

Centre for Quantum Technologies

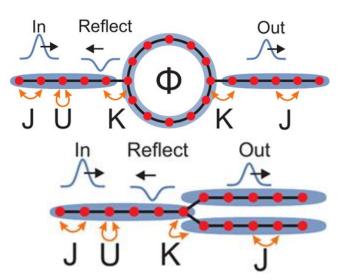


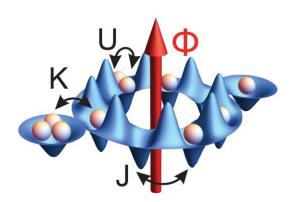
Andreev-reflection and Aharonov-Bohm dynamics in atomtronic circuits arXiv:1807.03616 arXiv:1706.05180

Tobias Haug, Rainer Dumke, Leong-Chuan

Kwek, Luigi Amico

02.08.2018



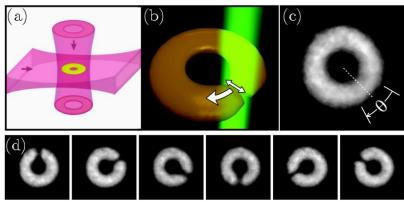




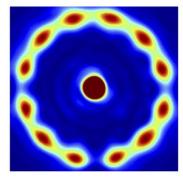


Cold atom ring and Y-junction

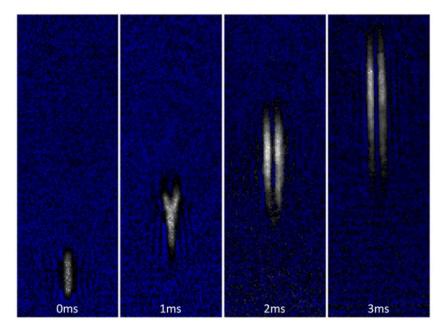
Create rotation/artificial magnetic field in atomic ring



G. Campbell, W. Phillips, C. Clark and co-workers@NIST, (2014–2015)

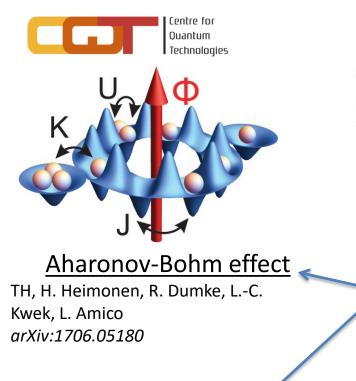


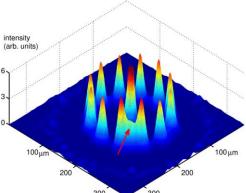
Control BEC in a Y-junction



C. Ryu and M. G. Boshier, New Journal of Physics 17, 092002 (2015).

Amico, Aghamalyan, Aukstol, Crepatz, Kwek, Dumke SREP 2014





Cold atoms + /

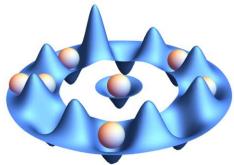
control currents

Out

Reflect

<u>Transport</u> TH, R. Dumke, L.-C. Kwek, L. Amico arXiv:1807.03616

AQUID read-out



TH, J. Tan, M. Theng, R. Dumke, L.C. Kwek, L. Amico *Phys. Rev. A* 97 (2018) *arXiv: 1707.09184*

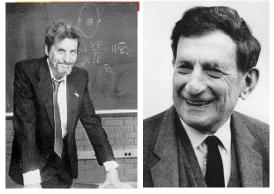
Quantum phases

gî ----

TH, L. Amico, R. Dumke, L.-C. Kwek Quantum Sci. Technol. (2018) arXiv:1612.09109

L. Amico, D. Aghamalyan, F. Auksztol, H. Crepaz, R. Dumke, L.C. Kwek Scientific reports 4 (2014)

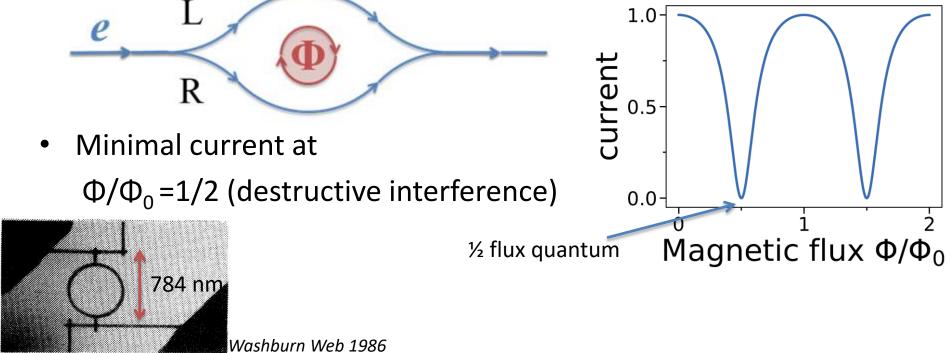
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Aharonov-Bohm effect

