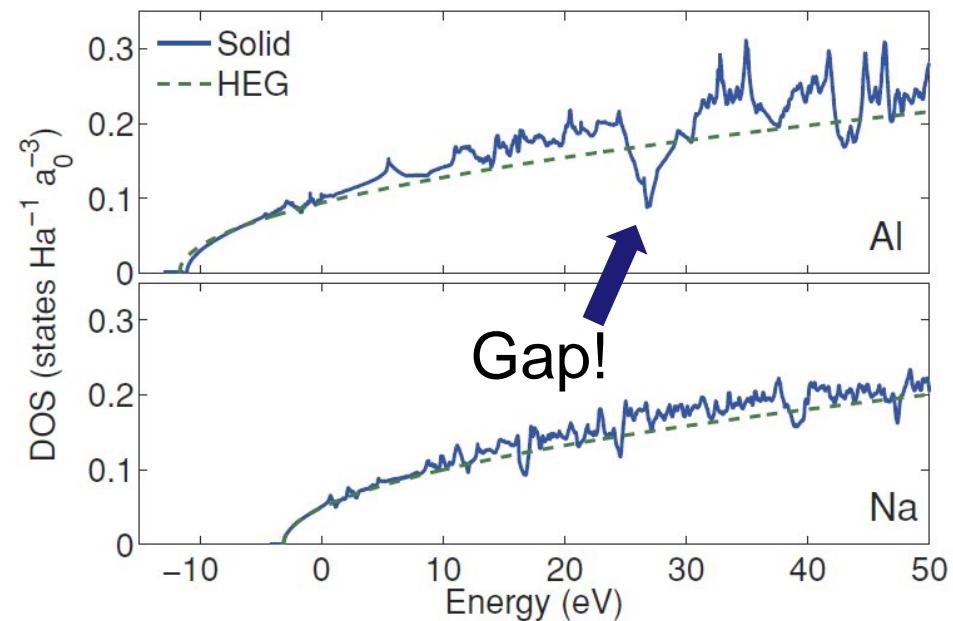


Gap-structure in Al



Winfried Schülke et al. PRB 47,
12426 (1993)

