

# MPS: (1+1)-D QED

Buyens, B., Haegeman, J., Van Acoleyen, K., Verschelde, H., & Verstraete, F., Phys. Rev. Lett. 113, 091601 (2014)

Boye Buyens,<sup>1</sup> Karel van Acoleyen,<sup>1</sup>  
Jutho Haegeman,<sup>1</sup> Henri Verschelde,<sup>1</sup> Frank Verstraete<sup>1,2</sup>

<sup>1</sup>Department of Physics and Astronomy, Ghent University, Krijgslaan 281, S9, 9000 Gent, Belgium

<sup>2</sup>Vienna Center for Quantum Science and Technology, Faculty of Physics, University of Vienna, Boltzmanngasse 5, 1090 Vienna, Austria



## Introduction

$SU(N)$  Yang-Mills-theories: Yang, C. N., & Mills, R. L. 1954, Physical Review, 96, 191

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}, D_\mu = \partial_\mu - igA_\mu, F_{\mu\nu} = \frac{i}{g}[D_\mu, D_\nu]$$

Aim: Tensor networks (TN) to tackle strong coupling regime numerically

- TN allow efficient representation of low-energy physics
- Hamiltonian formulation → suited for real-time evolution
- Construction in manifest gauge invariant way

## The Schwinger model

= (1+1)-D QED with 1 flavour J. S. Schwinger, Phys. Rev. 125 (1962) 397.

$$\frac{H}{g} = \frac{1}{2\sqrt{x}} \sum_n \left( L_n^2 + \frac{\mu}{2}(-1)^n(\sigma_n^z + (-1)^n) + x(\sigma_n^+ e^{i\theta_n} \sigma_{n+1}^- + h.c.) \right),$$

Banks, T. and Susskind, Leonard and Kogut, John, 1976, Phys. Rev. D, 13, 1043-1053

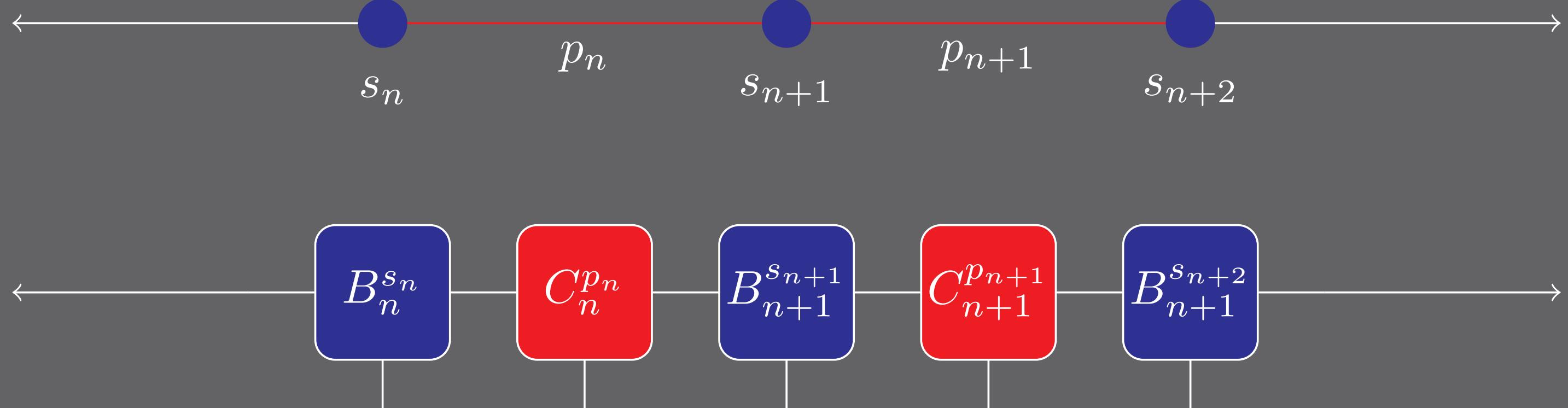
- Matter: spin operators  $\sigma_z$  on sites,  $\sigma^\pm$  ladder operators
- Gauge fields:  $L$  angular operator on links,  $e^{\pm i\theta}$  ladder operators
- $x = 1/(g^2 a^2)$ ,  $\mu = 2m/(g^2 a)$ ,  $a$  lattice spacing
- Continuum limit:  $x = +\infty$

Symmetries:

- Local  $U(1)$ :  $G_n = (L_n - L_{n-1} - [\sigma_n^z + (-1)^n]/2) = 0$ .
- Translation symmetry over two sites
- CT-symmetry  
→  $C$  charge conjugation,  $T$  translation

## Ground state

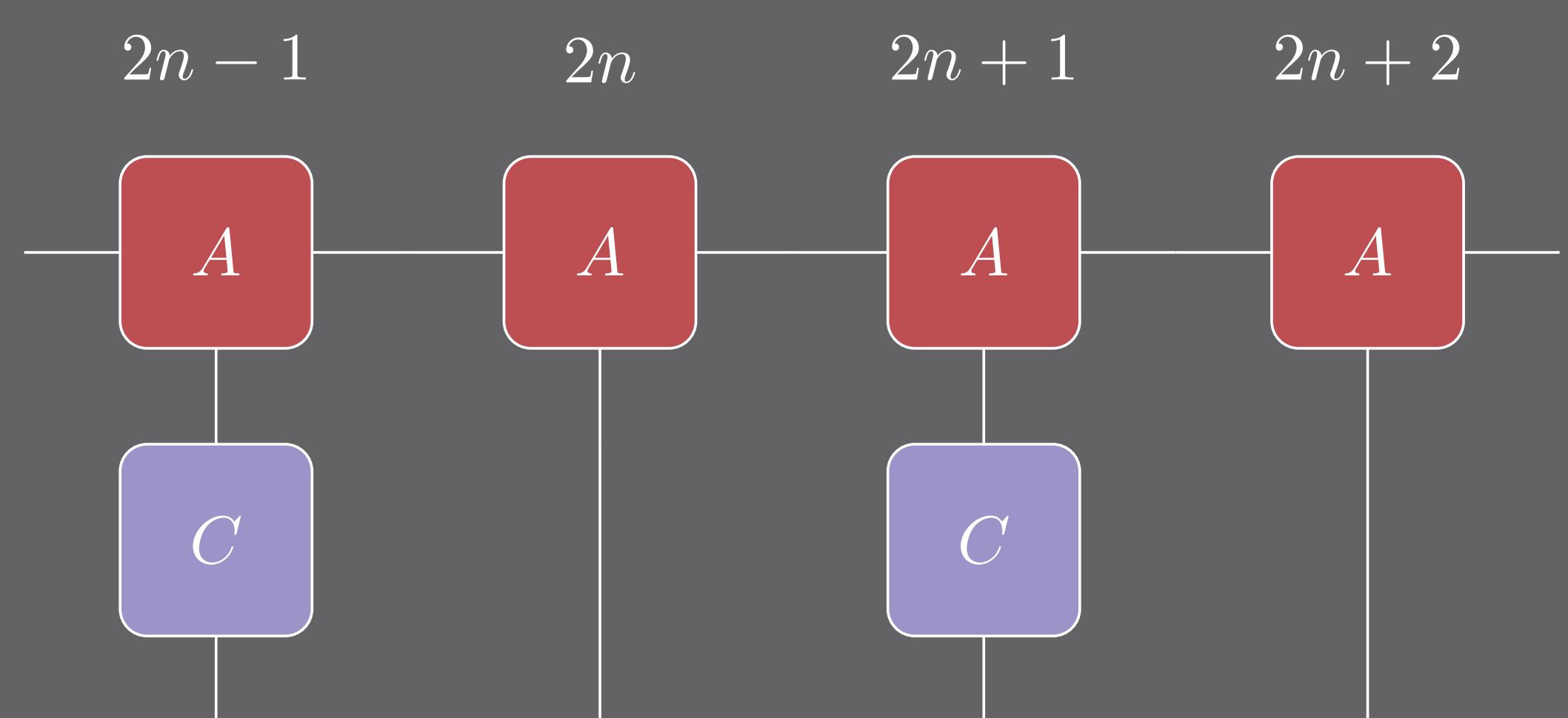
MPS-ansatz: M. Fannes, B. Nachtergaele, and R.F., Werner, Comm. Math. Phys. 144, 443 (1992)



