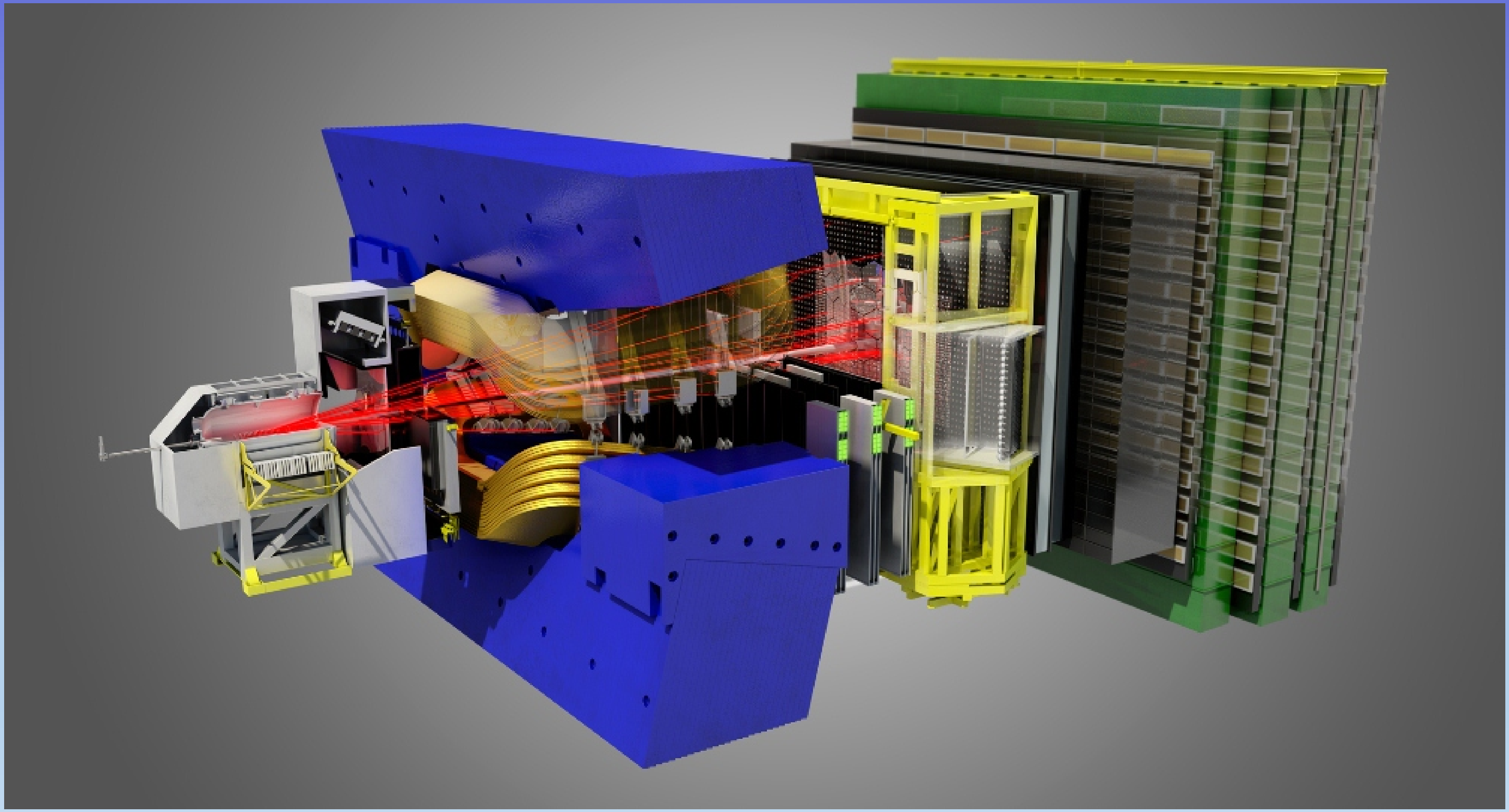
TAE 2019, 12th-13th September, Benasque, Spain

LHCb Physics



Arantza Oyanguren (IFIC – CSIC/UV)

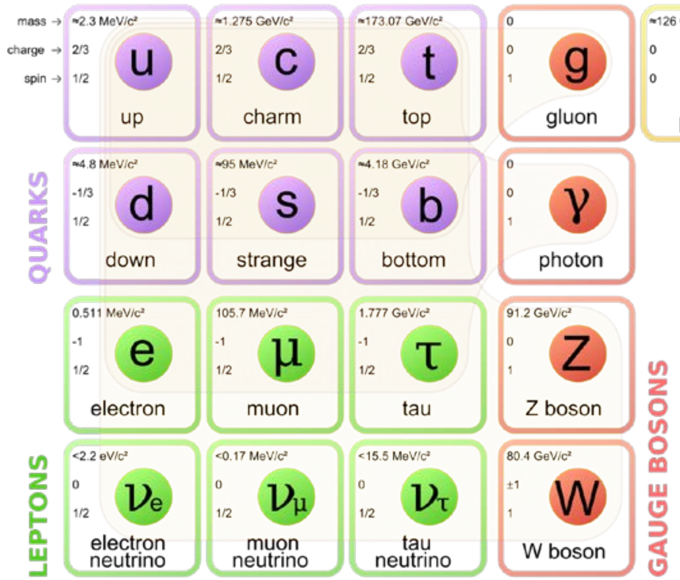

LHCb
ATLAS
CERN Meyrin
CERN Prévessin
SPS 7 km
SUISSE
FRANCE
ALICE
CMS
LHC 27 km

Outline

- Introduction
- The LHCb experiment
- The CKM matrix
- B mixing and CP violation
- Rare decays of B mesons
- Future plans

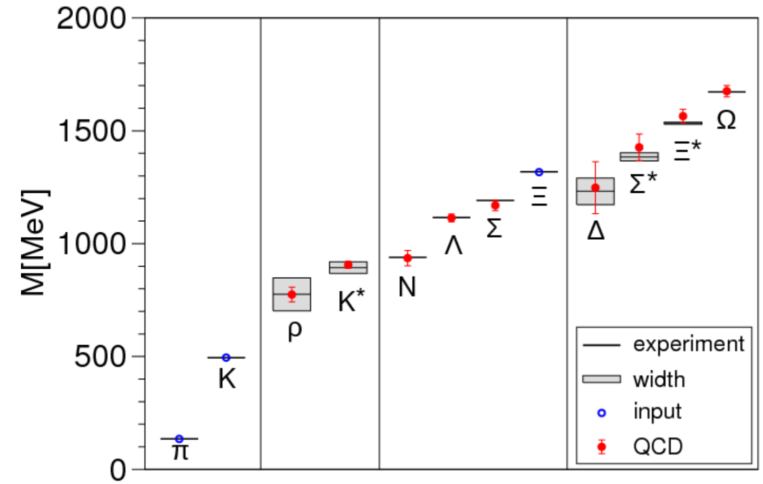
Introduction

Our Standard Model of Particle Physics:

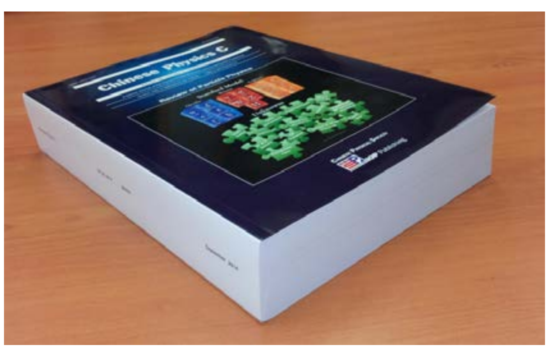


+ antiparticles

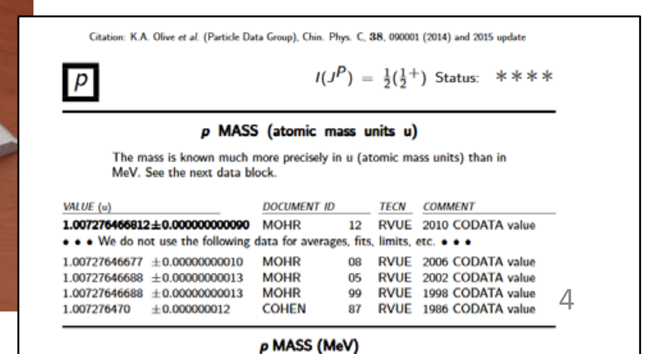
Hadrons:



Particle Data Book (PDG):



1675 pages!!



Citation: K.A. Olive *et al.* (Particle Data Group), Chin. Phys. C, **38**, 090001 (2014) and 2015 update

$\boxed{p}$ $I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$ Status: ∗∗∗∗

p MASS (atomic mass units u)

The mass is known much more precisely in u (atomic mass units) than in MeV. See the next data block.

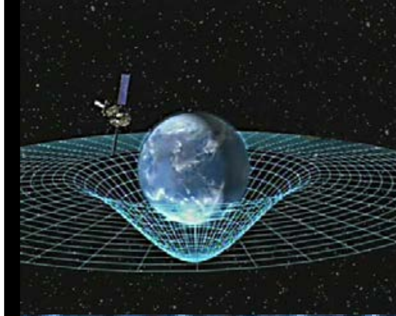
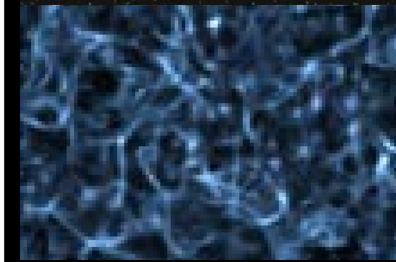
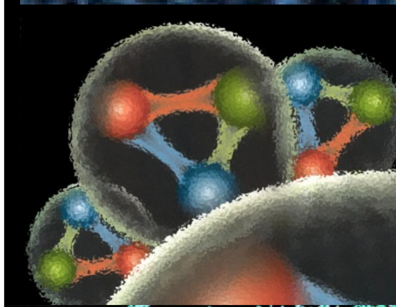
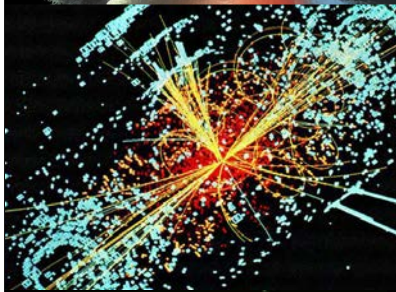
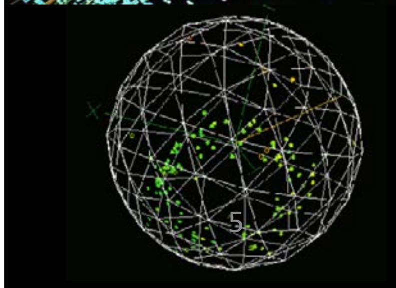
VALUE (u)	*DOCUMENT ID*		*TECN*	*COMMENT*
1.007276466812±0.000000000090	MOHR	12	RVUE	2010 CODATA value
• • • We do not use the following data for averages, fits, limits, etc. • • •				
1.00727646677 ±0.00000000010	MOHR	08	RVUE	2006 CODATA value
1.00727646688 ±0.00000000013	MOHR	05	RVUE	2002 CODATA value
1.00727646688 ±0.00000000013	MOHR	99	RVUE	1998 CODATA value
1.007276470 ±0.000000012	COHEN	87	RVUE	1986 CODATA value

p MASS (MeV)

Introduction

The problems of our Standard Model …

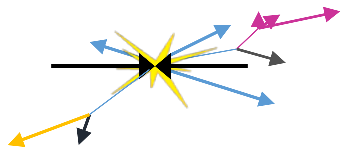
- Quantum Theory of Gravity
- Inflation?
- Quark/lepton generation masses: compositeness?
  Substructure? Strings?
  Common sub-elements quarks and leptons?
  Why three families?
- Matter-Antimatter asymmetry
  CPV in SM (K, B) + Big Bang ?
- Cosmological constant (dark energy …)
- Dark matter
- Higgs & EW symmetry breaking? Forces Unification?
- Neutrinos (mass?, hierarchy?...)

Introduction

Looking for New Physics…

Direct searches:

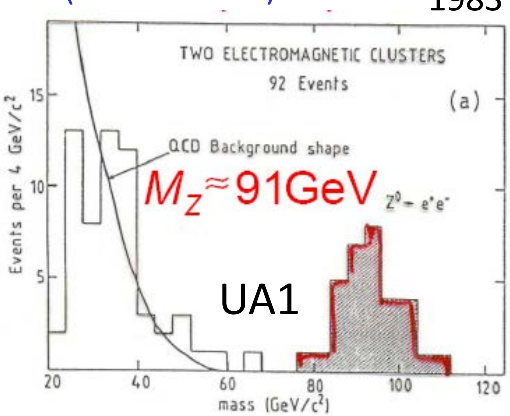


$$E^2 = M^2c^4 + p^2c^2$$

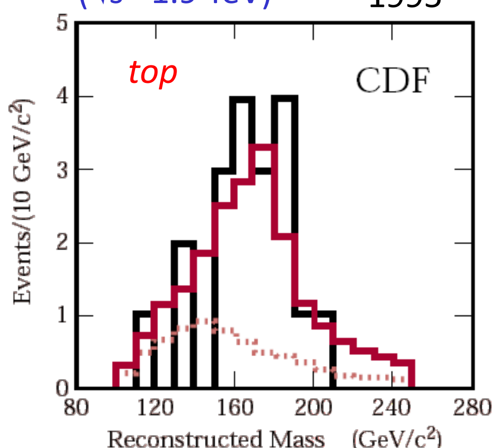
High energy

$\rightarrow$ **particles created *on-shell:* *Evidence in mass plots***

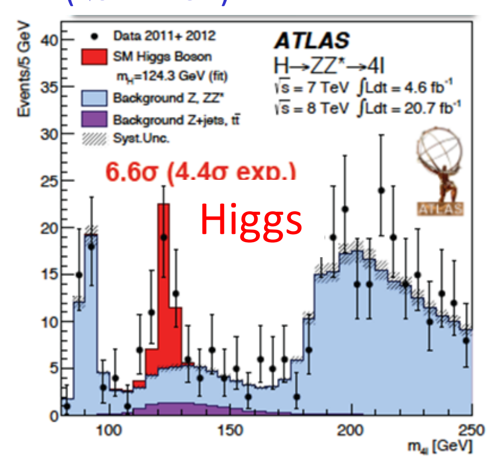
($\sqrt{s}$= 546 GeV) 1983



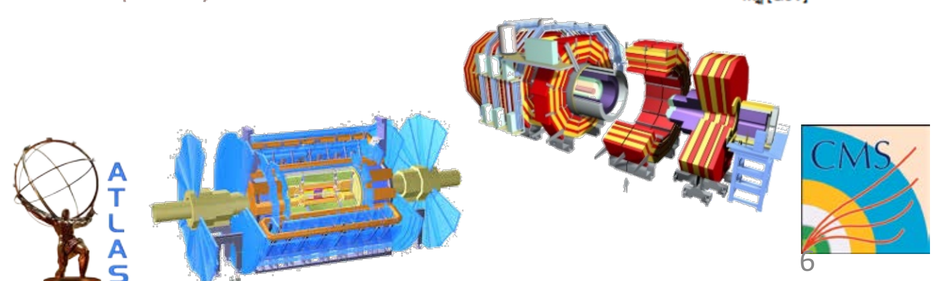
($\sqrt{s}$= 1.9 TeV) 1995



($\sqrt{s}$= 7 TeV) 2012



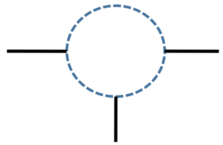
Higgs discovery, 2012



Introduction

Looking for New Physics…

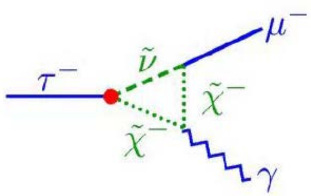
Indirect searches:

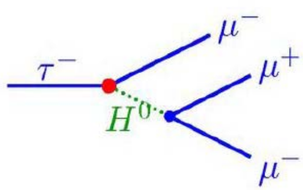


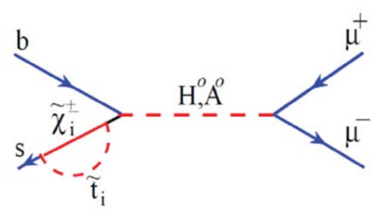
High precision

$\rightarrow$ **particles created *off-Shell: Evidence in quantum effects (loops)***

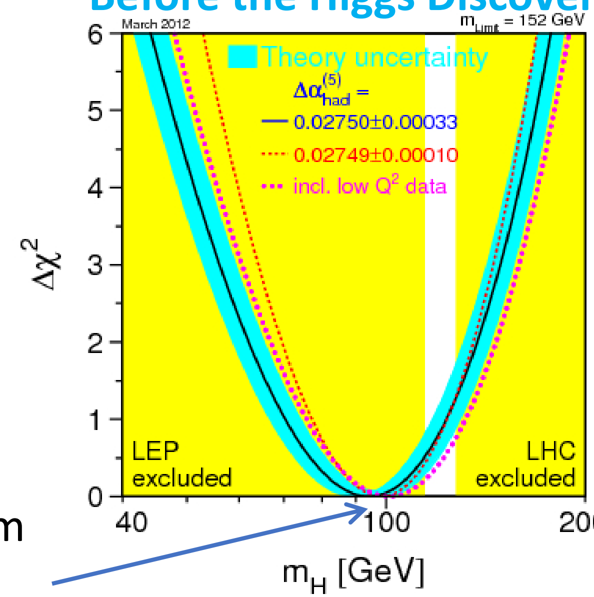
(BR's, asymmetries…)

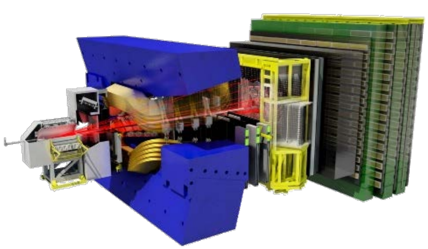






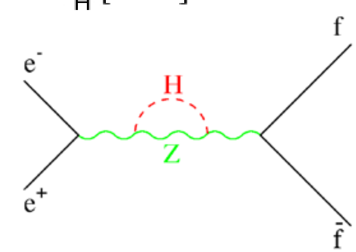
Before the Higgs Discovery…





LHCb

Predicted from electroweak measurements



Direct and indirect searches are complementary

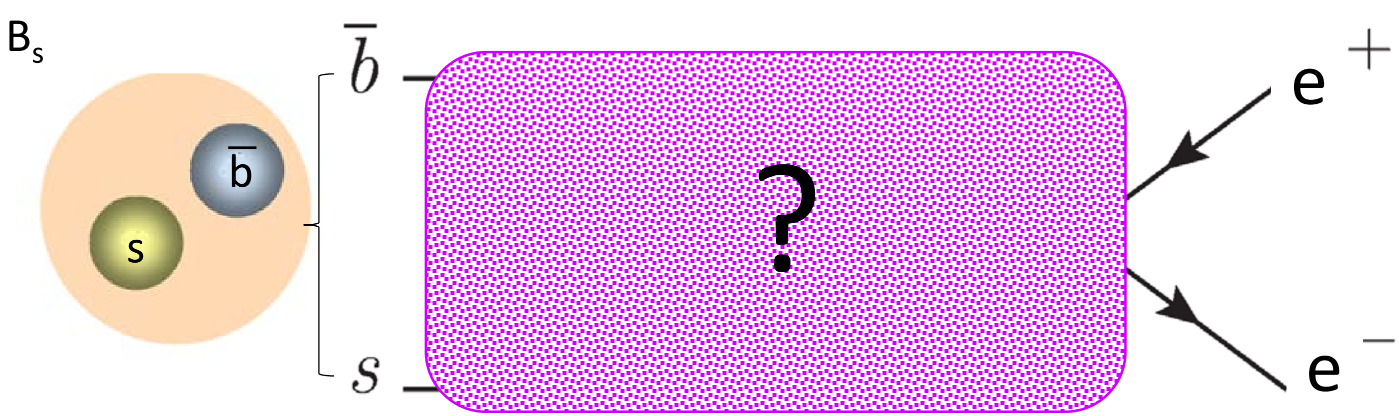
Introduction

What we see

What we think it is

What it is

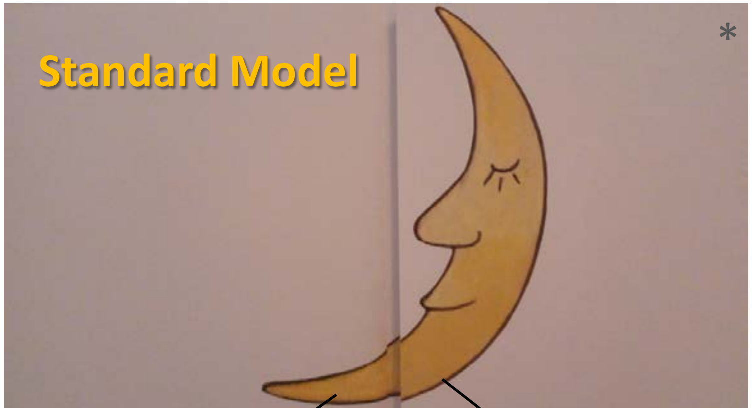
It can be tested by studying quantum effects:



Number of particles produced,
angular distributions,
origin point in the detector…

* ¡Oh!, Josse Goffin

Introduction

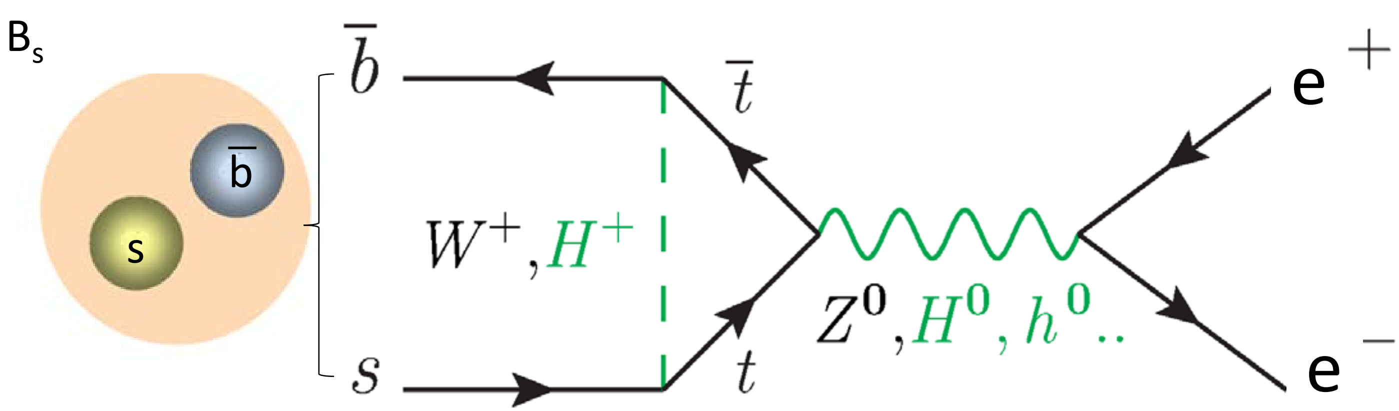


What we see

What we think it is

What it is

It can be tested by studying quantum effects:

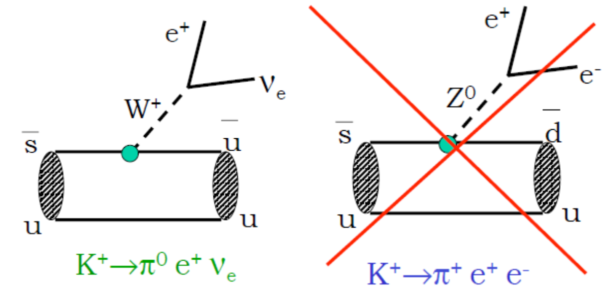


Number of particles produced, angular distributions, origin point in the detector…

New Particles

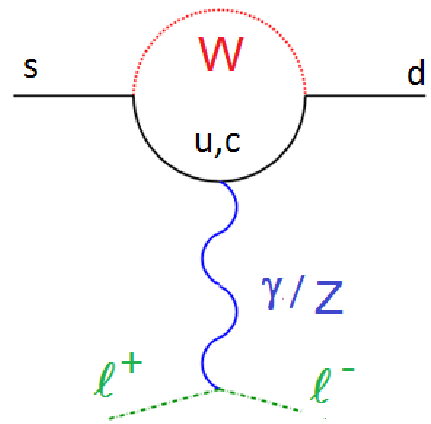
* ¡Oh!, Josse Goffin

Introduction

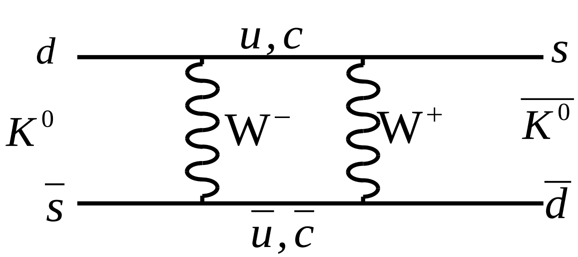


- The GIM mechanism:

In 1970's Glashow, Iliopoulos and Maini described the mechanism by which flavour-changing neutral currents (FCNCs) are suppressed, and predicted the existence of the c quark

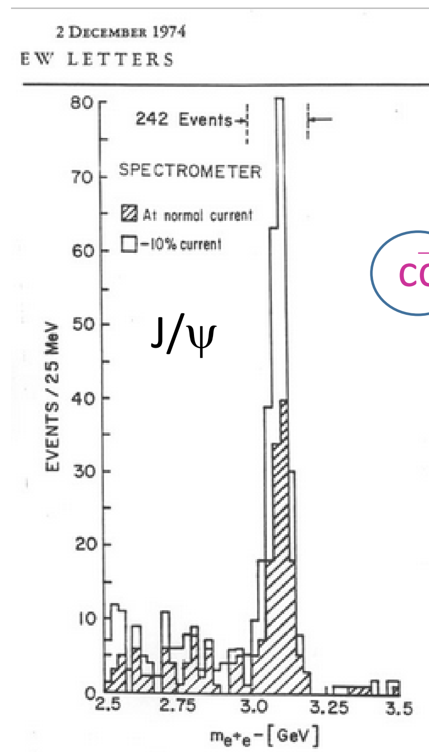


- Gaillard, Lee and Rosner :
  m_c~1.5 GeV from kaon mixing



$$\Delta m_K = \frac{G_F^2}{4\pi} m_K f_K^2 m_c^2 \cos^2\theta_c \sin^2\theta_c$$