Explained to be a band-structure effect: a gap in the DOS

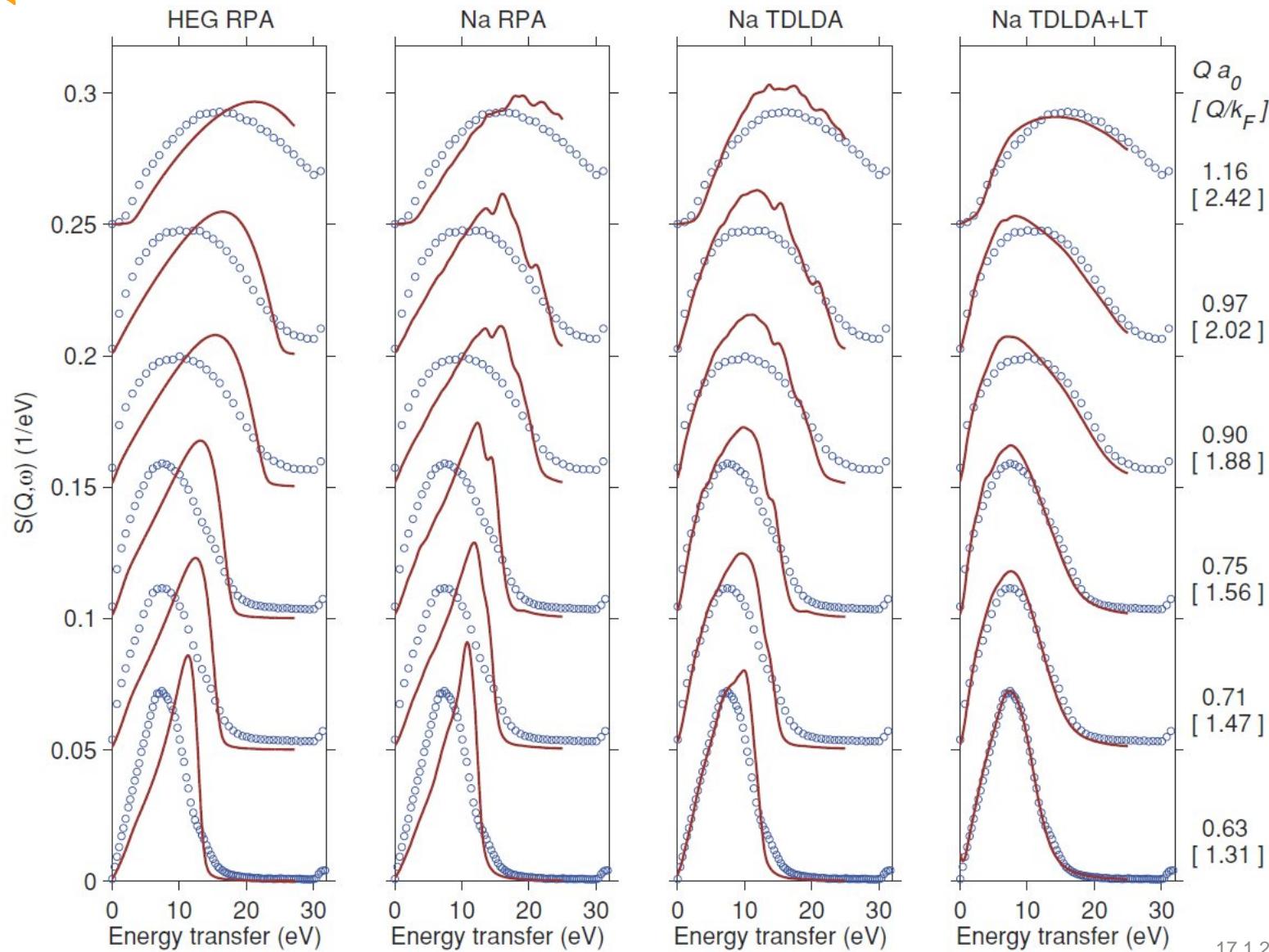


S. Huotari, M. Cazzaniga et al.,
PRB 84, 075108 (2011)



Results for sodium

S. Huotari et al., PRB
84, 075108 (2011)

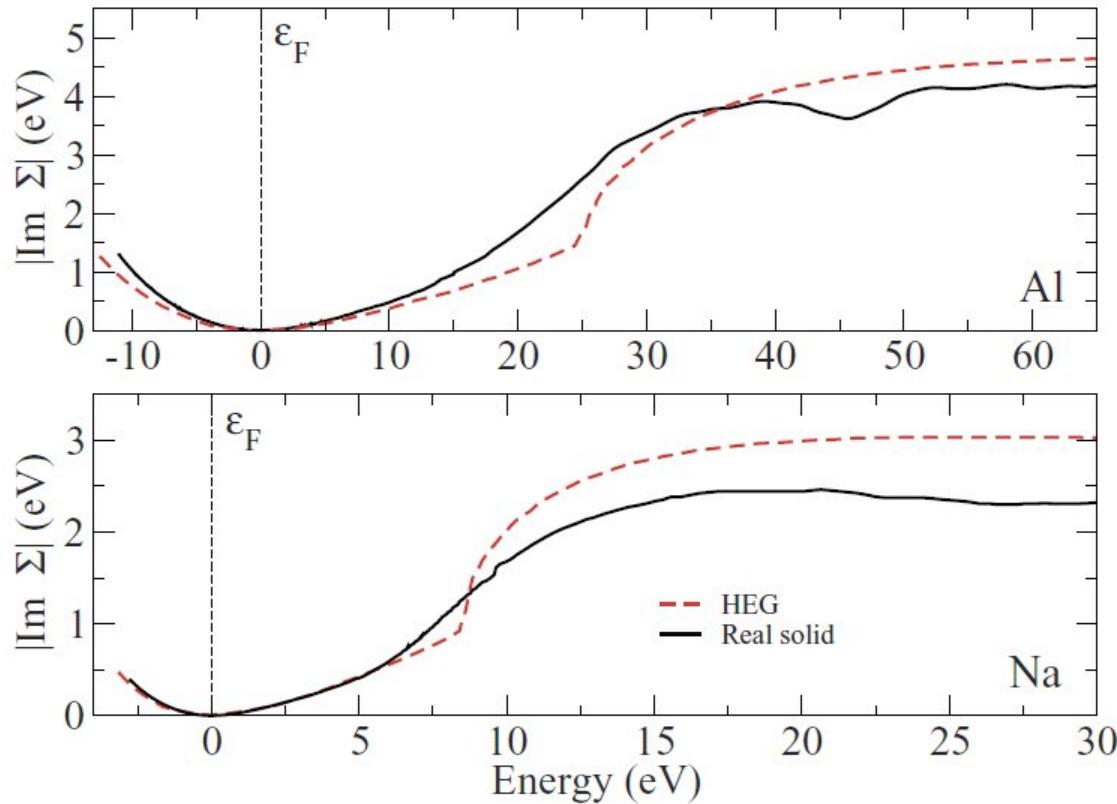




TDLDA+LifeTimes (LT)

H.-C. Weissker et al., Phys. Rev. Lett. 97, 237602 (2006)
M. Cazzaniga et al., Phys. Rev. B 84, 075109 (2011)

$$\chi_{\mathbf{G}, \mathbf{G}'}^0(\mathbf{q}, \omega) = -\frac{1}{V_{\text{BZ}}} \sum_{j, j'} \int_{\text{BZ}} d^3k [f(\epsilon_{j'}(\mathbf{k} + \mathbf{q})) - f(\epsilon_j(\mathbf{k}))] \frac{\langle \mathbf{k}, j | e^{-i(\mathbf{q} + \mathbf{G}) \cdot \hat{\mathbf{r}}} | \mathbf{k} + \mathbf{q}, j' \rangle \langle \mathbf{k} + \mathbf{q}, j' | e^{i(\mathbf{q} + \mathbf{G}') \cdot \hat{\mathbf{r}}} | \mathbf{k}, j \rangle}{\omega - [\epsilon_{j'}(\mathbf{k} + \mathbf{q}) - \epsilon_j(\mathbf{k})] + i\eta}$$



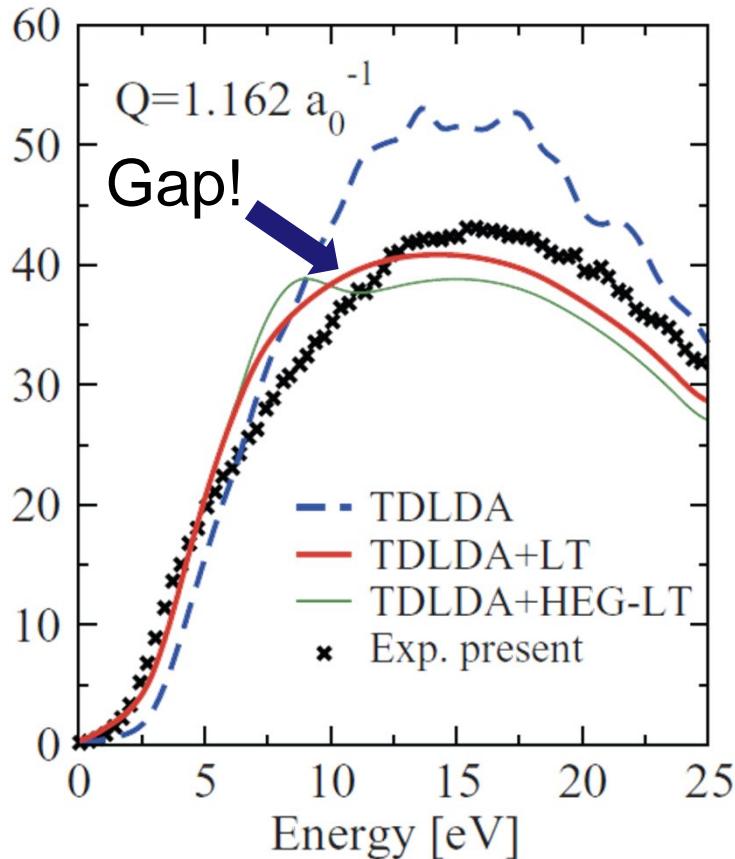
$i|\text{Im } \Sigma_{j'}(\mathbf{k} + \mathbf{q})| + i|\text{Im } \Sigma_j(\mathbf{k})|$

Inclusion of lifetimes
via a modified
independent particle
polarizability χ_0^{LT}



Results for sodium

M. Cazzaniga et al., PRB **84**, 075109 (2011)



Using HEG lifetimes to HEG or real-metal Na, does a gap structure in the $S(Q, E)$.

The gap is washed out if real-metal lifetimes are used.

