

Edge state dynamics of a bosonic fractional Chern insulator

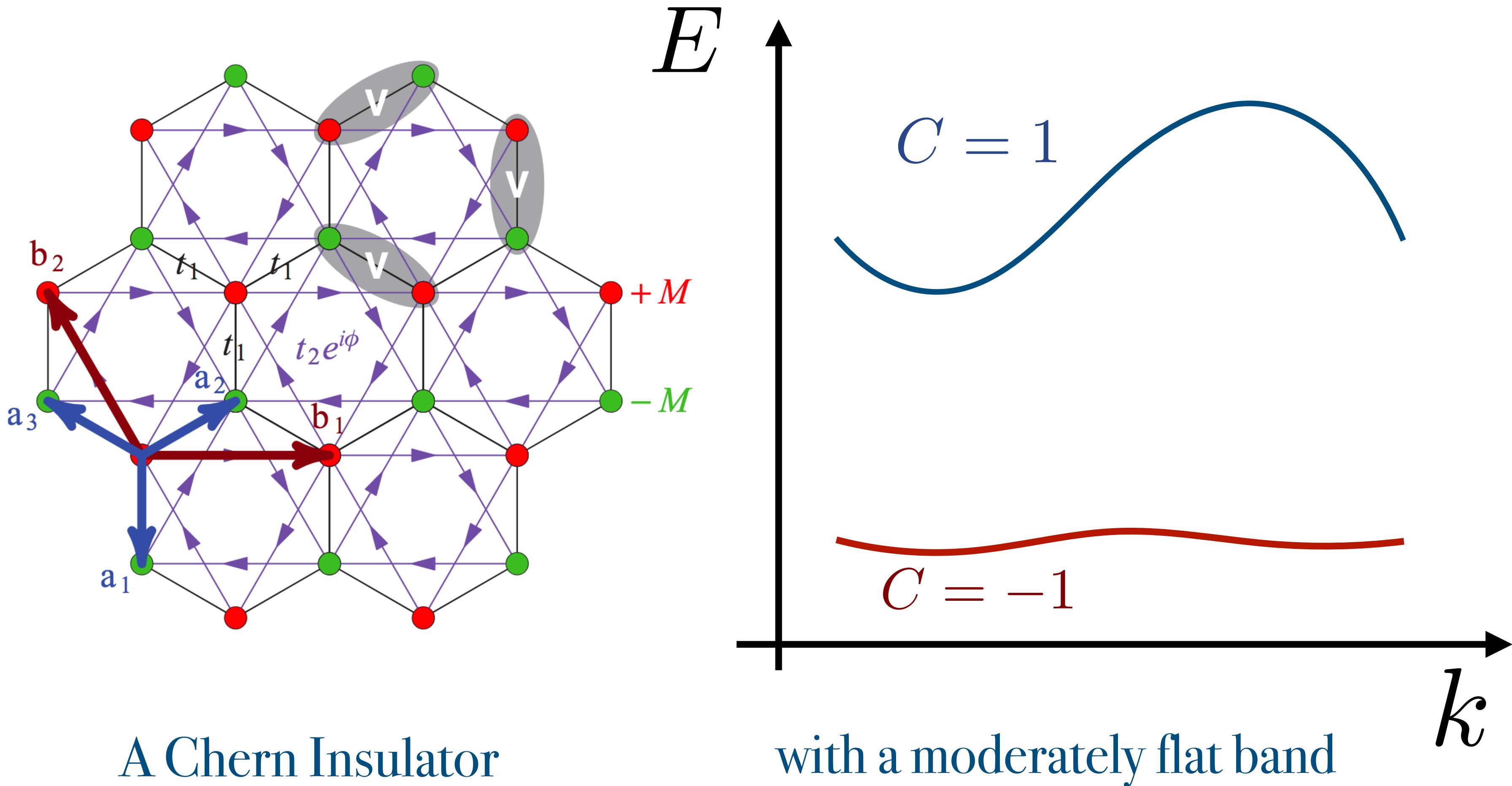
Adolfo G. Grushin, Néel Institute, CNRS

Benasque, 26/2/19



X. Y. Dong, AGG, J. Motruk, F.
Pollmann, Phys Rev. Lett. (2018)

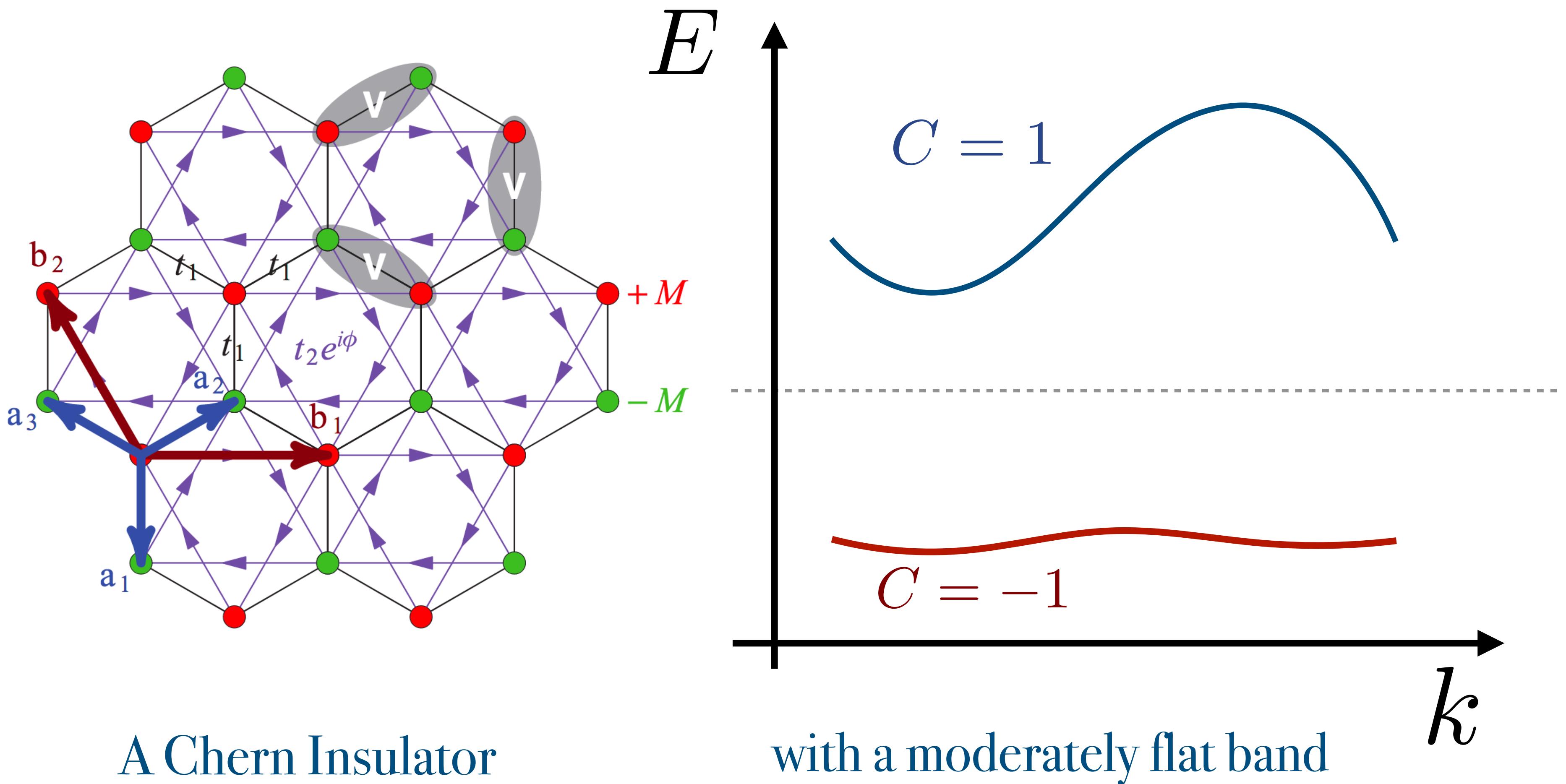
The promise of a fractional quantum Hall effect without Landau levels



Neupert et al. Phys. Scr. T164, 014005, (2015)

Emil J. Bergholtz, Zhao Liu Int. J. Mod. Phys. B 27, 1330017 (2013)

The promise of a fractional quantum Hall effect without Landau levels



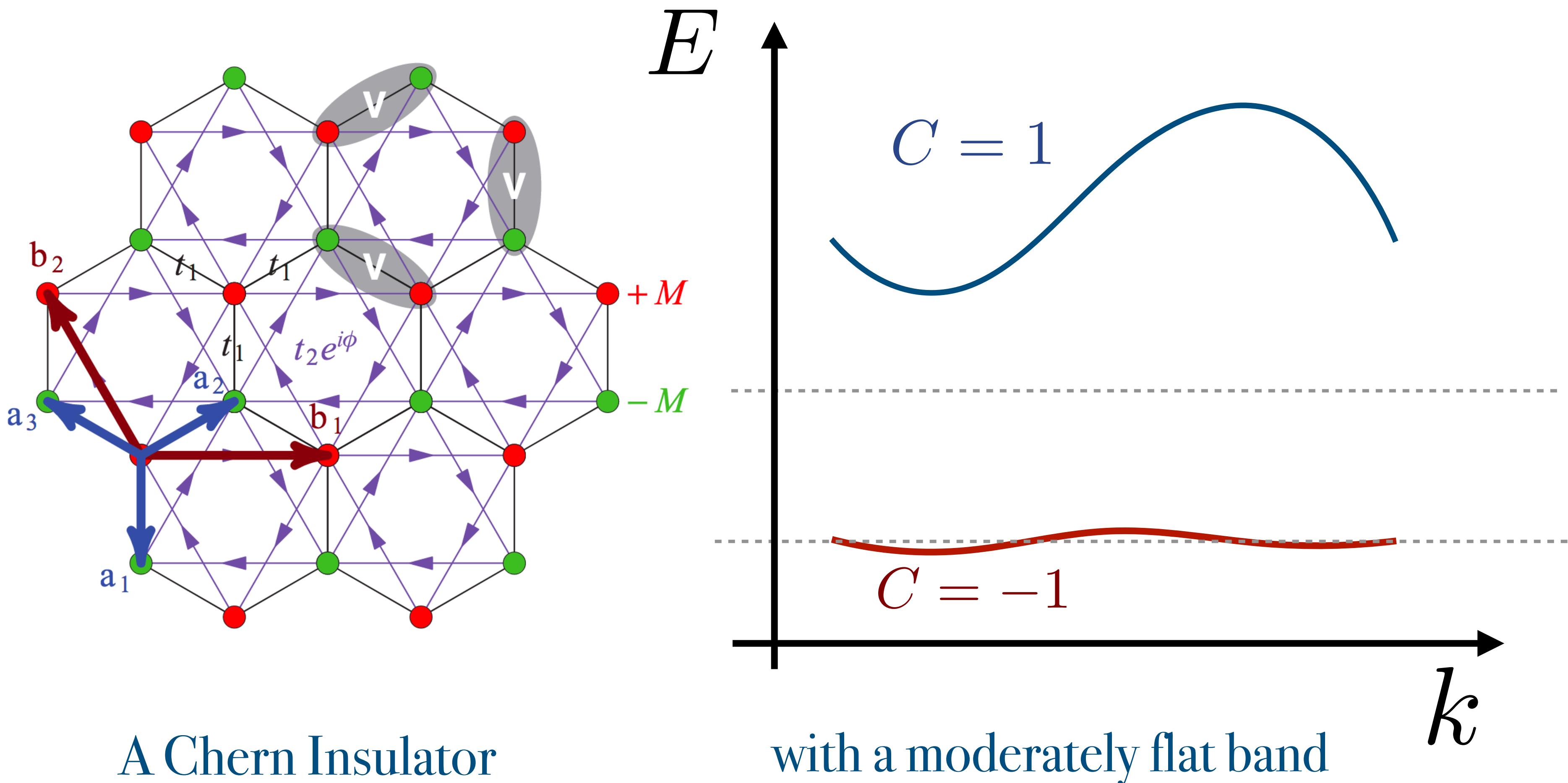
A Chern Insulator

$$\sigma_{xy} = C \frac{e^2}{h}$$

Neupert et al. Phys. Scr. T164, 014005, (2015)

Emil J. Bergholtz, Zhao Liu Int. J. Mod. Phys. B 27, 1330017 (2013)

The promise of a fractional quantum Hall effect without Landau levels



$$\sigma_{xy} = C \frac{e^2}{h}$$

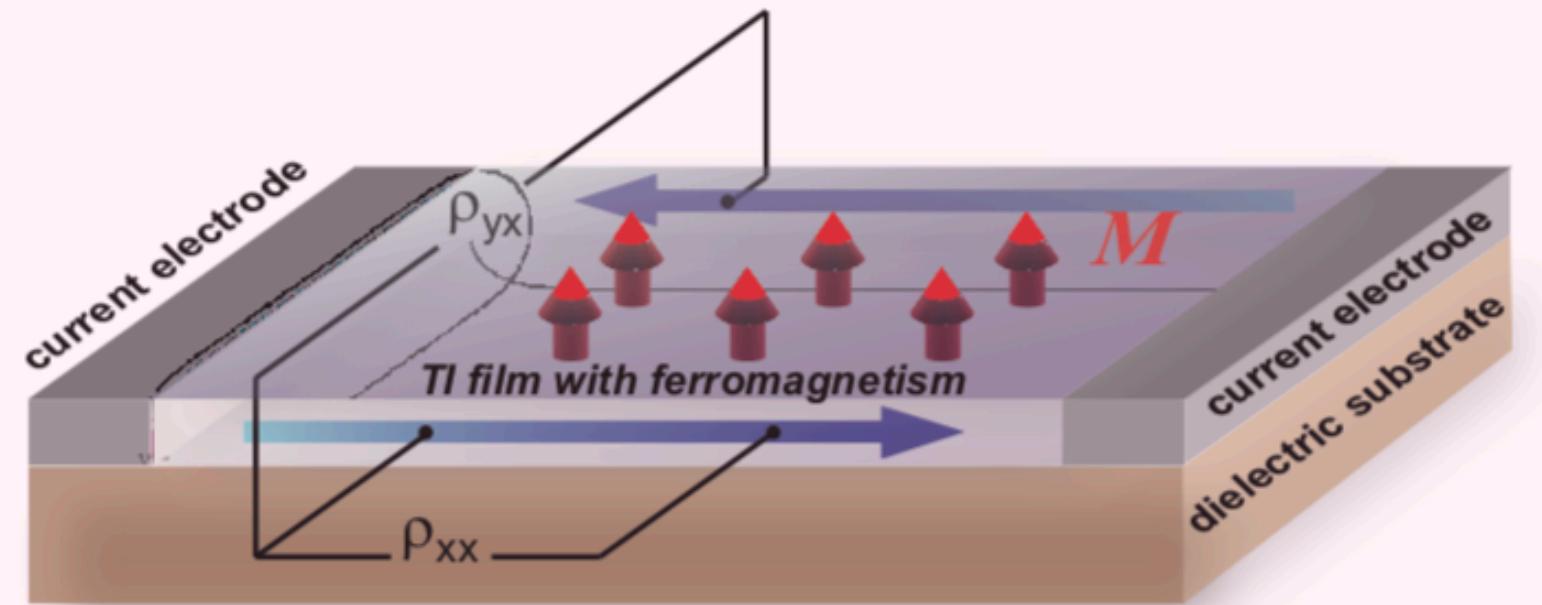
$$\sigma_{xy} = \frac{1}{m} \frac{e^2}{h}$$

at $1/m$ filling

Neupert et al. Phys. Scr. T164, 014005, (2015)

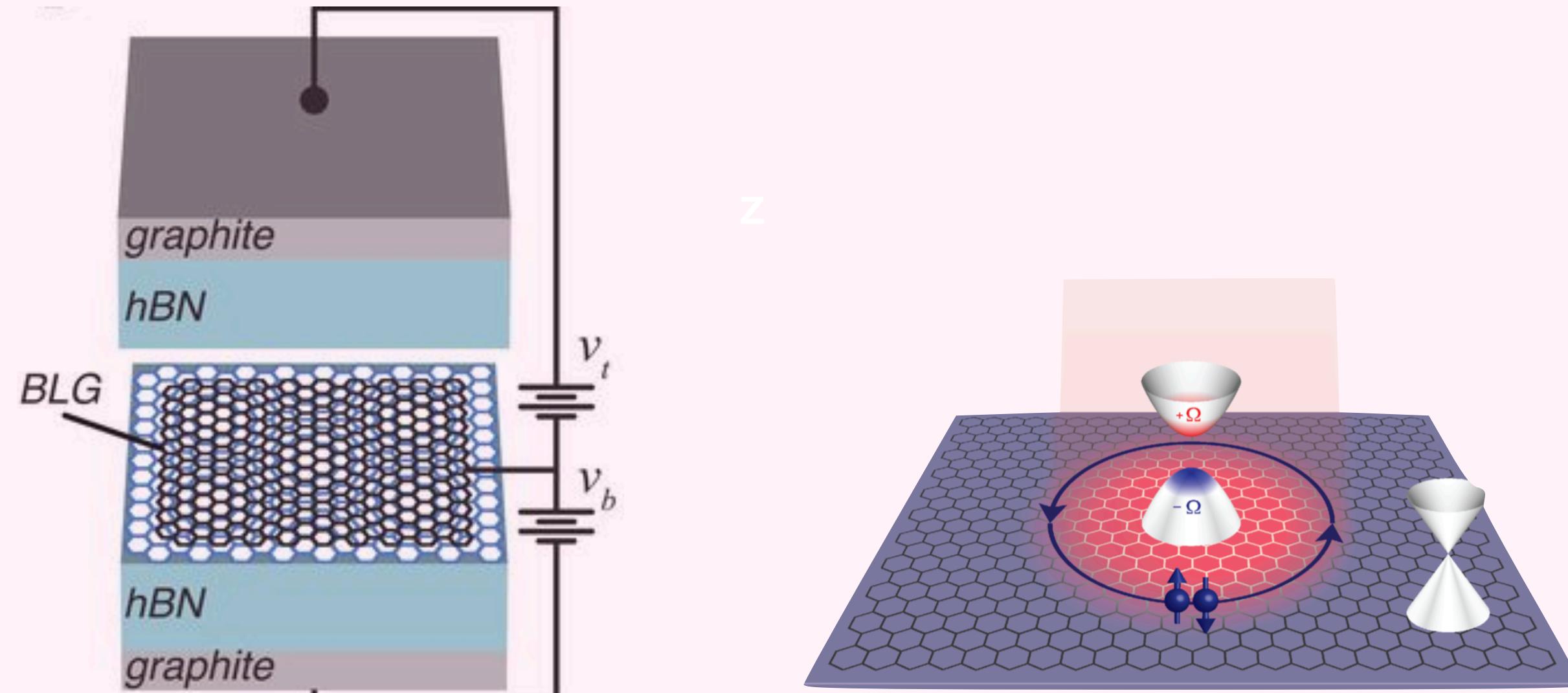
Emil J. Bergholtz, Zhao Liu Int. J. Mod. Phys. B 27, 1330017 (2013)

Chern insulators in the solid state



Doped topological insulators

C. Z. Chang Science (2013)



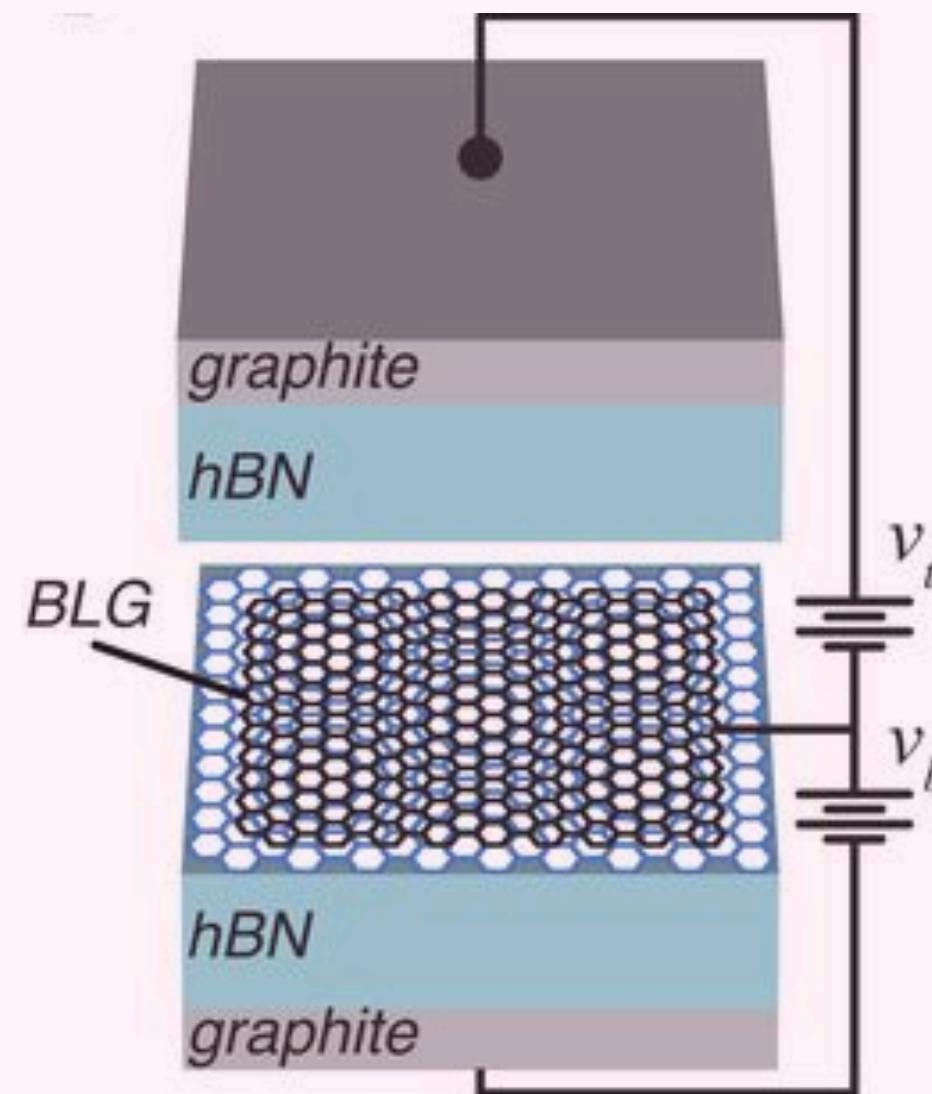
Moire Graphene

E. M. Spanton et. al Science (2018)

Irradiated Graphene

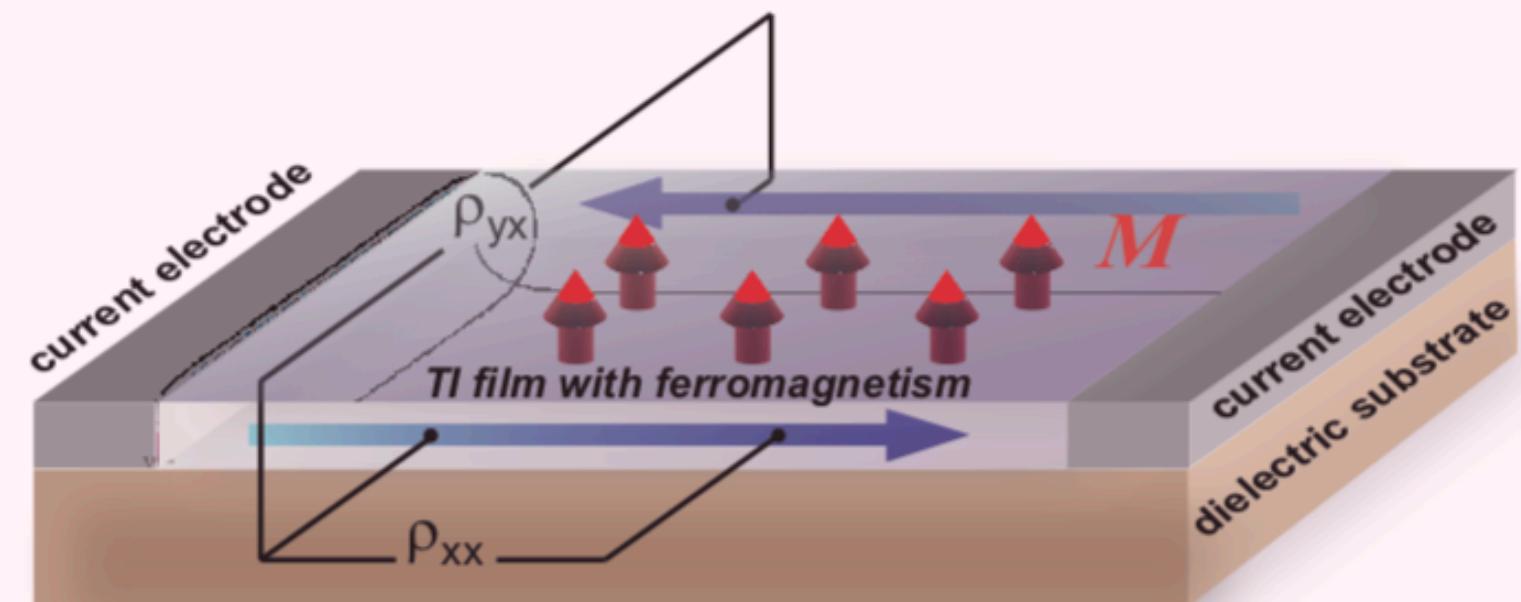
J.W. McIver et al. 1811.03522

Chern insulators in the solid state



Moire Graphene

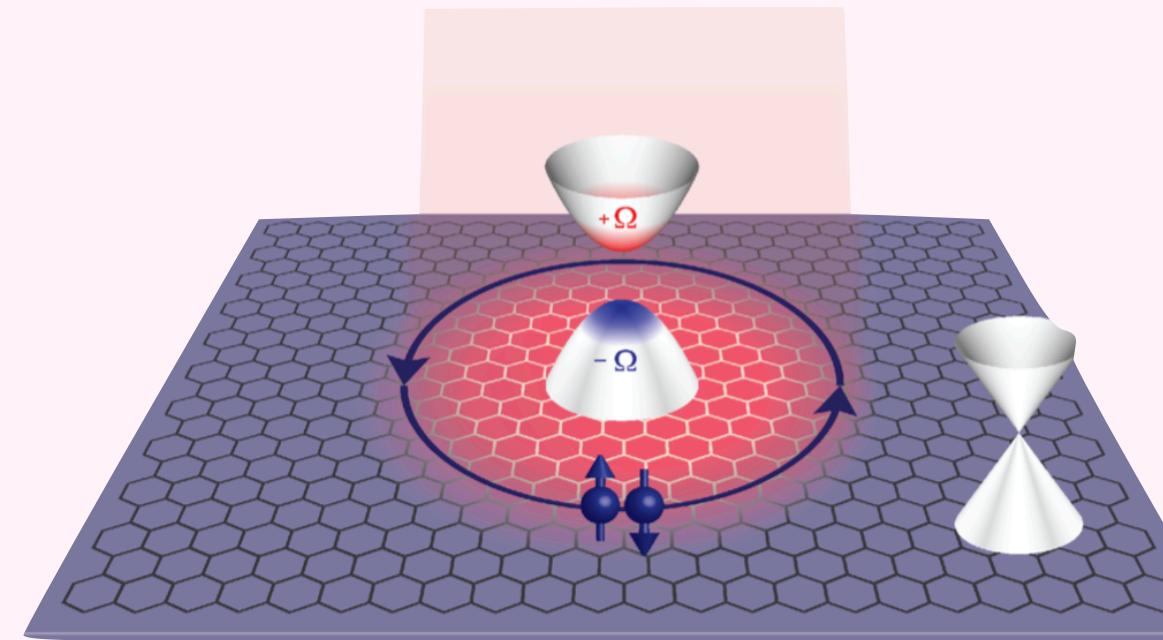
E. M. Spanton et. al Science (2018)



