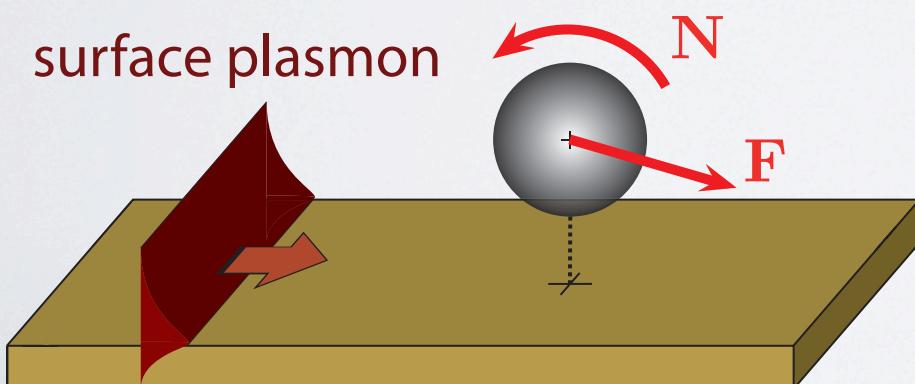


OPTICAL FORCES FROM PLASMONIC AND CHIRAL FIELDS

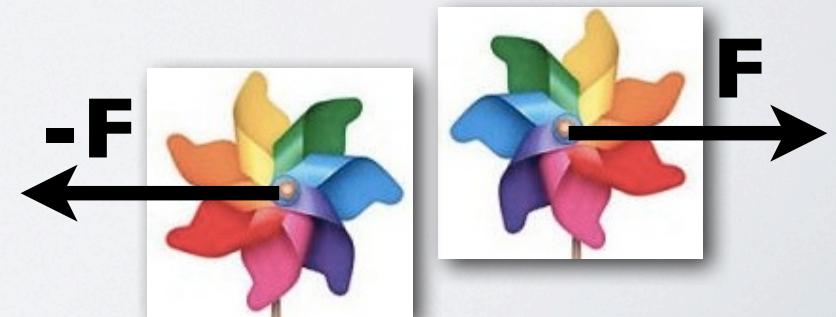
Antoine Canaguier-Durand - antoine.canaguier@gmail.com

ISIS, University of Strasbourg and CNRS

with Aurélien Cuche, James Hutchison, Éloïse Devaux,
Jino Georges, Cyriaque Genet and Thomas Ebbesen.

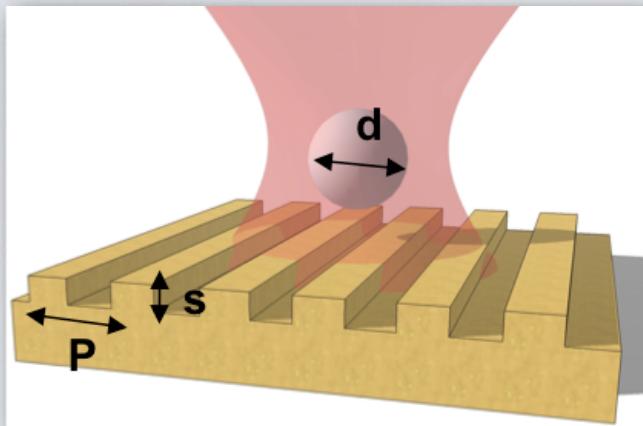


Chiral optical forces



Context: optical manipulation

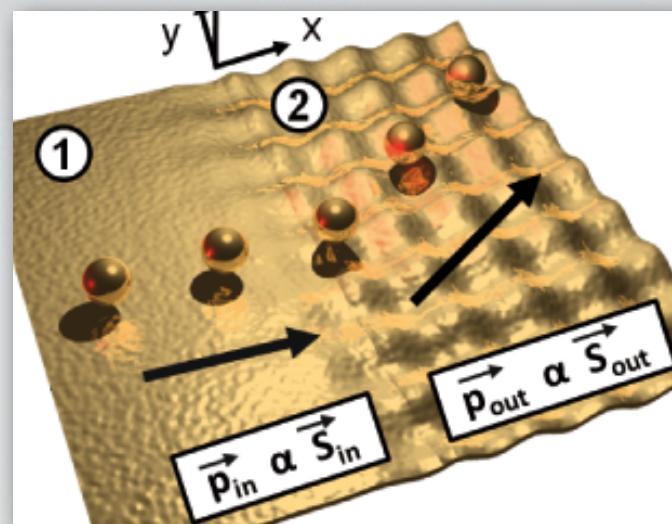
Optical tweezers
trap and move a NP



PRL **108**, 026801 (2012)

- biophysics (DNA, ...)
- optomechanical coupling

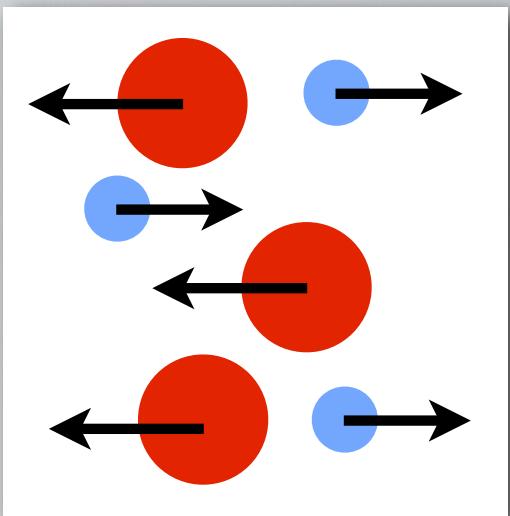
Optical routing
drive many NPs



Nano Lett. **12**, 4329 (2012)

- parallel manipulation
- micro/nano-fluidics

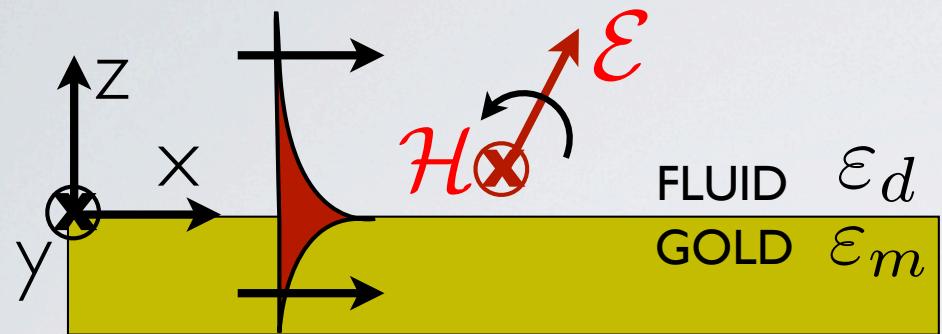
Optical sorting
separate NPs



Nano Lett. **13**,
4230 (2013)