Fictitious HEG does in fact keep the gap structure but Na is not close enough to the HEG!



Plasmons

Conclusions about plasmons in Na:

- Correlation effects beyond RPA surprisingly large in Na
- Quasiparticle lifetimes have to be taken into account properly
- Fine structure to the $S(Q,E)$ expected in HEG; washed out in Na

PHYSICAL REVIEW B **84**, 075108 (2011)

Dynamical response function in sodium studied by inelastic x-ray scattering spectroscopy

Simo Huotari,^{1,2,*} Marco Cazzaniga,^{3,4} Hans-Christian Weissker,^{4,5,6} Tuomas Pylkkänen,^{1,2} Harald Müller,² Lucia Reining,^{4,5} Giovanni Onida,^{3,4} and Giulio Monaco²

PHYSICAL REVIEW B **84**, 075109 (2011)

Dynamical response function in sodium and aluminum from time-dependent density-functional theory

Marco Cazzaniga,^{1,2,*} Hans-Christian Weissker,^{2,3,4} Simo Huotari,^{5,6} Tuomas Pylkkänen,^{5,6} Paolo Salvestrini,^{1,7} Giulio Monaco,⁵ Giovanni Onida,^{1,2} and Lucia Reining^{2,3}

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Outline

Part 1 Metal-to-insulator transition

Part 2 Plasmons

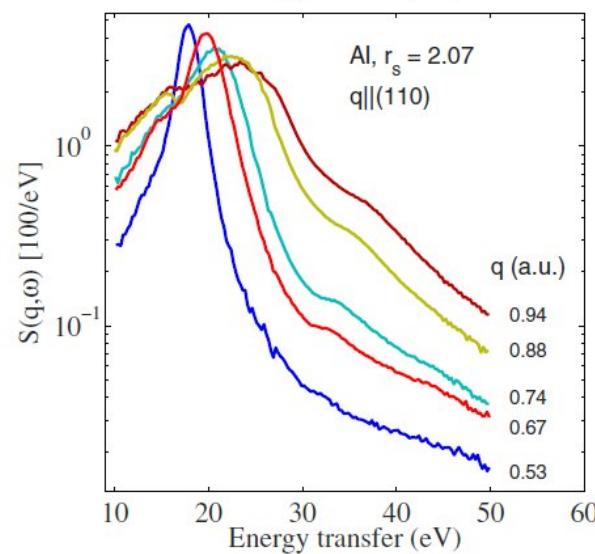
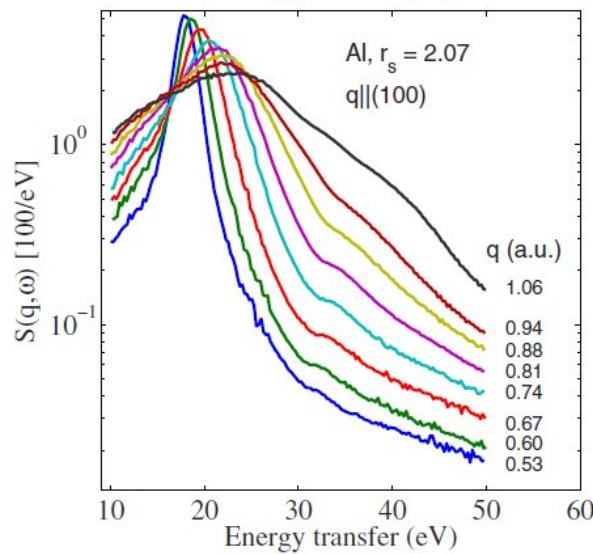
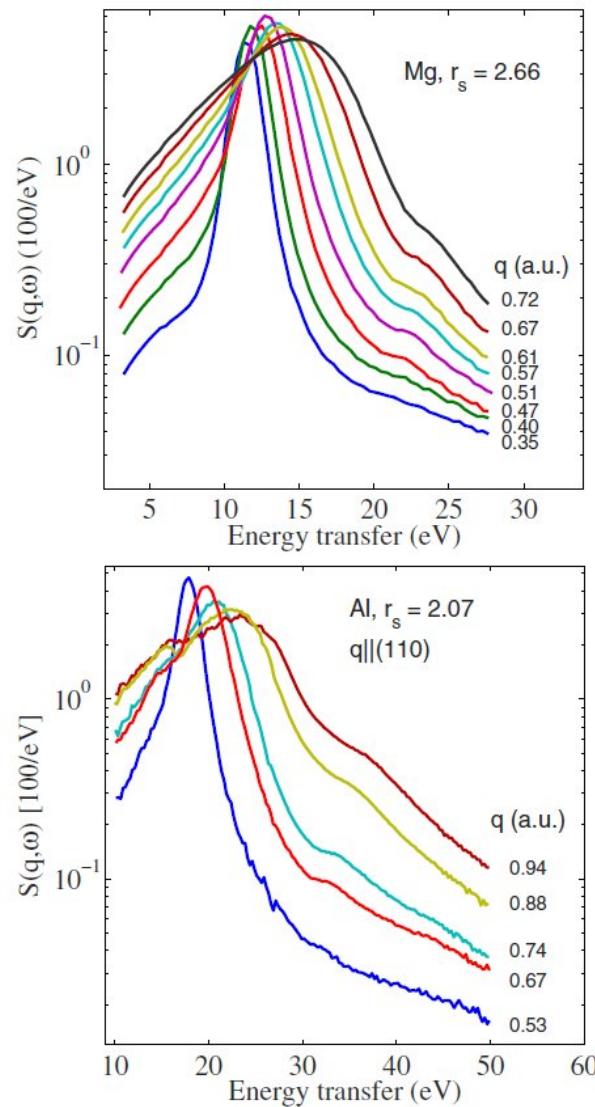
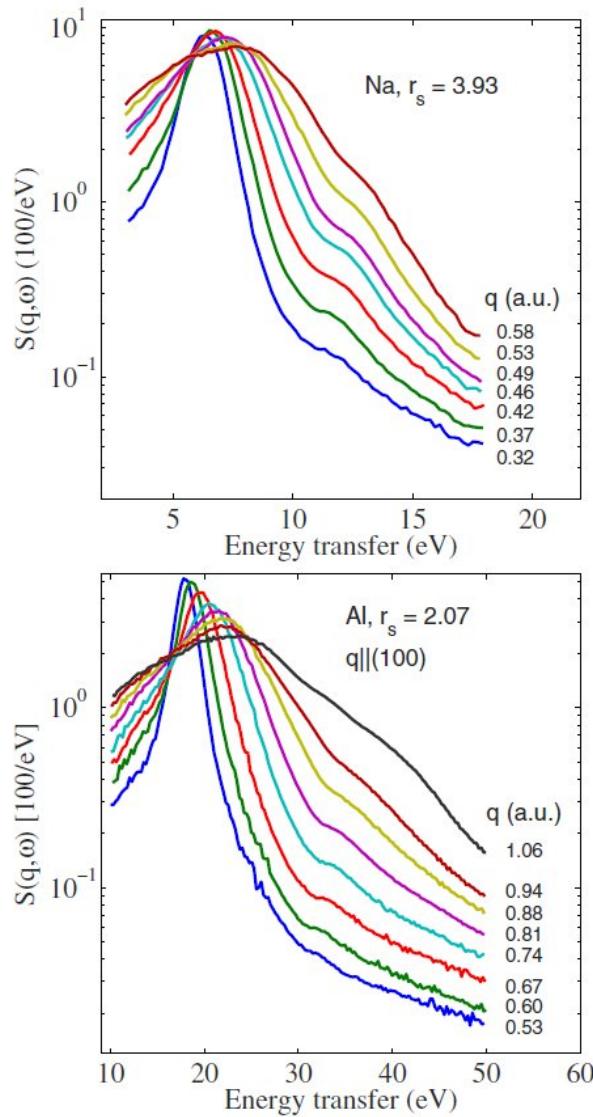
Part 3 Double plasmons