Doped topological insulators

C. Z. Chang Science (2013)

Floquet

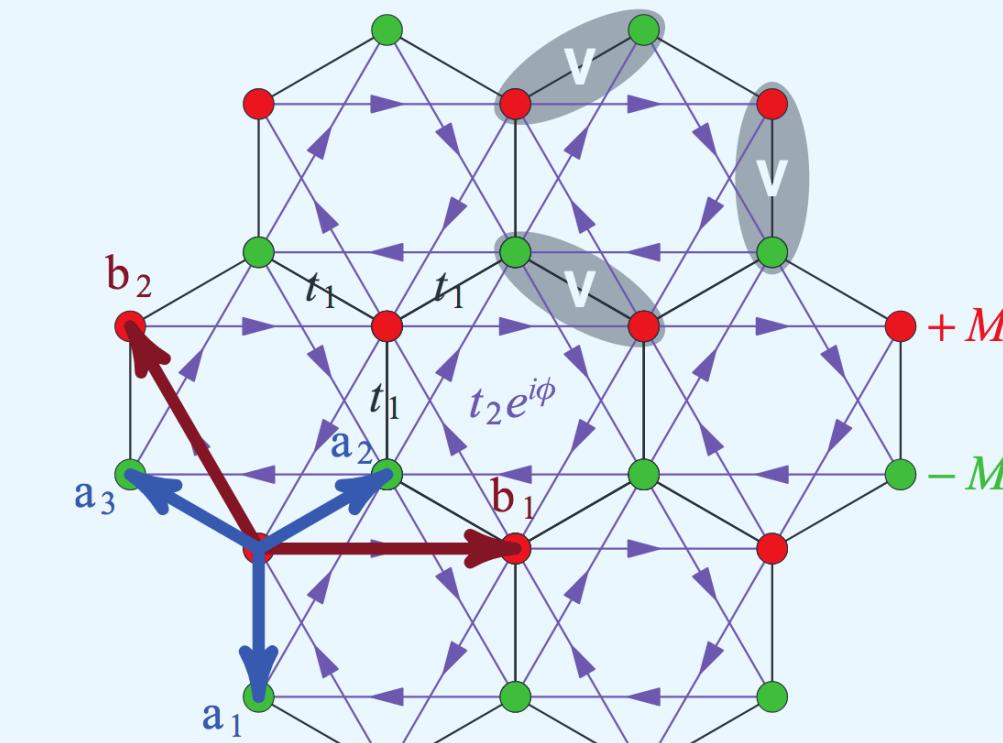


Irradiated Graphene

J.W. McIver et al. 1811.03522

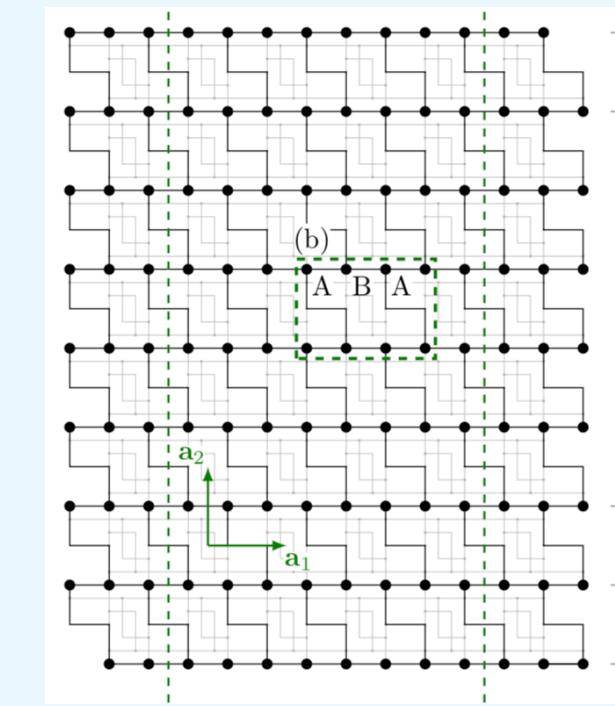
Chern insulators in synthetic matter

Ultra-cold fermions



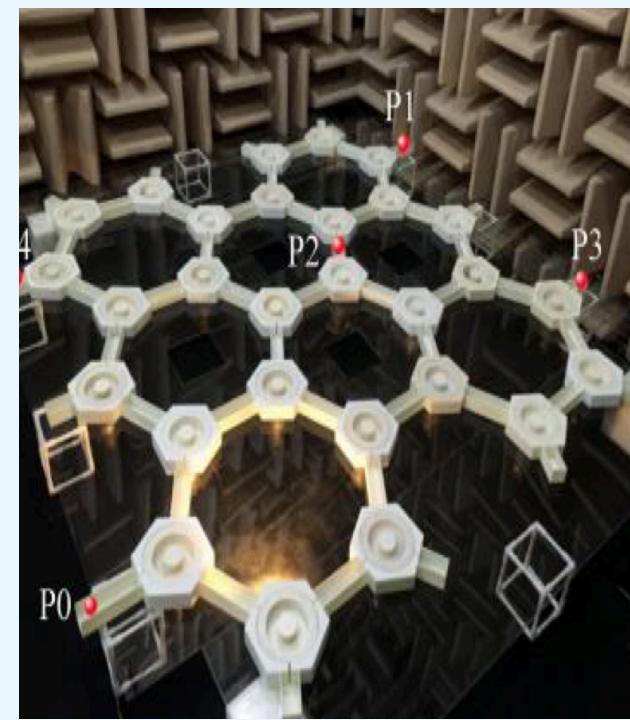
G. Jotzu et. al. Nature 2014

Circuit



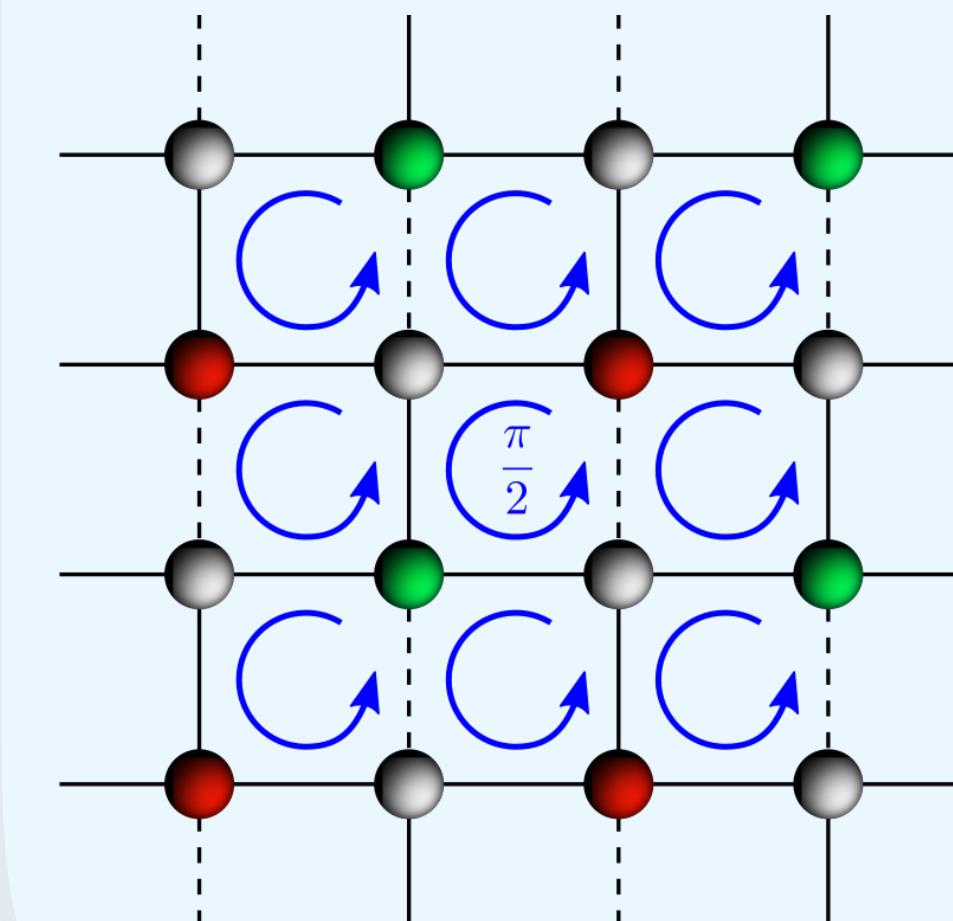
T. Hofmann et al.
1801.07942

Acoustic



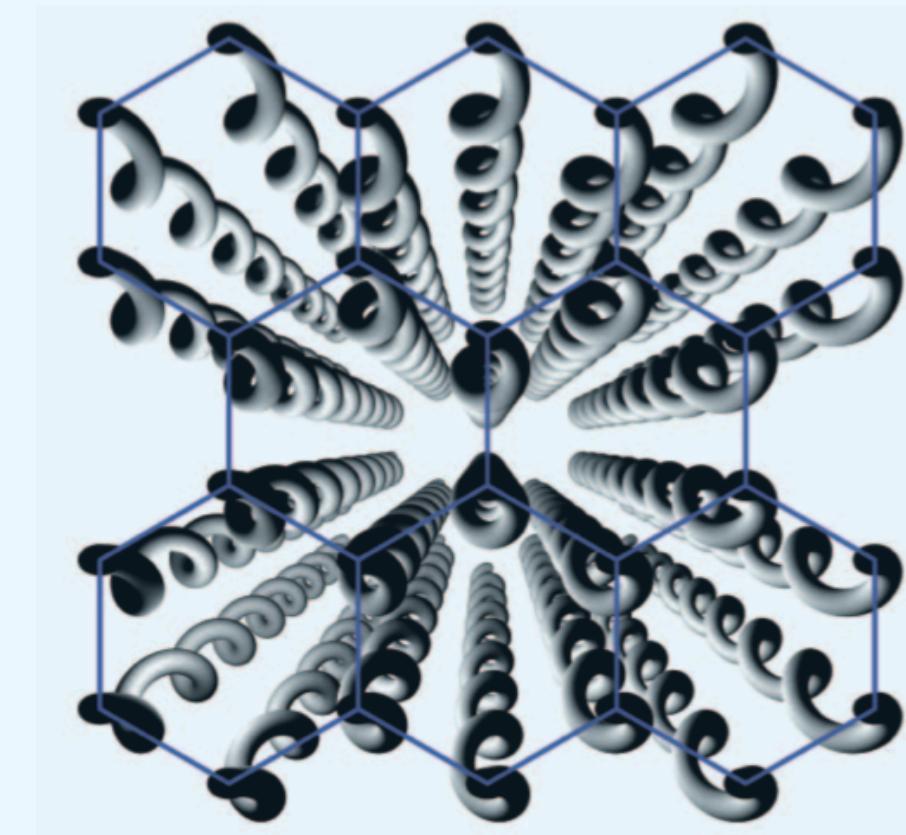
Y. Zhu et al.
1801.07942

Ultra-cold bosons



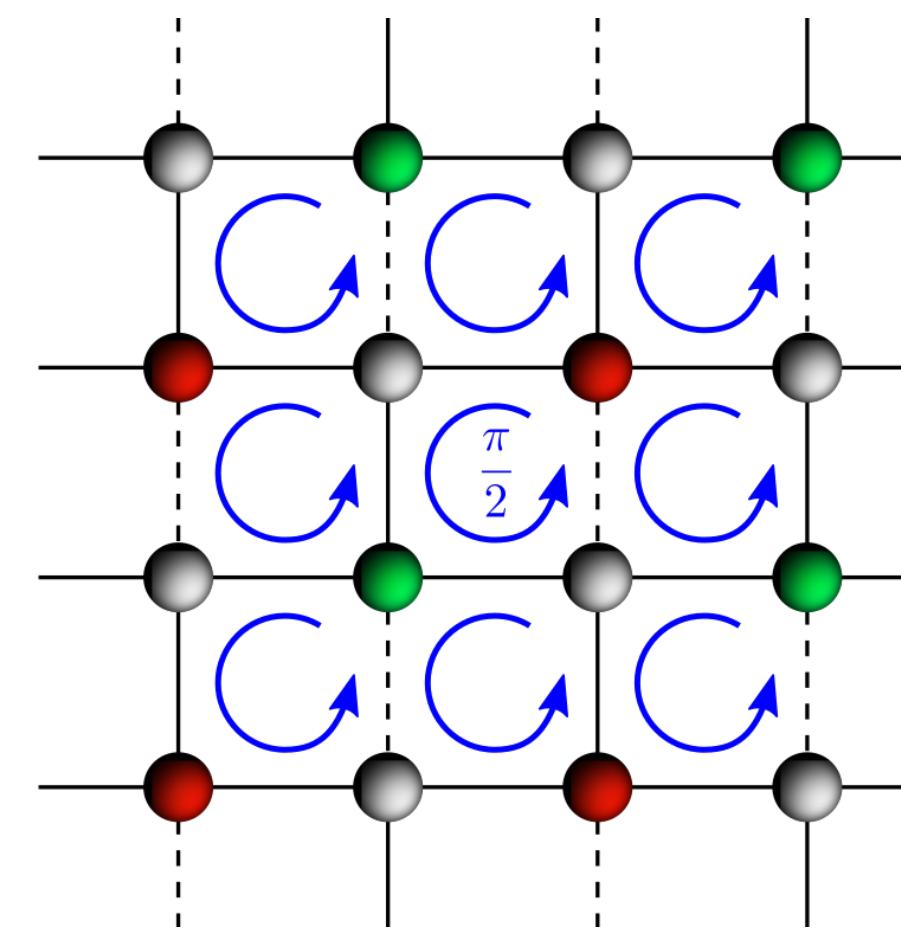
Aidelsburger et. al Nat. Phys (2013)

Photonic

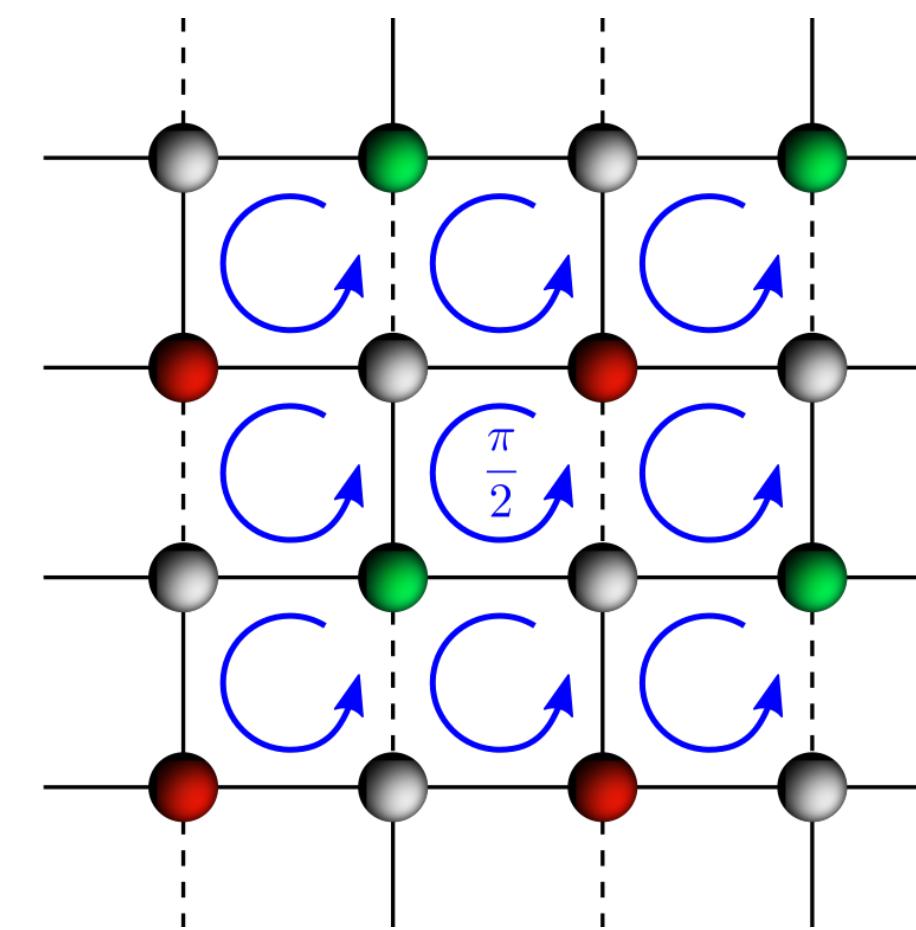


M. Rechtsman, Nature (2013)

Ultra-cold bosons

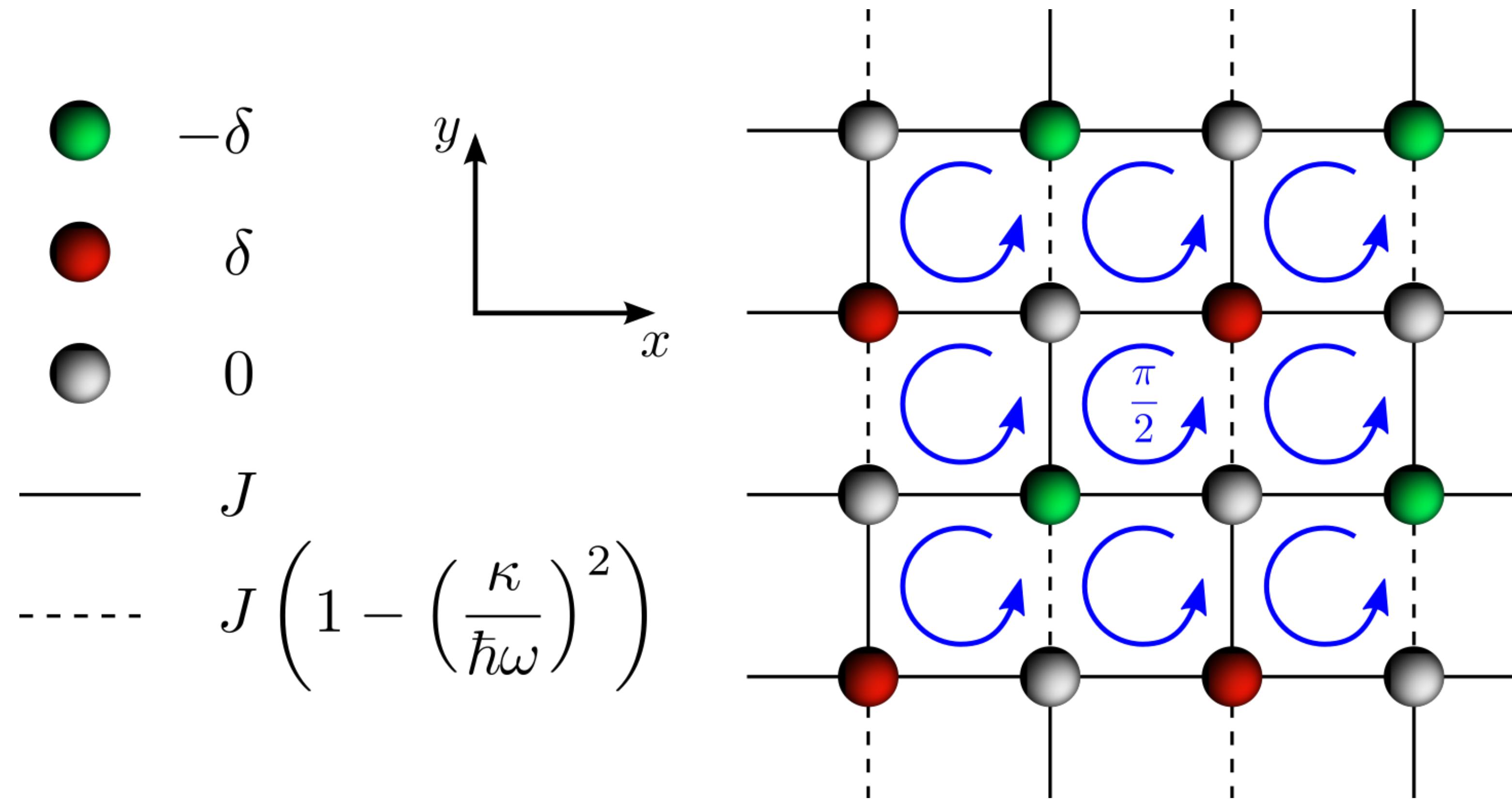


Ultra-cold bosons



Aidelsburger et. al Nat. Phys (2013)

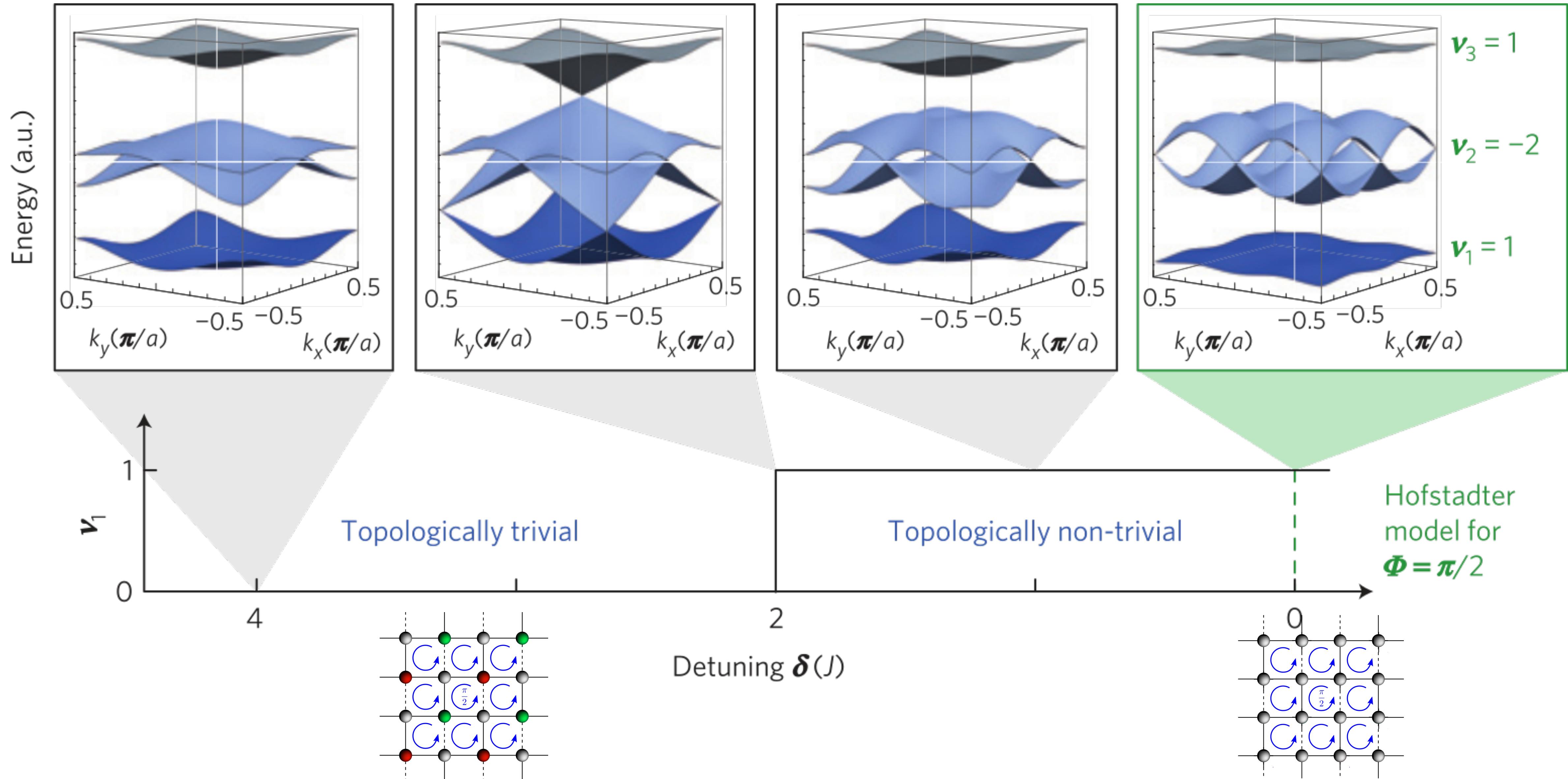
(Floquet) Chern insulator with ultra cold bosons



$$H_{\text{eff}} = -J \sum_{m,n} \left\{ \hat{a}_{m+1,n}^\dagger \hat{a}_{m,n} e^{i[\pi/2(m+n) - \phi_0]} + (1 + f_{m,n}) \hat{a}_{m,n+1}^\dagger \hat{a}_{m,n} + \text{h.c.} \right\} + \frac{\delta}{2} \sum_{m,n} [(-1)^m + (-1)^n] \hat{n}_{m,n}$$

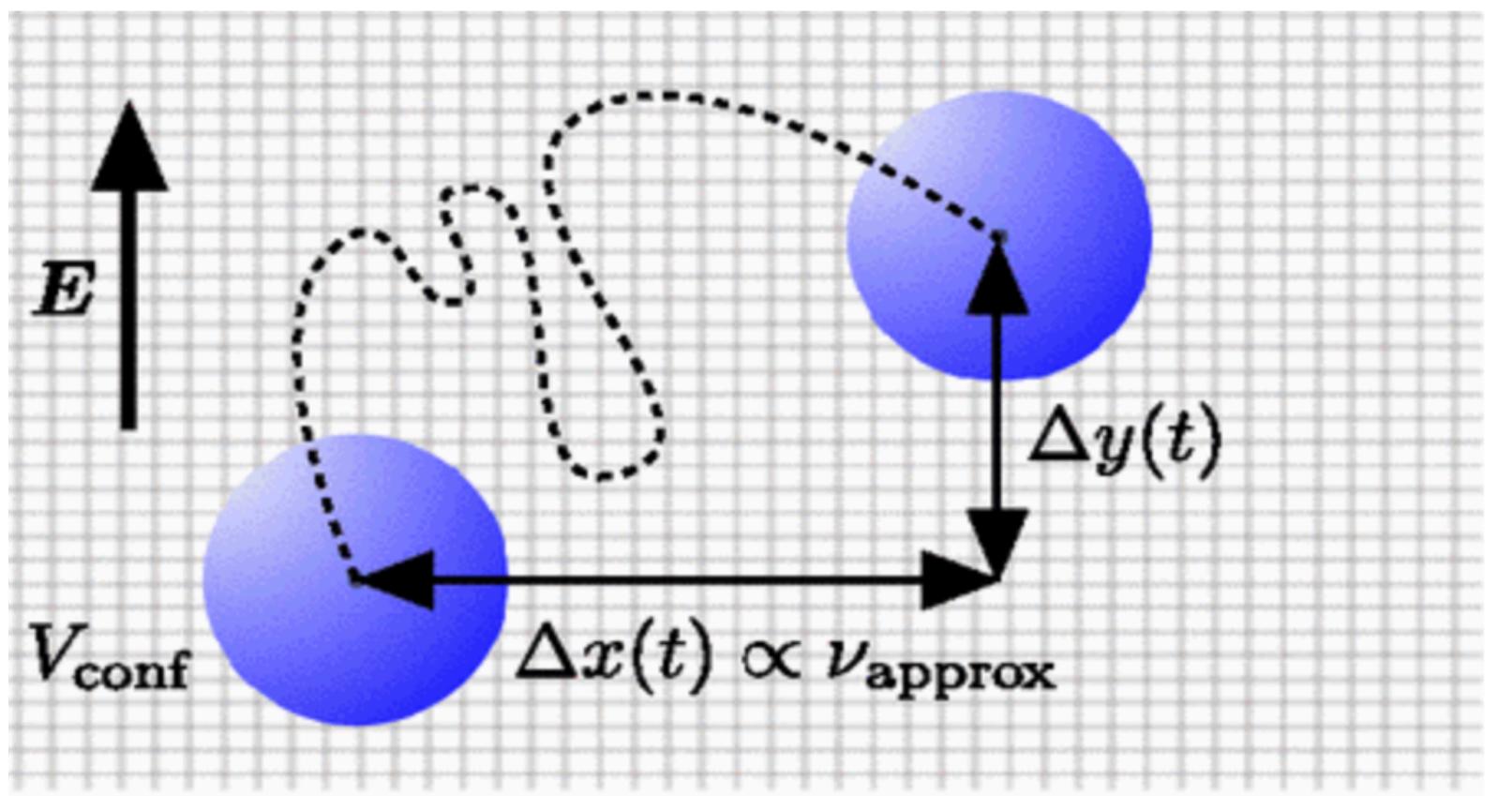
(Floquet) Chern insulator with ultra cold bosons

M. Aidelsburger et al. Nat. Phys (2015)



How does one know it is a Chern insulator?

