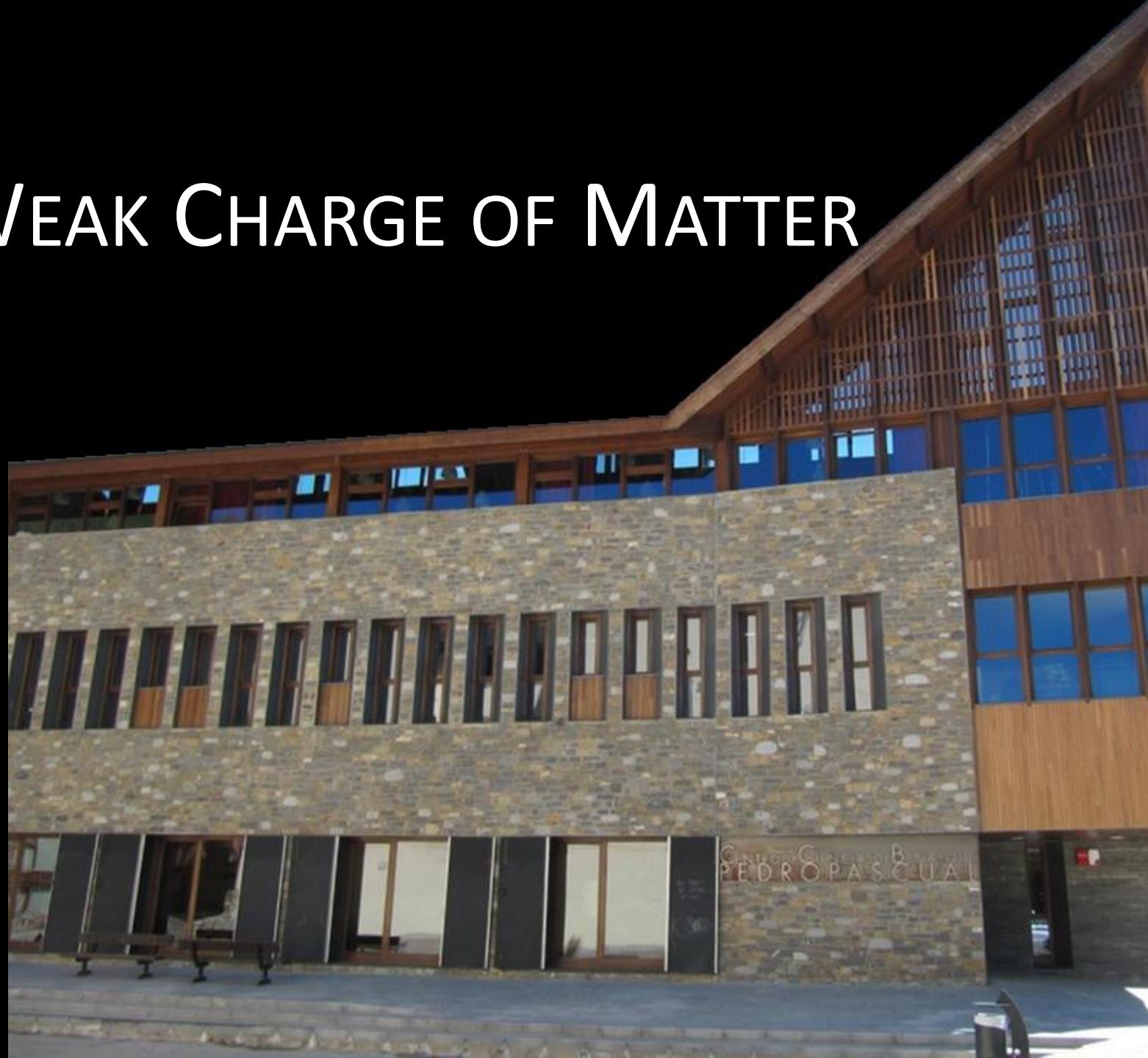


# THE COHERENT WEAK CHARGE OF MATTER

Alejandro Segarra  
IFIC (Univ. Valencia - CSIC)



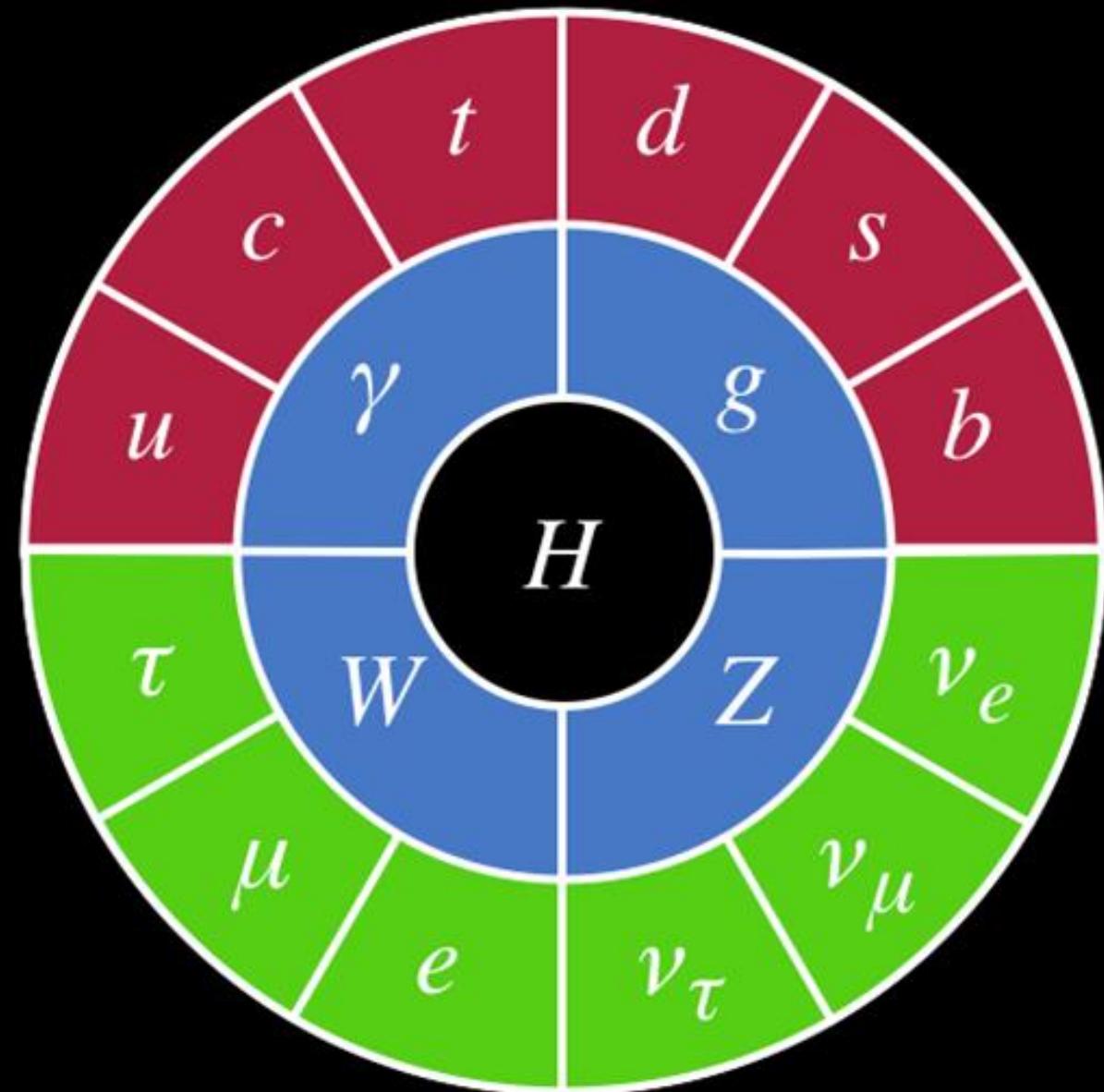
# LONG-RANGE WEAK FORCES

Alejandro Segarra  
IFIC (Univ. Valencia - CSIC)