- 1974 c quark discovered
  (B. Richter at SLAC
  and S. Ting at BNL)

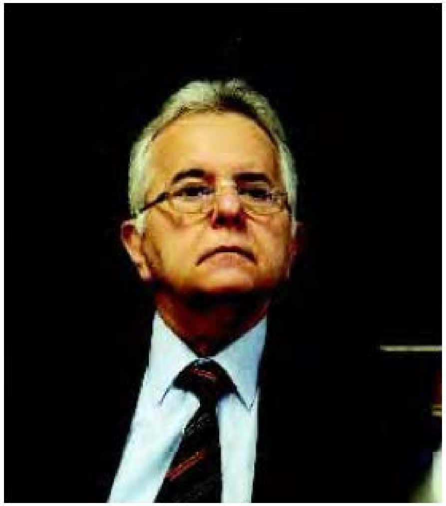


IG. 2. Mass spectrum showing the existence of J.
sults from two spectrometer settings are plotted
wing that the peak is independent of spectrometer
rents. The run at reduced current was taken two
nths later than the normal run.

cc̄

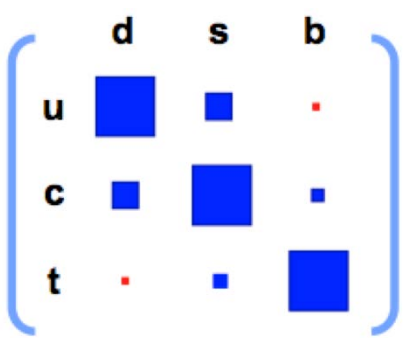
Introduction

- The CKM mechanism:



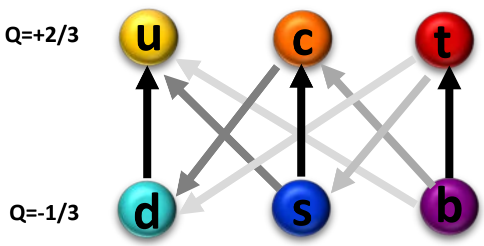

Cabibbo Kobayashi Maskawa

$V_{CKM} =$

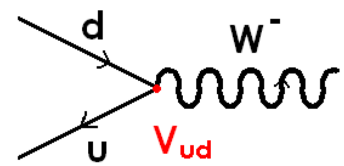


Introduction

- In the Standard Model of Particle Physics, transitions between different quarks are governed by the CKM mechanism:

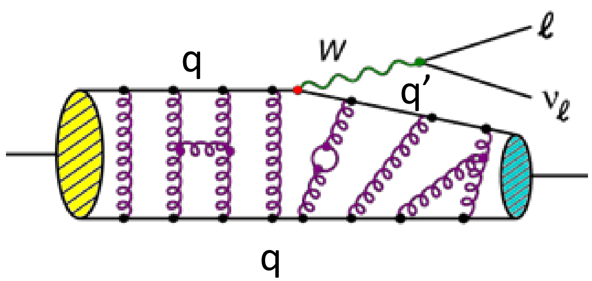


$$V_{\mathrm{CKM}} = \begin{pmatrix} V_{\mathrm{ud}} & V_{\mathrm{us}} & V_{\mathrm{ub}} \\ V_{\mathrm{cd}} & V_{\mathrm{cs}} & V_{\mathrm{cb}} \\ V_{\mathrm{td}} & V_{\mathrm{ts}} & V_{\mathrm{tb}} \end{pmatrix}$$



- The amplitude of a hadron decay process can be described using Effective Field Theories: Operator Product Expansion (OPE)

$$A(M \to F) = \langle F|\mathcal{H}_{eff}|M\rangle = \frac{G_F}{\sqrt{2}} \sum_i V^i_{CKM} C_i(\mu) \langle F|O_i(\mu)|M\rangle$$



CKM couplings

Wilson Coefficients (μ = scale)

Hadronic Matrix Elements

Introduction

$$A(M \to F) = \langle F|\mathcal{H}_{eff}|M\rangle = \frac{G_F}{\sqrt{2}} \sum_i V^i_{CKM} C_i(\mu) \langle F|O_i(\mu)|M\rangle$$

CKM couplings

Wilson Coefficients (μ = scale)

Hadronic Matrix Elements

$\rightarrow$ OPE: a series of effective vertices multiplied by effective coupling constants C_i.

Electroweak scale ~ $1/M_W$

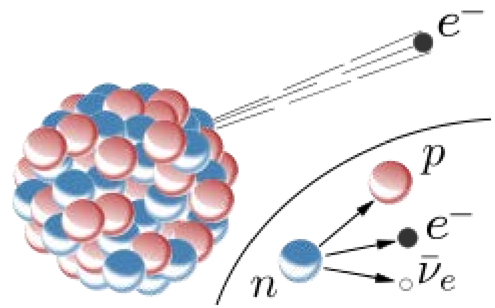
New Physics scale ~ $1/M_{NP}$

$C_i = C_i^{SM} + C_i^{NP}$

$C'_i = \cancel{C'^{SM}_i} + C'^{NP}_i$

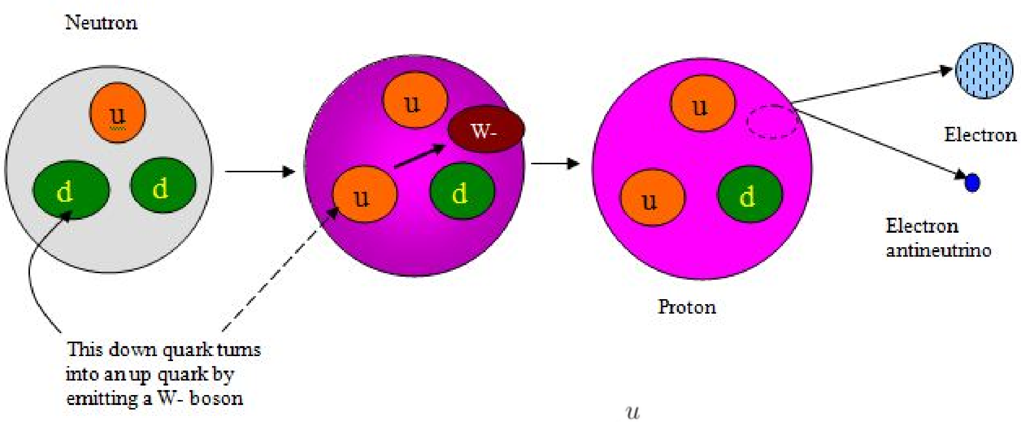
Primed $C'_i \rightarrow$ right handed currents: suppressed in SM

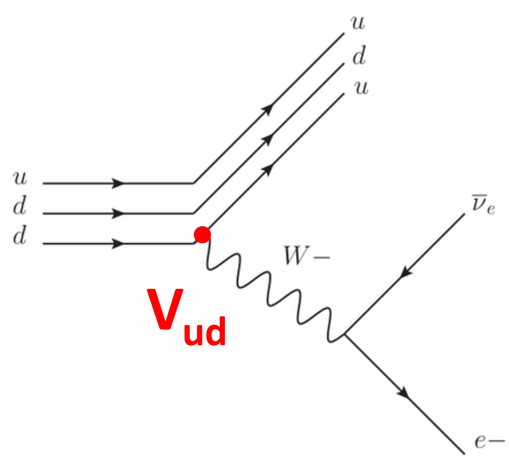
Introduction



Example: β -decay (very well known)

$$\boxed{{}^{14}_{6}C \rightarrow {}^{14}_{7}N + {}^{0}_{-1}\beta + \overline{\nu}_e}$$

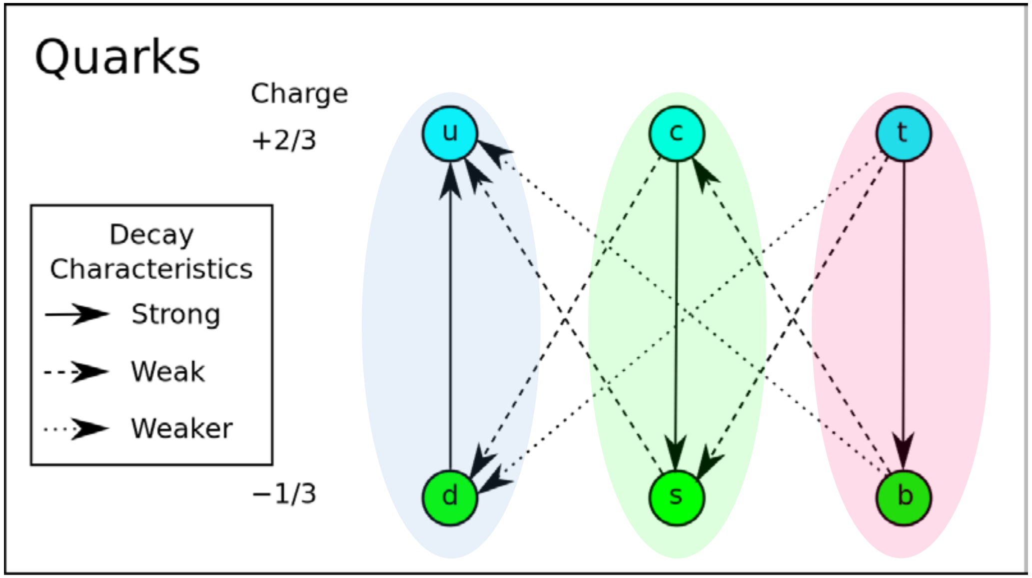


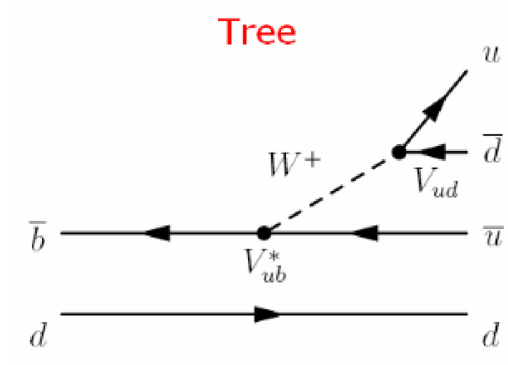


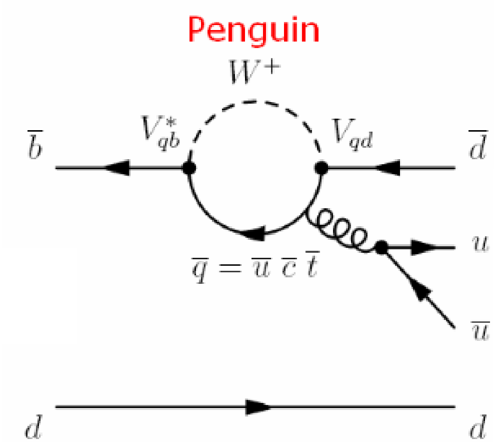
(V_{ud} is large ~ 0.97)

Introduction

CKM mechanism







$$V_{\mathrm{CKM}} = \begin{pmatrix} V_{\mathrm{ud}} & V_{\mathrm{us}} & V_{\mathrm{ub}} \\ V_{\mathrm{cd}} & V_{\mathrm{cs}} & V_{\mathrm{cb}} \\ V_{\mathrm{td}} & V_{\mathrm{ts}} & V_{\mathrm{tb}} \end{pmatrix} \approx \begin{bmatrix} 0.974 & 0.225 & 0.003 \\ 0.225 & 0.973 & 0.041 \\ 0.009 & 0.040 & 0.999 \end{bmatrix},$$

- Transitions between the same family are favored
- Some of them are very rare (ex: V_{ub})
- Need to change charge: FCNC not allowed at tree level, need to proceed via loop diagrams (CKM suppressed)
- If a transition occurs with larger probability than expected $\rightarrow$ **New Particle** (i.e. New Physics)

Introduction

In summary:

- We understand that the Standard Model cannot be the ultimate theory

  It should be a low-energy effective theory of a more fundamental theory at a higher energy scale (TeV range), but no new particles have been found so far at LHC by ATLAS and CMS!

- Flavour structure of the SM:
  $\rightarrow$ provide the suppression mechanism for FCNC processes already observed.
  $\rightarrow$ need to measure the flavour structure to distinguish between the NP models.

- The physics performed at LHCb (flavour physics) goes hand-in-hand with direct searches (ATLAS and CMS).

Introduction

Why the b of LHCb?

- The *b*-quark is the heaviest quark forming hadronic bound states (m~4.7 GeV)
- Must decay outside the 3rd family
  - $\rightarrow$ Long lifetime (~1.6 ps)
  - $\rightarrow$ Many accessible decay channels (small BR's)

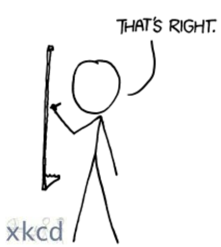
Good for experimentalists!

- Type of processes:

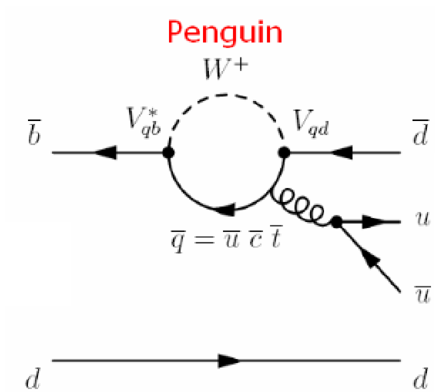
 Dominant: $b \rightarrow c$ (favoured) and $b \rightarrow u$ (suppressed)

 Rare: Flavour Changing Neutral Current (FCNC): $b \rightarrow s,d$

 Flavour oscillations and CP violation

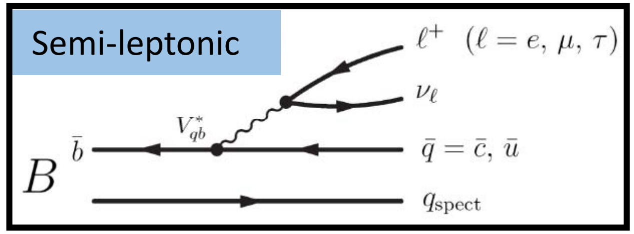


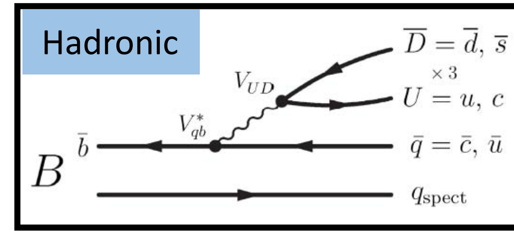
Ideal place to probe New Physics effects!



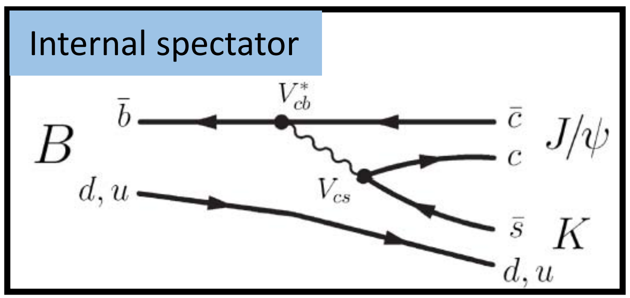
Introduction

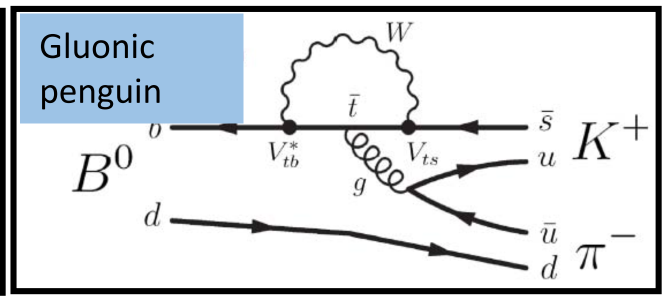
Dominant tree decays:

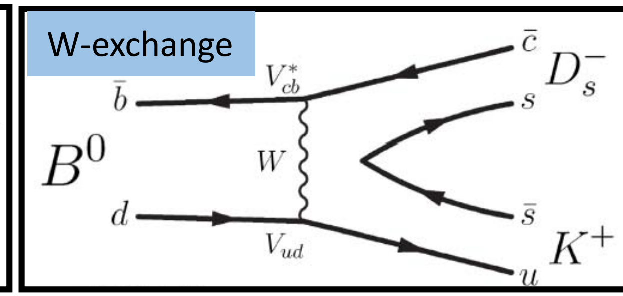




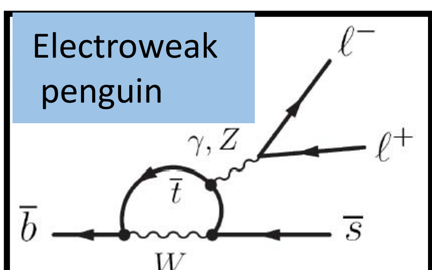
Rare hadronic decays

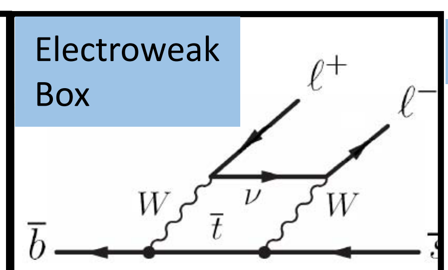


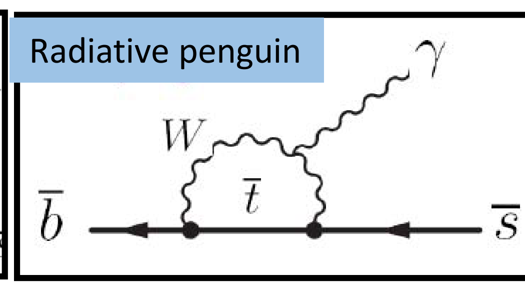


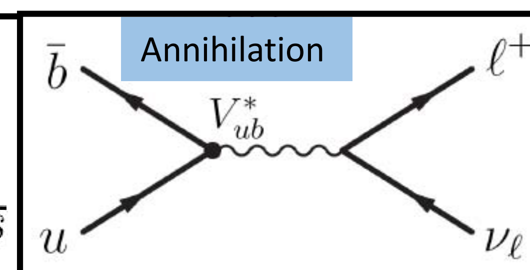


Radiative and leptonic decays









Introduction

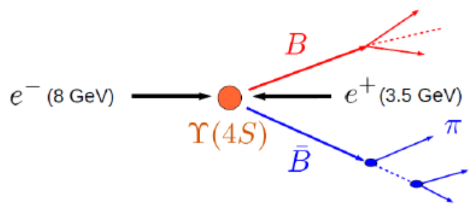
(1999 - 2008 / 2010)

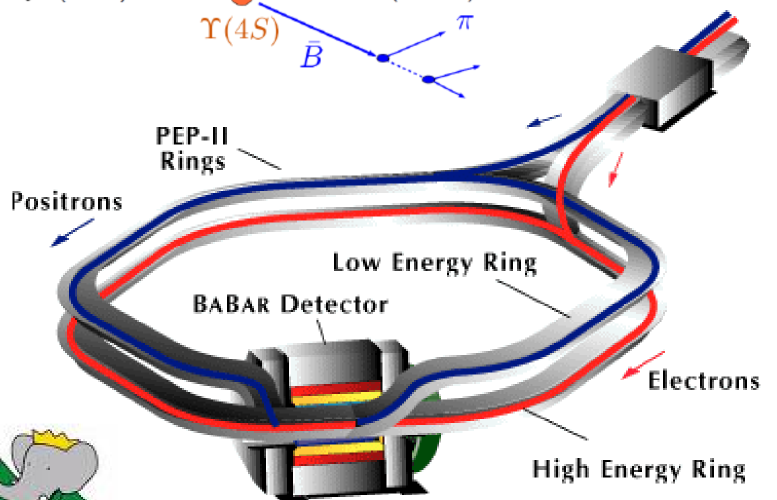
* First measurement of CPV in the B system
* High precision CKM matrix
* Discovery of η_b

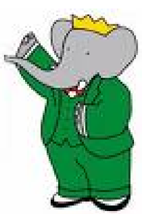
The precesors of LHCb, key in flavour physics:

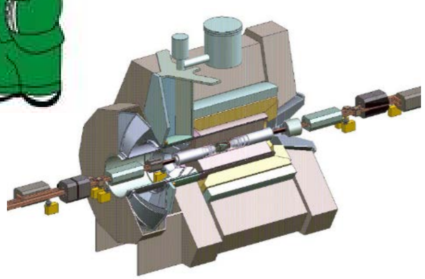
The b-factories: Belle at KEK (Japan) and BaBar at PEP-II (California)

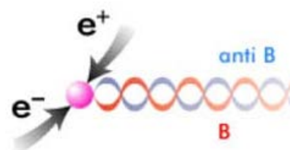
Asymmetric e+ e- colliders working at the Υ(4S) energy.

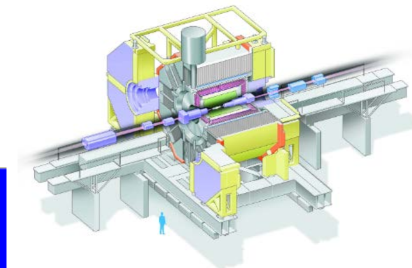


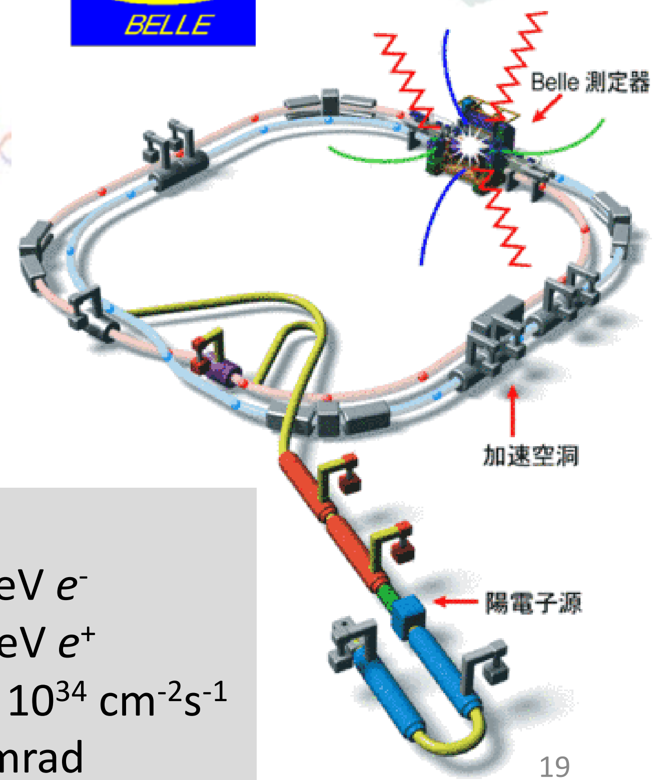












PEP-II /KEK-B

High Energy Ring : 9.0 / 8.0 GeV e^-

Low Energy Ring : 3.1 / 3.5 GeV e^+

Design luminosity : 3 x 10^{33} / 10^{34} $cm^{-2}s^{-1}$

Beam crossing angle : 0 / 22 mrad

Introduction

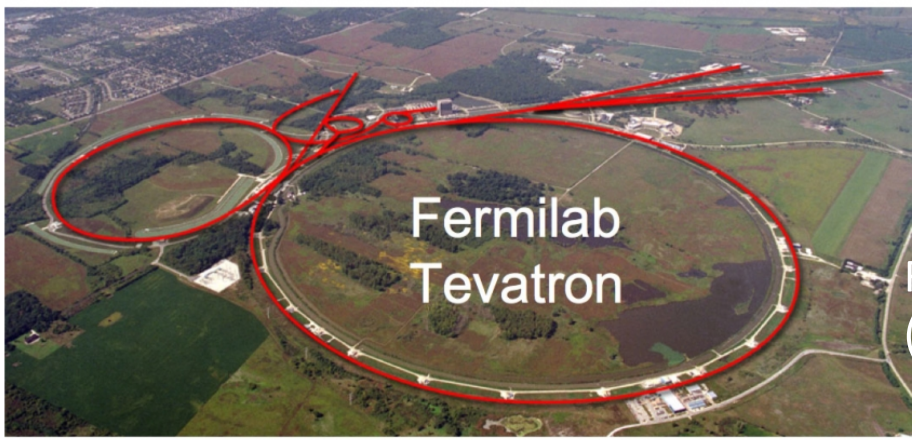
(1987- 2011)

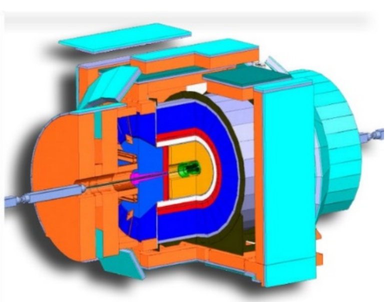
* Discovery of the top quark
* First measurement of B_s oscillations
* Discovery of the Ξ_b baryon

The precesors of LHCb, key in flavour physics

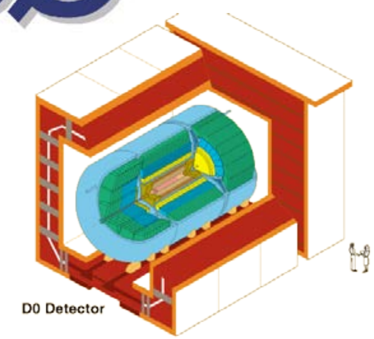
The Tevatron at Fermilab (Illinois): CDF and D0

$p\bar{p}$ collider working at center of mass energy (c.m.) of 1.96 TeV.