Part 4 Quasiparticle renormalisation



Double plasmons

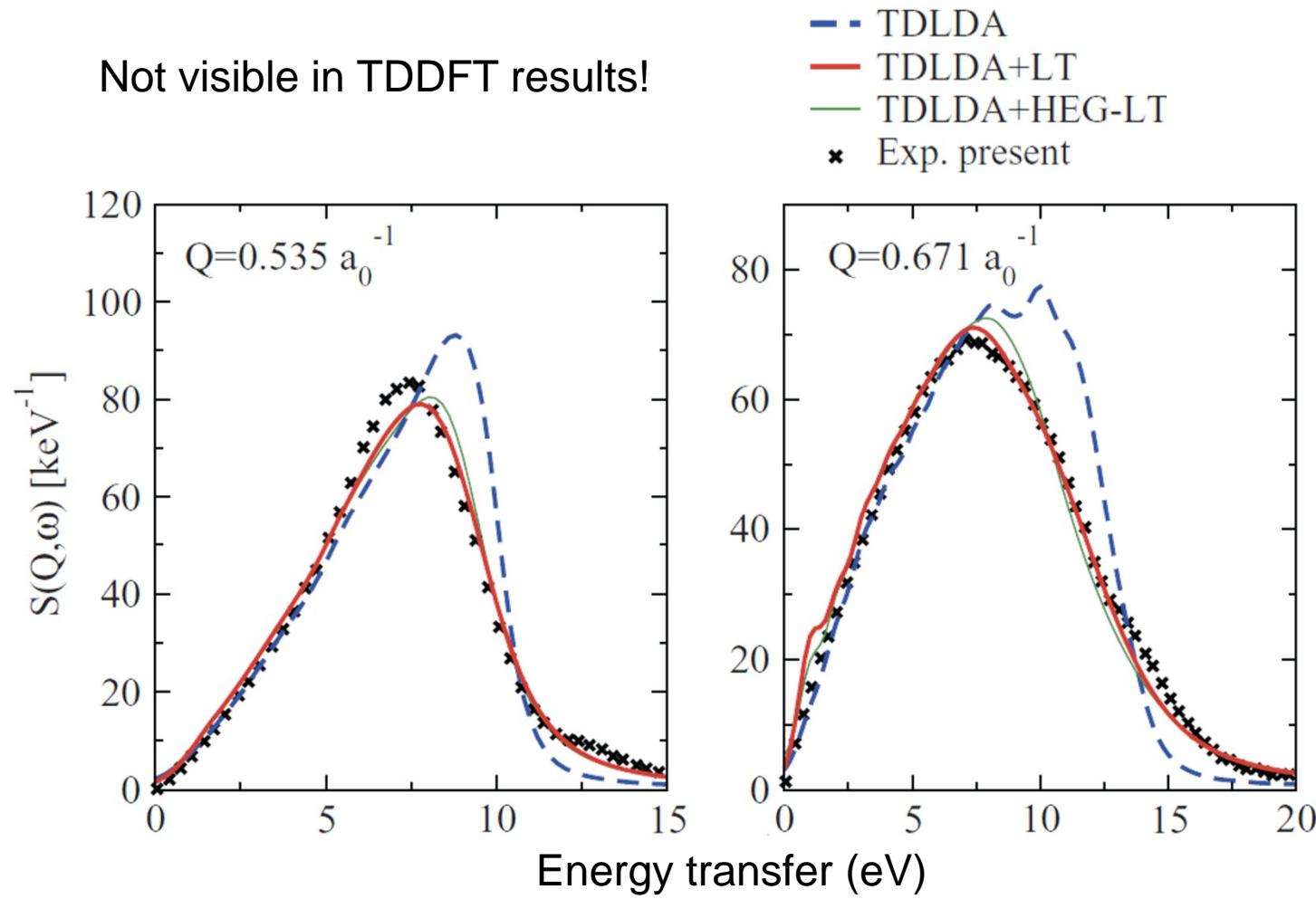
S. Huotari et al.,
PRB 77, 195125 (2008)





Double plasmons

Not visible in TDDFT results!



