

Grassmannian Geometry of Scattering Amplitudes

LECTURE 4

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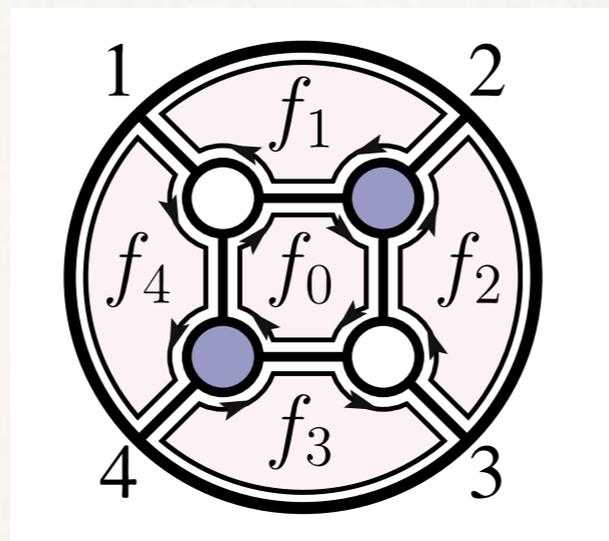
Center for Quantum Mathematics and Physics (QMAP)

University of California, Davis

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Positive Grassmannian

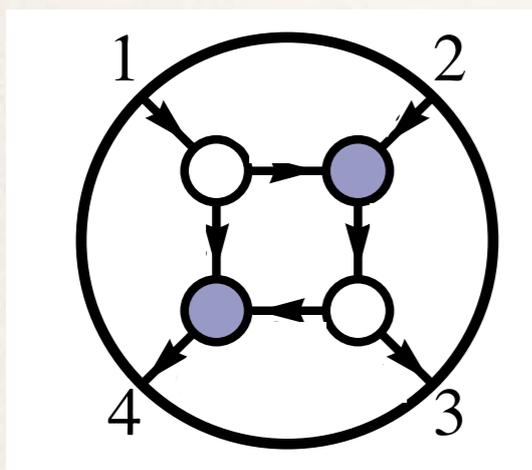
❖ Face variables



with the property

$$\prod_j f_j = -1$$

❖ Perfect orientation



Elements of $(k \times n)$ matrix

$$c_{ab} = - \sum_{\Gamma} \prod_j (-f_j)$$

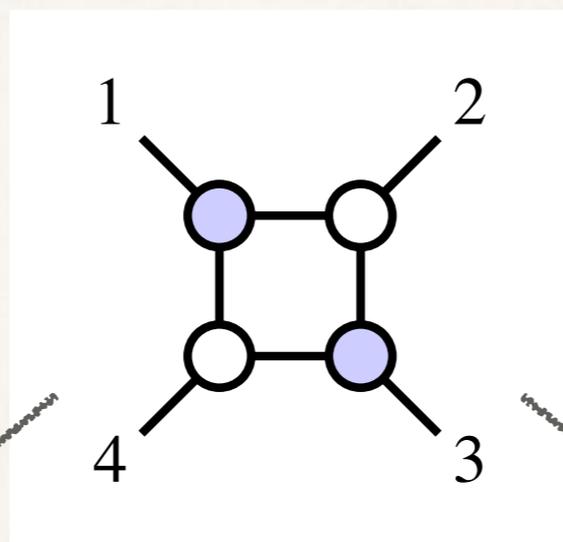
$$k \begin{pmatrix} * & * & * & \dots & * \\ * & * & * & \dots & * \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ * & * & * & \dots & * \end{pmatrix}^n$$

Cell in positive Grassmannian $G_+(k, n)$

$$k \begin{vmatrix} * & * & \dots & * \\ * & * & \dots & * \\ \vdots & \vdots & \vdots & \vdots \\ * & * & \dots & * \end{vmatrix}^k \geq 0$$

Surprising connection

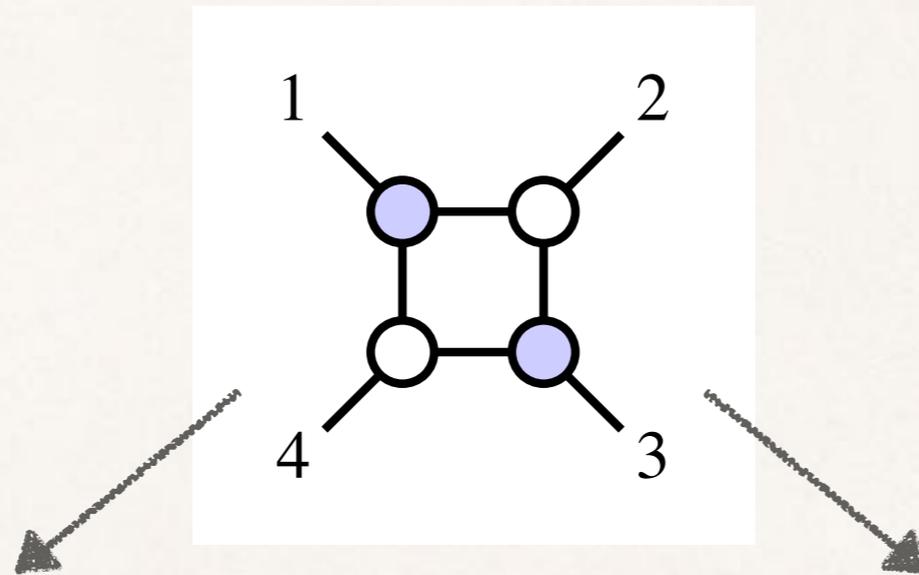
Connection



$$R = \mathcal{M}_1^{tree} \mathcal{M}_2^{tree} \mathcal{M}_3^{tree} \mathcal{M}_4^{tree}$$

$$C = \begin{pmatrix} 1 & 0 & f_0 f_3 f_4 & f_4(1 - f_0) \\ 0 & 1 & -f_0 f_1 f_3 f_4 & -f_0 f_1 f_4 \end{pmatrix}$$

Connection



$$R = \mathcal{M}_1^{tree} \mathcal{M}_2^{tree} \mathcal{M}_3^{tree} \mathcal{M}_4^{tree}$$

$$C = \begin{pmatrix} 1 & 0 & f_0 f_3 f_4 & f_4(1 - f_0) \\ 0 & 1 & -f_0 f_1 f_3 f_4 & -f_0 f_1 f_4 \end{pmatrix}$$

$$R = \int \frac{df_0}{f_0} \frac{df_1}{f_1} \frac{df_2}{f_2} \frac{df_3}{f_3} \delta(C \cdot Z)$$

$$Z = (\lambda, \tilde{\lambda}, \tilde{\eta})$$

Momentum conservation

$$\delta(C \cdot Z) = \delta(C \cdot \tilde{\lambda}) \delta(C^\perp \cdot \lambda)$$

- ❖ Simple motivation: linearize momentum conservation

$$\delta(P) = \delta \left(\sum_a \lambda_a \tilde{\lambda}_a \right)$$

- ❖ We want to write it as two linear factors

$$\delta \left(C_{ab} \tilde{\lambda}_b \right) \delta \left(D_{ab} \lambda_b \right)$$

and get the condition: $D_{ab} = C_{ab}^\perp$

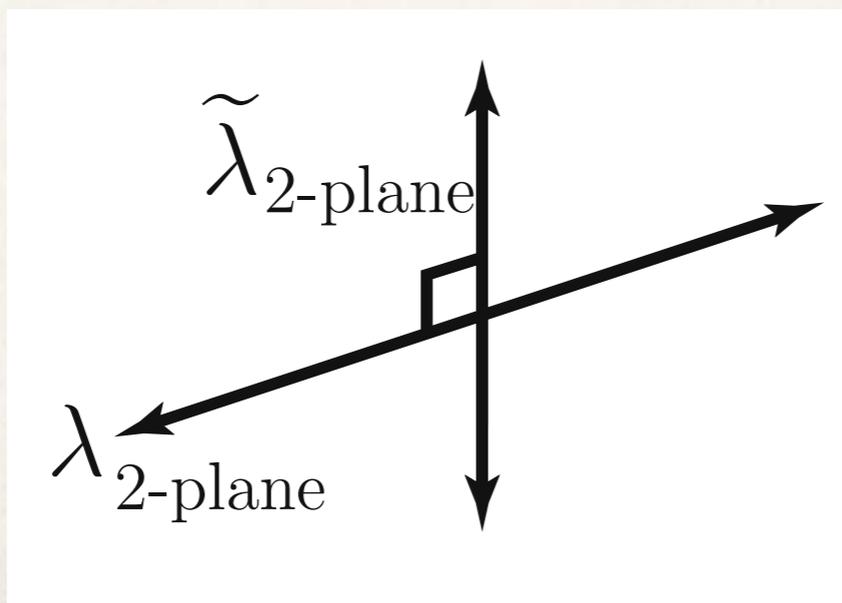
Geometry of delta function

- ❖ 2-planes for λ and $\tilde{\lambda}$ in n-dimensions

$$\lambda \equiv \begin{pmatrix} \lambda_1^1 & \lambda_2^1 & \cdots & \lambda_n^1 \\ \lambda_1^2 & \lambda_2^2 & \cdots & \lambda_n^2 \end{pmatrix} \Leftrightarrow (\lambda_1 \ \lambda_2 \ \cdots \ \lambda_n)$$

same for
 $\tilde{\lambda}$

- ❖ Momentum conservation



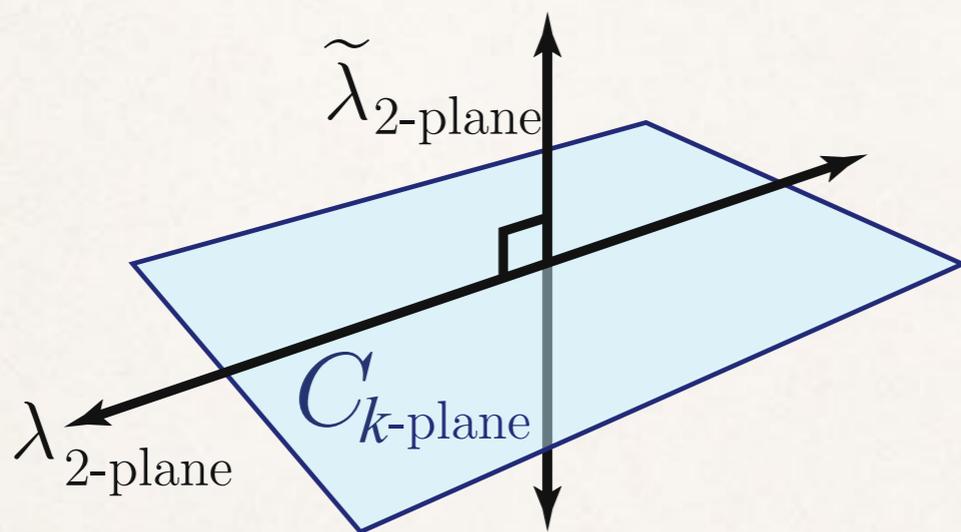
Two planes are orthogonal

$$\delta\left(\sum_k p_k\right) = \delta(\lambda \cdot \tilde{\lambda})$$

in the n-dimensional
space

Geometry of delta function

- ❖ Introduce an auxiliary k-plane C in n-dimensions



C is orthogonal to $\tilde{\lambda}$

$$\delta(C \cdot \tilde{\lambda})$$

C contains λ

= λ orthogonal to C^\perp

$$\delta(C^\perp \cdot \lambda)$$

- ❖ This forces λ and $\tilde{\lambda}$ to be orthogonal

$$\delta(\lambda \cdot \tilde{\lambda})$$

Logarithmic form

$$R = \int \frac{df_1}{f_1} \frac{df_2}{f_2} \cdots \frac{df_m}{f_m} \delta(C \cdot Z)$$

❖ Delta functions

$$\delta(C \cdot Z) \equiv \delta(C \cdot \tilde{\lambda}) \times \delta(C^\perp \cdot \lambda) \times \delta(C \cdot \tilde{\eta})$$


$$\delta(\lambda \cdot \tilde{\lambda})$$

Solves for some / all
parameters f_j

For example

$$\delta \left(f_1 - \frac{\langle 12 \rangle}{\langle 13 \rangle} \right)$$

Then use $\int \frac{df}{f} \delta(f - f_0) = \frac{1}{f_0}$

Logarithmic form

$$R = \int \frac{df_1}{f_1} \frac{df_2}{f_2} \cdots \frac{df_m}{f_m} \delta(C \cdot Z)$$

❖ Delta functions

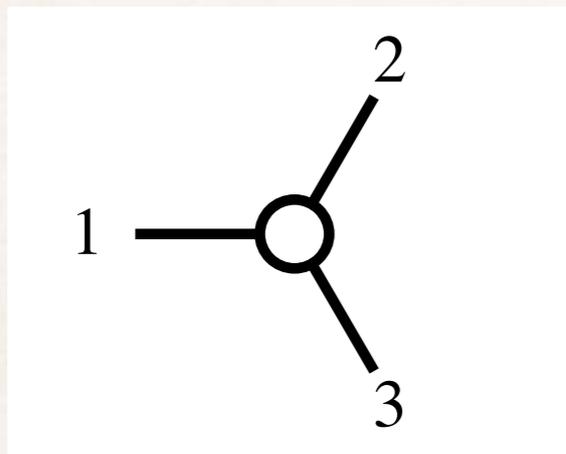
$$\delta(C \cdot Z) \equiv \delta(C \cdot \tilde{\lambda}) \times \delta(C^\perp \cdot \lambda) \times \delta(C \cdot \tilde{\eta})$$

Depending on the dimensionality

Polynomial in fermionic $\tilde{\eta}$

- We solve for all f_j
- We solve for all f_j and in addition delta functions for $\lambda, \tilde{\lambda}$
- Not enough delta functions to solve for all f_j - some left unspecified