TEVATRON
Superconducting $p\bar{p}$ ring
Energy : 1 TeV/beam
Detectors: CDF, D0
Luminosity: 10^{32} cm^{-2}s^{-1}
Physics: W, Z,Top Production
Higgs searches
B physics



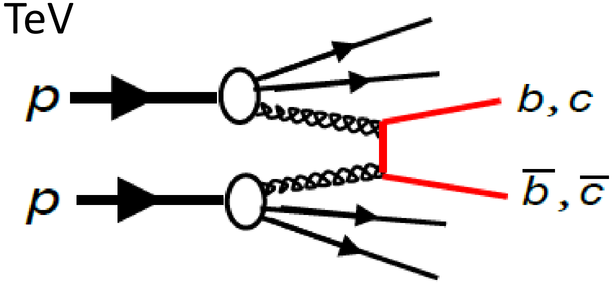
The LHCb experiment

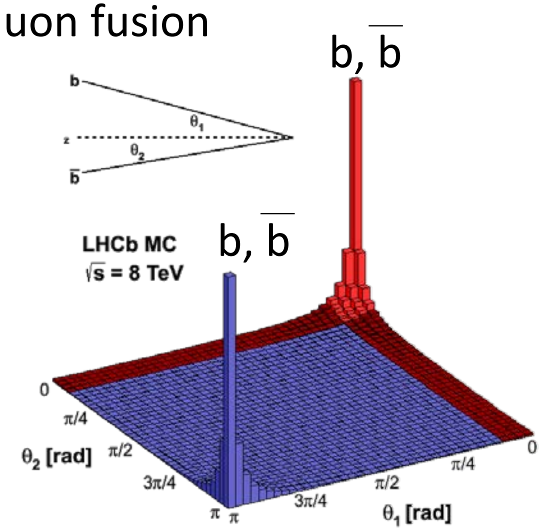
The LHCb experiment

- The $b\bar{b}$ cross section in pp collisions is large, mainly from gluon fusion

  ~250 μb @ $\sqrt{s}$=7 TeV

  ~500 μb @ $\sqrt{s}$=14 TeV





The b quarks hadronize in B, B_s, $B^*_{(s)}$, b-baryons…

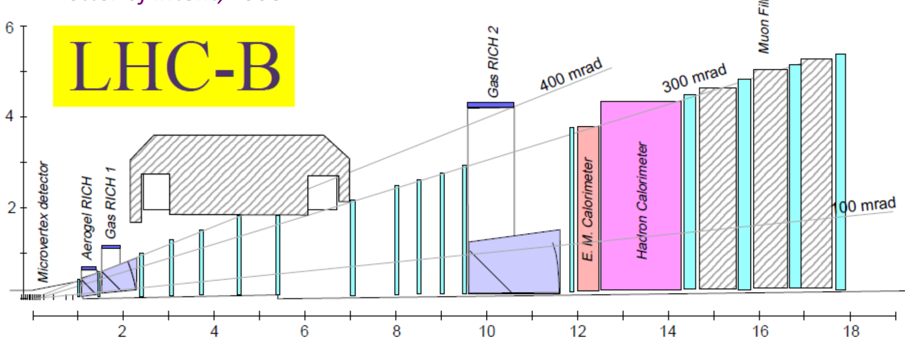
$\rightarrow$ average B meson momentum ~ 80 GeV

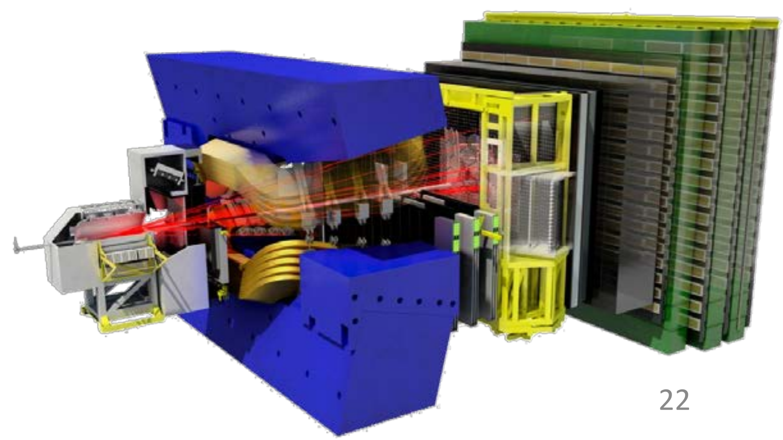
- The LHCb idea: to build a single-arm forward spectrometer:

  ~ 4% of the solid angle ($2 < \eta < 5$),

  ~30% of the *b* hadron production

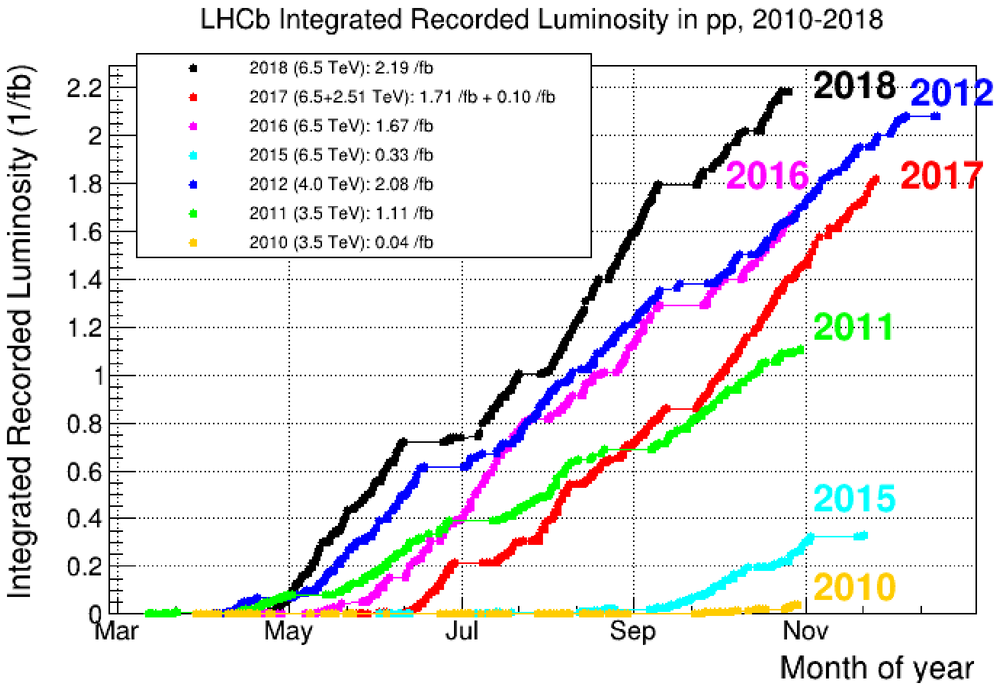
Letter of Intent, 1995





The LHCb experiment

- Very good performance:
  3 fb^{-1} accumulated in Run1 at 7 TeV, and **6 fb^{-1}** in Run2 at 13TeV



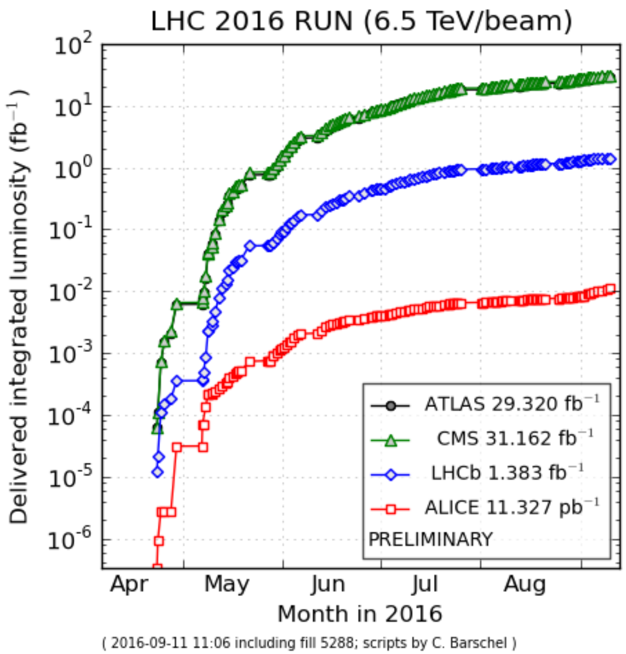
In terms of b-hadrons: $N=\int \mathcal{L}\sigma$

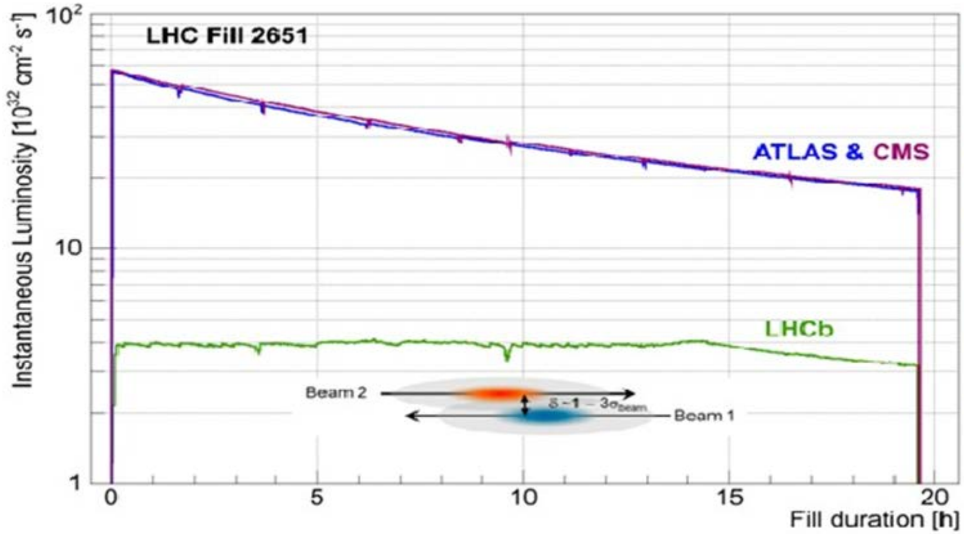
$\rightarrow$ σ ~ 500 μb at 13TeV, x 30% (due to the acceptance) = 150 μb

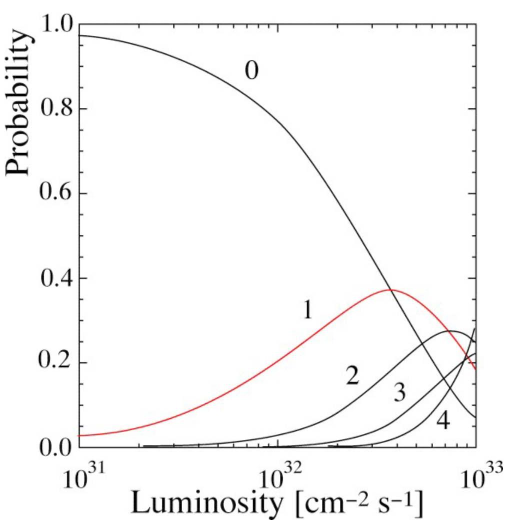
$\rightarrow$ $b\bar{b}$ pairs produced in *1 inverse femtobarn* (N/fb^{-1}) = 10^{15} * 150 x 10^{-6}

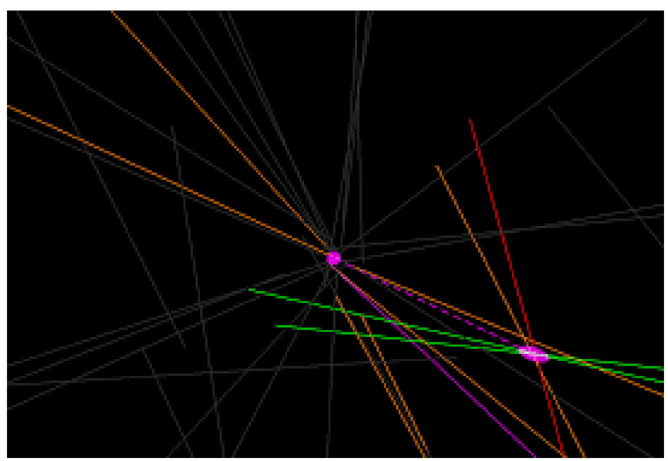
~ 1.5 x 10^{11}

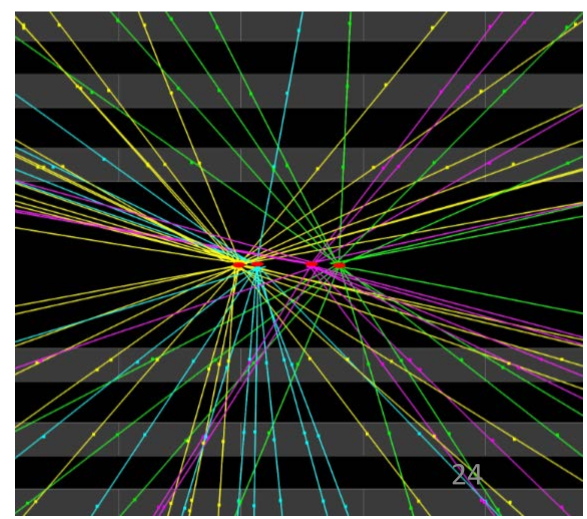
The LHCb experiment





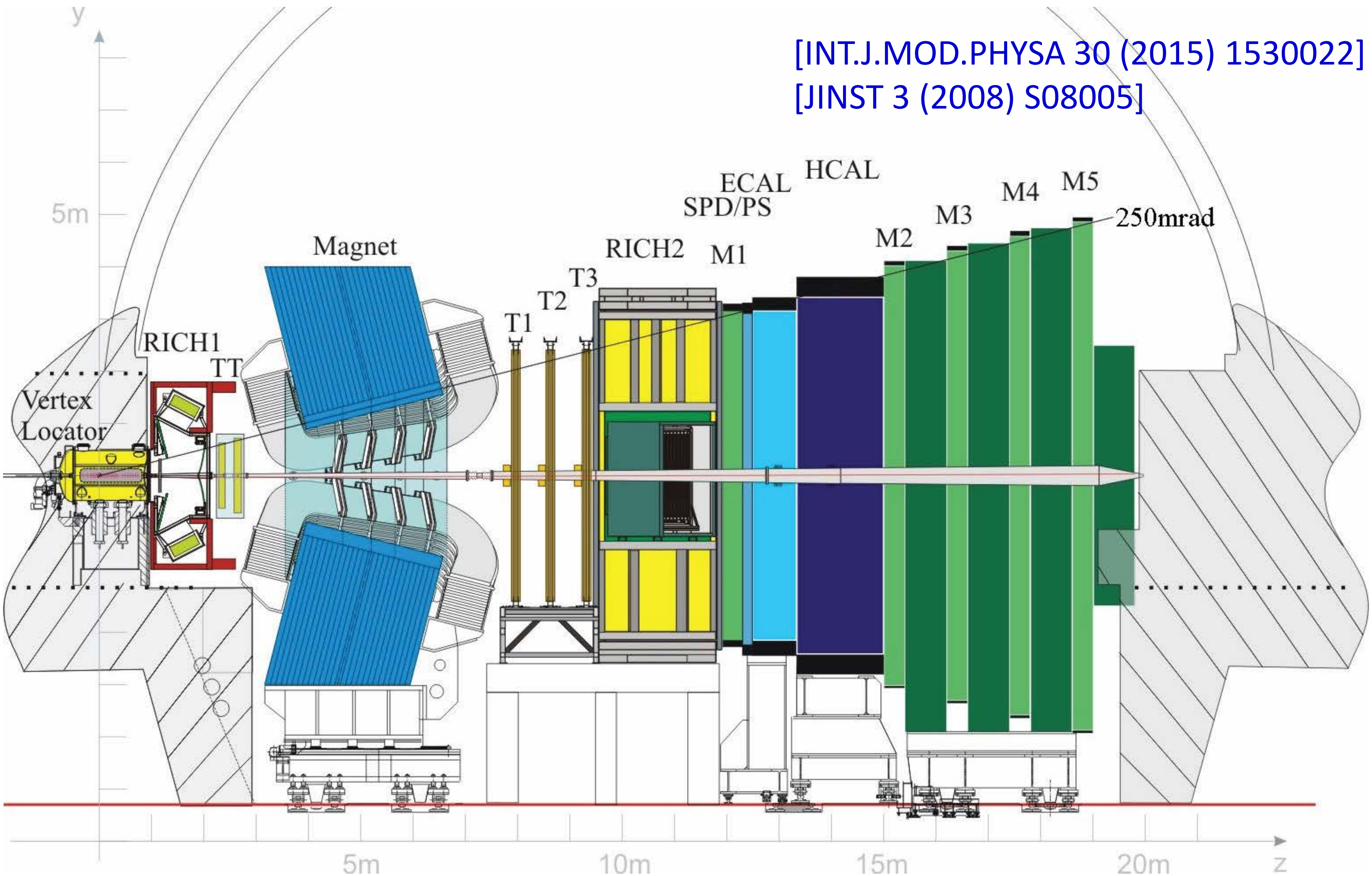






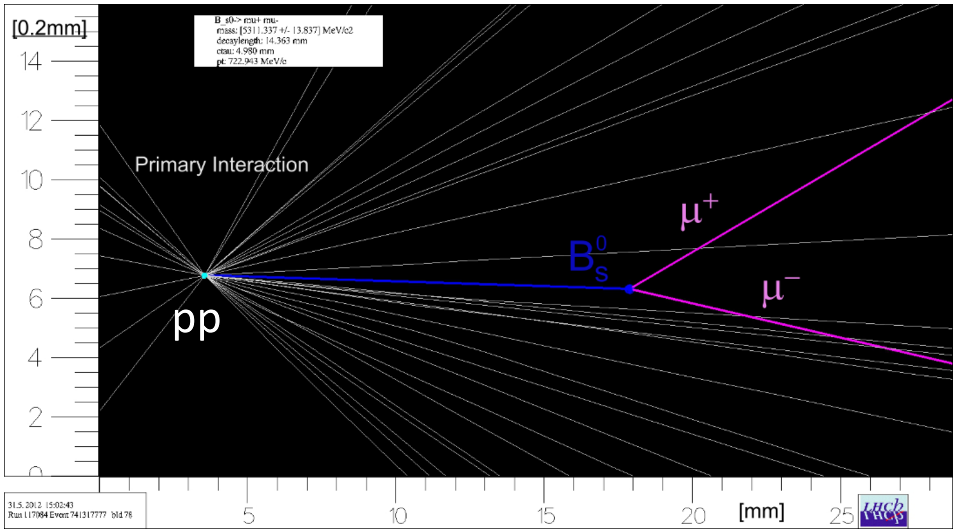
The LHCb experiment

[INT.J.MOD.PHYSA 30 (2015) 1530022]
[JINST 3 (2008) S08005]



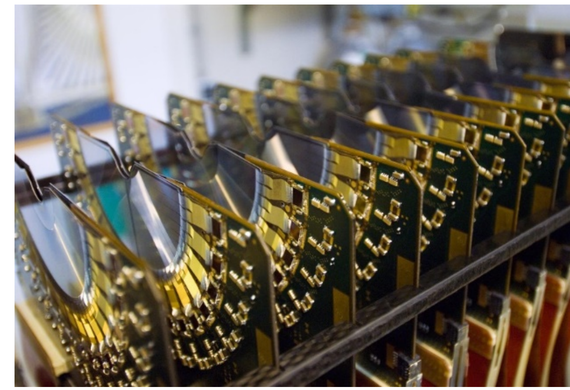
The LHCb experiment

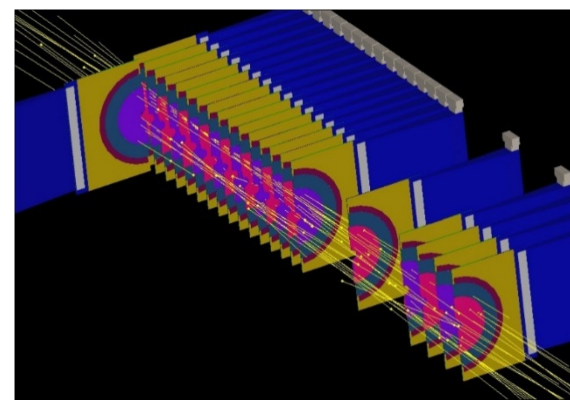
What do we need?



- To reconstruct production and decay vertices
  - $\rightarrow$ Good decay vertex resolution
  - $\rightarrow$ Good impact parameter resolution
- To reconstruct the particle trajectory
  - $\rightarrow$ Good momentum resolution

Vertex detector (VELO)



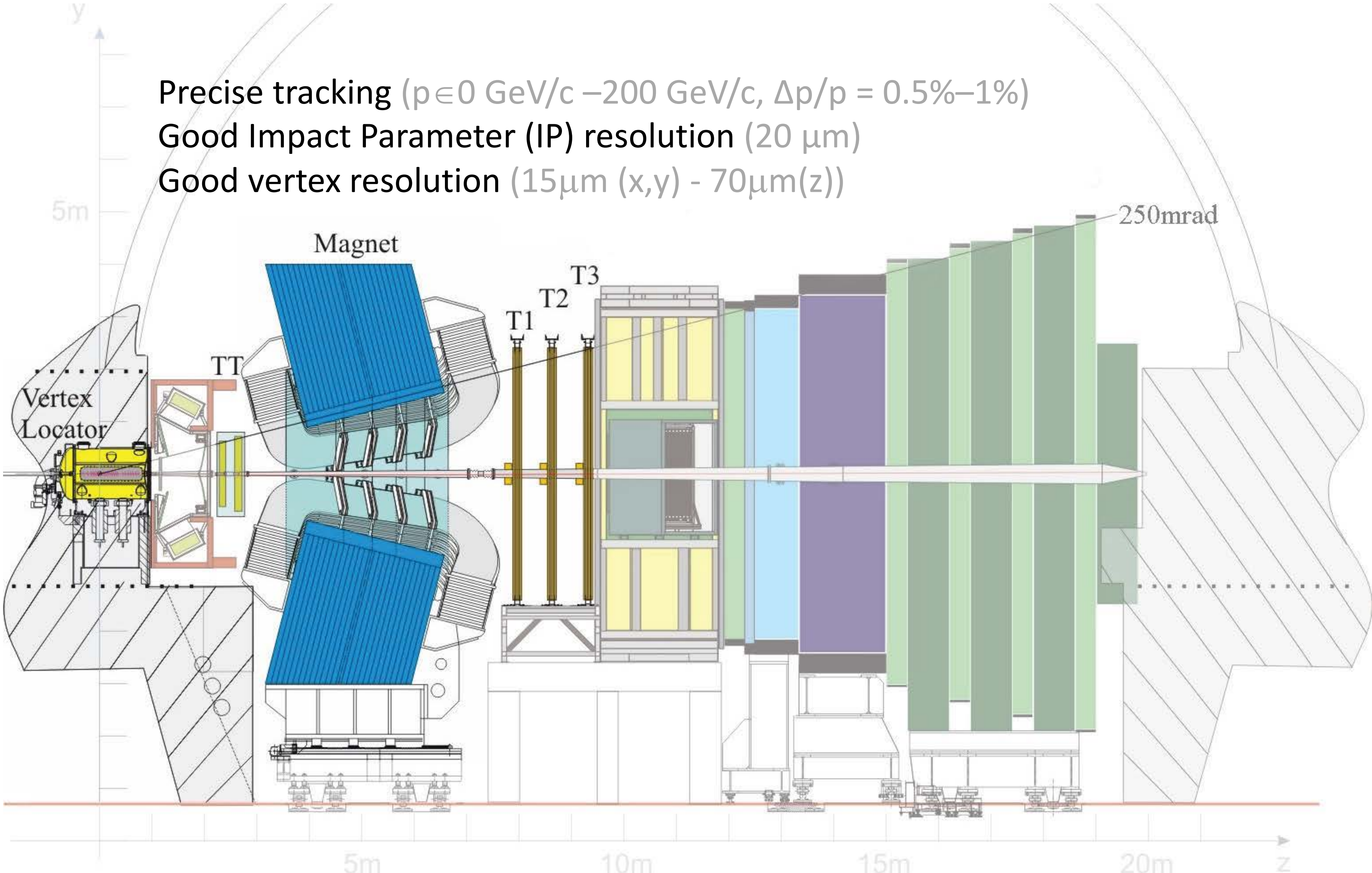


$\langle L \rangle = c\tau\beta\gamma$

Q : How long will a Λ_b baryon be travelling in the detector before decaying? ($\beta\gamma$ ~100)

The LHCb experiment

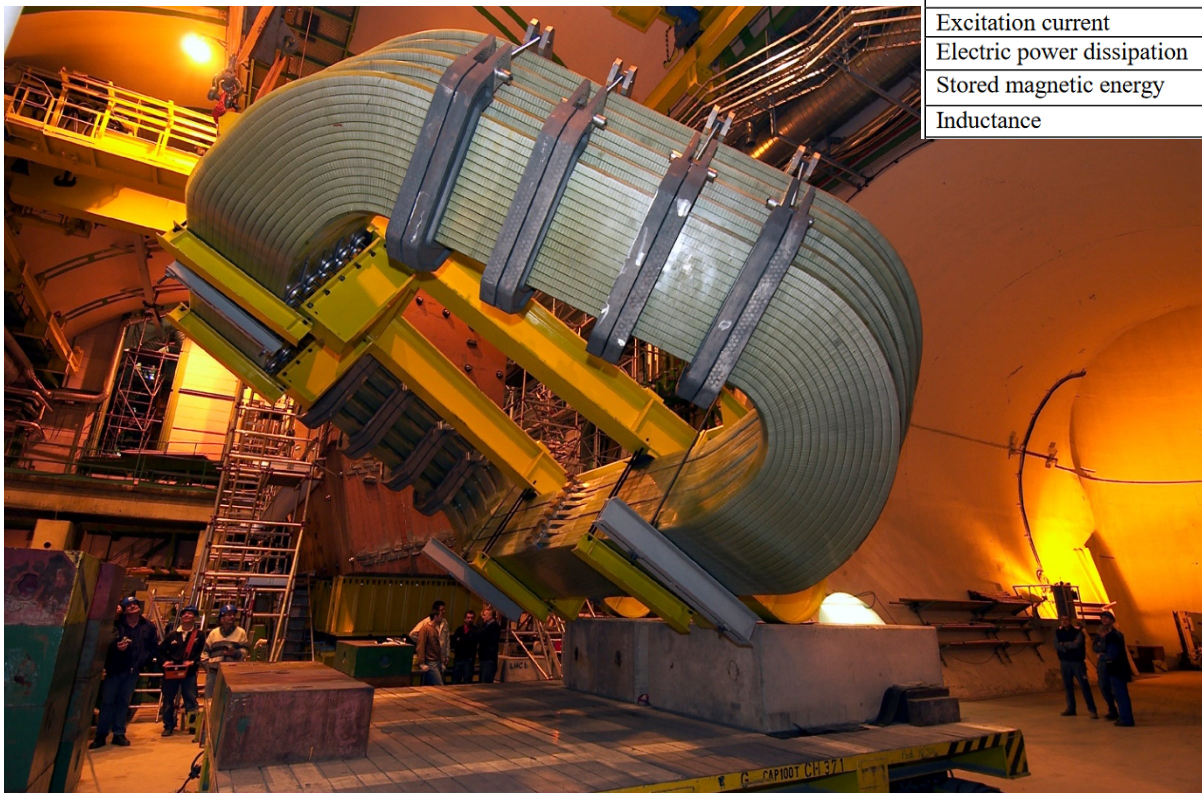
Precise tracking (p∈0 GeV/c –200 GeV/c, Δp/p = 0.5%–1%)
Good Impact Parameter (IP) resolution (20 μm)
Good vertex resolution (15μm (x,y) - 70μm(z))

