

# Majorana edge states and level crossings in the transverse field Ising model

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# Scope

- The 1D Kitaev model of p-wave superconductor
  - Majorana wave-function oscillations
- Chains of cobalt adatoms
  - XY model with in-plane magnetic field
  - Series of level crossings
- Transverse field Ising model with longitudinal coupling
  - Perturbative approach
  - Mapping on Kitaev
- Conclusions

# Kitaev p-wave SC chain

Kitaev, 2001

$$\mathcal{H} = -t \sum_{i=1}^{N-1} (c_i^\dagger c_{i+1} + \text{H.c.}) + \Delta \sum_{i=1}^{N-1} (c_{i+1}^\dagger c_i^\dagger + \text{H.c.}) - \mu \sum_{i=1}^N c_i^\dagger c_i$$

- $|\mu| < 2t$ 
  - topological phase
  - 2-fold degenerate ground state
- $|\mu| > 2t$ 
  - trivial phase
  - unique ground state

# Majorana formulation

## Majorana fermions

$$\begin{cases} \gamma'_i = c_i + c_i^\dagger \\ \gamma''_i = -i(c_i - c_i^\dagger) \end{cases}$$

$$\gamma''_i{}^\dagger = \gamma''_i$$

$$\gamma'_i{}^\dagger = \gamma'_i$$

$$\{\gamma'_i, \gamma''_j\} = 0$$

$$\{\gamma'_i, \gamma'_j\} = \{\gamma''_i, \gamma''_j\} = 2\delta_{ij}$$

$$\mathcal{H} = \frac{i}{2} \sum_{i=1}^{N-1} (\Delta + t) \gamma''_i \gamma'_{i+1} + (\Delta - t) \gamma'_i \gamma''_{i+1} - \frac{i\mu}{2} \sum_{i=1}^N \gamma'_i \gamma''_i$$

# Majorana edge states

$$\Delta = t, \mu = 0$$



$$\mathcal{H} = it \sum_{i=1}^{N-1} \gamma_i'' \gamma_{i+1}'$$

$\gamma_1'$  and  $\gamma_N''$  do not appear in  $\mathcal{H}$



Perfectly localized states

$$\begin{cases} c^\dagger = \gamma_1' + i\gamma_N'' \\ c = \gamma_1' - i\gamma_N'' \end{cases}$$