- Charged particle enclosing a region with magnetic field $\Delta \phi = \frac{e}{\hbar} \oint_C \mathbf{A}(\mathbf{r}) d\mathbf{r} \propto \Phi$
- Phase shift by magnetic field changes interference pattern/current







Many-body symmetry

- Two types of particles
- Fermion anti-symmetric
 - Exchange gives minus sign
 - Pauli principle
- Bosons symmetric
 - Exchange gives plus sign
 - Can occupy same space
- Bosons with infinite repulsion: Pauli principle, but still symmetric
- Jordan-Wigner to map fermion to hard-core boson

$$a_j^\dagger = e^{\left(+i\pi\sum_{k=1}^{j-1}f_k^\dagger f_k
ight)} \cdot f_j^\dagger \ a_j = e^{\left(-i\pi\sum_{k=1}^{j-1}f_k^\dagger f_k
ight)} \cdot f_j$$

- j-1 j j+1
- Mapping not exact beyond 1D \rightarrow break 1D a little?

 $|\Psi(r_a, r_b)\rangle = \pm |\Psi(r_b, r_a)\rangle$





Models

- Bose-Hubbard model (atoms on lattice with tunneling)
- Simulate with DMRG and TEBD
- **Y junction:** three 1D chains coupled with *K*

$$\mathcal{H}_{\rm S} = -\sum_{j=1}^{L_{\rm S}-1} \left(J \hat{s}_{j}^{\dagger} \hat{s}_{j+1} + \text{H.C.} \right) + \sum_{j=1}^{L_{\rm S}} \frac{U}{2} \hat{n}_{j}^{s} (\hat{n}_{j}^{s} - 1)$$
$$\mathcal{H}_{\rm I} = -K \hat{s}_{1}^{\dagger} \left(\hat{d}_{1}^{} + \hat{f}_{1}^{} \right) + \text{H.C.} \quad \text{(junction coupling)}$$

Ring with leads and artificial magnetic field Φ

$$\mathcal{H}_{\mathrm{R}} = -\sum_{j=1}^{L_{\mathrm{R}}} \left(J e^{i2\pi\Phi/L} \hat{a}_{j}^{\dagger} \hat{a}_{j+1} + \mathrm{H.C.} \right) + \frac{U}{2} \sum_{j=1}^{L_{\mathrm{R}}} \hat{n}_{j}^{a} (\hat{n}_{j}^{a} - 1)$$

• Add localized potential offset in source lead, calculate ground state, then switch off \rightarrow density wave $\mathcal{H}_{P} = -\epsilon_{D} \sum_{i=1}^{L_{S}} \exp\left(-\frac{(j-j_{0})^{2}}{2\sigma^{2}}\right) \hat{n}_{j}$





Out

tl(time)

Hard-core boson: Y junction

- On-site interaction $U \rightarrow \infty$
- 260 sites, half-filling (superfluid)
- Vary junction coupling K/J
- 0.02 K/J = 1.0drain B 250transmission *K/J*=0.5 drain •K/J = 0.2200drain A density 1 0 ی 0.00incoming junction reflection source 100 source -0.02 25 50 0 t (time) 20 40 60

Three regimes (J=1)

- K≈1: negative reflection (similar to Andreev reflection) (transmission 4/3, reflection -1/3)
- *K*≈0.5: no/positive-negative reflection
- *K*≈0.2: positive reflection

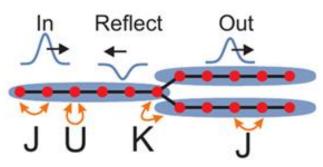
Also at bosonic interfaces: I. Zapata and F. Sols, Phys. Rev. Lett. 102, 180405 (2009)

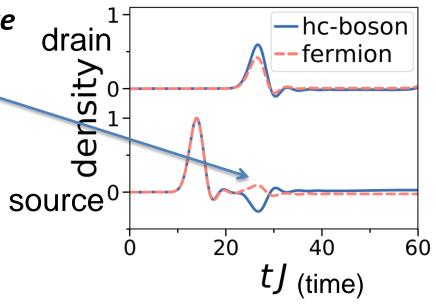




Fermion vs hard-core boson

- Spinless fermions *do not have Andreev-reflection peak* for
 strongly coupled junction
- Jordan-Wigner cannot be applied on Y-junction