Block site  $n$  and link  $n$ :

$$B_n \otimes C_n = A_n$$

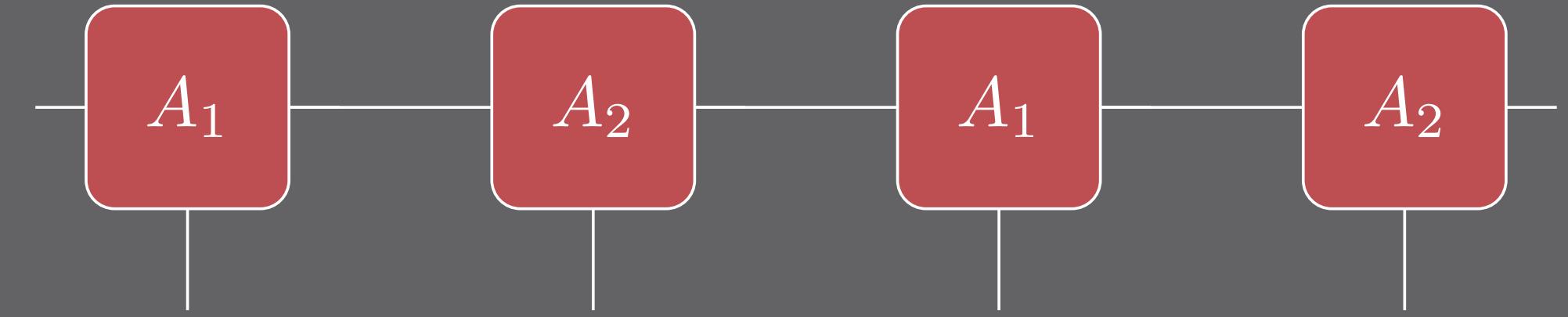
CT-symmetry:



- Gauge invariance:  $[A^{s,p}]_{(q,\alpha);(r,\beta)} = \delta_{p,q+(s+1)/2} \delta_{r,-p} [a^{s,p}]_{\alpha,\beta}$
- Optimization via TDVP Haegeman, J., Osborne, T. J., & Verstraete, F. 2013, Phys. Rev. B, 88, 075133

## Results

**Real-time evolution.** Quench  $\alpha$ :  $L_n \rightarrow L_n + \alpha$  in  $H$ : CT broken! MPS:



- Invariant under translations over two sites
- Gauge invariance:  $[A_n^{s,p}]_{(q,\alpha);(r,\beta)} = [a_n^{s,p}]_{\alpha,\beta} \delta_{p,q+(s+(-1)^n)/2} \delta_{r,p}$
- Real-time evolution with iTEBD Vidal, G. 2007, Physical Review Letters, 98, 070201

