

Ruben Verresen^(1,2)

With Matthias Gohlke⁽¹⁾, Roderich Moessner⁽¹⁾ and Frank Pollmann⁽²⁾

⁽¹⁾ Max-Planck-Institute for the Physics of Complex Systems (Dresden)

⁽²⁾ Technical University of Munich

@ Benasque 2017

Excited about Excitations

Quantum phases of matter \rightarrow interesting excitations

E.g. fractionalization

= quasi-particle cannot be created locally

Excited about Excitations

Quantum phases of matter \rightarrow interesting excitations

- E.g. fractionalization
 - = quasi-particle cannot be created locally



Spectral Function

Question:

"How many eigenstates with energy ω and momentum k are in $\mathcal{O} |0\rangle$?"



Answer:

$$\begin{split} \boldsymbol{\mathcal{S}}(\boldsymbol{k}, \boldsymbol{\omega}) &= \sum_{n} \delta(\boldsymbol{\omega} - \boldsymbol{\omega}_{n}) \; \delta(\boldsymbol{k} - \boldsymbol{k}_{n}) \; |\langle n | \mathcal{O} | 0 \rangle|^{2} \\ &= \; \text{Fourier transform of } \; \langle \mathcal{O}(\boldsymbol{r}, t) \; \mathcal{O}(\boldsymbol{0}, 0) \rangle \end{split}$$

Spectral Function

Question:

"How many eigenstates with energy ω and momentum k are in $\mathcal{O} |0\rangle$?"



Answer:

$$S(\mathbf{k}, \omega) = \sum_{n} \delta(\omega - \omega_{n}) \ \delta(\mathbf{k} - \mathbf{k}_{n}) \ |\langle n|\mathcal{O}|0\rangle|^{2}$$

= Fourier transform of $\langle \mathcal{O}(\mathbf{r}, t) \ \mathcal{O}(\mathbf{0}, 0) \rangle$
How to obtain it

Spectral Function: Examples $(\mathcal{O} = \sigma^y)$

$$H = \sum \left(-J\sigma_n^z \sigma_{n+1}^z + g\sigma_n^x \right)$$



Spectral Function: Examples $(\mathcal{O} = \sigma^y)$

$$H = \sum \left(-J\sigma_n^z \sigma_{n+1}^z + g\sigma_n^x \right)$$



Spectral Function: Summary

$$\mathcal{S}(\boldsymbol{k},\omega) = \frac{1}{2\pi} \sum_{\mathbf{r}} \int_{-\infty}^{\infty} e^{i(\omega t - \boldsymbol{k} \cdot \boldsymbol{r})} \langle \boldsymbol{S}_{\boldsymbol{r}}(t) \cdot \boldsymbol{S}_{\boldsymbol{0}}(0) \rangle \, \mathrm{d}t$$

Information about

 $\rightarrow \text{spectrum of the Hamiltonian (Lehmann repres.)}$ $\mathcal{S}(\boldsymbol{k},\omega) = \sum_{n} \delta(\omega - \omega_{n}) \ \delta(\boldsymbol{k} - \boldsymbol{k}_{n}) \ |\langle n|\boldsymbol{S}_{0}|0\rangle|^{2} \quad \checkmark$

 \rightarrow are quasi-particles local or not (fractionalization)

 \rightarrow quasi-particle lifetime, etc...

Spectral Function: Experiment

$$\mathcal{S}(\boldsymbol{k},\omega) = \frac{1}{2\pi} \sum_{\mathbf{r}} \int_{-\infty}^{\infty} e^{i(\omega t - \boldsymbol{k} \cdot \boldsymbol{r})} \langle \boldsymbol{S}_{\boldsymbol{r}}(t) \cdot \boldsymbol{S}_{\boldsymbol{0}}(0) \rangle \, \mathrm{d}t$$

Measurable through inelastic neutron scattering !



Spectral Function: Experiment

$$\mathcal{S}(\boldsymbol{k},\omega) = \frac{1}{2\pi} \sum_{\mathbf{r}} \int_{-\infty}^{\infty} e^{i(\omega t - \boldsymbol{k} \cdot \boldsymbol{r})} \langle \boldsymbol{S}_{\boldsymbol{r}}(t) \cdot \boldsymbol{S}_{\boldsymbol{0}}(0) \rangle \, \mathrm{d}t$$

Measurable through inelastic neutron scattering !



Banerjee et al. (2016)

A Very Short Outline

 Numerical method for calculating dynamic correlations for two-dimensional systems



Application to the Kitaev-Heisenberg model





2D in a 1D jacket

Real space:



Momentum space:



(cylinder by Adolfo Grushin)

Cylinder: 1D or 2D?



$$H = \sum_{\langle \boldsymbol{i}, \boldsymbol{j} \rangle} S_{\boldsymbol{i}}^{x} S_{\boldsymbol{j}}^{x} + S_{\boldsymbol{i}}^{y} S_{\boldsymbol{j}}^{y} + \Delta S_{\boldsymbol{i}}^{z} S_{\boldsymbol{j}}^{z}$$



Cylinder: 1D or 2D?



$$H = \sum_{\langle \boldsymbol{i}, \boldsymbol{j} \rangle} S_{\boldsymbol{i}}^{x} S_{\boldsymbol{j}}^{x} + S_{\boldsymbol{i}}^{y} S_{\boldsymbol{j}}^{y} + \Delta S_{\boldsymbol{i}}^{z} S_{\boldsymbol{j}}^{z}$$



Cylinder: 1D or 2D?