Gross-Pitaevskii

- Limit of weak interaction and many atoms
- →Gross-Pitaevskii equation (GPE)

$$i\partial_t \Psi(x,t) = \left[\frac{1}{2}(-i\partial_x - A(x,t))^2 + V(x) + g\left|\Psi(x,t)\right|^2\right]\Psi(x,t)$$

- Linearize equation $\partial_t^2 \delta n = g n_0 \partial_x^2 \delta n$ (small excitation \rightarrow sound waves)
- **Y-junction**: solve at a boundary between three chains, assume
 - current conserved
 - Density/ derivative of phase continuous at boundary

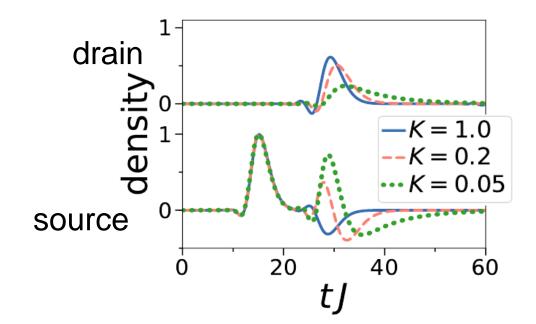
→ 4/3 transmission, -1/3 reflection for strong coupling (GPE).
Also valid for Bose-Hubbard in strong-coupling





Gross-Pitaevskii Y-junction

- For GPE: Total transmission is always 4/3, reflection -1/3, independent of coupling K (dynamics just stretched longer in time) → contrast to Bose-Hubbard (coefficient depend on K)
- Even for weakest coupling *K* there is no purely positive reflection \rightarrow No weak coupling regime
- Nonlinear wave (GPE) vs many-body dynamics (BHM)?

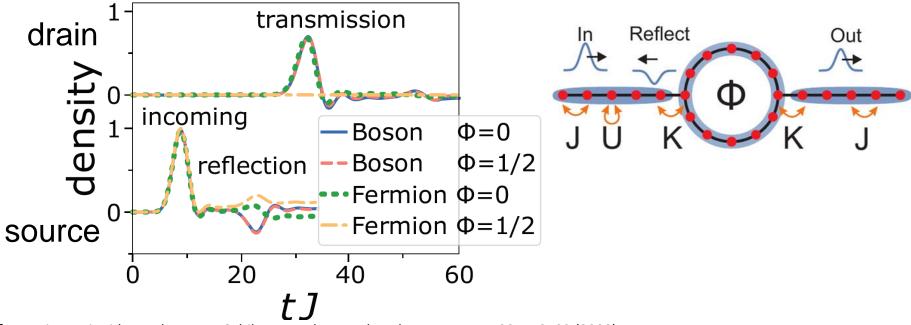






Ring dynamics

- Simulate full closed system ring + long leads (hard-core bosons)
- Initialize small density excitation in source lead
- Spinless fermions: Aharonov-Bohm effect, total reflection $\Phi=1/2$
- Interacting bosons independent of flux (no Aharonov-Bohm!)



Also for Luttinger Liquid: A. Tokuno, M. Oshikawa, and E. Demler, Phys. Rev. Lett. 100, 140402 (2008).





Gross-Pitaevskii

• Ring with Gross-Pitaesvkii equation (GPE) with flux A

$$i\partial_t \Psi(x,t) = \left[\frac{1}{2}(-i\partial_x - A(x,t))^2 + V(x) + g\left|\Psi(x,t)\right|^2\right]\Psi(x,t)$$

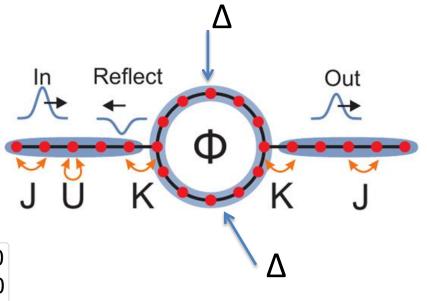
- Linearize equation on top of condensate $\partial_t^2 \delta n = \left[gn_0 - (A - v_0)^2 \right] \partial_x^2 \delta n + 2(A - v_0) \partial_t \partial_x \delta n$ $c_{\pm} = \sqrt{gn_0} \pm (A - v_0)$
- **Ring with flux A**: Flux causes a (direction dependent) shift in speed of sound, but no change in interference pattern
- →No Aharonov-Bohm effect for excitations on top of a bosonic condensate

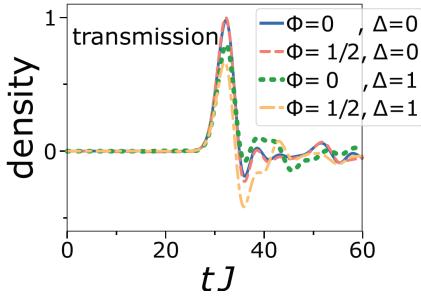




Breaking translational invariance

- Add two potential barriers Δ into the ring for Bose-Hubbard
- Restores flux dependence