→ fermion, can be empty or occupied

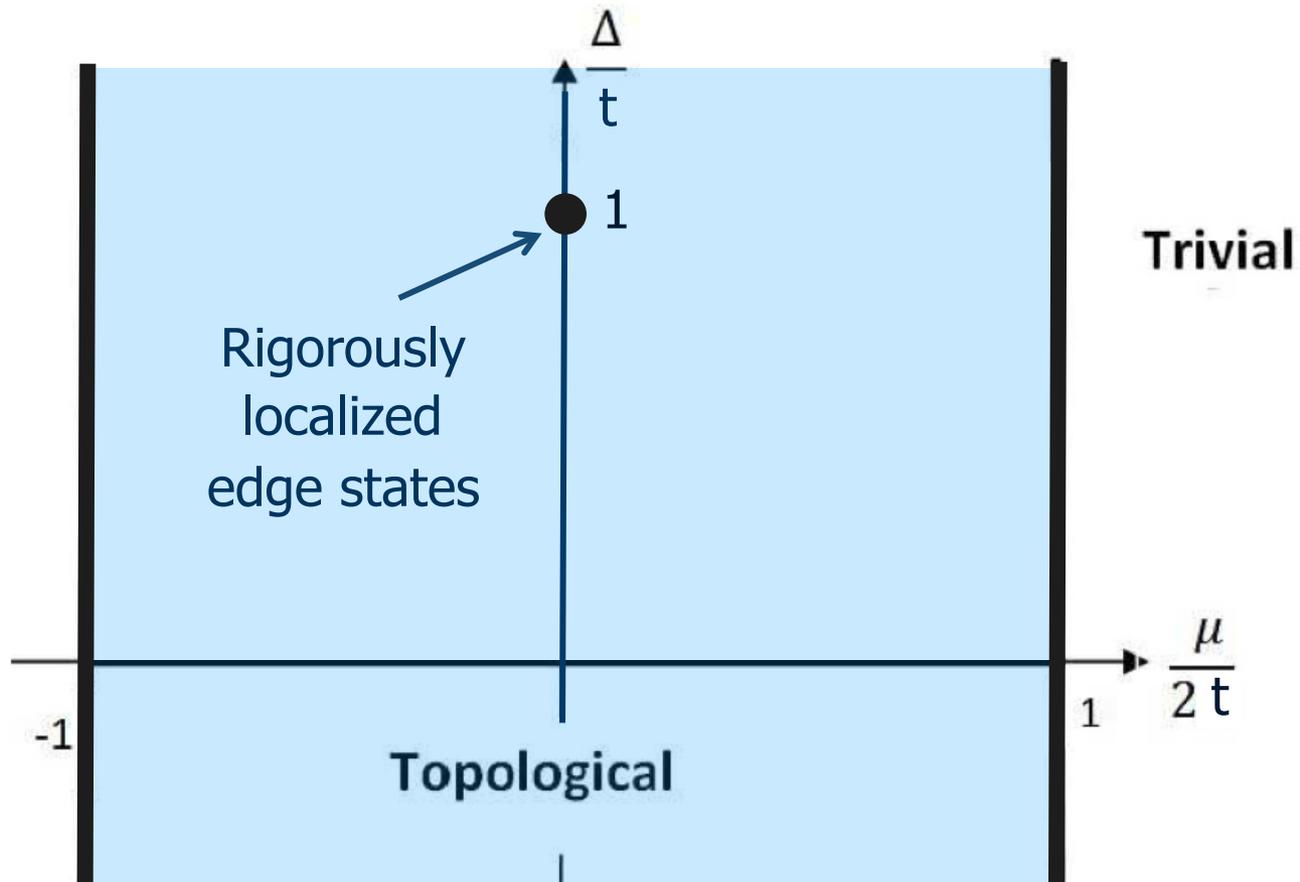


Two-fold GS degeneracy

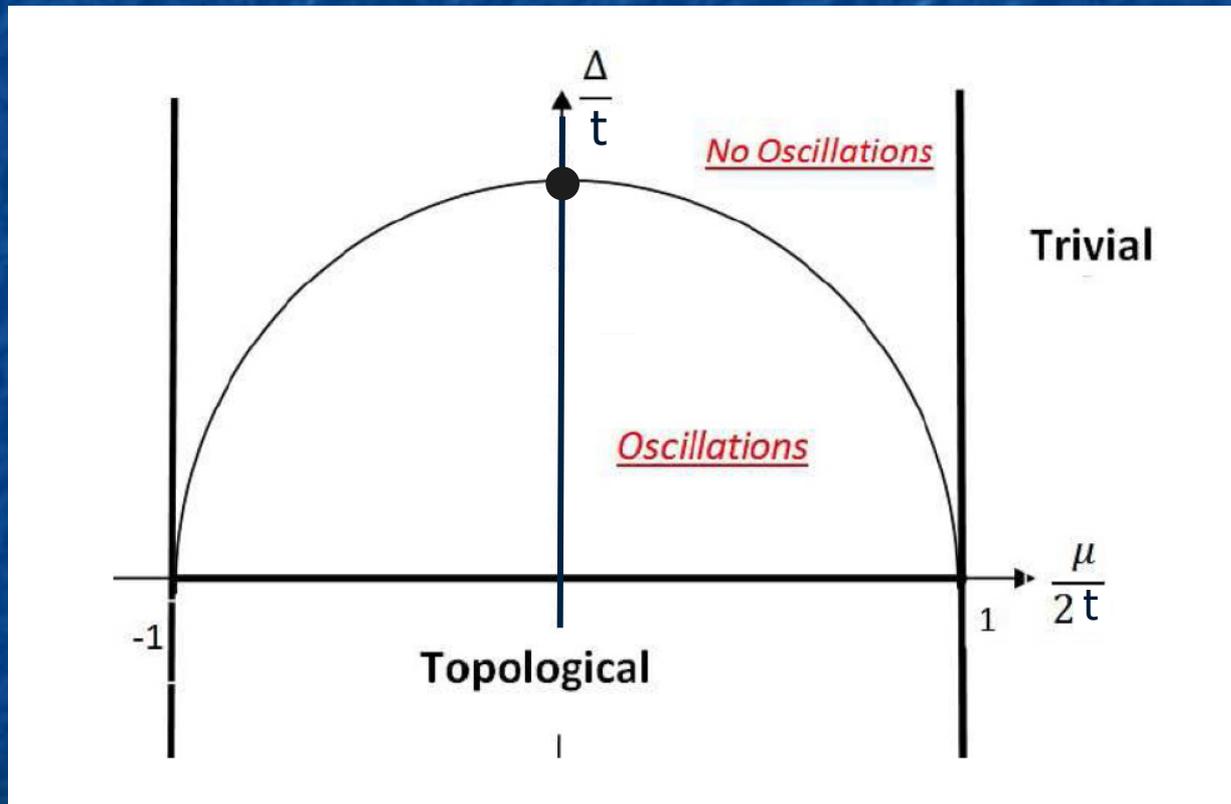
# General properties of topological phase

