

# Zitterbewegung in External Magnetic Field

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# The Dirac Equation

The Dirac equation describes spin  $\frac{1}{2}$  particles.

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$\psi = u(\mathbf{p})e^{-ip \cdot x}$$

$$Hu = \begin{pmatrix} m & \boldsymbol{\sigma} \cdot \mathbf{p} \\ \boldsymbol{\sigma} \cdot \mathbf{p} & -m \end{pmatrix} \begin{pmatrix} u_A \\ u_B \end{pmatrix} = E \begin{pmatrix} u_A \\ u_B \end{pmatrix}$$

$$\boldsymbol{\sigma} \cdot \mathbf{p} u_B = (E - m)u_A$$

$$\boldsymbol{\sigma} \cdot \mathbf{p} u_A = (E + m)u_B$$

$$u^{(s)} = N \begin{pmatrix} \chi^{(s)} \\ \frac{\boldsymbol{\sigma} \cdot \mathbf{p}}{E + m} \chi^{(s)} \end{pmatrix}, \quad E > 0 \quad u^{(s+2)} = N \begin{pmatrix} \frac{-\boldsymbol{\sigma} \cdot \mathbf{p}}{|E| + m} \chi^{(s)} \\ \chi^{(s)} \end{pmatrix}, \quad E < 0$$

# Zitterbewegung

Erwin Schrödinger in 1930

$$\vec{v} = \hbar \frac{\partial x_k}{\partial t} = i[H, x_k] = c\alpha_k$$

$$(\vec{\alpha})_k^2 = \mathbf{1}$$

$$v_k = \pm c$$

$$[(\vec{\alpha})_k, (\vec{\alpha})_\ell] \neq 0 \quad \text{for } k \neq \ell$$

$$\hbar \frac{\partial \alpha_k}{\partial t} = i[H, \alpha_k] = 2[i\gamma_k m - \sigma_{kl} p^l] = 2i[p_k - \alpha_k H]$$

$$\alpha_k(t) = \alpha_k(0) \exp[-2iHt/\hbar] + cp_k H^{-1}$$

$$x_k(t) = x_k(0) + c^2 p_k H^{-1} + \\ + \frac{1}{2} i \hbar c H^{-1} (\alpha_k(0) - cp_k H^{-1}) (\exp(-2iHt/\hbar) - 1)$$

$$\vec{L} \; = \; \vec{r} \; \times \; \vec{p}$$

$$\vec{S} \; = \; \frac{1}{2} \; \vec{\Sigma}$$

$$\frac{d}{dt} \; \vec{L} \; = \; i \left[ H, \vec{L} \right] \; = \; i(\vec{\alpha})_\ell \; \left[ p_\ell, \vec{L} \right]$$

$$\frac{d}{dt} \; \vec{L} \; = \; \vec{\alpha} \; \times \; \vec{p}$$

$$\frac{d}{dt} \; \vec{S} \; = \; -\vec{\alpha} \; \times \; \vec{p}$$

$$\vec{J} \; := \; \vec{L} \; + \; \vec{S}$$

$$\frac{d}{dt} \; \vec{J} \; = \; 0$$

$$\begin{aligned}
 \langle \vec{v} \rangle &= \langle \vec{\alpha} \rangle \\
 &= u^\dagger (\vec{p}) \vec{\alpha} u(\vec{p}) \\
 &= N^2 \, \eta^\dagger \, \frac{\vec{\sigma}(\vec{\sigma} \cdot \vec{p}) + (\vec{\sigma} \cdot \vec{p}) \, \vec{\sigma}}{p^0 + m} \, \eta \\
 &= N^2 \, \frac{2\vec{p}}{p^0 + m} \, \eta^\dagger \, \eta \; .
 \end{aligned}$$

$$\langle \vec{v} \rangle = \frac{\vec{p}}{p_0} = \vec{v}_{classical}$$

$$\langle \frac{d\vec{L}}{dt} \rangle = \vec{0} = \langle \frac{d\vec{S}}{dt} \rangle$$

It turns out that if one only looks at a the expectation value of only positive or only negative energy waves the Zitterbewegung term vanishes.

The most common interpretation of the Zitterbewegung-term is usually interpreted as an interference between positive- and negative-energy waves.

Phenomena usually mentioned in advanced books of Quantum Mechanics <sup>1</sup>

Often mentioned in passing to the interpretation of negative energy states and the " Dirac Sea" as antiparticles.