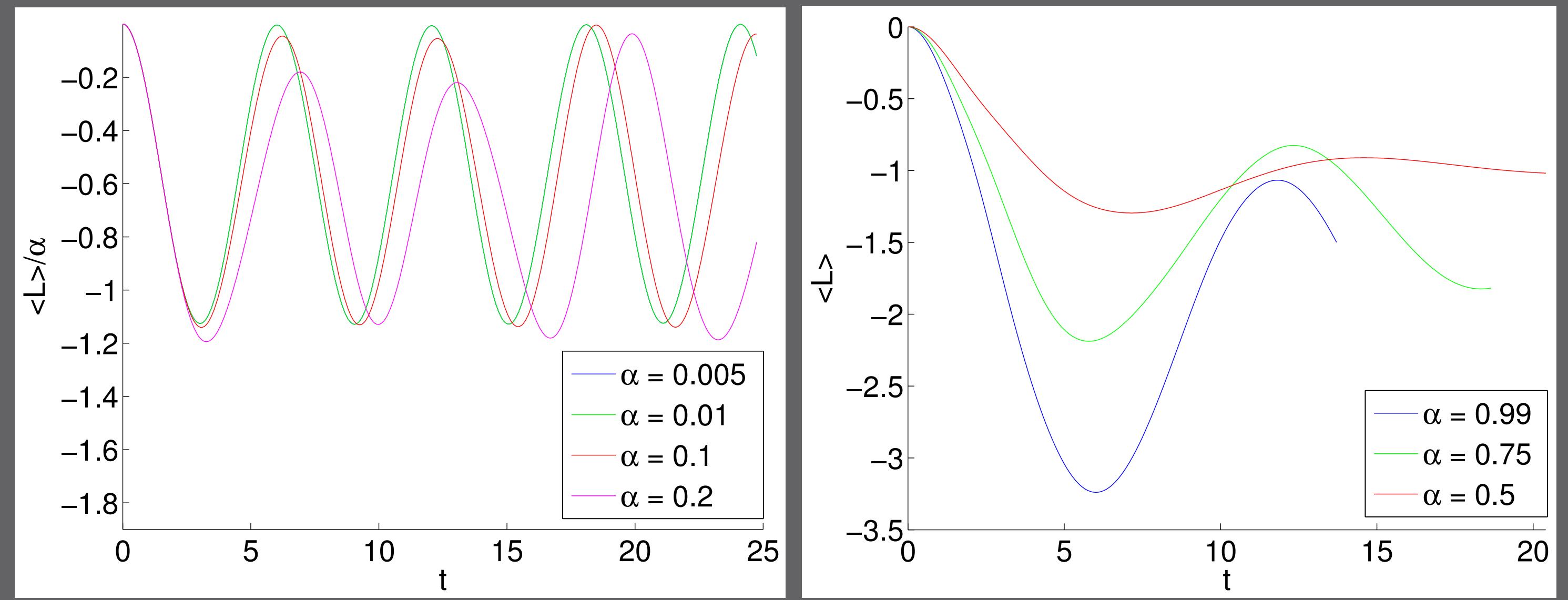


Figure: Real-time evolution ( $g = 1, x = 100, m = 0.25$ ): Evolution of electric field within linear response regime (left) and beyond (right).

**Excitations.** Ansatz with momentum  $k \in [-\pi/2a, \pi/2a]$  and C-number  $\gamma \in \{\pm 1\}$ . Haegeman, J., Pirvu, B., Weir, D. J., et al. 2012, Phys. Rev. B, 85, 100408