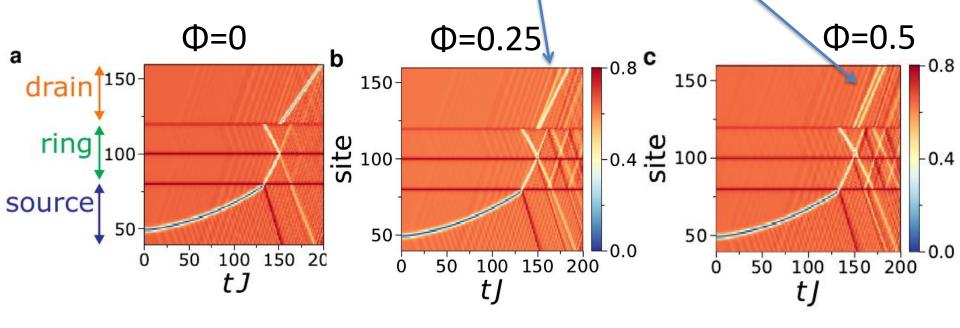






Grey soliton in ring

- Gray soliton passing through ring
- Soliton splits into two in transmission with flux
- \rightarrow Use to measure rotation

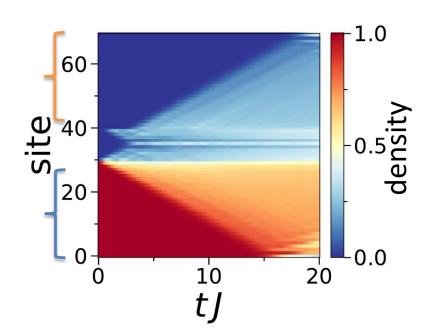


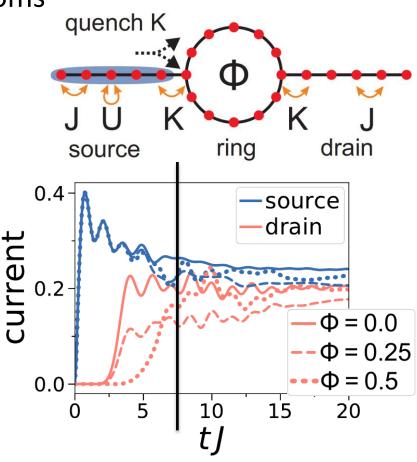


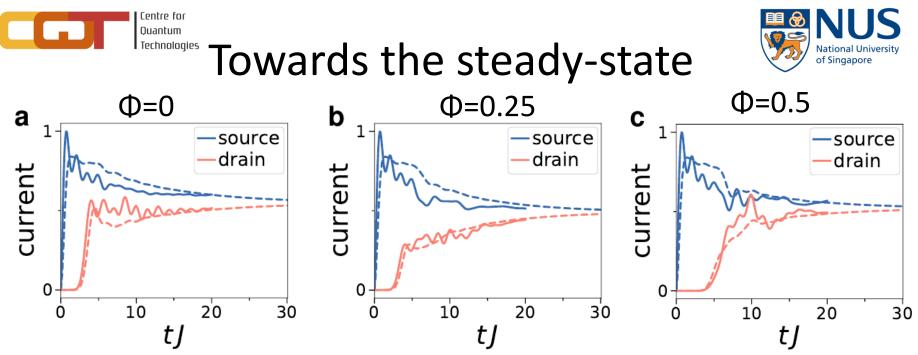


Quench of ring

- Simulate full closed system ring + long leads (hard-core bosons)
- Initially fill only source lead with atoms
- Switch on coupling K to ring (K=1)







- <u>Solid</u>: Full simulation of ring+lead with **tDMRG**
- <u>Dashed</u>: Approximate leads as Markovian baths: Open system with Lindblad + ring
- For $\Phi=0.5$, current into drain starts later
- For Φ =0.25, current increases slower
- Dynamics depends on flux, however steady-state nearly independent of flux





Conclusion

<u>Y-junction:</u>

- Different types of reflection (Andreev-like, positive-negative, positive) controlled by the coupling strength
- GPE shows different behavior than Bose-Hubbard (no weak coupling regime)

<u>Ring:</u>

- Aharonov-Bohm effect **suppressed for bosons**
- Can be restored by breaking translational invariance
- Solitons could measure flux in ring device
- Quench dynamics is **flux dependent**, **steady-state not**

arXiv:1807.03616 arXiv:1706.05180