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<sup>1</sup>For example: A. Messiah, Quantum Mechanics Volume II  
**J. J. Sakurai:** Advanced Quantum Mechanics

Shroedinger's interpretation of his results from the Dirac equation:

Local circulatory motion  $\Rightarrow$  Electron spin and magnetic moment

$$\omega_0 = 2mc^2/\hbar = 1.6 \cdot 10^{21} s^{-1}$$

$$\text{So } \lambda_0 = c/\omega_0 = \hbar/2mc = 1.9 \cdot 10^{-13} m$$



$$\psi(\mathbf{r}, t) = h^{-\frac{3}{2}} \int [C^+(\mathbf{p})e^{-i\omega t} + C^-(\mathbf{p})e^{i\omega t}] \exp(i\mathbf{p} \cdot \mathbf{r}/\hbar) d^3\mathbf{p},$$

$$C^+(\mathbf{p}) = a_1 u_1 + a_2 u_2$$

$$C^-(\mathbf{p}) = a_3 u_3 + a_4 u_4$$

$$u_1 = \begin{bmatrix} 1 \\ 0 \\ kp_3 \\ kp_+ \end{bmatrix}, \quad u_2 = \begin{bmatrix} 0 \\ 1 \\ kp_- \\ -kp_3 \end{bmatrix},$$

$$u_3 = \begin{bmatrix} -kp_3 \\ -kp_+ \\ 1 \\ 0 \end{bmatrix}, \quad u_4 = \begin{bmatrix} -kp_- \\ kp_3 \\ 0 \\ 1 \end{bmatrix},$$

$$\psi(\mathbf{r}, 0) = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} f(r/r_0)$$

$$f(r/r_0) = (2/\pi r_0^2)^{\frac{1}{2}} \exp[-(r/r_0)^2]$$

$$\begin{aligned} \psi(\mathbf{r}, t) = & h^{-\frac{3}{2}} \int \left[ \begin{bmatrix} 1 \\ 0 \\ kp_3 \\ kp_+ \end{bmatrix} e^{-i\omega t} \right. \\ & \left. + \begin{bmatrix} k^2 p^2 \\ 0 \\ -kp_3 \\ -kp_+ \end{bmatrix} e^{i\omega t} \right] \frac{f(p/\eta)}{1+k^2 p^2} \exp(i\mathbf{p} \cdot \mathbf{r}/\hbar) d^3\mathbf{p} \end{aligned}$$

$$\langle \dot{\mathbf{r}} \rangle = \bar{\mathbf{v}} + 2c \int \mathbf{K}(\mathbf{p}) \sin(2\omega t + \phi(\mathbf{p})) d^3\mathbf{p}$$

$$\mathbf{K}(\mathbf{p}) = |C^{-*} \boldsymbol{\alpha} C^+|,$$

$$\phi(\mathbf{p}) = \tan^{-1}\{\text{Re}(C^{-*} \boldsymbol{\alpha} C^+)/\text{Im}(C^{-*} \boldsymbol{\alpha} C^+)\}$$

$$\langle x_1 \rangle = -I \int_0^{2\pi} \lambda \sin(2\omega t + \varphi) d\varphi \quad \langle x_2 \rangle = -I \int_0^{2\pi} \lambda \cos(2\omega t + \varphi) d\varphi$$

$$I = -(32\pi)^{-\frac{1}{2}}(\lambda/r_0)$$

$$\langle x_1 \rangle_{\varphi} = -\lambda \sin(2\omega t + \varphi)$$

$$\langle x_2 \rangle_{\varphi} = -\lambda \cos(2\omega t + \varphi)$$

$$\langle \mathbf{r} \times \dot{\mathbf{r}} \rangle_1 = 0, \quad \langle \mathbf{r} \times \dot{\mathbf{r}} \rangle_2 = 0$$

$$\langle \mathbf{r} \times \dot{\mathbf{r}} \rangle_3 = (\hbar/m)(1 - \cos 2\omega t)$$

$$\langle \mu_1 \rangle = 0, \quad \langle \mu_2 \rangle = 0$$

$$\langle \mu_3 \rangle = (e\hbar/2mc)(1 - \cos 2\omega t)$$

### *Zitterbewegung*

in the presence of an external uniform static magnetic field

$$i\hbar\frac{\partial\Psi}{\partial t}=\left[c\alpha_i(p^i-\frac{e}{c}A^i)+\beta mc^2+eA_0\right]\Psi$$

$$E_{\rm ln}^2=m^2c^4+p_z^2c^2+ceB(n-l+1-2s_z)$$

$$\frac{E_{\rm B\pm}}{\hbar}\simeq\pm\left(\frac{mc^2}{\hbar}\pm\frac{|e|\,B\sigma_{\pm}}{2mc}\right)\equiv\pm(\omega\pm\Omega\sigma_{\pm})$$

$$\Psi(\mathbf{r}, t) = h^{-3/2} \int \left\{ \mathcal{A} \begin{pmatrix} \phi^\uparrow \\ k_+^\uparrow \boldsymbol{\sigma} \cdot \boldsymbol{\pi} \phi^\uparrow \end{pmatrix} \exp[-i(\omega + \Omega)t] + \mathcal{B} \begin{pmatrix} \phi^\downarrow \\ k_+^\downarrow \boldsymbol{\sigma} \cdot \boldsymbol{\pi} \phi^\downarrow \end{pmatrix} \exp[-i(\omega - \Omega)t] \right. \\ \left. + \mathcal{C} \begin{pmatrix} k_-^\uparrow \boldsymbol{\sigma} \cdot \boldsymbol{\pi} \chi^\uparrow \\ \chi^\uparrow \end{pmatrix} \exp[i(\omega - \Omega)t] + \mathcal{D} \begin{pmatrix} k_-^\downarrow \boldsymbol{\sigma} \cdot \boldsymbol{\pi} \chi^\downarrow \\ \chi^\downarrow \end{pmatrix} \exp[i(\omega + \Omega)t] \right\} \exp[i(p_y y + p_z z)/\hbar] d^3\pi,$$