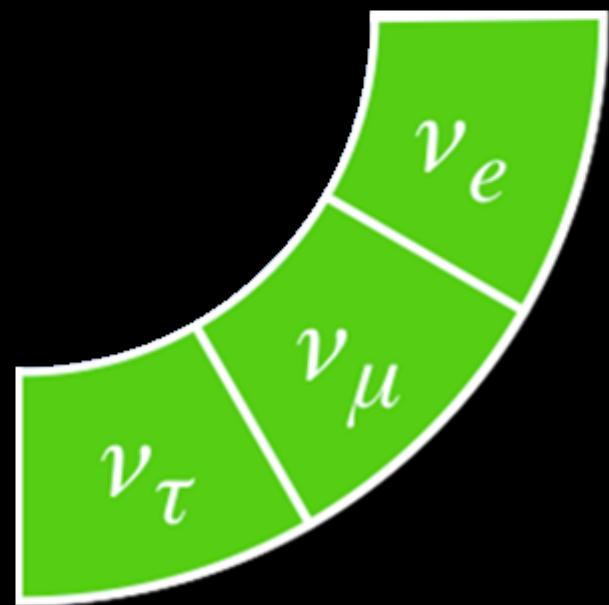
# MOTIVATION

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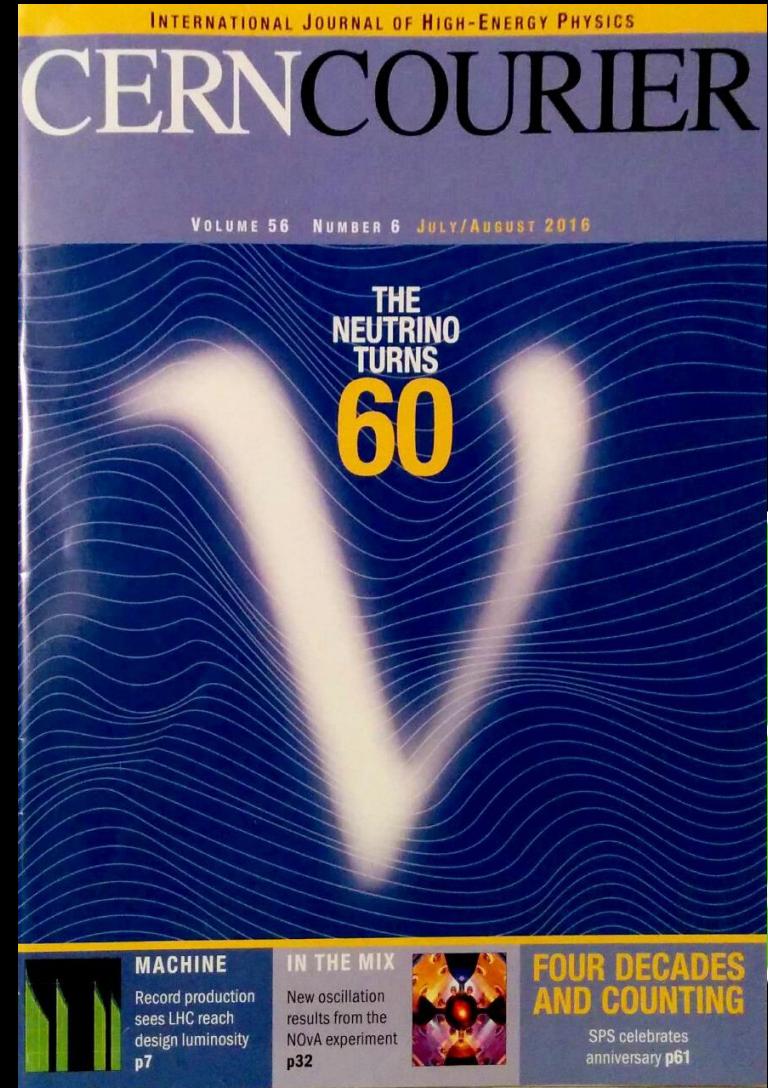
# MOTIVATION

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# MOTIVATION

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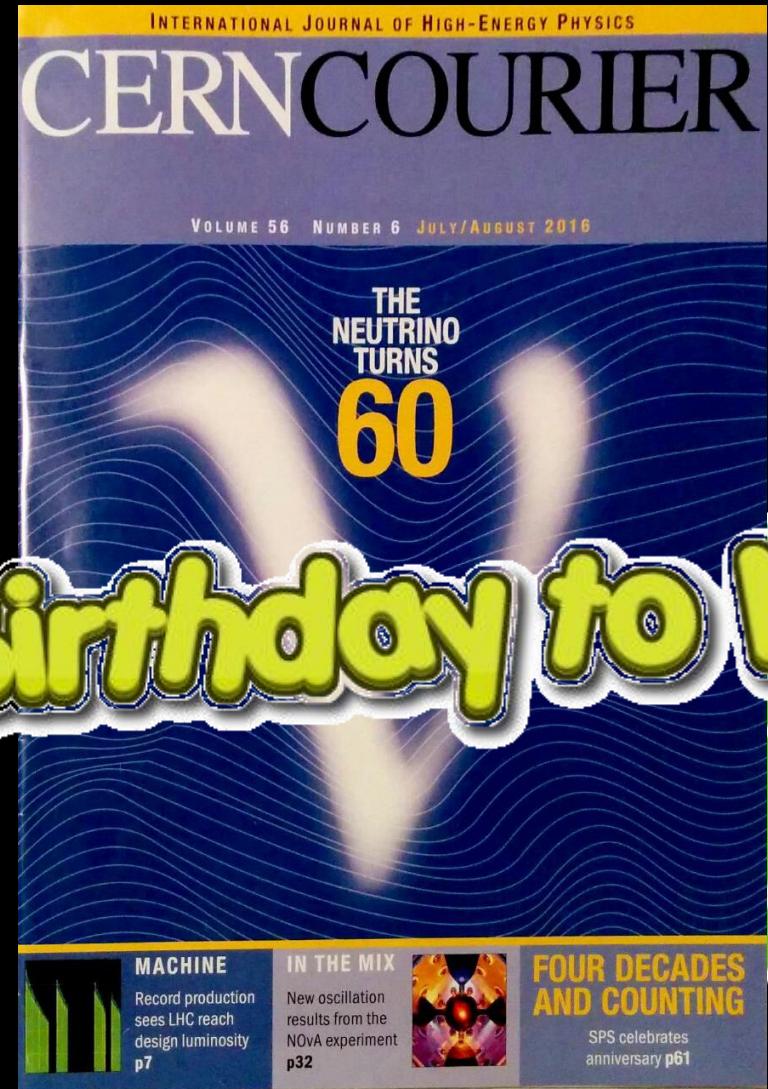


# MOTIVATION

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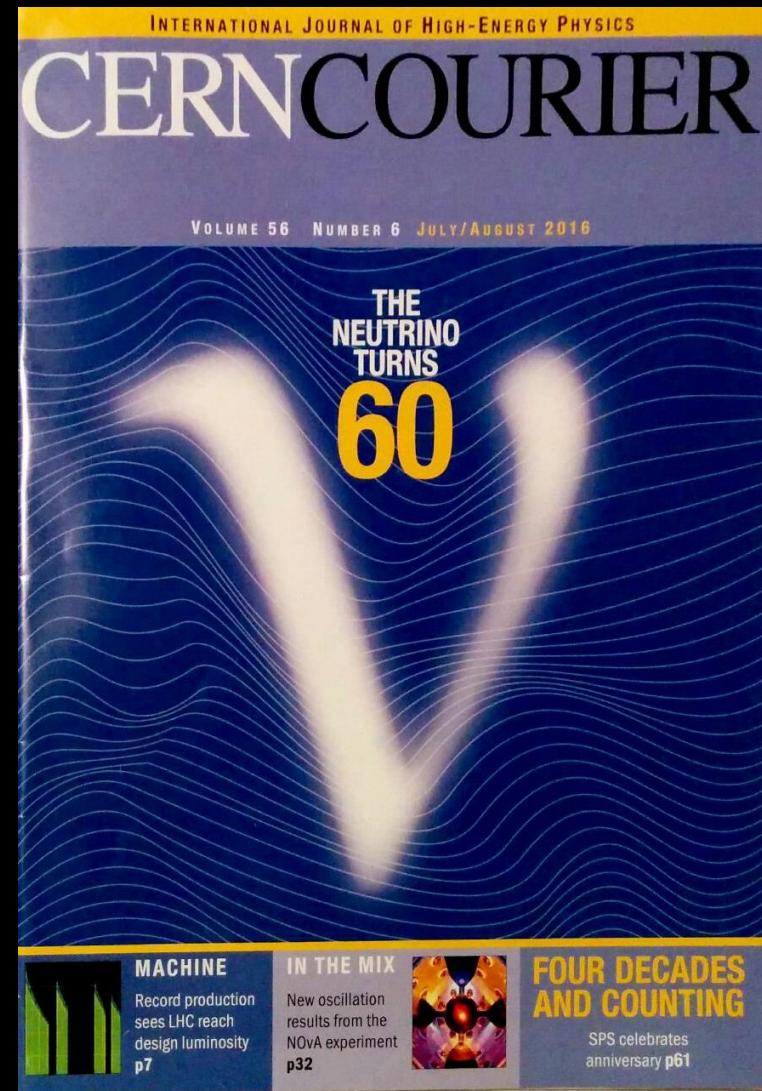
Happy Birthday to Nu



