

Cosmological interaction of vacuum energy and dark matter

Clara Hormigos Feliu

Universitat de Barcelona

14th September 2016



UNIVERSITAT DE
BARCELONA

Outline

- Introduction
- Approach
- Results
- Conclusions

Current densities

- Barionic matter: $\Omega_b \simeq 0.05$
- Dark matter (DM): $\Omega_c \simeq 0.24$
- Dark energy (DE): $\Omega_D \simeq 0.72$

WMAP Collaboration, 2013

Equation of state of DE ($p_D = \omega_D \rho_D$): $\omega_D = -1.006 \pm 0.045$

Planck Collaboration, 2015

Dark energy

Possible descriptions:

- Cosmological constant, Λ ($\omega_\Lambda = -1$)
 Λ CDM description
- Modified gravity
- Scalar fields:
 - Quintessence ($\omega_\phi > -1$)
 - Phantom fields ($\omega_\phi < -1$)
- **Dynamic vacuum**
Evolving Λ , $\omega_D(z)$

Dark energy as a dynamic vacuum

- Friedmann's equations take the known form:

$$3H^2 = 8\pi G(a)(\rho_m + \rho_r + \rho_\Lambda(a)) \quad (1)$$

$$3H^2 + 2\dot{H} = -8\pi G(a)\left(\frac{\rho_r}{3} - \rho_\Lambda(a)\right), \quad (2)$$

with $\rho_\Lambda(a) = \frac{\Lambda(a)}{8\pi G(a)}$.

- Need to postulate the evolution of $\rho_\Lambda(a)$
- Possible coupling of ρ_Λ to ρ_m and G

Approach

- Find evolution of densities, H and G
- Identify possible anomalous conservation laws

and also...

- Find an effective equation of state (EoS) for dark energy

Studied models

$$\rho_{\Lambda}(H; \nu, \alpha) = \frac{3}{8\pi G} \left(c_0 + \nu H^2 + \frac{2}{3} \alpha \dot{H} \right) \quad (3)$$

Solà *et al*, 2015 & 2016

- Type G: evolving G , Λ CDM evolution of ρ_m, ρ_r
- Type A: $G = \text{const.}$, modified evolution of ρ_m, ρ_r

$$\dot{\rho}_{\Lambda} = -q_V H \rho_{\Lambda} \quad (4)$$

Salvatelli *et al*, 2014

- Type S, S' (different fits)

Effective equation of state

- Comparison with scalar field EoS (not interacting with dark matter)
- Impose equality of the models' Hubble functions,

$$E^2(z) = \Omega_m^0 f_m(z, r_i) (1+z)^{\alpha_m} + \Omega_\Lambda^0 f_\Lambda(z, r_i), \quad (5)$$

with theoretical scalar field Hubble function:

$$E_D^2(z) = \tilde{\Omega}_m^0 (1+z)^{\alpha_m} + \tilde{\Omega}_D^0 \zeta(z). \quad (6)$$

- Obtain $\omega_D(z)$

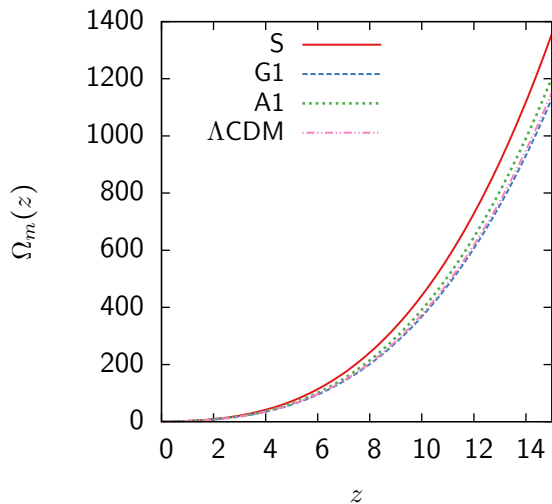
Basilakos & Solà, 2013

Conservation laws

- Type G: $\dot{G}(\rho_m + \rho_r + \rho_\Lambda) + G\dot{\rho}_\Lambda = 0$
 - ρ_m and ρ_r follow conventional conservation laws
 - Coupling to G is essential for a dynamic vacuum

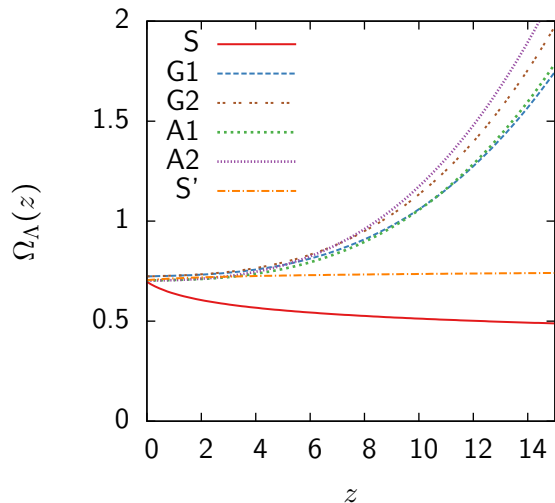
- Type A and S: $\dot{\rho}_m + \dot{\rho}_r + 3H(\rho_m + \rho_r) = -\dot{\rho}_\Lambda$
 - Anomalous matter conservation laws
 - Dynamic vacuum couples to matter

