



Vacuum expectation value of the surface energy-momentum tensor in braneworlds for a complex scalar field

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Bulk + surface action

In GR, when the underlying manifold has boundaries, correct variational problem is achieved by supplementing the action with **Gibbons-Hawking** surface term

$$S_{GR} = \frac{1}{2} \int d^{D+1}x \sqrt{|g|} R \quad \longrightarrow \quad S_{GH} = -\int d^D x \sqrt{|h|} \varepsilon K$$

A complex scalar field on the background of a manifold with boundaries
The standard bulk action with coupling (to the background curvature) constant

$$S_b = \frac{1}{2} \int d^{D+1}x \sqrt{|g|} (\nabla_i \varphi \nabla^i \varphi^* - m^2 \varphi \varphi^* - \xi R \varphi \varphi^*)$$

The surface action with the analog of Gibbons-Hawking term and surface mass term

$$S_{s} = -\int d^{D}x \sqrt{|h|} \varepsilon (\xi K \varphi \varphi^{*} + m_{s} \varphi \varphi^{*}) \qquad h_{ik} = g_{ik} - \varepsilon n_{i} n_{k}$$
$$K_{ik} = h_{i}^{l} h_{k}^{m} \nabla_{l} n_{m}$$

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Metric energy-momentum tensor

] The EMT yielded by variational problem with respect to metric: $T_{ik} = T_{ik}^{(vol)} + T_{ik}^{(surf)}$

$$T_{ik}^{(vol)} = \frac{1}{2} (\nabla_i \varphi \nabla_k \varphi^* + \nabla_k \varphi \nabla_i \varphi^*) + \left[\left(\xi - \frac{1}{4} \right) g_{ik} g^{lm} \nabla_l \nabla_m - \xi R_{ik} - \xi \nabla_i \nabla_k \right] \varphi \varphi^*$$
$$T_{ik}^{(surf)} = \delta(x; \partial M_s) \left[\xi \varphi \varphi^* K_{ik} - \left(\xi - \frac{1}{4} \right) h_{ik} n^l \nabla_l (\varphi \varphi^*) \right]$$

Important notes

Equations of motion

 $g^{lm}\nabla_l\nabla_m\varphi + m^2\varphi + \xi R\varphi = 0$

 $2\left(\xi K + m_s\right)\varphi + n^l \nabla_l \varphi = 0$

Boundary condition of Robin type.

Surface EMT is orthogonal to the boundary

Integral conservation law for EMT

The model under discussion

- Randall-Sundrum braneworlds
 - □ There are natural timelike boundaries (branes) in AdS spacetime
 - The background in the original RS models is 5D AdS spacetime and the branes are located perpendicular to the additional (timelike) dimension
 - In the 2-brane model the hierarchy between electroweak and Planck scales is resolved at large inter-brane distances
- The model at hand
 - □ D+1 dimensional spacetime
 - \Box Local AdS geometry $ds^2 = e^{-2y/a} \eta_{\mu\nu} dx^\mu dx^\nu dy^2$

$$ds^{2} = \left(\frac{a}{z}\right)^{2} \left(\eta_{\mu\nu}dx^{\mu}dx^{\nu} - dz^{2}\right)$$

One or two D dimensional brane(s) located at y1 and y2

The model hand (continued)

 \square Globally the topology is different from AdS: $R^p imes T^q$

 R^1 ———

The proper length of compact dimensions

T¹

$$L_{(p)i} = (a/z)L_i = e^{-y/a}L_i$$

Presence of constant background gauge field

Quasi-periodicity condition

 $\varphi(t, x^1, \dots, x^l + L_l, \dots, y) = e^{i\alpha_l}\varphi(t, x^1, \dots, x^l, \dots, y), \quad l = p+1, \dots, D-1$

VEV of surface EMT

The vacuum expectation value (VEV) of the surface energy-momentum tensor is proportional to the VEV of field squared

$$\langle 0 | T_{\mu\nu}^{(surf)} | 0 \rangle = -\delta(x; \partial M_s) g_{\mu\nu} \left[\xi \frac{n^{(j)}}{a} - \frac{2\xi - \frac{1}{2}}{\beta^{(j)}} \right] \langle 0 | \varphi \varphi^{\dagger} | 0 \rangle$$

Cosmological constant type contribution on the brane

$$\langle 0 | \tau_i^{(j)k} | 0 \rangle_{y=y_j} = -\frac{c_j}{a} \delta_i^k \langle 0 | \varphi \varphi^{\dagger} | 0 \rangle_{y=y_j}$$