# MOTIVATION

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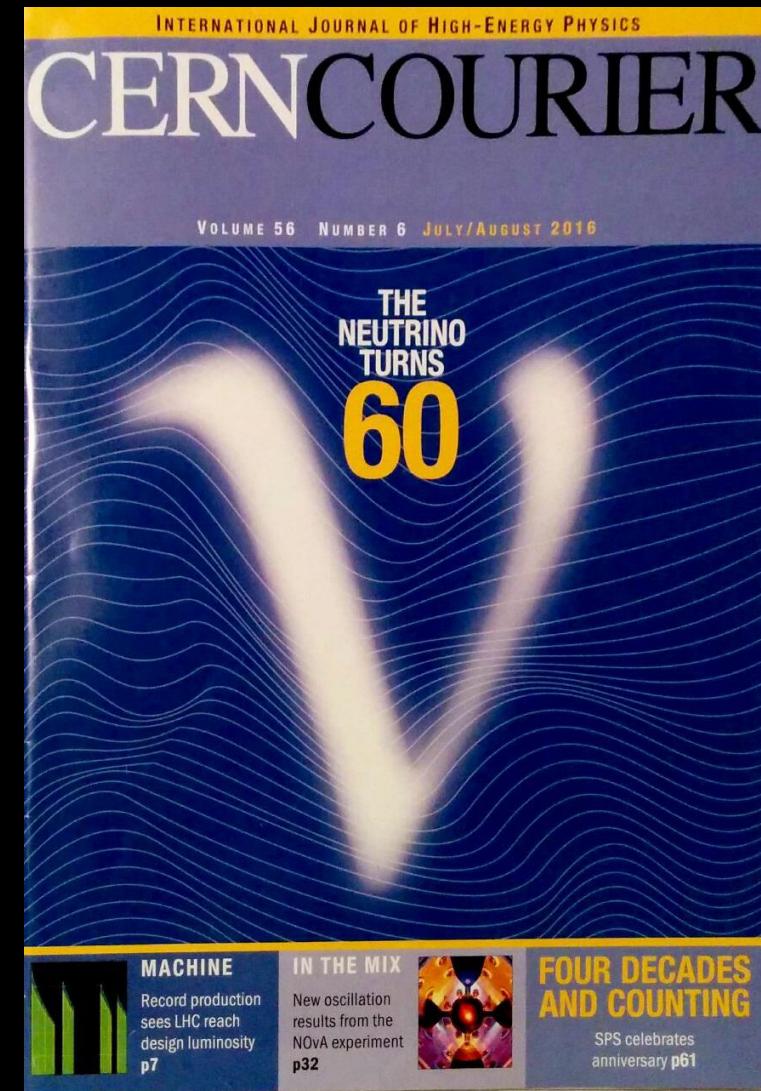
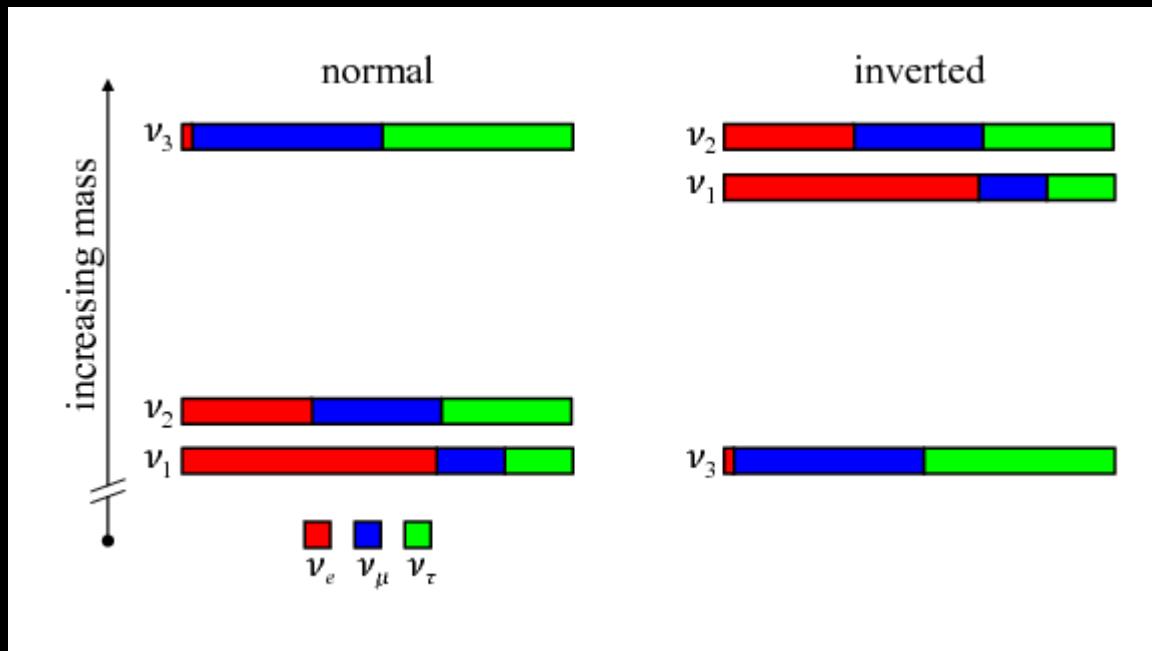
- Absolute mass scale: only  $\Delta m^2$



# MOTIVATION

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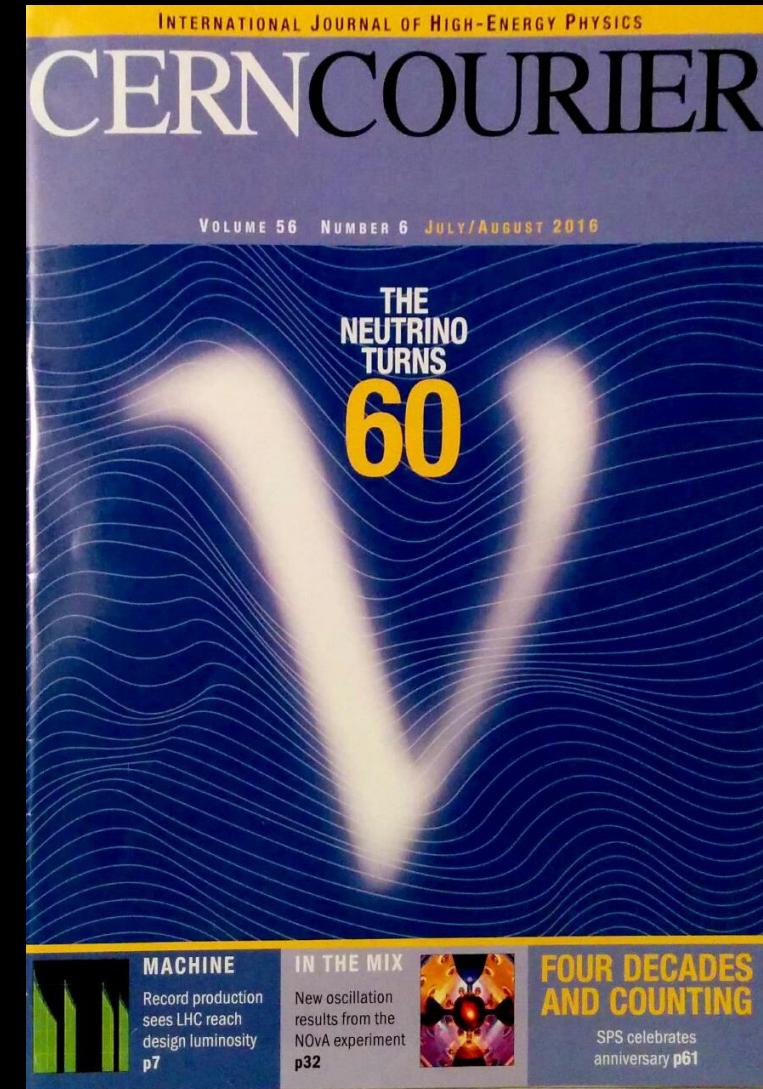
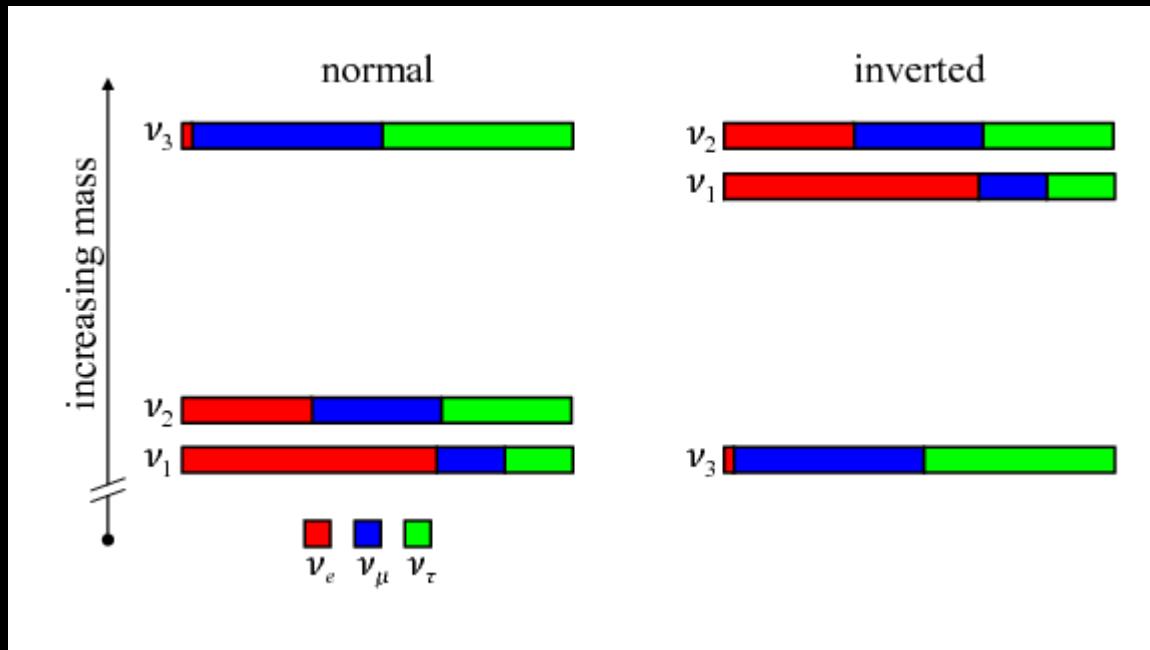
- Absolute mass scale: only  $\Delta m^2$
- Mass Hierarchy