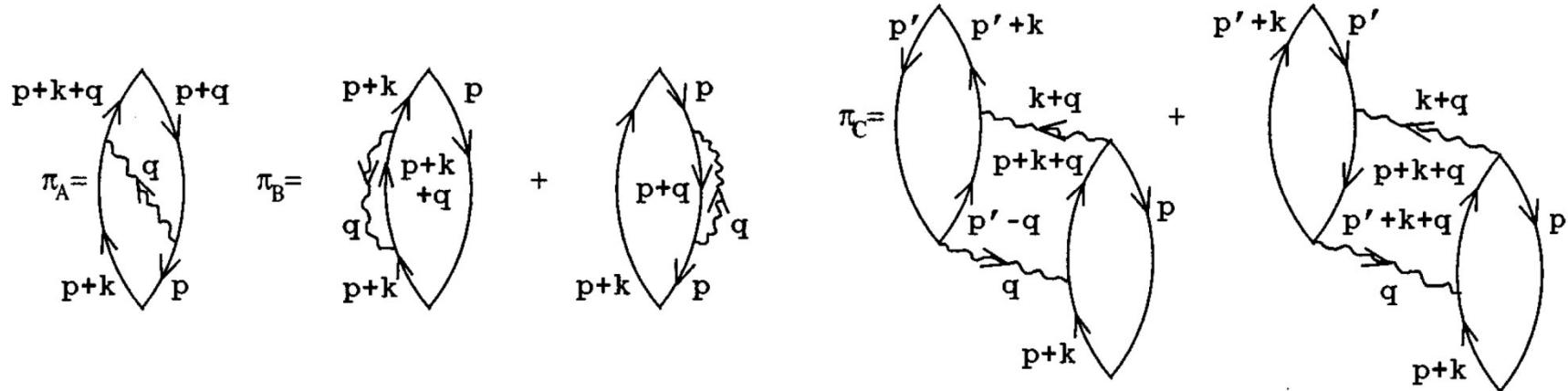
M. Cazzaniga et al., PRB **84**, 075109 (2011)



Double plasmons

Need to go many-body!

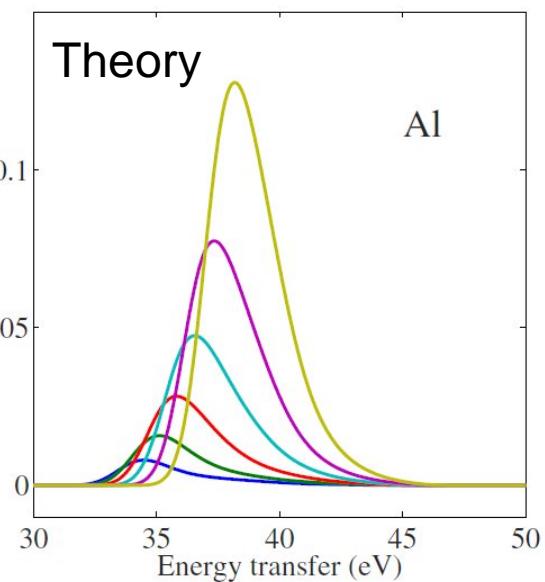
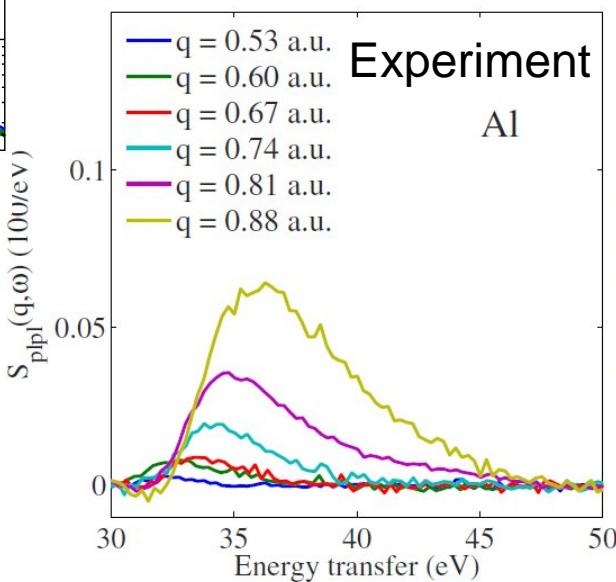
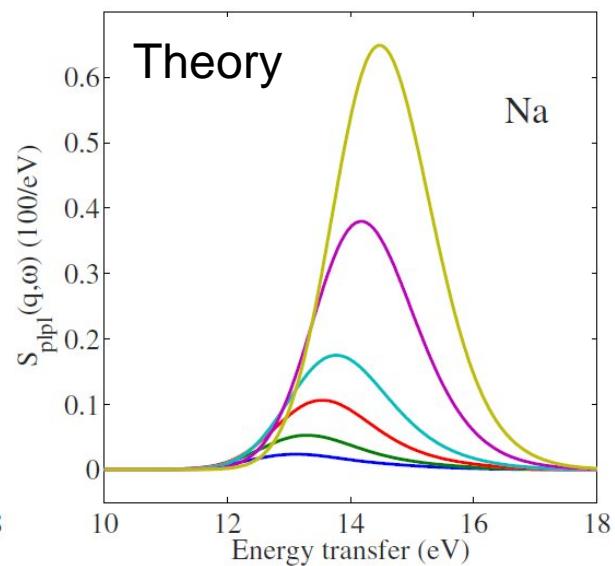
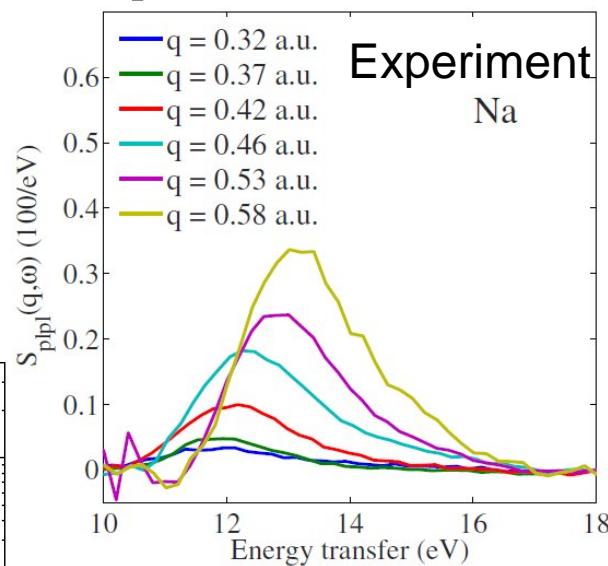
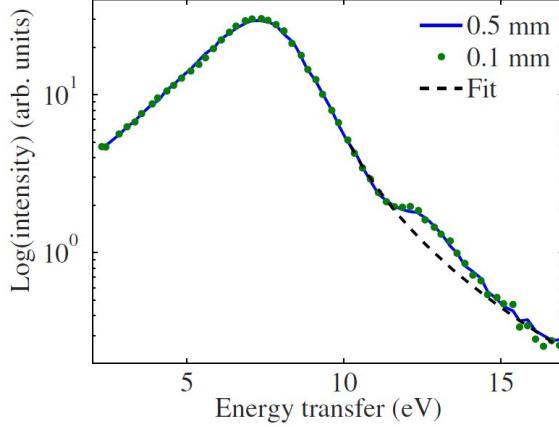
$$S(\mathbf{q}, \omega) = \frac{q^2}{4\pi^2 n_0} \text{Im} \left(\frac{-1}{\epsilon(\mathbf{q}, \omega)} \right) \approx \frac{1}{\pi n_0} \times \frac{\text{Im}[\pi_0(\mathbf{q}, \omega) + \pi_A(\mathbf{q}, \omega) + \pi_B(\mathbf{q}, \omega) + \pi_C(\mathbf{q}, \omega)]}{\epsilon_L^2(\mathbf{q}, \omega)}$$