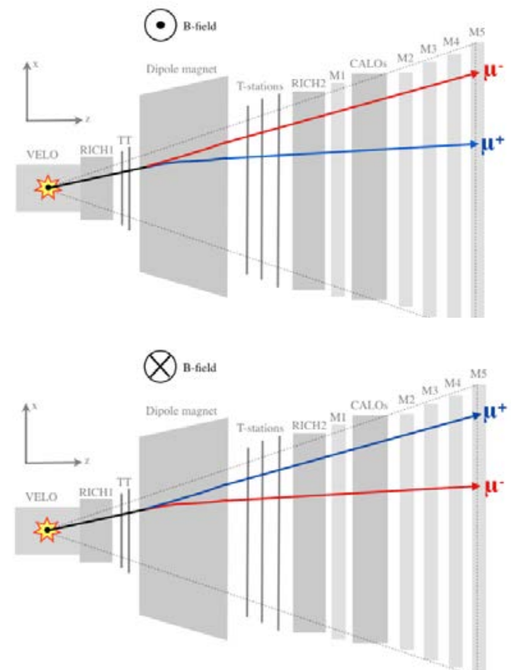


The LHCb experiment

The LHCb magnet:

Magnetic Parameters	
Bending power	$\int B\,dl = 4$ Tm (10 m track length)
Non-uniformity of $\int B\,dl$	$\leq \pm 5\%$ in acceptance (hor.: ± 300 mrad, vert.: ± 250 mrad)
Excitation current	NI = 2 x 1.3 MA
Electric power dissipation	$P_e = 4.2$ MW
Stored magnetic energy	$W_m \approx 32$ MJ
Inductance	$L \approx 2$ H





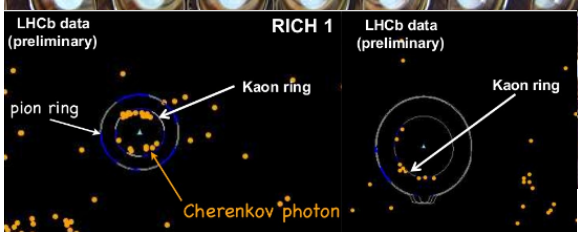
→ Inversion of polarity to study detector asymmetries

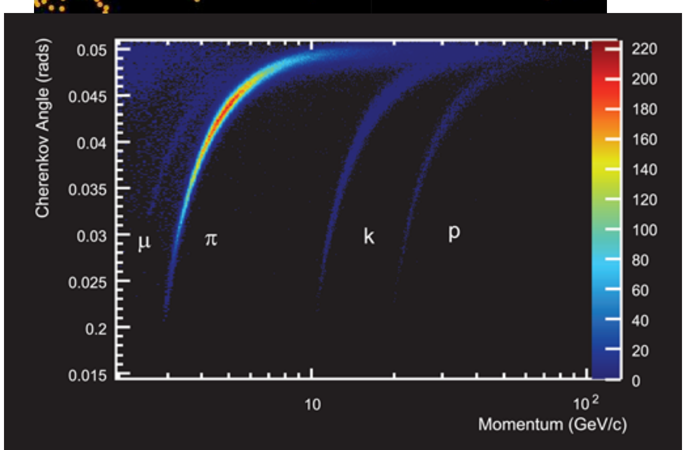
The LHCb experiment

- To recognize the type of particles

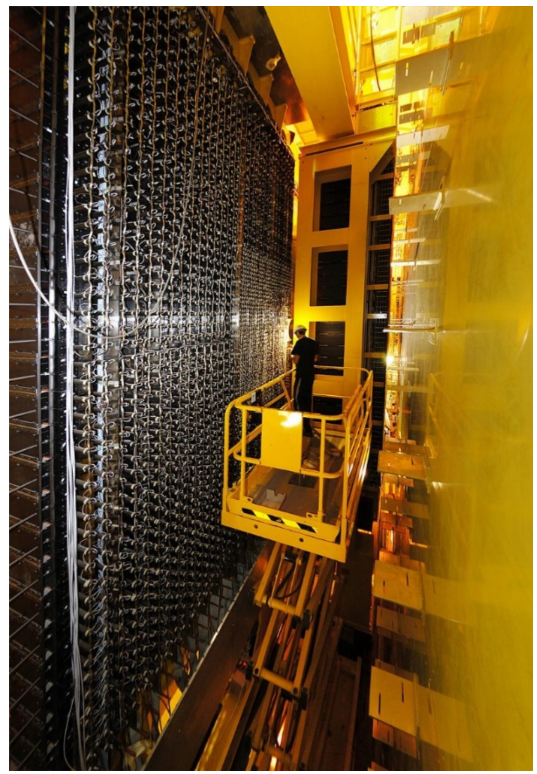
  $\rightarrow$ Good particle identification systems (PID)

Cherenkov detectors (RICH)

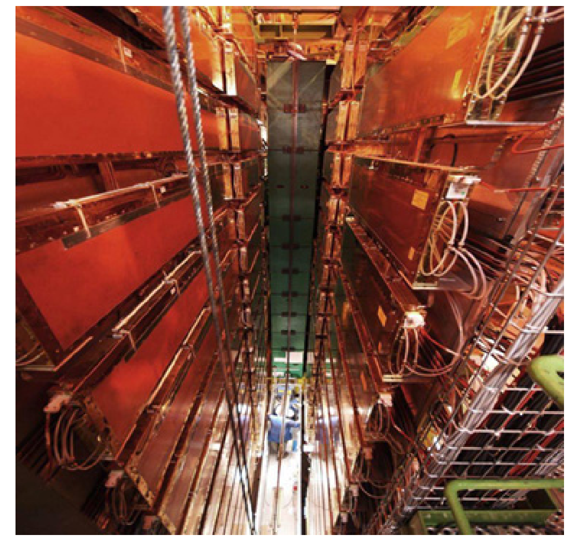


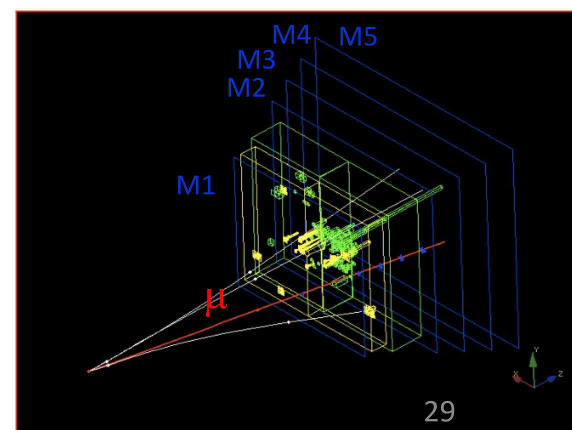


Calorimeters (ECAL, HCAL)



Muon chambers



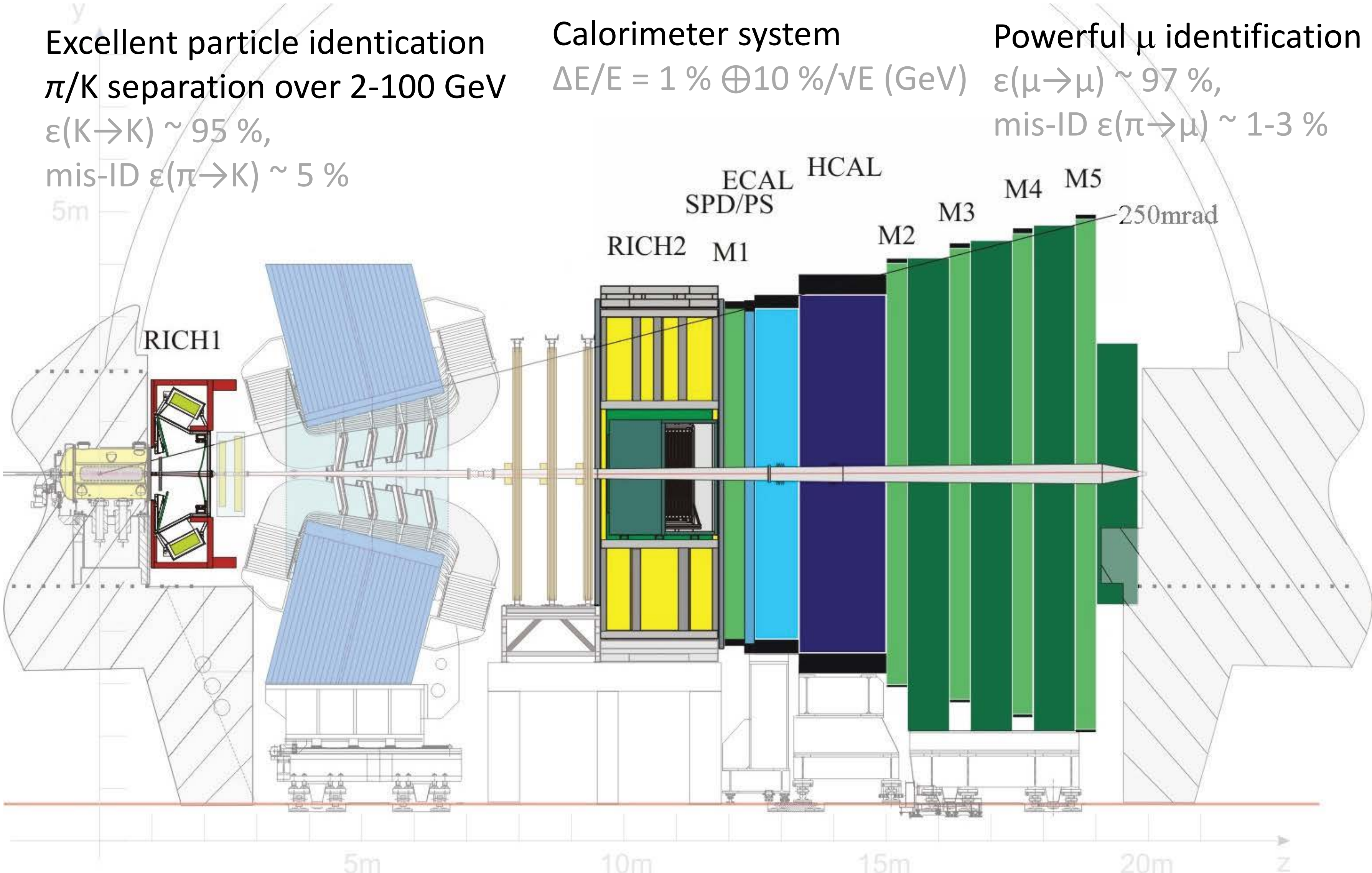


The LHCb experiment

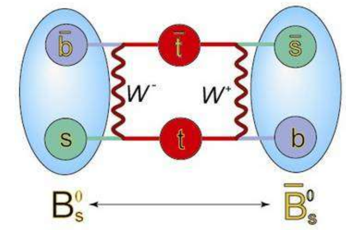
Excellent particle identication
π/K separation over 2-100 GeV
ε(K→K) ~ 95 %,
mis-ID ε(π→K) ~ 5 %

Calorimeter system
ΔE/E = 1 % ⊕10 %/√E (GeV)

Powerful μ identification
ε(μ→μ) ~ 97 %,
mis-ID ε(π→μ) ~ 1-3 %



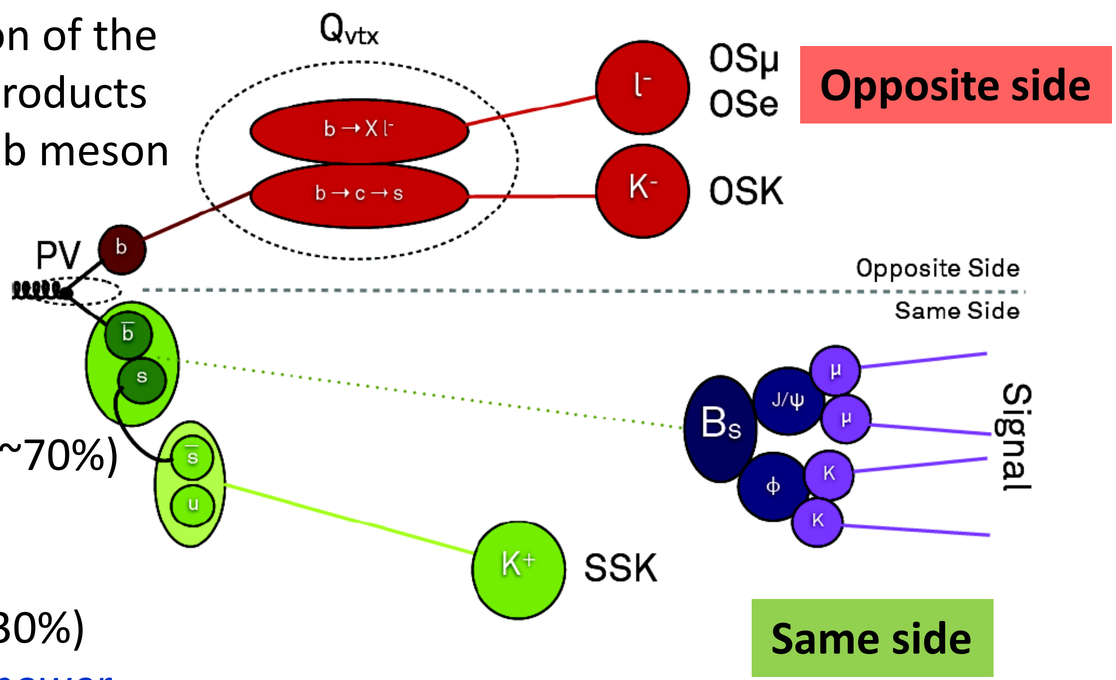
The LHCb experiment



- B mesons oscillates between particle and antiparticle $B^0 \rightarrow \bar{B}^0$
- We need to know the flavour of the particle at the production point

Flavour tagging

Use different algorithms that make use of the characteristics of the fragmentation of the b quark, the charge of the decay products or the charge related to the other b meson produced in $pp \rightarrow X\, b\bar{b}$



Tagging efficiency ε_{tag} : fraction of events with a flavour tag decision (~70%)

Wrong-tag fraction ω:
fraction of tagged events for which tagging decision is wrong (~30%)

Figure of merit: *effective tagging power*

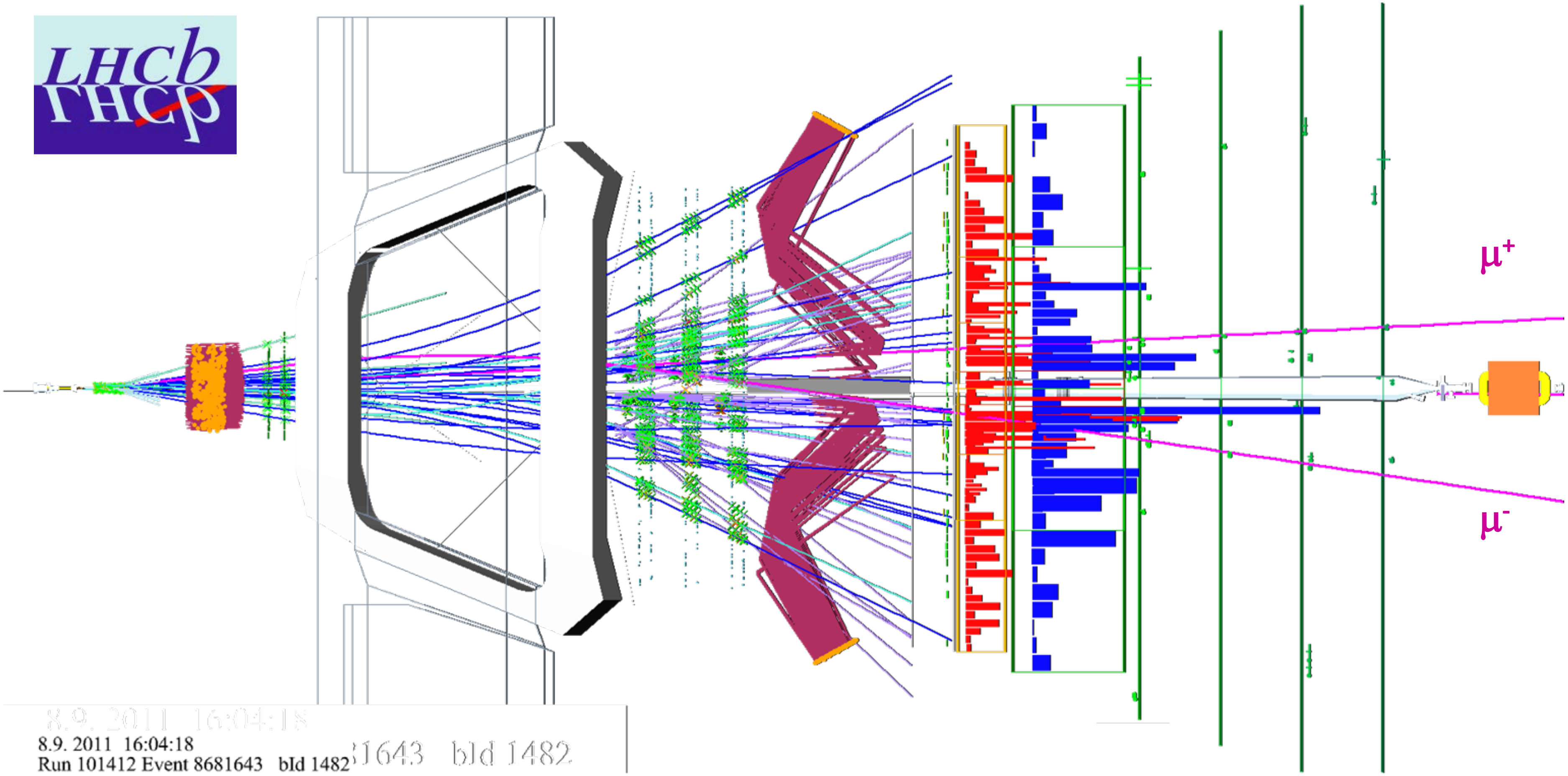
$$\varepsilon_{eff} = \varepsilon_{tag} D^2 = \varepsilon_{tag} (1 - 2\omega)^2$$

$$\boxed{N_{evts} \times \varepsilon_{tag}(1\text{-}2\omega)^2}$$

$D^2 \equiv$ dilution factor

The LHCb experiment

$B_s \to \mu^+\mu^-$ event



The LHCb experiment

- Comparison between facilities:

	$e^+e^- \to \Upsilon(4S) \to B\bar{B}$ PEP-II, KEKB	$p\bar{p} \to b\bar{b}X$ ($\sqrt{s} = 2\,\mathrm{TeV}$) Tevatron	$pp \to b\bar{b}X$ ($\sqrt{s} = 14\,\mathrm{TeV}$) LHC
Production cross-section	1 nb	$\sim 100\,\mu\mathrm{b}$	$\sim 500\,\mu\mathrm{b}$
Typical $b\bar{b}$ rate	10 Hz	$\sim 100\,\mathrm{kHz}$	$\sim 500\,\mathrm{kHz}$
Pile-up	0	1.7	0.5–20
b hadron mixture	B^+B^- (50%), $B^0\bar{B}^0$ (50%)	B^+ (40%), B^0 (40%), B_s^0 (10%), Λ_b^0 (10%), others ($< 1\%$)	
b hadron boost	small ($\beta\gamma \sim 0.5$)	large ($\beta\gamma \sim 100$)	
Underlying event	$B\bar{B}$ pair alone	Many additional particles	
Production vertex	Not reconstructed	Reconstructed from many tracks	
B^0–$\bar{B}^0$ pair production	Coherent (from $\Upsilon(4S)$ decay)	Incoherent	
Flavour tagging power	$\varepsilon D^2 \sim 30\%$	$\varepsilon D^2 \sim 5\%$	

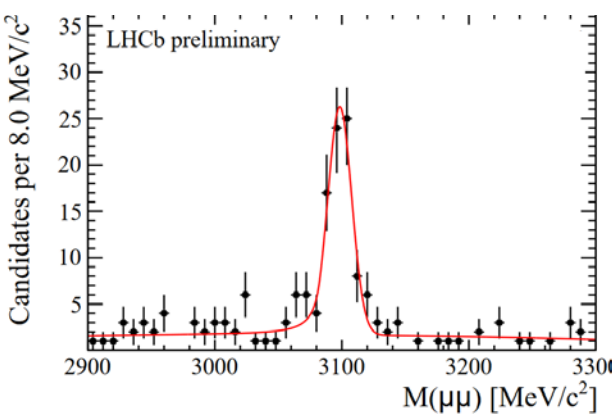
Q : Which is the maximum momentum of the pion in the B $\to \pi\ell\nu$ decay in the lab frame in BaBar at PEP-II and LHCb at LHC experiments ? Which will be easier to measure?

The LHCb experiment

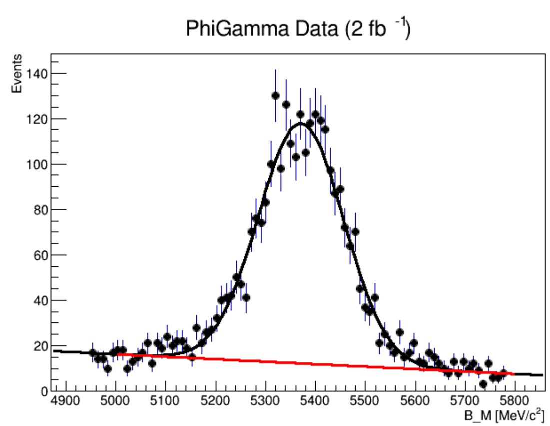
What do we measure?