# MOTIVATION

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- Absolute mass scale: only  $\Delta m^2$
- Mass Hierarchy
- Dirac/Majorana



# MOTIVATION

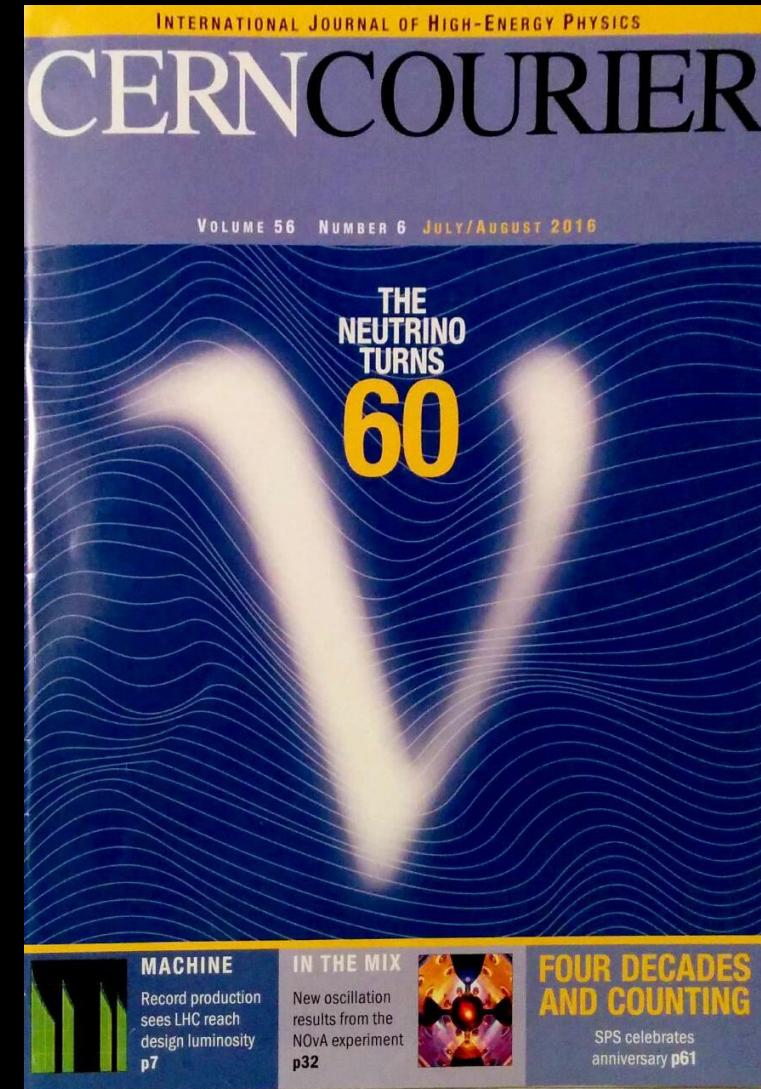
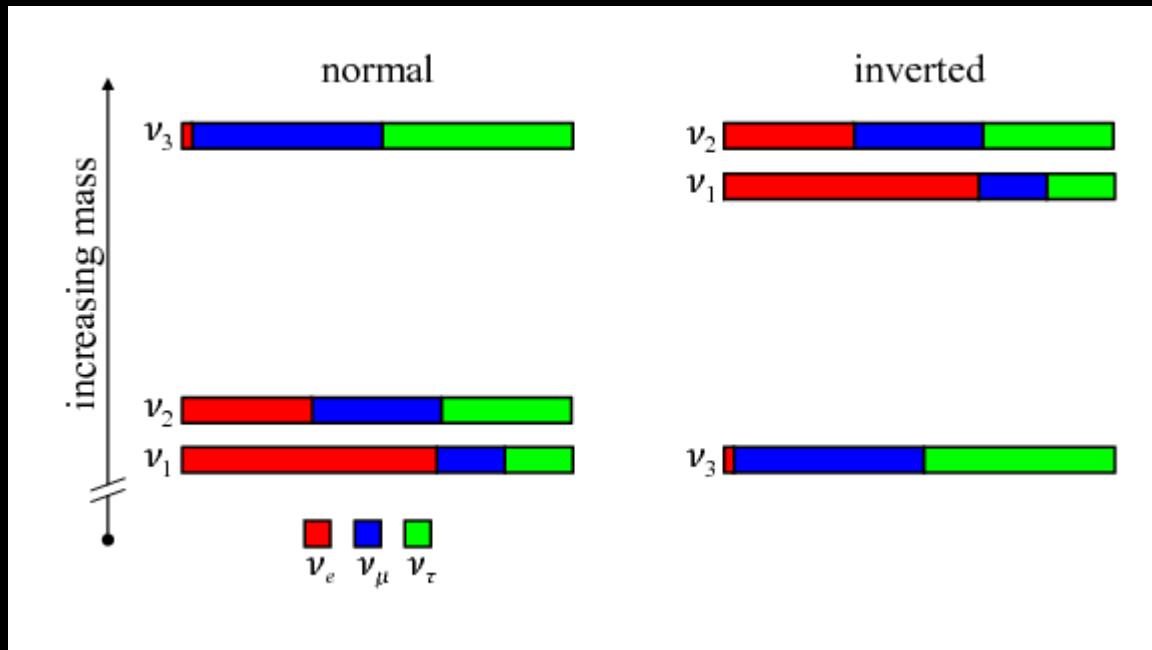
## Additional generations | lepton sector

- If neutrinos were massless we could redefine the (LH) fields  $\Rightarrow$  no lepton mixing  
But they have (tiny) masses because there are neutrino oscillations!
- Neutrinos are special:  
they *may* be their own antiparticle (Majorana) since they are neutral
- If they are Majorana:
  - Mass terms are different to Dirac case  
(neutrino and antineutrino *may* mix)
  - Intergenerational mixings are richer (more CP phases)