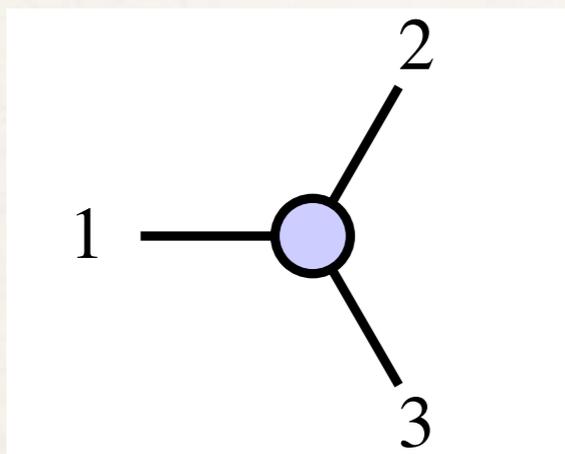
Derivation: starting with 3pt



$$C = \begin{pmatrix} 1 & z_1 & z_2 \end{pmatrix}$$

$$A_3^{(1)} = \int \frac{dz_1 dz_2}{z_1 z_2} \delta^{(1)}(C \cdot Z)$$

$$\delta^{(1)}(C \cdot Z) = \delta^{1 \times 2}(C \cdot \tilde{\lambda}) \delta^{2 \times 2}(\lambda \cdot C^\perp) \delta^{1 \times 4}(C \cdot \tilde{\eta})$$



$$C = \begin{pmatrix} 1 & 0 & z_1 \\ 0 & 1 & z_2 \end{pmatrix}$$

$$A_3^{(2)} = \int \frac{dz_1 dz_2}{z_1 z_2} \delta^{(2)}(C \cdot Z)$$

$$\delta^{(2)}(C \cdot Z) = \delta^{2 \times 2}(C \cdot \tilde{\lambda}) \delta^{2 \times 1}(\lambda \cdot C^\perp) \delta^{2 \times 4}(C \cdot \tilde{\eta})$$

New procedure

- ❖ Write the amplitude as a sum of on-shell diagrams using recursion relations
- ❖ For each on-shell diagram construct the C-matrix using the boundary measurement
- ❖ Write logarithmic form which calculates the diagram

Definition of the theory

- ❖ Why is this for N=4 SYM? What about other theories?
- ❖ Diagrams and connection to Grassmannian is general
- ❖ Specific for theory: differential form

N=4 SYM:

$$\Omega = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \dots \frac{d\alpha_n}{\alpha_n} \delta(C \cdot Z)$$

Definition of the theory

- ❖ Why is this for N=4 SYM? What about other theories?
- ❖ Diagrams and connection to Grassmannian is general
- ❖ Specific for theory: differential form

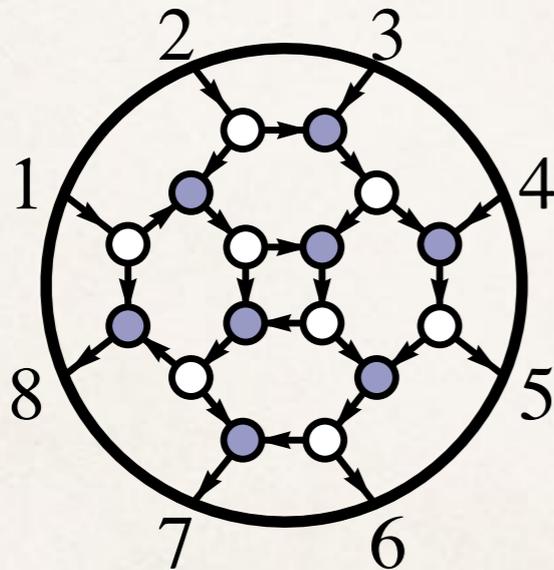
General QFT:

$$\Omega = F(\alpha) \delta(C \cdot Z)$$

- ❖ In a sense $F(\alpha)$ defines a theory (as Lagrangian does)

Exploring space of theories

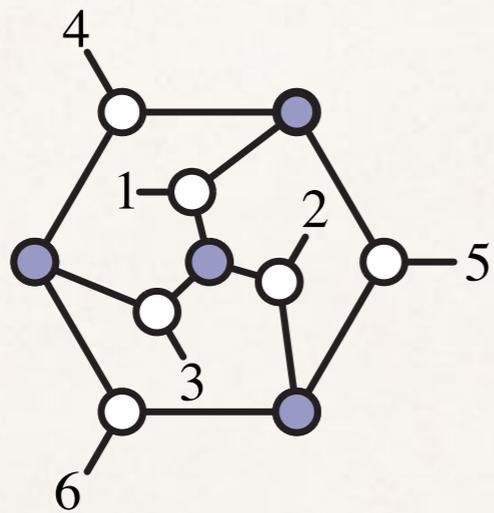
- ❖ One step at a time away from N=4 SYM
- ❖ **Case 1:** planar N<4 SYM (including QCD)



$$\Omega = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \dots \frac{d\alpha_n}{\alpha_n} \cdot \mathcal{J}(\alpha) \delta(C \cdot Z)$$

Exploring space of theories

❖ Case 2: Non-planar N=4 SYM



same form as planar

$$\Omega = \frac{d\alpha_1}{\alpha_1} \frac{d\alpha_2}{\alpha_2} \cdots \frac{d\alpha_n}{\alpha_n} \delta(C \cdot Z)$$

less is known on mathematical side

❖ Case 3: N=8 Supergravity

Both have natural extensions to lower SUSY

$$\Omega = \frac{d\alpha_1}{\alpha_1^3} \frac{d\alpha_2}{\alpha_2^3} \cdots \frac{d\alpha_m}{\alpha_m^3} \prod_v \Delta_v \cdot \delta(C \cdot Z)$$

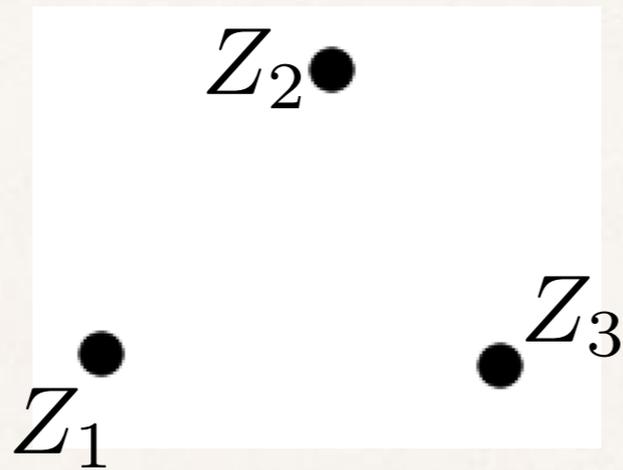
Back to planar $N=4$ SYM theory

- ❖ Each on-shell diagram has a nice geometric interpretation
- ❖ To get the amplitude we use the recursion relations = consequence of the factorization of tree-level amplitudes
- ❖ From geometric point of view: why this particular sum?
- ❖ Goal: to find the geometric formulation for the full amplitude

Inside of the triangle

Inside of the triangle

- ✦ Let us consider three points in a projective plane



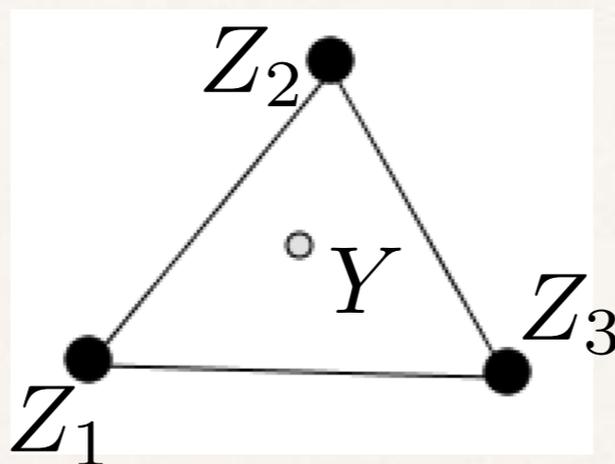
$$Z_j = \begin{pmatrix} * \\ * \\ * \end{pmatrix} \quad Z_j \sim tZ_j$$

We can also fix

$$Z_j = \begin{pmatrix} 1 \\ a_j \\ b_j \end{pmatrix}$$

Inside of the triangle

- ❖ Point inside the triangle



$$Z_j = \begin{pmatrix} * \\ * \\ * \end{pmatrix} \quad Z_j \sim tZ_j$$

We can also fix

$$Z_j = \begin{pmatrix} 1 \\ a_j \\ b_j \end{pmatrix}$$

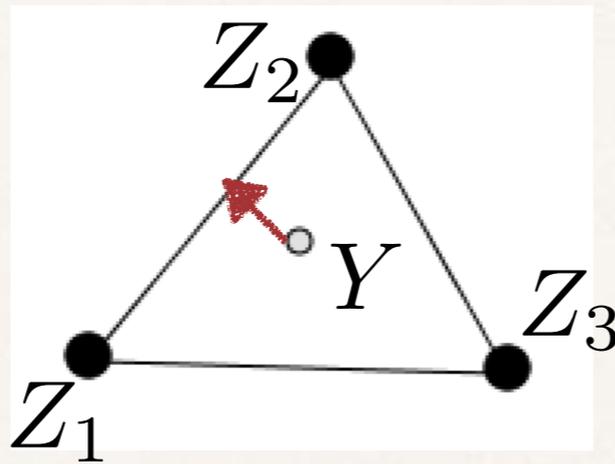
- ❖ Point inside the triangle

$$Y = c_1 Z_1 + c_2 Z_2 + c_3 Z_3 \quad c_1, c_2, c_3 > 0$$

Projective: one of c_j can be fixed to 1

Inside of the triangle

- ❖ Point inside the triangle



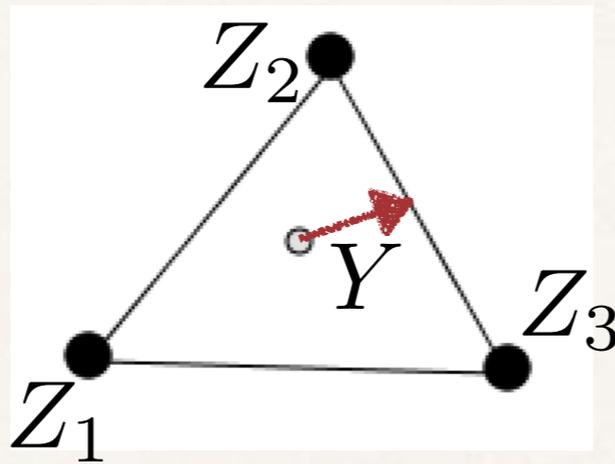
$$Y = c_1 Z_1 + c_2 Z_2 + c_3 Z_3$$

On the boundary

$$c_3 = 0$$

Inside of the triangle

- ❖ Point inside the triangle



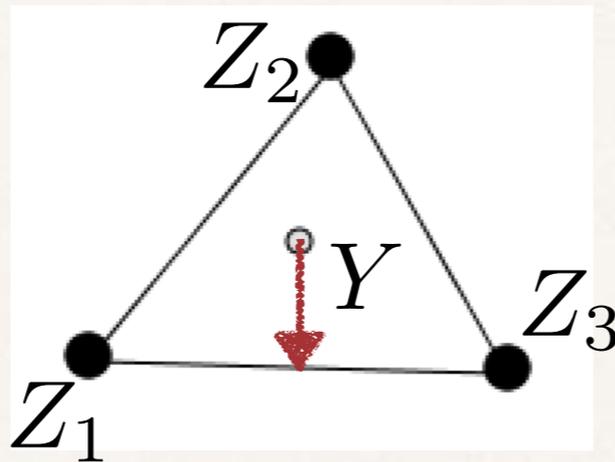
$$Y = c_1 Z_1 + c_2 Z_2 + c_3 Z_3$$

On the boundary

$$c_1 = 0$$

Inside of the triangle

- ❖ Point inside the triangle



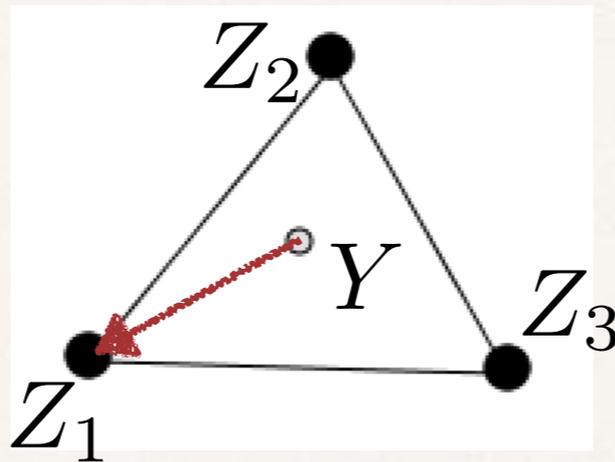
$$Y = c_1 Z_1 + c_2 Z_2 + c_3 Z_3$$

On the boundary

$$c_2 = 0$$

Inside of the triangle

- ❖ Point inside the triangle



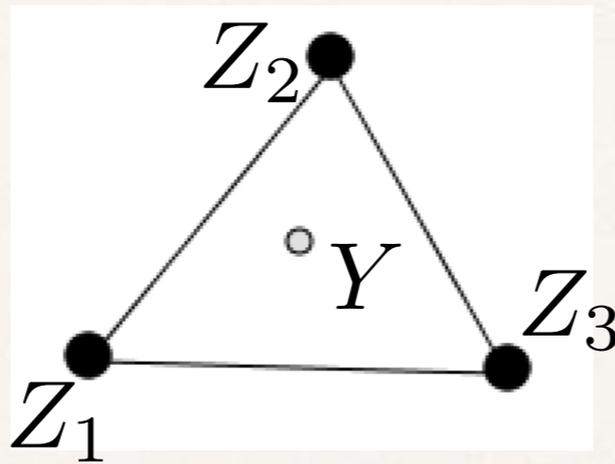
$$Y = c_1 Z_1 + c_2 Z_2 + c_3 Z_3$$

On the boundary

$$c_2 = c_3 = 0$$

Logarithmic form

- ❖ Point inside the triangle



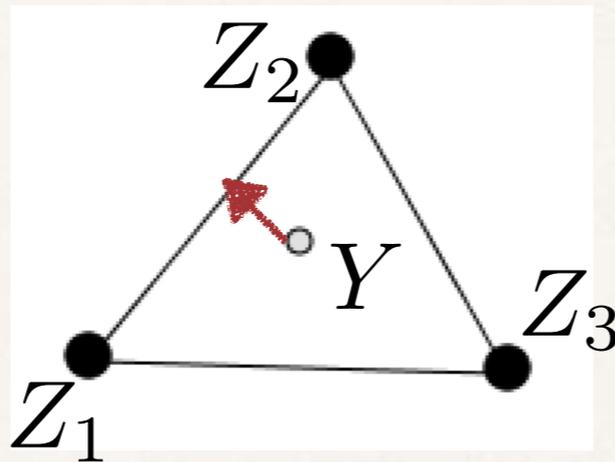
$$Y = Z_1 + c_2 Z_2 + c_3 Z_3$$

- ❖ Form with **logarithmic singularities on boundaries**

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3}$$

Logarithmic form

- ❖ Point inside the triangle



$$c_3 = 0$$

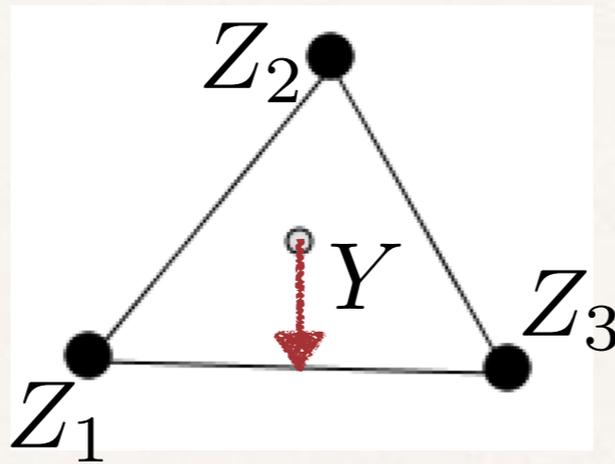
$$Y = Z_1 + c_2 Z_2 + \cancel{c_3 Z_3}$$