K. Sturm and A. Gusarov, PRB 62, 16474 (2000)



Double plasmons



S. Huotari et al.,
PRB 77, 195125 (2008)



Double plasmons

Conclusions about double plasmons in Al, Na and Mg:

- Double plasmons observed in S(Q,E), not explained by TDDFT
- Many-body effects explain the spectra to ~2 eV and factor of 2 in intensity
- Should one try to incorporate the appropriate interaction in f_{xc} in TDDFT?

PRL **95**, 157401 (2005)

PHYSICAL REVIEW LETTERS

week ending
7 OCTOBER 2005

Correlation-Induced Double-Plasmon Excitation in Simple Metals Studied by Inelastic X-Ray Scattering

C. Sternemann,¹ S. Huotari,² G. Vankó,² M. Volmer,¹ G. Monaco,² A. Gusarov,^{3,4} H. Lustfeld,³ K. Sturm,³ and W. Schülke¹

PHYSICAL REVIEW B **77**, 195125 (2008)

Electron-density dependence of double-plasmon excitations in simple metals

S. Huotari,¹ C. Sternemann,² W. Schülke,² K. Sturm,³ H. Lustfeld,³
H. Sternemann,² M. Volmer,² A. Gusarov,⁴ H. Müller,¹ and G. Monaco¹



Outline

Part 1 Metal-to-insulator transition

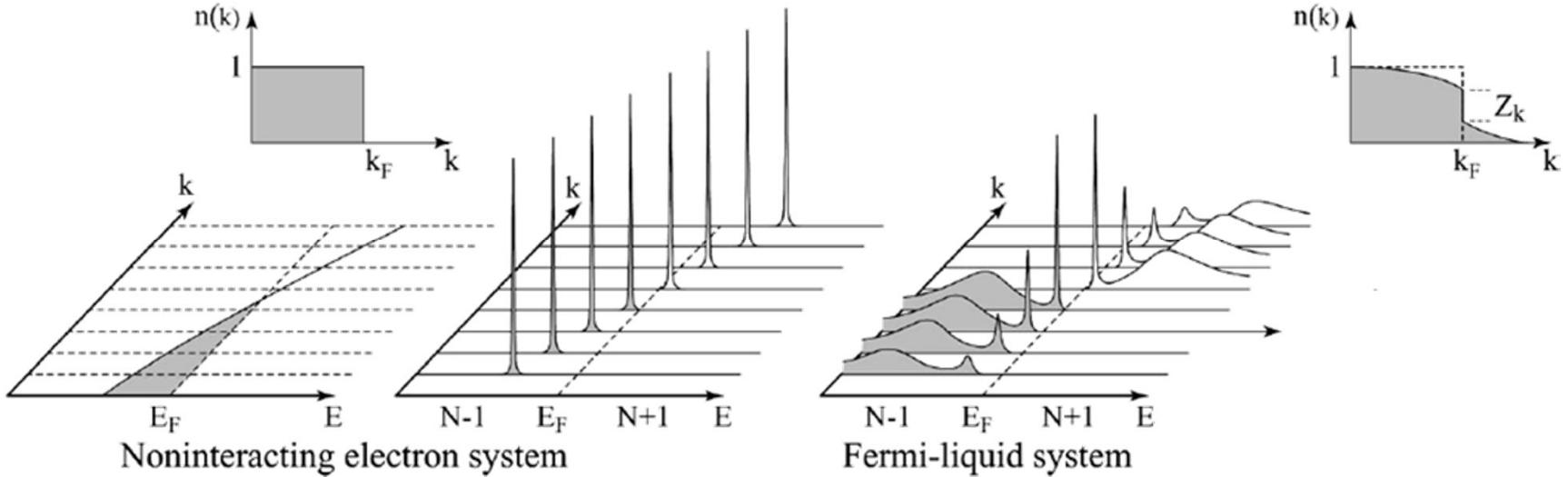
Part 2 Plasmons

Part 3 Double plasmons

Part 4 Quasiparticle renormalisation



Quasiparticle renormalisation

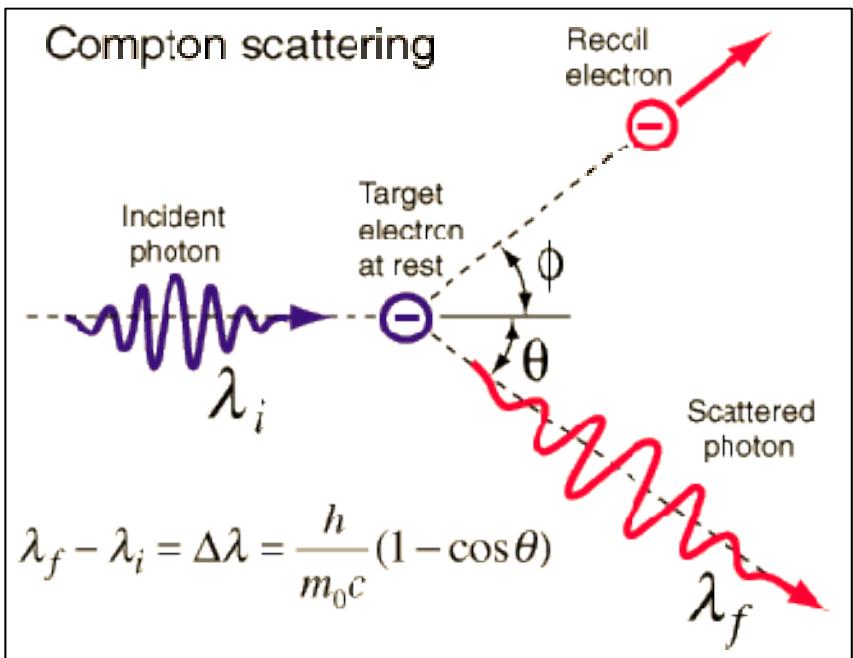


$$\text{PES} \propto \sum_{f,i} \left| M_{f,i}^{\mathbf{p}} \right|^2 A(\mathbf{p}, E) \delta(E_K + E_m^{N-1} - E_i^N - \omega_1)$$