- Localized Majorana edge states
- Quasi-degeneracy
  - **splitting** of the order of  $\exp(-Na/\xi)$
  - $\xi$  = localization length of Majorana edge states
- Detection of Majorana edge states
  - STM at the end of the chain
    - Mourik et al, Science 2012; Nadj-Perge et al, Science 2014
- Other way of detecting them? **YES!**

# Phase diagram



# Majorana oscillations



Kao, PRB 2014; Hedge and Vishveshwara, PRB 2016

# Majorana wave-functions

Diagonalization of  $H \longrightarrow$

$$\mathcal{H} = \frac{i}{2} \sum_{k=1}^N \epsilon_k \tilde{\gamma}'_k \tilde{\gamma}''_k$$

$$\tilde{\gamma}'_k = \sum_j (a_+ x_+^j + b_+ x_+^{N+1-j} + a_- x_-^j + b_- x_-^{N+1-j}) \gamma'_j$$

$$\tilde{\gamma}''_k = \sum_j (a_+ x_+^{N+1-j} + b_+ x_+^j + a_- x_-^{N+1-j} + b_- x_-^j) \gamma''_j$$

Smallest  
eigenvalue

$$\left(\frac{\mu}{2t}\right)^2 + \left(\frac{\Delta}{t}\right)^2 < 1 \rightarrow x_{\pm} = r e^{\pm i\phi}, \phi > 0$$

# Boundary conditions for zero mode

Smallest eigenvalue

$$\epsilon_1 = 0 \quad \longrightarrow \quad \begin{cases} a_+ + a_- = 0 \\ a_+ x_+^{N+1} + a_- x_-^{N+1} = 0 \end{cases}$$