for the spin-up states

$$\begin{cases} \exp(i\omega t) \rightarrow \exp[i(\omega - \Omega)t] & \text{for negative energy} \\ \exp(-i\omega t) \rightarrow \exp[-i(\omega + \Omega)t] & \text{for positive energy} \end{cases}$$

for the spin-down states

$$\begin{cases} \exp(i\omega t) \rightarrow \exp[i(\omega + \Omega)t] & \text{for negative energy} \\ \exp(-i\omega t) \rightarrow \exp[-i(\omega - \Omega)t] & \text{for positive energy} \end{cases}$$

$$\Psi(\boldsymbol{r},0)=\begin{pmatrix}1\\0\\0\\0\end{pmatrix}f\left(\frac{r}{r_o}\right)$$

$$f(\frac{\pi}{\pi_o})=(\frac{2}{\bar{\pi}\pi_o^2})^{3/4}\exp\Big[-(\frac{\pi}{\pi_o})^2\Big]$$

$$\begin{aligned}\Psi(\boldsymbol{r},t)\simeq h^{-3/2}\int\bigg\{\begin{pmatrix}1\\0\\K\pi_z\\K\pi_+\end{pmatrix}\exp\big[-i(\omega+\Omega)t\big]&+\begin{pmatrix}0\\0\\-K\pi_z\\0\end{pmatrix}\exp\big[i(\omega-\Omega)t\big] \\ +\begin{pmatrix}0\\0\\0\\-K\pi_+\end{pmatrix}\exp\big[i(\omega+\Omega)t\big]\bigg\}f(\pi/\pi_o)\exp[i(p_yy+p_zz)/\hbar]d^3\pi\,.\end{aligned}$$

$$< \dot{\boldsymbol{r}} > = \int \Psi^*(\boldsymbol{r}, t) (c\boldsymbol{\alpha}) \Psi(\boldsymbol{r}, t) d^3r$$

$$< x > \simeq I^\uparrow \frac{\lambda_c}{2} \int_0^{2\pi} \sin \left[ (\omega_{\text{zbw}} + \omega_c) t + \varphi \right] d\varphi$$

$$< y > \simeq I^\uparrow \frac{\lambda_c}{2} \int_0^{2\pi} \cos \left[ (\omega_{\text{zbw}} + \omega_c) t + \varphi \right] d\varphi$$

$$< z > \simeq J \lambda_c \sin (\omega_{\text{zbw}} t)$$

$$I^\uparrow \equiv -(8\pi)^{-\frac{1}{2}} \frac{\lambda_c}{r_o} \left( 1 - \frac{\omega_c}{\omega_{\text{zbw}}} \right)$$

$$J \equiv 0$$

$$< x >_{\varphi_o} \simeq \frac{\lambda_c}{2} \sin \left[ (\omega_{\text{zbw}} + \omega_c) t + \varphi_o \right]$$

$$< y >_{\varphi_o} \simeq -\frac{\lambda_c}{2} \cos \left[ (\omega_{\text{zbw}} + \omega_c) t + \varphi_o \right]$$

The same procedure for the spin-down

$$\langle x \rangle \simeq I \downarrow \frac{\lambda_c}{2} \int_0^{2\pi} \sin \left[ (\omega_{\text{zbw}} - \omega_c) t - \varphi \right] d\varphi$$

$$\langle y \rangle \simeq I \downarrow \frac{\lambda_c}{2} \int_0^{2\pi} \cos \left[ (\omega_{\text{zbw}} - \omega_c) t - \varphi \right] d\varphi$$

$$\langle z \rangle \simeq J \lambda_c \sin (\omega_{\text{zbw}} t)$$

$$\begin{aligned} I \downarrow &\equiv -(8\pi)^{-\frac{1}{2}} \frac{\lambda_c}{r_o} \left( 1 + \frac{\omega_c}{\omega_{\text{zbw}}} \right) \\ J &\equiv 0 \end{aligned}$$



$$\langle \mu_x^\uparrow \rangle = 0, \quad \langle \mu_y^\uparrow \rangle = 0 \quad \text{and} \quad \langle \mu_z^\uparrow \rangle = -\frac{|e| \lambda_c}{2} \left[ 1 - \cos(\omega_{\text{zbw}} + \omega_c)t \right]$$

$$\langle \mu_x^\downarrow \rangle = 0, \quad \langle \mu_y^\downarrow \rangle = 0 \quad \text{and} \quad \langle \mu_z^\downarrow \rangle = \frac{|e| \lambda_c}{2} \left[ 1 - \cos(\omega_{\text{zbw}} - \omega_c)t \right]$$

$$v_{\text{zbw}} = \left[ \frac{\lambda_c}{2} (1 \pm \varepsilon) \right] [\omega_{\text{zbw}} (1 \mp \varepsilon)] = c(1 - \varepsilon^2) \simeq c$$

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