- drug industry

The surface plasmon: a spinning field



$$\mathbf{E}_0 = E_0 e^{ikx} e^{iqz} \left(\tilde{q}, 0, -\tilde{k} \right)^t$$

$$\mathbf{H}_0 = H_0 e^{ikx} e^{iqz} (0, 1, 0)^t$$

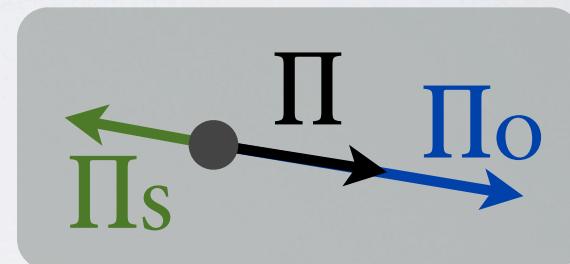
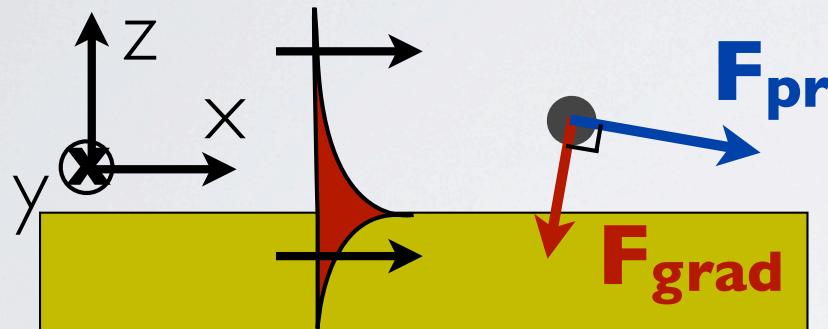
- Complex wavevector: $k = \frac{\omega}{c} \sqrt{\frac{\epsilon_m \epsilon_d}{\epsilon_d + \epsilon_m}}$
- Evanescent field
- Electric ellipticity: $\Phi_E = \mathcal{E} \times \frac{\dot{\mathcal{E}}}{\omega}$
- Decomposition of the time-averaged Poynting vector:

$$\Pi = \underbrace{\Pi_O}_{\text{orbital part}} + \underbrace{\frac{\nabla \times \Phi_E}{2\omega \mu_0}}_{\text{spin part}}$$

→ non-zero spin part

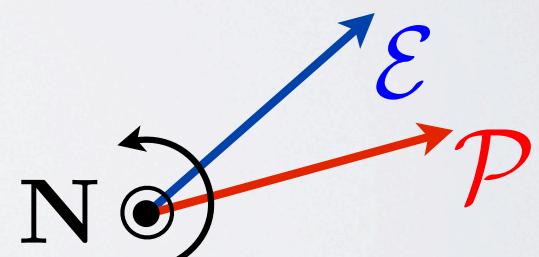
Plasmonic force and torque

Dipolar approximation: $\langle \mathbf{F} \rangle_T = \underbrace{\frac{\text{Re}[\alpha]}{4} \nabla \|\mathbf{E}_0\|^2}_{\text{gradient force}} + \underbrace{\text{Im}[\alpha] \omega \mu_0 \Pi_O}_{\text{radiation pressure}}$

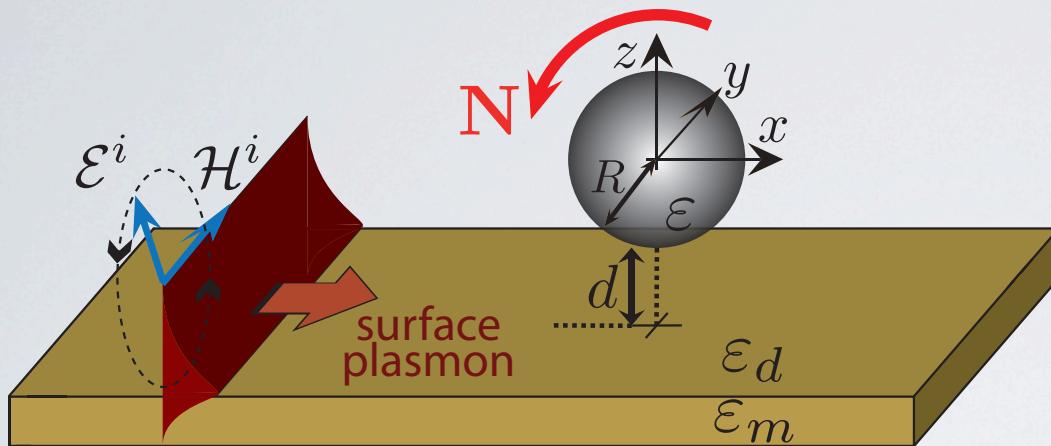


- Optical torque: $\mathbf{N} = \mathcal{P} \times \mathcal{E} = \text{Im}[\alpha] \Phi_E$

→ genuine mechanical rotation?

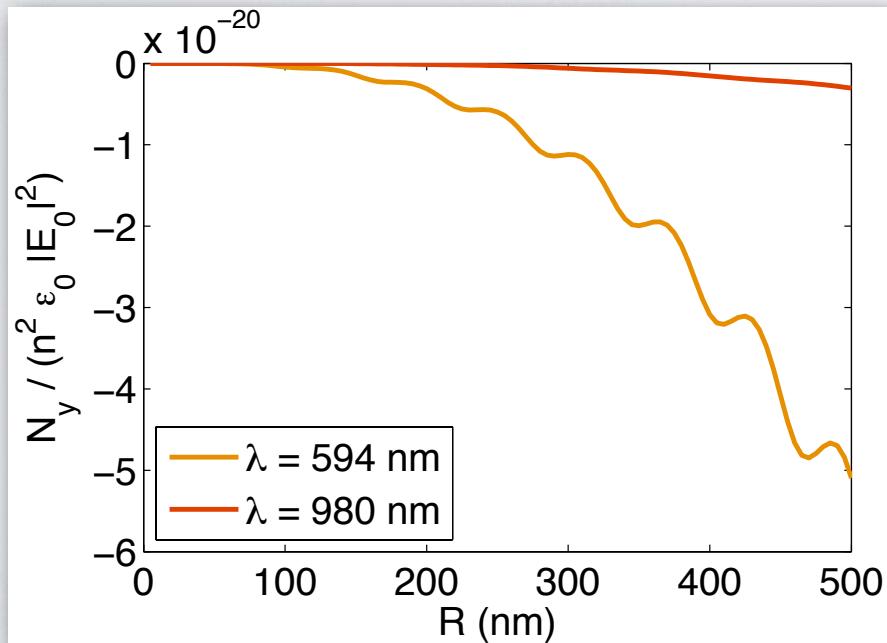


Plasmonic torque on a sphere



Method: multipolar calculation

- scattered field (Mie theory)
- optical force and torque from Maxwell stress tensor