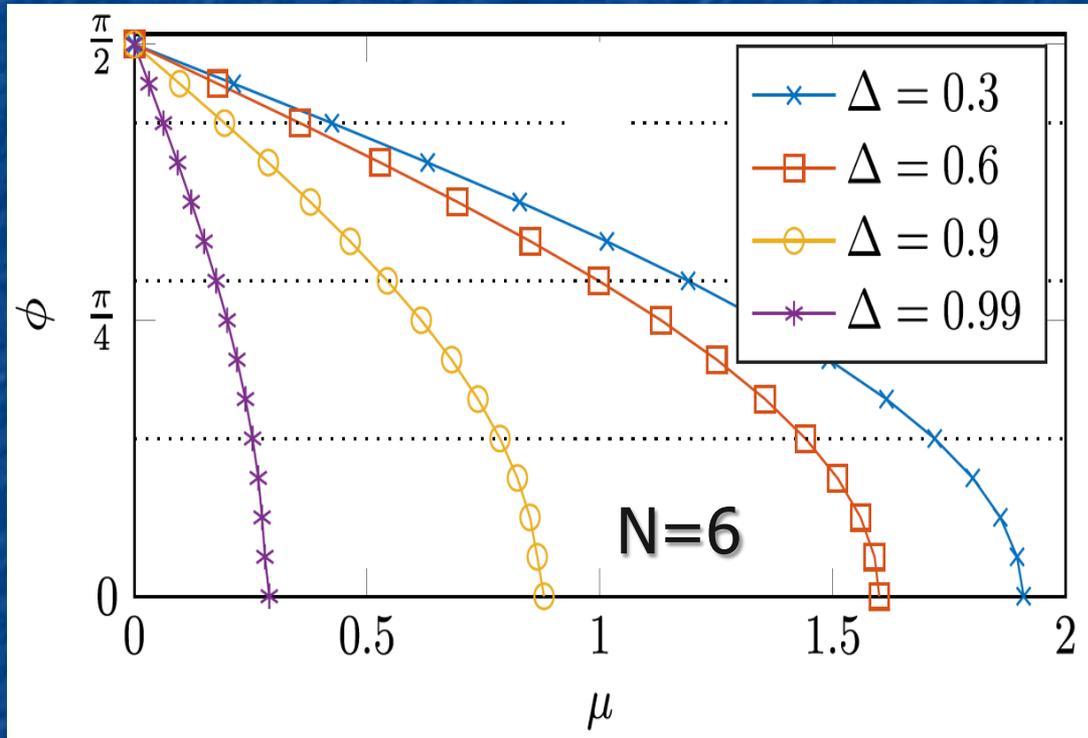
Possible if

$$(x_+^{N+1} - x_-^{N+1}) \propto r^{N+1} \sin[(N+1)\phi] = 0$$

or

$$\phi = \frac{\pi m}{N+1}$$

# Phase condition

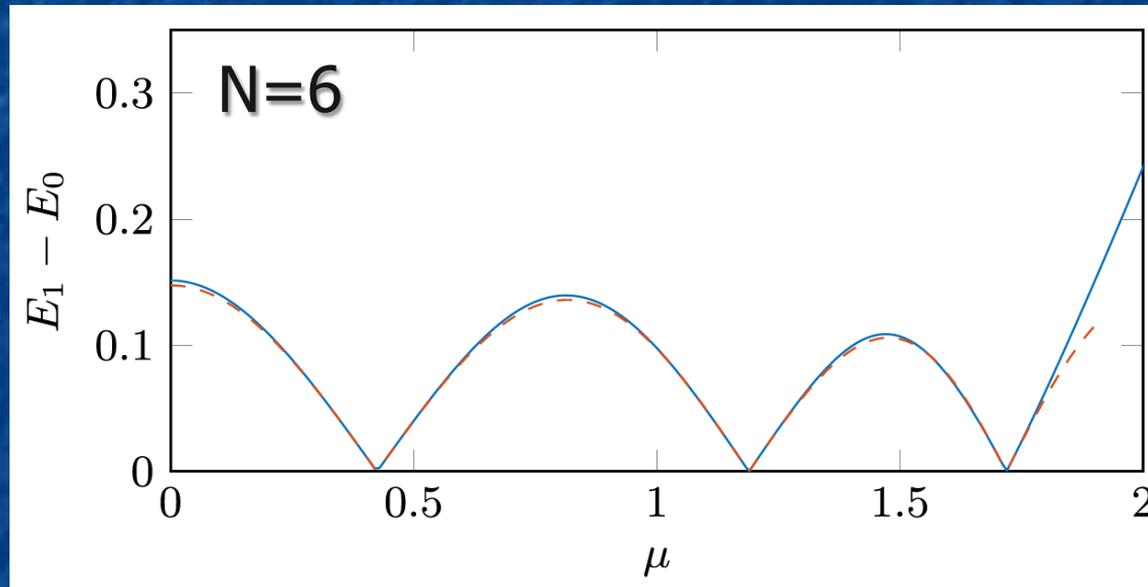


N solutions  
as soon as

$$\left(\frac{\Delta}{t}\right)^2 < 1$$

G. Vionnet, B. Kumar, FM, arXiv:1701.08057

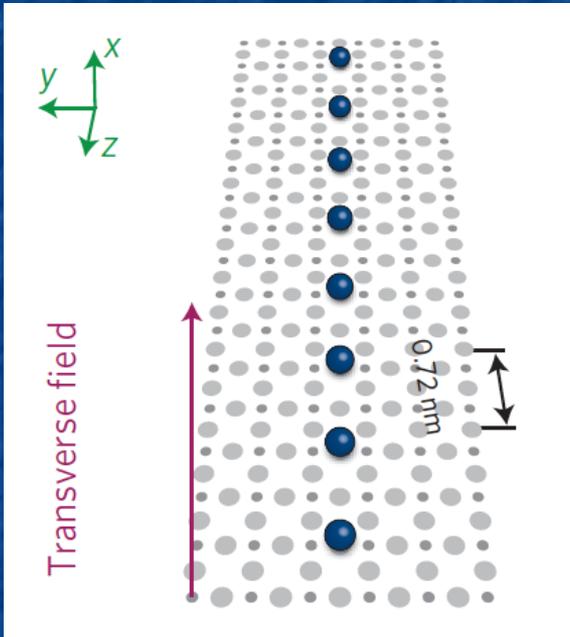
# Zero modes



N zero modes in the interval  $-2t < \mu < 2t$

NB: **level crossings** because low-lying states have different **fermion parity**

# Chain of Co adatoms on Cu<sub>2</sub>N



$$\mathcal{H}_{3/2} = J \sum_{i=1}^{N-1} \mathbf{S}_i \cdot \mathbf{S}_{i+1} + D \sum_{i=1}^N (S_i^z)^2 - g\mu_B B_x \sum_{i=1}^N S_i^x$$