Wave packet motion

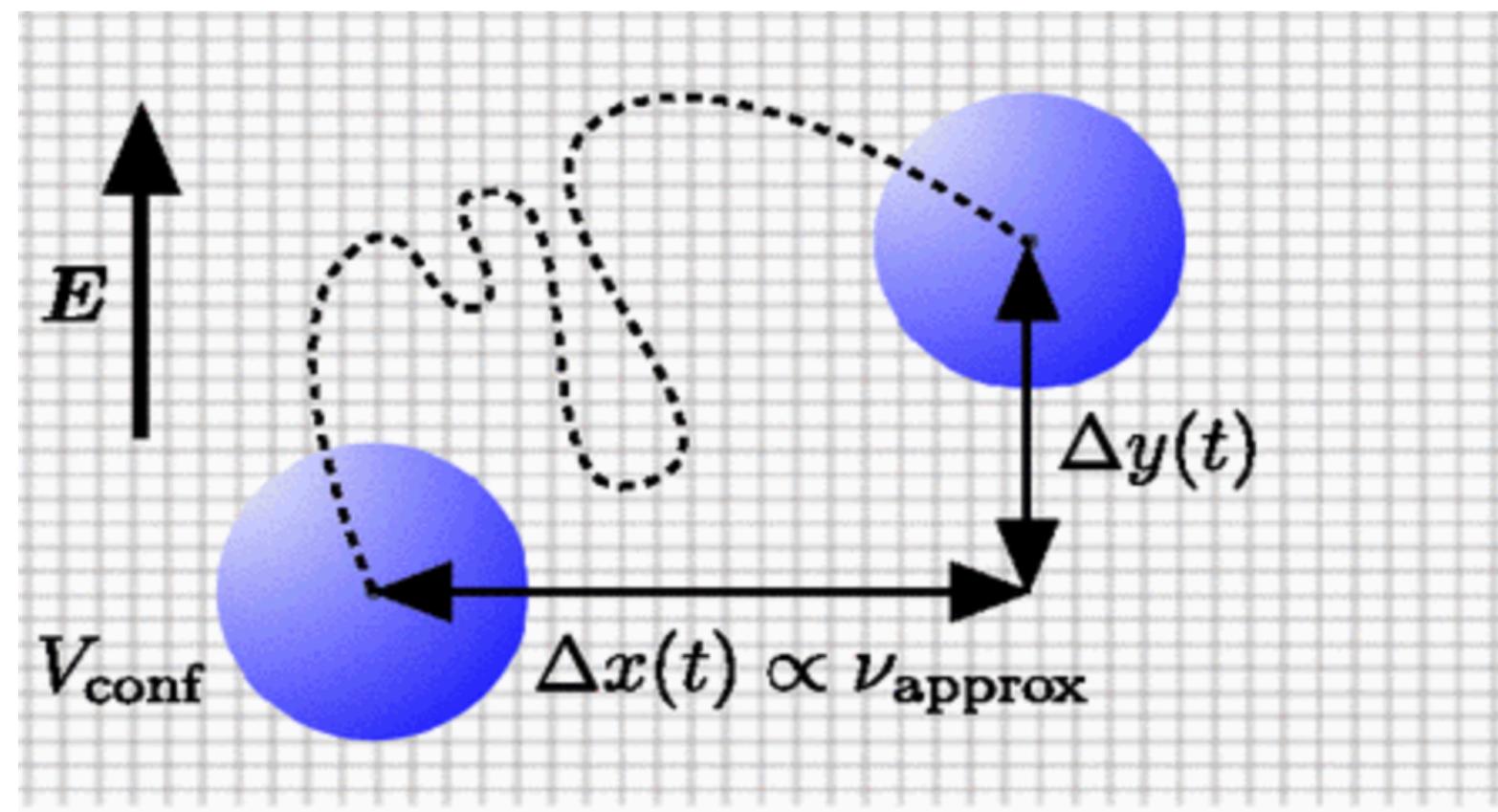


A. Dauphin and N. Goldman PRL 2013

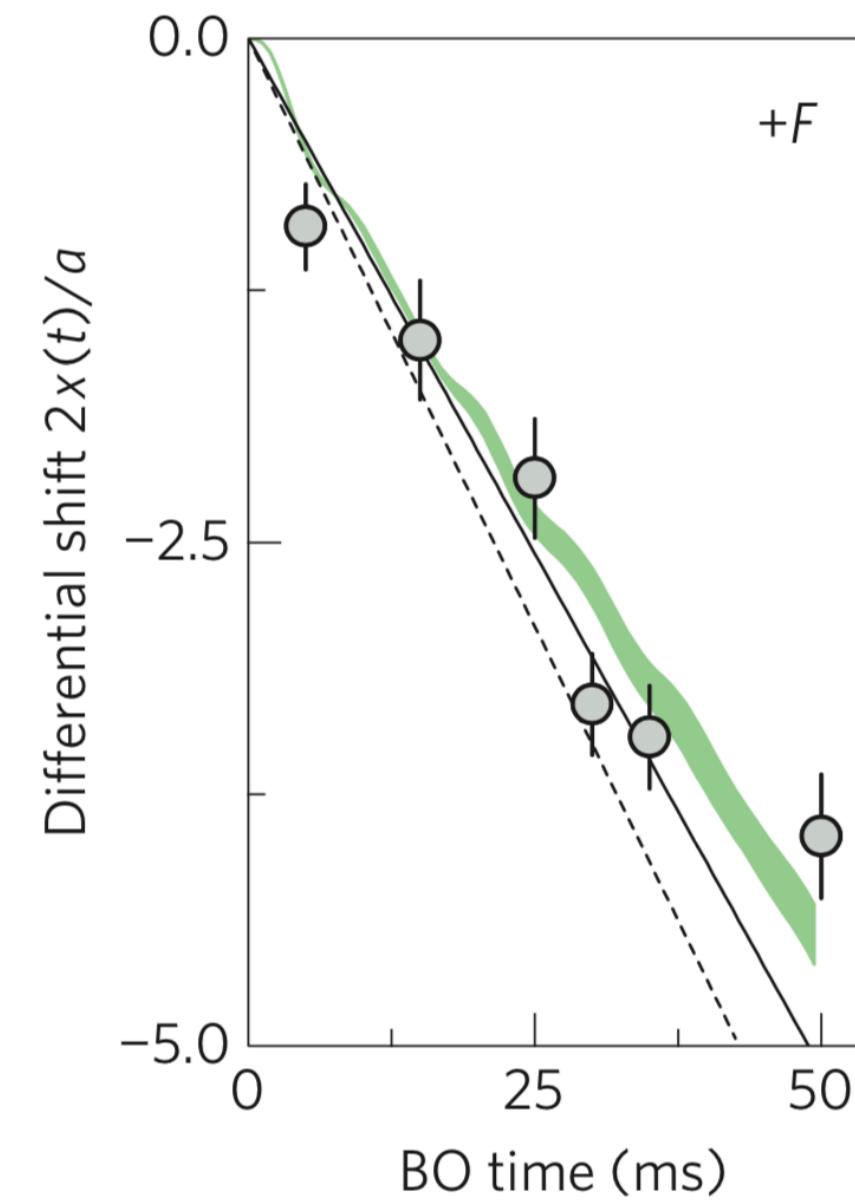
$$x(t) = -(a^2 t E_y / \pi \hbar) \nu_{\text{approx}}$$

How does one know it is a Chern insulator?

Wave packet motion



A. Dauphin and N. Goldman PRL 2013

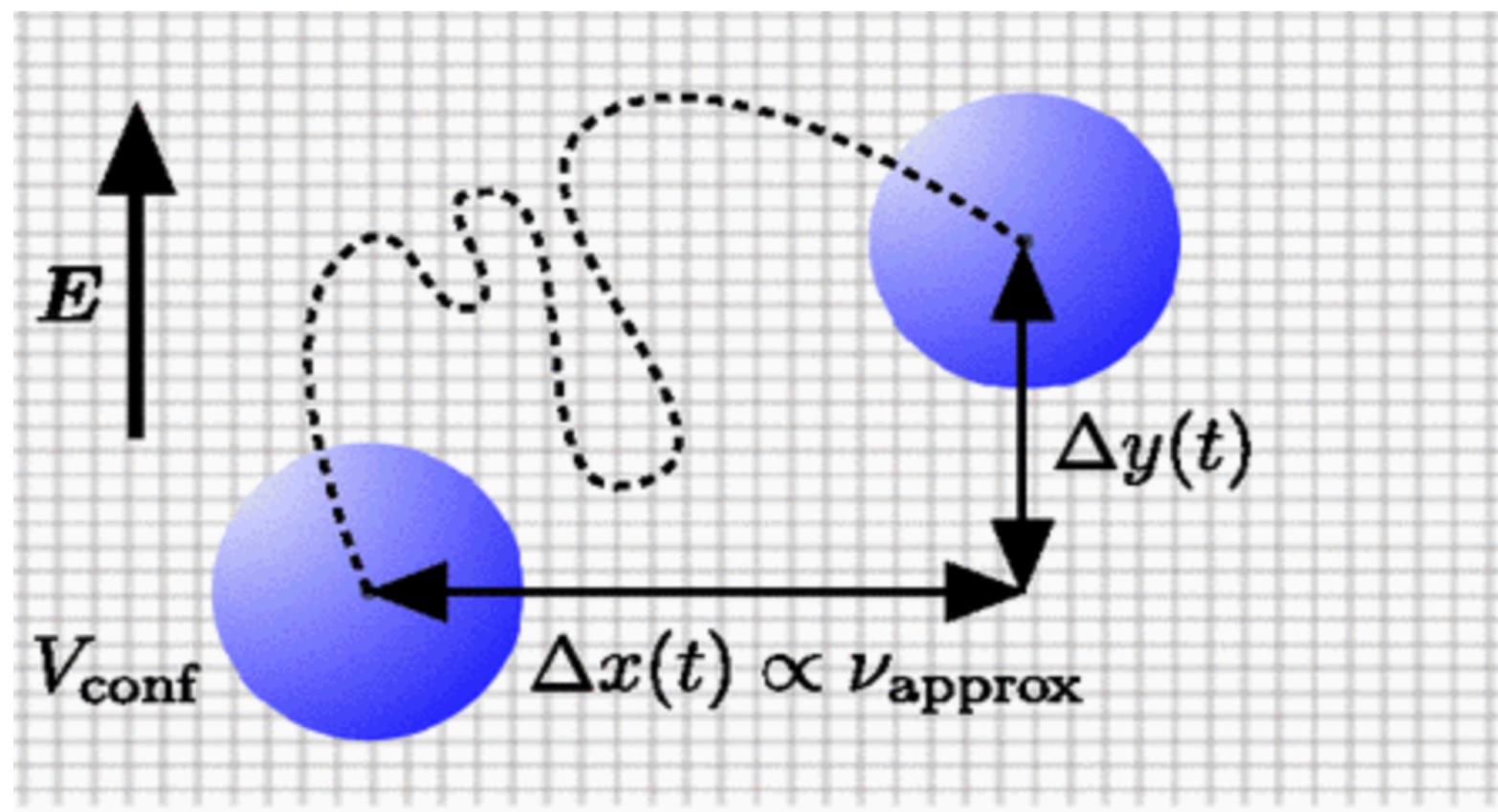


M. Aidelsburger et al. Nat. Phys (2015)

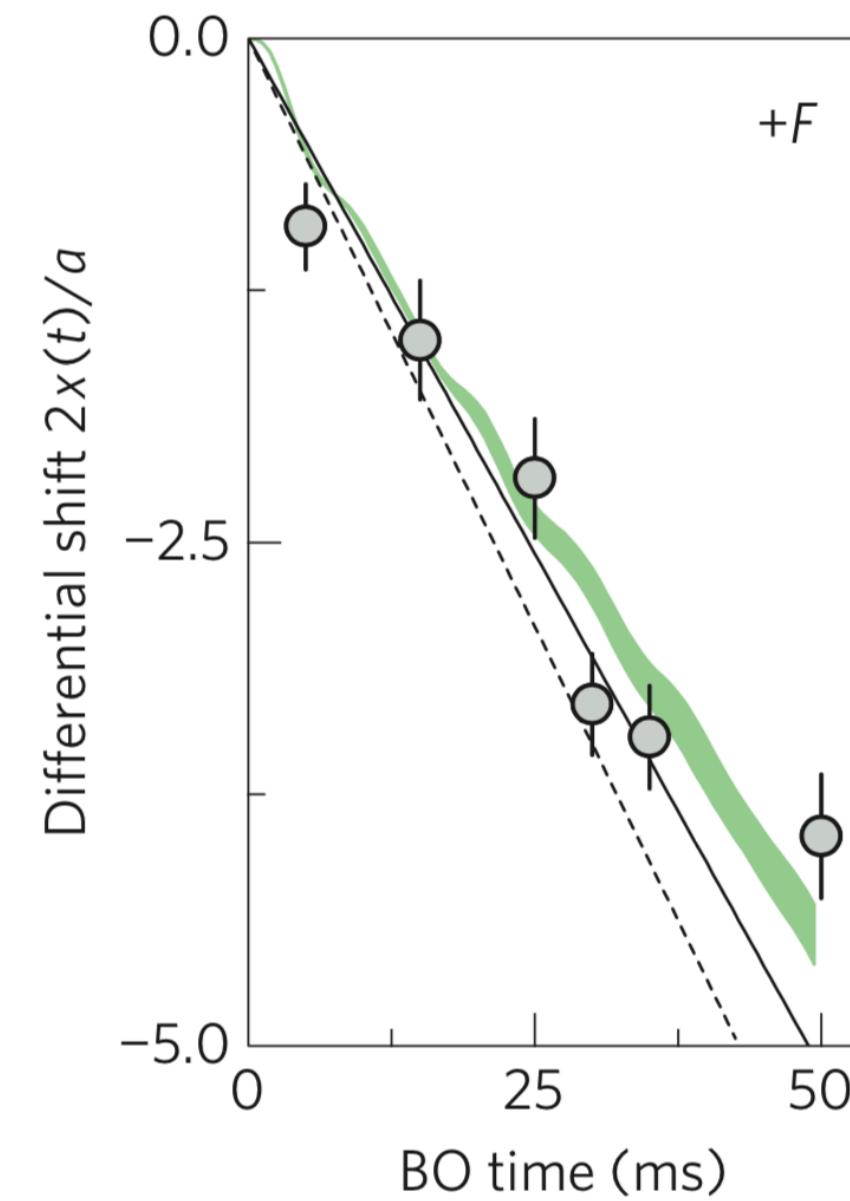
$$x(t) = -(a^2 t E_y / \pi \hbar) \nu_{\text{approx}}$$

How does one know it is a Chern insulator?

Wave packet motion

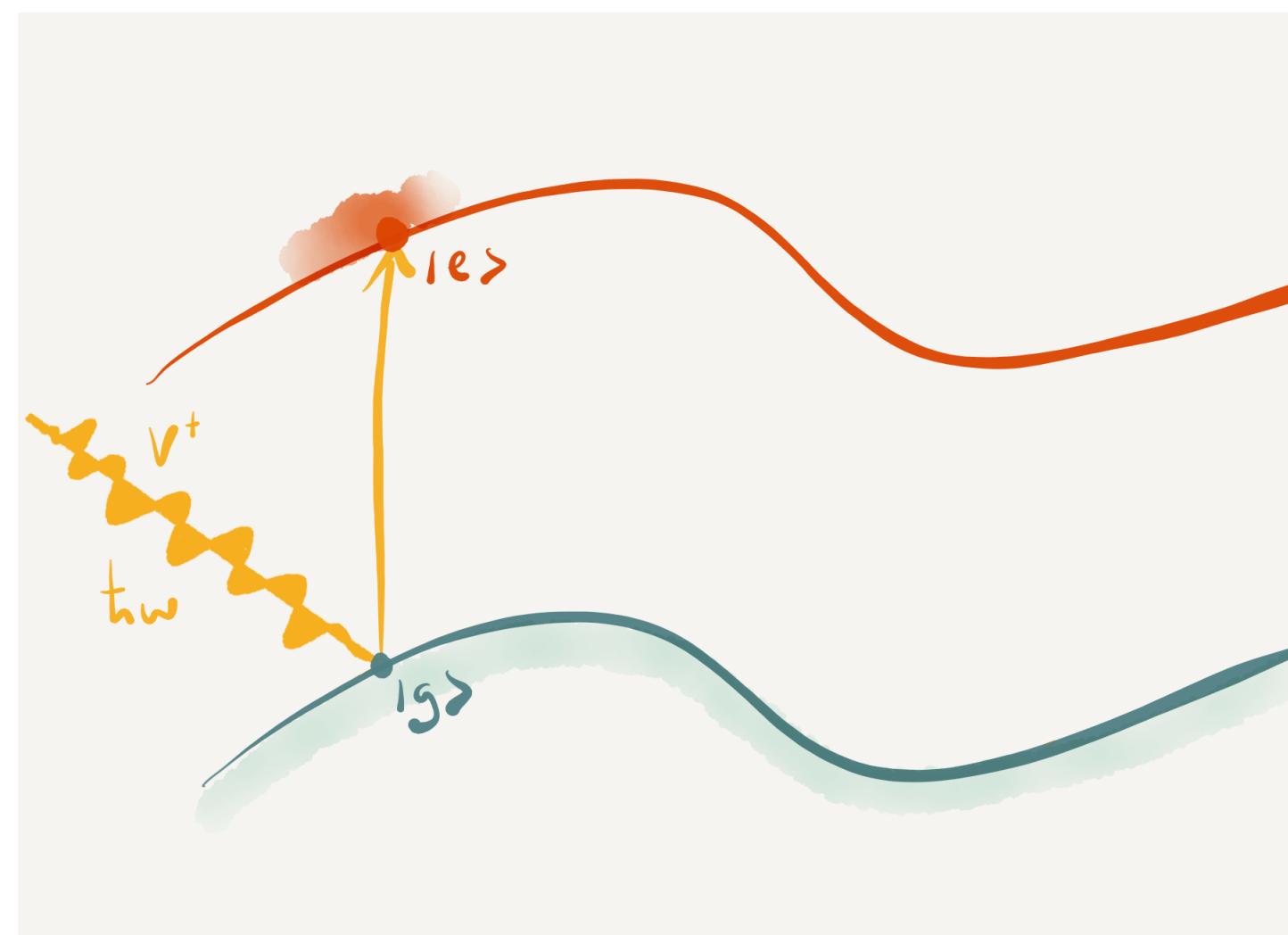


A. Dauphin and N. Goldman PRL 2013



$$x(t) = -(a^2 t E_y / \pi \hbar) \nu_{\text{approx}}$$

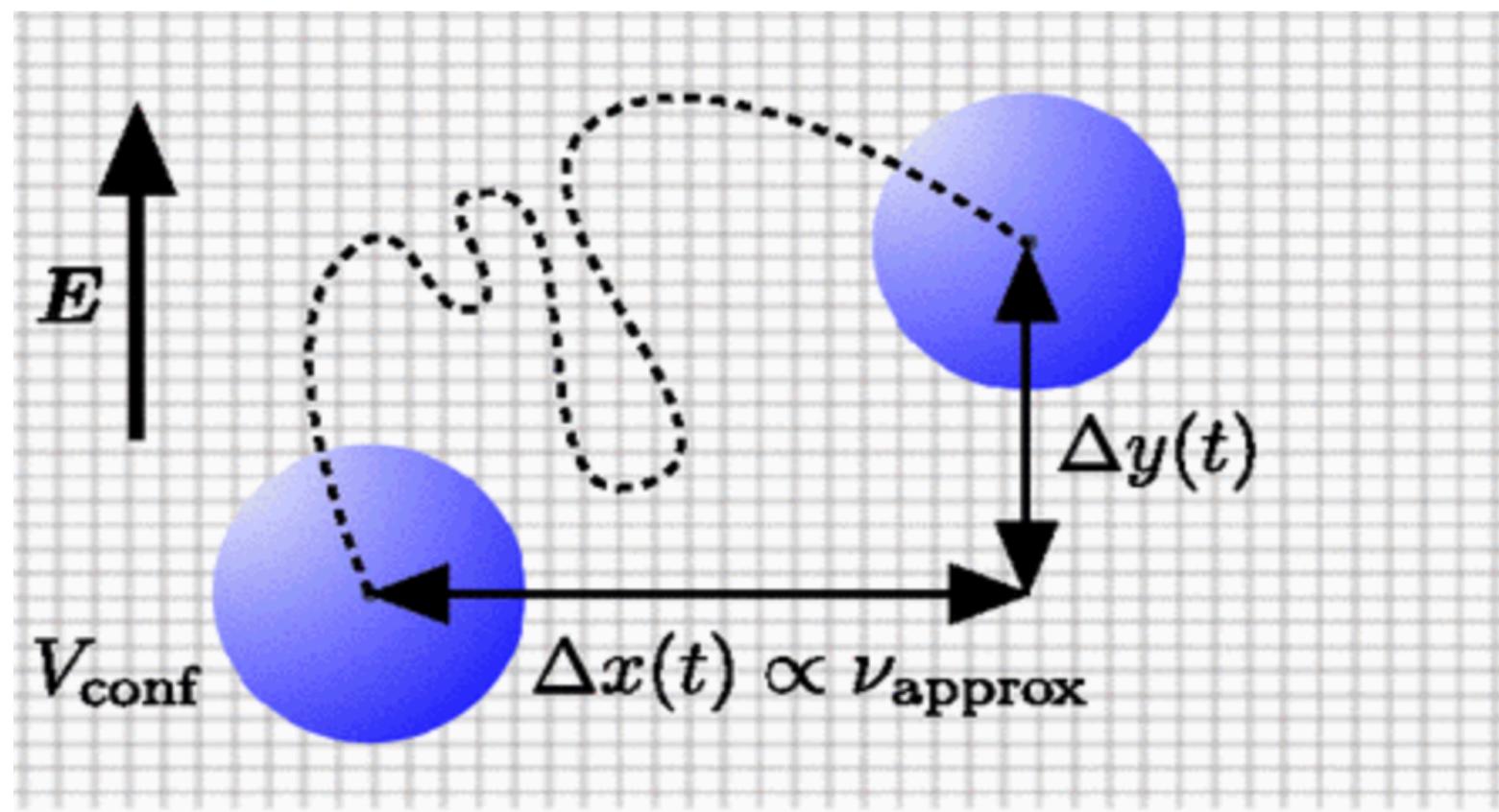
M. Aidelsburger et al. Nat. Phys (2015)



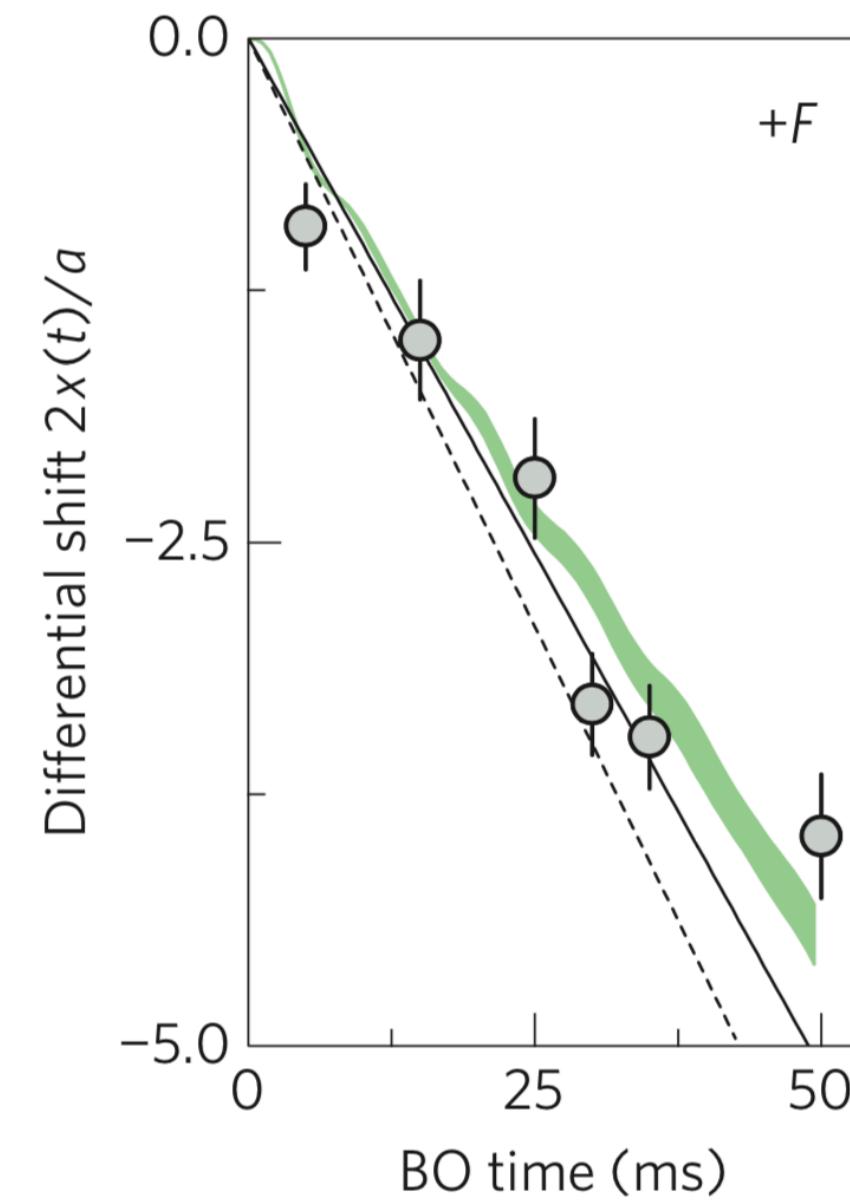
$$(\Gamma_+^{\text{int}} - \Gamma_-^{\text{int}})/2A_{\text{cell}} = (E_{\text{sp}}/\hbar)^2 C$$

How does one know it is a Chern insulator?