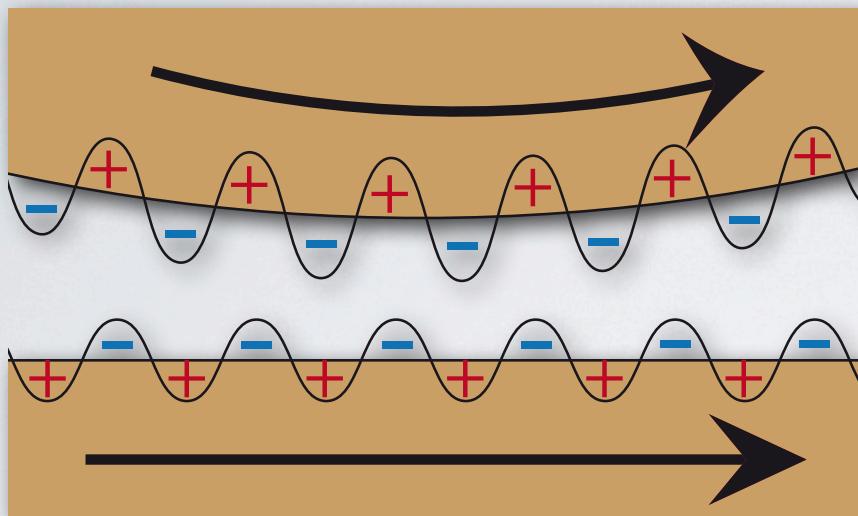
Results for gold nanospheres:

- strong dependence on R and λ
- transverse spinning

→ new tool for optical manipulation

Role of dissipation in the torque

Electronic point of view:



- induced movement of electrons
- dissipation transfers momentum

→ Non-contact gears

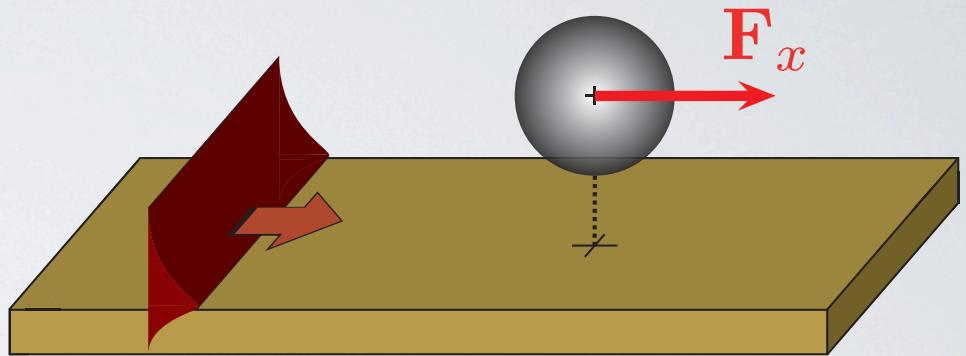
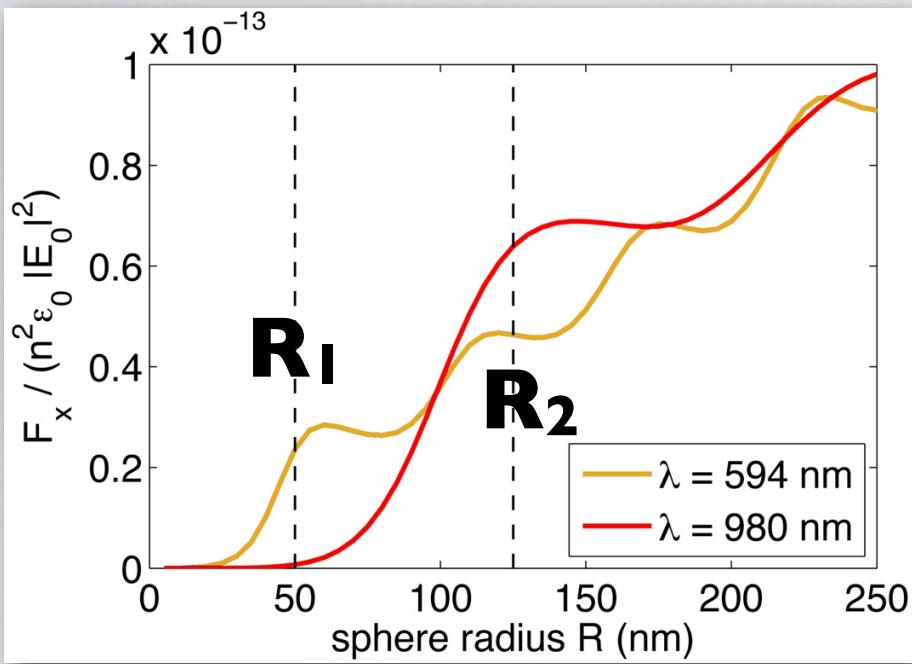
Test:

Case of a dielectric sphere with refractive index $n = 1.5 + in''$

→ torque is proportional to dissipation

Plasmonic force on a sphere

Results for the longitudinal force:



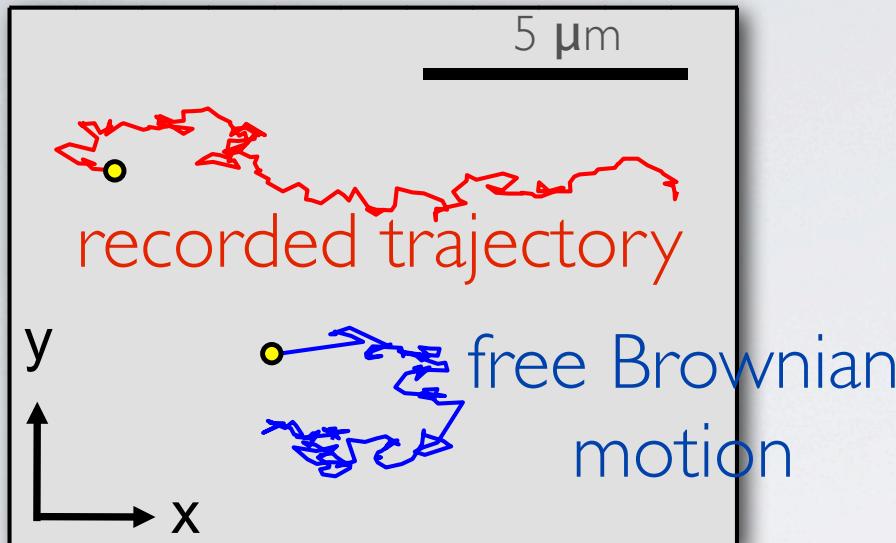
- confirms dipolar result for small R
- resonance effects (gold sphere)

potential for optical sorting R_1-R_2

counter-propagating scheme:



Plasmonic sorting demonstrated

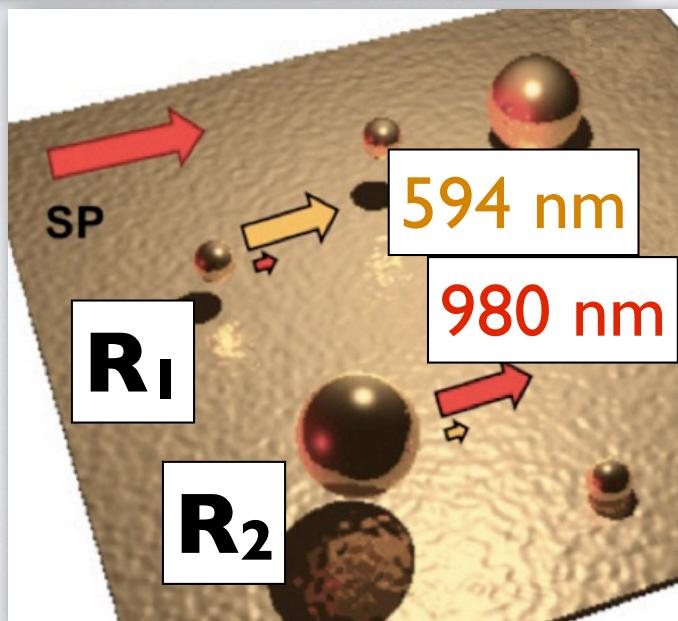


Experiment in the group at ISIS:
force observed through statistical
analysis of particles' motion.

$$F_x = \mu \frac{\gamma}{\Delta t}$$

← friction by the liquid

μ : mean displacement during Δt



Experimental confirmation for the potential of SP for sorting NPs.

Chiral objects

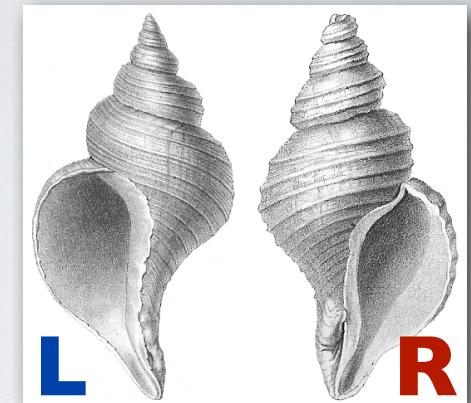
→ objects without spatial symmetry on a plane

In everyday life:



L: Neptunea angulata
R: Neptunea despecta

In gastropods:

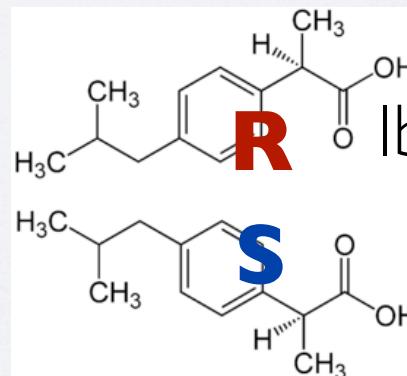


(10%) (90%)

In fundamental science:

- DNA, proteins
- weak interactions
- creation of life
- homochirality

In drug industry:

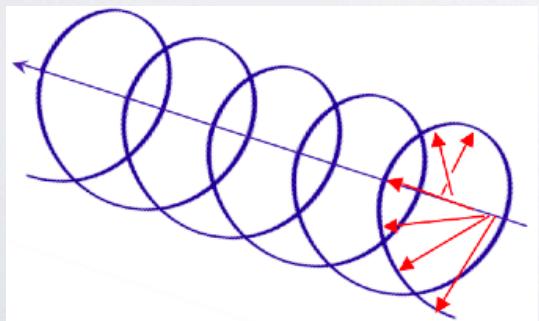


Ibuprofene (pain reliever)
S: active **R:** inactive
→ sold as **racemate**

Optical chirality

	Energy	Chirality
Density	$W(\mathbf{r}, t) = \frac{\varepsilon_0}{2} \mathcal{E} \cdot \mathcal{E} + \frac{\mu_0}{2} \mathcal{H} \cdot \mathcal{H}$	$K(\mathbf{r}, t) = \frac{\varepsilon_0}{2} \mathcal{E} \cdot (\nabla \times \mathcal{E}) + \frac{\mu_0}{2} \mathcal{H} \cdot (\nabla \times \mathcal{H})$
Flux	$\mathbf{S}(\mathbf{r}, t) = \mathcal{E} \times \mathcal{H}$	$\Phi(\mathbf{r}, t) = \frac{\mathcal{E} \times (\nabla \times \mathcal{H}) - \mathcal{H} \times (\nabla \times \mathcal{E})}{2}$
Conservation	$\nabla \cdot \mathbf{S} + \frac{\partial W}{\partial t} = 0$	$\nabla \cdot \Phi + \frac{\partial K}{\partial t} = 0$

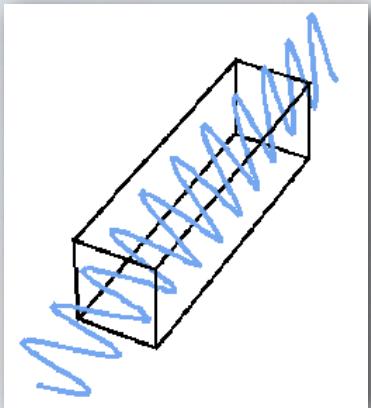
Circularly polarized wave



$$\Phi = \frac{\omega I_0}{2c} \hat{\mathbf{z}} \quad K = \frac{\omega I_0}{2c^2}$$

Chirality in light-matter interactions

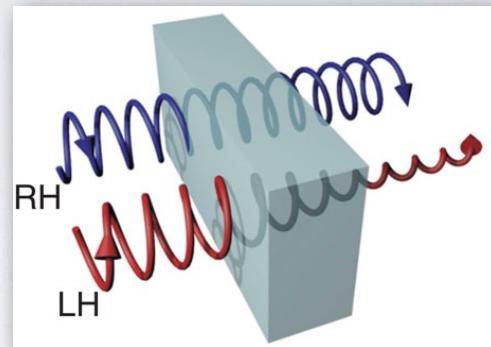
- Chiral matter on light:



Optical
rotation

$\text{Re}[\chi]$

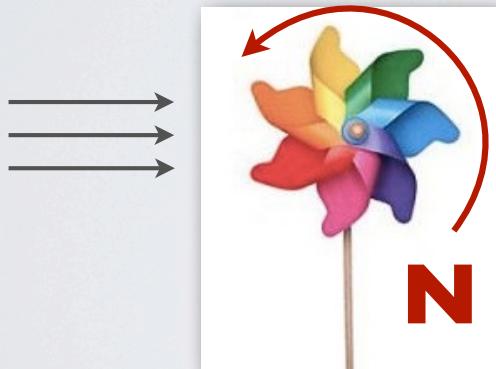
- Chiral matter on chiral light:



Circular
dichroism

$\text{Im}[\chi]$

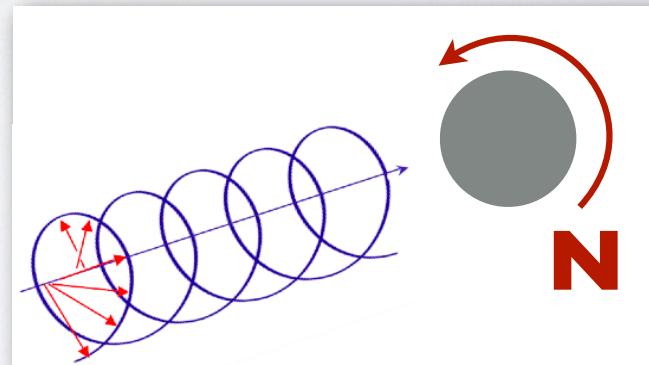
- Light on chiral matter:



Torque

$\text{Im}[\chi]$

- Chiral light on matter:



Torque

$\text{Im}[\alpha]$

→ Optical force of chiral light on chiral objects?

Optical force on a chiral dipole

- Chiral dipole with moments: $\begin{cases} \mathbf{p}_0 = \alpha \mathbf{E}_0 + i\chi \mathbf{H}_0 \\ \mathbf{m}_0 = -i\chi \mathbf{E}_0 + \beta \mathbf{H}_0 \end{cases}$
- Optical force: $\mathbf{F} = (\mathcal{P} \cdot \nabla) \mathcal{E} + (\mathcal{M} \cdot \nabla) \mathcal{H} + \mu_0 \dot{\mathcal{P}} \times \mathcal{H} + \varepsilon_0 \dot{\mathcal{M}} \times \mathcal{E}$

$$\rightarrow \langle \mathbf{F} \rangle_T = \operatorname{Re} [\alpha \mathbf{f}_0 + \beta \mathbf{g}_0 + \chi \mathbf{h}_0] / 2$$

chiral optical force

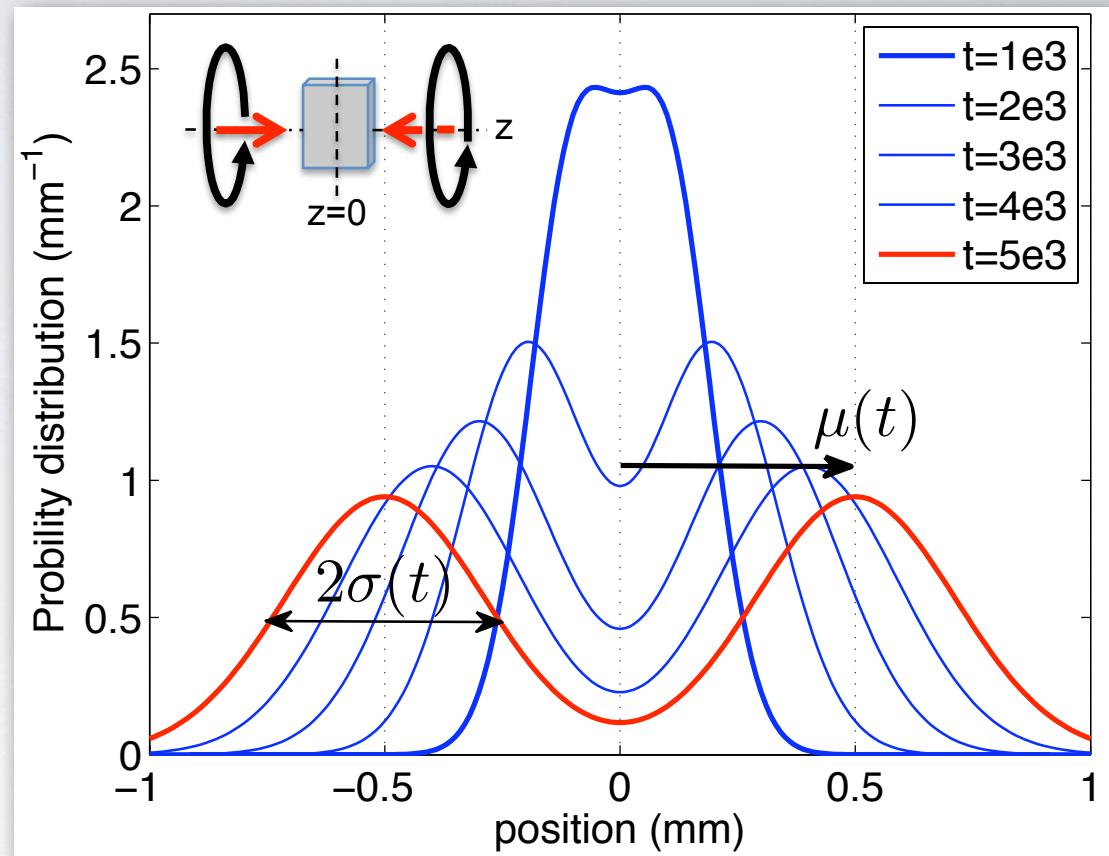
- $\operatorname{Re}[\chi] \longleftrightarrow \operatorname{Re}[\mathbf{h}_0] = \frac{2c^2}{\omega} \nabla K \rightarrow \text{reactive part}$
- $\operatorname{Im}[\chi] \longleftrightarrow \operatorname{Im}[\mathbf{h}_0^*] = 4 \left(\Phi - \frac{\nabla \times \mathbf{\Pi}}{2} \right) \rightarrow \text{dissipative part}$