# MOTIVATION

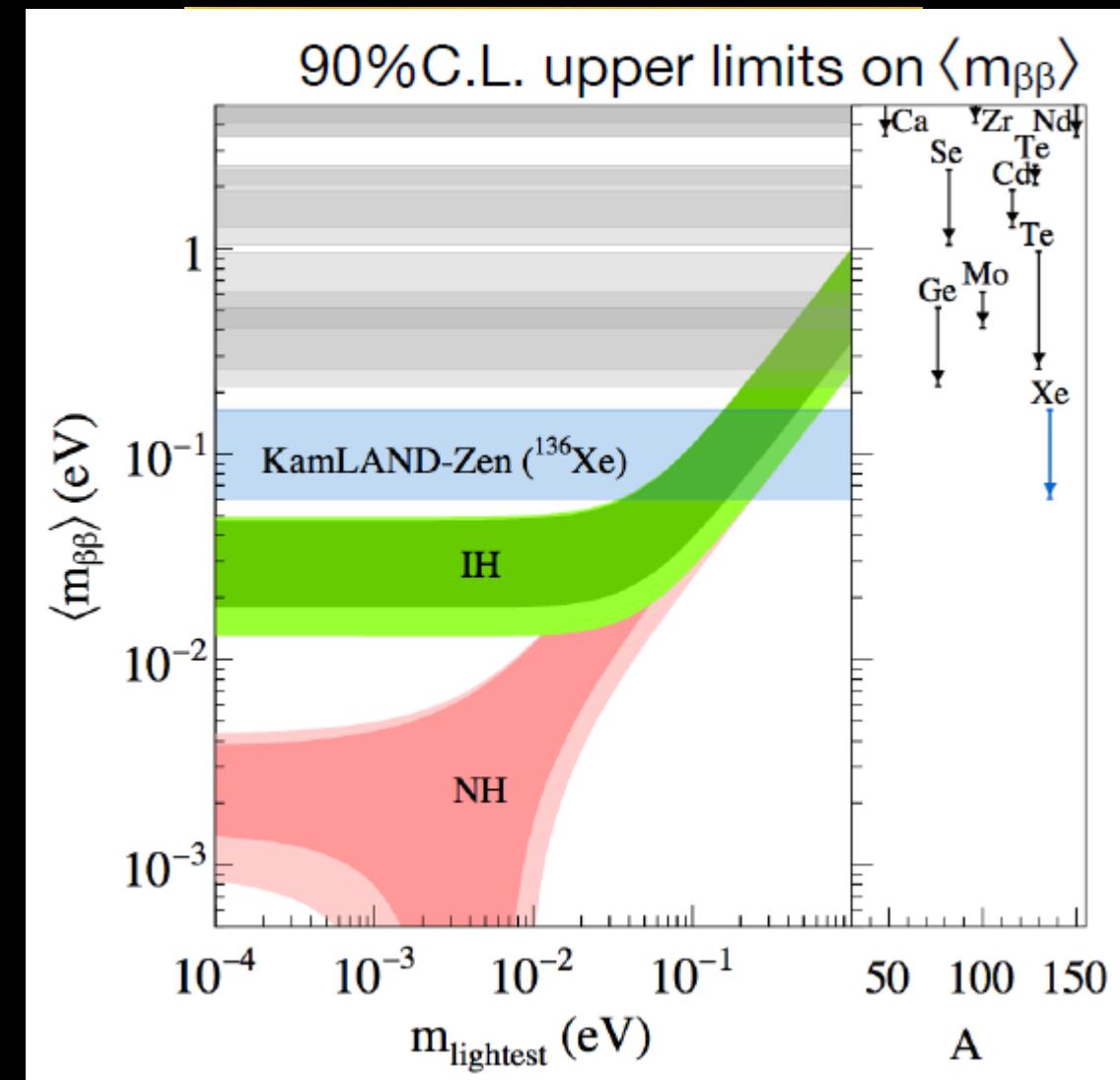
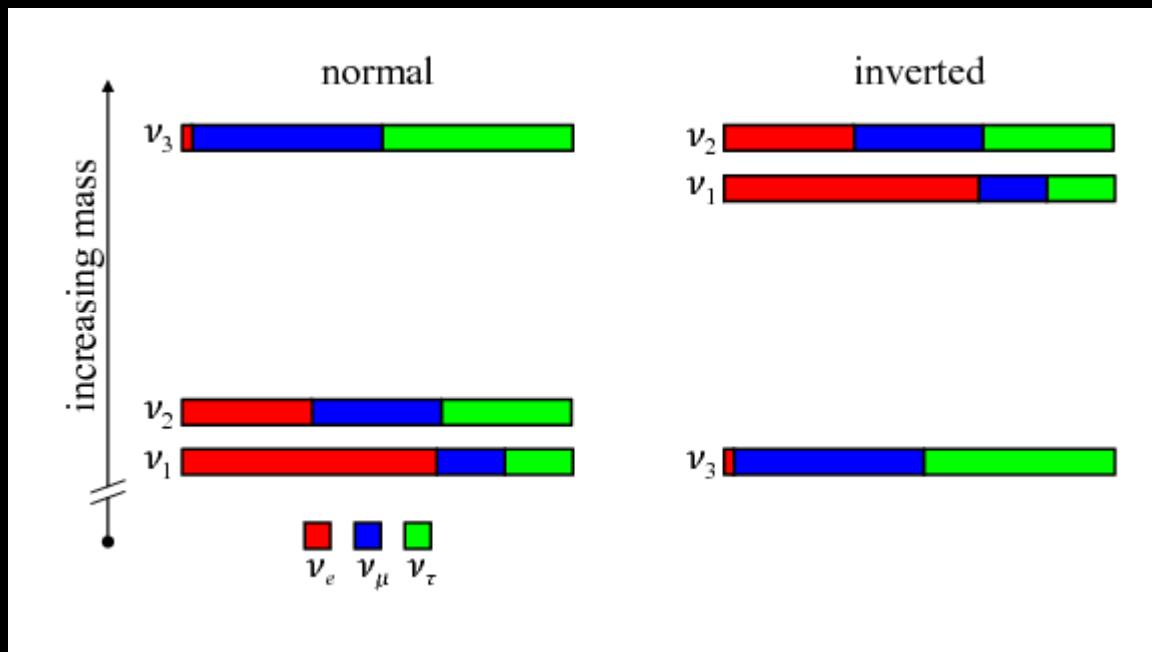
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- Absolute mass scale: only  $\Delta m^2$
- Mass Hierarchy
- Dirac/Majorana



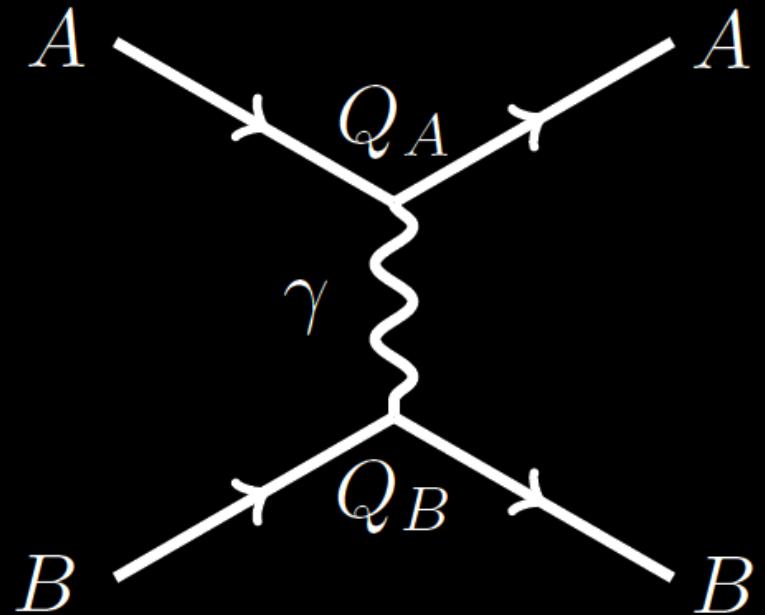
# MOTIVATION

- Absolute mass scale: only  $\Delta m^2$
- Mass Hierarchy
- Dirac/Majorana

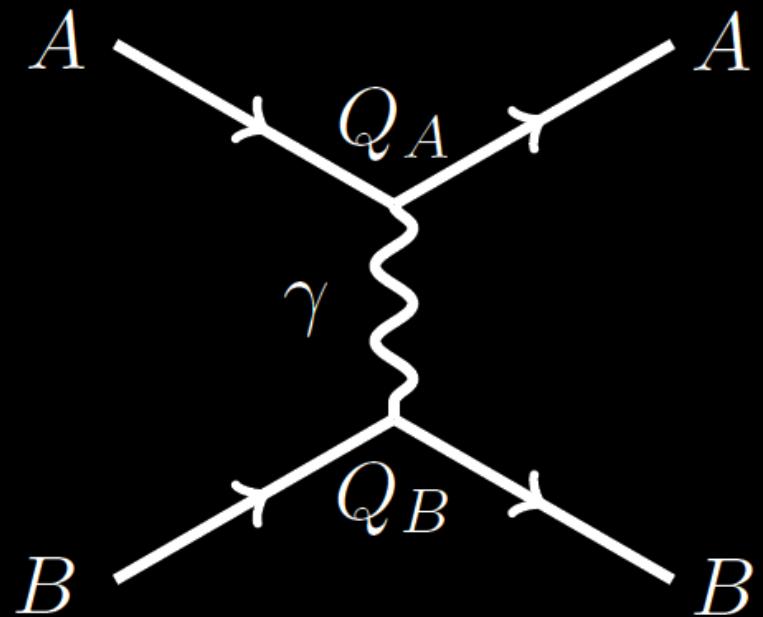


# QED

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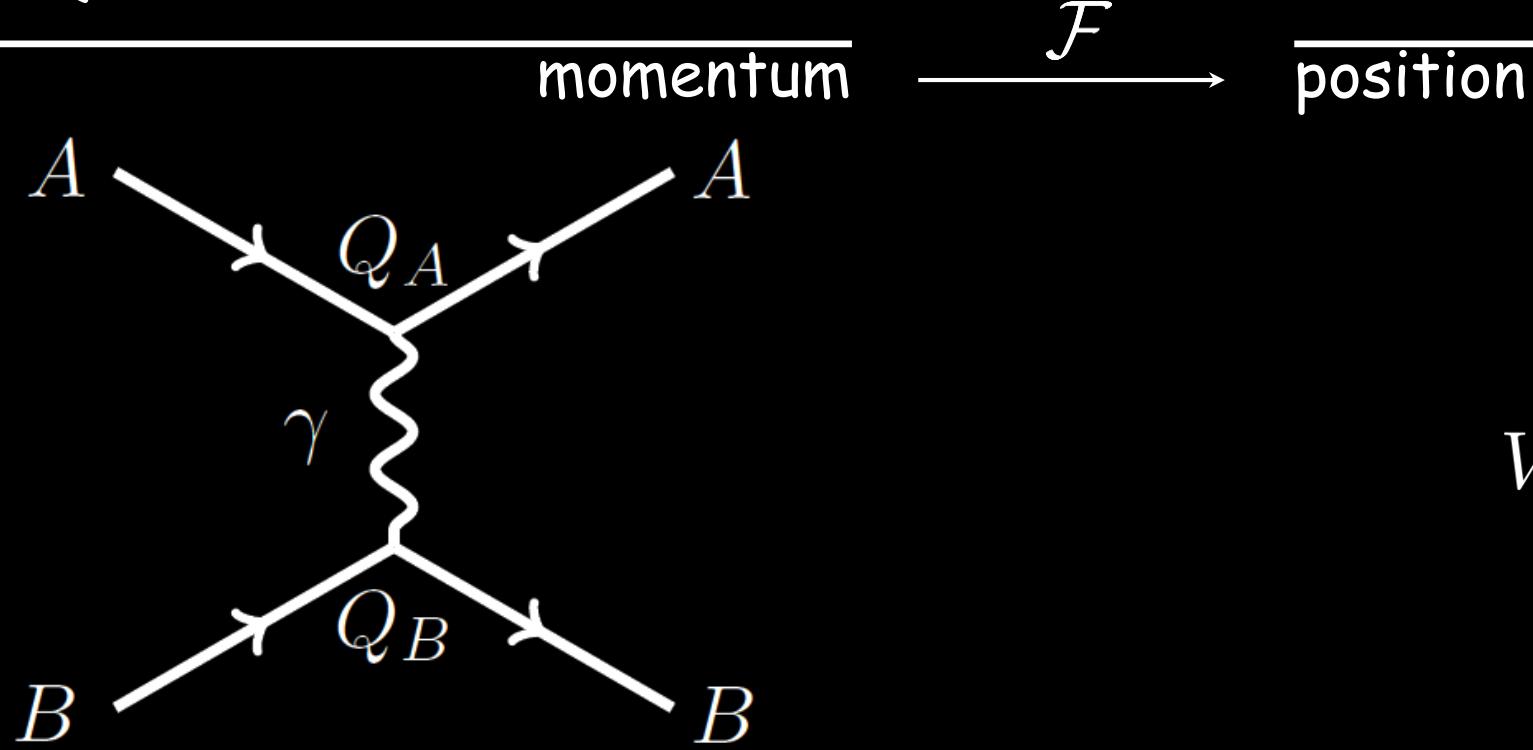
$$M(q) = e^2 Q_A Q_B \frac{\vec{j}_A \cdot \vec{j}_B}{q^2}$$