A. Damascelli et al.,
Rev. Mod. Phys 75, 473 (2003)



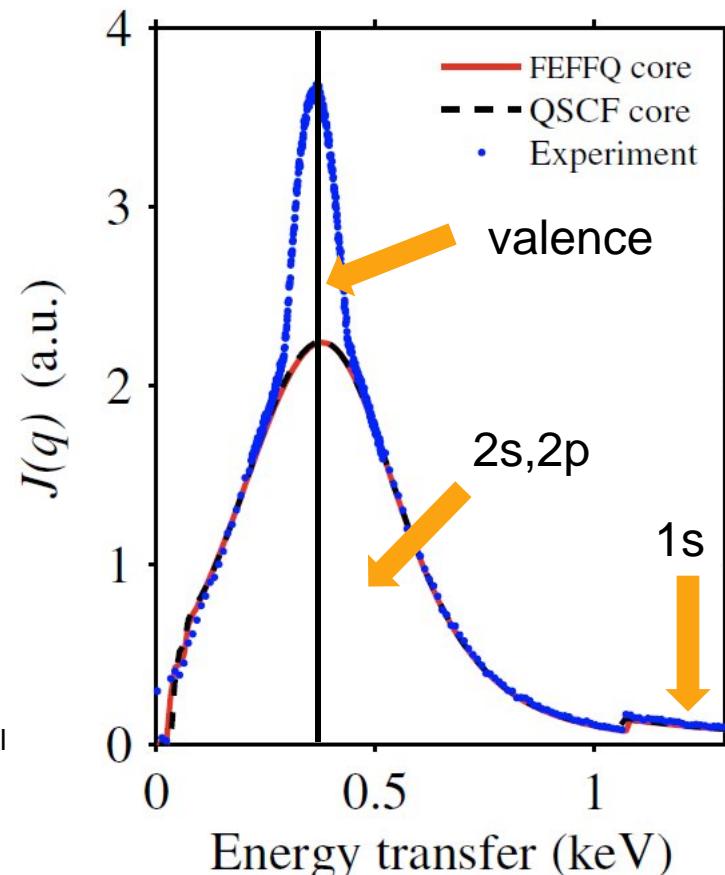
Compton spectroscopy



<http://hyperphysics.phy-astr.gsu.edu/hbase/quantum/compton.html>

Compton profile $J(q)$

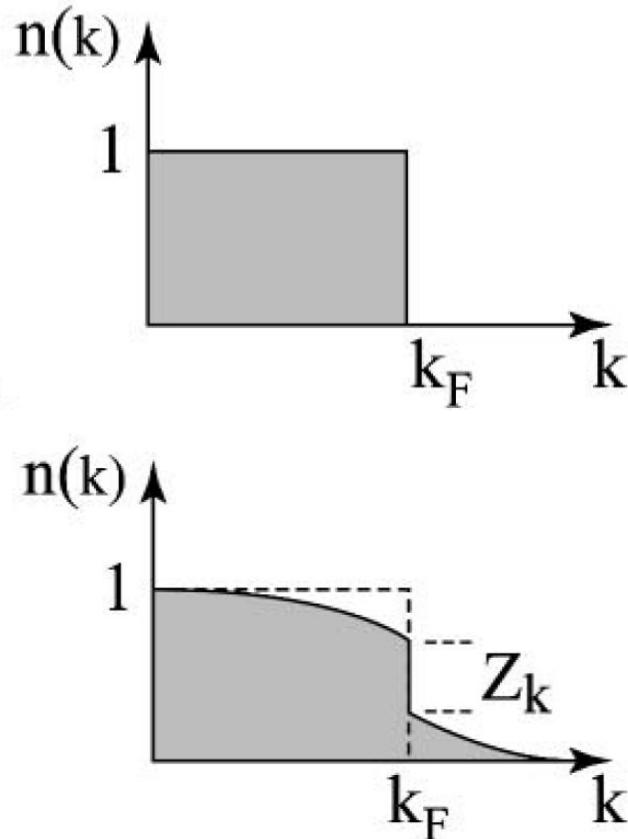
$$J(q) = \frac{3}{8\pi p_F^3} \int_{4\pi} d\Omega \int_{|q|}^\infty p n(\mathbf{p}) dp$$