Compared to exact diagonalization:

Similar care needed when extrapolating

Advantages:

- \rightarrow larger system sizes
- \rightarrow continuous momentum
- \rightarrow (much) longer times

Method to Obtain $\mathcal{S}({\pmb k},\omega)$ in 2D

Physics: $\Phi_y \qquad (L \le 12)$

Numerical ingredients:

 \rightarrow For ground state:

Density Matrix Renormalization Group (DMRG) (*)

 \rightarrow For time evolution:

Matrix Product Operator (MPO) method (**)

We obtain: $C(\mathbf{r},t) \xrightarrow{\text{Fourier}} \mathcal{S}(\mathbf{k},w)$



(*) White, PRL 69, 2863 (1992) (**) Zaletel, Mong, Karrasch, Moore, Pollmann; PRB 91, 165112 (2015)

A Very Short Outline

 Numerical method for calculating dynamic correlations for two-dimensional systems



Application to the Kitaev-Heisenberg model





Kitaev (Honeycomb) Model
$$H = \sum_{\langle \mathbf{i}, \mathbf{j} \rangle_x} S^x_{\mathbf{i}} S^x_{\mathbf{j}} + \sum_{\langle \mathbf{i}, \mathbf{j} \rangle_y} S^y_{\mathbf{i}} S^y_{\mathbf{j}} + \sum_{\langle \mathbf{i}, \mathbf{j} \rangle_z} S^z_{\mathbf{i}} S^z_{\mathbf{j}}$$

7 N1

Exactly soluble model (2006):

$$S^{\gamma} = c_a^{\dagger} \sigma_{ab}^{\gamma} c_b \quad \dashrightarrow \quad c_a = \gamma_a + i \tilde{\gamma}_a \quad \stackrel{\text{nolf}}{\longrightarrow} \quad \underset{\text{ball}}{\text{majoranas}}$$

- \rightarrow spin liquid ground state
- \rightarrow fractionalization:

one spin = two gapped fluxes and gapless majorana \rightarrow gapped $\mathcal{S}({\pmb k},w)$ (Knolle et al., `13) Dynamics of the gapped Kitaev Model

$$H = -\sum_{\langle \mathbf{i}, \mathbf{j} \rangle_{\gamma}} J_{\gamma} S_{\mathbf{i}}^{\gamma} S_{\mathbf{j}}^{\gamma} \qquad (J_x = 6J_y = 6J_z)$$

Benchmark: gapped Kitaev model $(L_{\rm circ} = 10)$



N₂

7 N1

Benchmark: gapless Kitaev model ($L_{\rm circ} = 6$)



Kitaev-Heisenberg Model

$$H = 2 \sin \alpha \sum_{\langle i, j \rangle_{\gamma}} S_{i}^{\gamma} S_{j}^{\gamma} + \cos \alpha \sum_{i, j} S_{i} \cdot S_{j}$$

→ relevance for $A_2 IrO_3$ and α -RuCl₃ (Chaloupka, Jackeli, Khaliullin, `10)

\rightarrow various works on ground state phase diagram

(Chaloupka et al., Iregui et al., Jiang et al., Sela et al., Shinjo et al.)

 $\rightarrow S(\mathbf{k}, w)$? Perturbative results, low energies (Song, You, Balents, `16; also recent E.D.: Gotfryd et al, `16)

→ experiments: "proximate spin liquids" (Banerjee et al., `16; Suzuki et al., `15)



Kitaev-Heisenberg Model

$$H = 2 \sin \alpha \sum_{\langle i, j \rangle_{\gamma}} S_{i}^{\gamma} S_{j}^{\gamma} + \cos \alpha \sum_{i, j} S_{i} \cdot S_{j}$$



Ferromagnet with $J_{\text{Kitaev}} = 0.65 J_{\text{Heis}}$ $(L_{\text{circ}} = 12)$







 k_1^{\perp}/π

Zigzag phase: proximate spin liquid? $(L_{circ} = 12)$



 k_1^\perp/π

 k_1^\perp/π





Zigzag phase: proximate spin liquid? $(L_{circ} = 12)$







Zigzag phase: proximate spin liquid? $(L_{circ} = 12)$







Zigzag phase: proximate spin liquid? $(L_{circ} = 12)$





 $\alpha - \operatorname{RuCl}_3$



(Banerjee et al., `16)

Summary

Obtaining $S(\mathbf{k},\omega)$ for two-dimensional systems

Application to Kitaev-Heisenberg model \rightarrow anomalous features of α -RuCl₃

Exciting questions waiting to be attacked...

How can a spin wave fractionalize? Spin wave anomalies on other lattices? Signatures of non-abelian anyons? How does 1D become 2D?

...



Ruben Verresen

With Matthias Gohlke, Roderich Moessner and Frank Pollmann

Gapless Cylinder



Coleman and 1D vs. 2D



Broadness in Zig-Zag Phase