Examples of observables:

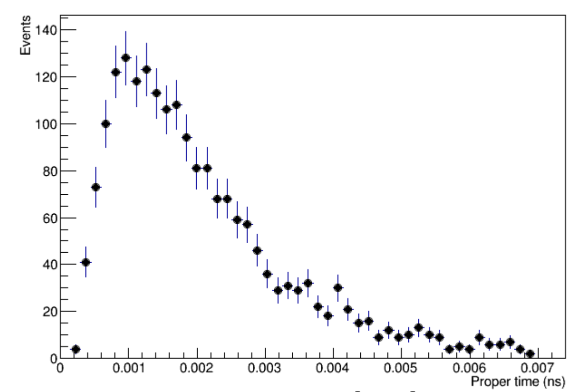
- Invariant masses: from momentum and PID hypothesis of the detected particles
- Decay times: from distance between the origin and decay vertices (and using information of the particle momentum)
- Angular distributions: from directions of the decay products (momentum, vertices)
- Branching fractions: from the mass distributions, counting the number of events
- Time dependent asymmetries (needed flavour tagging!)
- Ratios of observables: cancellation of systematic uncertainties



$M_{J/\psi}$ (MeV)



M_{Bs}(MeV) ($\phi\gamma$)

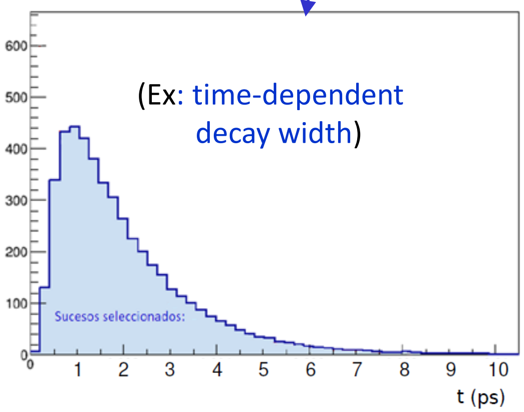


t_{Bs} (ns)

The LHCb experiment

Including experimental effects:

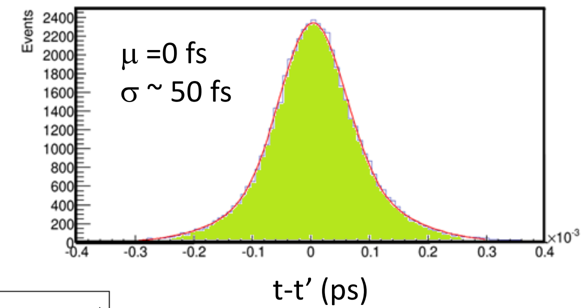
$$\boxed{\text{Observable}(X')} = \boxed{\text{Physics }(X, P)} \cdot \boxed{\text{Acceptance}(X)} \times \boxed{\text{Resolution }(X, X')}$$

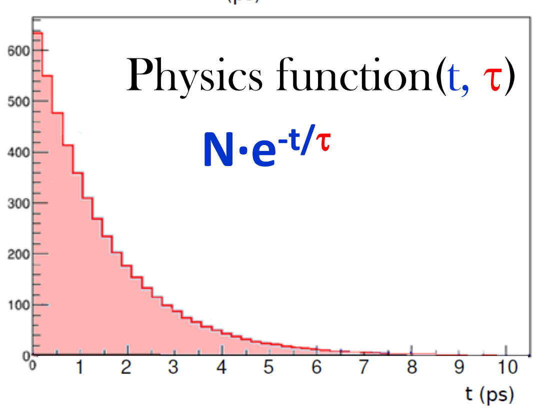


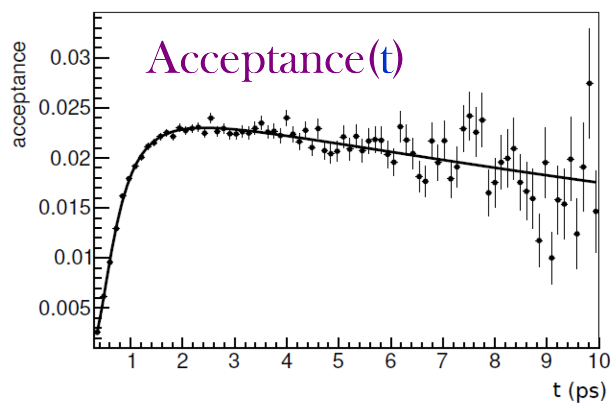
Efficiency dependent of X

$$\frac{\#\ \text{detected}\ (X)}{\#\ \text{produced}\ (X)}$$

R(t,t')







- One can use MC simulations to study the acceptance and resolution functions
- Better: Use control samples from data (similar to the signal channel) to extract them

The CKM matrix

The CKM matrix V_{CKM} describes rotation for quarks between the weak eigenstates (d',s',b') and mass eigenstates (d,s,b)

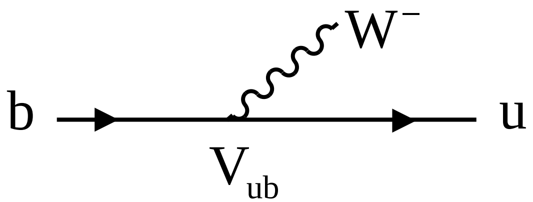
Quarks

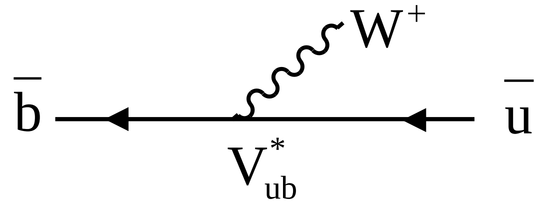
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Antiquarks

$$\begin{pmatrix} \overline{d}' \\ \overline{s}' \\ \overline{b}' \end{pmatrix} = \begin{pmatrix} V_{ud}^* & V_{us}^* & V_{ub}^* \\ V_{cd}^* & V_{cs}^* & V_{cb}^* \\ V_{td}^* & V_{ts}^* & V_{tb}^* \end{pmatrix} \begin{pmatrix} \overline{d} \\ \overline{s} \\ \overline{b} \end{pmatrix}$$

V_{ij} governs the transition from quark j to quark i





CP violation arises in the SM due to complex phases of CKM matrix elements

The CKM matrix

- The CKM matrix is complex and unitary $\hat{V}^{+}_{CKM}\hat{V}_{CKM} = 1$
- 4 independent parameters
  - $\rightarrow$ Fundamental constants of the Standard Model
  - $\rightarrow$ Must be determined from experiment

- **Standard parametrization (PDG):**
  - 3 angles: $\theta_{12}, \theta_{23}, \theta_{13}$ and 1 phase δ

$$V_{CKM} = R_{23} \times R_{13} \times R_{12}$$

$$R_{12} = \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad R_{23} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \quad R_{13} = \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}$$

$$s_{ij} = \sin\theta_{ij} \quad c_{ij} = \cos\theta_{ij}$$

The CKM matrix

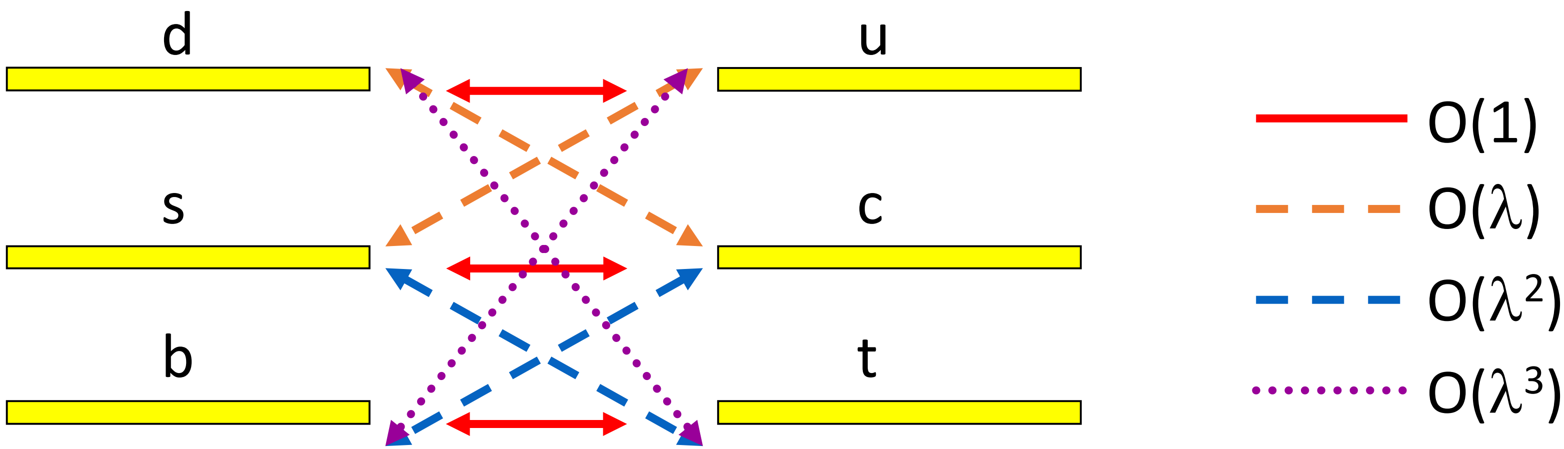
- **Wolfenstein parameterization:** $s_{12} \sim 0.2$, $s_{23} \sim 0.04$, $s_{23} \sim 0.004$

- Perturbative, reflects the hierarchy of the matrix elements in terms of λ

$$\lambda = sin\,\theta_{12} \approx 0.23 \qquad (V_{us})$$

- The four parameters are defined as:

$$\boxed{\lambda} = s_{12} \qquad \boxed{A} = \frac{s_{23}}{s_{12}^2} \qquad \boxed{\rho} = \frac{s_{13}\cos\delta}{s_{12}s_{23}} \qquad \boxed{\eta} = \frac{s_{13}\sin\delta}{s_{12}s_{23}}$$



The CKM matrix

Wolfenstein parameterization to O(λ^3):

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1-\lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1-\lambda^2/2 & A\lambda^2 \\ A\lambda^3(1-\rho-i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

(but next-to leading order corrections in λ may be important at LHC)

$$V_{CKM} = \begin{pmatrix} 1-\lambda^2/2-\lambda^4/8 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + A^2\lambda^5(1/2-\rho-i\eta) & 1-\lambda^2/2-\lambda^4/8\,(1+4A^2) & A\lambda^2 \\ A\lambda^3(1-\bar{\rho}-i\bar{\eta}) & -A\lambda^2 + A\lambda^4(1/2-\rho-i\eta) & 1-A^2\lambda^4/2 \end{pmatrix} + O(\lambda^6)$$

$$(\bar{\rho},\bar{\eta}) \equiv (1-\lambda^2/2)(\rho,\eta)$$

The CKM matrix

- CP Violation in the Standard Model:

  - Requirements for CP violation

$$\left(m_t^2 - m_c^2\right)\left(m_t^2 - m_u^2\right)\left(m_c^2 - m_u^2\right)$$
$$\times\left(m_b^2 - m_s^2\right)\left(m_b^2 - m_d^2\right)\left(m_s^2 - m_d^2\right)\times J_{CP} \neq 0$$

$$J_{CP} = \left|\mathrm{Im}\left\{ V_{i\alpha} V_{j\beta} V_{i\beta}^* V_{j\alpha}^* \right\}\right| \quad \left(i \neq j, \alpha \neq \beta\right)$$

  - *Jarlskog* invariant:

$$J_{CP} = s_{12} s_{13} s_{23} c_{12} c_{23} c_{13} \sin\delta = \lambda^6 A^2 \eta = O\left(10^{-5}\right)$$

$\rightarrow$ CP violation is small in the Standard Model

(and cannot explain the observed baryon asymmetry in the Universe)

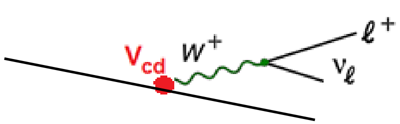
J_{CP} (exp)= (3.172 $\pm$ 0.094) 10^{-5}

The CKM matrix

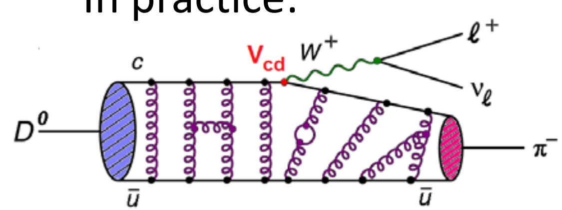
- PDG 2018:

$$
\begin{pmatrix}
\begin{gathered} 0.97446 \pm 0.00010 \\ \text{superallowed } 0^+ \to 0^+ \ \beta \text{ decays} \end{gathered} & \begin{gathered} 0.22452 \pm 0.00044 \\ \text{semileptonic / leptonic kaon decays hadronic} \\ \text{tau decays} \end{gathered} & \begin{gathered} (3.65 \pm 0.12) \times 10^{-3} \\ \text{semileptonic / leptonic B decays} \end{gathered} \\
\begin{gathered} 0.22438 \pm 0.00044 \\ \text{semileptonic charm decays} \\ \text{charm production in neutrino beams} \end{gathered} & \begin{gathered} 0.97359 \pm 0.00011 \\ \text{semileptonic / leptonic charm decays} \end{gathered} & \begin{gathered} (42.14 \pm 0.76) \times 10^{-3} \\ \text{semileptonic B decays} \end{gathered} \\
\begin{gathered} (8.96 \pm 0.24) \times 10^{-3} \\ B_d \text{ oscillations} \end{gathered} & \begin{gathered} (41.33 \pm 0.74) \times 10^{-3} \\ B_s \text{ oscillations} \end{gathered} & \begin{gathered} 0.999105 \pm 0.000032 \\ \text{single top production} \end{gathered}
\end{pmatrix}
$$

In theory:



In practice:



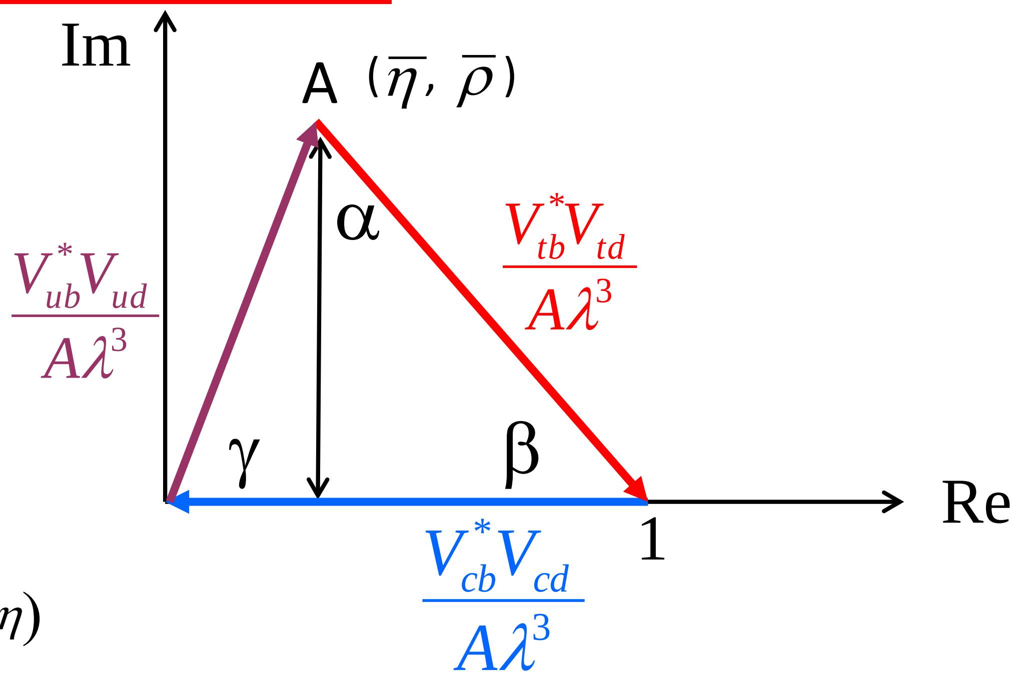
$\rightarrow$ Need theory to describe QCD effects (lattice QCD)

The CKM matrix

The Unitarity Triangle $\rightarrow$ CKM is unitary:

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} V_{ud} & V_{cd} & V_{td} \\ V_{us} & V_{cs} & V_{ts} \\ V_{ub} & V_{cb} & V_{tb} \end{pmatrix}^{*} = \mathrm{I}$$

$$\boxed{V_{ud}V_{ub}^{*} + V_{cd}V_{cb}^{*} + V_{td}V_{tb}^{*} = 0}$$



$$(\bar{\rho}, \bar{\eta}) \equiv (1 - \lambda^{2}/2)(\rho, \eta)$$

$$\boxed{\gamma \equiv \arg\left[-\frac{V_{ub}^{*}V_{ud}}{V_{cb}^{*}V_{cd}}\right] = \tan^{-1}\frac{\eta}{\rho} \sim 70^{\circ}}$$

$$\boxed{\beta \equiv \arg\left[-\frac{V_{cb}^{*}V_{cd}}{V_{tb}^{*}V_{td}}\right] = \tan^{-1}\frac{\bar{\eta}}{1-\bar{\rho}} \sim 21^{\circ}}$$

$$\boxed{\alpha \equiv \pi - \beta - \gamma}$$

The CKM matrix

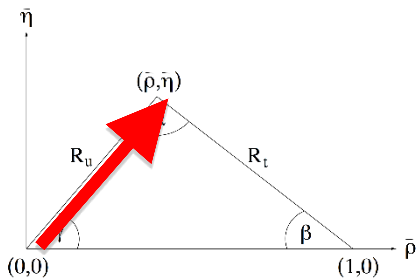
The Unitarity Triangle

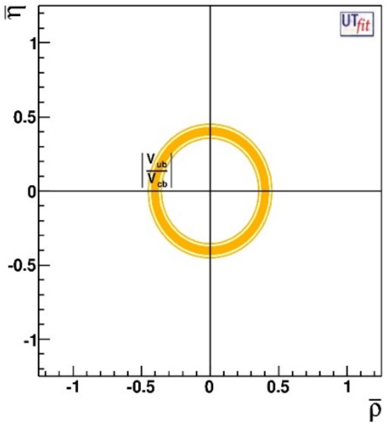
The idea: try to measure as many flavour observables as possible
overconstraint the unitarity triangle

Ex: Measuring the b→ u $\ell\nu$ vs the b→ c $\ell\nu$ transition

$$\frac{\Gamma(b \to u \ell \nu)}{\Gamma(b \to c \ell \nu)} \sim \left|\frac{V_{ub}}{V_{cb}}\right|^2 \sim \frac{1}{50}$$

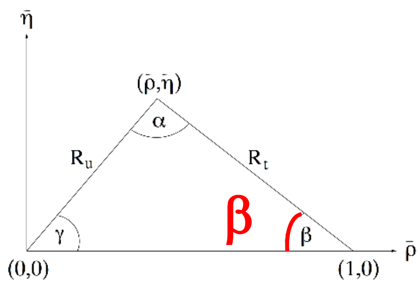
$$\left|\frac{V_{ub}}{V_{cb}}\right| = \frac{\lambda}{1-\lambda^2/2}\sqrt{\bar{\rho}^2 + \bar{\eta}^2}$$

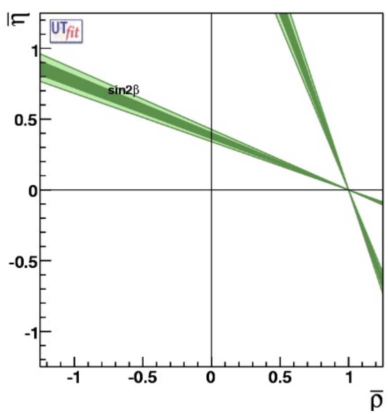




Ex: Measuring time-dependent asymmetries in b→c$\bar{c}$ s decays
(effect from interference of mixing and decay)

$$\beta = \tan^{-1}\frac{\bar{\eta}}{1-\bar{\rho}}$$

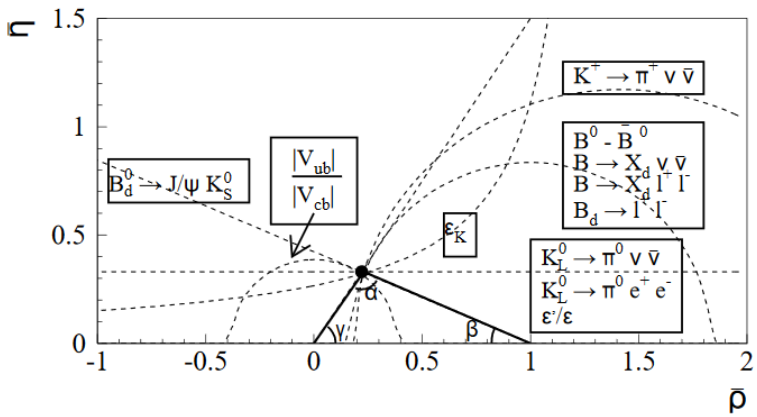


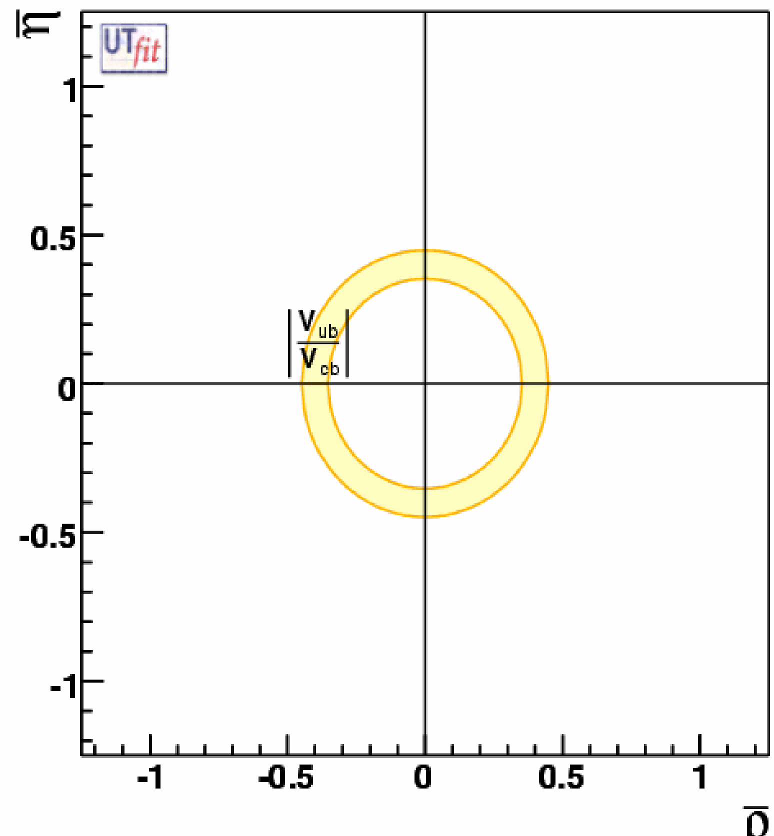


The CKM matrix

The Unitarity Triangle

The idea: try to measure as many flavour observables as possible
overconstraint the unitarity triangle





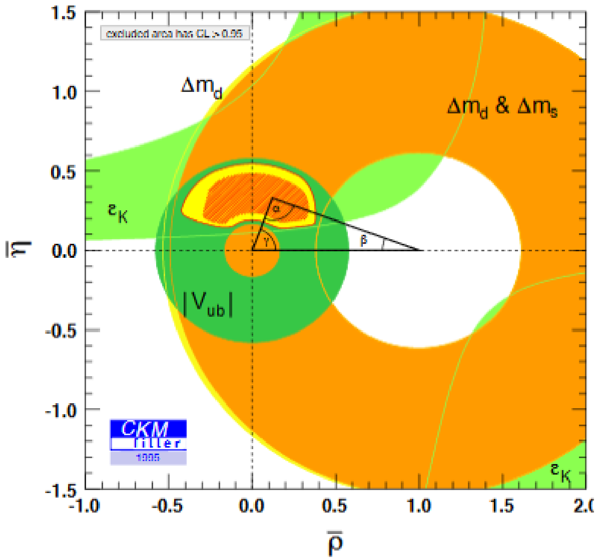
- If all measurements meet in the same apex$\rightarrow$ good understanding of the flavour structure of the SM

- If not $\rightarrow$ New Physics

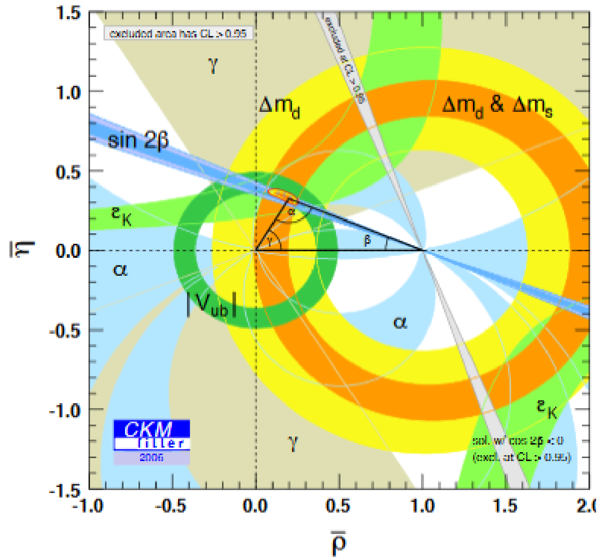
The CKM matrix

The Unitarity Triangle

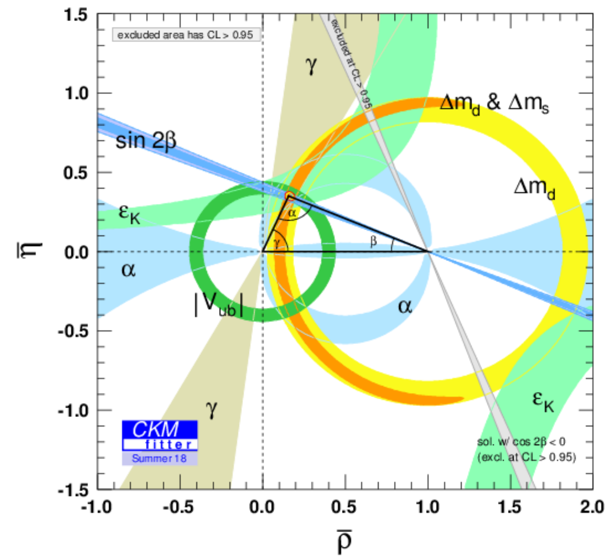
1995



2006

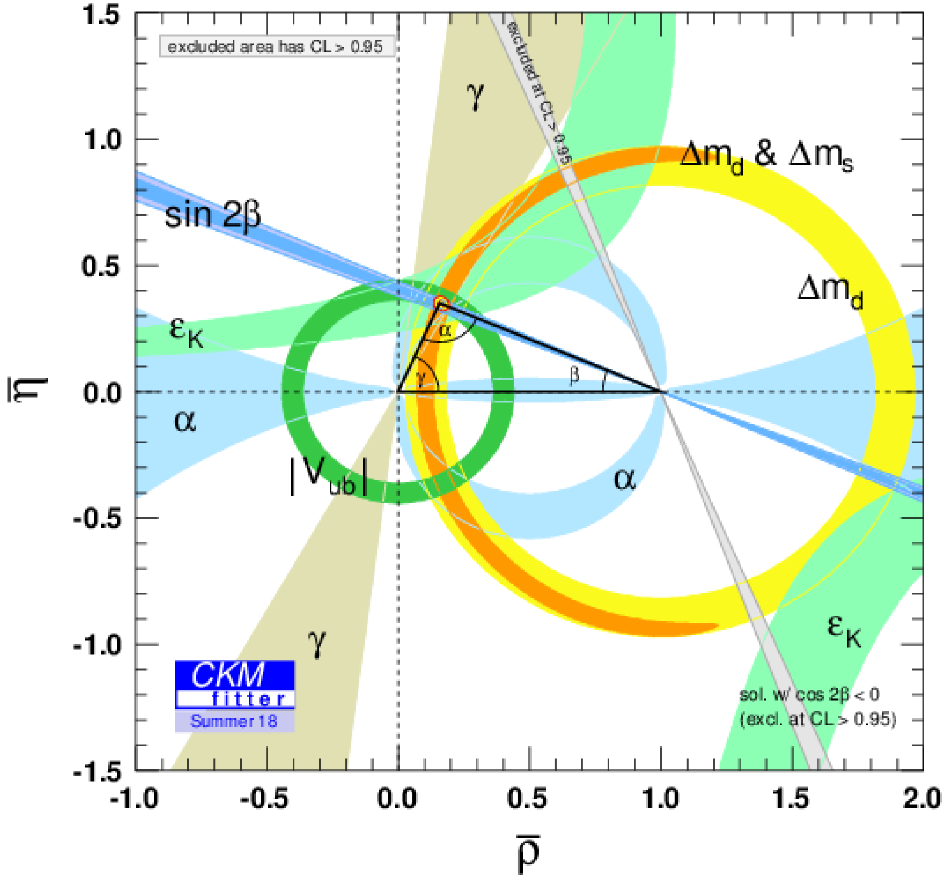


2018



- Precision measurements of CKM elements (fundamental parameters!)
- Measure all angles and sides in many different ways and look for inconsistencies $\rightarrow$ quantum effects from new particles
- Compare tree level processes (new physics is not expected) with loop processes sensitive to new particles
- With more precision the new physics scale has to be higher.