Optical enantioseparation

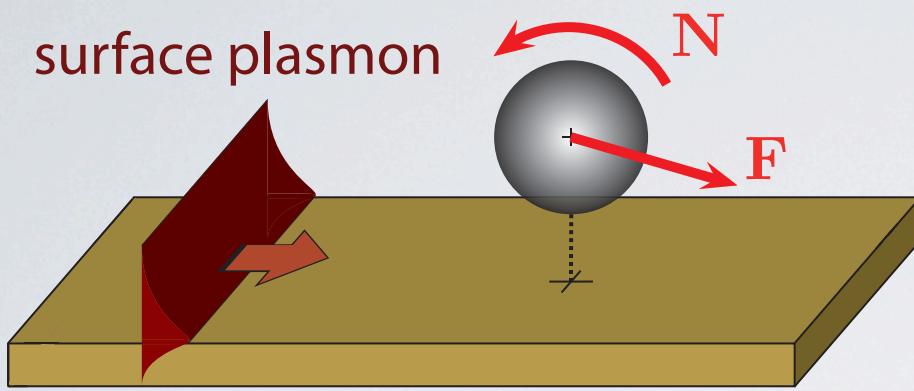
- counter-propagating CPL's
 - opposite handedness
 - incoherent fields
- only dissipative chiral force

$$\langle F \rangle_T = \frac{\omega I_0}{c} \text{Im}[\chi] \hat{\mathbf{z}}$$

- separation $\mu(t) \propto t$
- Brownian motion causes a variance $\sigma(t) \propto \sqrt{t}$



→ separation of $\sim 1\text{ mm}$ in one hour could be achieved
(50 mW on $1\text{ }\mu\text{m}^2$, for a 50 nm object with high chirality)



Chiral optical forces



conclusions:

- radiation pressure given by Π_0 , not Π !
- plasmonic torque able to spin NP transversely
- chiral light exerts chiral optical force on chiral objects