$$V(r) = \frac{e^2}{4\pi} \frac{Q_A Q_B}{r}$$

$$M(q) = e^2 Q_A Q_B \frac{\vec{j}_A \cdot \vec{j}_B}{q^2}$$

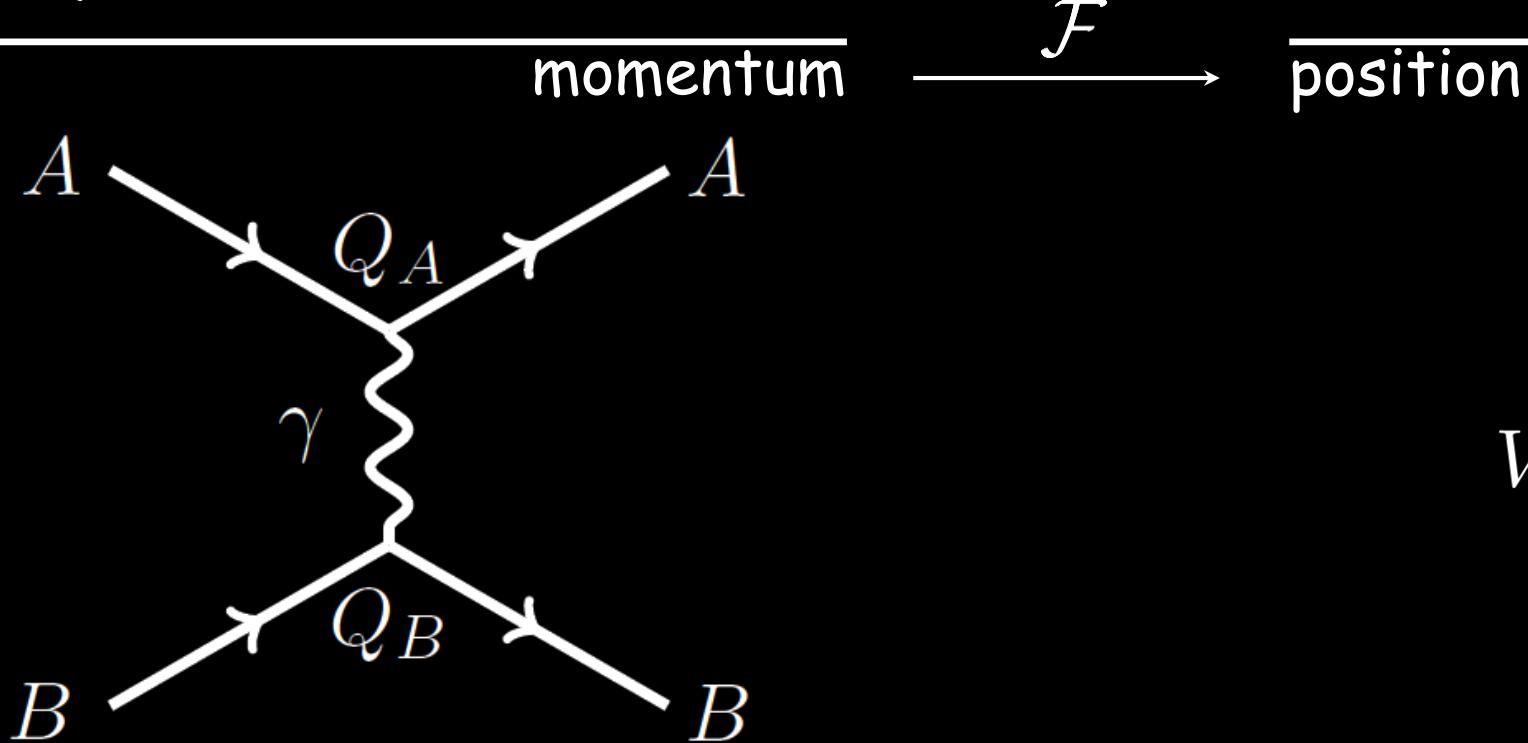
QED



COULOMB POTENTIAL

$$V(r) = \frac{e^2}{4\pi} \frac{Q_A Q_B}{r}$$

$$M(q) = e^2 Q_A Q_B \frac{j_A \cdot j_B}{q^2}$$



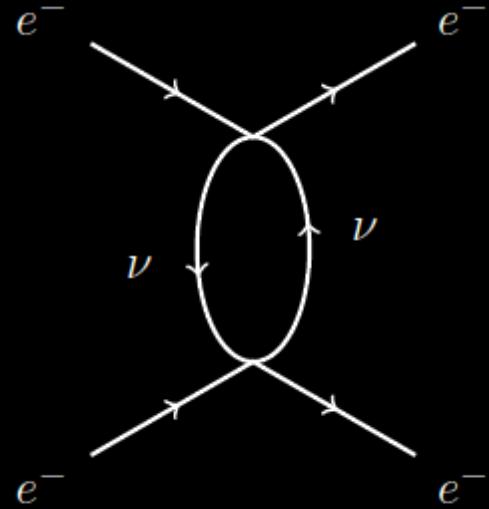
$$V(r) = \frac{e^2}{4\pi} \frac{Q_A Q_B}{r}$$

$$M(q) = e^2 Q_A Q_B \frac{j_A \cdot j_B}{q^2}$$

Photon is massless  $\longrightarrow$  range is  $\infty$

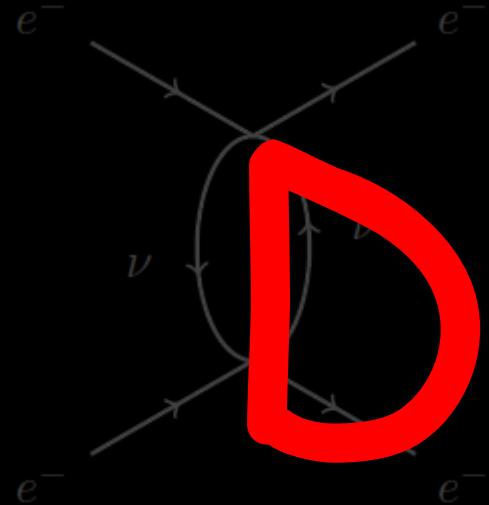
# ALREADY ON THE LITERATURE...