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \rightarrow \frac{dc_2}{c_2}$$

Logarithmic form

- ❖ Point inside the triangle



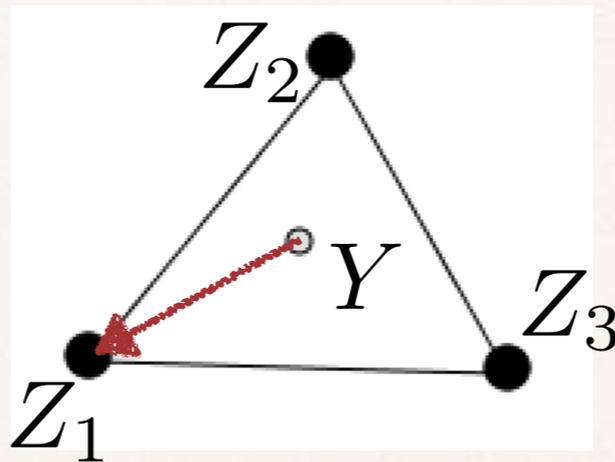
$$Y = Z_1 + \overset{c_2 = 0}{\cancel{c_2 Z_2}} + c_3 Z_3$$

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \rightarrow \frac{dc_3}{c_3}$$

Logarithmic form

- ❖ Point inside the triangle



$$Y = Z_1 + \cancel{c_2 Z_2} + \cancel{c_3 Z_3}$$

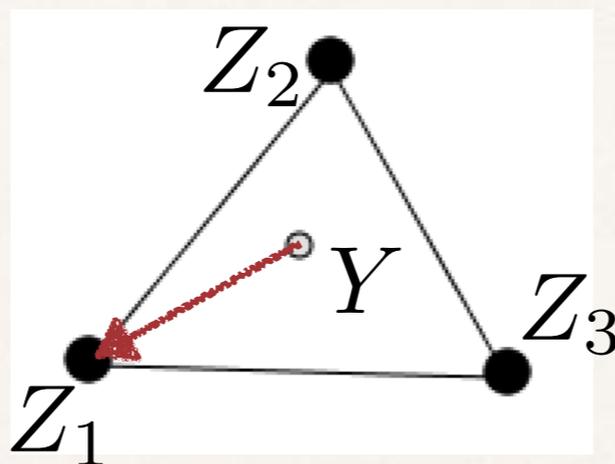
$c_2 = c_3 = 0$

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \rightarrow \frac{dc_3}{c_3} \rightarrow 1$$

Logarithmic form

- ❖ Point inside the triangle



$$Y = Z_1 + \cancel{c_2 Z_2} + \cancel{c_3 Z_3} \quad c_2 = c_3 = 0$$

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \rightarrow \frac{dc_3}{c_3} \rightarrow 1$$

- ❖ Other boundaries can correspond to $c_2, c_3 \rightarrow \infty$

Logarithmic form

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \quad \langle X_1 X_2 X_3 \rangle = \epsilon_{abc} X_1^a X_2^b X_3^c$$
$$d^2 Y = dc_2 dc_3 Z_2 Z_3$$

- ❖ Solve for c_2, c_3 from $Y = Z_1 + c_2 Z_2 + c_3 Z_3$

$$\begin{array}{c} \downarrow \\ c_2 = \frac{\langle Y 13 \rangle}{\langle Y 23 \rangle} \quad c_3 = \frac{\langle Y 12 \rangle}{\langle Y 23 \rangle} \quad dc_2 dc_3 = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 23 \rangle^3} \end{array}$$

Logarithmic form

- ❖ Form with logarithmic singularities on boundaries

$$\Omega = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \quad \langle X_1 X_2 X_3 \rangle = \epsilon_{abc} X_1^a X_2^b X_3^c$$
$$d^2 Y = dc_2 dc_3 Z_2 Z_3$$

- ❖ Solve for c_2, c_3 from $Y = Z_1 + c_2 Z_2 + c_3 Z_3$



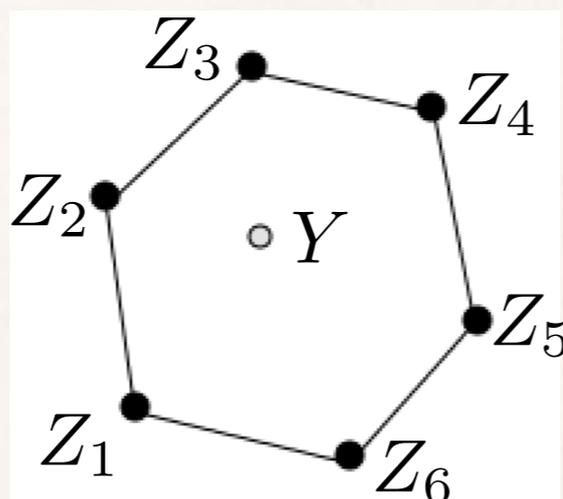
$$\Omega = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 31 \rangle}$$

Projective in all variables

Polygon

Point inside the polygon

- ✦ Consider a point inside a polygon in projective plane



$$Y = c_1 Z_1 + c_2 Z_2 + \dots + c_n Z_n$$

$$\downarrow$$

$$c_j > 0$$

interior of the polygon

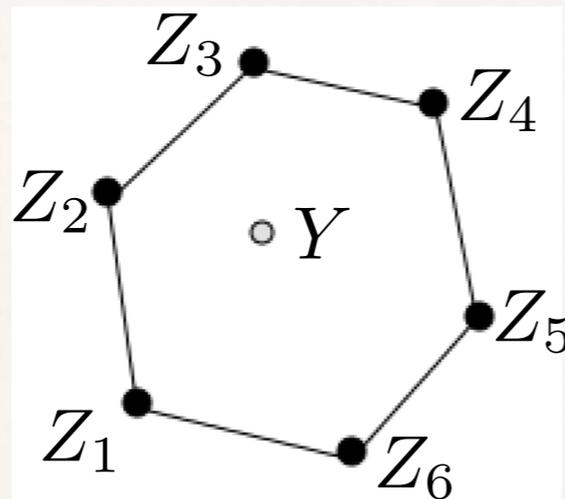
- ✦ Convex polygon: condition on points Z_i

$$Z = \begin{pmatrix} \uparrow & \uparrow & \uparrow & \uparrow & \uparrow \\ Z_1 & Z_2 & Z_3 & \dots & Z_n \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{pmatrix} \quad \text{All main minors positive}$$

$$\begin{vmatrix} * & * & * \\ * & * & * \\ * & * & * \end{vmatrix} > 0$$

Point inside the polygon

- ✦ Consider a point inside a polygon in projective plane



$$Y = c_1 Z_1 + c_2 Z_2 + \dots + c_n Z_n$$

Space of all points
inside convex polygon

More formally:

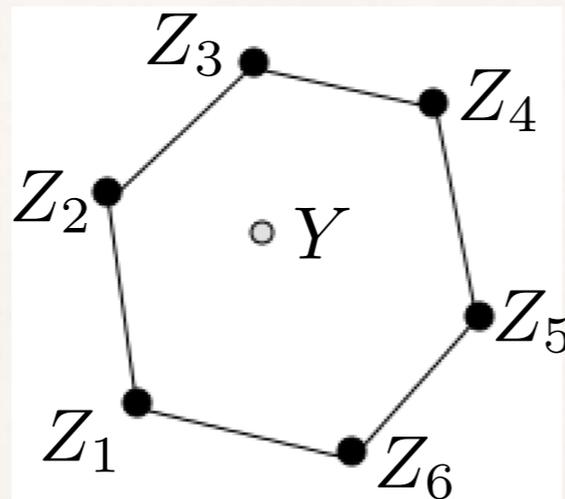
$$Y = C \cdot Z$$

$$C = \begin{pmatrix} c_1 & c_2 & c_3 & \dots & c_n \end{pmatrix} \in G_+(1, n)$$

$$Z = \begin{pmatrix} \uparrow & \uparrow & \uparrow & \uparrow & \uparrow \\ Z_1 & Z_2 & Z_3 & \dots & Z_n \\ \downarrow & \downarrow & \downarrow & \downarrow & \downarrow \end{pmatrix} \in M_+(3, n)$$

Logarithmic form

- ❖ Form with logarithmic singularities on boundaries



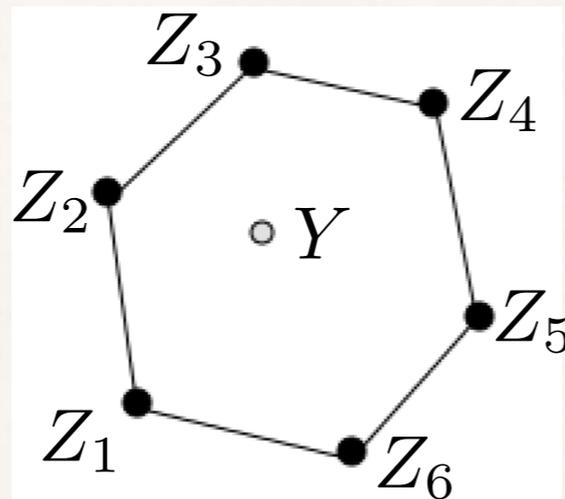
$$Y = c_1 Z_1 + c_2 Z_2 + \dots + c_n Z_n$$

First guess

$$\Omega = \frac{dc_1}{c_1} \frac{dc_2}{c_2} \dots \frac{dc_n}{c_n}$$

Logarithmic form

- ❖ Form with logarithmic singularities on boundaries



$$Y = c_1 Z_1 + c_2 Z_2 + \dots + c_n Z_n$$

First guess

~~$$\Omega = \frac{dc_1}{c_1} \frac{dc_2}{c_2} \dots \frac{dc_n}{c_n}$$~~

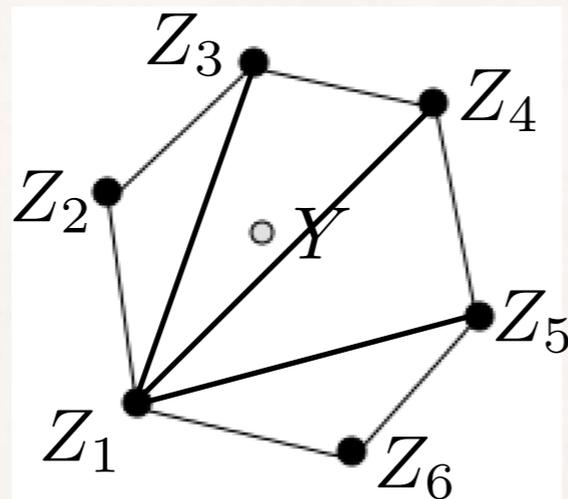
Space of all Y is only two-dimensional

- Two-form with n poles

$$\Omega \sim \frac{dc_1 dc_2 N(c_1, c_2)}{D(c_1, c_2)}$$

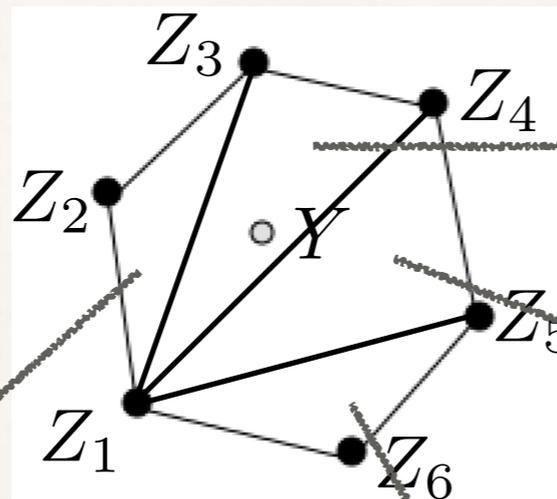
Logarithmic form

- ❖ Easiest way how to write the form is to triangulate



Logarithmic form

- ❖ Easiest way how to write the form is to triangulate



$$Y = Z_1 + c_2 Z_2 + c_3 Z_3$$
$$\Omega_1 = \frac{dc_2}{c_2} \frac{dc_3}{c_3} \quad c_2, c_3 \geq 0$$

$$Y = Z_1 + c_3 Z_3 + c_4 Z_4$$
$$\Omega_2 = \frac{dc_3}{c_3} \frac{dc_4}{c_4} \quad c_3, c_4 \geq 0$$

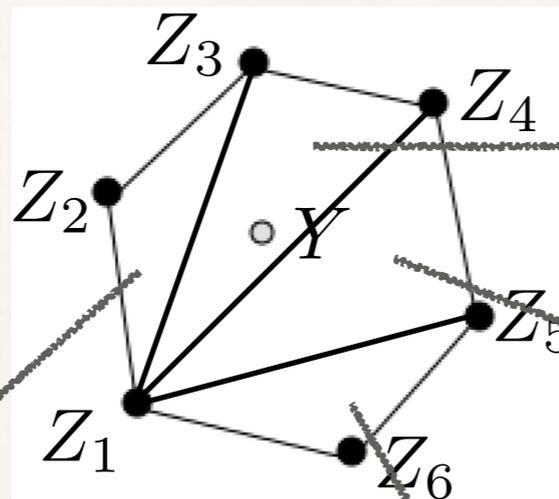
$$Y = Z_1 + c_4 Z_4 + c_5 Z_5$$
$$\Omega_3 = \frac{dc_4}{c_4} \frac{dc_5}{c_5} \quad c_4, c_5 \geq 0$$

$$Y = Z_1 + c_5 Z_5 + c_6 Z_6$$
$$\Omega_4 = \frac{dc_5}{c_5} \frac{dc_6}{c_6} \quad c_5, c_6 \geq 0$$

How to sum them?