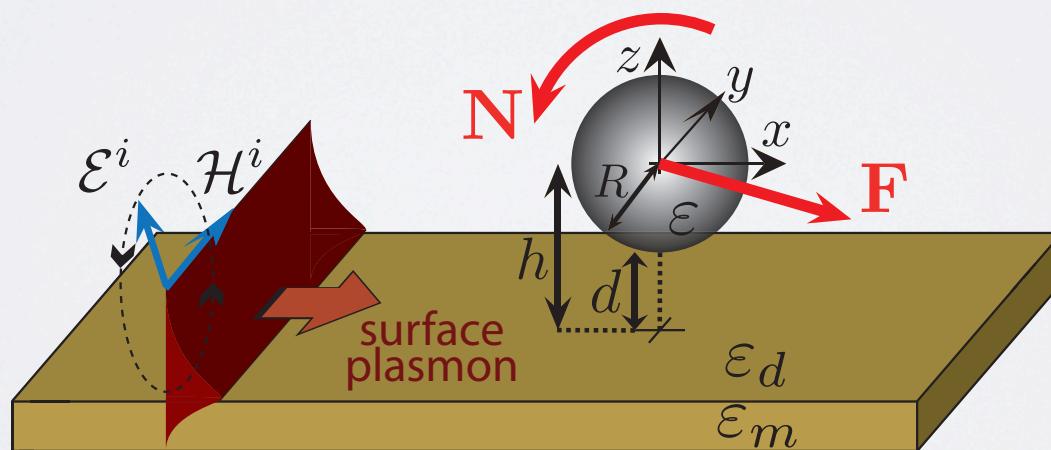
perspectives: towards the near-field

→ surface plasmons for efficient optical separation

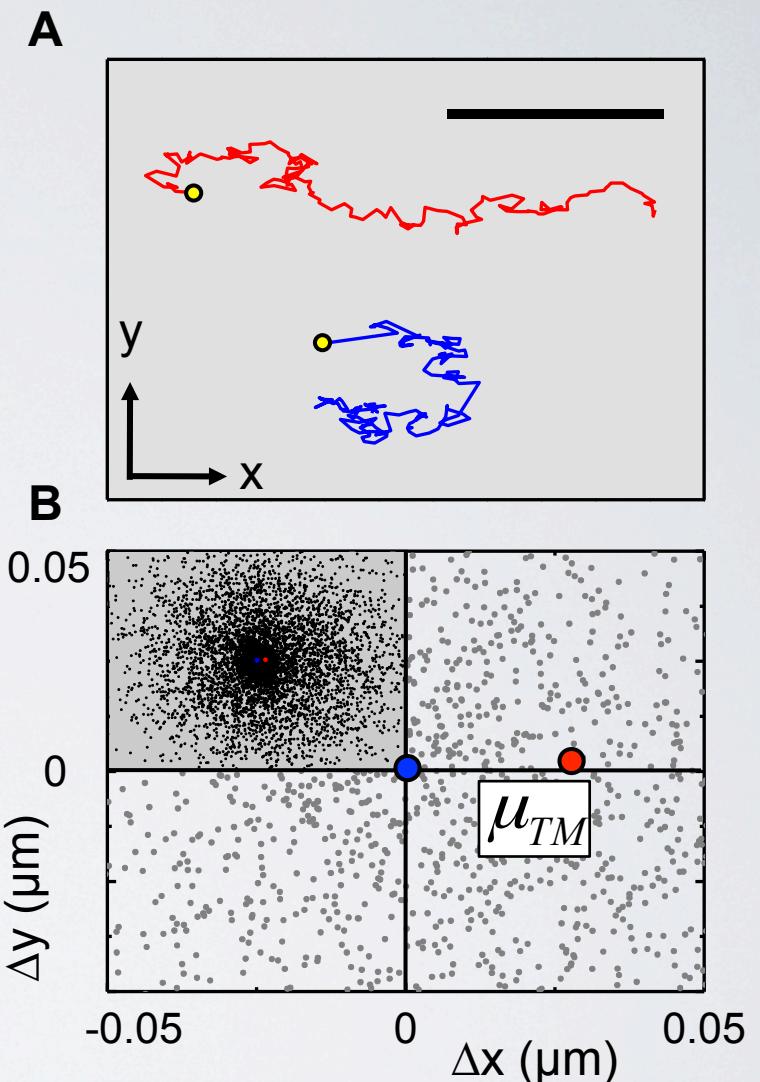
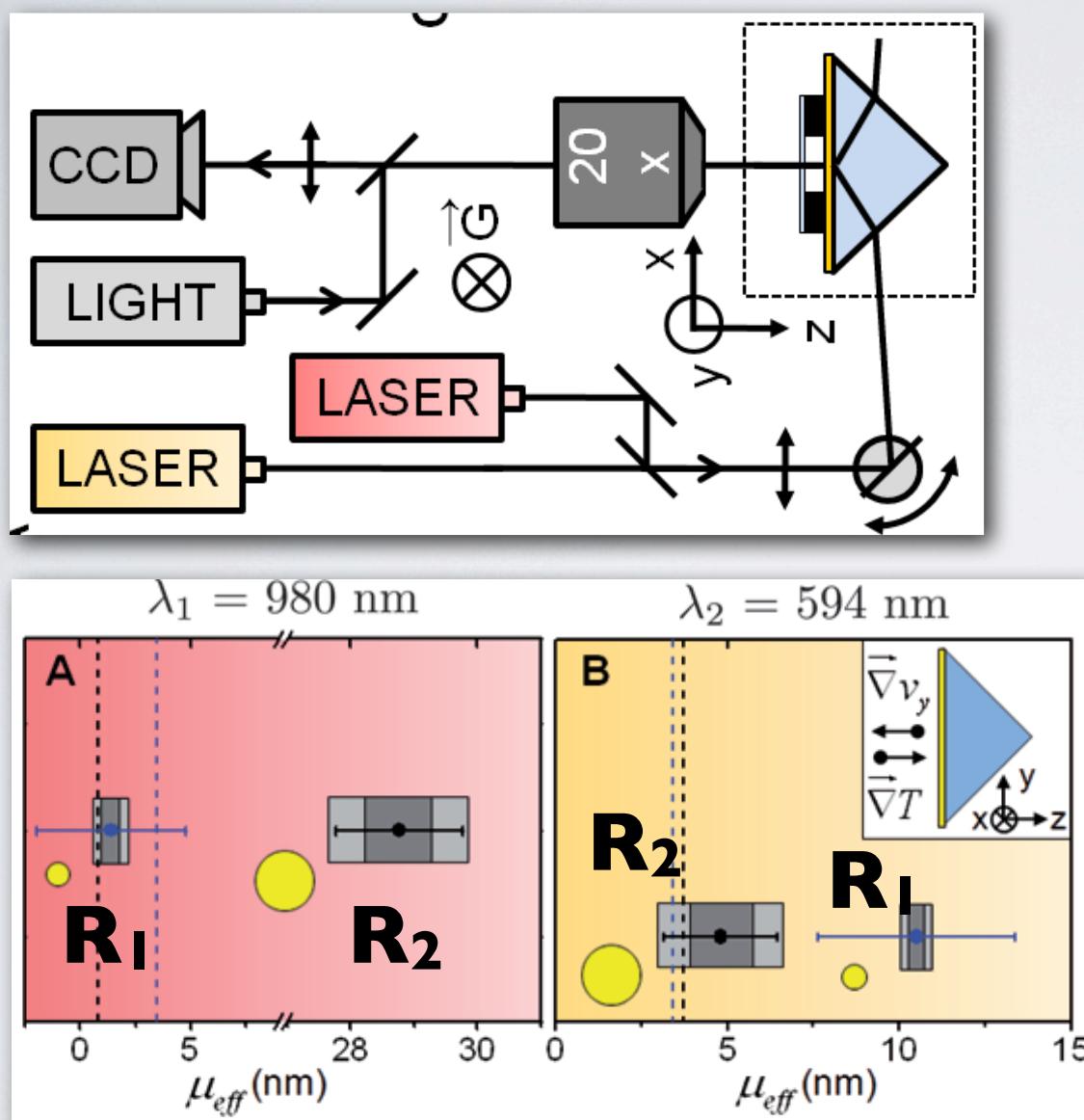
OPTICAL FORCES IN THE NEAR-FIELD

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Jino Georges, Cyriaque Genet and [Thomas Ebbesen](#).



Plasmonic sorting: experiment



A. Cuche, A. Canaguier-Durand, É. Devaux, J.A. Hutchison, C. Genet and T.W. Ebbesen, *Nano Letters* **13**, 4230-4235 (2013).

Plasmonic sorting: details

- Gold nanospheres of radii 50 and 125 nm in a 70 μm thick fluidic cell.
- Incident laser: $10^5\text{-}10^6 \text{ W.m}^{-2}$ @ 594 and 980 nm.
- 50 nm gold film evaporated on nanostructured ITO coated glass substrate.
- 5 nm SiO_2 layer evaporated on top of gold film to prevent electrostatic adhesion.
- Kretschmann configuration with metal/fluid interface (gold/water).
- Illumination angles $\theta = 64.3^\circ$ [594] and 76.5° [980], TM polarization.
- Evanescent tail of ~ 200 nm in the fluid.
- NP positions recorded with a high speed CCD camera ($f = 100$ Hz).

No natural generalization of the phase-gradient

$$\langle \mathbf{F} \rangle_T = \frac{1}{2} \operatorname{Re} [\alpha \mathbf{f}_0] , \quad \mathbf{f}_0 = \mathbf{E}_0 \cdot (\nabla) \quad \mathbf{E}_0^* = \sum_i E_i \nabla E_i^*$$

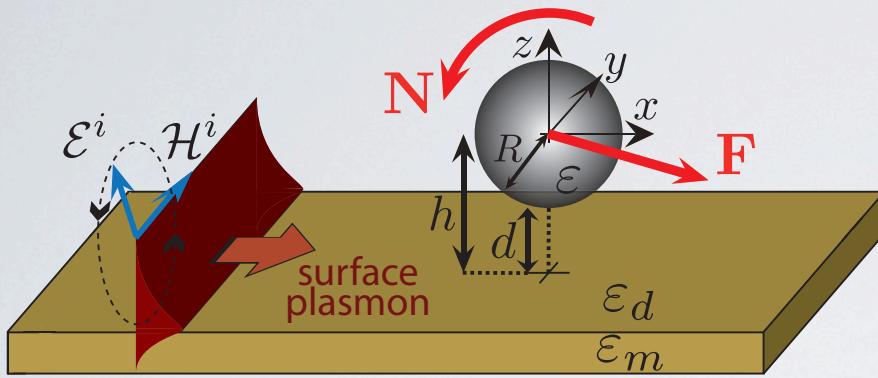
- linear polarization: $\boldsymbol{\Pi} = \boldsymbol{\Pi}_O$

$$\begin{aligned} \mathbf{E}_0(\mathbf{r}) = \rho(\mathbf{r}) e^{i\phi(\mathbf{r})} \hat{\mathbf{x}} &\longrightarrow \operatorname{Re}[\mathbf{f}_0] = \frac{1}{2} \nabla \rho^2 \\ &\longrightarrow \operatorname{Im}[\mathbf{f}_0^*] = 2\omega \mu_0 \boldsymbol{\Pi}_O = \rho^2 \nabla \phi \end{aligned}$$