Wave packet motion

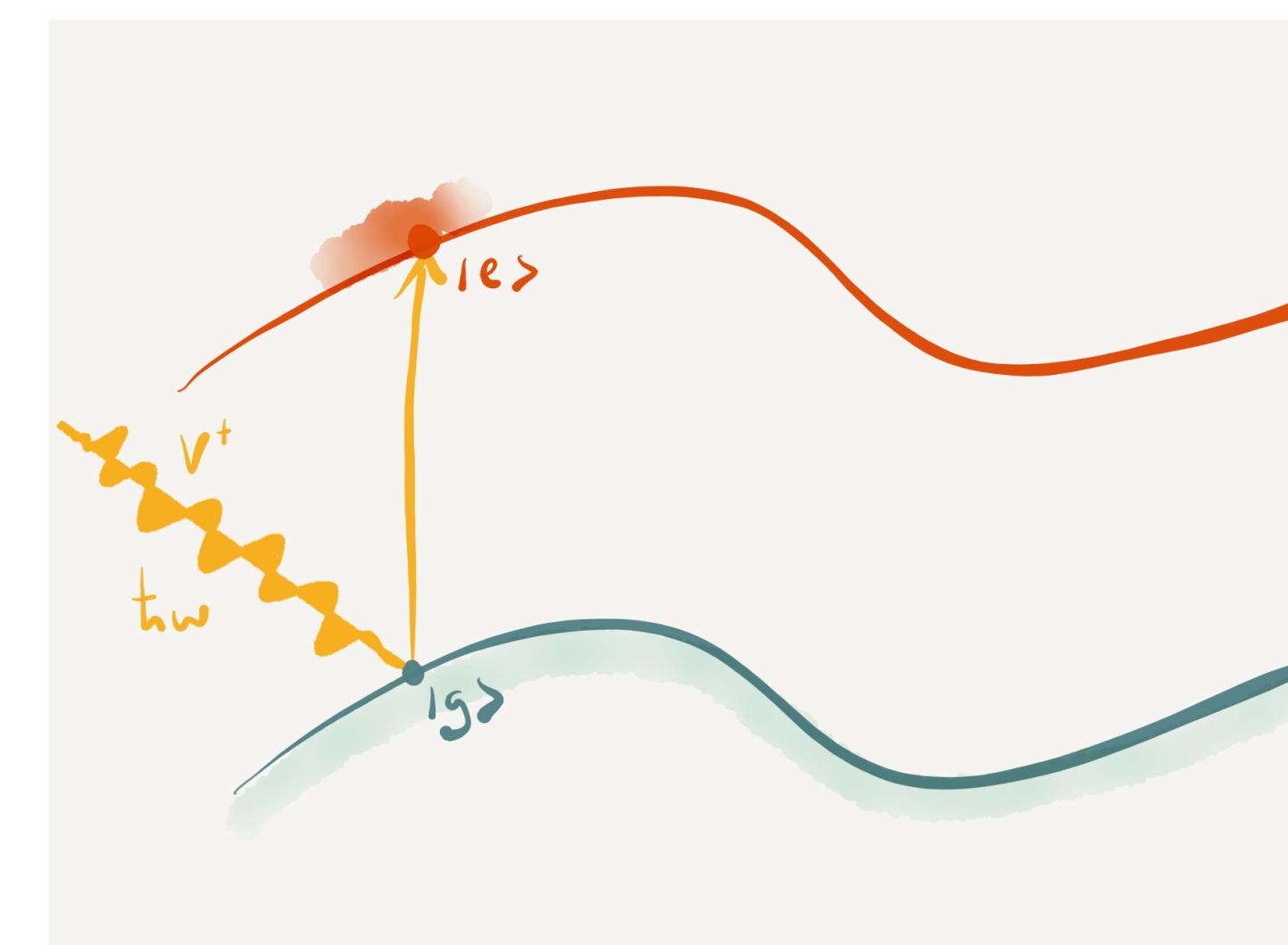


A. Dauphin and N. Goldman PRL 2013



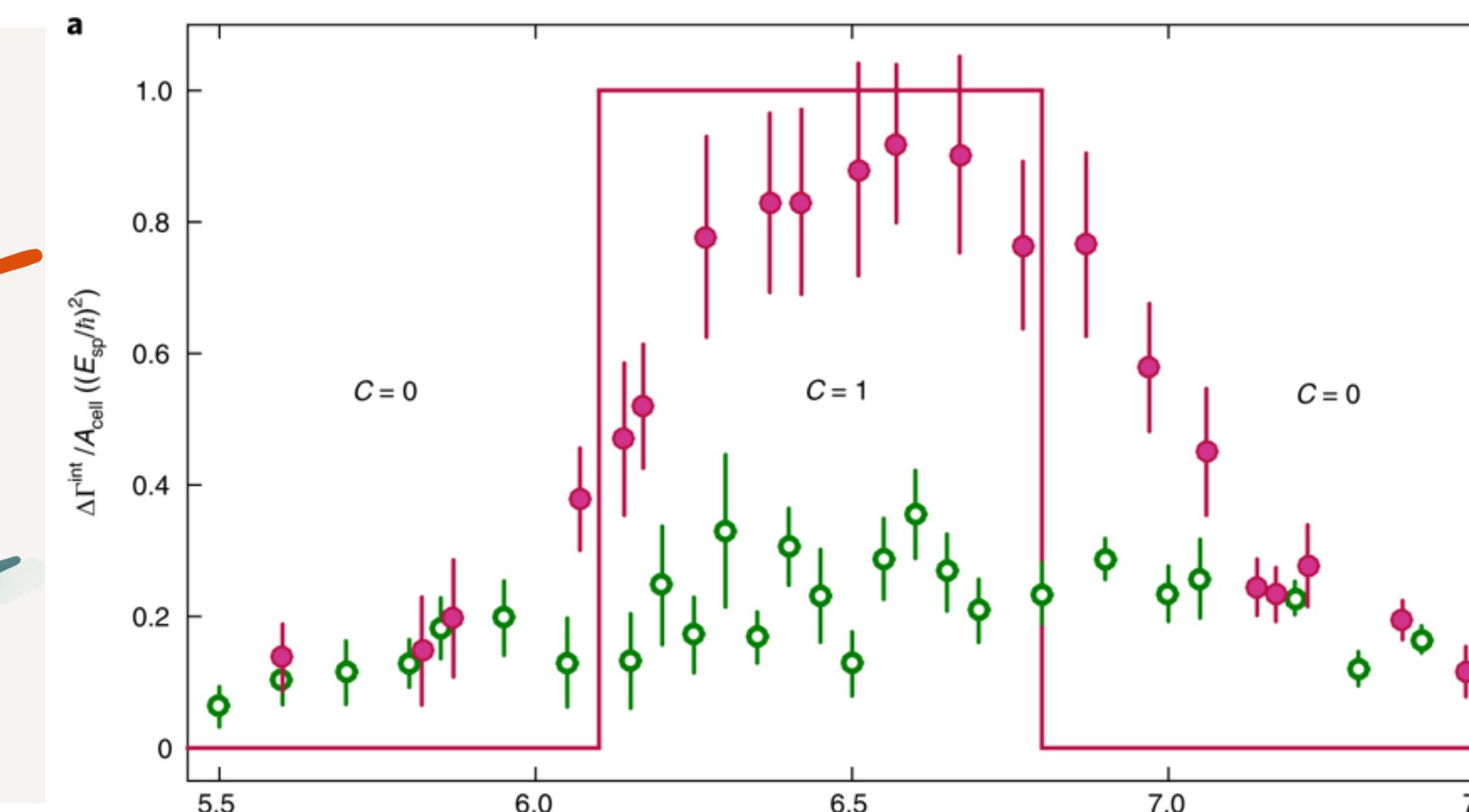
$$x(t) = -(a^2 t E_y / \pi \hbar) \nu_{\text{approx}}$$

Circular dichroism



D. T. Tran, A. Dauphin, AGG, N. Goldman, P. Zoller Sci. Adv. (2018)

M. Aidelsburger et al. Nat. Phys (2015)

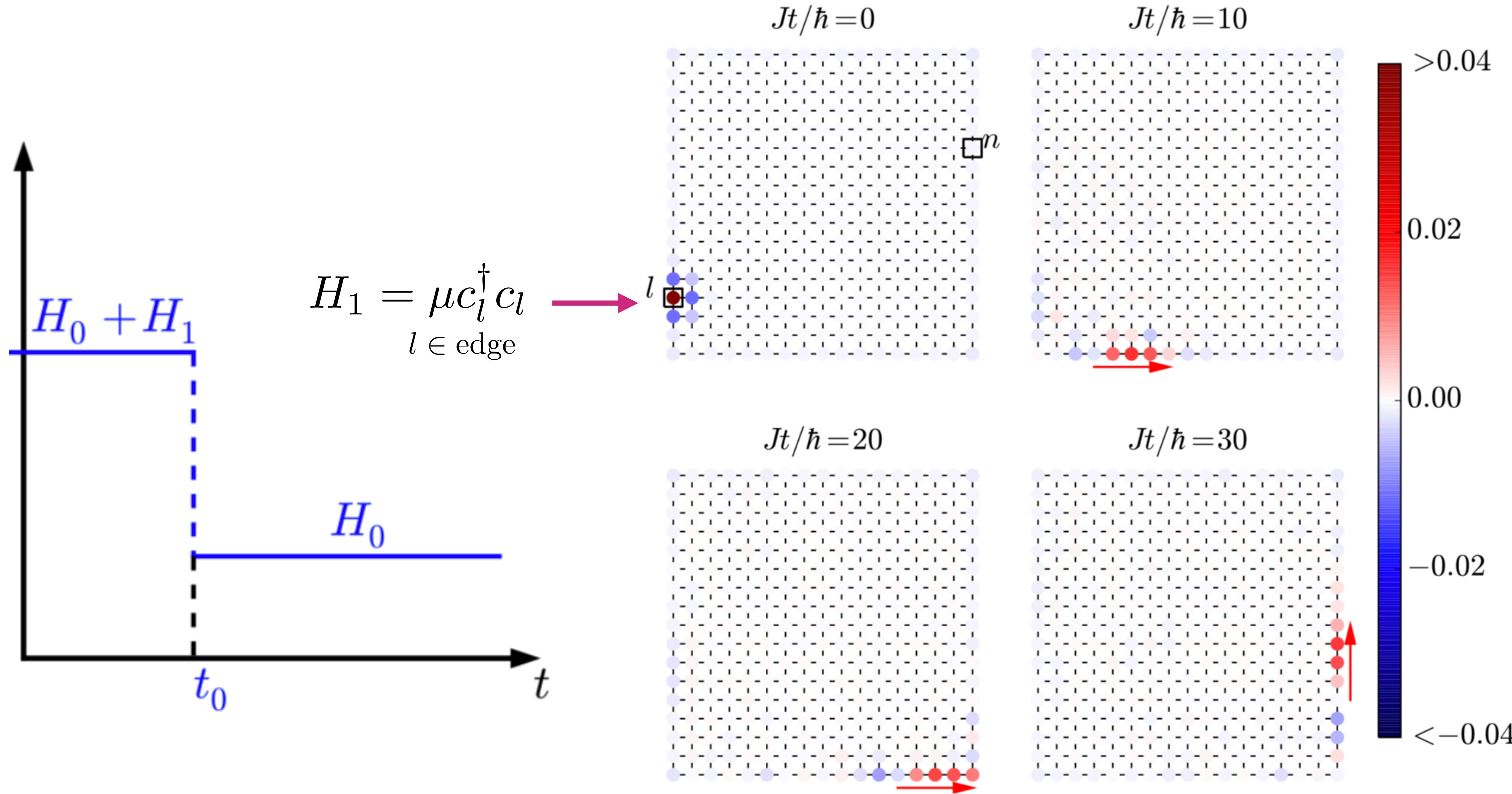


$$(\Gamma_+^{\text{int}} - \Gamma_-^{\text{int}})/2A_{\text{cell}} = (E_{\text{sp}}/\hbar)^2 C$$

L. Asteria, et. Al Nat. Phys, 2019

Can we distinguish a fractional Chern insulator by looking at edge state dynamics?

Edge state dynamics of a Chern insulator



N. Goldman et. al PNAS 2013

A. Grushin et. al J. Stat. Mech. (2014)

Edge state dynamics of a bosonic fractional Chern insulator

Density Matrix Renormalization Group

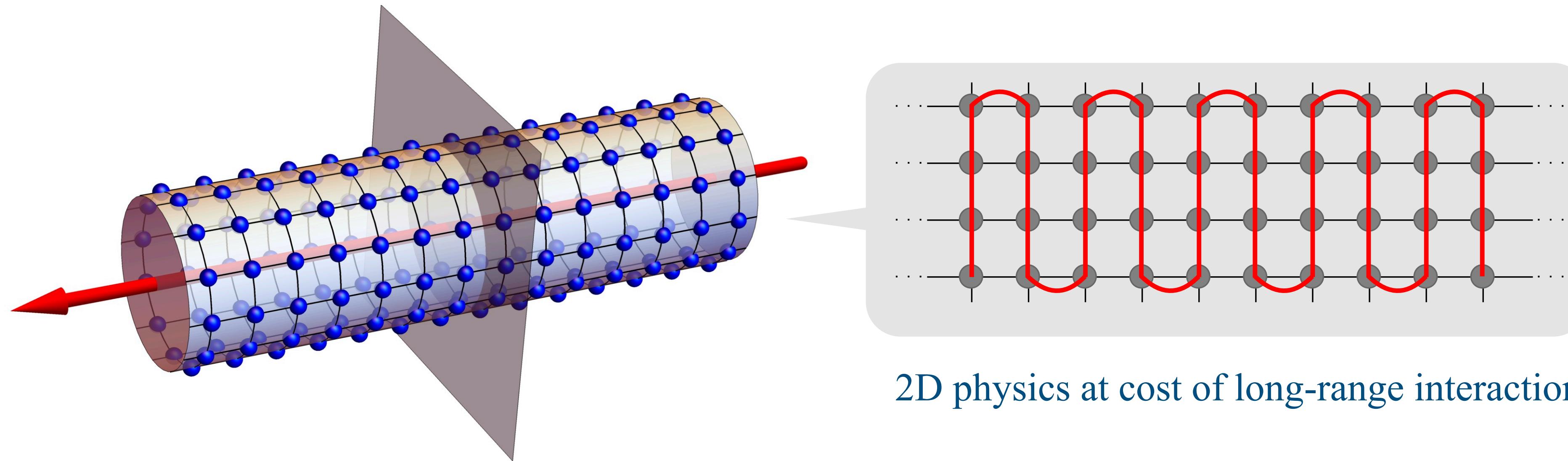
Matrix-Product State representation of the ground state

M. Fannes et al Comm Math. Phys. '92, Schollwoeck Ann. Phys.'11

$$|\psi_0\rangle : \cdots \xrightarrow{B} \xrightarrow{B} \xrightarrow{B} \xrightarrow{B} \xrightarrow{B} \xrightarrow{B} \xrightarrow{B} \cdots \xrightarrow{B_{\alpha\beta}^{i_n}}$$

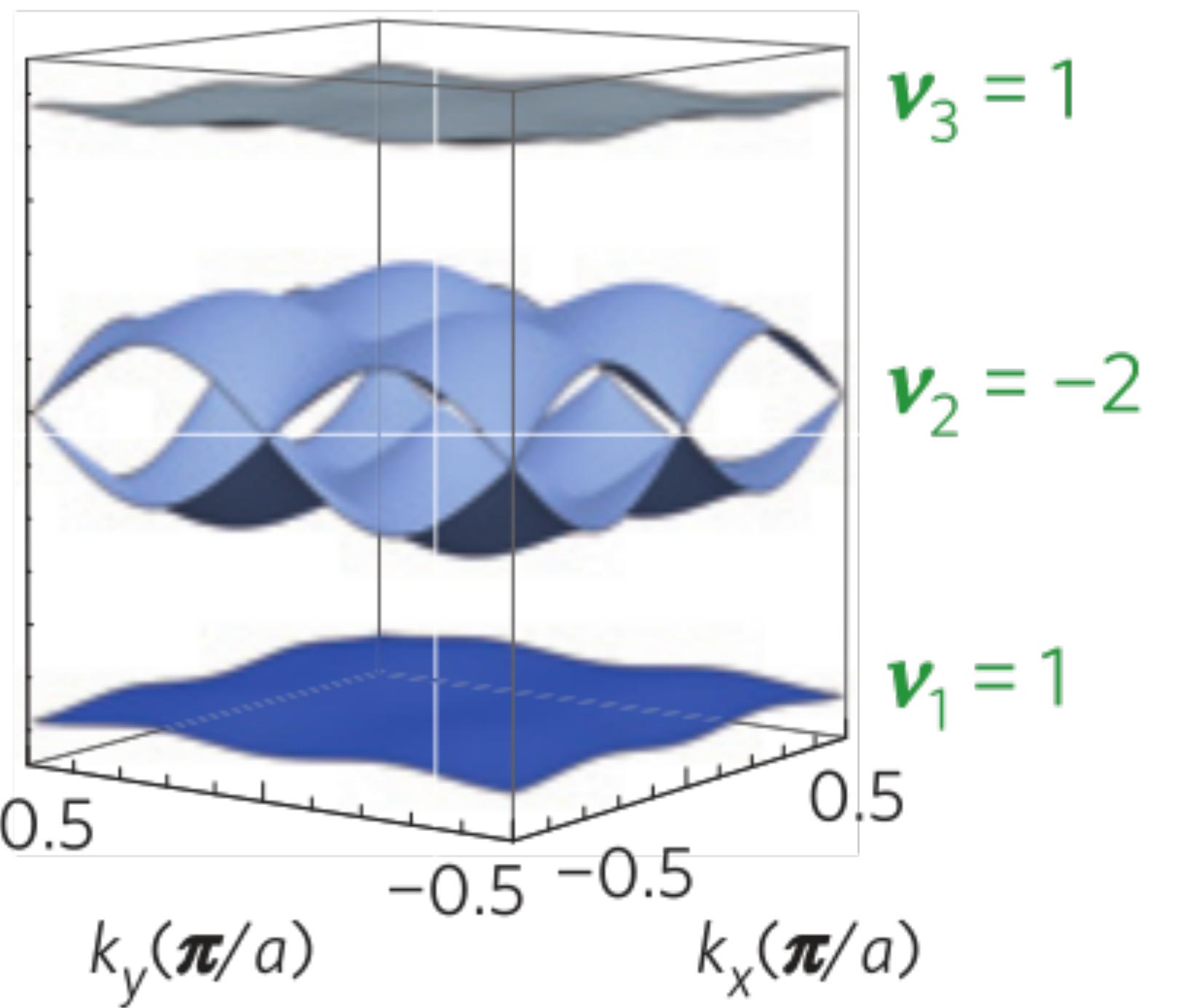
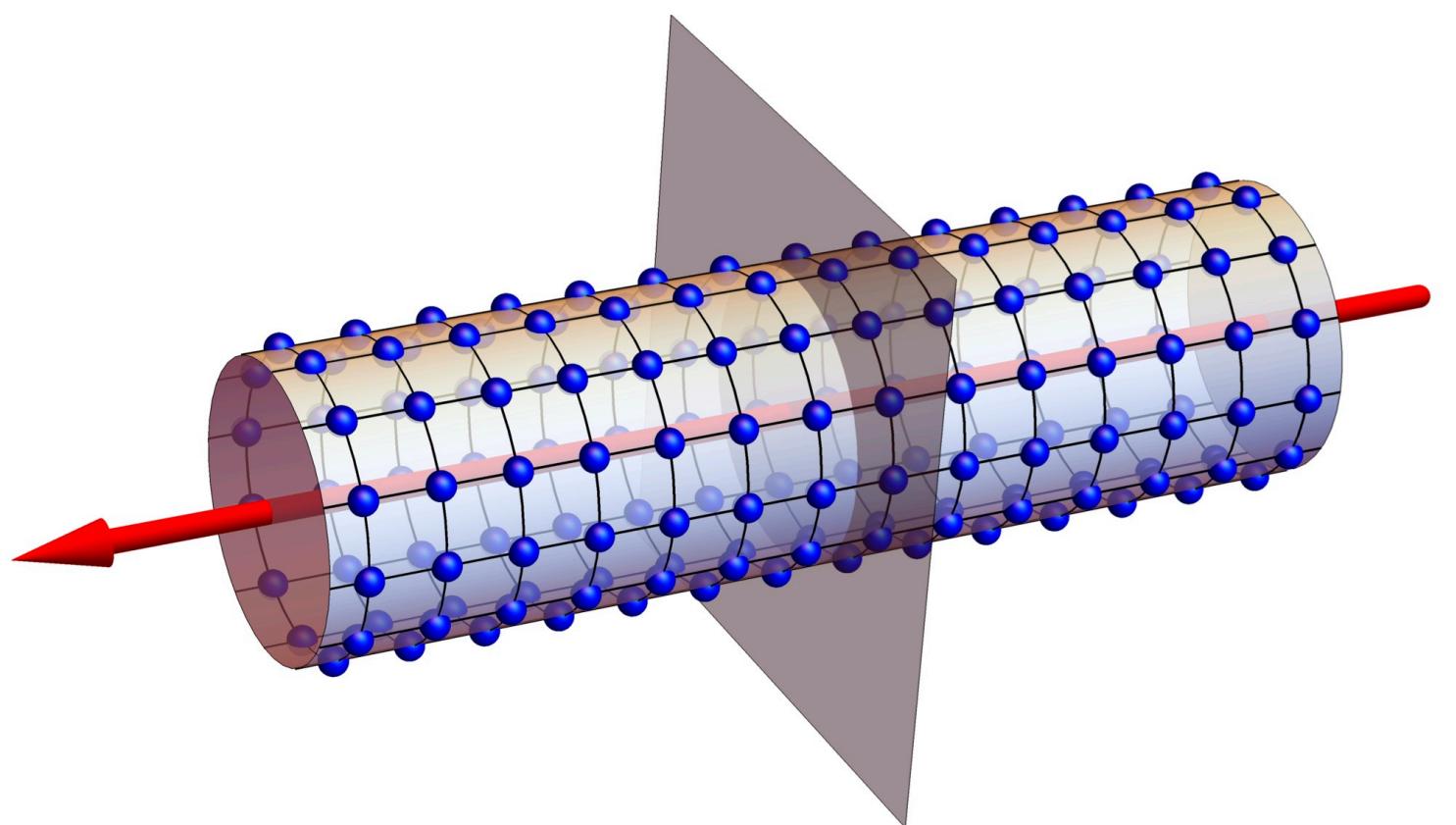
DMRG on cylinders with circumference up to L= 12

S. R. White PRL '92

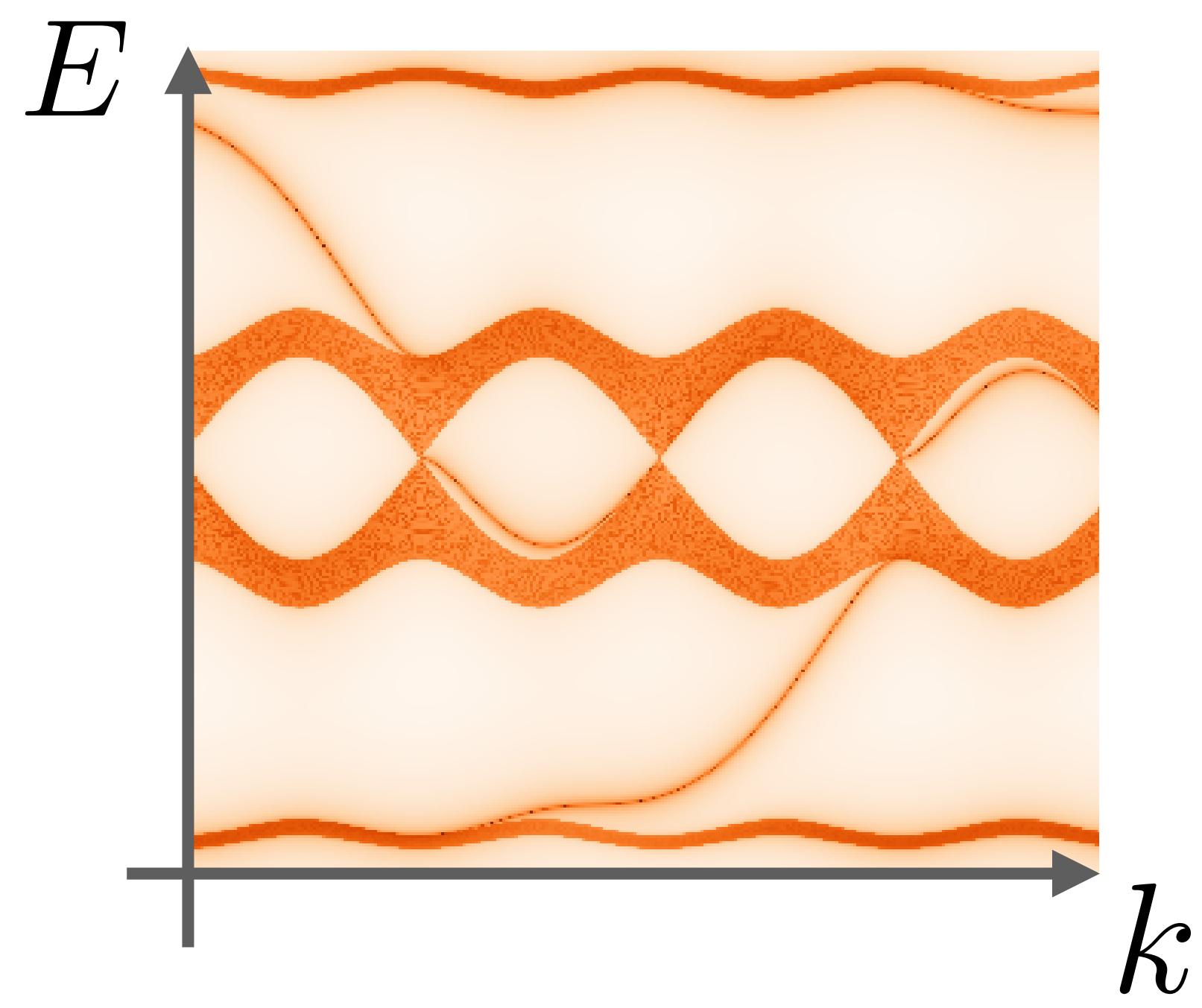
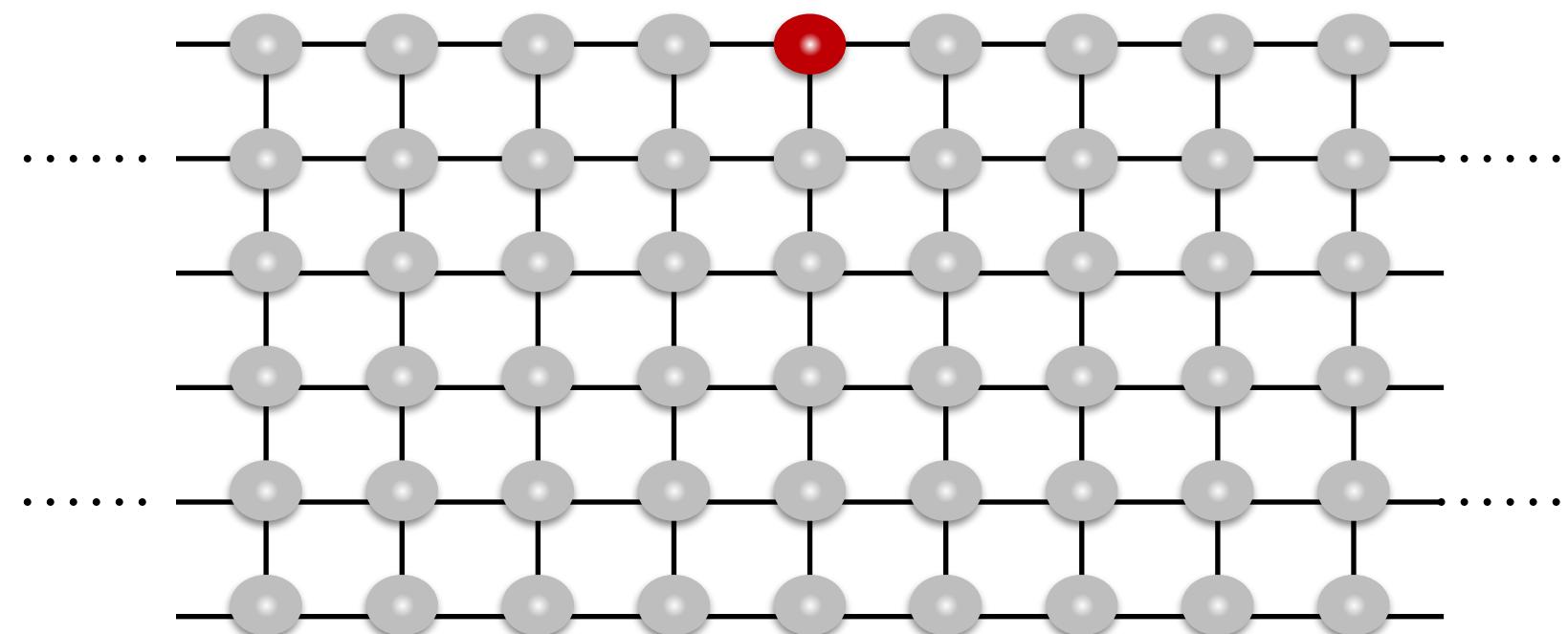


2D physics at cost of long-range interaction in 1D representation

Infinite cylinder



Finite cylinder



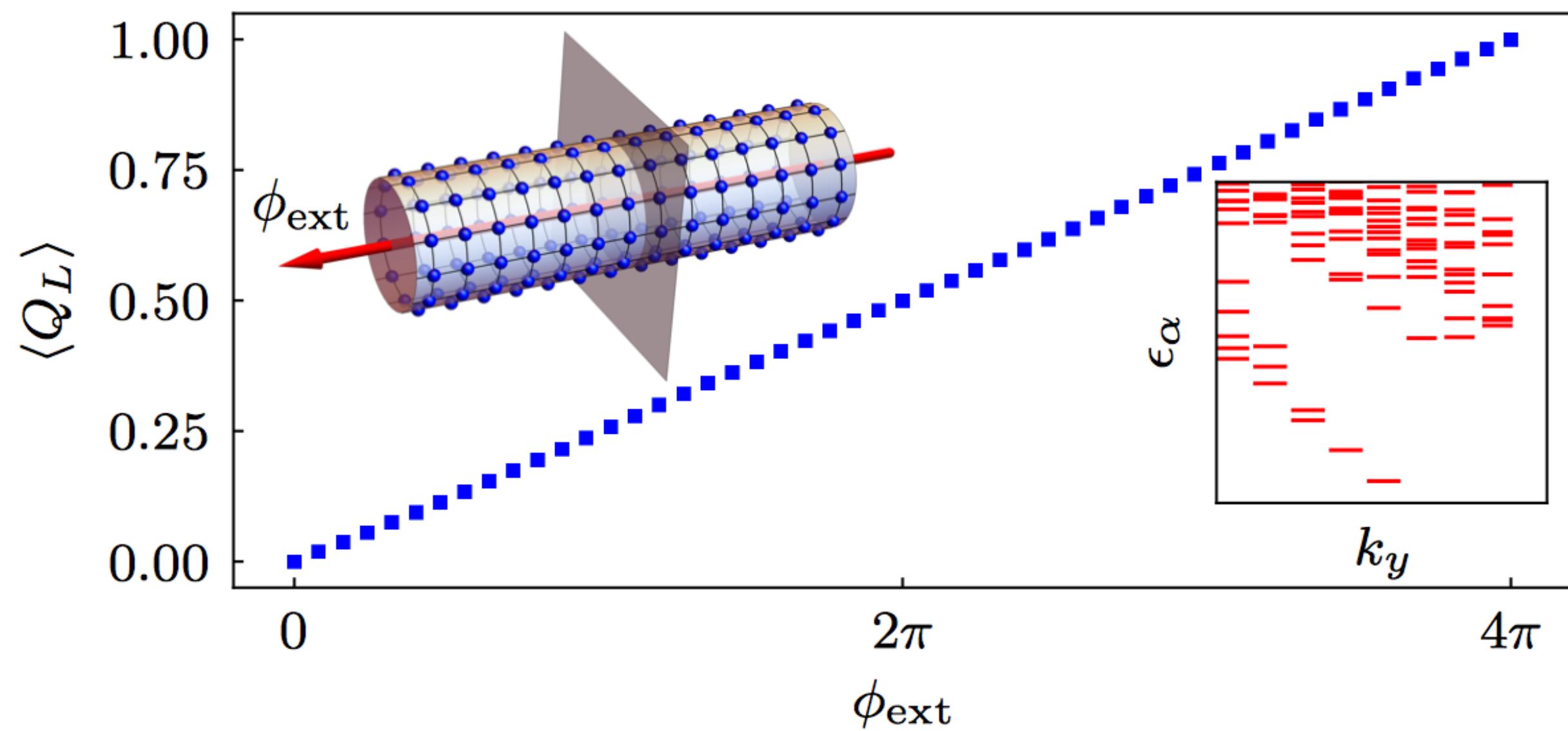
Static ground state: fractional Chern insulator at 1/8 filling

Its been a while...