Logarithmic form

- ❖ Easiest way how to write the form is to triangulate



$$Y = Z_1 + c_3 Z_3 + c_4 Z_4$$

$$\Omega_2 = \frac{\langle Y d^2 Y \rangle \langle 134 \rangle^2}{\langle Y 13 \rangle \langle Y 34 \rangle \langle Y 41 \rangle}$$

$$Y = Z_1 + c_4 Z_4 + c_5 Z_5$$

$$\Omega_3 = \frac{\langle Y d^2 Y \rangle \langle 145 \rangle^2}{\langle Y 14 \rangle \langle Y 45 \rangle \langle Y 51 \rangle}$$

$$Y = Z_1 + c_2 Z_2 + c_3 Z_3$$

$$\Omega_1 = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 31 \rangle}$$

$$Y = Z_1 + c_5 Z_5 + c_6 Z_6$$

$$\Omega_4 = \frac{\langle Y d^2 Y \rangle \langle 156 \rangle^2}{\langle Y 15 \rangle \langle Y 56 \rangle \langle Y 61 \rangle}$$

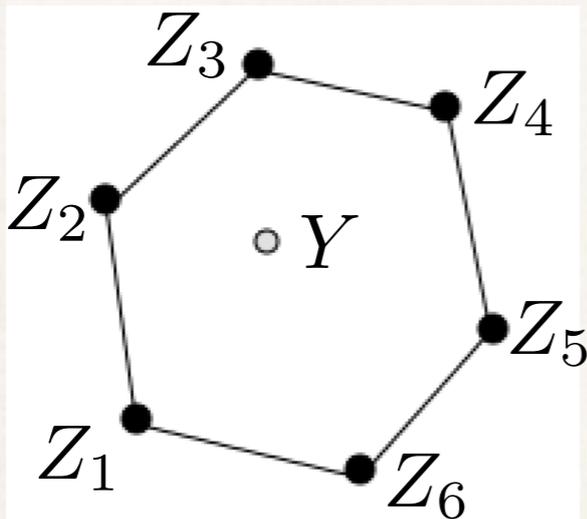
Write in projective form

Logarithmic form

- ❖ Now it makes sense to sum them

$$\Omega = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 31 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 134 \rangle^2}{\langle Y 13 \rangle \langle Y 34 \rangle \langle Y 41 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 145 \rangle^2}{\langle Y 14 \rangle \langle Y 45 \rangle \langle Y 51 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 156 \rangle^2}{\langle Y 15 \rangle \langle Y 56 \rangle \langle Y 61 \rangle}$$

- ❖ Boundaries of the polygon are $\langle Y i i + 1 \rangle = 0$

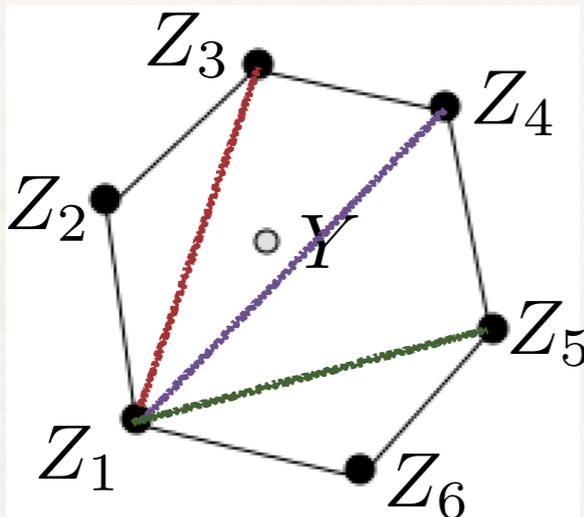


Logarithmic form

- Now it makes sense to sum them

$$\Omega = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 31 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 134 \rangle^2}{\langle Y 13 \rangle \langle Y 34 \rangle \langle Y 41 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 145 \rangle^2}{\langle Y 14 \rangle \langle Y 45 \rangle \langle Y 51 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 156 \rangle^2}{\langle Y 15 \rangle \langle Y 56 \rangle \langle Y 61 \rangle}$$

- Boundaries of the polygon are $\langle Y i i + 1 \rangle = 0$



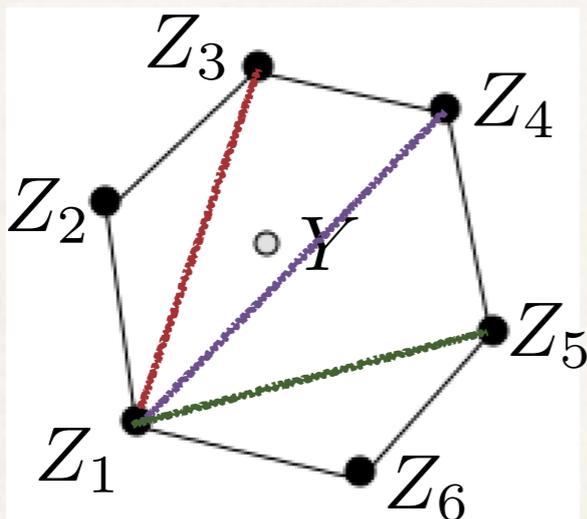
Spurious poles
Cancel in the sum

Logarithmic form

- Now it makes sense to sum them

$$\Omega = \frac{\langle Y d^2 Y \rangle \langle 123 \rangle^2}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 31 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 134 \rangle^2}{\langle Y 13 \rangle \langle Y 34 \rangle \langle Y 41 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 145 \rangle^2}{\langle Y 14 \rangle \langle Y 45 \rangle \langle Y 51 \rangle} + \frac{\langle Y d^2 Y \rangle \langle 156 \rangle^2}{\langle Y 15 \rangle \langle Y 56 \rangle \langle Y 61 \rangle}$$

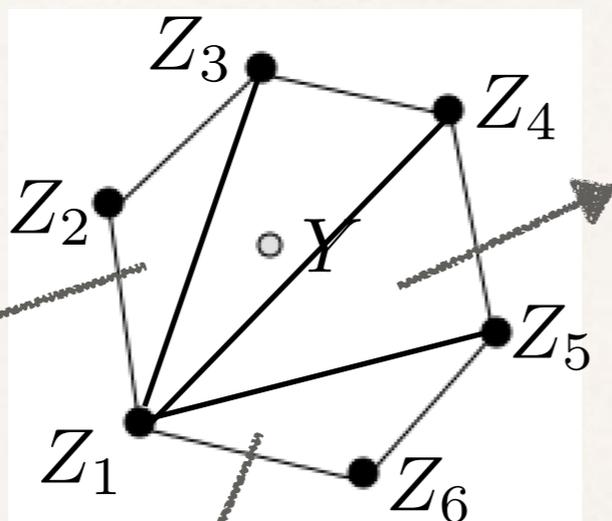
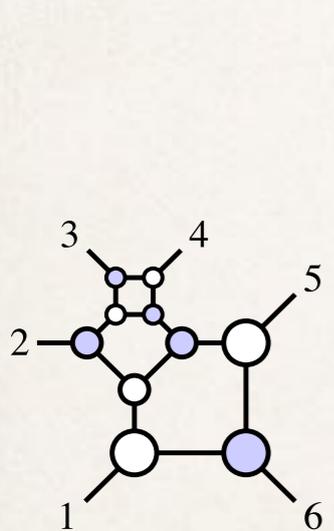
- Boundaries of the polygon are $\langle Y i i + 1 \rangle = 0$



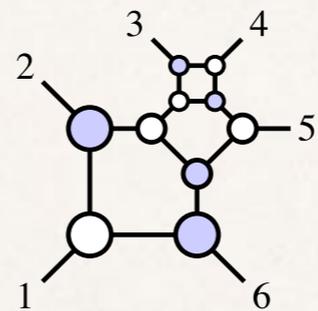
$$\Omega = \frac{\langle Y d^2 Y \rangle \mathcal{N}(Y, Z_j)}{\langle Y 12 \rangle \langle Y 23 \rangle \langle Y 34 \rangle \langle Y 45 \rangle \langle Y 56 \rangle \langle Y 61 \rangle}$$

Similarities with on-shell diagrams

- Notice some similarities with recursion relations and on-shell diagrams



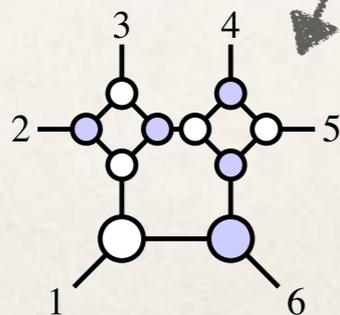
$$C = \begin{pmatrix} 1 & 0 & 0 & c_4 & c_5 & 0 \end{pmatrix} \in G_+(1,6)$$

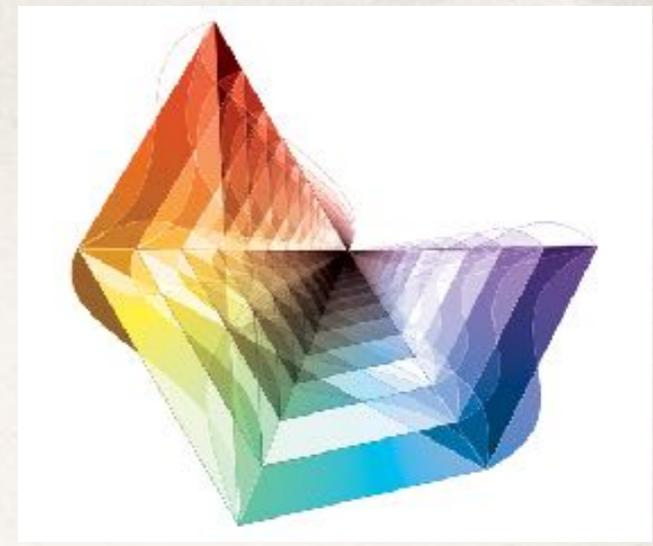


$$\Omega = \frac{dc_4}{c_4} \frac{dc_5}{c_5} \text{ looks similar}$$

$$Y = C \cdot Z \quad \delta(C \cdot Z)$$

$$\Omega = \frac{\langle Y d^2 Y \rangle \langle 145 \rangle^2}{\langle Y 14 \rangle \langle Y 45 \rangle \langle Y 51 \rangle}$$





Amplituhedron

Amplituhedron

- ❖ On-shell diagrams triangulate the bigger object which represents the scattering amplitude: Amplituhedron

$$\begin{aligned} \text{tree } Y_{\alpha}^I &= C_{\alpha a} Z_a^I \\ C_{\alpha a} &\in G_+(k, n) \\ Z_a^I &\in M_+(n, n+4) \\ Y &\in G(k, k+4) \end{aligned}$$

- ❖ Z - kinematical variables (momentum twistors)
- ❖ Extension to the loop integrand

Amplitude from Amplituhedron

- ❖ Amplitude is associated with the form with logarithmic singularities on the boundaries of the Amplituhedron

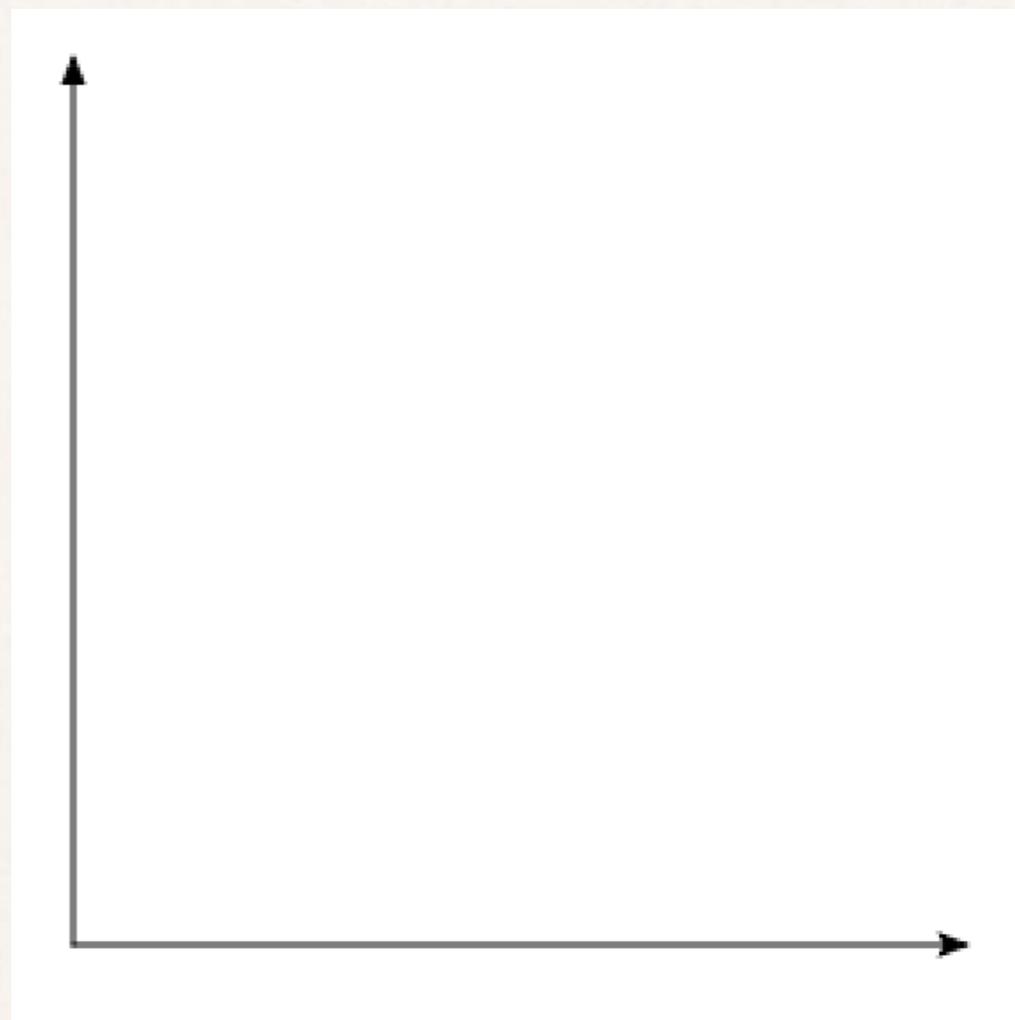
$$\Omega \sim \frac{dx}{x} \quad \text{near } x = 0$$