Normalized matter density



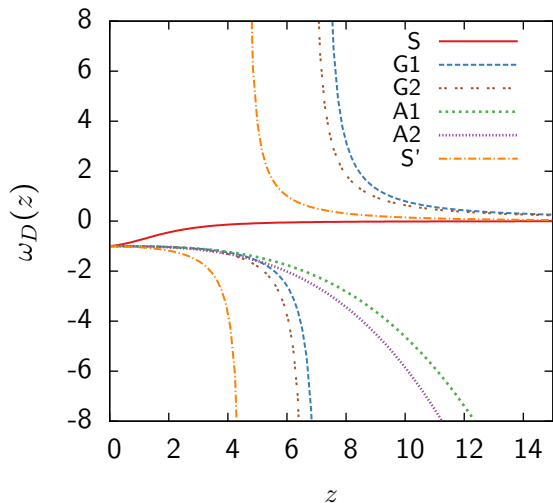
- Evolution close to Λ CDM

Normalised vacuum density



- Evolution differs from Λ CDM ($\Omega_\Lambda \simeq 0.7$)

Effective equation of state



Effective behavior:

- Phantom field-like: G, A (S')
- Quintessence-like: S

Conclusions

- Dynamical vacuum models provide an alternative explanation of dark energy with
 - No need for external fields
 - Significant improvement when fit to data
- Need to postulate an evolution of ρ_Λ
- Introducing a dynamical vacuum leads to anomalous conservation laws
- Effective EoS can differ from Λ CDM and explain possible $\omega_D \neq -1$

Thank you for your attention

Normalised vacuum density

Λ CDM:

$$\Omega_{\Lambda}(z) = \Omega_{\Lambda}^0 \simeq 0.7$$

Type G:

$$\Omega_{\Lambda}(z) = (1+z)^3 \left((1+z)^{-3\xi} \left(1 - \frac{\Omega_m^0}{\xi} \right) + \Omega_m^0 (\xi^{-1} - 1) \right)$$

Type A:

$$\Omega_{\Lambda}(z) = \Omega_{\Lambda}^0 + \Omega_m^0 (\xi^{-1} - 1) ((1+z)^{3\xi} - 1)$$

Type S:

$$\Omega_{\Lambda}(z) = \Omega_{\Lambda}^0 (1+z)^{q_V}$$

Akaike and Bayesian criteria with respect to Λ CDM

Model	Δ AIC	Δ BIC
RVM (G,A)	27.66	25.44
Q_Λ (S)	17.13	14.91

Solà, De Cruz Pérez, Gómez-Valent, Nunes (2016)

Δ AIC, Δ BIC = 6 – 10 \rightarrow "strong evidence" against Λ CDM

Δ AIC, Δ BIC > 10 \rightarrow "very strong evidence" against Λ CDM

Evolution of ρ_Λ

- "Running constant", QFT

$$\frac{d\rho_\Lambda}{d\ln\mu^2} = \frac{1}{4\pi} \sum_i \left[a_i M_i^2 \mu^2 + b_i \mu^4 + \frac{c_i \mu^6}{M_i^2} + \dots \right] \quad (7)$$

$$\mu \sim H$$

Solà, 2016. arXiv:1601.01668v2

FLRW metric:

$$ds^2 = -dt^2 + a^2(t) \left(\frac{dr^2}{1 - Kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right) \quad (8)$$

Field equations:

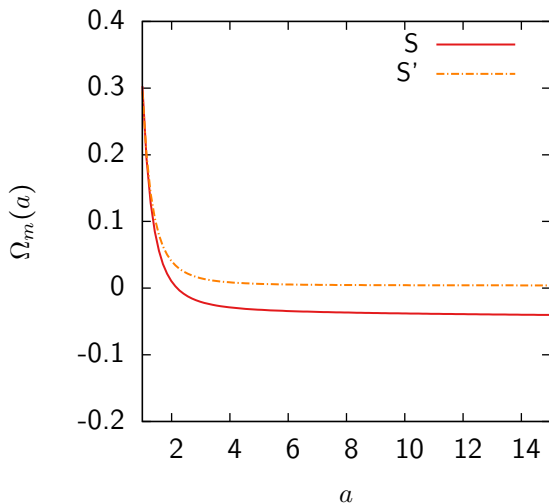
$$R_{\mu\nu} - \frac{1}{2}R = 8\pi G \tilde{T}_{\mu\nu}, \quad \tilde{T}_{\mu\nu} \equiv T_{\mu\nu} + g_{\mu\nu} \rho_\Lambda \quad (9)$$

Bianchi identities:

$$\nabla^\mu (G \tilde{T}_{\mu\nu}) = 0 \quad (10)$$

Local conservation law:

$$\frac{d}{dt}(G\rho) + 3GH(\rho + p) = 0 \quad (11)$$



Evolution of the normalized non-relativistic matter density with the scale factor for model S.