[Hafezi et al. '07; Möller and Cooper '09; ...]

Quantized Hall conductivity

$$\sigma_{xy} = 1/2$$



Static ground state: fractional Chern insulator at 1/8 filling

Its been a while...

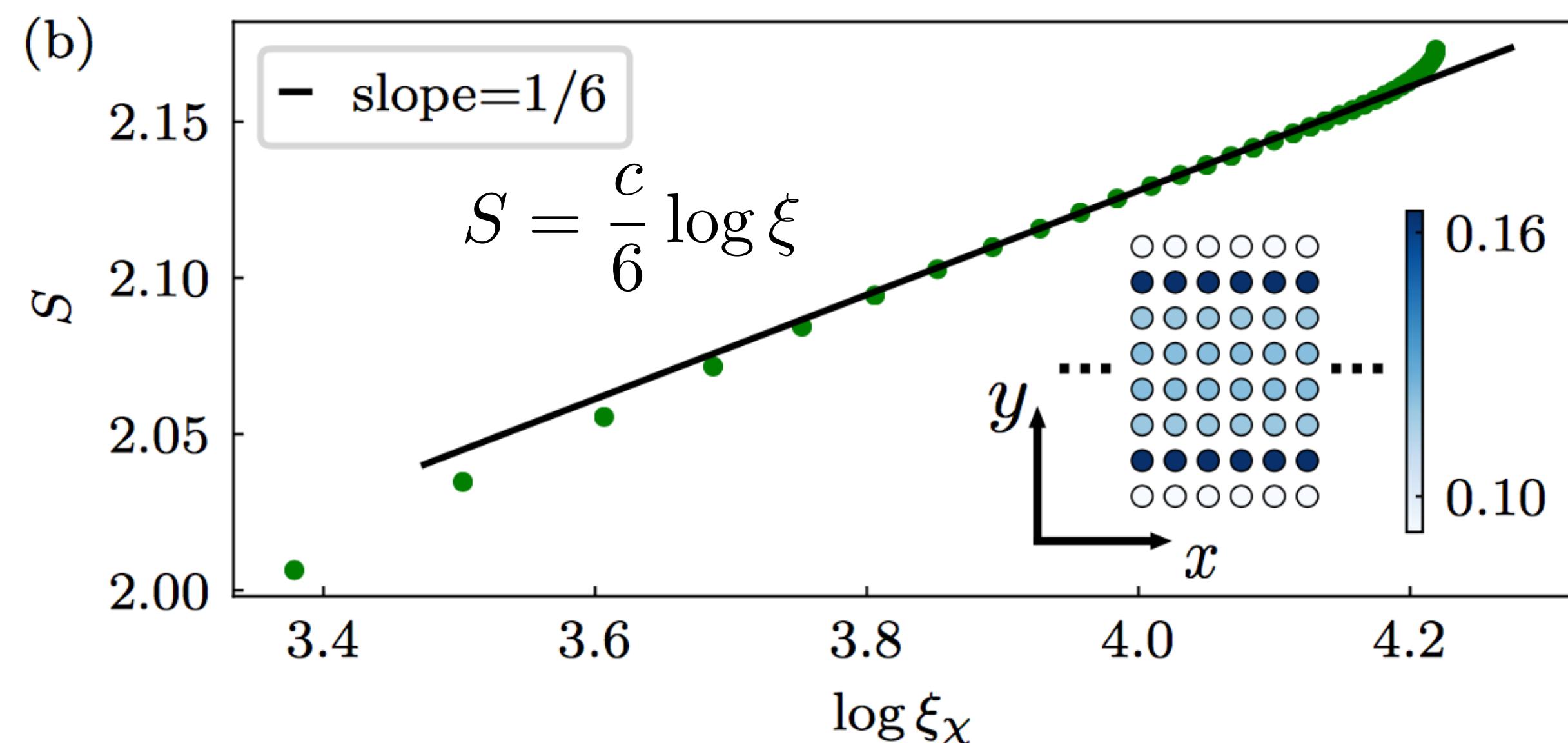
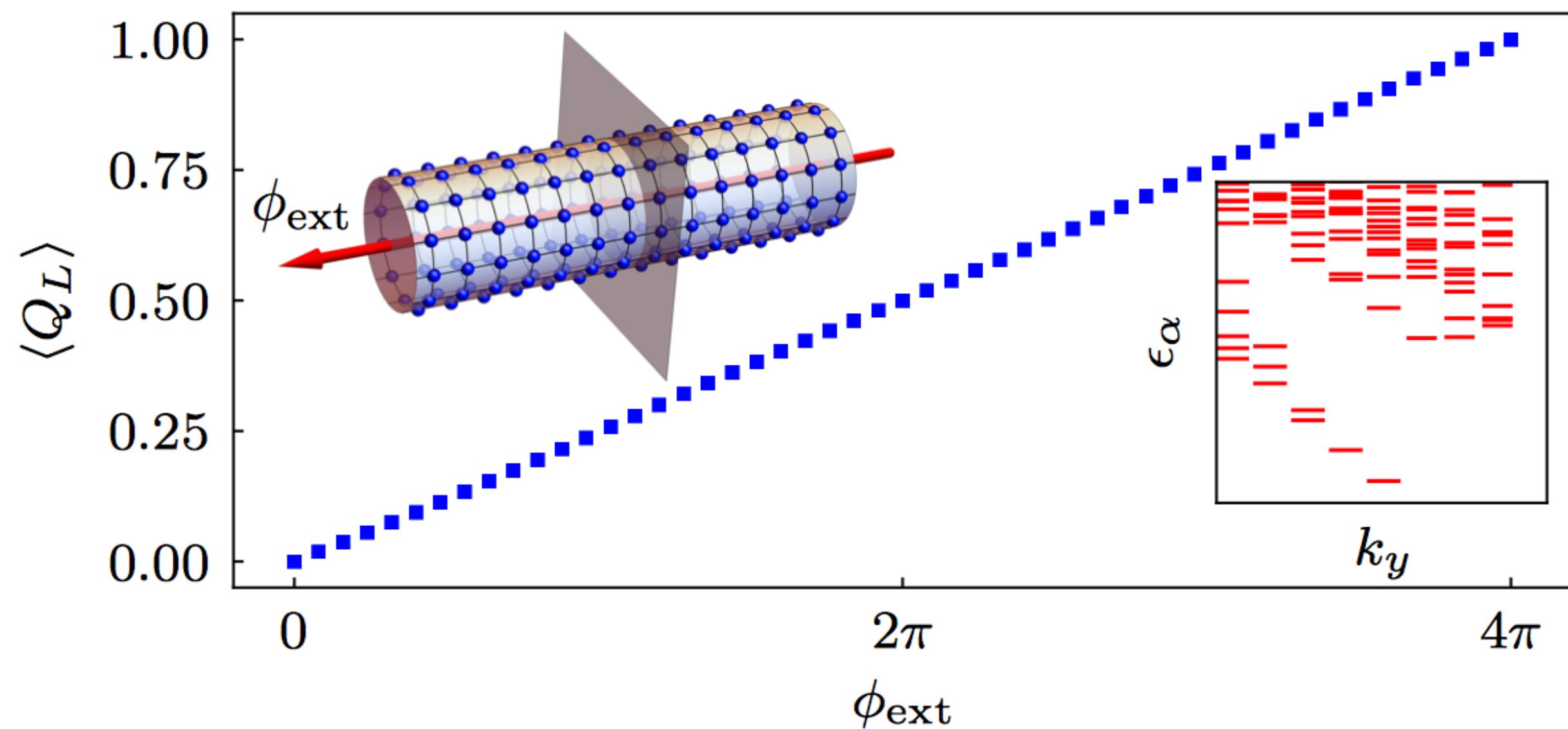
[Hafezi et al. '07; Möller and Cooper '09; ...]

Quantized Hall conductivity

$$\sigma_{xy} = 1/2$$

Gapless edge states

$$c = 1$$



Time evolution

Hamiltonian expressed as a sum of local terms

$$H = \sum_x H_x$$

Expand $U = \exp(-itH)$ for $t \ll 1$:

$$U(t) = 1 + t \sum_x H_x + \frac{1}{2} t^2 \sum_{x,y} H_x H_y + \dots$$

Compact Matrix Product Operator (MPO)

$$W_{\alpha\beta}^{[n]j_n j'_n} = \alpha \begin{array}{c} j'_n \\ \diamond \\ j_n \end{array} \beta$$

Time evolution

Hamiltonian expressed as a sum of local terms

$$H = \sum_x H_x$$

Expand $U = \exp(-itH)$ for $t \ll 1$:

$$U(t) = 1 + t \sum_x H_x + \frac{1}{2} t^2 \sum_{x,y} H_x H_y + \dots$$

$$\approx 1 + t \sum_x H_x + t^2 \underbrace{\sum_{x < y} H_x H_y}_{\epsilon \sim \underline{L} t^2}$$

Neglect overlapping
terms in expansion

Compact Matrix Product Operator (MPO)

$$W_{\alpha\beta}^{[n]j_n j'_n} = \alpha \begin{array}{c} j'_n \\ \diamond \\ j_n \end{array} \beta$$

Time evolution

Hamiltonian expressed as a sum of local terms

$$H = \sum_x H_x$$

Expand $U = \exp(-itH)$ for $t \ll 1$:

$$U(t) = 1 + t \sum_x H_x + \frac{1}{2} t^2 \sum_{x,y} H_x H_y + \dots$$

$$\approx 1 + t \sum_x H_x + t^2 \underbrace{\sum_{x < y} H_x H_y}_{\epsilon \sim \underline{L} t^2}$$

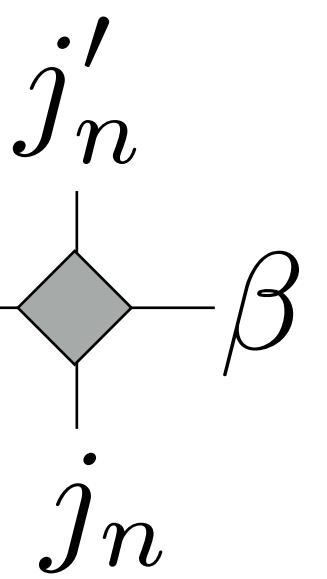
$$1 + t \sum_x H_x \rightarrow \underbrace{\prod_x (1 + t H_x)}_{\epsilon \sim \underline{L}^2 t^2}$$

Neglect overlapping
terms in expansion

M. P. Zaletel et al 'PRB 15

Compact Matrix Product Operator (MPO)

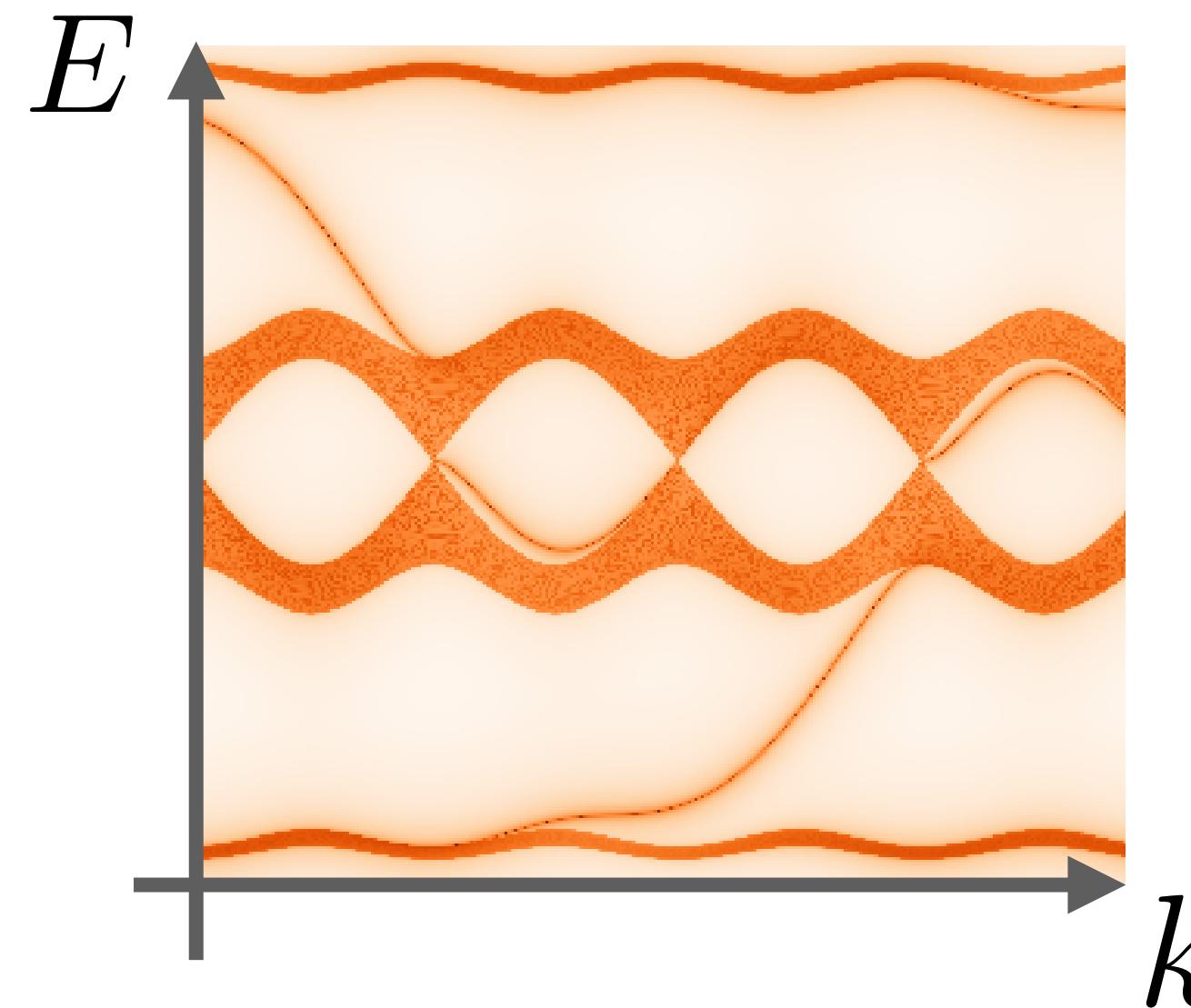
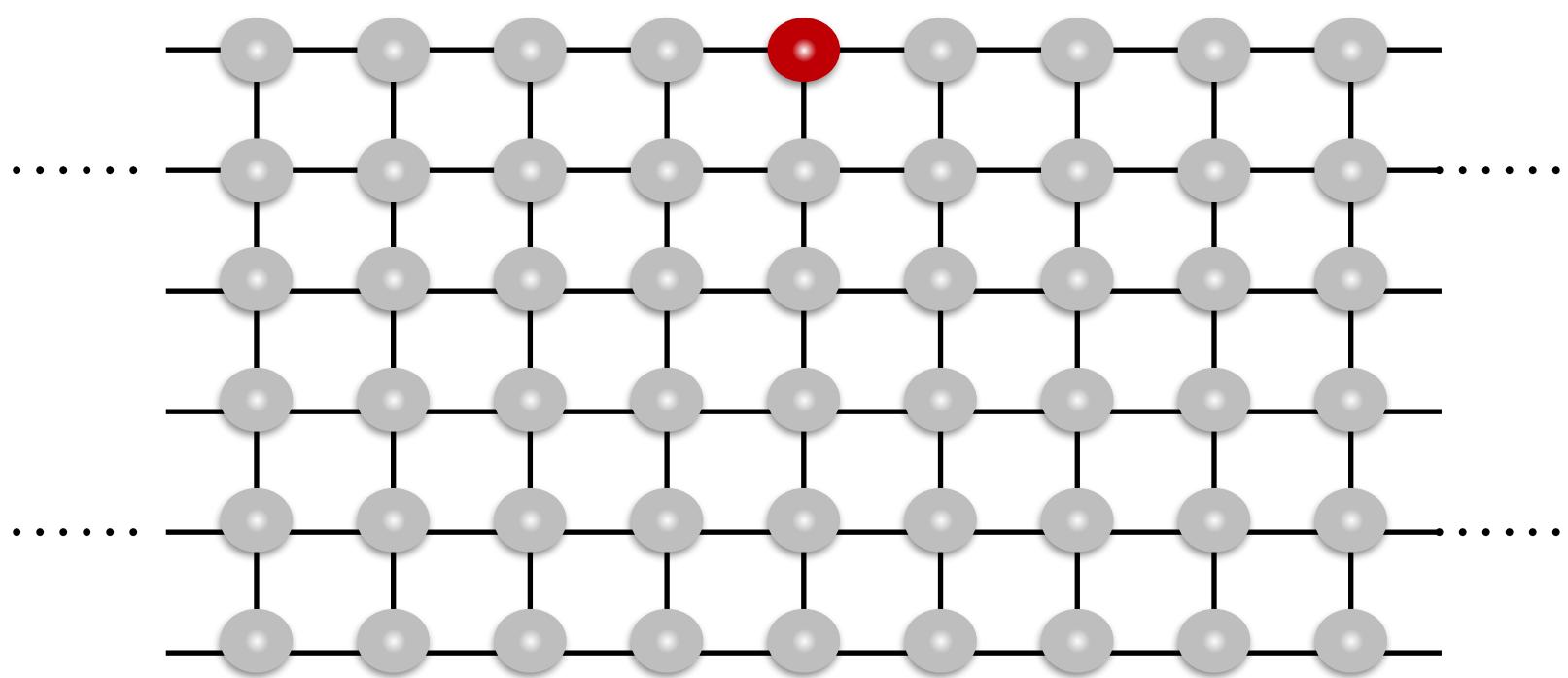
$$W_{\alpha\beta}^{[n]j_n j'_n} = \alpha \begin{array}{c} j'_n \\ \diamond \\ j_n \end{array} \beta$$



Dynamical signatures of the FCI phase

Dynamical correlation function

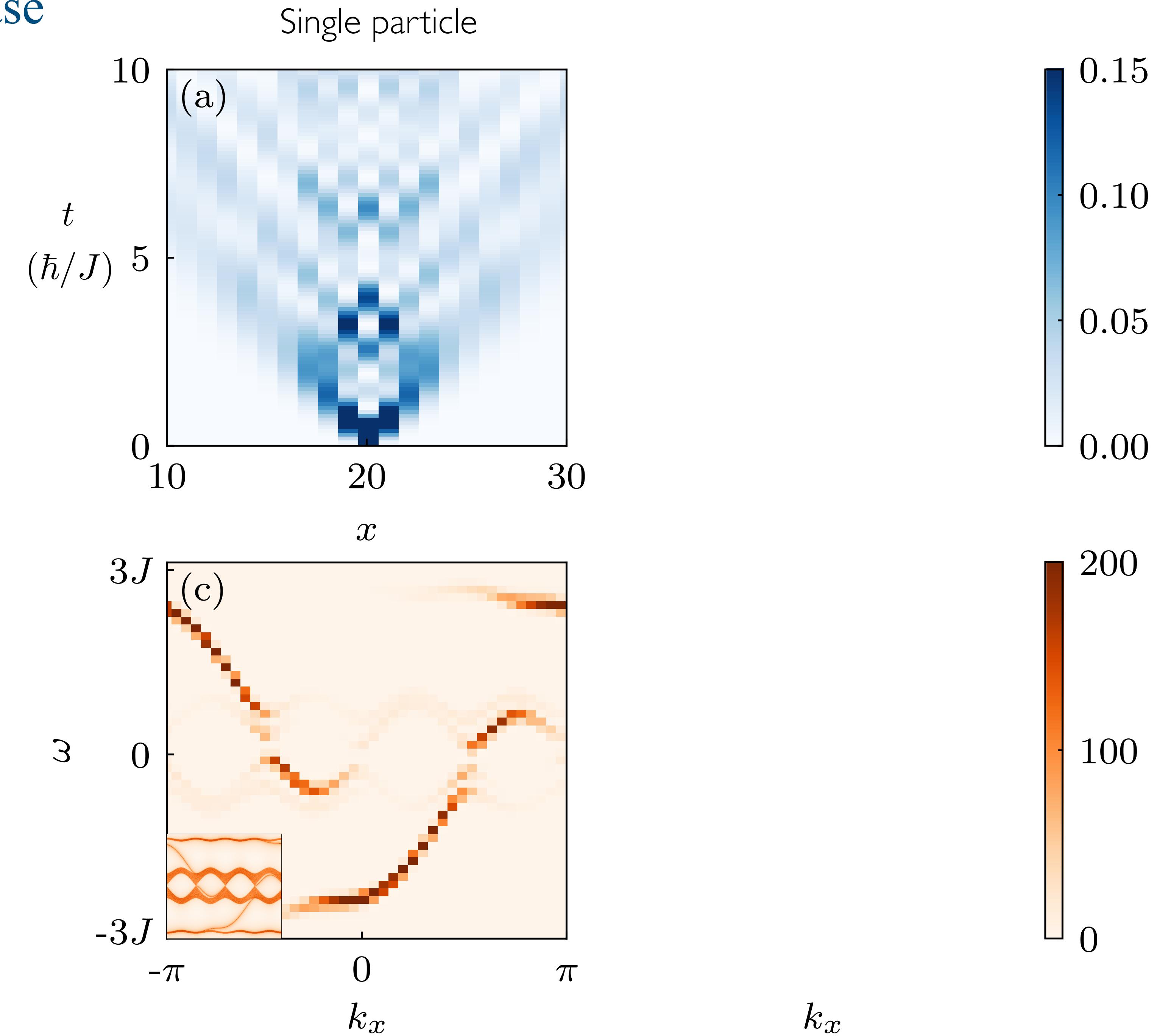
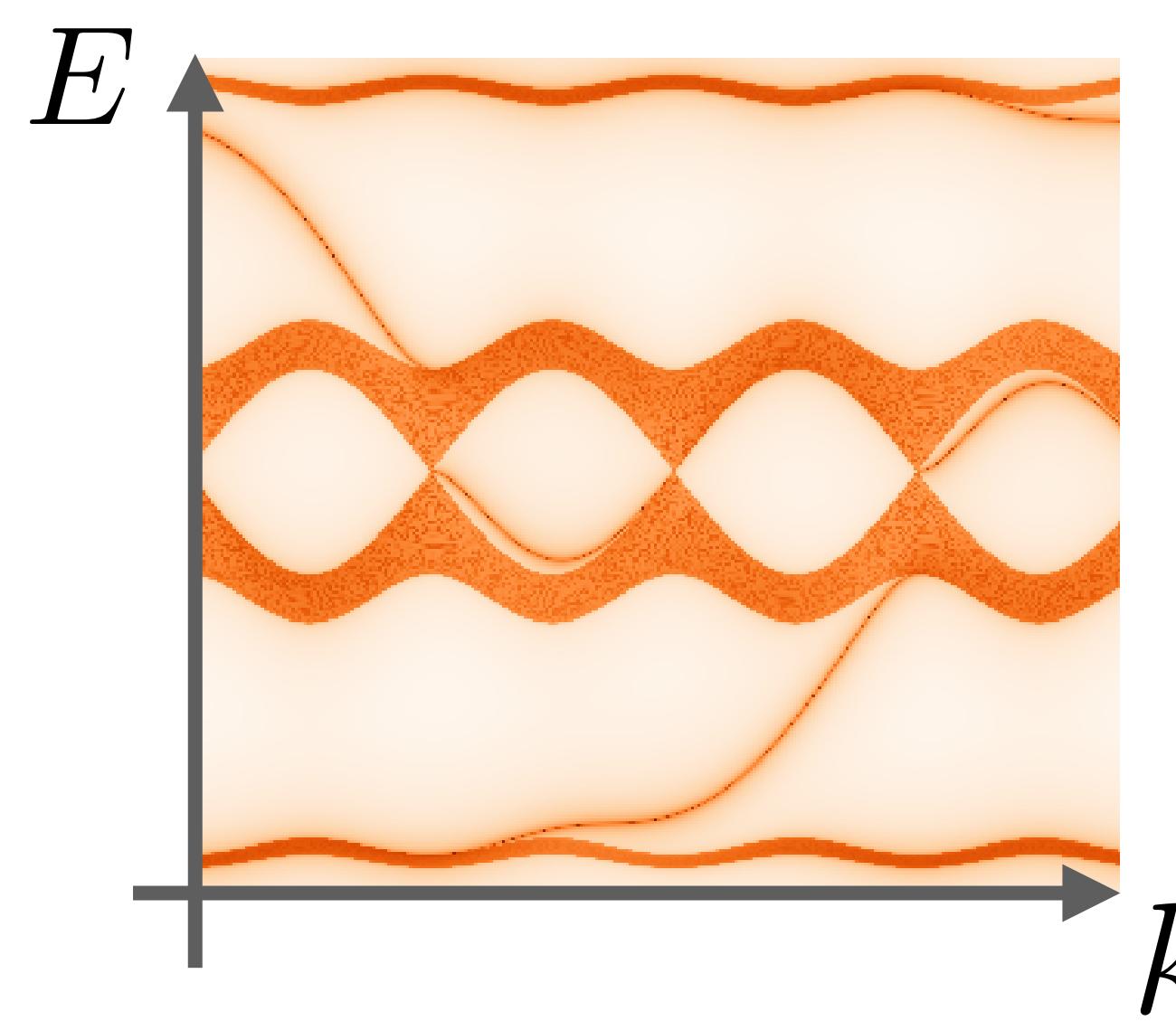
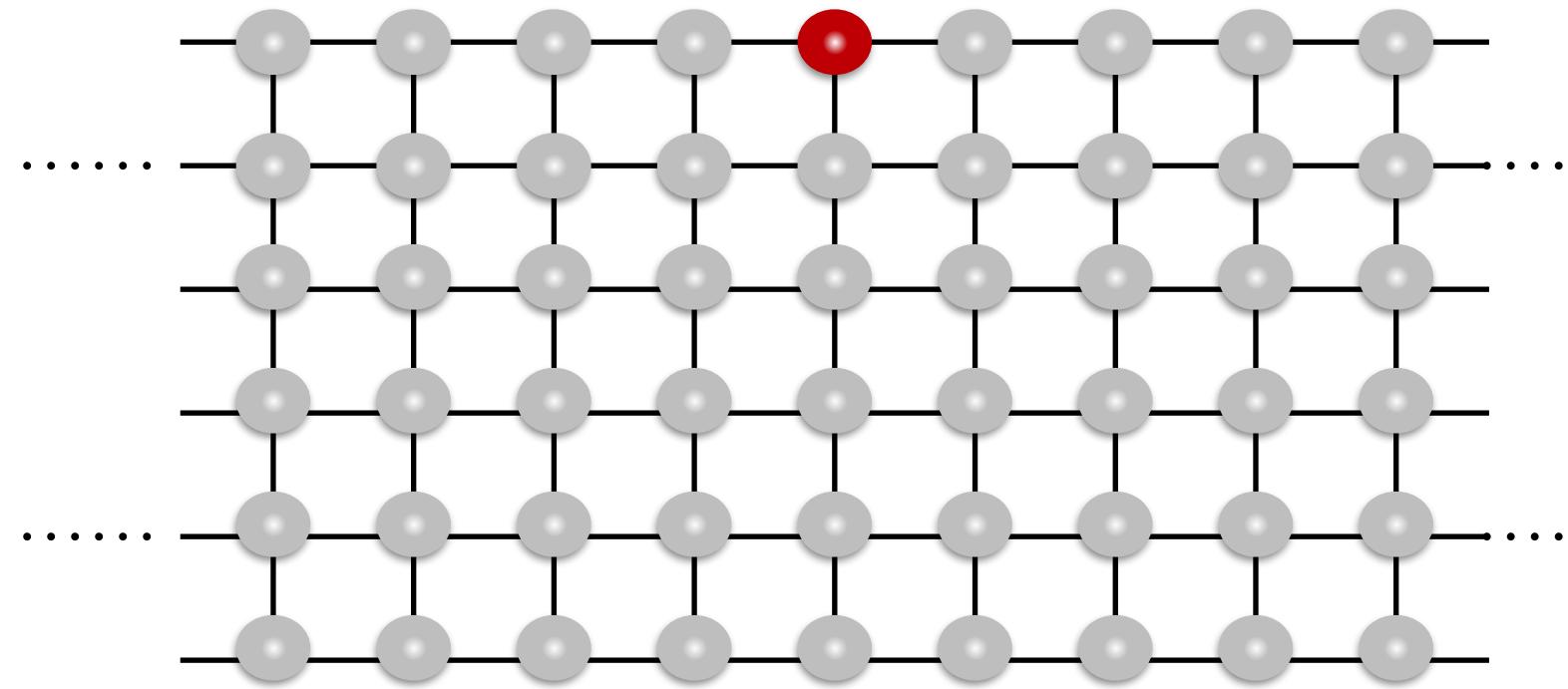
$$G(x, t) = \langle b(x, t) b^\dagger(0, 0) \rangle$$