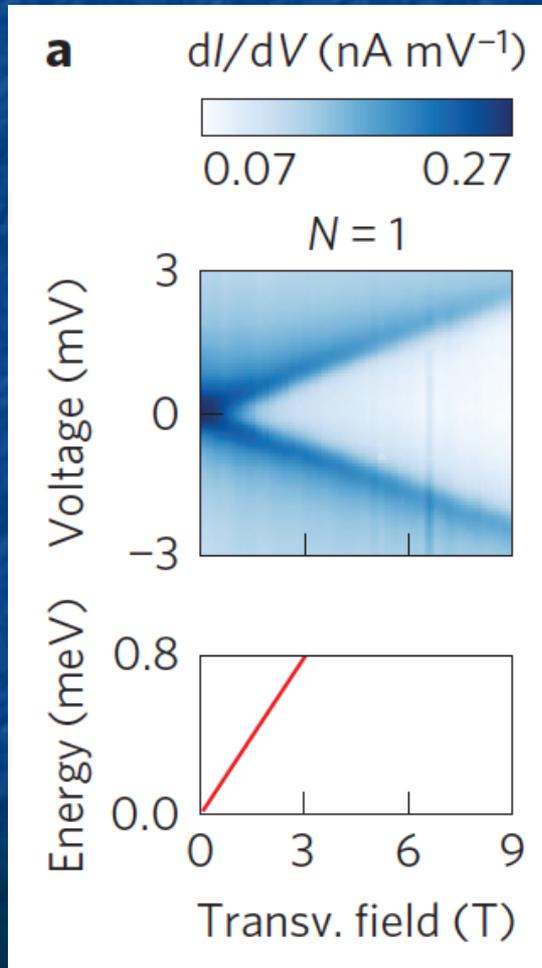
$$J = 0.24 \text{ meV}, \quad D = 2.75 \text{ meV}$$

J/D expansion

$$\mathcal{H}_{1/2} \simeq \sum_{i=1}^{N-1} J_{\perp} (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y) - \mu_B B_x \sum_{i=1}^N g_i S_i^x$$