- general field: $\boldsymbol{\Pi} \neq \boldsymbol{\Pi}_O$

$$\begin{aligned} E^j = \rho_j e^{i\phi_j} &\longrightarrow \operatorname{Re}[\mathbf{f}_0] = \frac{1}{2} \nabla \sum_j \rho_j^2 \\ &\longrightarrow \operatorname{Im}[\mathbf{f}_0^*] = 2\omega \mu_0 \boldsymbol{\Pi}_O = \sum_j \rho_j^2 \nabla \phi_j \end{aligned}$$

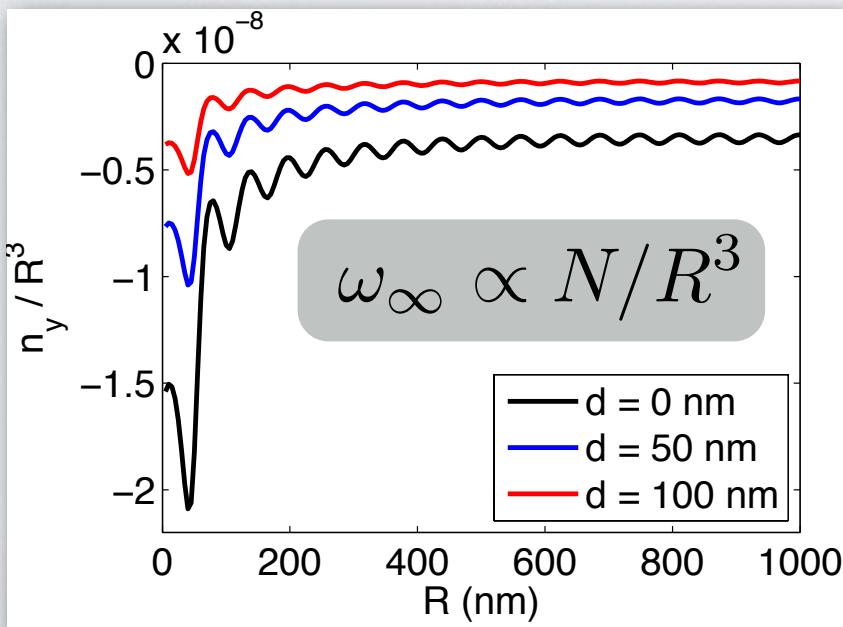
Transverse spinning of a NP in a fluid



$$I\dot{\omega} = N + N_{\text{drag}}$$

applied torque friction torque:
 $N_{\text{drag}} = -\gamma\omega$
 $\gamma = 8\pi\eta R^3$

→ final angular frequency: $\omega_\infty = N/\gamma$ reached in time I/γ .



- independent on R for $R \geq 200$ nm
- for 1 mW on $1 \mu\text{m}^2$, $\lambda = 594$ nm:
up to 20 Hz in water, 400 Hz in air
- transverse spinning for optical manipulation

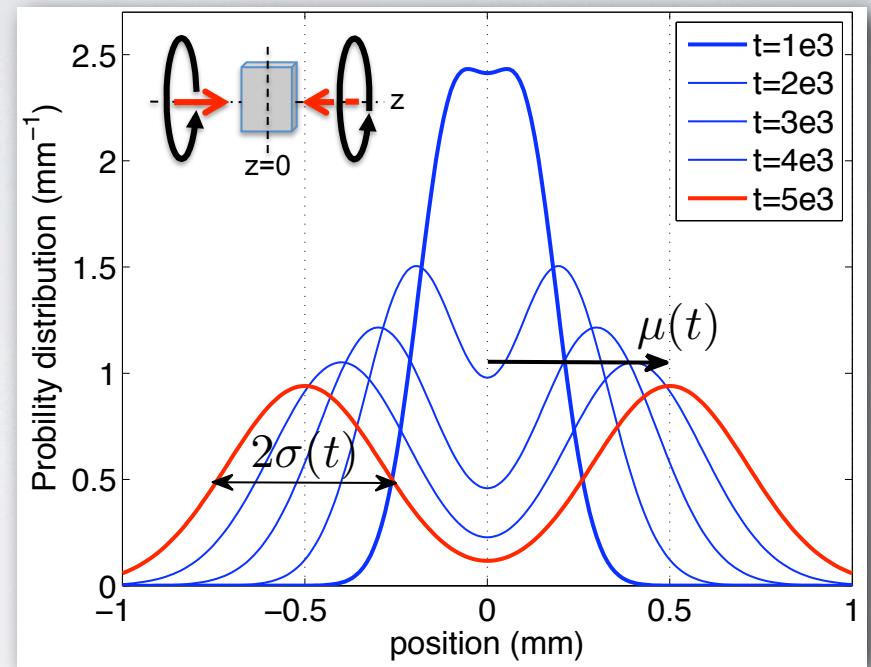
Optical enantioseparation

- counter-propagating CPL's
- opposite handedness
- incoherent fields

$$\langle F \rangle_T = \frac{\omega I_0}{c} \text{Im}[\chi] \hat{\mathbf{z}}$$

separation of $\sim 1\text{ mm}$ in one hour
could be achieved

- incident intensity: 50 mW on $1\text{ }\mu\text{m}^2 \rightarrow I_0 = 5 \cdot 10^{10} \text{ W} \cdot \text{m}^{-2}$
- very chirality object: $\text{Im}[c\chi] = 10^{-2} \cdot \text{Im}[\alpha/\epsilon_0]$ (W.Yan et. al. 2012)
- 20 nm radius gold nanoparticles: $\text{Im}[\alpha/\epsilon_0] \sim 10^{-22}$
- optical chiral force: $F \sim |fN| \rightarrow \Delta x = F \cdot \Delta t / \gamma \sim 10^{-6} \Delta t$



Optical torque on a chiral dipole

- Chiral dipole with moments: $\begin{cases} \mathbf{p}_0 = \alpha \mathbf{E}_0 + i\chi \mathbf{H}_0 \\ \mathbf{m}_0 = -i\chi \mathbf{E}_0 + \beta \mathbf{H}_0 \end{cases}$
- Optical torque: $\boldsymbol{\Gamma} = \mathcal{P} \times \mathcal{E} + \mathcal{M} \times \mathcal{H}$

$$\rightarrow \langle \boldsymbol{\Gamma} \rangle_T = \text{Im}[\alpha] \boldsymbol{\Phi}_E + \text{Im}[\beta] \boldsymbol{\Phi}_H + 2\text{Im}[\chi] \boldsymbol{\Pi}$$