Dynamical signatures of the FCI phase

Dynamical correlation function

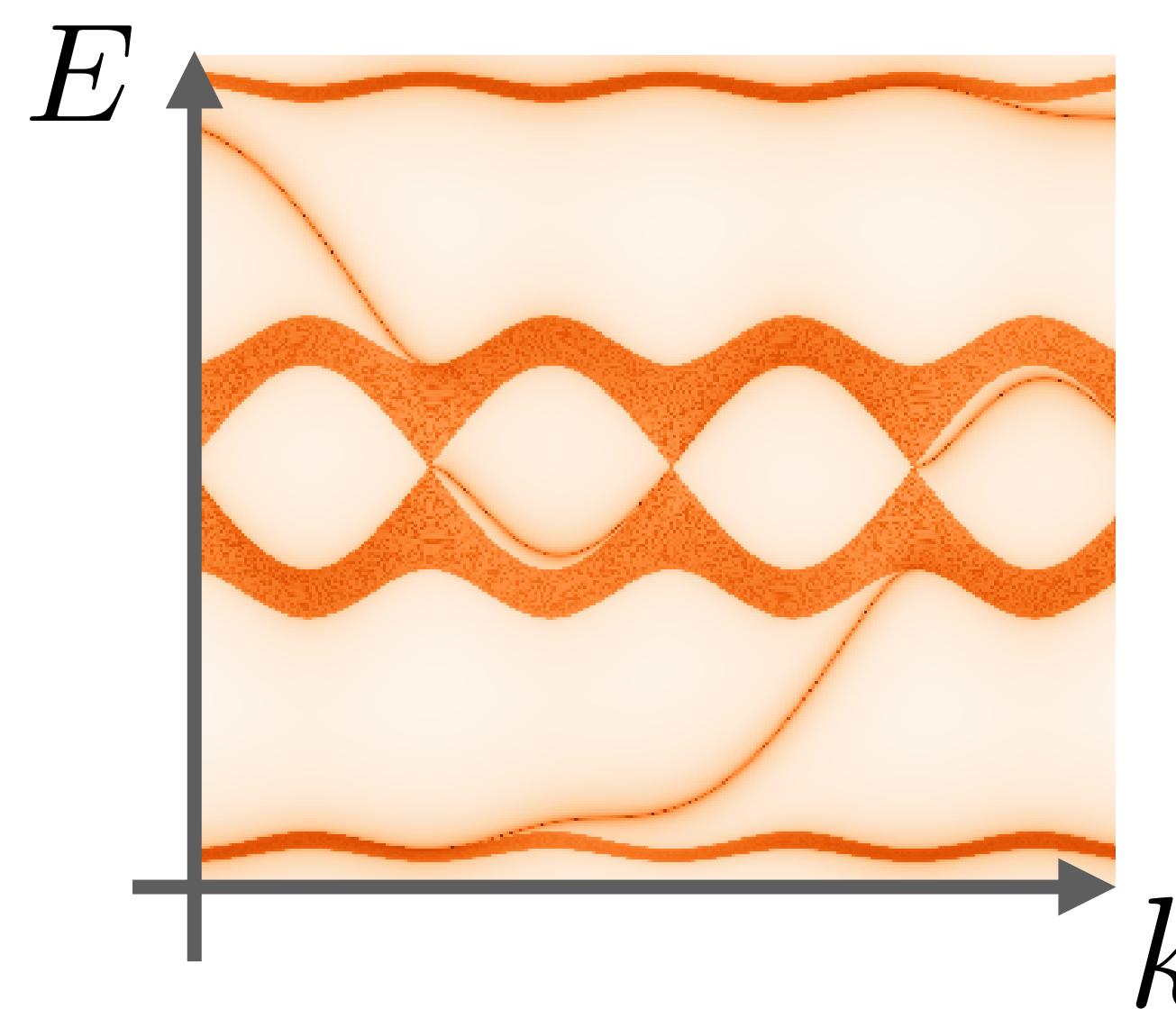
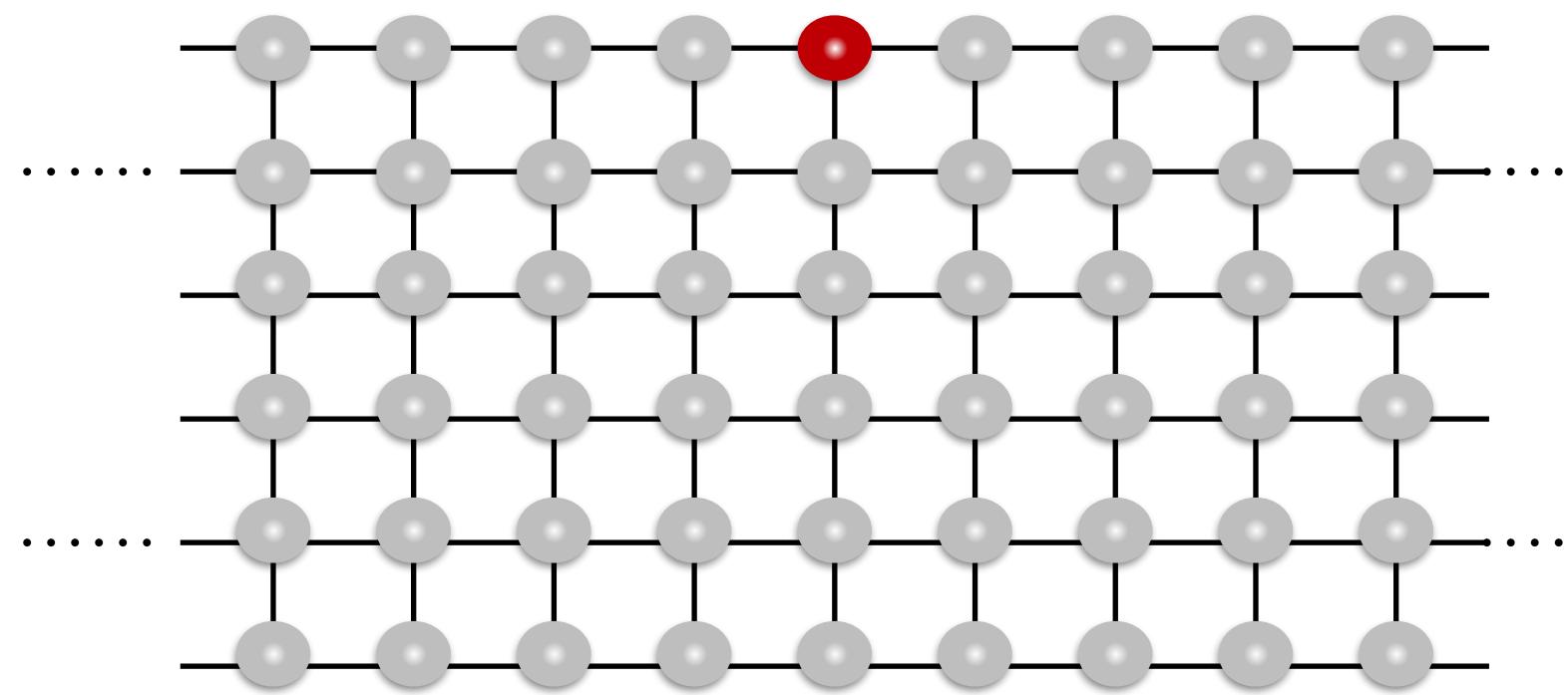
$$G(x, t) = \langle b(x, t) b^\dagger(0, 0) \rangle$$



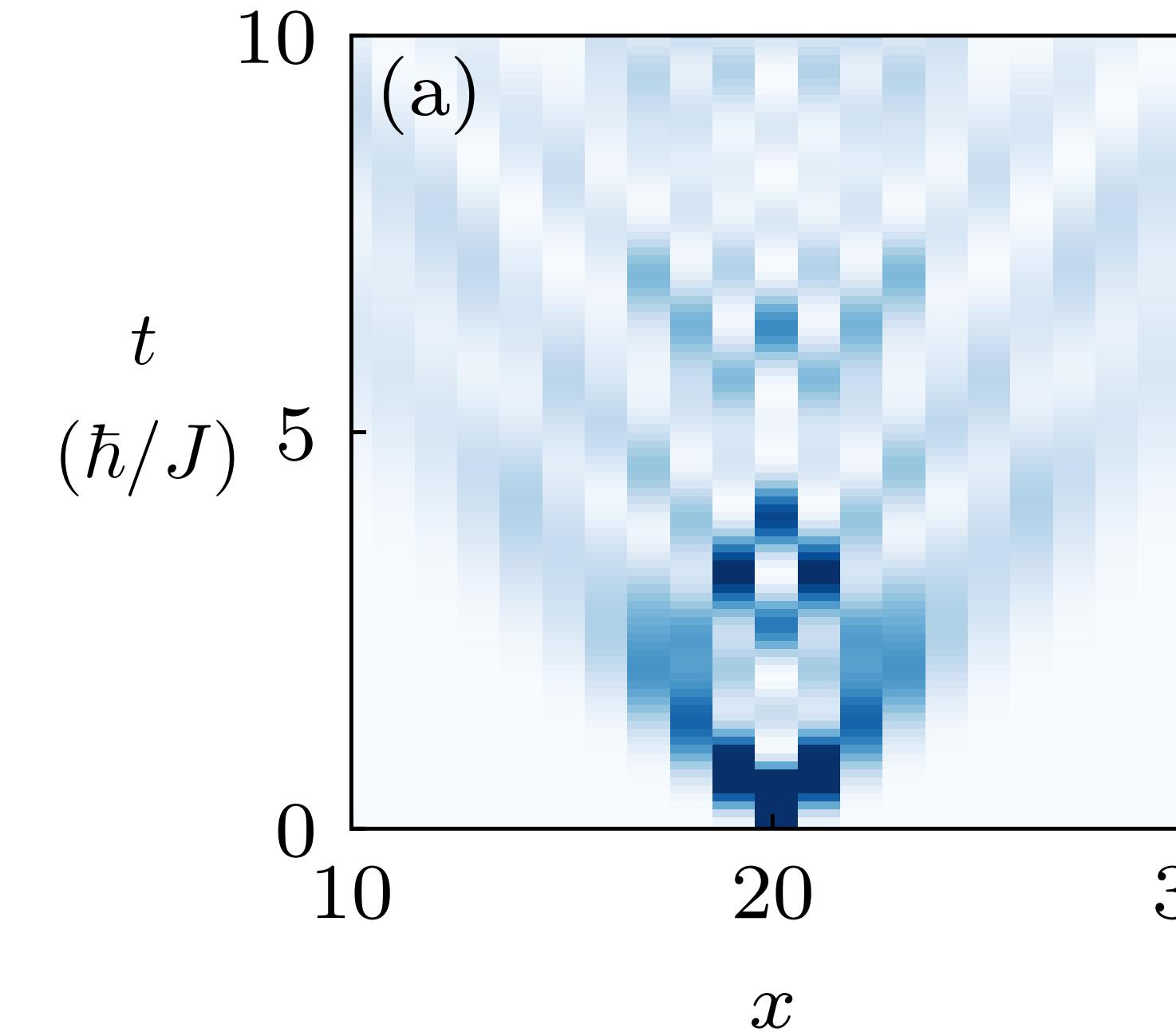
Dynamical signatures of the FCI phase

Dynamical correlation function

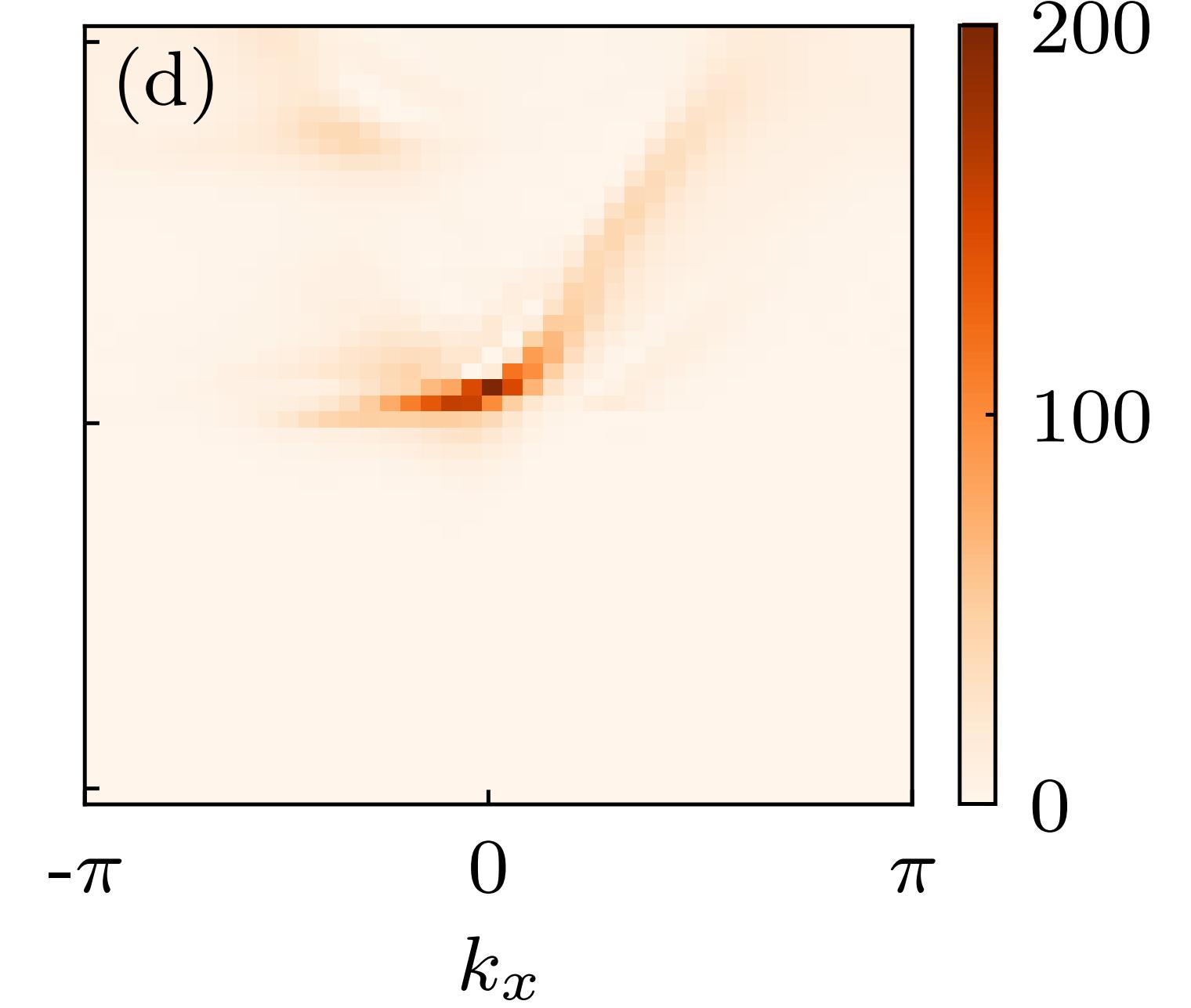
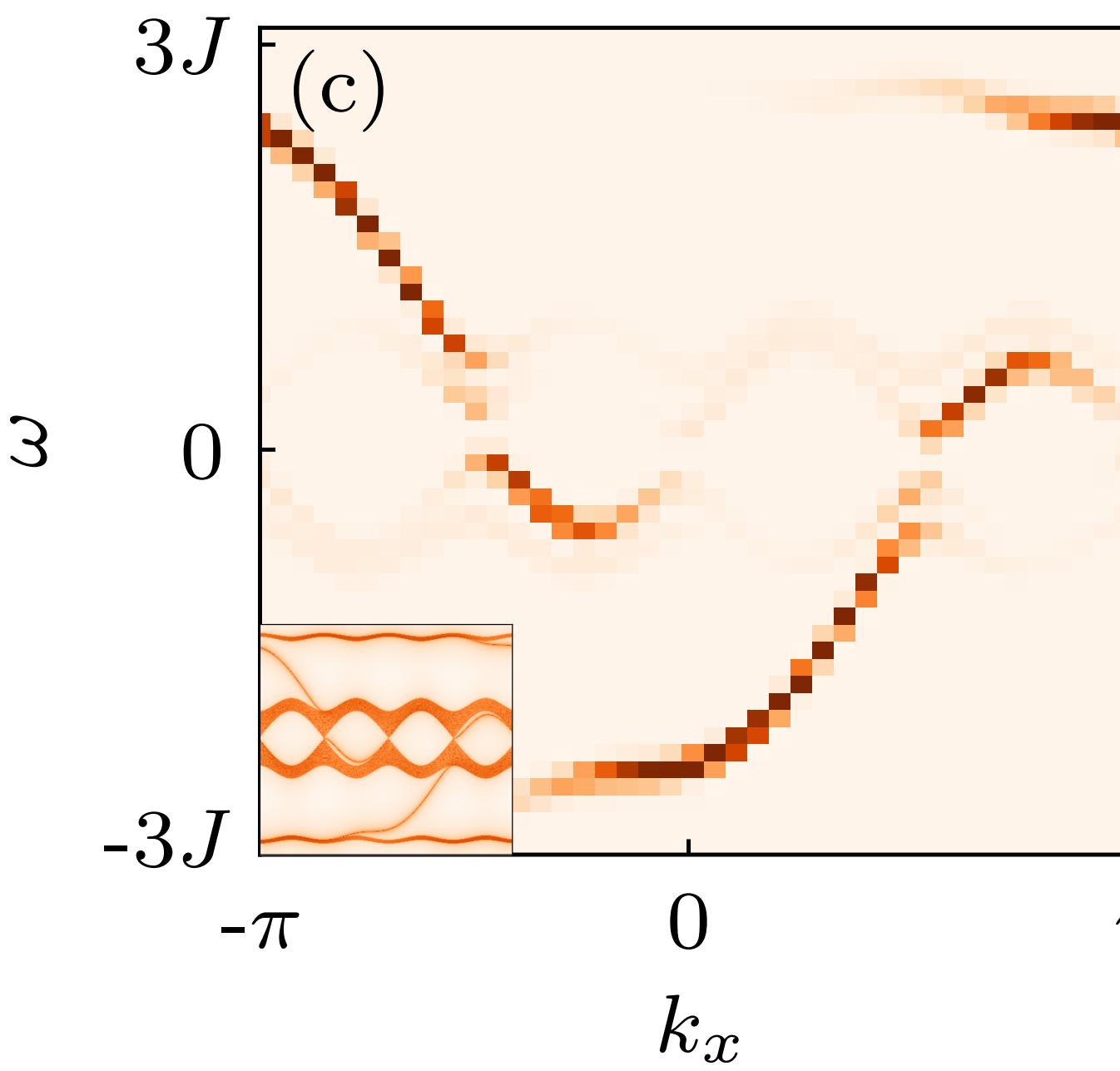
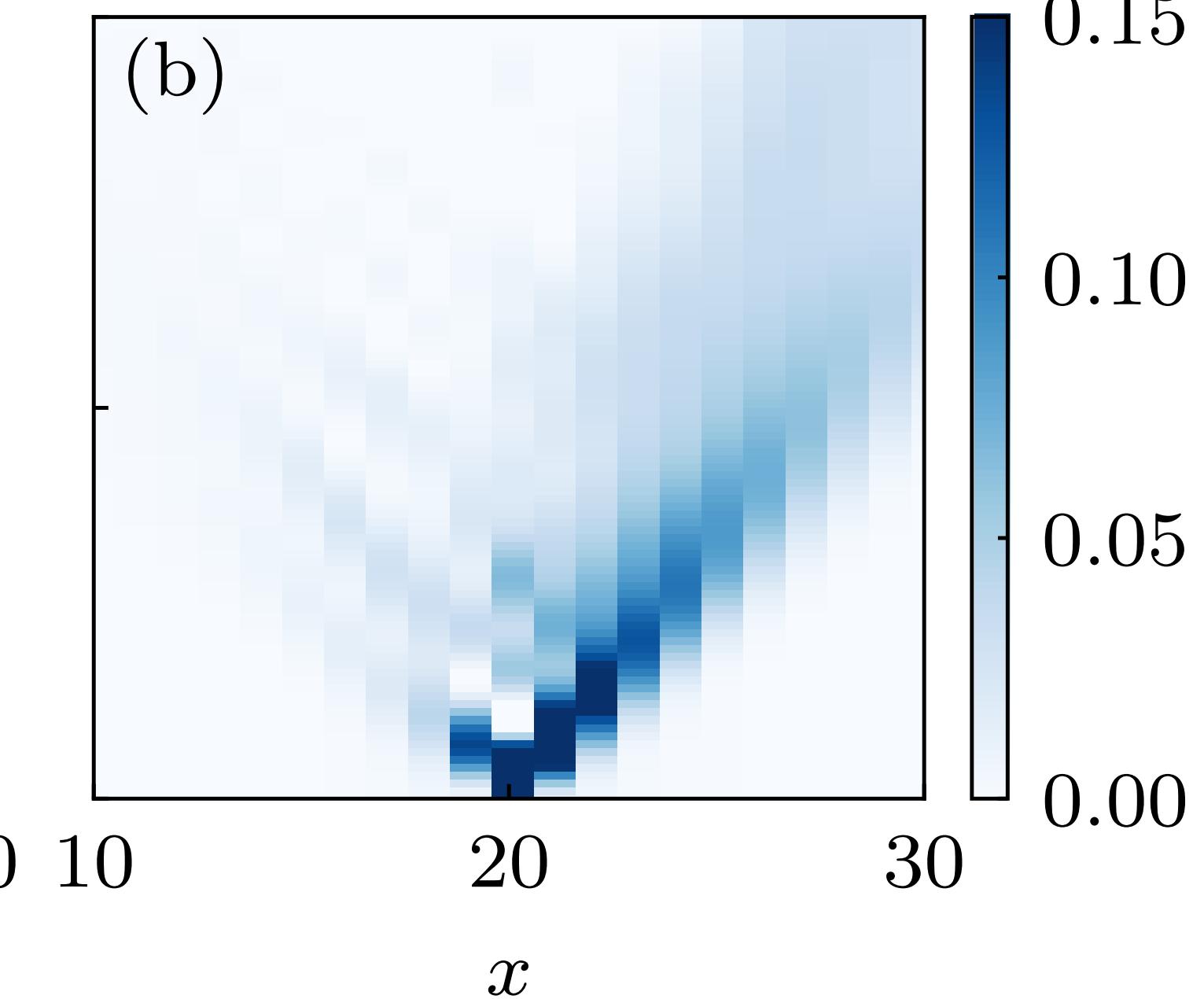
$$G(x, t) = \langle b(x, t) b^\dagger(0, 0) \rangle$$