- ❖ More than just simple poles: $\frac{dx dy}{xy(x+y)} \xrightarrow{x=0} \frac{dy}{y^2}$
- ❖ The main work is to triangulate the space
- ❖ Instead of explaining the precise definition: example of the 4pt scattering to all loops

High school problem $gg \rightarrow gg$

High school problem $gg \rightarrow gg$

- ❖ Positive quadrant

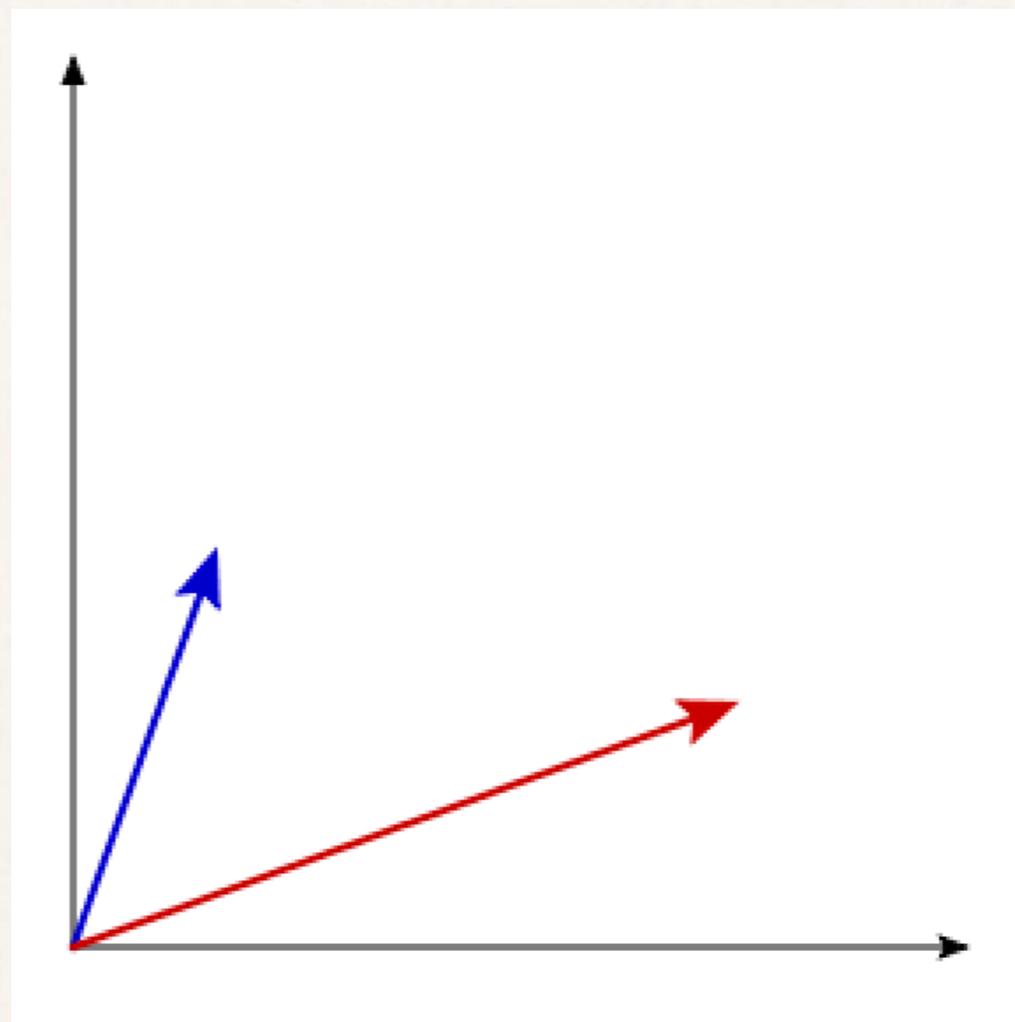


High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$



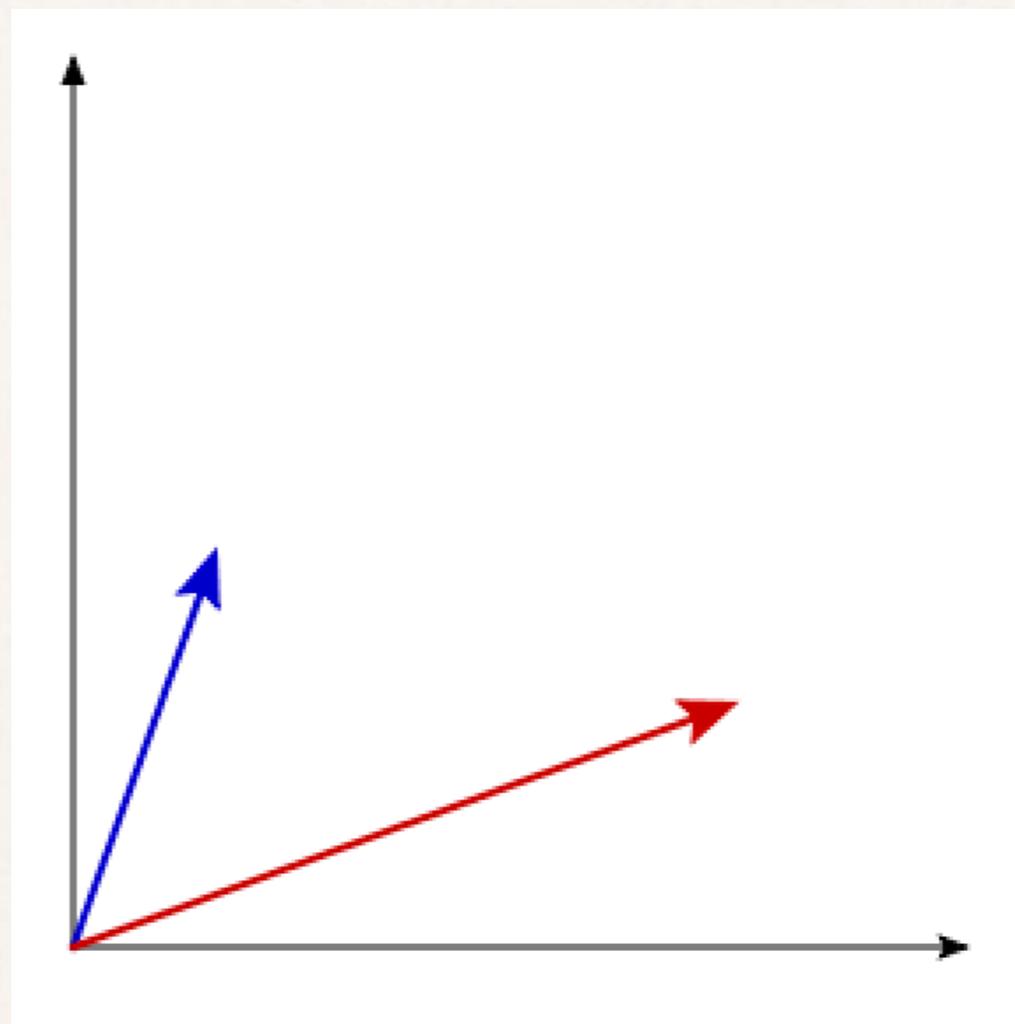
$$\text{Vol}(1) = \frac{dx_1}{x_1} \frac{dy_1}{y_1} \frac{dz_1}{z_1} \frac{dw_1}{w_1}$$

High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$



$$\text{Vol}(1) = \frac{dx_1}{x_1} \frac{dy_1}{y_1} \frac{dz_1}{z_1} \frac{dw_1}{w_1} = \text{Diagram of a square with outward-pointing lines at each corner}$$

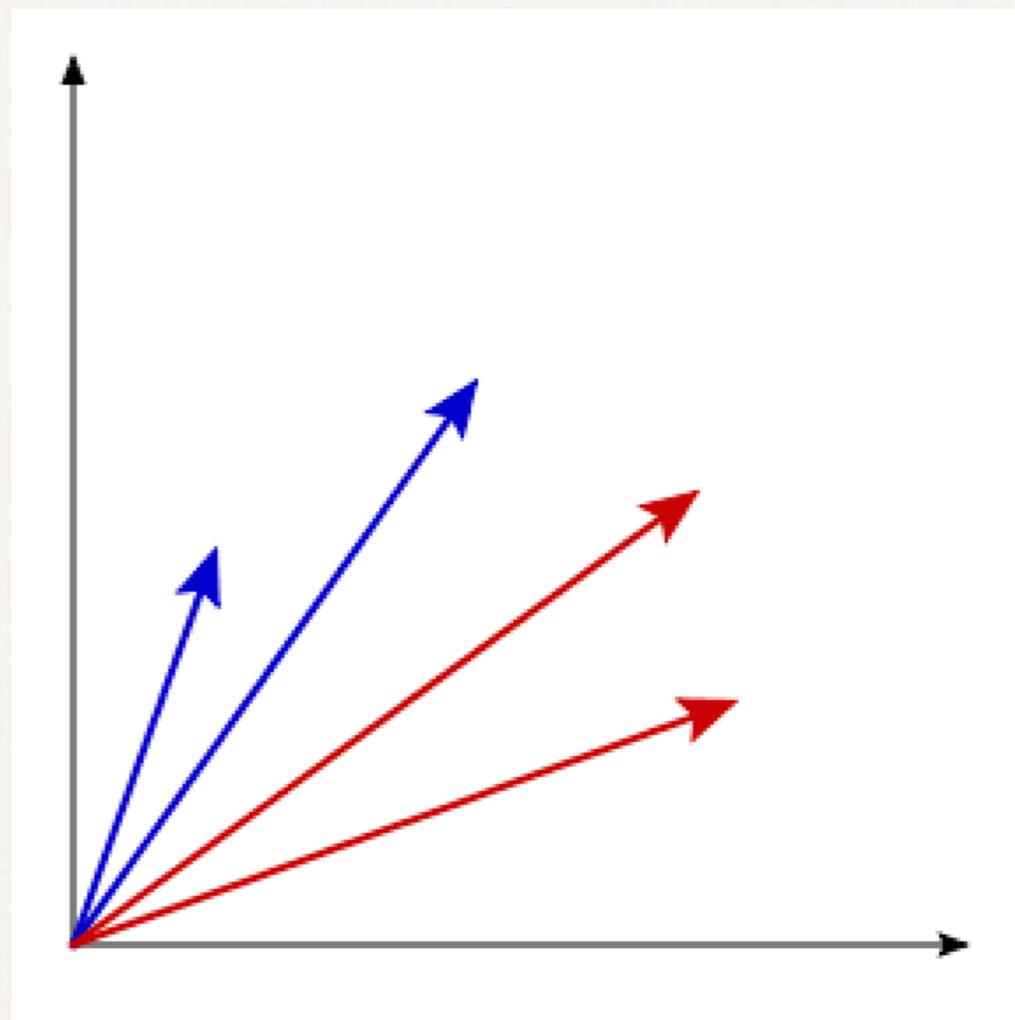
High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$

$$\vec{a}_2 = \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \quad \vec{b}_2 = \begin{pmatrix} z_2 \\ w_2 \end{pmatrix}$$



$$[\text{Vol}(1)]^2 = \frac{dx_1}{x_1} \frac{dy_1}{y_1} \frac{dz_1}{z_1} \frac{dw_1}{w_1} \frac{dx_2}{x_2} \frac{dy_2}{y_2} \frac{dz_2}{z_2} \frac{dw_2}{w_2} = \begin{array}{c} \diagup \quad \diagdown \\ \square \\ \diagdown \quad \diagup \end{array} \times \begin{array}{c} \diagup \quad \diagdown \\ \square \\ \diagdown \quad \diagup \end{array}$$