Hsu, Sikivie, arXiv:hep-ph/9211301



$$V(r) = \frac{G_F^2}{8\pi^3} \frac{1}{r^5}$$

- Only  $ee \rightarrow ee$  scattering considered
- SM: massless neutrinos

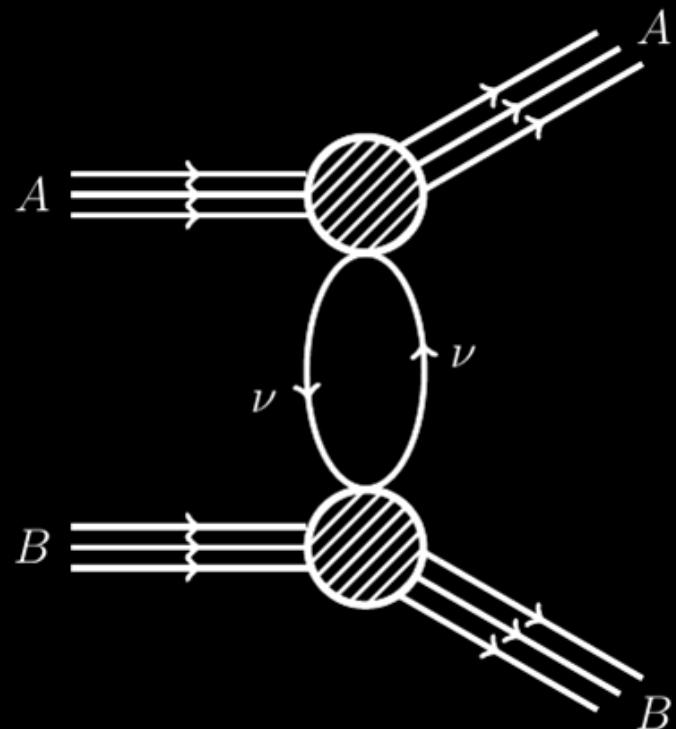


# Dismissed!

- Only  $ee \rightarrow ee$  scattering considered
- SM: massless neutrinos

# LONG-RANGE WEAK INTERACTION

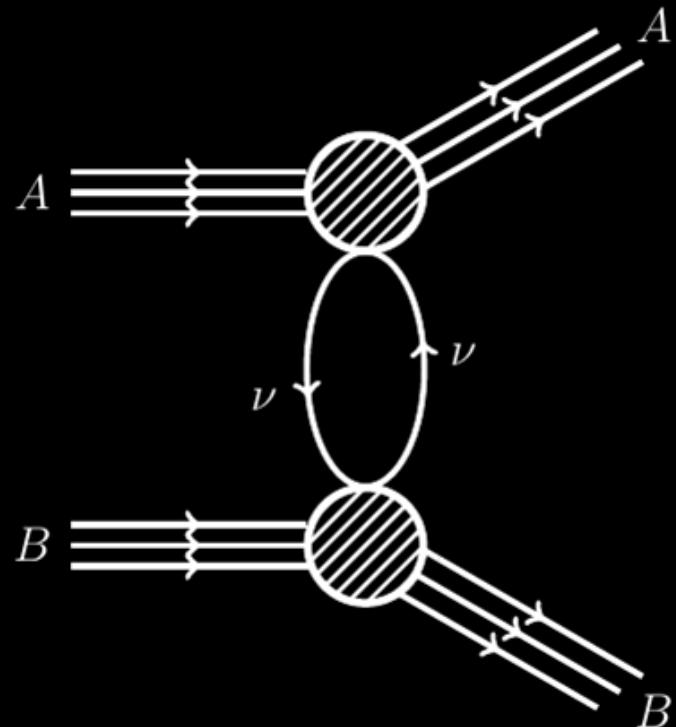
AS, arXiv:hep-ph/1606.05087



$$V(r) = \frac{G_F^2}{8\pi^3} \left[ (2Z - N)^2 + 2N^2 \right] \frac{1}{r^5}$$

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AS, arXiv:hep-ph/1606.05087



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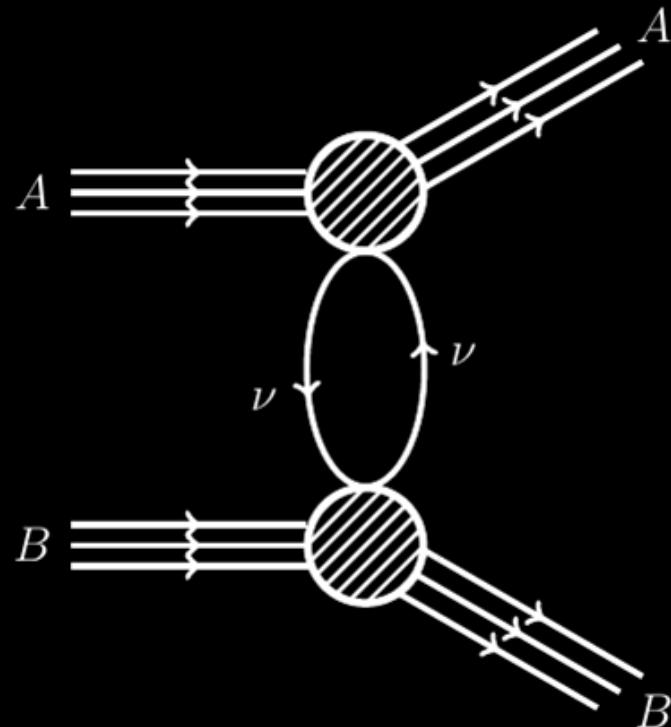
*A* and *B* are *anything*

$Z = \#$  protons  $= \#$  electrons

$N = \#$  neutrons

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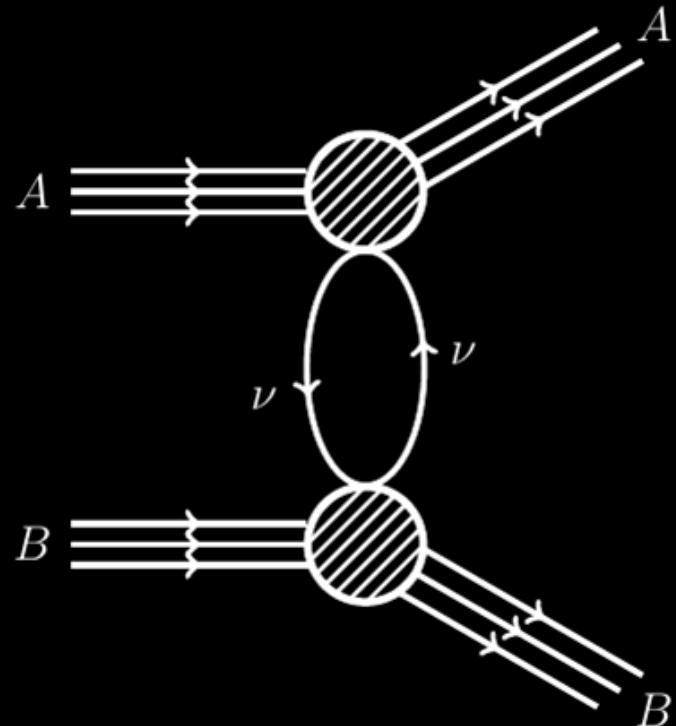
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Dispersion Relation

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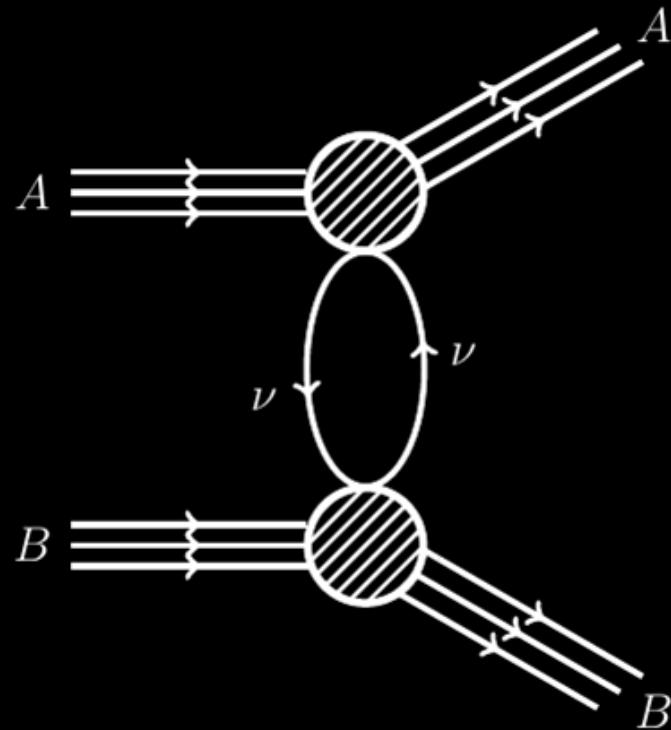
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Dispersion Relation  
 $S$ -matrix Unitarity

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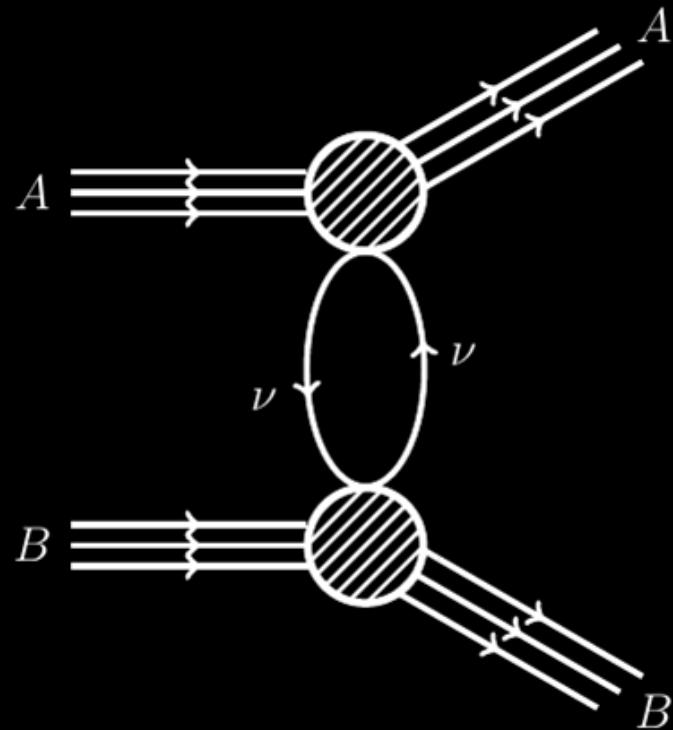
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Massless Neutrinos