- In 1-brane model regularized quantities are obtained via zeta function regularization, with the help of extended Chowla-Selberg formula.
- In 2-brane model the zeta function is represented as sum of one brane induced and interference terms

$$\zeta_{j\boldsymbol{k}_{q}}(s) = \zeta_{j\boldsymbol{k}_{q}}^{(J)}(s) - \frac{B_{j}}{\Gamma\left(-\frac{s}{2}\right)\Gamma\left(\alpha_{s}+1\right)\mu^{1+s}} \int_{k_{q}}^{\infty} \mathrm{d}\lambda\lambda \left(\lambda^{2}-k_{q}^{2}\right)^{\alpha_{s}}\Omega_{j\nu}(\lambda z_{1},\lambda z_{2})$$

The interference term

The energy density corresponding to the interference term is finite

-(2)

Asymptotics for visible brane in the case of large inter-brane distance

 \Box Large KK modes: $k_q z_2 \gg 1$

$$\Delta \epsilon_{2,\boldsymbol{k}_{q}} \approx \frac{2\pi c_{2} B_{2}^{2} z_{2}^{D-\frac{p-1}{2}}}{a^{D} (4\pi)^{\frac{p+1}{2}} V_{q}} \frac{k_{q}^{\frac{p+3}{2}}}{\left[A_{2}+B_{2} k_{q} z_{2}\right]^{2}} \frac{\bar{I}_{\nu}^{(1)}(k_{q} z_{1})}{\bar{K}_{\nu}^{(1)}(k_{q} z_{1})} e^{-2k_{q} z_{2}}$$

 \square Small KK modes: $k_q z_1 \ll 1$

$$\Delta \epsilon_{2,\boldsymbol{k_q}} = \frac{z_2^q}{a^D d_1(\nu) V_q} \left(\frac{z_1}{z_2}\right)^{2\nu} f_{2\nu}^{(k_q)}$$
$$f_{2\nu}^{(k_q)} = \frac{2^{2-2\nu} c_2 B_2^2}{(4\pi)^{\frac{p+1}{2}} \Gamma\left(\frac{p+1}{2}\right) \nu \Gamma^2(\nu)} \int_{k_q z_2}^{\infty} \mathrm{d}u \frac{u^{1+2\nu} \left(u^2 - k_q^2 z_2^2\right)^{\frac{p-1}{2}}}{\bar{I}_{\nu}^{(2)2}(u)}$$

Cosmological constant

The cosmological constant due to the interference term on the visible brane, in the case of large inter-brane distances in Planck units

$$h_2 \approx (D-2)^{-\bar{\nu}} (aM_{D+1})^{\bar{\nu}-(p+1)} (V'_{2,q} M^q_{D+1})^{\bar{\nu}} \left(\frac{M_{D+1}}{M_{p+1;2}}\right)^{p+1+\bar{\nu}(p-1)} \frac{1}{d_1(\nu)} \sum_{\boldsymbol{k}_q} f_{2\nu}^{(k_q)}$$

Observed cosmological constant (Plank units): ~ 10^{-123}

■ When the hierarchy problem between electroweak and Planck scales is resolved □ The fundamental scales in the problem are TeV: $1/a \sim M_{D+1} \sim 1$ TeV □ Planck scale for the visible brane is 10^{16} TeV: $M_{p+1;2} = M_{Pl} \sim 10^{16}$ TeV □ Proper volume of compact subspace is of Planck order: $V'_{2,q}M^q_{D+1} \sim 1$ □ The cosmological constant due to the interference term is as follows $h_2 \approx 10^{-32(2+\bar{\nu})} \frac{(D-2)^{-\bar{\nu}}}{V_{2,q}} \sum f^{(k_q)}_{2m}$

$$u_2 \approx 10^{-32(2+\bar{\nu})} \frac{(D-2)}{d_1(\nu)} \sum_{k_q} f_{2\nu}^{(k_q)}$$

Cosmological constant

□ E.g. for a massless, minimally coupled scalar field, and a single compact dimension the power of ten alone in the coefficient is 10⁻¹¹⁷

The rest of the expression for cosmological constant



□ The sign of cosmological constant is not dependent of the gauge flux

Conclusion

- Surface terms for scalar field with coupling to the background curvature
 - Achieve correctly posed variational (with respect to metric tensor) problem in the presence of scalar field
 - □ Naturally arising Robin boundary conditions
 - Integral conservation of total energy density
- The cosmological constant problem
 - □ Natural mechanism for obtaining small cosmological constant from VEV
 - The cosmological constant problem is solved with the same configuration when the hierarchy problem is solved
 - Sign of cosmological constant on visible brane is determined by the boundary conditions on both branes

Thank you!