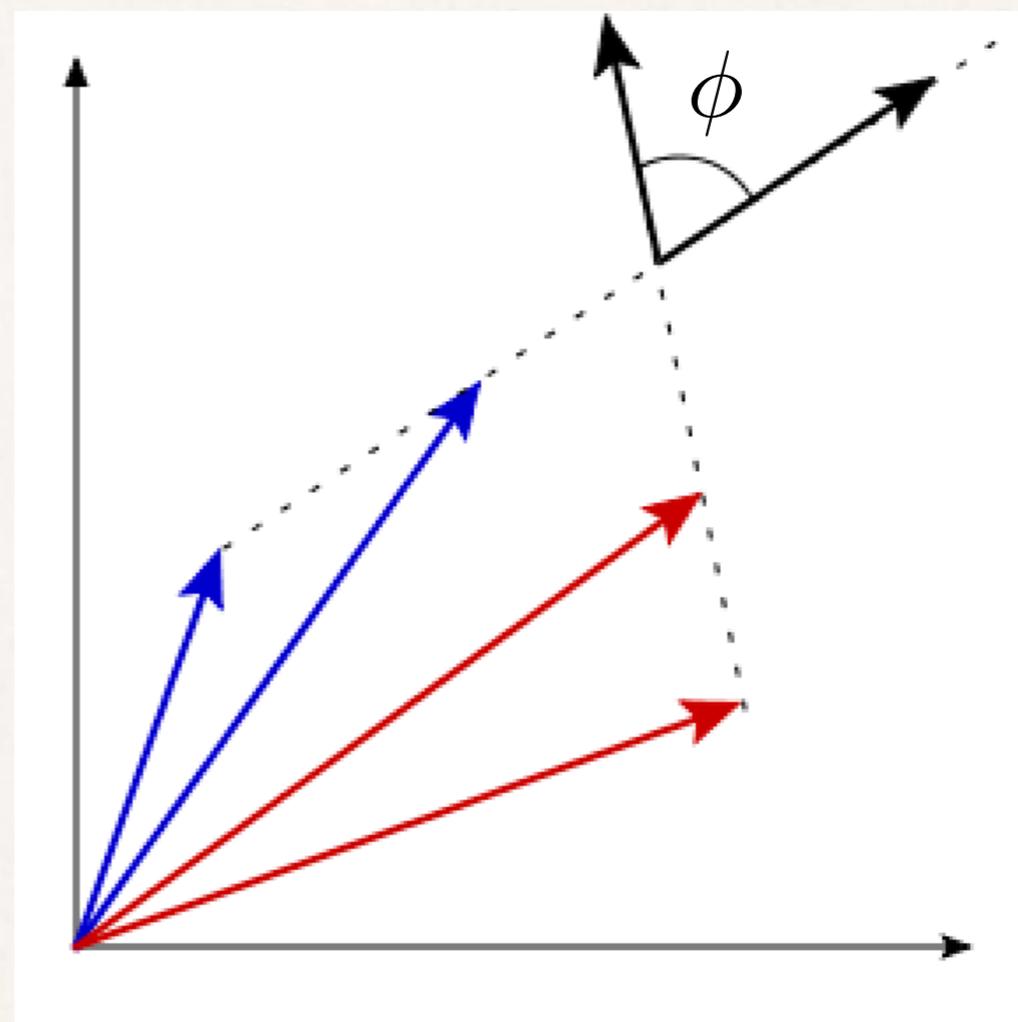
High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$

$$\vec{a}_2 = \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \quad \vec{b}_2 = \begin{pmatrix} z_2 \\ w_2 \end{pmatrix}$$



❖ Impose: $(\vec{a}_2 - \vec{a}_1) \cdot (\vec{b}_2 - \vec{b}_1) \leq 0 \quad \phi > 90^\circ$

Subset of configurations allowed: triangulate

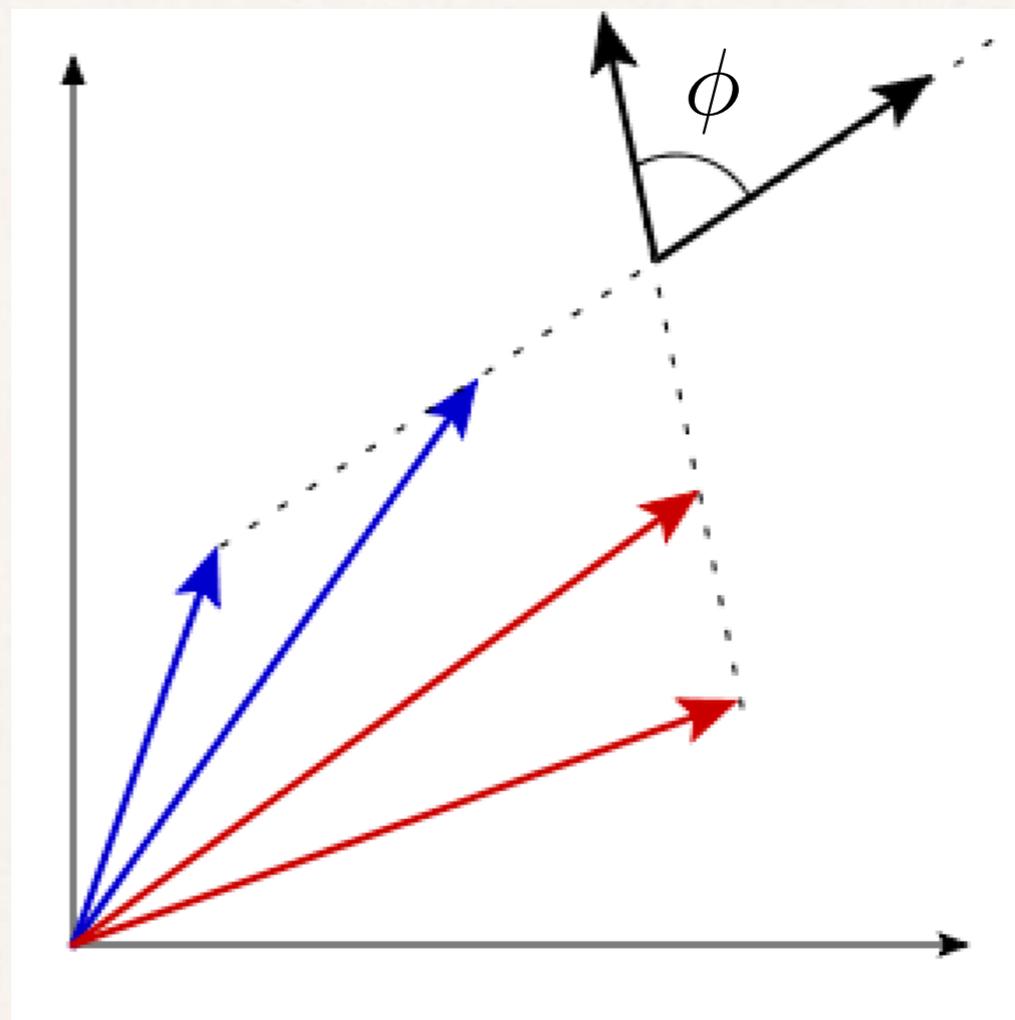
High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$

$$\vec{a}_2 = \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \quad \vec{b}_2 = \begin{pmatrix} z_2 \\ w_2 \end{pmatrix}$$



$$\text{Vol (2)} = \frac{dx_1}{x_1} \frac{dy_1}{y_1} \frac{dz_1}{z_1} \frac{dw_1}{w_1} \frac{dx_2}{x_2} \frac{dy_2}{y_2} \frac{dz_2}{z_2} \frac{dw_2}{w_2} \left[\frac{\vec{a}_1 \cdot \vec{b}_2 + \vec{a}_2 \cdot \vec{b}_1}{(\vec{a}_2 - \vec{a}_1) \cdot (\vec{b}_2 - \vec{b}_1)} \right]$$

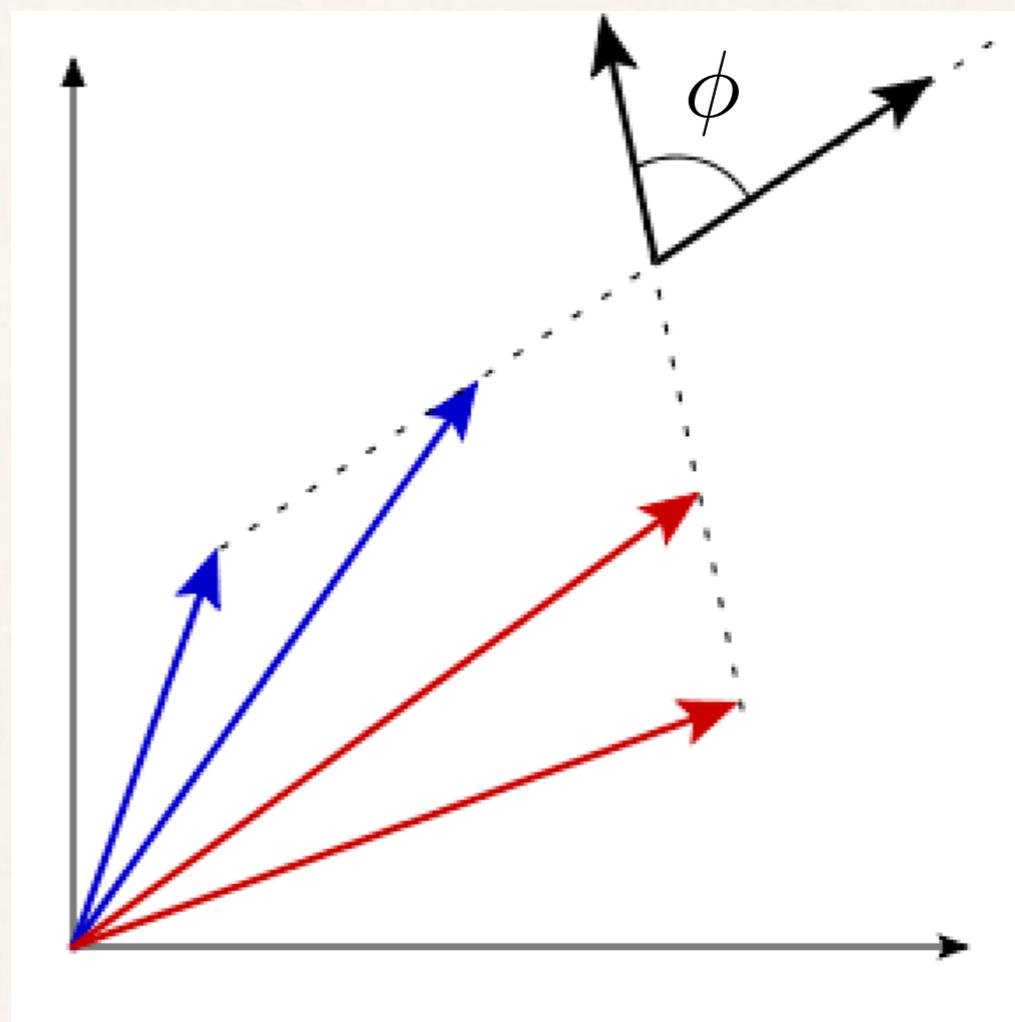
High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

$$\vec{a}_1 = \begin{pmatrix} x_1 \\ y_1 \end{pmatrix} \quad \vec{b}_1 = \begin{pmatrix} z_1 \\ w_1 \end{pmatrix}$$

$$\vec{a}_2 = \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \quad \vec{b}_2 = \begin{pmatrix} z_2 \\ w_2 \end{pmatrix}$$



$$\text{Vol}(2) = \begin{array}{|c|} \hline \text{---} \\ \hline \end{array} \begin{array}{|c|} \hline \text{---} \\ \hline \end{array} \begin{array}{|c|} \hline \text{---} \\ \hline \end{array}$$

High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

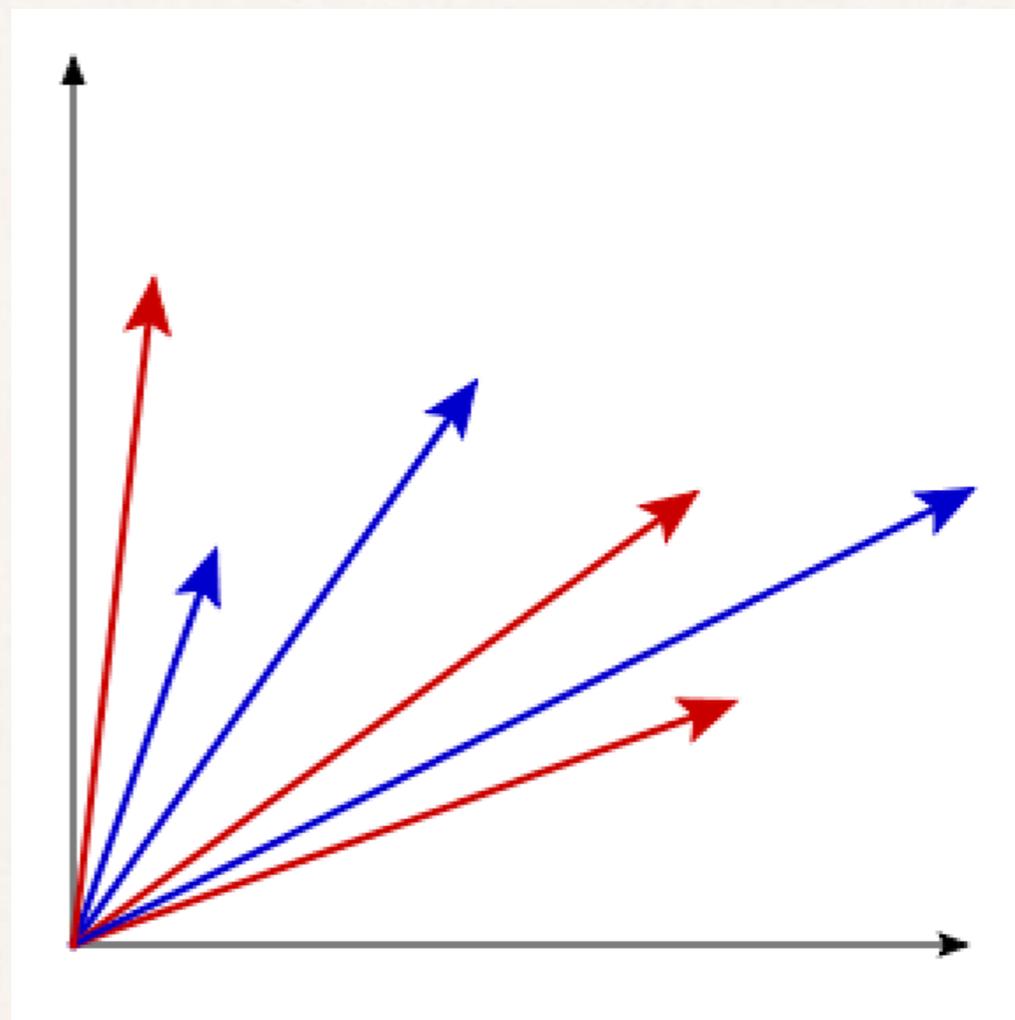
$$\vec{a}_1, \vec{a}_2, \vec{a}_3 \quad \vec{b}_1, \vec{b}_2, \vec{b}_3$$

❖ Conditions

$$(\vec{a}_1 - \vec{a}_2) \cdot (\vec{b}_1 - \vec{b}_2) \leq 0$$

$$(\vec{a}_1 - \vec{a}_3) \cdot (\vec{b}_1 - \vec{b}_3) \leq 0$$

$$(\vec{a}_2 - \vec{a}_3) \cdot (\vec{b}_2 - \vec{b}_3) \leq 0$$



$$\text{Vol}(3) = \begin{array}{ccc} \begin{array}{|c|} \hline \square \\ \hline \end{array} & \begin{array}{|c|} \hline \square \\ \hline \end{array} & \begin{array}{|c|} \hline \square \\ \hline \end{array} \\ \begin{array}{|c|} \hline \square \\ \hline \end{array} & \begin{array}{|c|} \hline \square \\ \hline \end{array} & \begin{array}{|c|} \hline \square \\ \hline \end{array} \end{array}$$