The CKM matrix



- Good agreement between experimental measurements

- Validation of Standard Model in the flavour sector

- Understanding from QCD is crucial

- Still room for New Physics, need more precision!

$$\bar{\rho} = 0.1577 \pm 0.0085$$
$$\bar{\eta} = 0.3493 \pm 0.0085$$

http://ckmfitter.in2p3.fr/

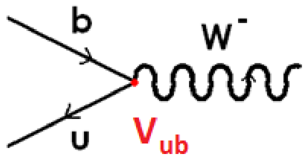
http://www.utfit.org/UTfit/

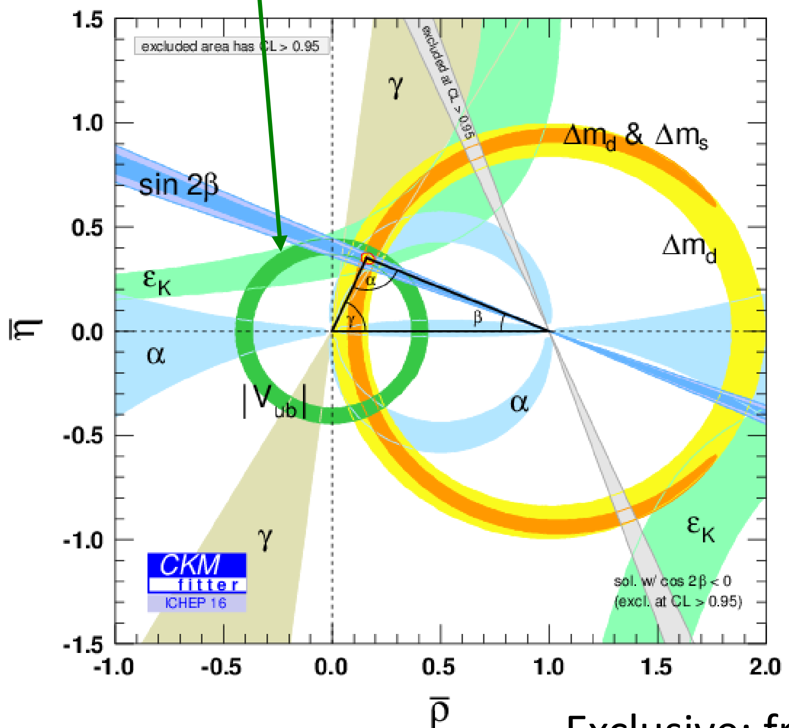
The CKM matrix

- Measuring the CKM matrix element V_{ub} at LHCb:

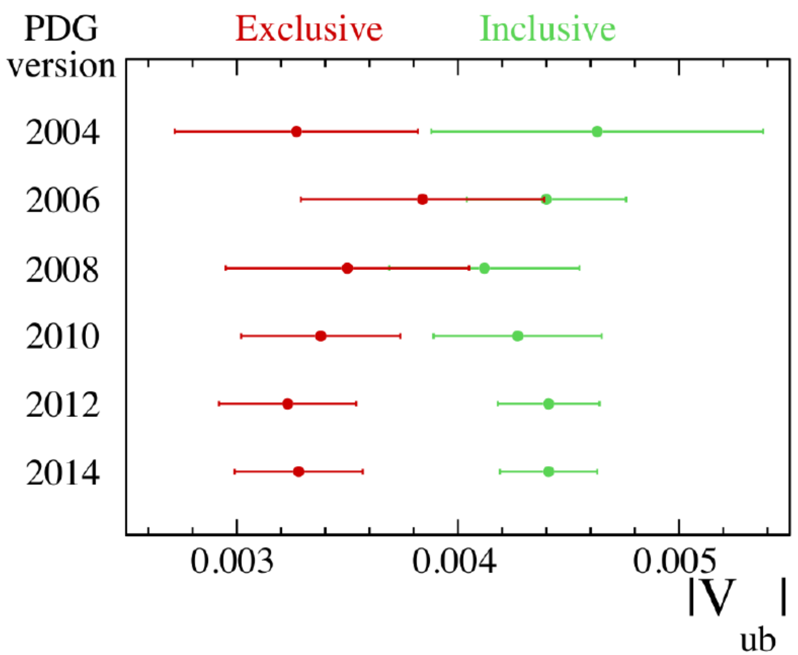
Key constraint in the flavour picture since it is extracted from decays at tree-level

(No loops $\rightarrow$ no New Physics expected, it's a reference to be compared with)





Discrepancy between ways of measuring this element during years:



Exclusive: from specific channels: $B\rightarrow\pi\ell\nu$, $B\rightarrow\rho\ell\nu$ …

Inclusive: from total rate $B \rightarrow X_u\, \ell\nu$

The CKM matrix

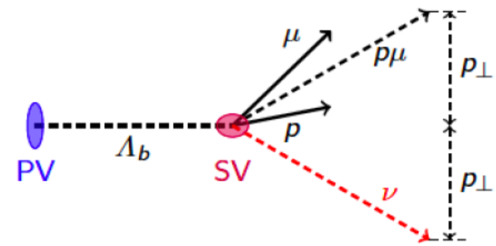
At LHCb very challenging due to the missing neutrino:

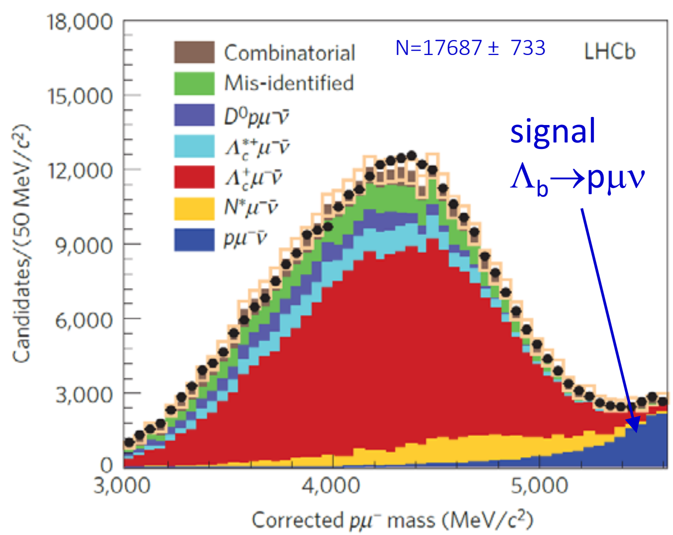
Using semileptonic decays of b-baryons:

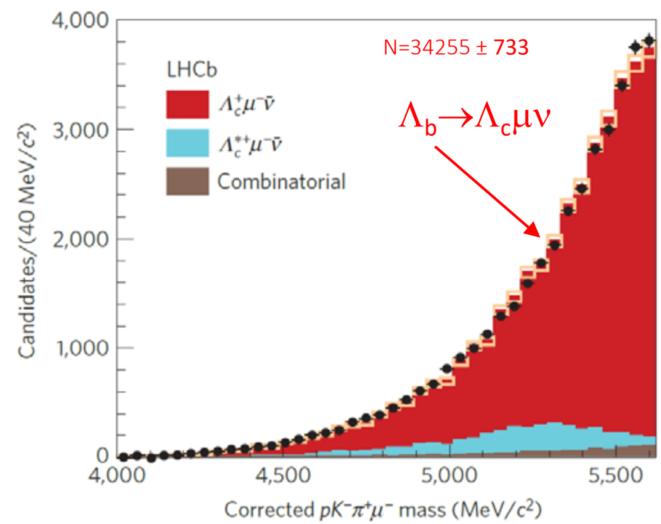
$$\frac{\mathcal{B}(\Lambda_b^0 \to p\mu^-\overline{\nu}_\mu)}{\mathcal{B}(\Lambda_b^0 \to \Lambda_c^+\mu^-\overline{\nu}_\mu)} = \frac{|V_{ub}|^2}{|V_{cb}|^2}$$ x Ratio of form factors

(few % accuracy from LQCD)

- Use information from displaced vertex
- Select high q^2 region (theory more precise)
- Corrected mass: $m_{\rm corr} = \sqrt{m_{h\mu}^2 + p_\perp^2} + p_\perp$







[Nature Physics 10 (2015) 1038]

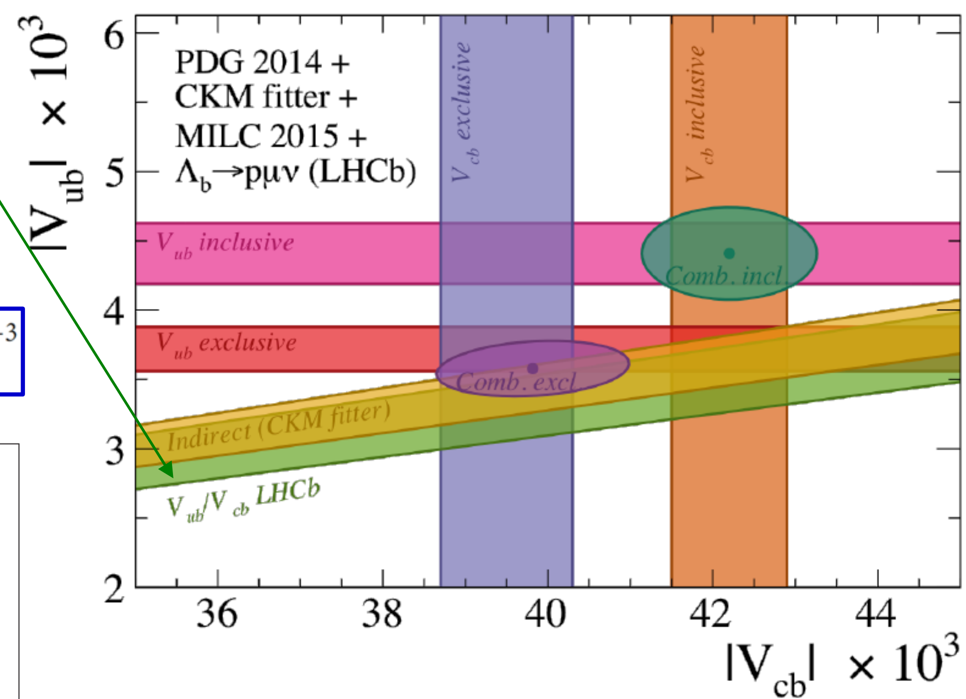
The CKM matrix

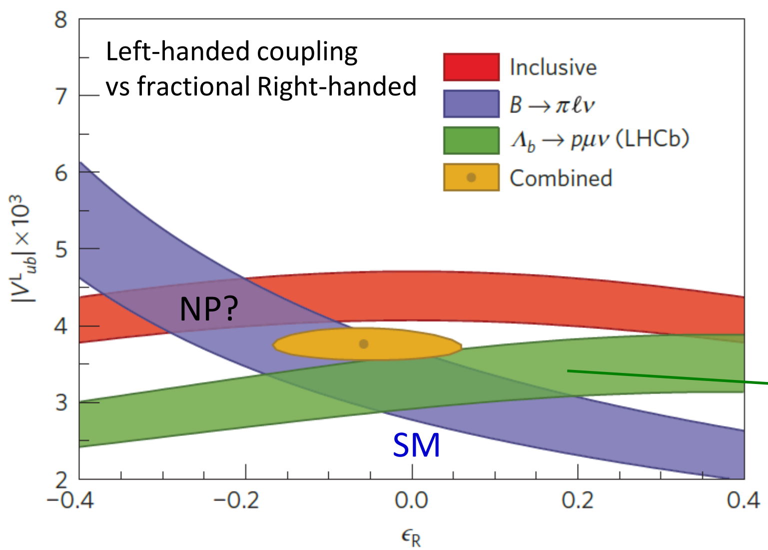
[Nature Physics 10 (2015) 1038]

$$\frac{|V_{ub}|}{|V_{cb}|} = 0.083 \pm 0.004 \pm 0.004$$

Using the world average from exclusive V_{cb}:

$$\boxed{|V_{ub}| = (3.27 \pm 0.15 \pm 0.16 \pm 0.06) \times 10^{-3}}$$

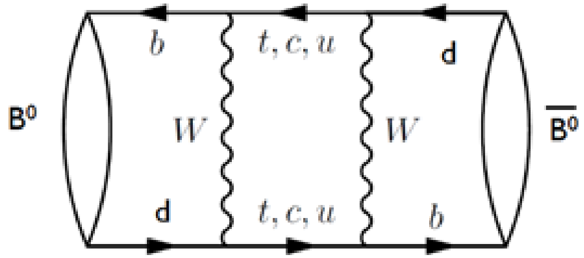




Disfavours New Physics models with Right-handed currents

B mixing and CP violation

Neutral mesons (K,D,B) oscillate between matter and anti-matter states:

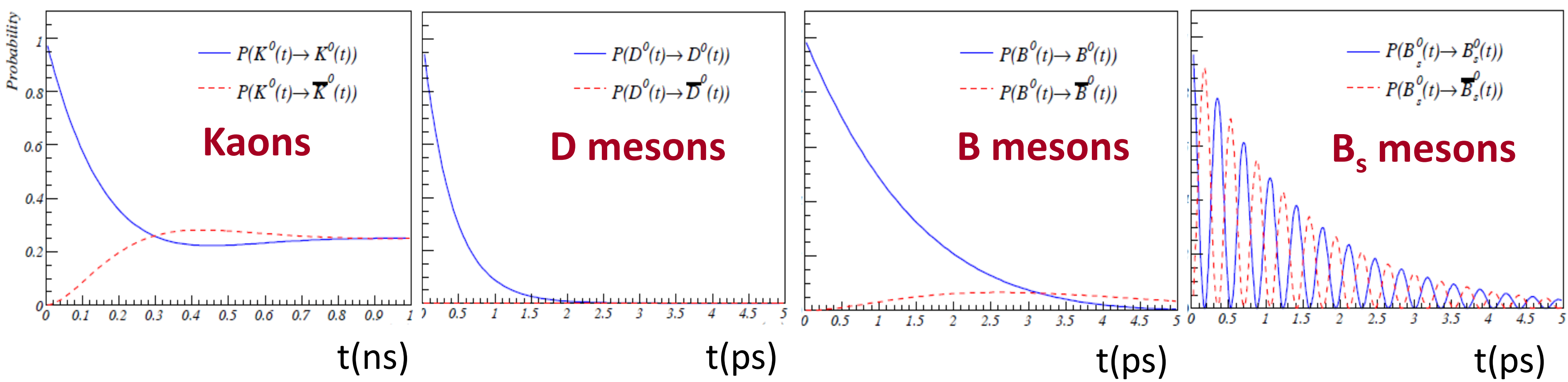


$$i\frac{\partial\psi}{\partial t} = H\psi = (M - \frac{i}{2}\Gamma)\psi$$

The physical mass states (m_H, m_L) are an admixture of weak states

$$\mathrm{Prob}(t) = \frac{e^{-\Gamma t}}{2}\left(\cosh\tfrac{1}{2}\Delta\Gamma t \pm \cos\Delta m t\right)$$

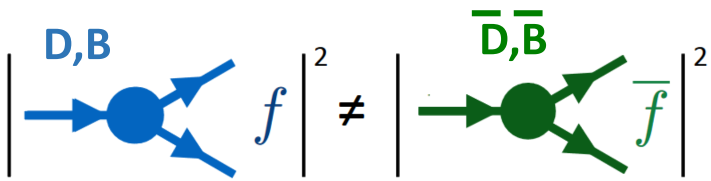
with mass difference $\Delta m = m_H - m_L$ and decay-width difference $\Delta\Gamma = \Gamma_L - \Gamma_H$



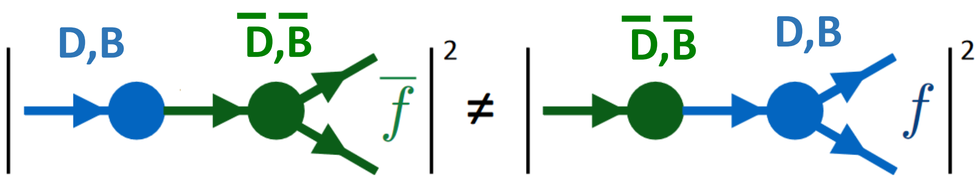
B mixing and CP violation

Three different types of CP violation:

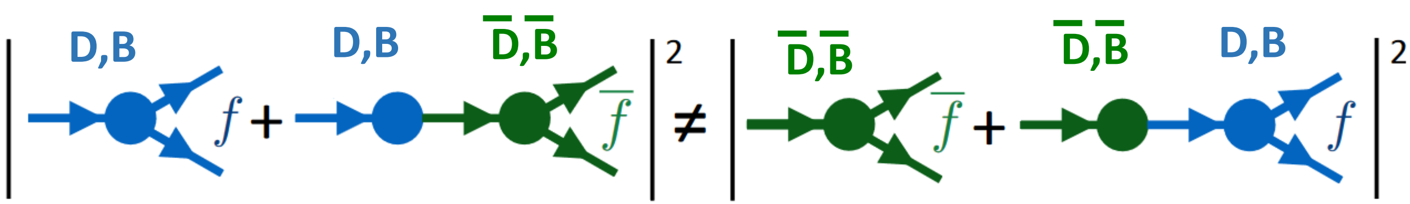
Direct CPV



CPV in mixing



CPV in interference between mixing and decay



B mixing and CP violation

Mixing of neutral B mesons governed by

$$i\frac{\partial}{\partial t}\begin{pmatrix}a\\b\end{pmatrix} = H\begin{pmatrix}a\\b\end{pmatrix} = \begin{pmatrix} M_{11}-\frac{i}{2}\Gamma_{11} & M_{12}-\frac{i}{2}\Gamma_{12} \\ M_{12}^{*}-\frac{i}{2}\Gamma_{12}^{*} & M_{22}-\frac{i}{2}\Gamma_{22} \end{pmatrix}\begin{pmatrix}a\\b\end{pmatrix}$$

Mass eigenstates:

$$\left|B_{L,H}\right\rangle = p\left|B^0\right\rangle \pm q\left|\overline{B^0}\right\rangle$$

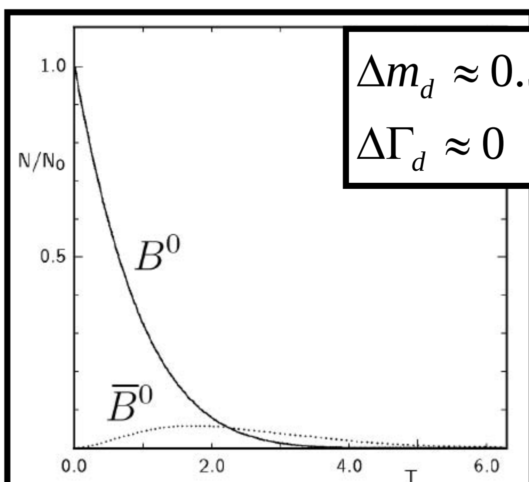
p and q represent the amount of state mixing

$$|p|^2 + |q|^2 = 1$$

$$|q/p| = 1$$

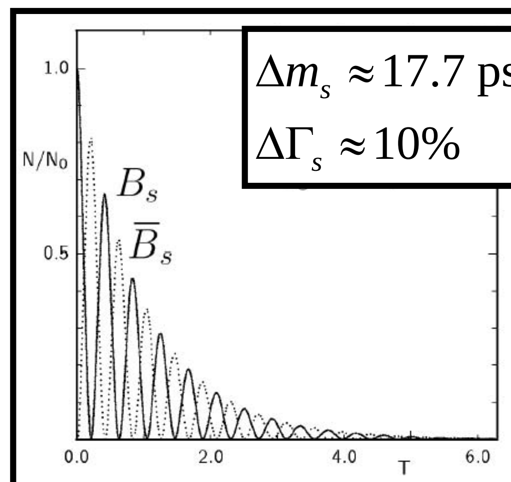
$$\frac{q}{p} = \sqrt{\frac{M_{12}^{*}-\frac{i}{2}\Gamma_{12}^{*}}{M_{12}-\frac{i}{2}\Gamma_{12}}}$$

$$\Delta m = m_H - m_L = 2|M_{12}|$$
$$\Delta\Gamma = \Gamma_L - \Gamma_H = 2|\Gamma_{12}|$$



$\Delta m_d \approx 0.505\ \mathrm{ps}^{-1}$

$\Delta\Gamma_d \approx 0$



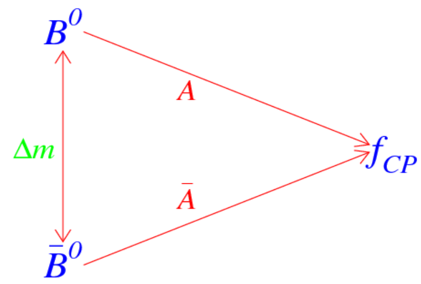
$\Delta m_s \approx 17.7\ \mathrm{ps}^{-1}$

$\Delta\Gamma_s \approx 10\%$

B mixing and CP violation

Decay amplitudes of flavour states decaying to the same final state *f*

$$A_f = \langle f | H | B^0 \rangle \quad \overline{A_f} = \langle f | H | \overline{B^0} \rangle$$



One can define $\lambda_f = \frac{q}{p} \frac{\overline{A_f}}{A_f}$

$$\tau \equiv 1/\Gamma$$

$$x \equiv \Delta m/\Gamma$$

$$y \equiv \Delta\Gamma/2\Gamma$$

Time dependence of decay rate for initially pure flavour states:

$$\Gamma_f \equiv \left|\langle f | H | B^0(t) \rangle\right|^2 = \frac{1+|\lambda_f|^2}{2} |A_f|^2 e^{-t/\tau} \left[\cosh yt/\tau + A_{\Delta f} \sinh yt/\tau + C_f \cos xt/\tau - S_f \sin xt/\tau\right]$$

$$\overline{\Gamma}_f \equiv \left|\langle f | H | \overline{B}^0(t) \rangle\right|^2 = \frac{1+|\lambda_f|^2}{2} \left|\frac{p}{q} A_f\right|^2 e^{-t/\tau} \left[\cosh yt/\tau + A_{\Delta f} \sinh yt/\tau - C_f \cos xt/\tau + S_f \sin xt/\tau\right]$$

$$S_f \equiv \frac{2\,\mathrm{Im}\lambda_f}{1+|\lambda_f|^2} \qquad C_f = \frac{1-|\lambda_f|^2}{1+|\lambda_f|^2} \qquad A_{\Delta f}^2 + S_f^2 + C_f^2 = 1$$

B mixing and CP violation

CP Violation $\rightarrow\ \Gamma_f \neq \overline{\Gamma}_f$

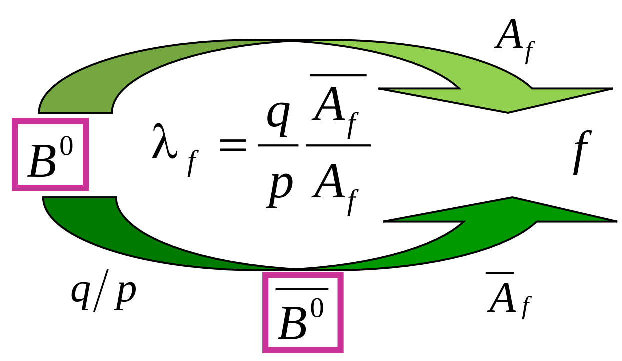
- **CPV in Decay (direct):** $B^0 \rightarrow f \quad \neq \quad \overline{B^0} \rightarrow \overline{f}$ $\quad \boxed{\left|\dfrac{\overline{A}_{\overline{f}}}{A_f}\right| \neq 1}$

- **CPV in Mixing:** $B^0 \rightarrow \overline{B^0} \quad \neq \quad \overline{B^0} \rightarrow B^0$ $\quad \boxed{\left|\dfrac{q}{p}\right| \neq 1} \quad \boxed{\mathrm{Im}\left\{\Gamma_{12}^{*} M_{12}\right\} \neq 0}$

- **CPV in Interference between mixing and decay:**

$$\boxed{\left|\lambda_f\right| = 1, \quad \mathrm{Im}\left\{\lambda_f\right\} \neq 0}$$

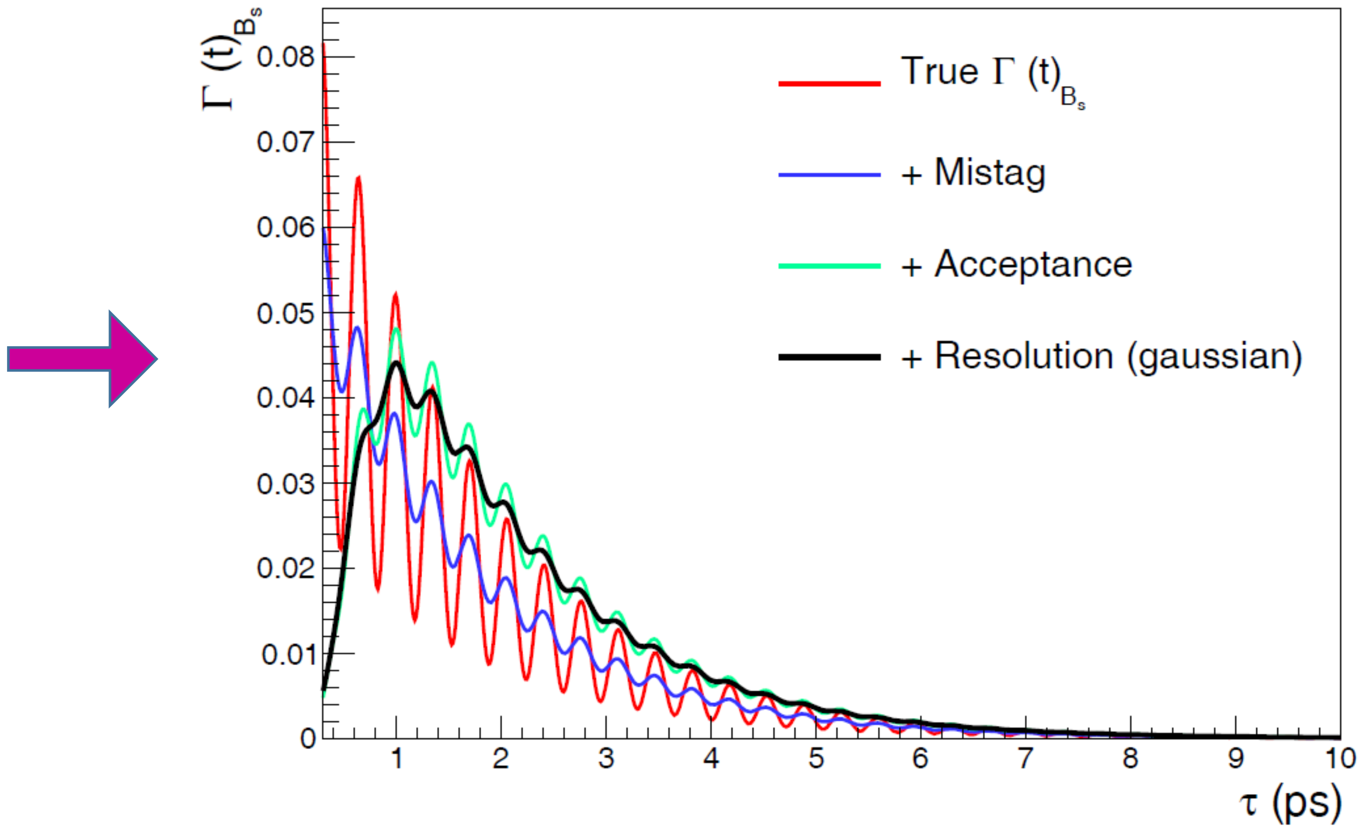
$$\boxed{A_f^{CP}(t) = \frac{\Gamma_f(t) - \overline{\Gamma}_f(t)}{\Gamma_f(t) + \overline{\Gamma}_f(t)} = \frac{-C_f \cos(\Delta m t) + S_f \sin(\Delta m t)}{\cosh(\Delta\Gamma t/2) + A_{\Delta f} \sinh(\Delta\Gamma t/2)}}$$



