# Tutorial: Octopus + BerkeleyGW

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- Supports a large set of Mean-Field codes: PARATEC, Quantum ESPRESSO, Octopus, PARSEC, SIESTA, EPM (TBPW)
- Supports 3D, 2D, 1D and Molecular Systems. Coulomb Truncation
- Support for Semiconductor, Metallic and Semi-Metallic Systems
- Efficient Algorithms and Use of Libraries. (BLAS, FFTW, LAPACK, SCALAPACK, FFTW) (OpenMP, FFTW3, HDF5 in BGW 1.1)
- Massively Parallel. Scales to 100,000 CPUs, distributed Memory.
- Efficient accurate solution to BSE via k-point Interpolation
- Support for LDA/GGA/Hybrid/HF/COHSEX starting points as well as offdiagonal Σ calculations







-The relative accuracy of Full-Frequency vs. Generalized Plasmon Pole (GPP) calculations is somewhat contentious.

$$\langle n\mathbf{k} | \Sigma_{\rm CH}(E) | n'\mathbf{k} \rangle = \frac{i}{2\pi} \sum_{n''} \sum_{\mathbf{q} \mathbf{G} \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}') \quad (20)$$
$$\times \int_0^\infty dE' \, \frac{\left[\epsilon_{\mathbf{G} \mathbf{G}'}^{\rm r}\right]^{-1}(\mathbf{q}; E') - \left[\epsilon_{\mathbf{G} \mathbf{G}'}^{\rm a}\right]^{-1}(\mathbf{q}; E')}{E - E_{n''\mathbf{k}-\mathbf{q}} - E' + i\delta} \quad v(\mathbf{q} + \mathbf{G}')$$

GPP is significantly faster, the integral over frequencies can be performed analytically if assume the dielectric response is dominated by a single plasmon pole.

BerkeleyGW supports both. With full-frequency you can compute spectral functions, lifetimes and weights.

# Practical issues

- 1. Screening models for Epsilon
- 2. Construction of **k**-grids
- 3. Symmetry and degeneracy
- 4. Real and complex version
- 5. Convergence

### Screening models: How do we use $\epsilon$ ?



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Sigma integrates over **q** with  $\varepsilon^{-1}(\mathbf{q})$ 

$$\begin{aligned} \langle n\mathbf{k} | \Sigma(E) | n'\mathbf{k} \rangle &= \frac{i}{2\pi} \sum_{n''} \sum_{\mathbf{q} \mathbf{G} \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}') \\ &\times \int_{-\infty}^{\infty} dE' e^{-i\delta E'} \frac{[\epsilon_{\mathbf{G} \mathbf{G}'}]^{-1}(\mathbf{q}; E')}{E - E_{n''\mathbf{k}-\mathbf{q}} - E' - i\delta_{n''\mathbf{k}-\mathbf{q}}} v(\mathbf{q} + \mathbf{G}') \end{aligned}$$

Absorption interpolates kernel over **q** with  $W(\mathbf{q}) = \varepsilon^{-1}(\mathbf{q}) v(\mathbf{q})$ 

$$\langle vc\mathbf{k}|K^{\mathrm{d}}|v'c'\mathbf{k}'\rangle = \sum_{\mathbf{G}\mathbf{G}'} M_{c'c}^{*}(\mathbf{k},\mathbf{q},\mathbf{G})W_{\mathbf{G}\mathbf{G}'}(\mathbf{q};0)M_{v'v}(\mathbf{k},\mathbf{q},\mathbf{G}')$$

## Problem 1: Non-smooth behavior



### **Problem 2: Divergent behavior**



# Solution: Screening models

Calculate at  $\mathbf{q}_0 \approx 0.001$  in periodic direction use to parametrize screening model

Sigma: Integrate over region around **q** = 0



$\epsilon_{{f G}{f G}'}^{-1}$	head	wing	wing'	body
Semiconductor	const	q	$\mathbf{q}/q^2$	$\operatorname{const}$
Metal	$q^2$	$q^2$	$\operatorname{const}$	$\operatorname{const}$
$W_{\mathbf{G}\mathbf{G}'}$	head	wing	wing'	body
Semiconductor	$1/q^{2}$	$\mathbf{q}/q^2$	$\mathbf{q}/q^2$	$\operatorname{const}$
Metal	const	$\operatorname{const}$	$\operatorname{const}$	$\operatorname{const}$

# Truncation of Coulomb potential

• GW and BSE utilize the Coulomb and screened Coulomb interaction

$$W = \varepsilon^{-1} V_c$$

• Long-range interactions make it computationally infeasible to increase lattice vectors until periodic images do not interact

Truncation Schemes within BerkeleyGW

- Cell box: 0D
- Cell wire: 1D
- Cell slab: 2D
- Spherical: Define radius of truncation
- **Cell truncation**: at half lattice vector length
  - Analytical form for Coulomb potential in k-space
- **Spherical truncation**: convenient, available in many packages

$$v_{
m t}({f r})=rac{\Theta(f({f r}))}{r}$$

Example: BN sheets



0

0

0.1

0.2

Convergence improved with truncation

Ismail-Beigi PRB 73 233103 (2006)

|q| (a.u.<sup>-1</sup>)

0.4

0.5

0.6

0.7

0.3

# Regular k-grids

Epsilon  

$$\chi_{\mathbf{GG'}}(\mathbf{q};0) = \sum_{n}^{\mathrm{occ}} \sum_{n'}^{\mathrm{emp}} \sum_{\mathbf{k}} M_{nn'}^*(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G'}) \frac{1}{E_{n\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}}}.$$

where

$$M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G}) = \langle n\mathbf{k} + \mathbf{q} | e^{i(\mathbf{q} + \mathbf{G}) \cdot \mathbf{r}} | n'\mathbf{k} \rangle$$