Single particle

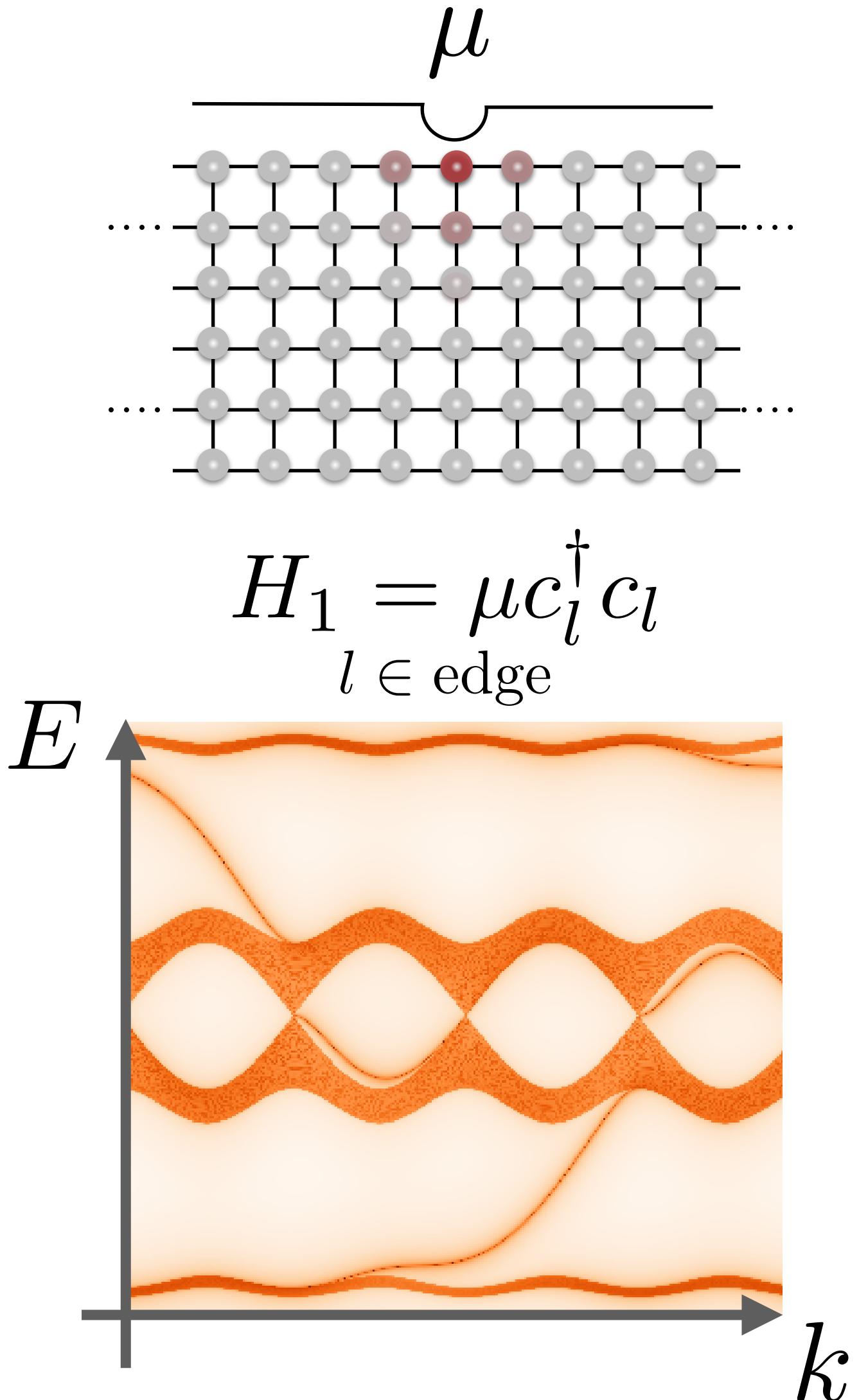


FCI



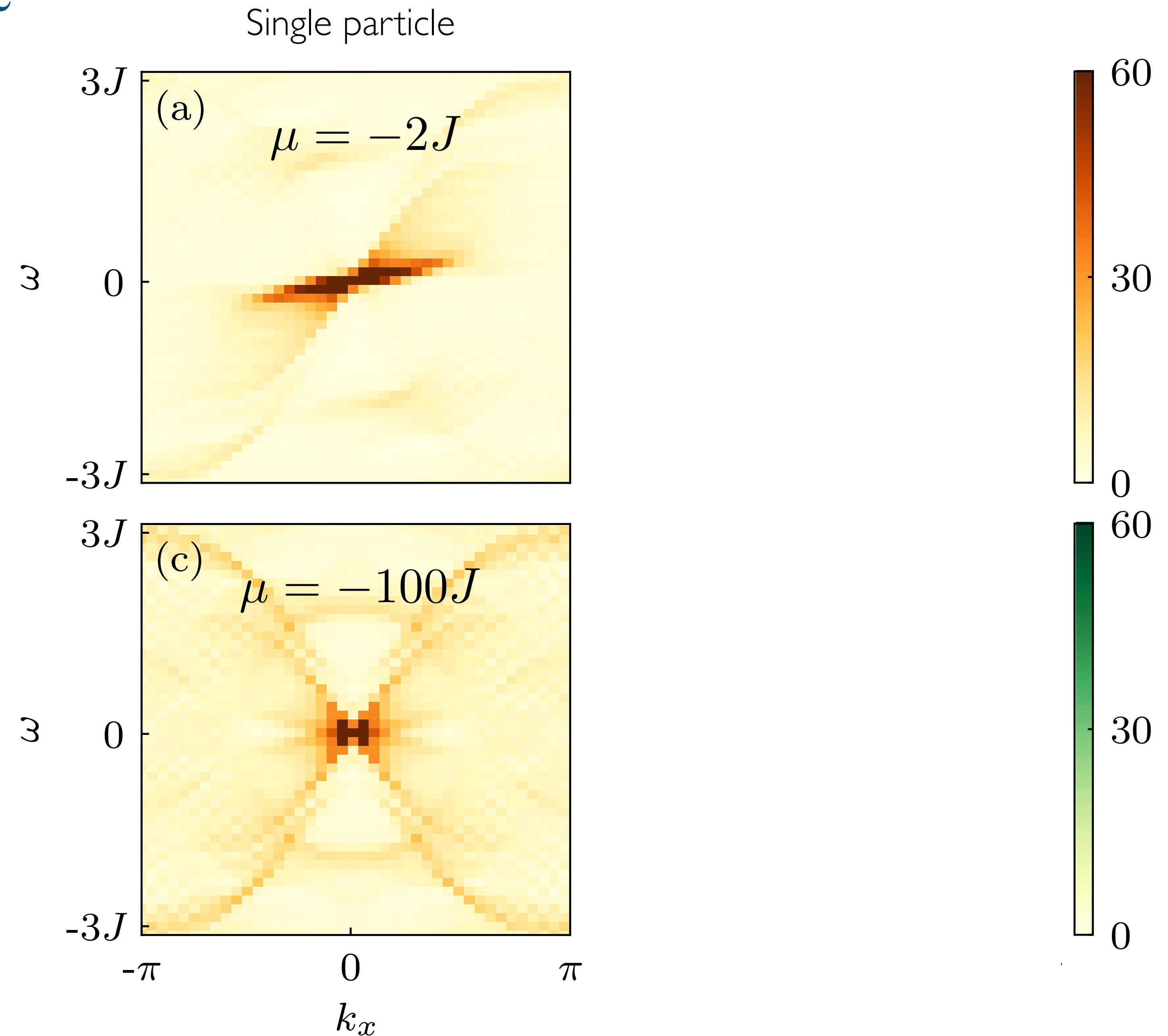
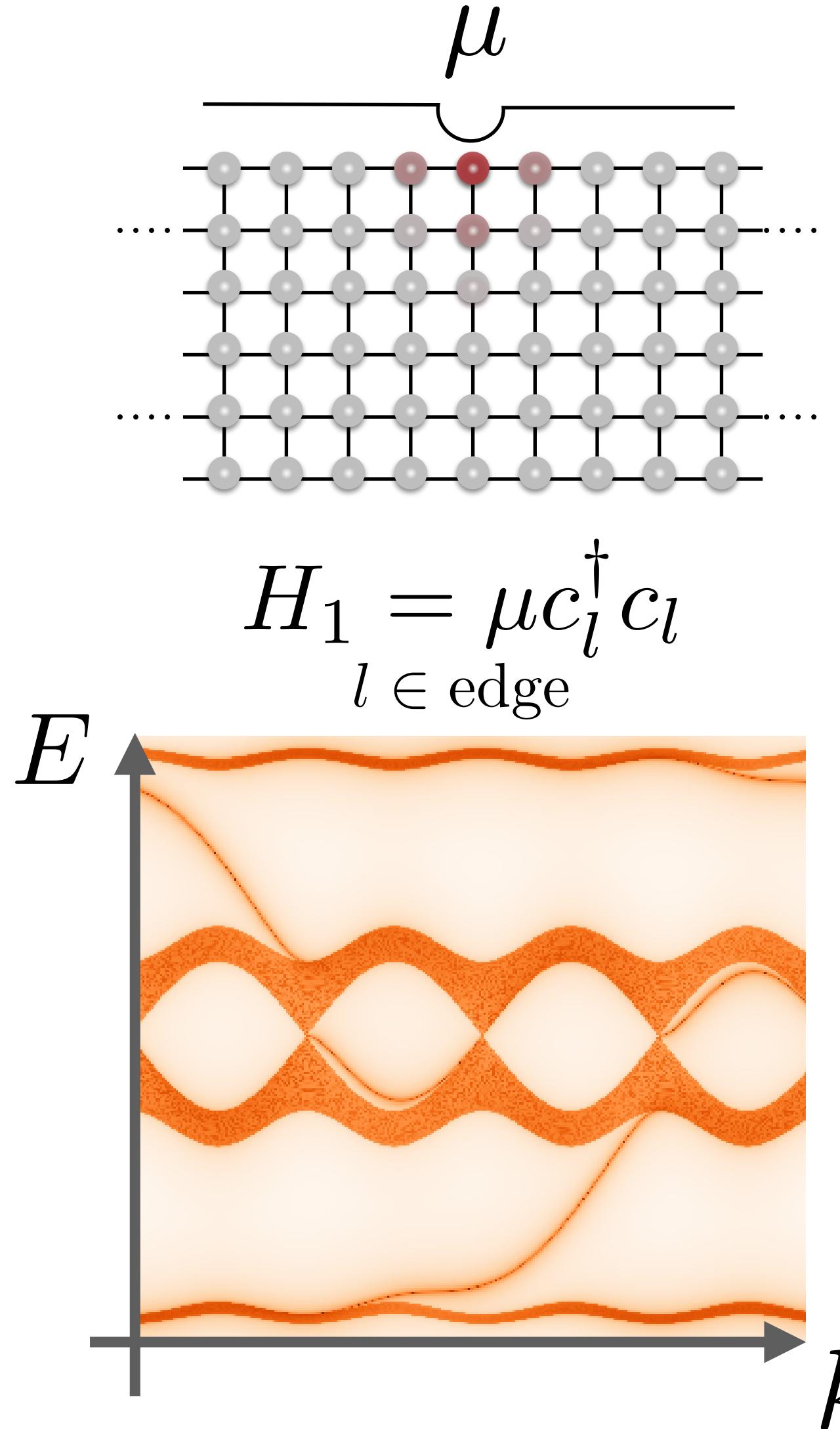
Dynamical signatures of the FCI phase

Fourier transformation of the density evolution
following local quench:



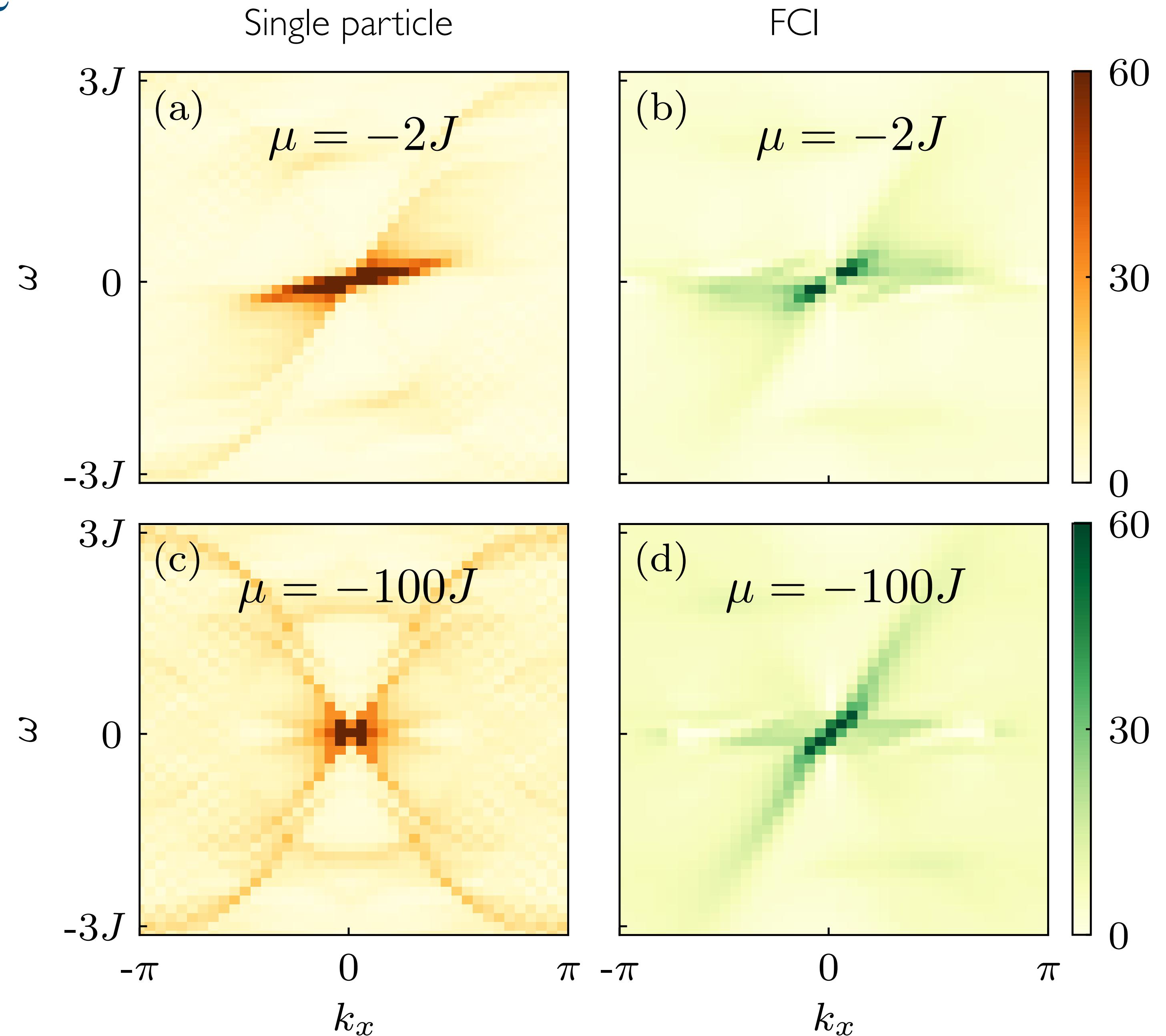
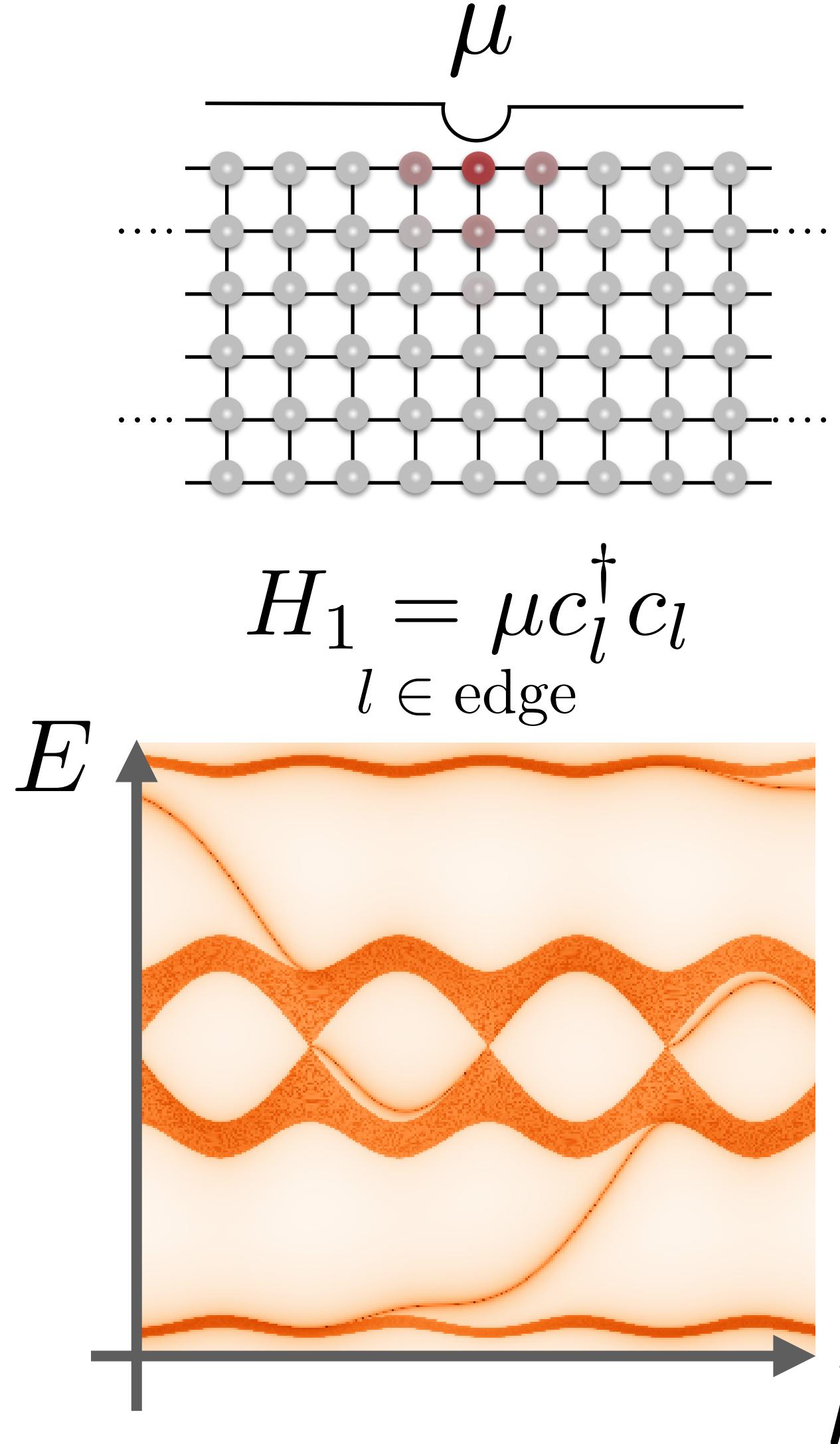
Dynamical signatures of the FCI phase

Fourier transformation of the density evolution
following local quench:



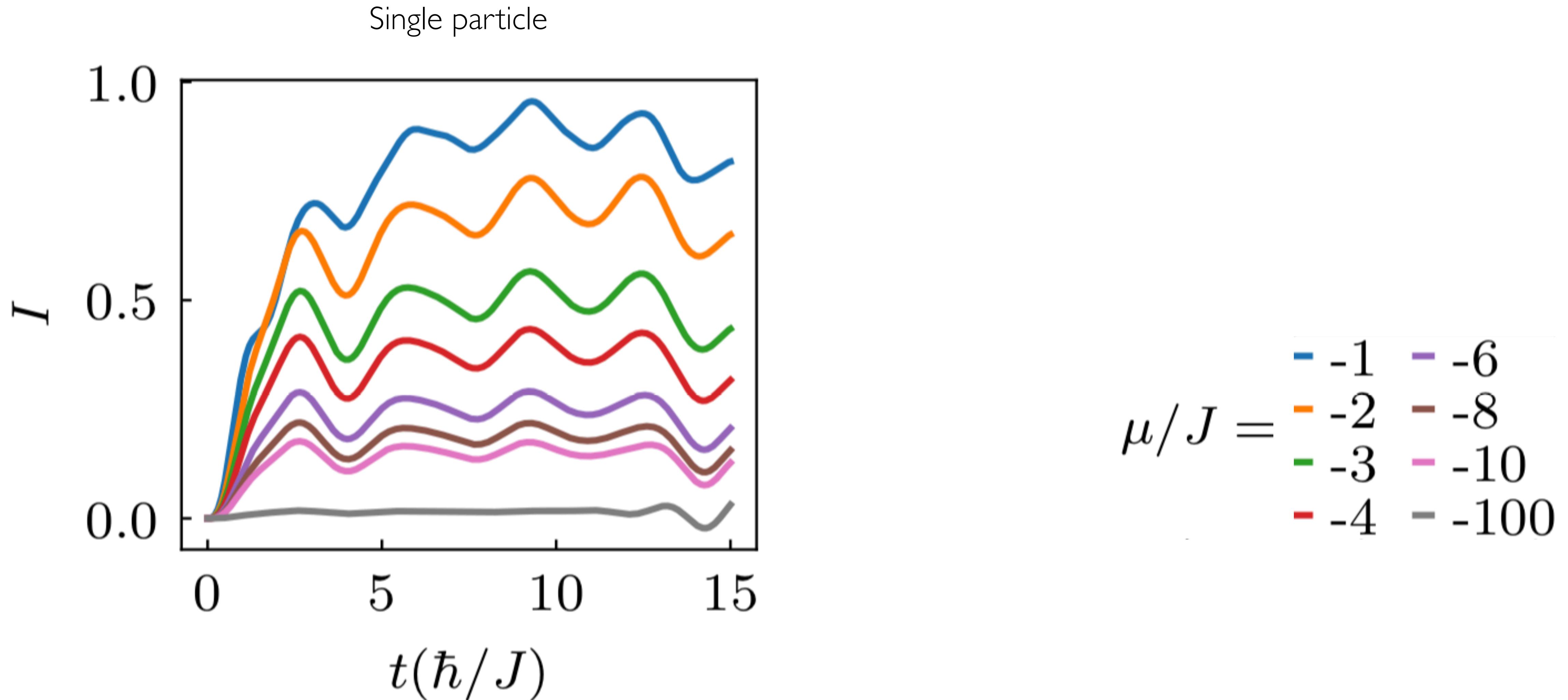
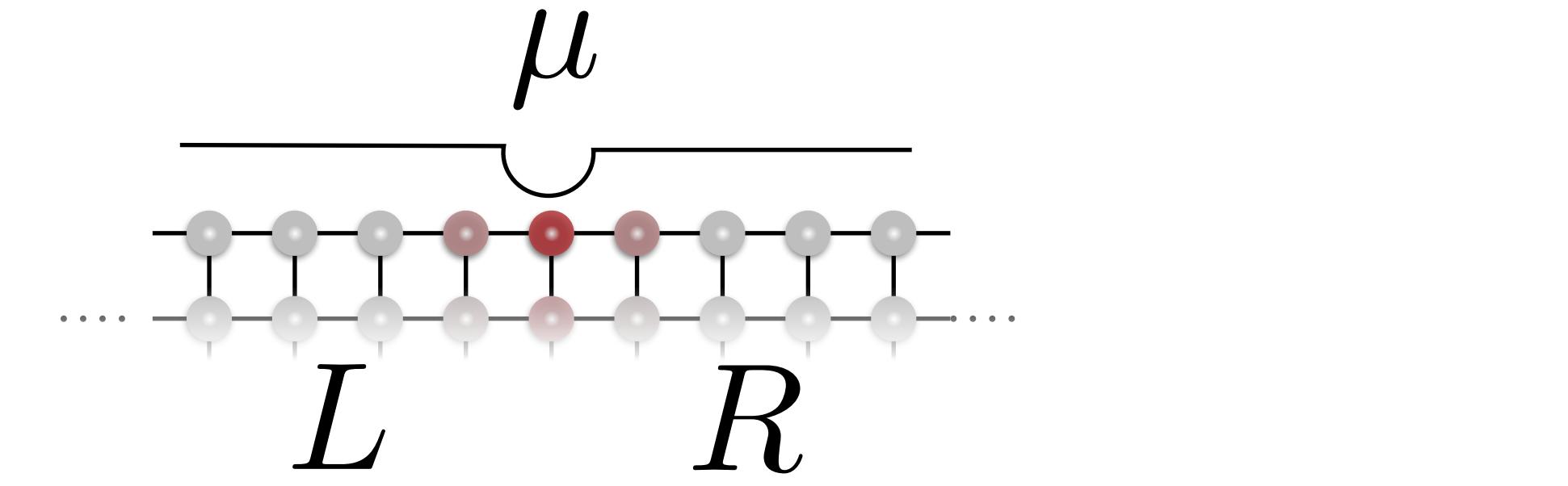
Dynamical signatures of the FCI phase

Fourier transformation of the density evolution
following local quench:



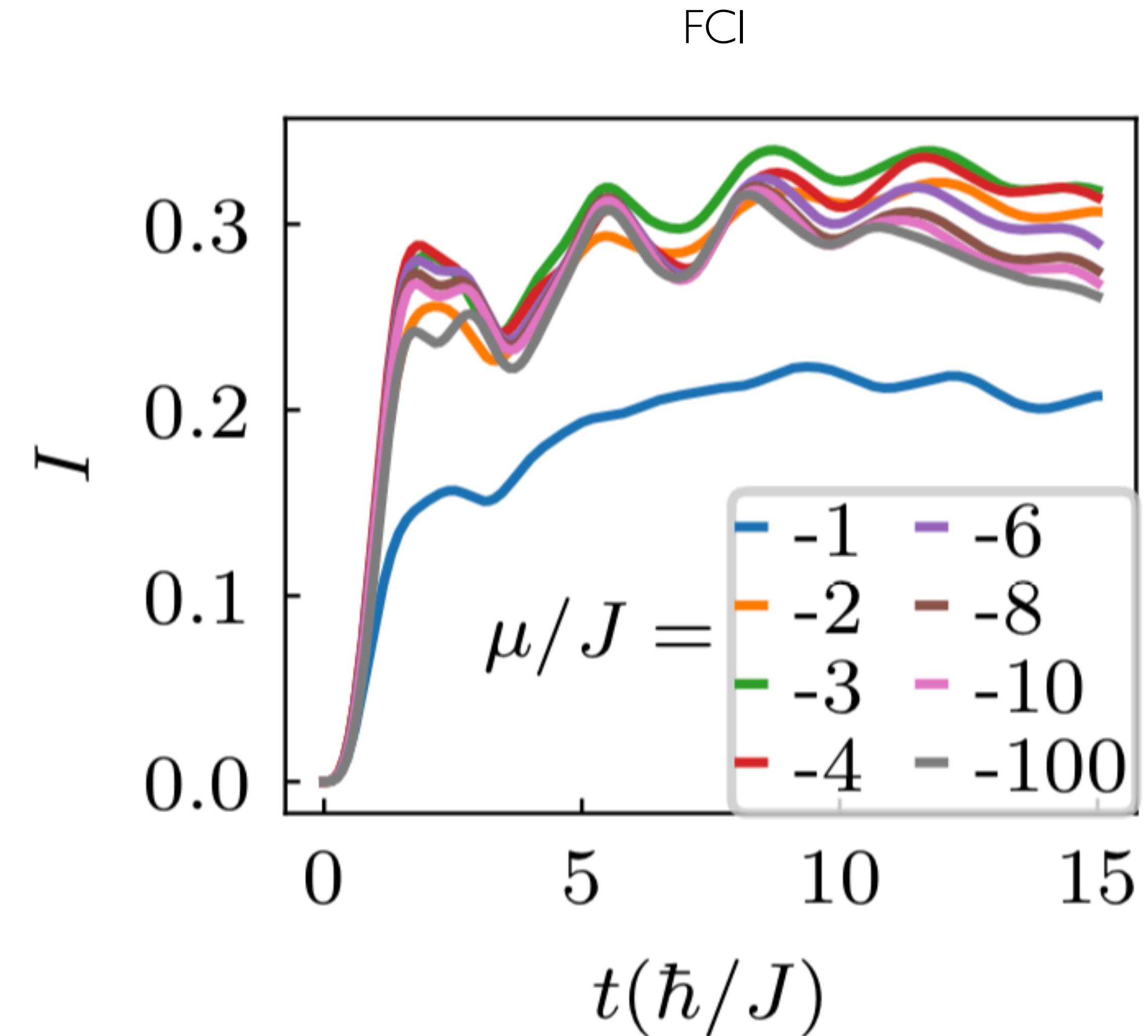
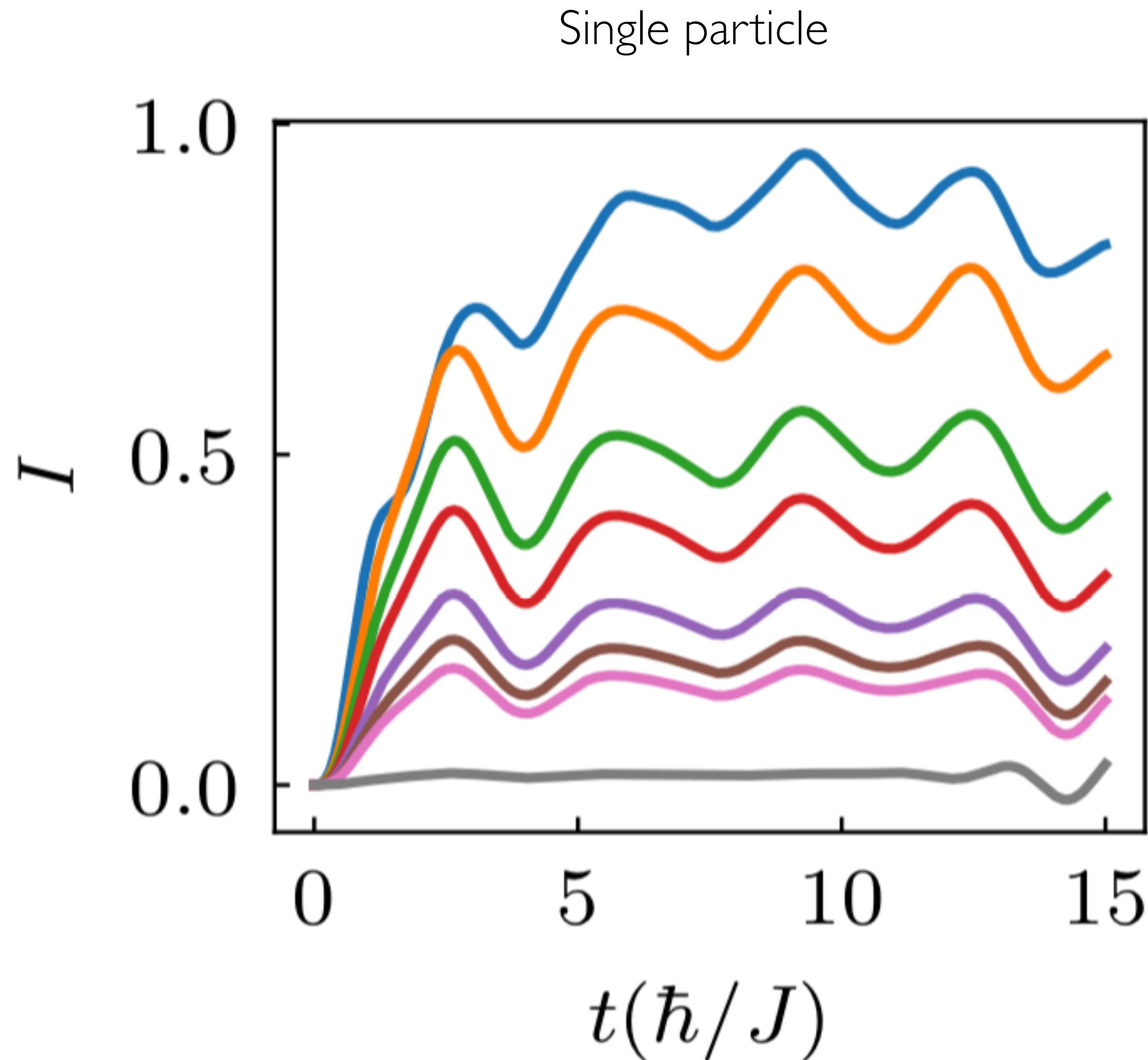
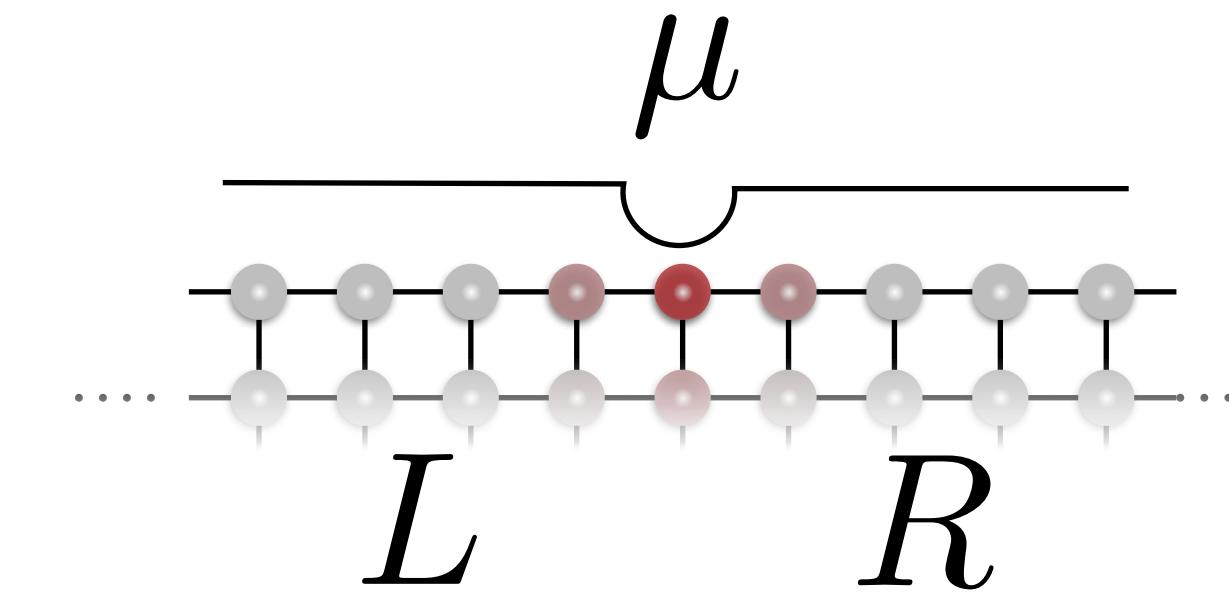
Dynamical signatures of the FCI phase