B mixing and CP violation

- Experimental effects:

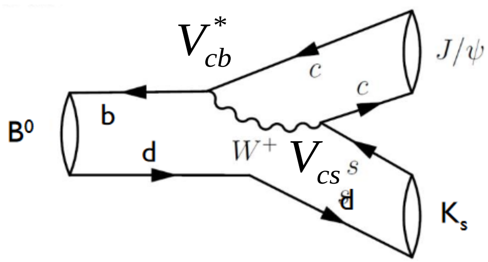
Dilution of the oscillation (lost of sensitivity of the oscillation parameters) due to reconstruction effects

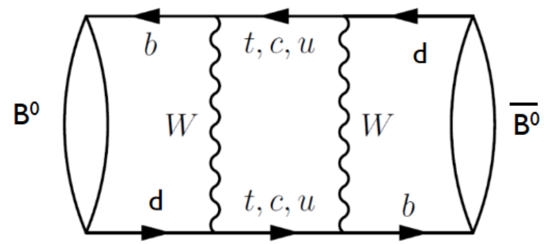


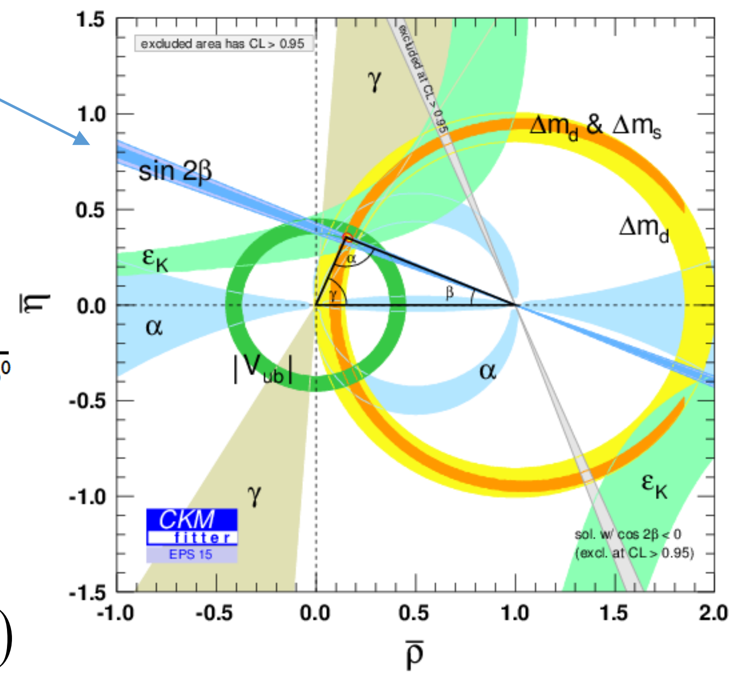
B mixing and CP violation

Measuring the β angle of the CKM triangle

Golden channel: $B^0 \to J/\psi\, K^0$







$$A_f^{CP}(t) = \frac{\Gamma_f(t) - \overline{\Gamma}_f(t)}{\Gamma_f(t) + \overline{\Gamma}_f(t)} = \frac{-C_f\cos(\Delta m t) + S_f \sin(\Delta m t)}{\cosh(\Delta\Gamma t/2) + A_{\Delta f}\sinh(\Delta\Gamma t/2)}$$

$$S_f \equiv \frac{2\,\mathrm{Im}\lambda_f}{1+|\lambda_f|^2}$$

$$\sim\ -C_f\cos(\Delta m t) + S_f\sin(\Delta m t)$$

$$\mathrm{Im}\{\lambda_f\} = \mathrm{Im}\{-e^{-2i\beta}\} = \sin 2\beta$$

C_f ~0 in the SM

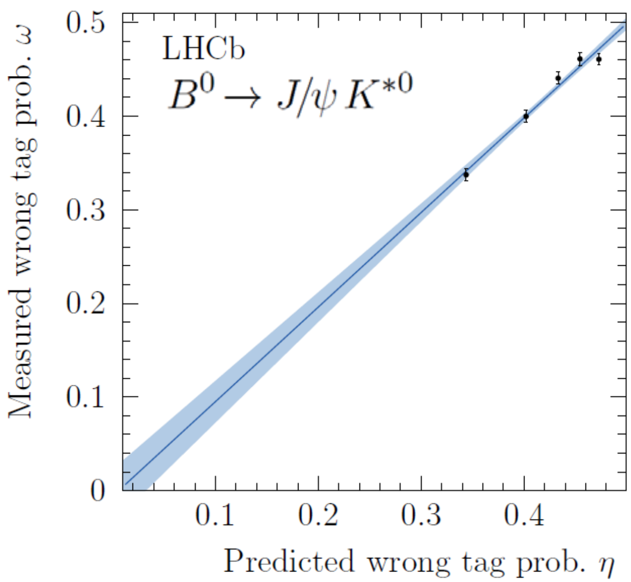
$$\boxed{A_f^{CP}(t) \sim \sin 2\beta \sin(\Delta m t)}$$

B mixing and CP violation

Count number of tagged signal events reconstructed as function of time

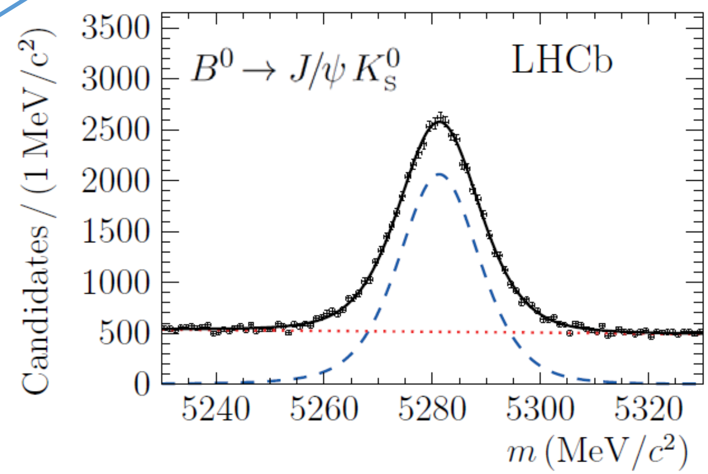
$$A_{CP}(t) = \frac{N\left(B_{tag} = B^0\right) - N\left(B_{tag} = \overline{B}^0\right)}{N\left(B_{tag} = B^0\right) + N\left(B_{tag} = \overline{B}^0\right)} \approx \pm\left\{(1-2\omega)\times\sin2\beta\times\sin(\Delta m\Delta t)\right\}\otimes \mathrm{R}(\Delta t)$$

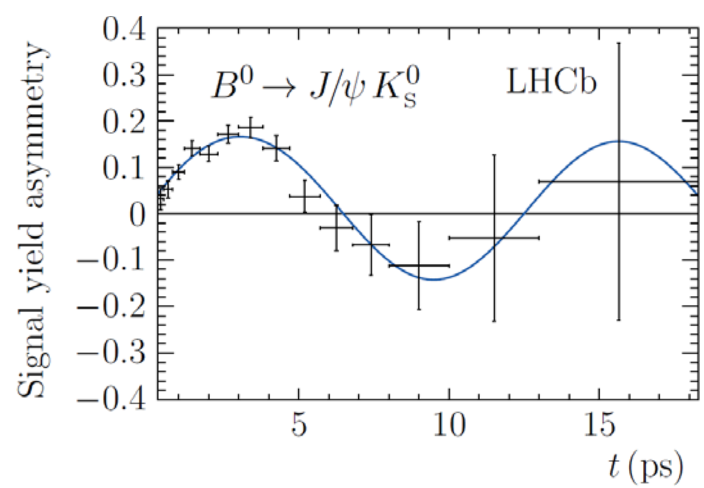
Flavour mistag calibrated using a control sample flavour specific ($K^{*0} \to K^-\pi^+$)



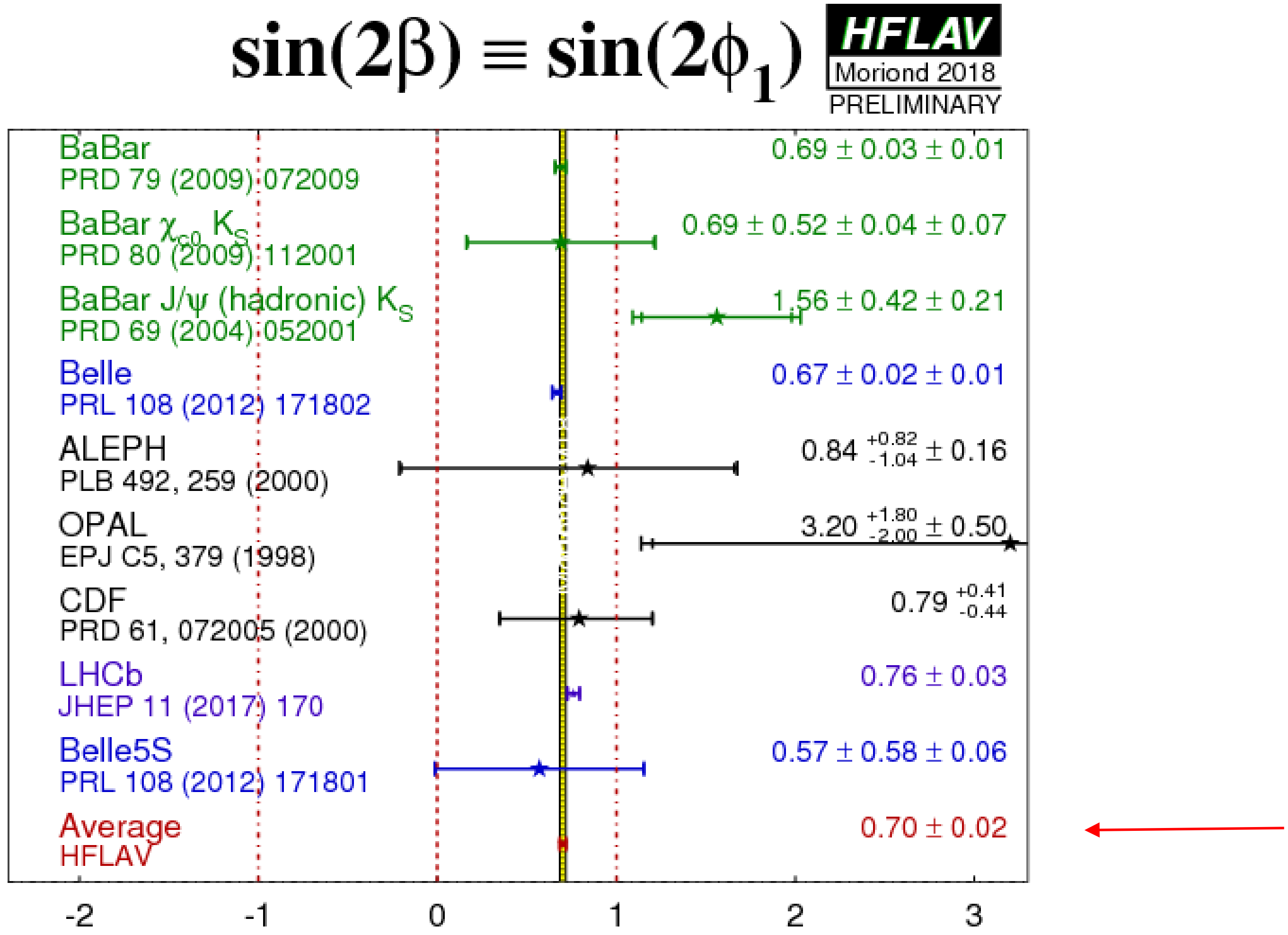
$\omega = (35.62 \pm 0.12)\,\%$

$\varepsilon_{\mathrm{tag}} = (36.54 \pm 0.14)\,\%$





B mixing and CP violation



B mixing and CP violation

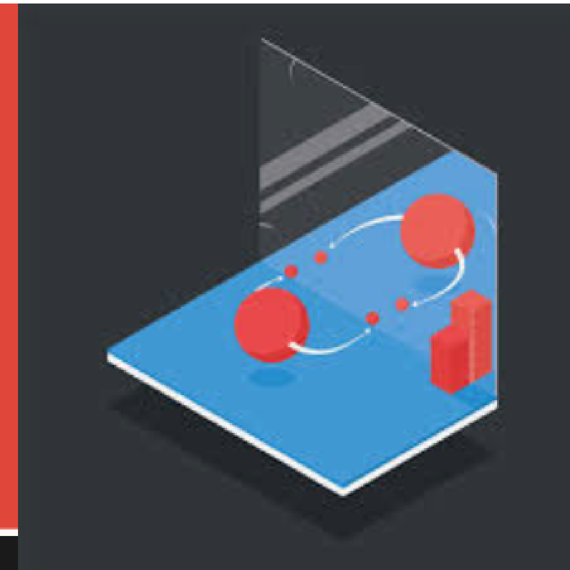
New results (March 2019):

symmetry
dimensions of particle physics
topics
follow
A joint Fermilab/SLAC publication

LHCb discovers matter-antimatter asymmetry in charm quarks

03/21/19 | By Sarah Charley

A new observation by the LHCb experiment finds that charm quarks behave differently than their antiparticle counterparts.



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News › Press release › Topic: Physics

Voir en français

LHCb sees a new flavour of matter–antimatter asymmetry

The LHCb collaboration has observed a phenomenon known as CP violation in the decays of a particle known as a D0 meson for the first time

21 MARCH, 2019

B mixing and CP violation

A lot of charm at LHCb: σ(pp→c$\bar{c}$) ~ 20σ(pp→b$\bar{b}$)

we can measure the asymmetry:

$$A_{h^+h^-} = \frac{N(D^0 \to h^+h^-) - N(\bar{D}^0 \to h^+h^-)}{N(D^0 \to h^+h^-) + N(\bar{D}^0 \to h^+h^-)}$$

Experimentally:

$$A_{h^+h^-} = A_{CP}(h^+h^-) + A_D + A_P$$

Detection asymmetry **Production asymmetry**

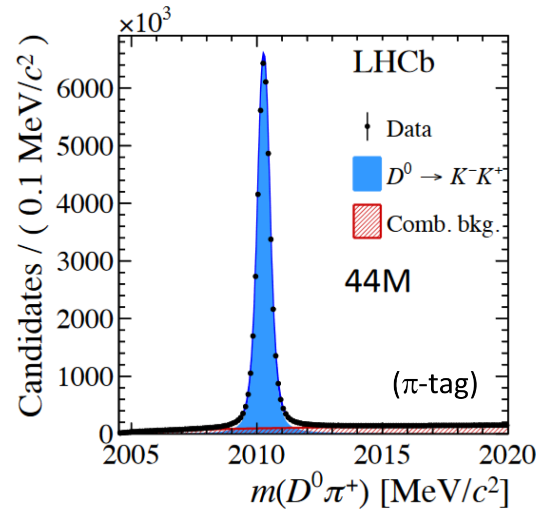
By taking differences detector and production asym. cancel :

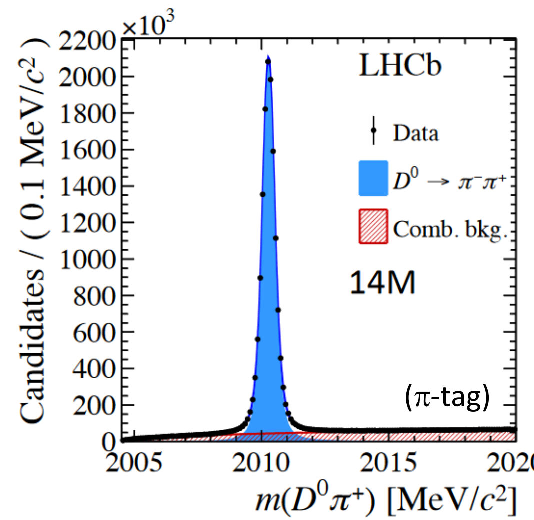
$$\Delta A_{CP} = A_{K^+K^-} - A_{\pi^+\pi^-}$$
$$= A_{CP}(K^+K^-) - A_{CP}(\pi^+\pi^-)$$

Using D^0 mesons coming from D^{*+} (π-tag) and B mesons (μ-tag)

$$\boxed{\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}}$$

First observation of CP violation in charm decays at 5.3σ !!



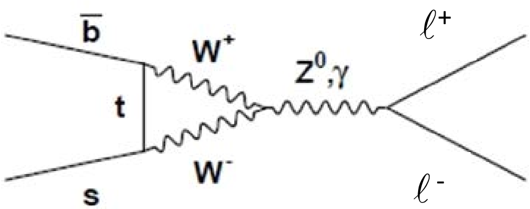


[PRL122(2019)211803]

Rare decays: leptonic

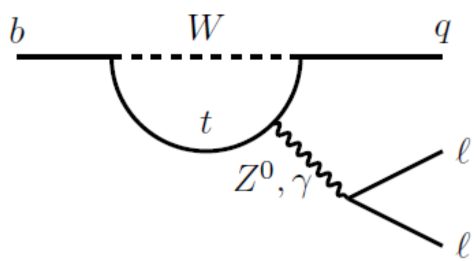
- *b→s,d* quark transitions are Flavor Changing Neutral Currents (FCNCs),
  → in the SM they only can occur through loops (*penguin and box diagrams*), excellent probe for physics beyond the SM

leptonic



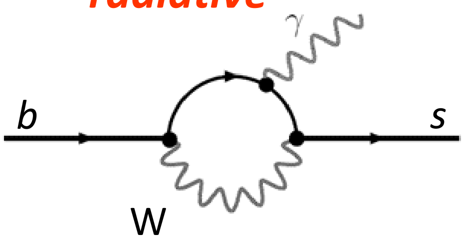
BR ~ 10^{-9}

semileptonic



BR ~ 10^{-7}

radiative



BR ~ 10^{-5}

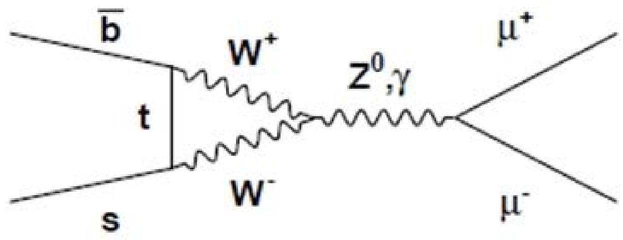
Experimentally → leptons/photons with high transverse momenta

Theoretically → observables can be calculated in terms of Wilson coefficients

Ex: $\Gamma\left(B_s^0 \to \mu^+\mu^-\right) \sim \frac{G_F^2\alpha^2}{64\pi^3} m_{Bs}^2 f_{Bs}^2 \left|V_{tb}V_{ts}\right|^2 \left|2m_\mu C_{10}\right|^2$

Hadronic uncertainties in decay constants or form factors

Rare decays: leptonic $B_s \to \mu^+\mu^-$

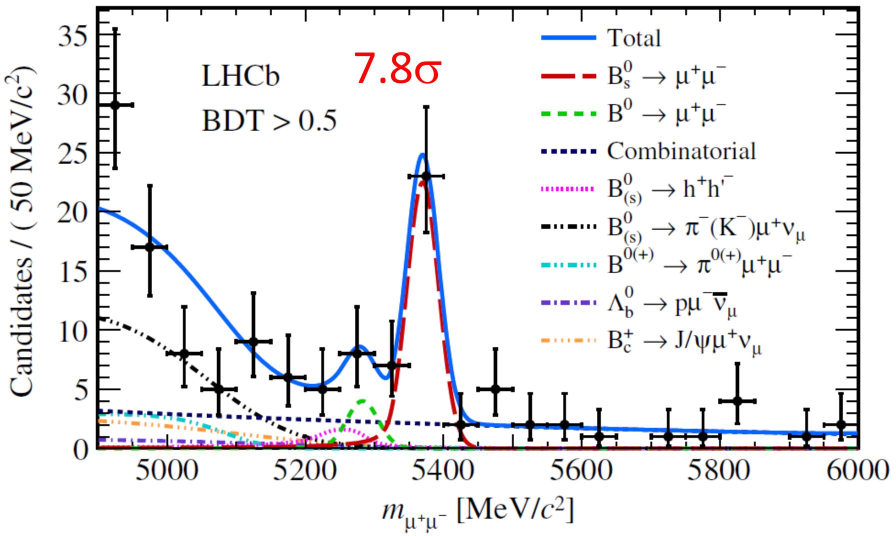


- Very rare decay:
FCNC and helicity suppressed
BR_{SM} = 3.66(23) x 10^{-9}

- Searched for over the last 30 years, observed by LHCb and CMS
[Nature 522 (2015) 68]

- Updated analysis by LHCb, including Run2 data
[PRL 118 (2017) 191801]

4.4 fb^{-1} [PRL 118 (2017) 191801]



$$\boxed{\mathcal{B}(B_s^0 \to \mu^+\mu^-) = (3.0 \pm 0.6\,^{+0.3}_{-0.2}) \times 10^{-9}}$$

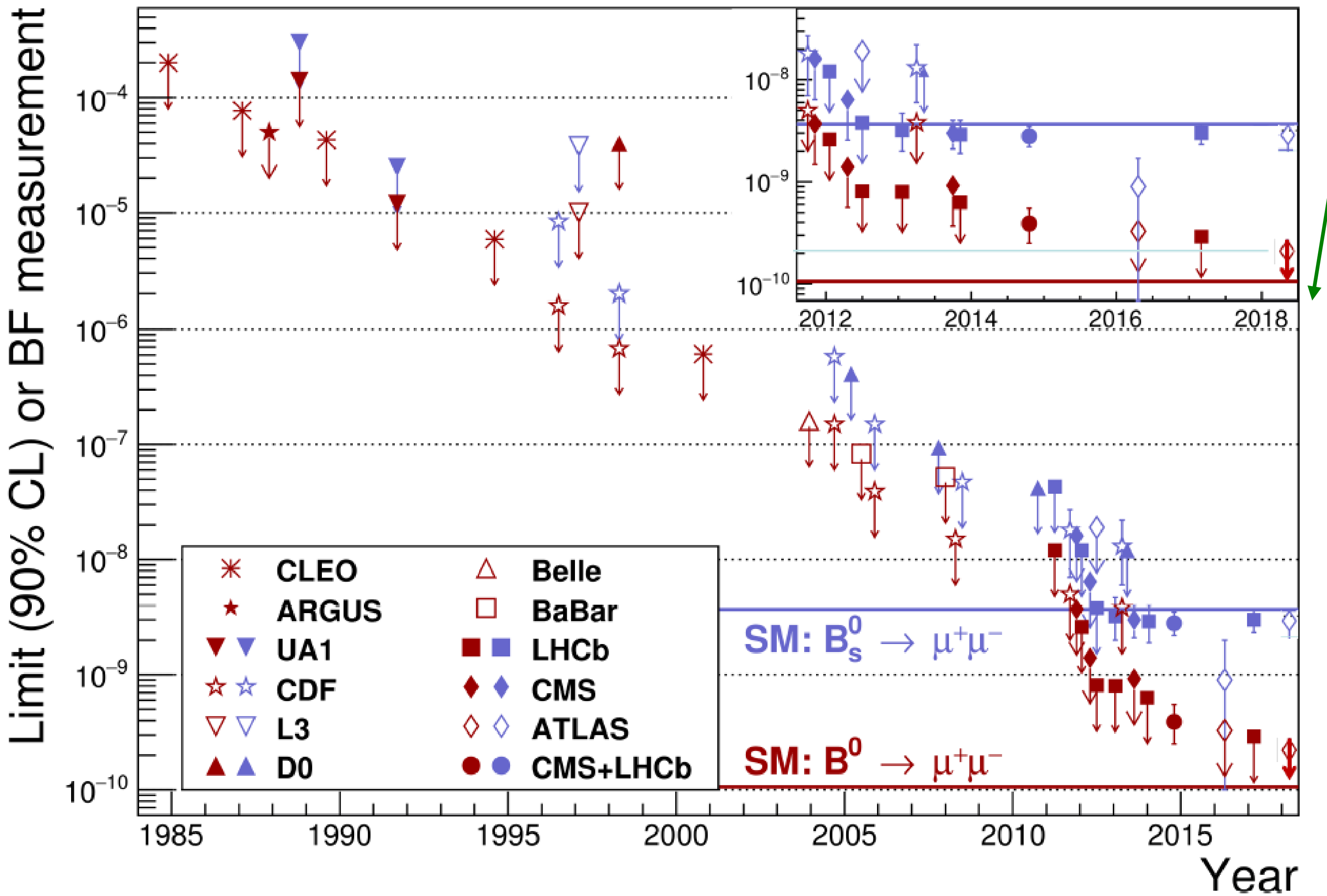
$$\boxed{\mathcal{B}(B^0 \to \mu^+\mu^-) < 3.4 \times 10^{-10} \text{ at } 95\%}$$

→ **Measurements in agreement with the SM**

→ Theoretical uncertainties ($f_{B(s)}$, V_{CKM}) well below statistical uncertainty

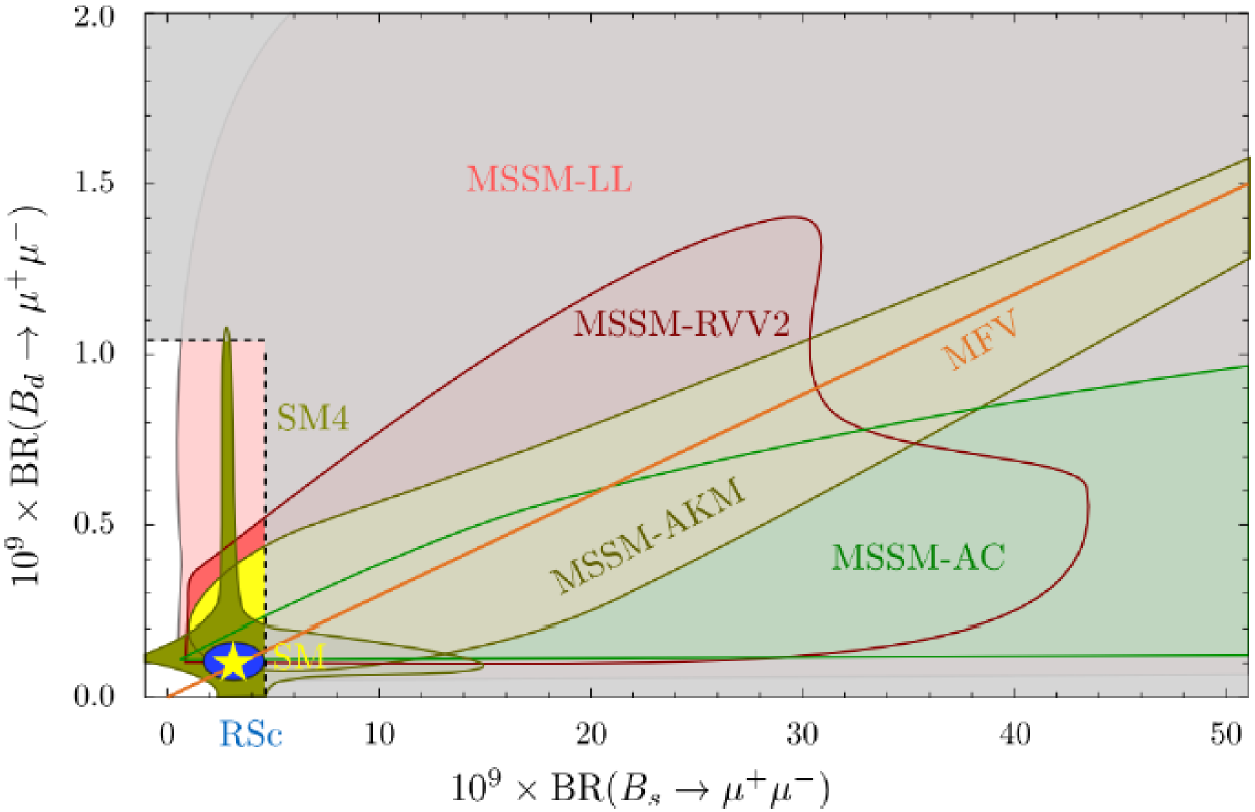
Rare decays: leptonic $B_s \to \mu^+\mu^-$

We are here!