Sigma

$$\langle n\mathbf{k} | \Sigma_{\mathrm{SX}}(E) | n'\mathbf{k} \rangle = -\sum_{n''}^{\mathrm{occ}} \sum_{\mathbf{q} \mathbf{G} \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}') \\ \times [\epsilon_{\mathbf{G} \mathbf{G}'}]^{-1} (\mathbf{q}; E - E_{n''\mathbf{k}-\mathbf{q}}) v(\mathbf{q} + \mathbf{G}')$$

Kernel

$$\begin{aligned} \langle vc\mathbf{k} | K^{\mathrm{d}} | v'c'\mathbf{k}' \rangle &= \sum_{\mathbf{G}\mathbf{G}'} M^*_{c'c}(\mathbf{k}, \mathbf{q}, \mathbf{G}) W_{\mathbf{G}\mathbf{G}'}(\mathbf{q}; 0) M_{v'v}(\mathbf{k}, \mathbf{q}, \mathbf{G}') \\ &\epsilon^{-1}(\mathbf{q}) \text{ for } \mathbf{q} = \mathbf{k} - \mathbf{k}'. \end{aligned}$$

# **k**-grids and bands

#### recommended approach

	k-grid	# bands	Comments
SCF	Uniform, 0.5 shift	occupied	as usual in DFT
WFN	Uniform, 0.5 shift	many	
WFNq	WFN + <b>q</b> -shift	occupied	
epsilon.inp <b>q</b> -points	WFN but no shift, <b>q</b> <sub>0</sub>	many	bands to sum over
WFN_inner	WFN but no shift	many	bands to sum over
sigma.inp <b>k</b> -points	subset of WFN_inner	few	can choose to calculate Sigma just for bands of interest
WFN_co	WFN_inner	few	
WFN_fi (absorption)	Uniform, random shift	few	
WFNq_fi	WFN_fi + <b>q</b> -shift	occupied	
WFN_fi (inteqp)	anything	few	whatever is of interest

### epsilon.inp



#### **k**-grid construction: 4×4 grid for graphene



#### **k**-grid construction: 4×4 grid for graphene



### Degeneracy

Epsilon, Sigma: symmetry of Hamiltonian

$$\langle n\mathbf{k} | \Sigma_{\mathrm{SX}}(E) | n'\mathbf{k} \rangle = -\sum_{n''}^{\mathrm{occ}} \sum_{\mathbf{q} \mathbf{G} \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}')$$
$$\times \left[ \epsilon_{\mathbf{G} \mathbf{G}'} \right]^{-1} (\mathbf{q}; E - E_{n''\mathbf{k} - \mathbf{q}}) v(\mathbf{q} + \mathbf{G}')$$

Absorption: symmetry of e-h basis

$$\left(E_{c\mathbf{k}}^{\rm QP} - E_{v\mathbf{k}}^{\rm QP}\right)A_{vc\mathbf{k}}^{S} + \sum_{v'c'\mathbf{k}'}\left\langle vc\mathbf{k}|K^{\rm eh}|v'c'\mathbf{k}'\right\rangle = \Omega^{S}A_{vc\mathbf{k}}^{S}$$

Summing over only some of a degenerate space will break symmetry.

Degeneracy in mean-field => broken in GW!

Results depends on arbitrary linear combinations in mean-field. Not reproducible! Incorrect oscillator strengths in absorption!

### Degeneracy check utility

```
$ degeneracy check.x WFN
Reading eigenvalues from file WFN
Number of spins:
                               1
Number of bands:
                              35
                              8
Number of k-points:
== Degeneracy-allowed numbers of bands (for epsilon and sigma) ==
            4
            8
           14
           18
           20
           32
Note: cannot assess whether or not highest band 35 is degenerate.
```

```
So, use number bands 32 in Epsilon.
```

# Real or complex flavor?

e.g. bin/epsilon.real.x, bin/epsilon.cplx.x

Complex is general, but real is faster, uses less memory and disk space

Real: only with inversion symmetry about the origin  $u(-\mathbf{r}) = au(\mathbf{r})$ 

and time-reversal symmetry  $u^{*}(\mathbf{r}) = bu(\mathbf{r})$ 

a, b each equal to  $\pm 1$ 

What breaks time-reversal? Magnetic fields, spin-polarization, spinors Plane-wave codes generally just use complex wavefunctions. Conditions for reality depends on the basis! Real-space: k = 0, time-reversal.

Real output not implemented in Octopus yet.

There are many convergence parameters in a GW calculations : convergence with each must be checked



# Coupled convergence parameters



B. Shih et al., ZnO

### Octopus interface to BerkeleyGW

Real space transformed to plane-waves for GW.

Can only produce complex wavefunctions currently.

Periodic systems must use orthogonal unit cells (*i.e.* not hcp, fcc, ...) so build a supercell to match this condition.

Can treat rigorously finite systems, unlike plane-wave codes.

Domain parallelization for real-space scales better than plane waves.

Not yet in released version of Octopus! Hence, warnings about experimental features.

# The tutorial

Two examples: (1) hexagonal boron nitride sheet (2) benzene molecule

Today: GW Tomorrow: Bethe-Salpeter equation

Instructions at <a href="http://www.tddft.org/programs/octopus/wiki/index.php/Tutorial:BerkeleyGW">http://www.tddft.org/programs/octopus/wiki/index.php/Tutorial:BerkeleyGW</a>

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