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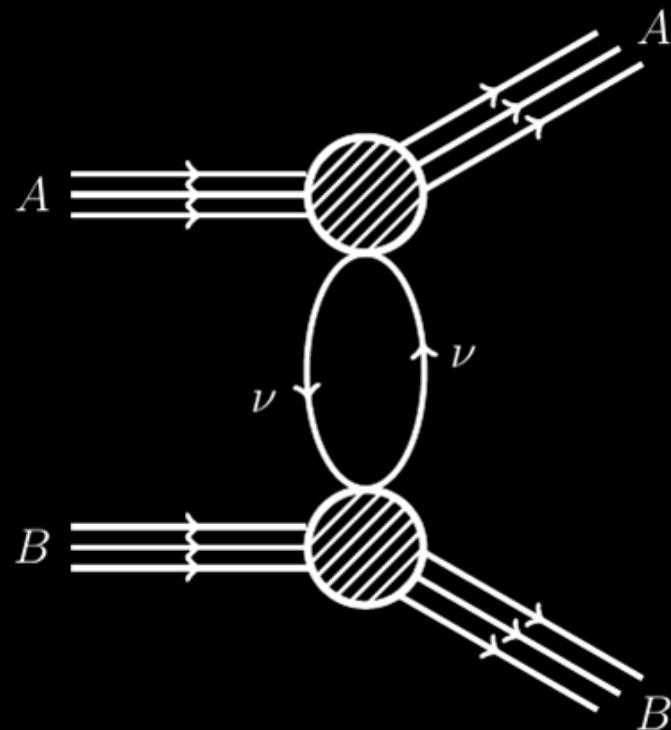
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Dispersion Relation  
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Massless Neutrinos  
Low-energy limit  
Coherent limit

# THE WEAK CHARGES OF MATTER

---

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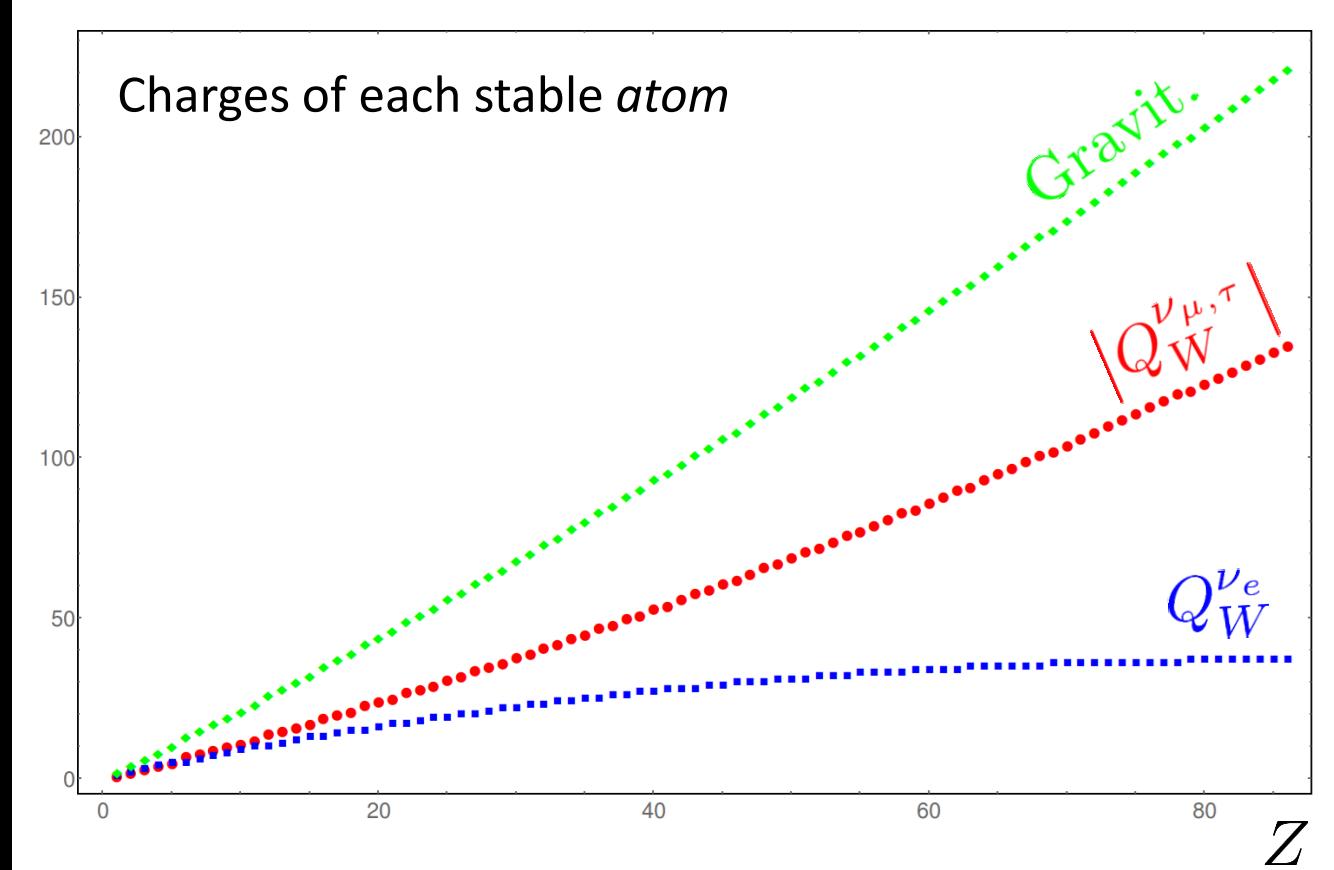
AS, arXiv:hep-ph/1606.05087

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# THE WEAK CHARGES OF MATTER (CONT'D)

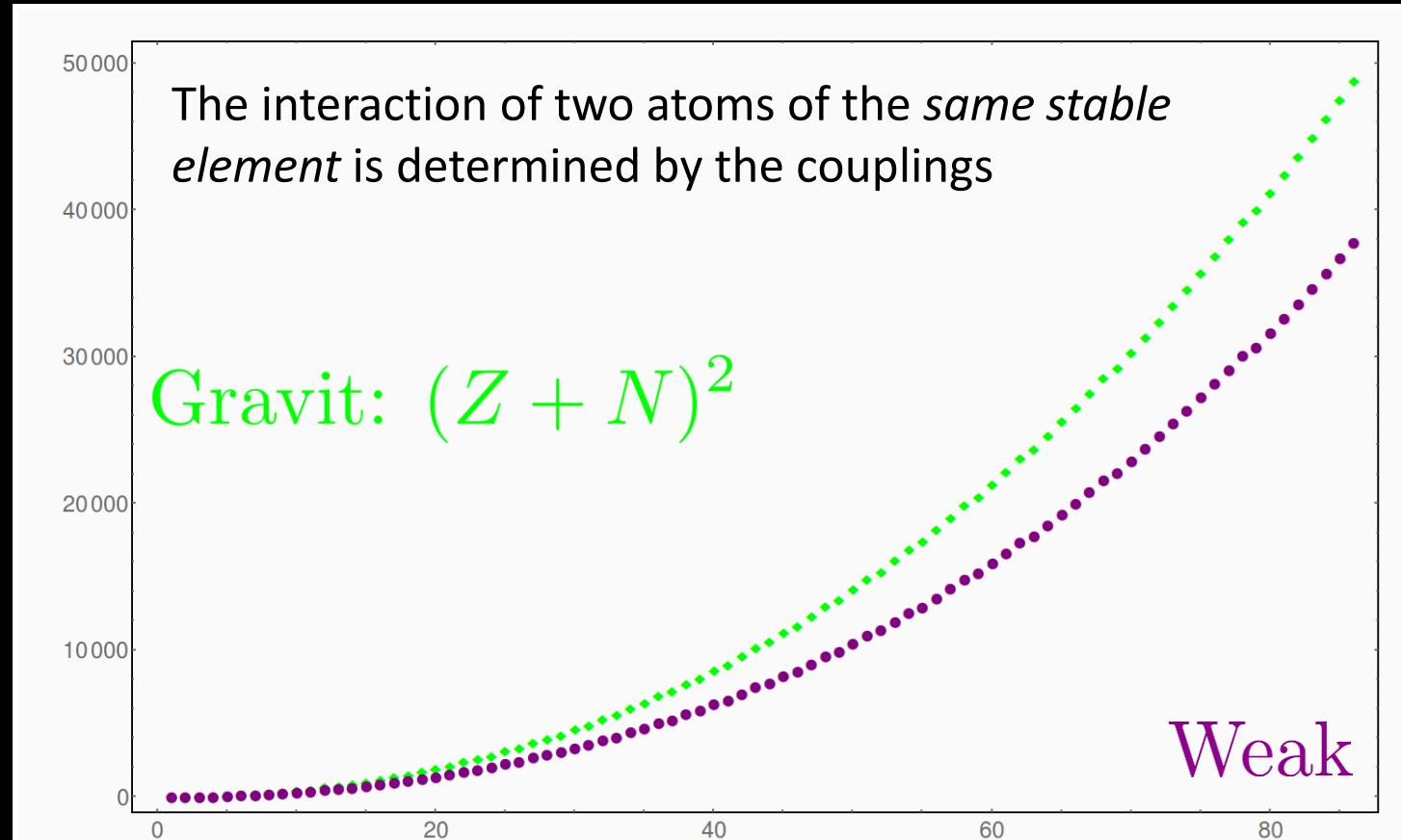
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# FINAL REMARKS

---

- ✓ Interesting ranges

$$r_{\min} \sim 1 \text{ nm}$$

$$r_{\max} \sim m_\nu^{-1} \sim (0.1 \text{ eV})^{-1} \sim 1 \mu\text{m}$$

$$\frac{a_0}{e^{-2mr}}$$

Need to recalculate with finite mass!

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Residual EM forces

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can be shielded  
repulsive!

deviations from Equivalence Principle

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Dirac or Majorana?