Teskovic et al, Nat. Phys. 2016

# STM results: single site



Zeeman effect



Splitting of Kondo resonance



V shape spectrum

# STM results: up to 9 sites

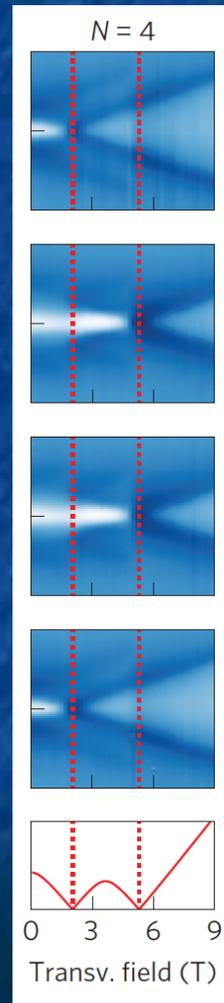
4 sites

Site 1

Site 2

Site 3

Site 4

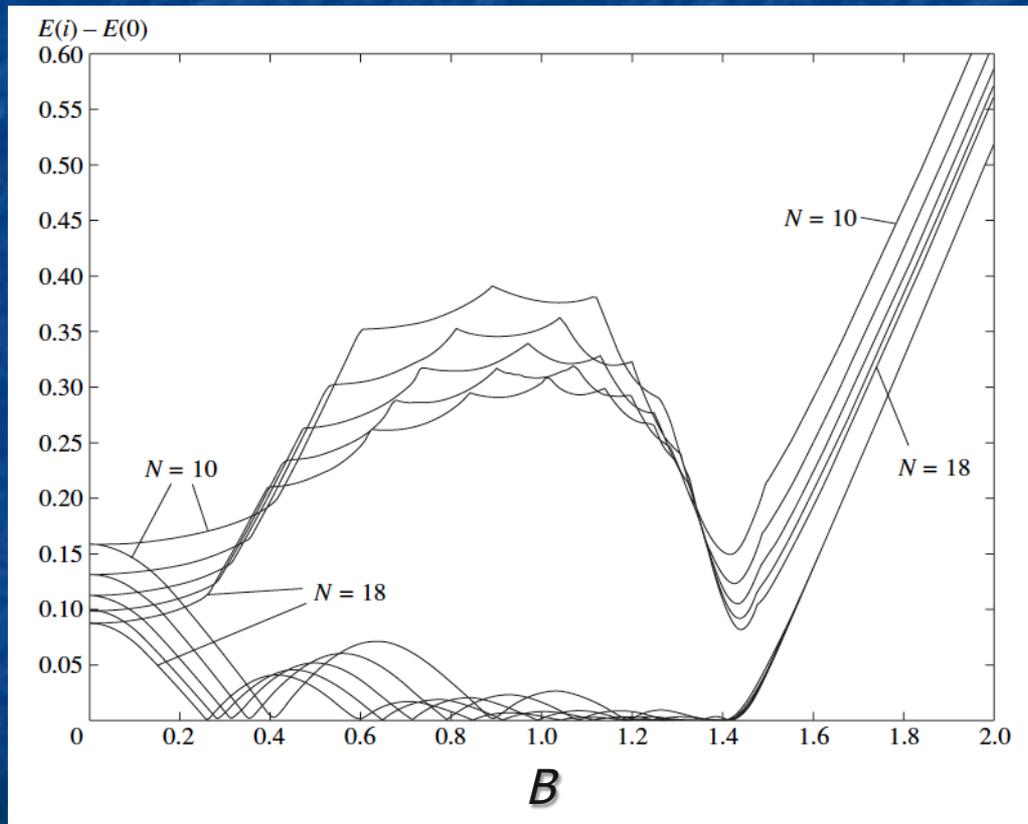


V shape



Level crossing at  
the tip of the V

# Exact diagonalization



$N/2$  level crossings  
for  $B > 0$



$N$  level crossings

Dmitriev, Krivnov, Ovchinnikov, Langari, JETP 2002

# Link with Majorana edge states

FM, Nat. Phys. 2016

Transverse field  
Ising model



Jordan-Wigner

Kitaev chain  
with  $\Delta=t$

Longitudinal  
coupling



Kitaev chain  
with  $\Delta<t$

# Mapping of TFI onto Kitaev

Transverse field  
Ising model

$$H = J_x \sum_{i=1}^{N-1} S_i^x S_{i+1}^x - h \sum_{i=1}^N S_i^z$$

Jordan-Wigner transformation

$$\begin{cases} S_i^x = \frac{1}{2}(c_i^\dagger + c_i) \exp\left(i\pi \sum_{j<i} c_j^\dagger c_j\right) \\ S_i^y = \frac{1}{2i}(c_i^\dagger - c_i) \exp\left(i\pi \sum_{j<i} c_j^\dagger c_j\right) \\ S_i^z = c_i^\dagger c_i - \frac{1}{2}, \end{cases}$$


$$\mathcal{H} = t \sum_{i=1}^{N-1} (c_i^\dagger c_{i+1} + \text{H.c.}) - \Delta \sum_{i=1}^{N-1} (c_{i+1}^\dagger c_i^\dagger + \text{H.c.}) - \mu \sum_{i=1}^N c_i^\dagger c_i$$

with  $\Delta = t = J_x/4$ ,  $\mu = h$

Pfeuty, 1970

How to get  $\Delta < t$ ?

# XY model in transverse field

$$H = J_x \sum_{i=1}^{N-1} S_i^x S_{i+1}^x + J_y \sum_{i=1}^{N-1} S_i^y S_{i+1}^y - h \sum_{i=1}^N S_i^z$$



$$\mathcal{H} = t \sum_{i=1}^{N-1} (c_i^\dagger c_{i+1} + \text{H.c.}) - \Delta \sum_{i=1}^{N-1} (c_{i+1}^\dagger c_i^\dagger + \text{H.c.}) - \mu \sum_{i=1}^N c_i^\dagger c_i$$

$$\text{with } t = (J_x + J_y)/4, \Delta = (J_x - J_y)/4, \mu = h$$

Oscillations of correlation functions

Barouch and McCoy, 1971

# TFI with longitudinal coupling

$$H = J_x \sum_{i=1}^{N-1} S_i^x S_{i+1}^x + J_z \sum_{i=1}^{N-1} S_i^z S_{i+1}^z - h \sum_{i=1}^N S_i^z$$


$$\mathcal{H} = \mathcal{H}(J_z = 0) + J_z \sum_{i=1}^{N-1} (c_i^\dagger c_i - \frac{1}{2})(c_{i+1}^\dagger c_{i+1} - \frac{1}{2})$$