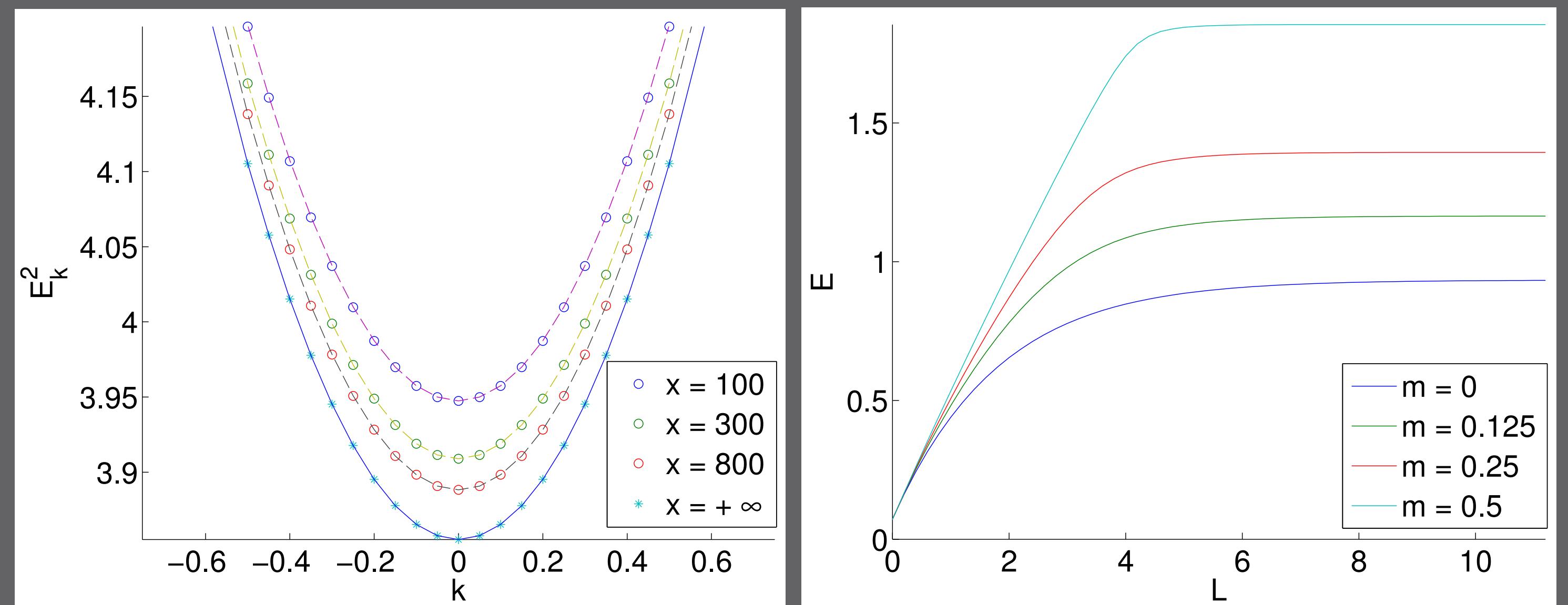
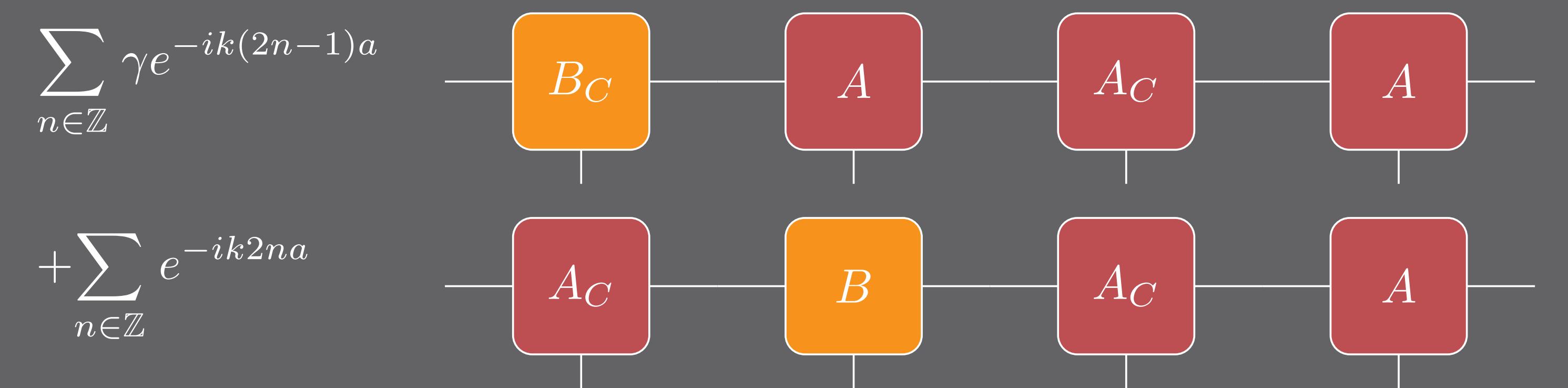
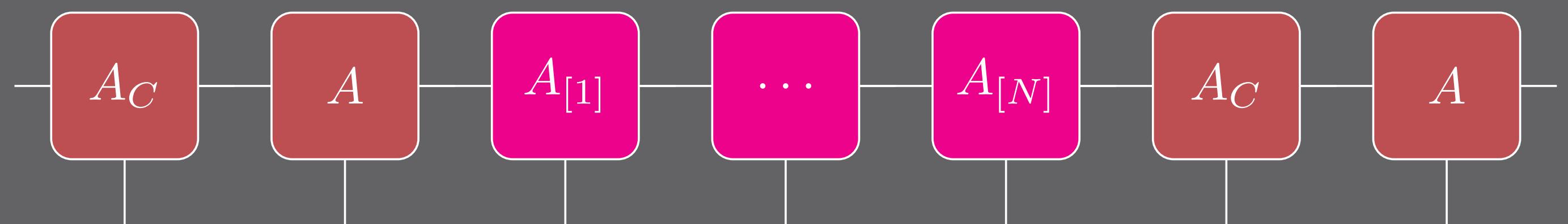


Figure: Left (a): Lorentz-dispersion relation for excitations. Energy squared vs. momentum ( $g = 1, m = 0.75, \gamma = -1$ ). Right (b): String breaking. Energy vs. separation length  $L$  between fermion and antifermion ( $x = 100, g = 1$ ).

**String breaking.** Massive fermion-antifermion pair in the vacuum.

- Anti-fermion on site  $j_-$ , fermion on site  $j_+$ , charge  $q^{ext}$
- Not dynamical, influence Gauss' law:  $G_n + q^{ext}(\delta_{n,j_+} - \delta_{n,j_-}) = 0$

MPS-ansatz:



- Asymptotically: ground state
- Around fermion-antifermion pair: new variational freedom  $A_{[n]}$   
→ DMRG White, S. R. 1992, Physical Review Letters, 69, 2863
- $[A_{[n]}^{s,p}]_{(q,\alpha);(r,\beta)} = [a_{[n]}^{s,p}]_{\alpha,\beta} \delta_{p,q-\delta_{n,j_-}+\delta_{n,j_+}+(s+(-1)^n)/2}$