S. Huotari et al.,
PRL 105, 086403 (2010)

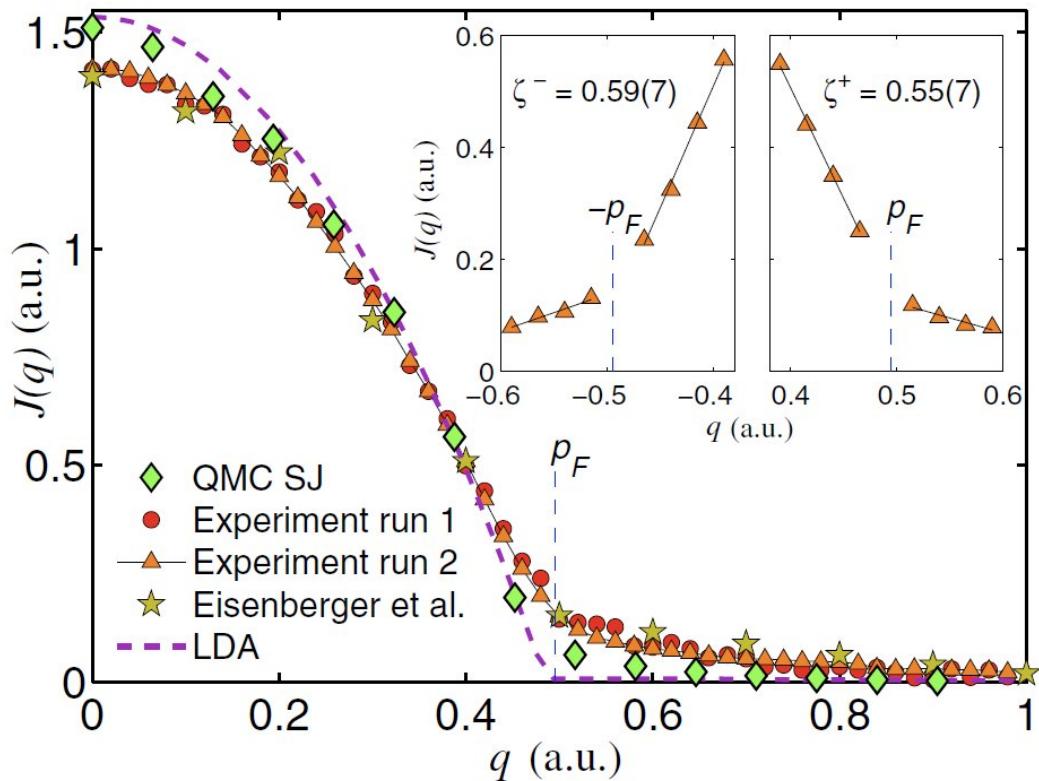


Compton spectroscopy



A. Damascelli et al.,
Rev. Mod. Phys 75, 473 (2003)

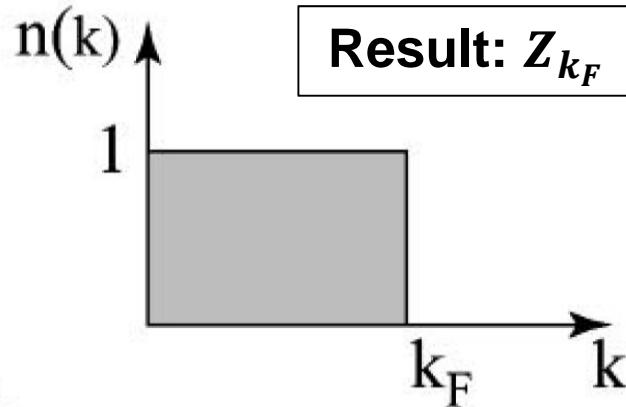
Independent electrons:
 $J(q) \propto (p_F^2 - q^2)$ for $|q| < p_F$



S. Huotari et al.,
PRL 105, 086403 (2010)

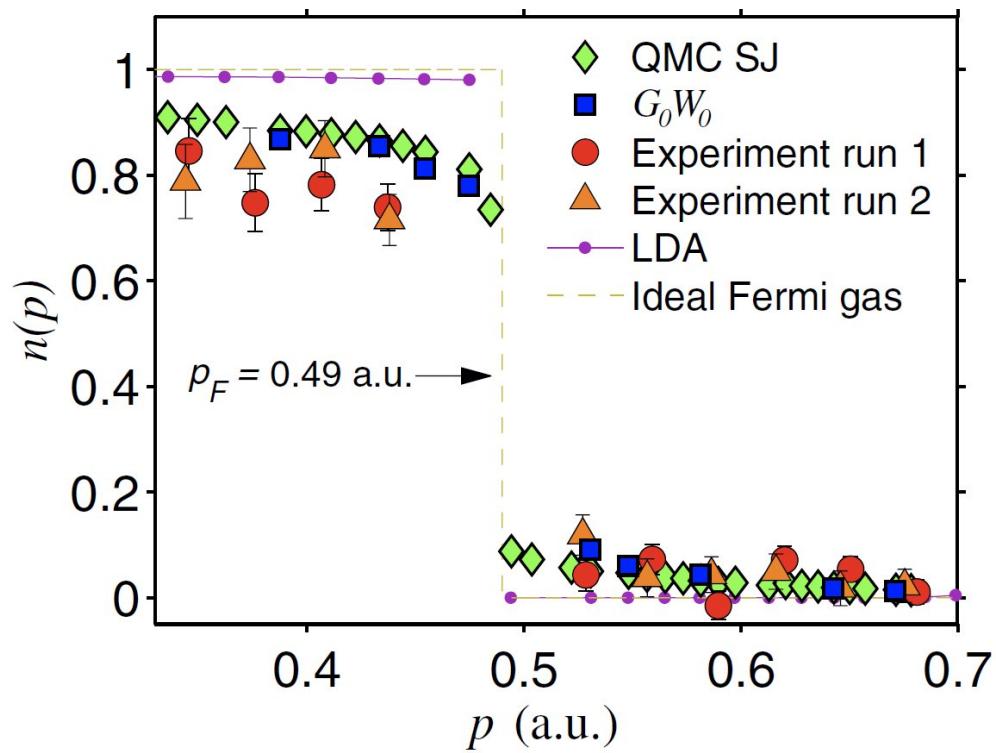
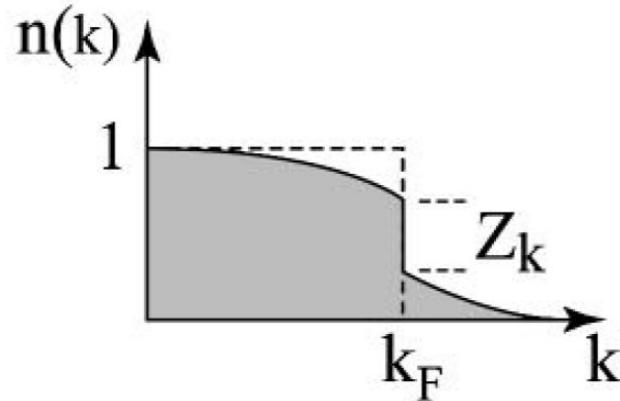


Momentum density of the electron gas



Result: $Z_{k_F} = 0.58 \pm 0.07$ for Na

$$n(p) = -\frac{2p_F^3}{3p} \frac{dJ(q)}{dq} \Big|_{q=p}$$





Quasiparticle renormalisation

Conclusions about the QP renormalisation factor:

- Is in principle the area under the coherent peak in photoemission
- Can be independently measured by momentum-density experiment
- Good results and agreement with theory obtained for Na

PRL 105, 086403 (2010)

PHYSICAL REVIEW LETTERS

week ending
20 AUGUST 2010

Momentum Distribution and Renormalization Factor in Sodium and the Electron Gas

Simo Huotari,^{1,2} J. Aleksi Soininen,² Tuomas Pylkkänen,^{1,2} Keijo Hämäläinen,² Arezki Issolah,³ Andrey Titov,⁴ Jeremy McMinis,⁵ Jeongnim Kim,⁵ Ken Esler,⁵ David M. Ceperley,⁵ Markus Holzmann,⁶ and Valerio Olevano⁴

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J. Aleksi Soininen (Univ. Helsinki)

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V. Olevano, M. Holtzmann (CNRS, Grenoble) (+everybody from QMC)

A. Eguiluz, A.V.Kozhevnikov, M.C.Troparevsky, T.C.Schulthess et al.