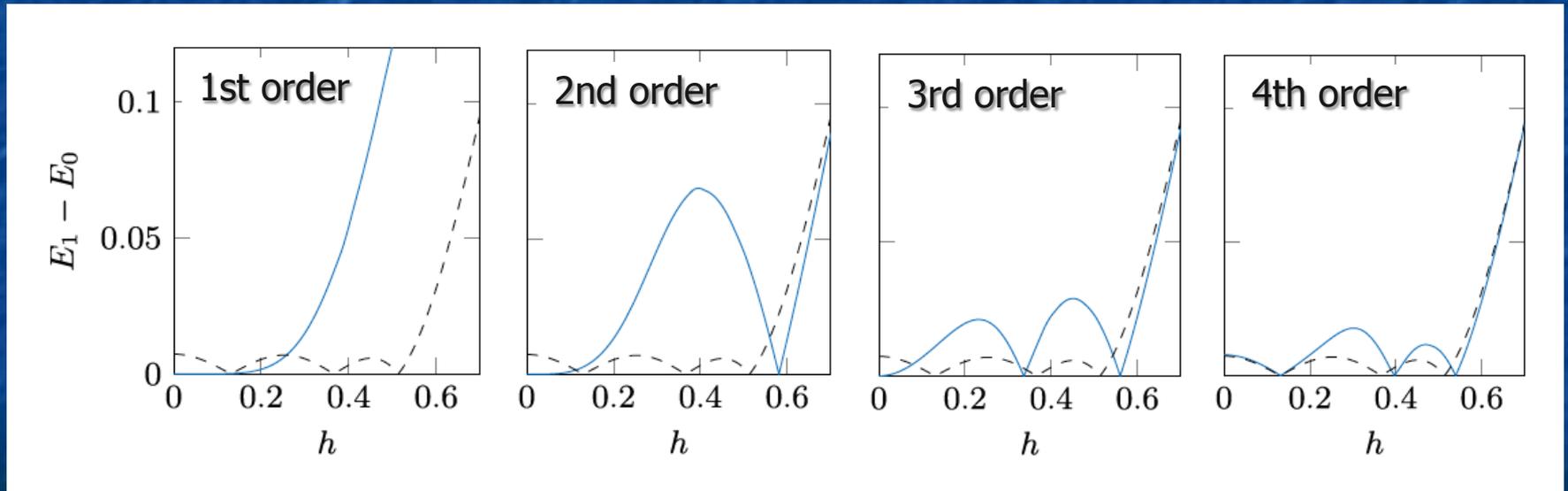
**Model of interacting fermions**

Symmetry: rotation by  $\pi$  around z

$\Leftrightarrow$  fermion parity

# Perturbation theory in $J_z$

- Different parity sectors  
→ non-degenerate perturbation theory



$$J_z = 0.25, N = 6$$

# Perturbation theory in $J_z$

- N level crossings seem to appear immediately when  $J_z$  is switched on
- The number of level crossings is limited by the order of perturbation theory  
→ hopeless for large values of N

G. Vionnet, B. Kumar, FM, arXiv:1701.08057

# Mean-field theory

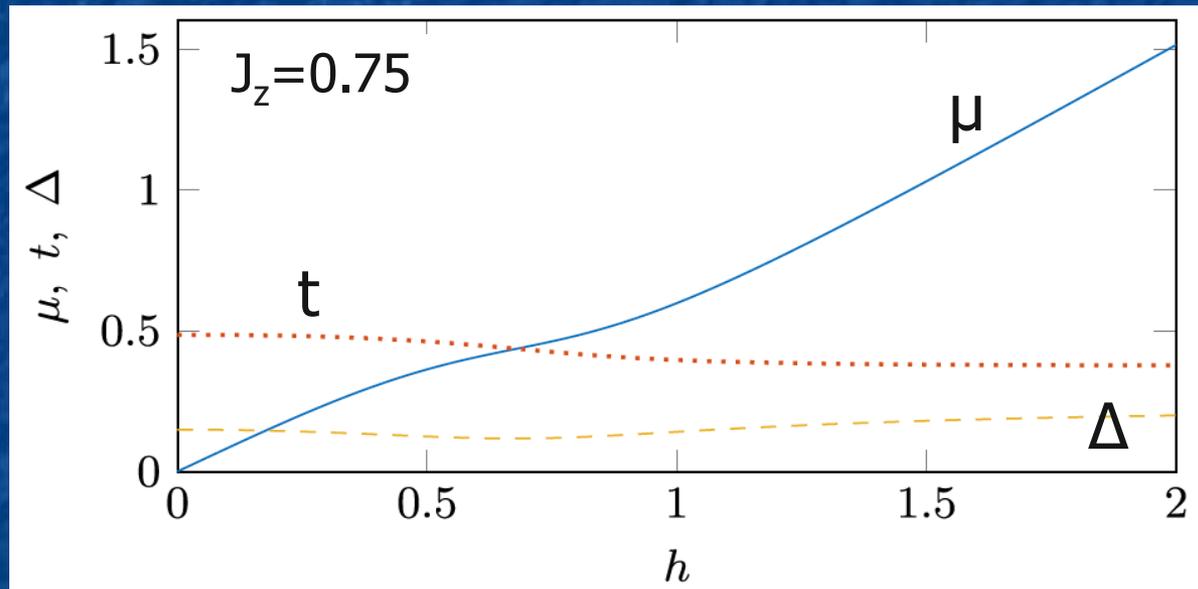
- Simultaneous decoupling in all channels

$$\begin{aligned}c_i^\dagger c_i c_{i+1}^\dagger c_{i+1} &\simeq \langle c_i^\dagger c_i \rangle c_{i+1}^\dagger c_{i+1} + \langle c_{i+1}^\dagger c_{i+1} \rangle c_i^\dagger c_i - \langle c_i^\dagger c_i \rangle \langle c_{i+1}^\dagger c_{i+1} \rangle \\ &\quad - \langle c_i^\dagger c_{i+1}^\dagger \rangle c_i c_{i+1} - \langle c_i c_{i+1} \rangle c_i^\dagger c_{i+1}^\dagger + \langle c_i^\dagger c_{i+1}^\dagger \rangle \langle c_i c_{i+1} \rangle \\ &\quad + \langle c_i^\dagger c_{i+1} \rangle c_i c_{i+1}^\dagger + \langle c_i c_{i+1}^\dagger \rangle c_i^\dagger c_{i+1} - \langle c_i^\dagger c_{i+1} \rangle \langle c_i c_{i+1}^\dagger \rangle\end{aligned}$$

