Semilinear structurally damped wave equations: supercritical case

Well-posedness and existence of finite number of determining Fourier modes

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The problem

We consider the initial-boundary value problem

$$\begin{cases} u_{tt} - \Delta u + 2\beta(-\Delta)^{\alpha} u_{t} + f(u) = 0, & x \in \Omega, \ t > 0, \\ u|_{\partial\Omega} = 0, & (1) \\ u(x,0) = u_{0}(x), & u_{t}(x,0) = u_{1}(x), \ x \in \Omega, \end{cases}$$

where $\alpha \in [0,1]$, $\beta > 0$ are given numbers, and $\Omega \subset \mathbb{R}^n$ is a bounded domain with smooth enough boundary $\partial \Omega$.

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$$(-\Delta)^{\alpha}e_{i}=\lambda_{i}e_{i}, i\geq 1,$$

where $\{\lambda_i\}_{i\geq 1}$, $\{e_i\}_{i\geq 1}$ are the eigenvalues and the eigenvectors of the Laplacian operator $-\Delta$ under Dirichlet boundary conditions.



The linear problem - The spectrum

The roots of the characteristic polynomial for the linear problem are:

$$\gamma_n^{\pm} := -\beta \lambda_n^{\alpha} \pm \left(\beta^2 \lambda_n^{2\alpha} - \lambda_n\right)^{\frac{1}{2}} < 0, \quad n \ge 1.$$

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$$\gamma_n^+ \gamma_n^- = \lambda_n,$$

$$\gamma_n^+ = -\frac{\lambda_n}{\beta \lambda_n^\alpha + (\beta^2 \lambda_n^{2\alpha} - \lambda_n)^{\frac{1}{2}}}.$$

The linear problem - Controllability properties

Gets worse as $\alpha \nearrow 1$. In fact, when $\alpha = 1$

$$\gamma_n \to \frac{1}{2\beta}$$
, as $n \nearrow \infty$.

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Another fact we see for $\alpha \in [\frac{1}{2}, 1]$ and large n:

$$\gamma_n^+ \le C\lambda_n^{1-\alpha} \le C\lambda_n^{\frac{1}{2}}.$$

This and Weyl's theorem imply

$$\sum \frac{1}{\gamma_n^+} = \infty.$$

 \Rightarrow No spectral controllability via a scalar control through a profile (Müntz theorem).

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- Uniqueness:
- if $q \in (0,4]$ and $\alpha \in [0,\frac{1}{2}]$. (A. Savostianov, Adv.D.E., 2014)

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Uniqueness:

- if $q \in (0,4]$ and $\alpha \in [0,\frac{1}{2}]$. (A. Savostianov, Adv.D.E., 2014)
- $q \in (0, \frac{8\alpha}{3-4\alpha}]$ if $\alpha \in [\frac{1}{2}, \frac{3}{4})$, and q > 0 if $\alpha \in [\frac{3}{4}, 1]$. (V.

Kalantarov, S. Zelik, J.D.E., 2009)

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But, as a consequence

$$(u, u_t) \in L^{\infty}(0, T; \mathcal{E}),$$

that is the continuous evolution in the phase space $\mathcal E$ is not known for the supercritical exponents.

Finitely many determining modes

Definition

We say that the problem (1) has finitely many determining modes in \mathcal{E} , if for any two solutions u, v, that satisfy

$$\|P_N w(t)\|_{\mathcal{E}} o 0 \text{ as } t o +\infty,$$

we have

$$||w(t)||_{\mathcal{E}} \to 0 \text{ as } t \to +\infty,$$

where w = u - v, and P_N is the projection operator on to the space spanned by the first N eigenfunctions of the Laplacian under Dirichlet b.c.'s.

Finitely many weakly determining modes

Theorem

Suppose that $q \in (0, \frac{8\alpha}{3-4\alpha}]$ if $\alpha \in [\frac{1}{2}, \frac{3}{4})$, and q > 0 if $\alpha \in [\frac{3}{4}, 1]$. Then, for any two solutions u, v, that satisfy

$$\|P_N w(t)\|_{\mathcal{E}} \to 0 \text{ as } t \to +\infty,$$

we have

$$\|w(t)\|_{\mathcal{E}_{\alpha}} \to 0 \text{ as } t \to +\infty,$$

and

$$\int\limits_{t}^{t+1}\|w(au)\|_{\mathcal{E}}d au o 0 \ ext{as } t o +\infty.$$

Thank You!