High school problem $gg \rightarrow gg$

❖ Positive quadrant

❖ Vectors

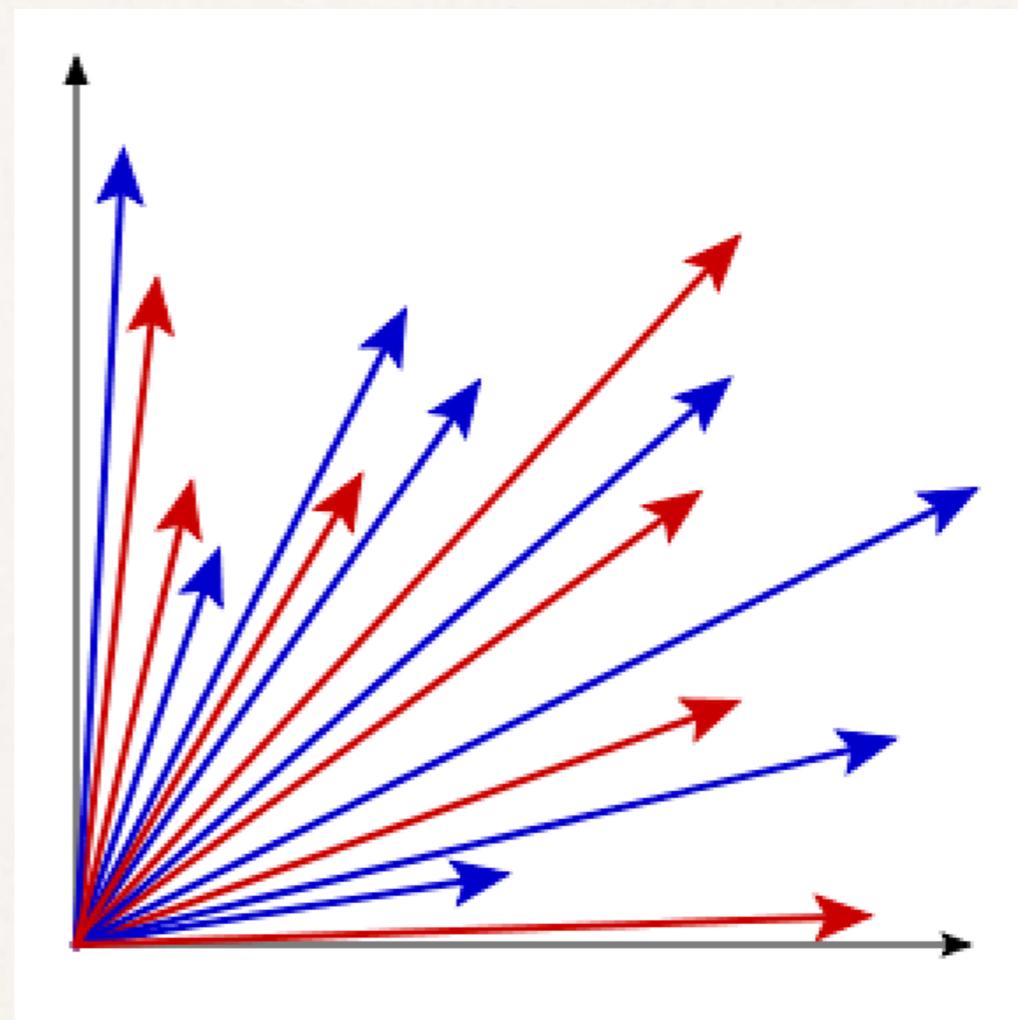
$$\vec{a}_1, \vec{a}_2, \dots, \vec{a}_\ell \quad \vec{b}_1, \vec{b}_2, \dots, \vec{b}_\ell$$

❖ Conditions

$$(\vec{a}_i - \vec{a}_j) \cdot (\vec{b}_i - \vec{b}_j) \leq 0$$

for all pairs i, j

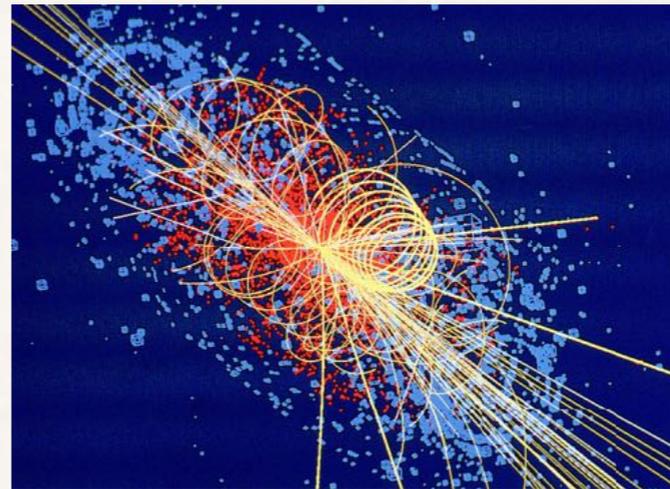
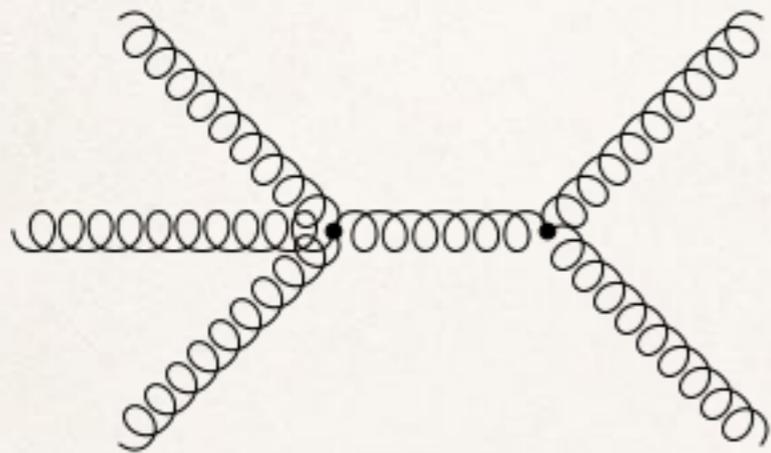
Let me know if you solve it!



$$\text{Vol}(\ell) = \dots\dots\dots$$

Physics vs geometry

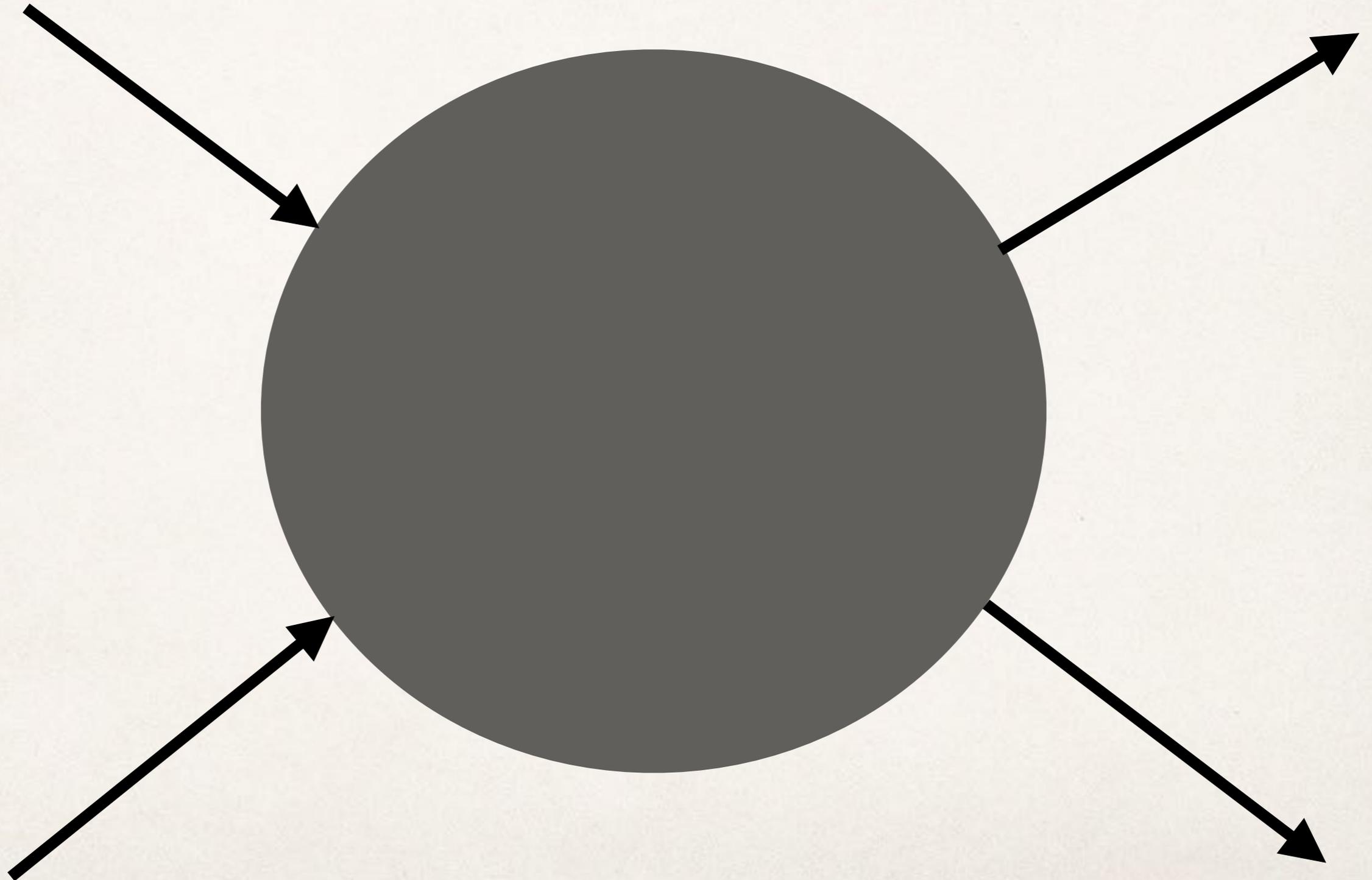
- ❖ Dynamical particle interactions in 4-dimensions



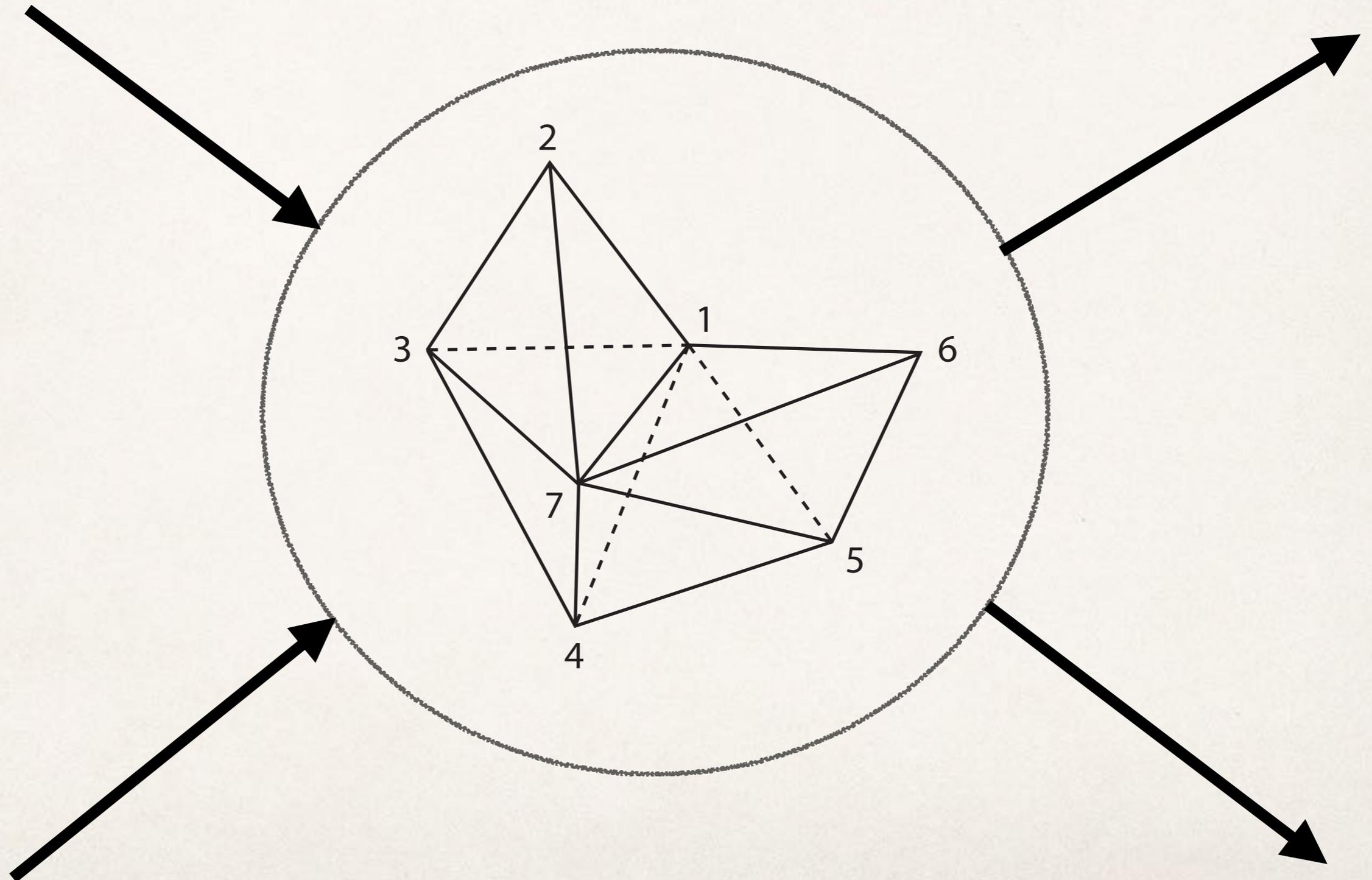
- ❖ Static geometry in high dimensional space



What is scattering amplitude?