→ Kitaev model for periodic BC with

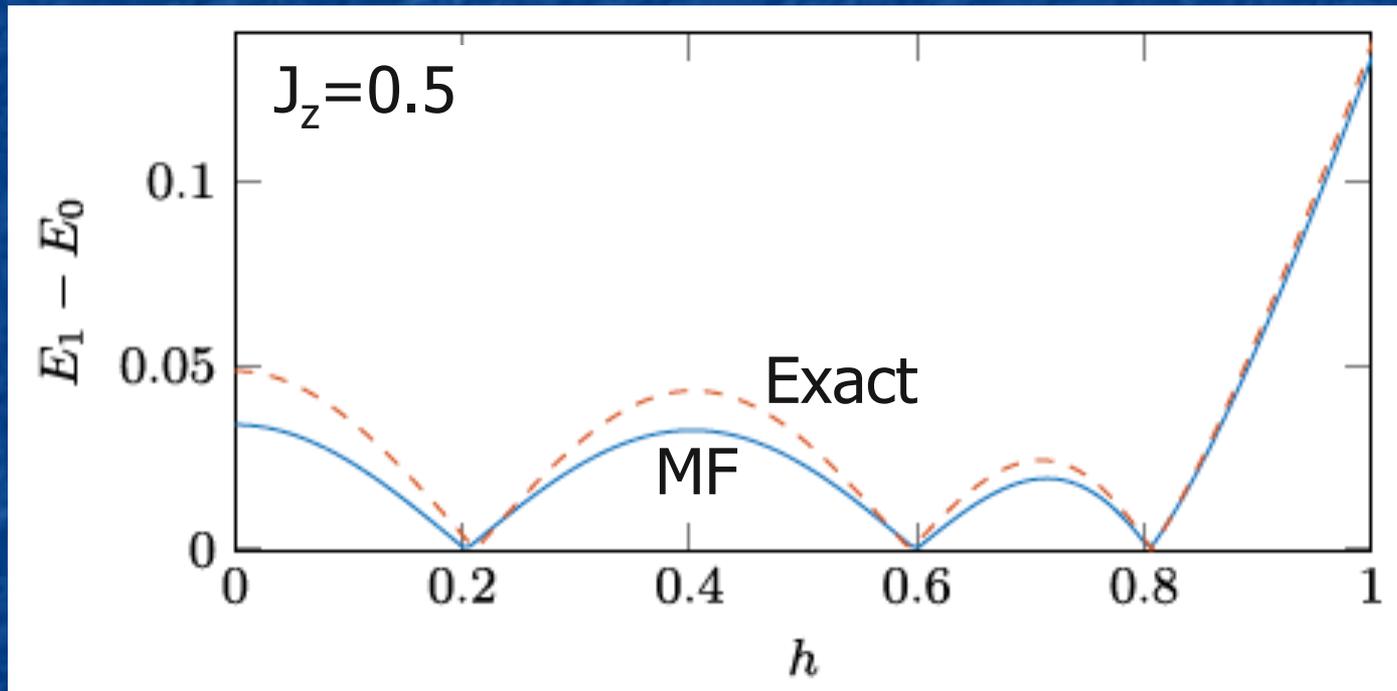
$$\mu = h + J_z(1 - 2\langle c_i^\dagger c_i \rangle), \quad t = \frac{J_x}{4} - J_z \langle c_i^\dagger c_{i+1} \rangle, \quad \text{and} \quad \Delta = \frac{J_x}{4} - J_z \langle c_i c_{i+1} \rangle$$

# Mean-field parameters



$t$  and  $\Delta$  nearly constant,  $\Delta < t$ ,  $\mu$  roughly linear in  $h$   
→ qualitatively equivalent to Kitaev

# MF theory of level crossings



Good agreement between ED and MF Kitaev model

# Conclusion/Perspectives

- **Level crossings in Co adatoms**
  - evidence of two quasi-degenerate GS
  - mechanism: Majorana oscillations
- **Other geometries** (perpendicular field, longer chains, 2D arrangements,...)
- **Observe them in SC wires**
  - see Albrecht et al, Nature 2016