The time evolution of the imbalance : $I = N_R - N_L$



Dynamical signatures of the FCI phase

The time evolution of the imbalance : $I = N_R - N_L$



Edge state dynamics of a quantum Hall edge



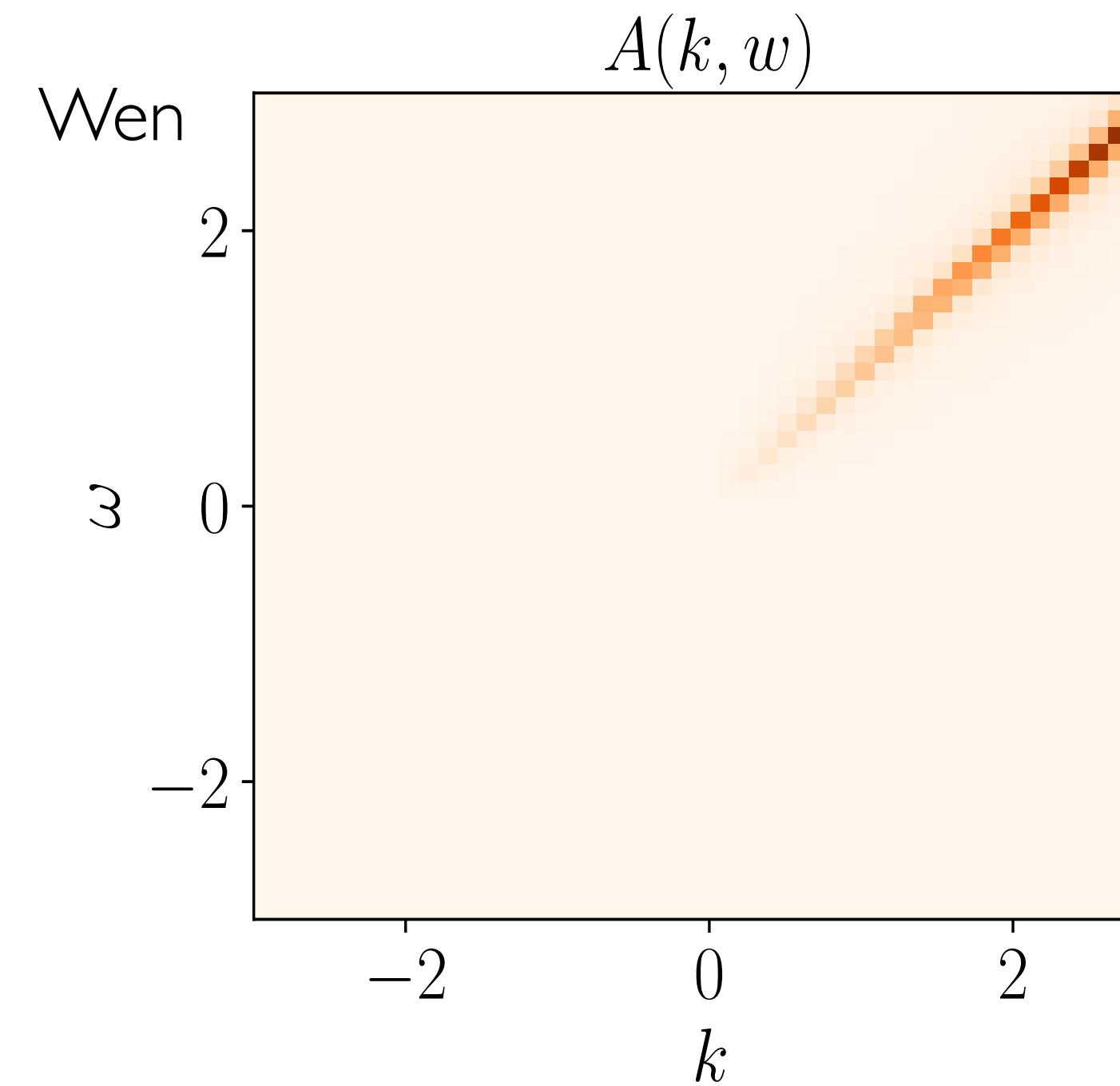
X. G. Wen PRB (1990)

Edge state spectral function knows about fractional excitations

For a state with $\sigma_{xy} = \frac{1}{m} \frac{e^2}{h}$

The spectral function is* $A(k, \omega) \propto (\omega + vk)^{m-1} \delta(\omega - vk)$
and the DOS is $N(\omega) \propto \omega^{m-1}$

*Assumptions: thermodynamic limit of a 1D isolated edge



Edge state dynamics of a quantum Hall edge



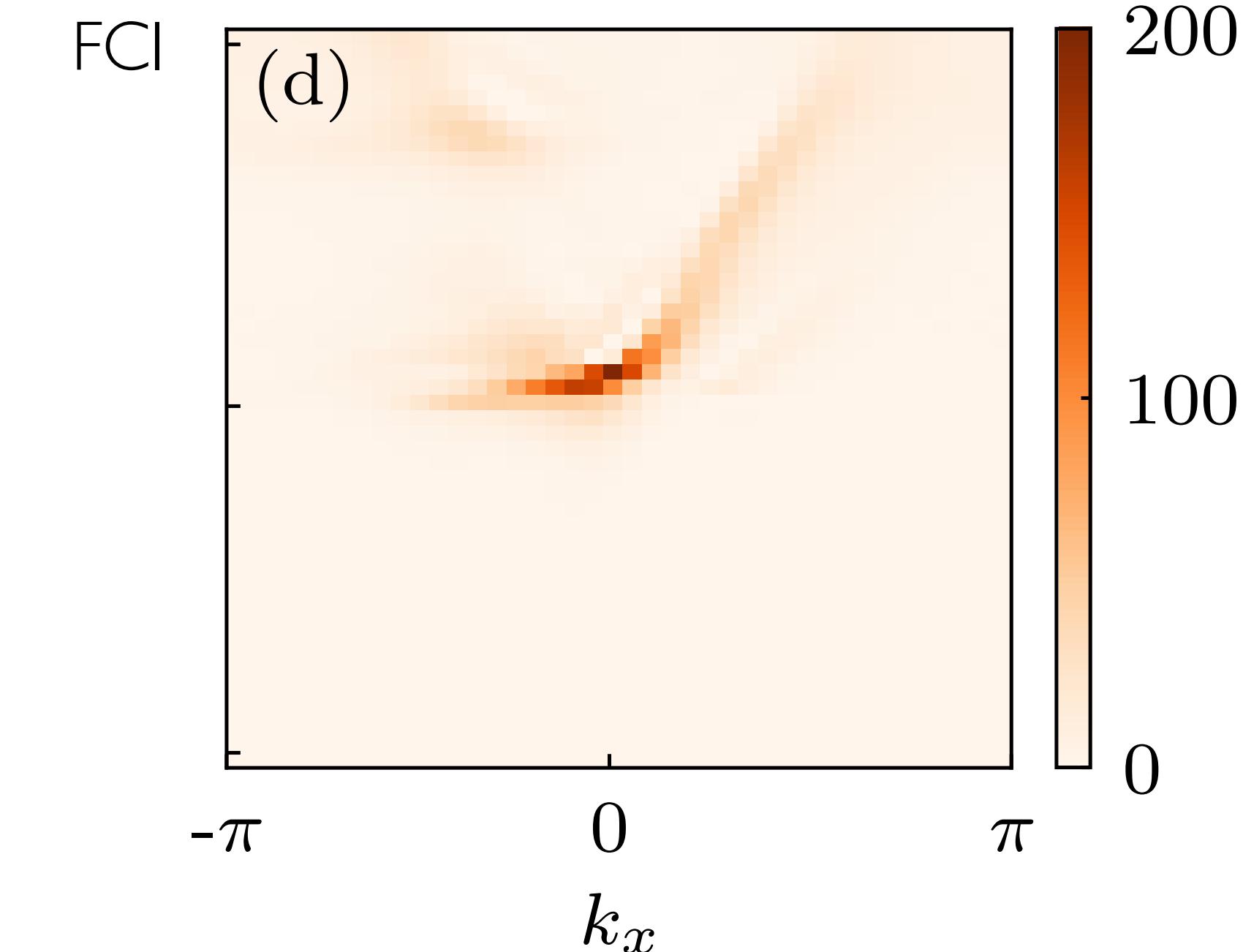
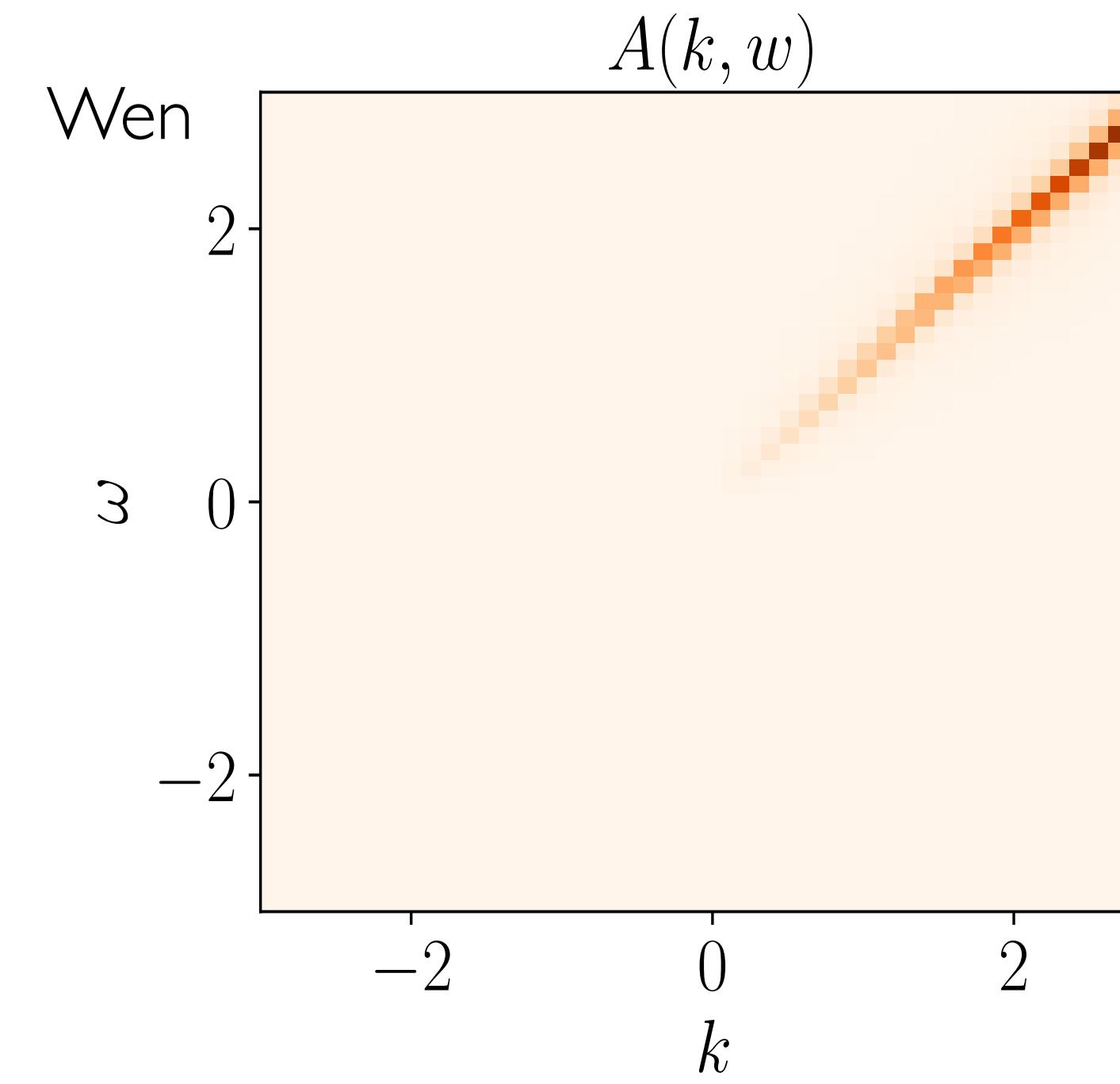
X. G. Wen PRB (1990)

Edge state spectral function knows about fractional excitations

For a state with $\sigma_{xy} = \frac{1}{m} \frac{e^2}{h}$

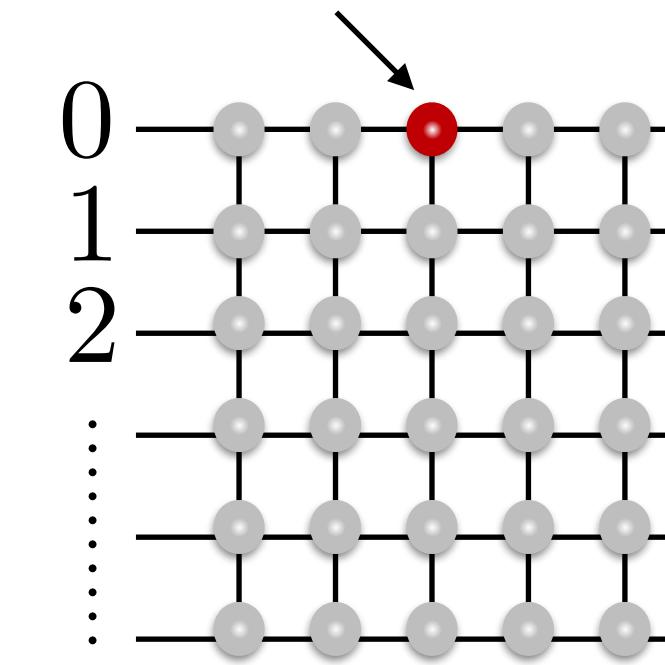
The spectral function is* $A(k, \omega) \propto (\omega + vk)^{m-1} \delta(\omega - vk)$
and the DOS is $N(\omega) \propto \omega^{m-1}$

*Assumptions: thermodynamic limit of a 1D isolated edge

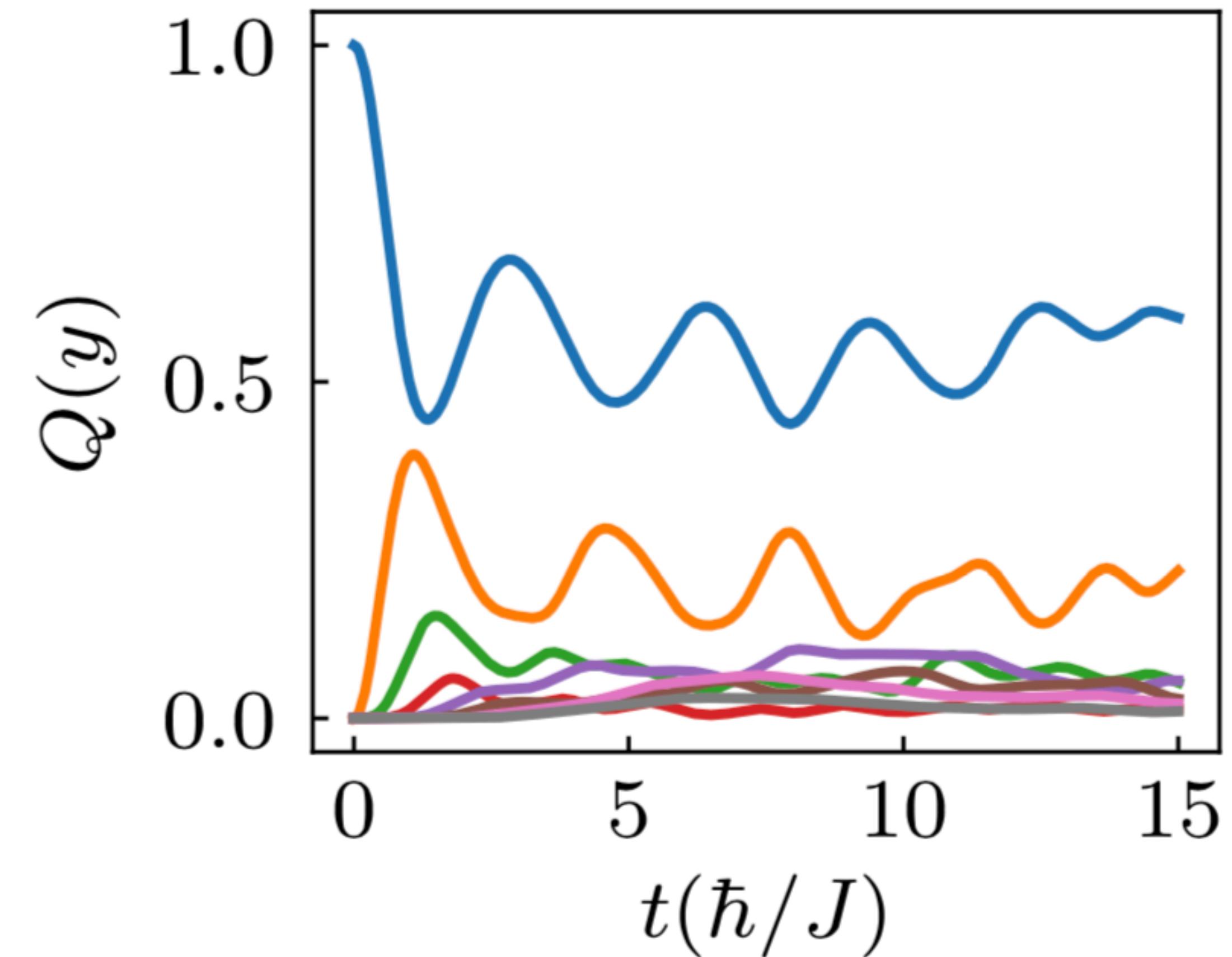


Dynamical signatures of the FCI phase

The edge does not behave as an isolated Luttinger liquid:

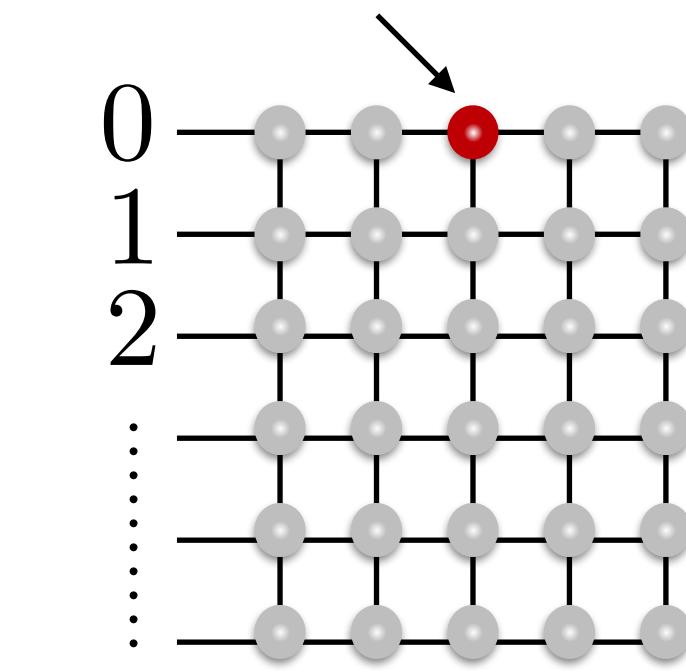


Single particle

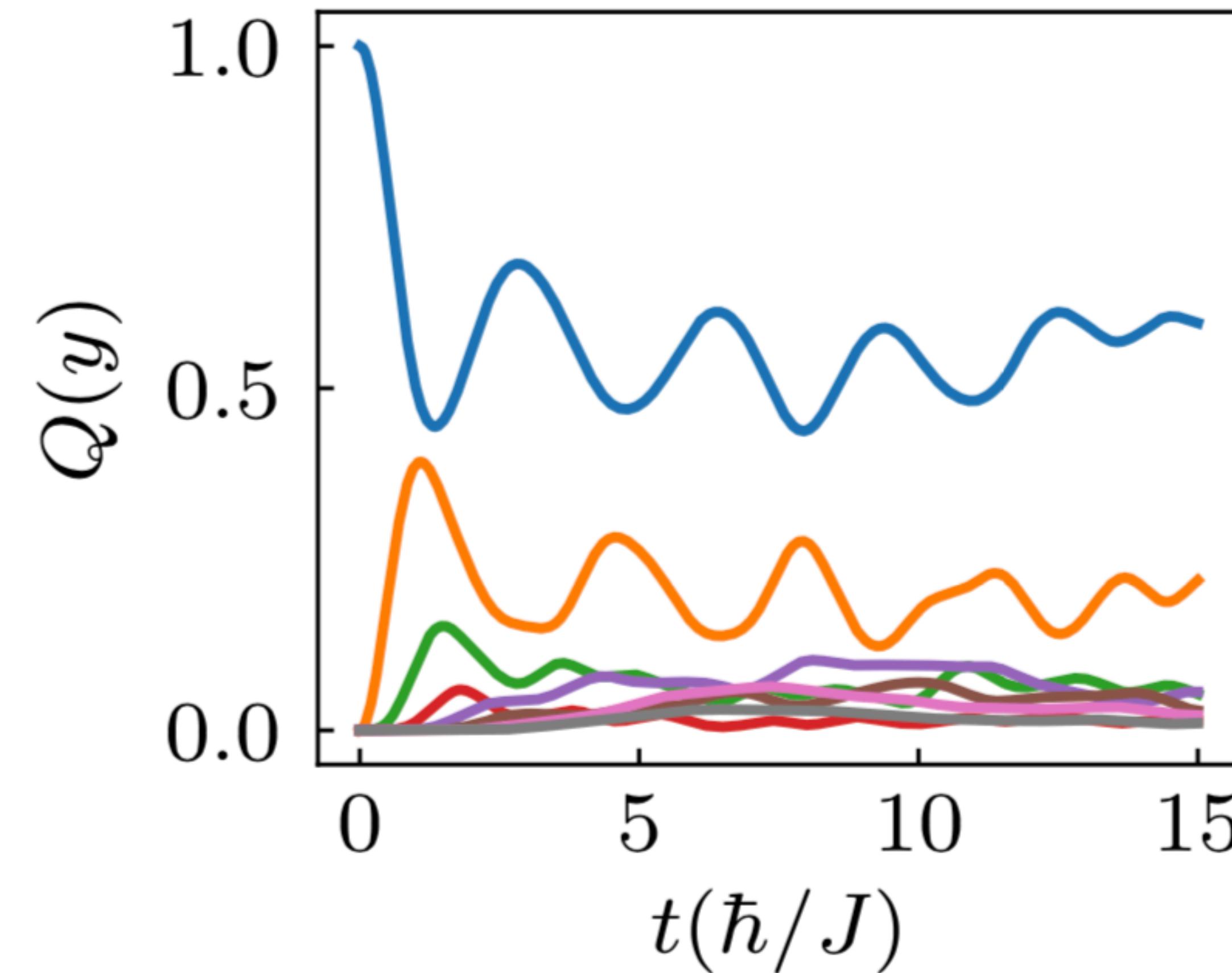


Dynamical signatures of the FCI phase

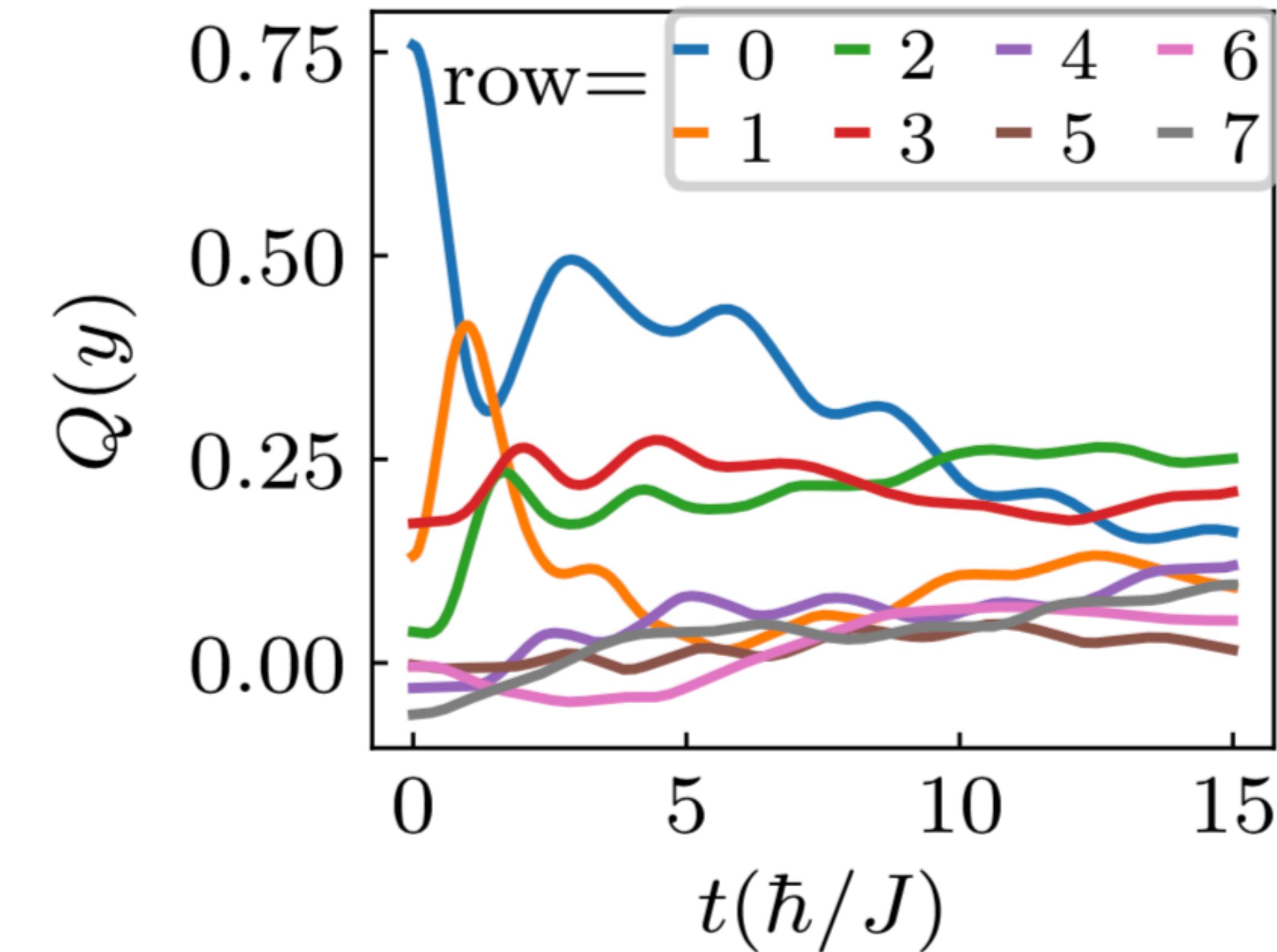
The edge does not behave as an isolated Luttinger liquid:



Single particle



FCI



Harper-Hofstadter

Edge state dynamics of a bosonic fractional Chern insulator

is chiral but not that of an isolated Luttinger liquid.

It is insensitive to the strength of a perturbation, unlike the Chern insulator



X. Y. Dong, AGG, J. Motruk, F.
Pollmann, Phys Rev. Lett. (2018)



Dynamical signatures of the Chern insulator