What is scattering amplitude?

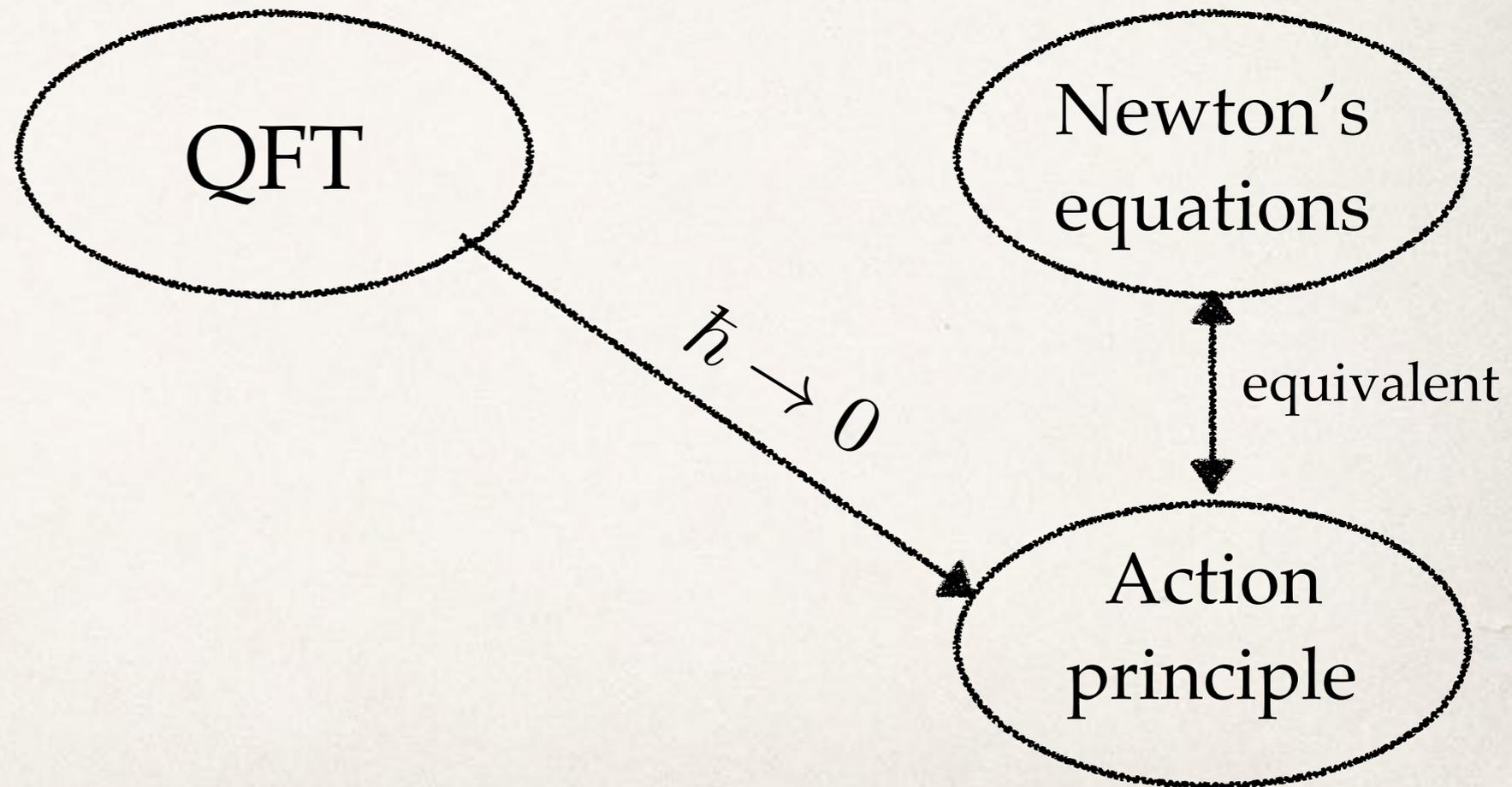


Step 1.1.1.

- ❖ It is very early to say if/how this can generalize
- ❖ Some encouraging news but more work needed
- ❖ Extend to other theories, beyond the integrand,....

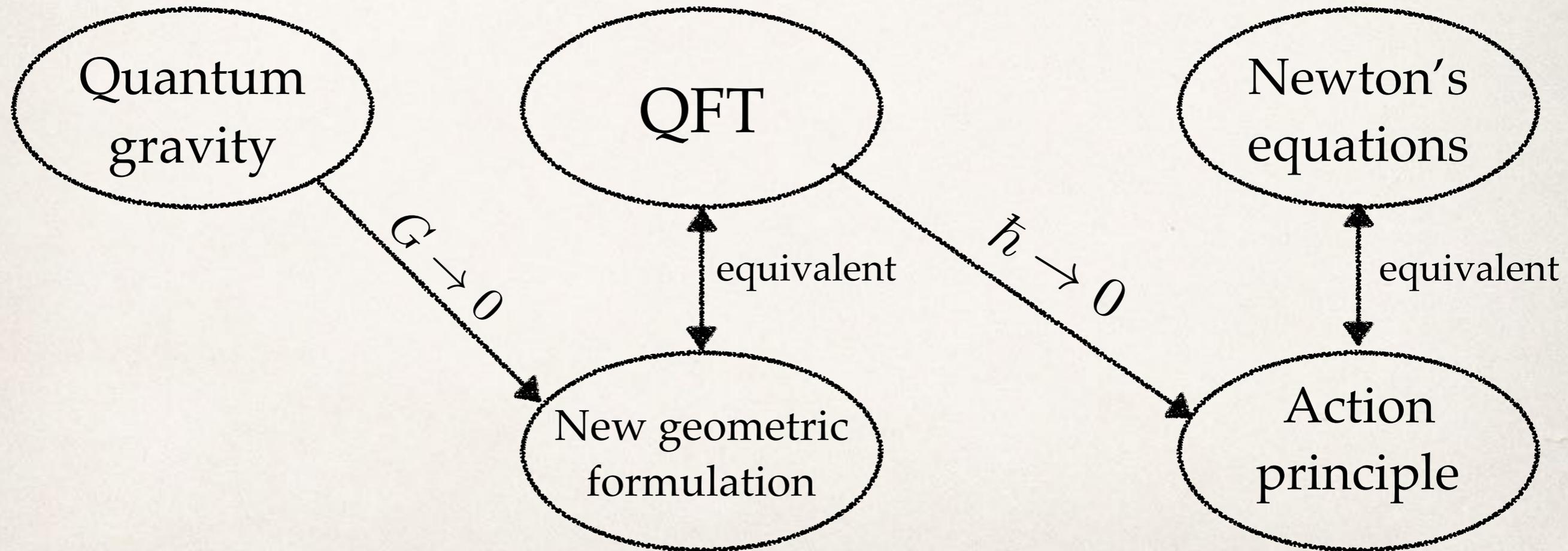
Fantasy

- ❖ Beyond understanding QFT better there is one more motivation



Fantasy

- ❖ Beyond understanding QFT better there is one more motivation

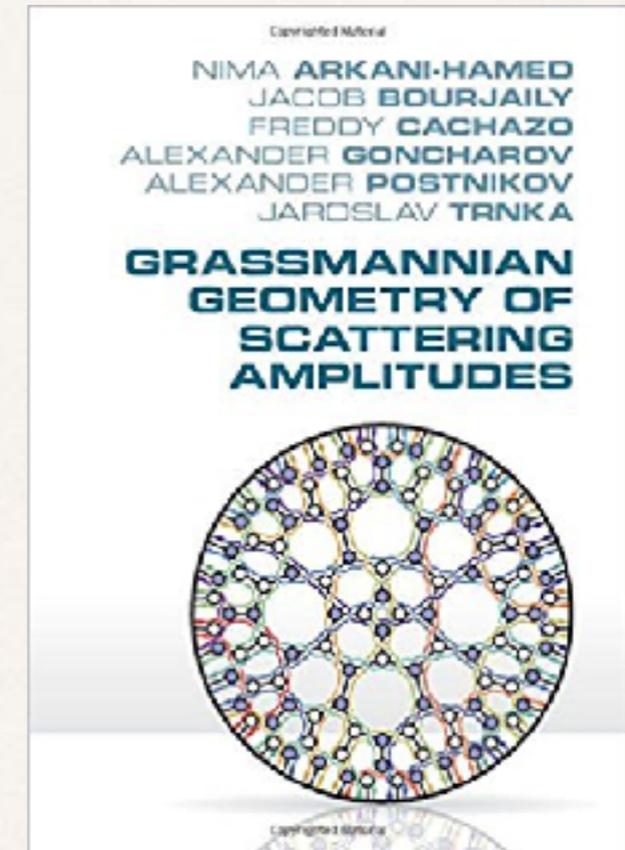
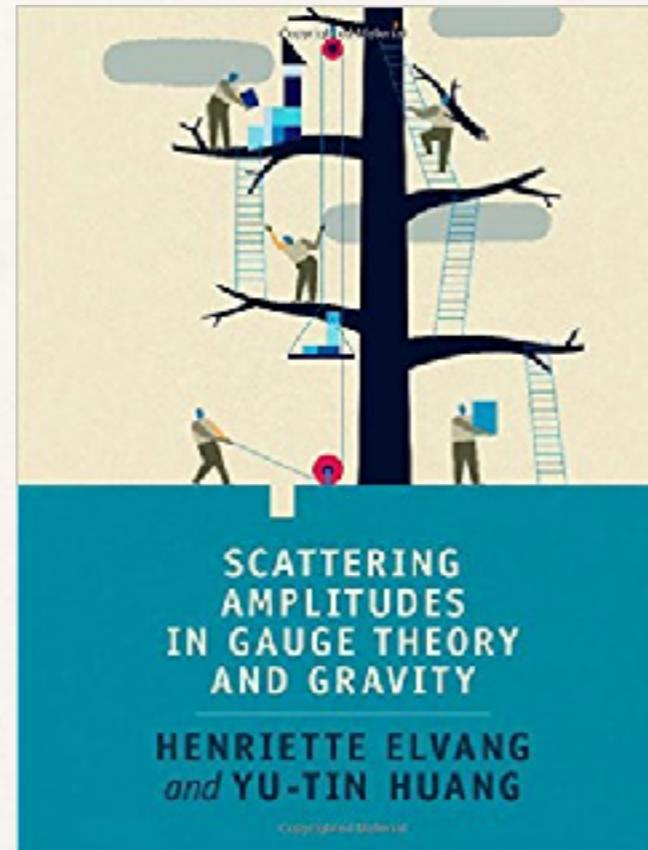
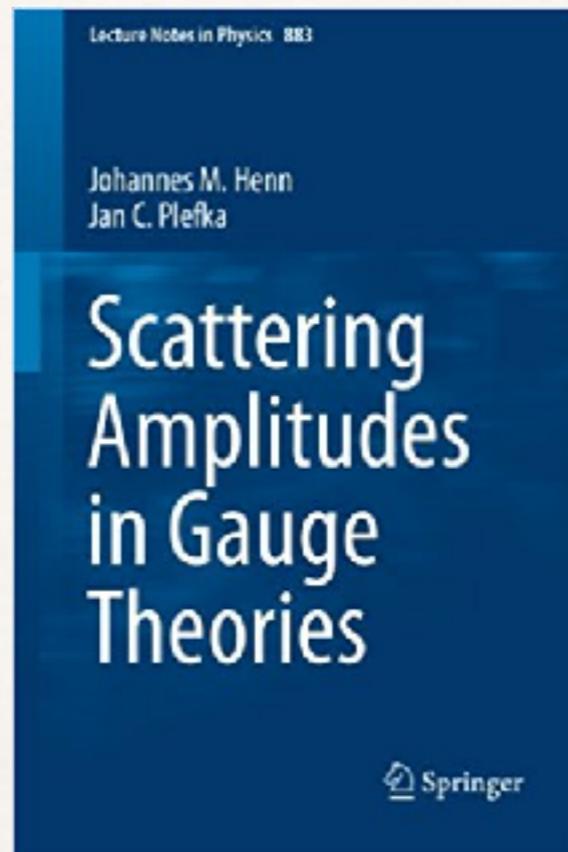


Amplitudes as a new field

- ❖ This is one of the directions in fast developing field
- ❖ Scattering equations, BCJ duality, string amplitudes, supergravity amplitudes, hexagon bootstrap, cluster polylogarithms, worldsheet models, integration techniques, soft theorems, LHC calculations,.....
- ❖ Zeroth order problems open, many chances for young people to make big discoveries!

Resources for amplitudes

❖ Books



❖ Many conferences, workshops, summer schools,....

Videos on youtube

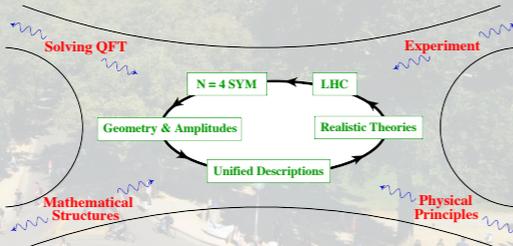
❖ In June 2018: QMAP Amplitudes Summer School at UC Davis

CENTER FOR QUANTUM MATHEMATICS AND PHYSICS
QMAP
UNIVERSITY OF CALIFORNIA
DAVIS
JUNE 11-15, 2018



AMPLITUDES SUMMER SCHOOL

LECTURERS INCLUDE:
NIMA ARKANI-HAMED
ZVI BERN
JACOB BOURJAILY
CLAUDE DUHR
SONG HE
ERIC D'HOKER
YUTIN HUANG
ALEXANDER POSTNIKOV



Followed by Amplitudes 2018 at SLAC, June 18-22

Local organizers:
Lance Dixon, Enrico Herrmann, Jaroslav Trnka, Andrew Waldron

QMAP.UCDAVIS.EDU/EVENTS/AMPLITUDES-SUMMER-SCHOOL

QUESTIONS: TRNKA@UCDAVIS.EDU MATHEMATICAL SCIENCES BUILDING, ROOM 1147 **UCDAVIS** UNIVERSITY OF CALIFORNIA

Lectures online:
35 hours on various
topics in amplitudes
including 11-hour
marathon lecture by
Nima Arkani-Hamed



Thank you!