## 1. Relativistic Boltzmann equation

Show

$$\left(p^{\mu}\partial_{\mu}+mF^{\mu}\frac{\partial}{\partial p^{\mu}}\right)\theta(p^{0})\delta(p^{2}+m^{2})=0$$

provided that either  $F^\mu p_\mu = 0$  or  $F^\mu = -\partial^\mu m = 0$ 

## 2. Entropy production

Starting from the fluid-fleld coupling model

$$\partial_{\mu}T_{\mathrm{f}}^{\mu
u}+\partial^{
u}\phirac{\partial V_{T}(\phi)}{\partial\phi}= ilde{\eta}rac{\phi^{2}}{T}(U\cdot\partial\phi)\partial^{
u}\phi$$

show that the entropy current ( $S^{\mu} = sU^{\mu}$ , s = dp/dT) satisfies

$$\partial_{\mu} S^{\mu} = \tilde{\eta} (\beta \phi)^2 (U \cdot \partial \phi)^2$$

Don't forget that  $T^{\mu\nu}=wU^{\mu}U^{\nu}+pg^{\mu\nu}$ , w=Ts, and  $p=g_{\rm eff}\frac{\pi^2}{90}T^4-V_T(\phi)$ .

## 3. Junction conditions

Starting from EM conservation across the bubble wall:

$$W_{-}\gamma_{-}^{2}V_{-} = W_{+}\gamma_{+}^{2}V_{+}, \quad W_{-}\gamma_{-}^{2}V_{-}^{2} + p_{-} = W_{+}\gamma_{+}^{2}V_{+}^{2} + p_{+}$$

show

$$v_+v_-=rac{
ho_+-
ho_-}{e_+-e_-}, \quad rac{v_+}{v_-}=rac{e_-+
ho_+}{e_++
ho_-}$$

and hence

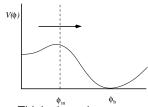
$$v_+v_-=rac{1-(1-3lpha_+)r}{3-3(1+lpha_+)r}, rac{v_+}{v_-}=rac{3+(1-3lpha_+)r}{1+3(1+lpha_+)r}$$

where

• 
$$\epsilon_{\pm} = \frac{1}{4} (e_{\pm} - 3p_{\pm}), \epsilon = \epsilon_{+} - \epsilon_{-}$$

$$r = w_+/w_-$$

## 4. Quantum tunnelling vs. thermal activation



- Position φ
- momentum π
- Hamiltonian  $H = \frac{1}{2}\pi^2 + V(\phi)$

Think about the quantum version of the thermal activation problem.

- Quantum-Statistical metastability, I. Affleck, PRL 46, 388 (1981)
- ► The uses of instantons, S. Coleman, in Aspects of Symmetry

$$\Gamma = \frac{2}{\hbar} \text{Im} E_0$$

where the imaginary part of the ground state energy is roughly

$${
m Im} E_0 \sim rac{\hbar \omega}{2} e^{-S_0/\hbar}$$

and  $S_0$  is the extremal value of the functional

$$S = \int d au \left( rac{1}{2} \left( rac{d\phi}{d au} 
ight)^2 + V(\phi) 
ight)$$