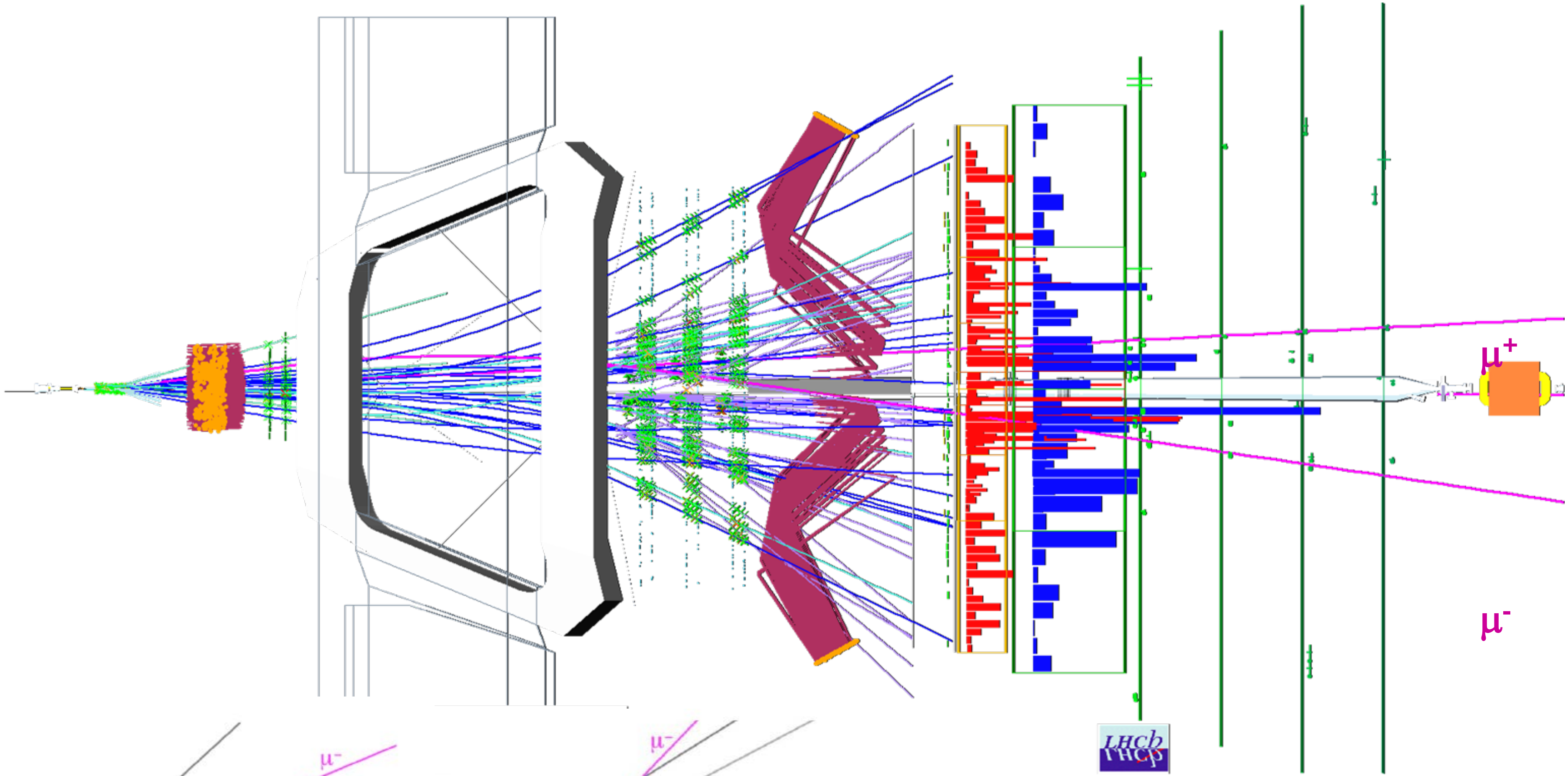
Rare decays: leptonic $B_s \to \mu^+\mu^-$

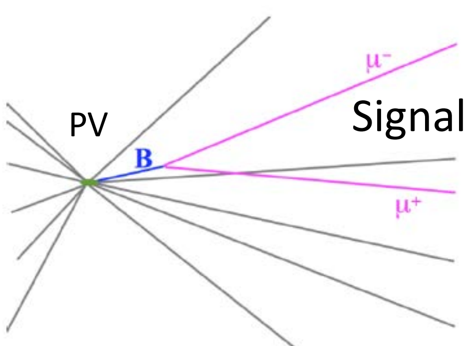
- Known as “the New Physics killer”:

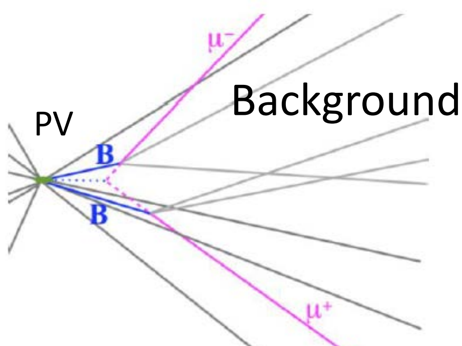


Rare decays: leptonic $B_s \to \mu^+\mu^-$

$B_s \to \mu^+\mu^-$ event



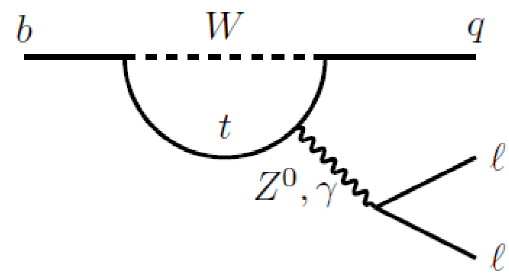




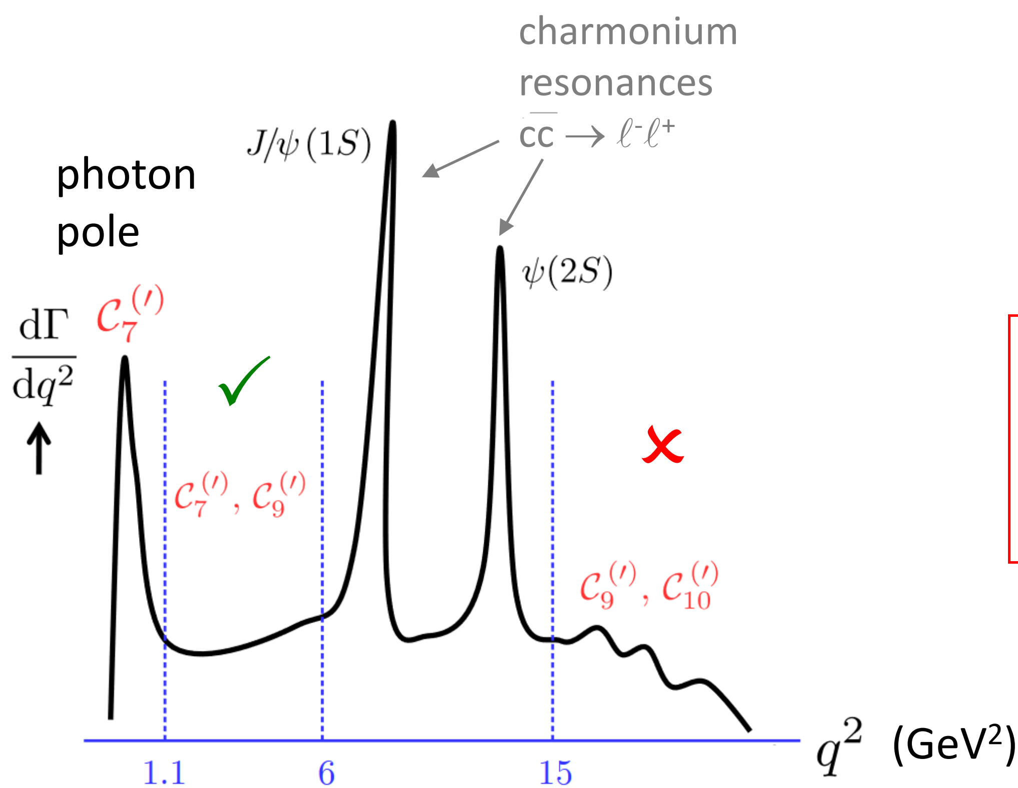
Rare decays: semileptonic $B \to K^{*}\mu^{+}\mu^{-}$

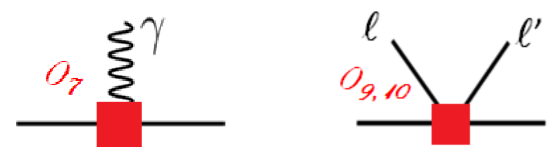
Differential decay width: $d\Gamma/dq^2$

Each q^2 region probes different processes



$q^2 = (p_{\ell^+} + p_{\ell^-})^2$





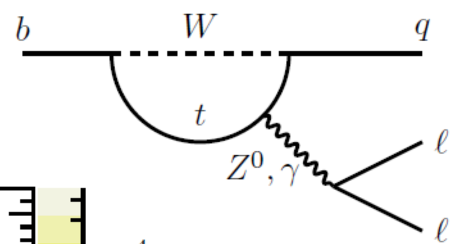
SM values ($\mu = m_b$): $C_7 \sim -0.33$
$C_9 \sim 4.27$
$C_{10} \sim -4.17$

(Everything else small or negligible)

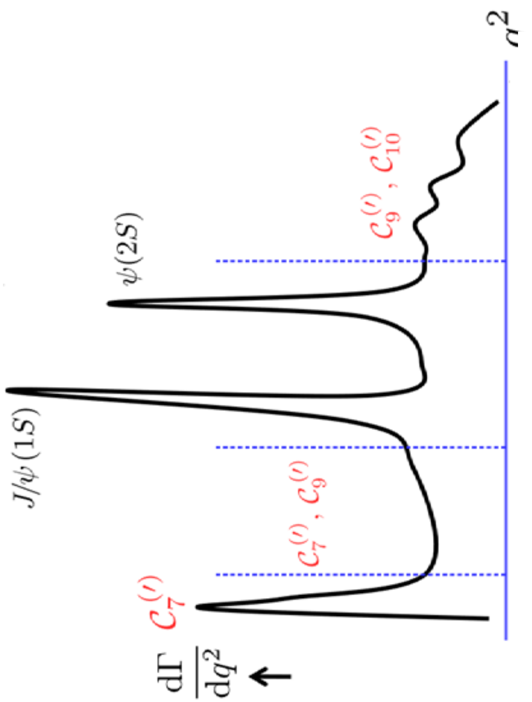
$C_i = C_i^{SM} + C_i^{NP}$

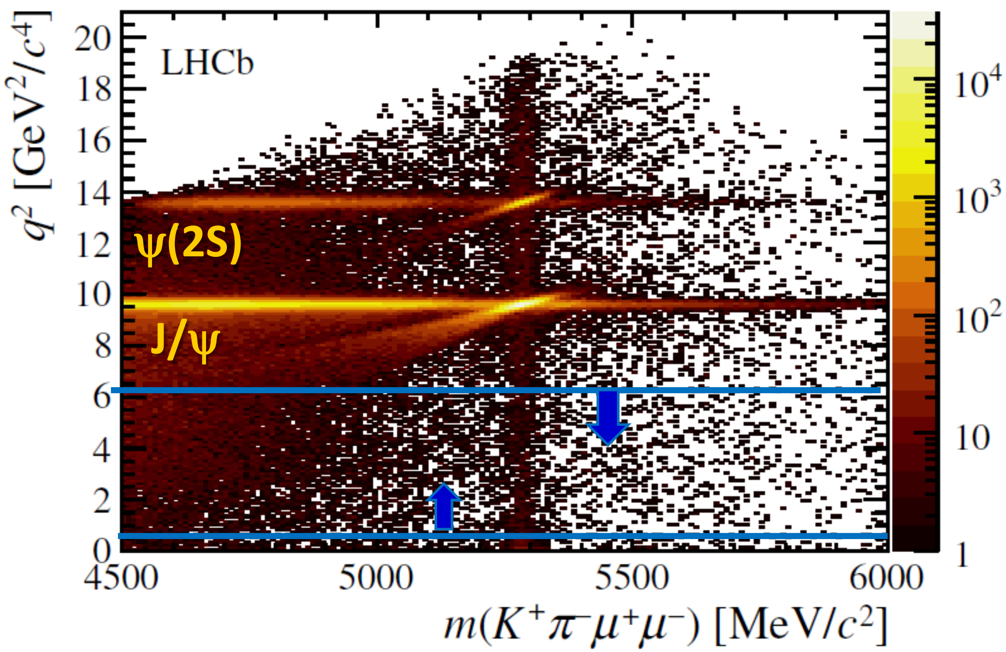
(Primed $C'_i \to$ right handed currents: suppressed in SM)

Rare decays: semileptonic B $\rightarrow K^{*}\mu^{+}\mu^{-}$



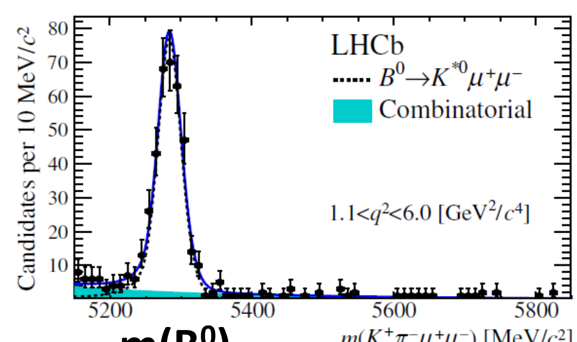
B mass vesus q^2 for $B^0 \rightarrow K^{*0}\mu^{+}\mu^{-}$





[JHEP 08 (2017) 055]

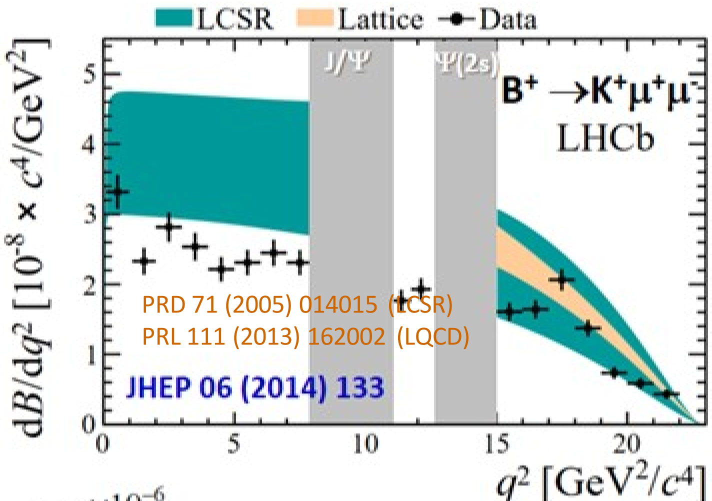
$B^0 \rightarrow K^{*0}\mu^{+}\mu^{-}$

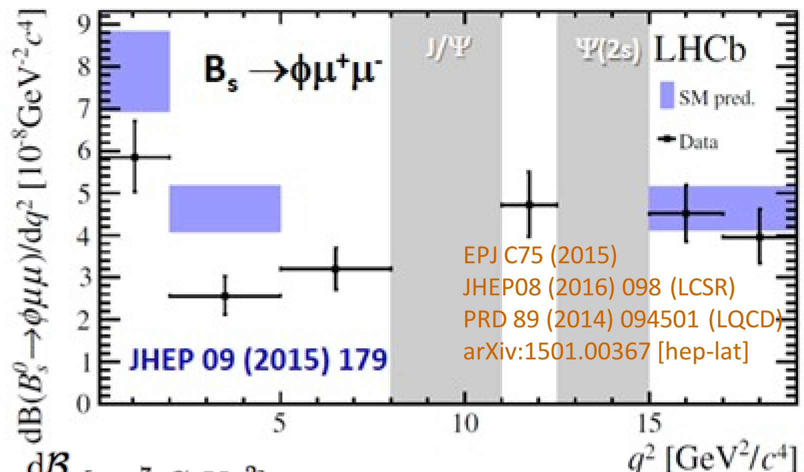


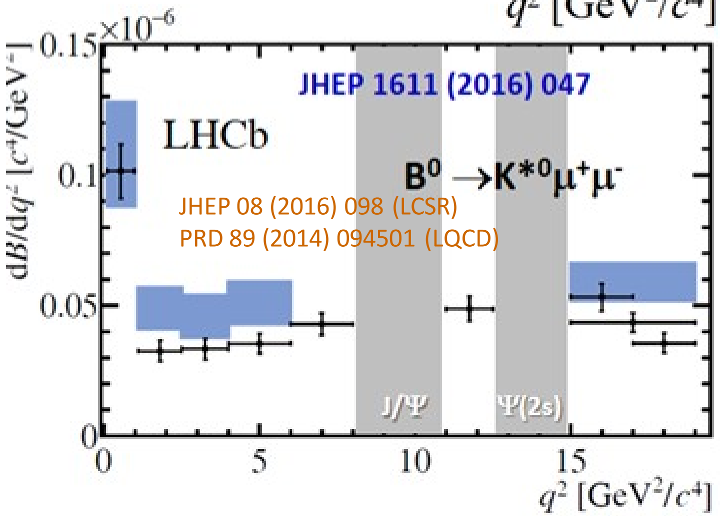
m(B^0)

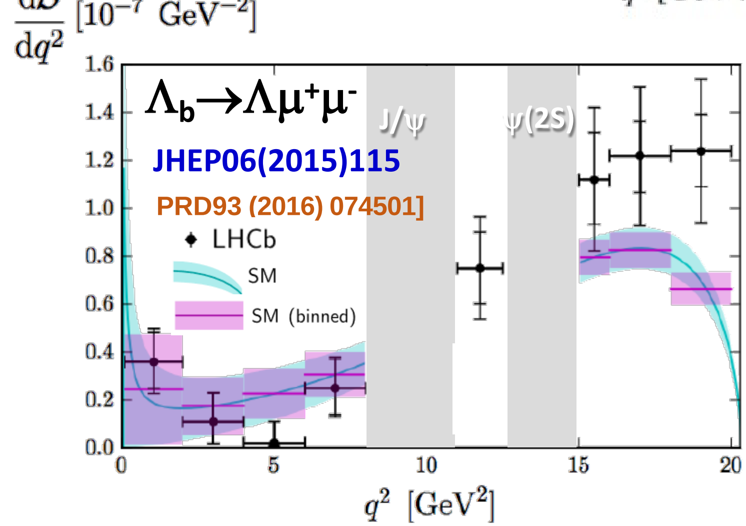
Rare decays: semileptonic

- Differential decay width as function of $q^2 = m^2_{\mu\mu}$ at **LHCb**, using 3fb^{-1}





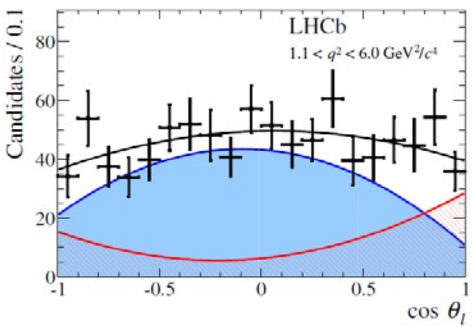


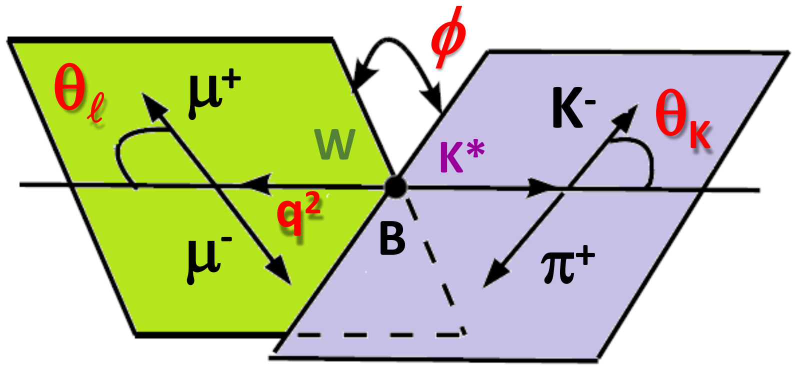


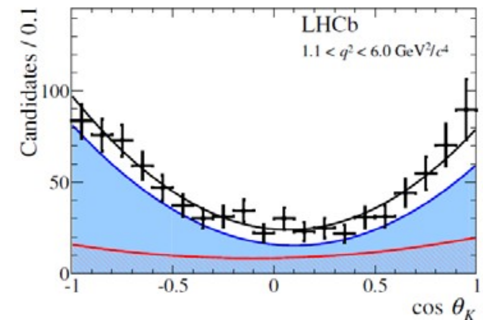
→ Smaller branching fractions than the SM predictions

Rare decays: semileptonic $B \to K^*\mu^+\mu^-$

- Angular distribution in $B \to K^*\ell^-\ell^+$: q^2 and three angles







$$\frac{1}{d\Gamma/dq^2}\frac{d^4\Gamma}{d\cos\theta_\ell\, d\cos\theta_K\, d\phi\, dq^2} = \frac{9}{32\pi}\Big[\frac{3}{4}(1-F_L)\sin^2\theta_K + F_L\cos^2\theta_K + \frac{1}{4}(1-F_L)\sin^2\theta_K\cos 2\theta_\ell$$
$$- F_L\cos^2\theta_K\cos 2\theta_\ell + S_3\sin^2\theta_K\sin^2\theta_\ell\cos 2\phi + S_4\sin 2\theta_K\sin 2\theta_\ell\cos\phi$$
$$+ S_5\sin 2\theta_K\sin\theta_\ell\cos\phi + S_6\sin^2\theta_K\cos\theta_\ell + S_7\sin 2\theta_K\sin\theta_\ell\sin\phi$$
$$+ S_8\sin 2\theta_K\sin 2\theta_\ell\sin\phi + S_9\sin^2\theta_K\sin^2\theta_\ell\sin 2\phi\Big]$$

$\rightarrow$ In the lepton massless limit there are **eight** independent observables:

F_L = fraction of the longitudinal polarization of the K*
$S_6 = 4/3\ A_{FB}$, the forward-backward asymmetry of the dimuon system
$S_{3,4,5,7,8,9}$ are the remaining CP-averaged observables

Rare decays: semileptonic $B \to K^* \mu^+ \mu^-$

!?

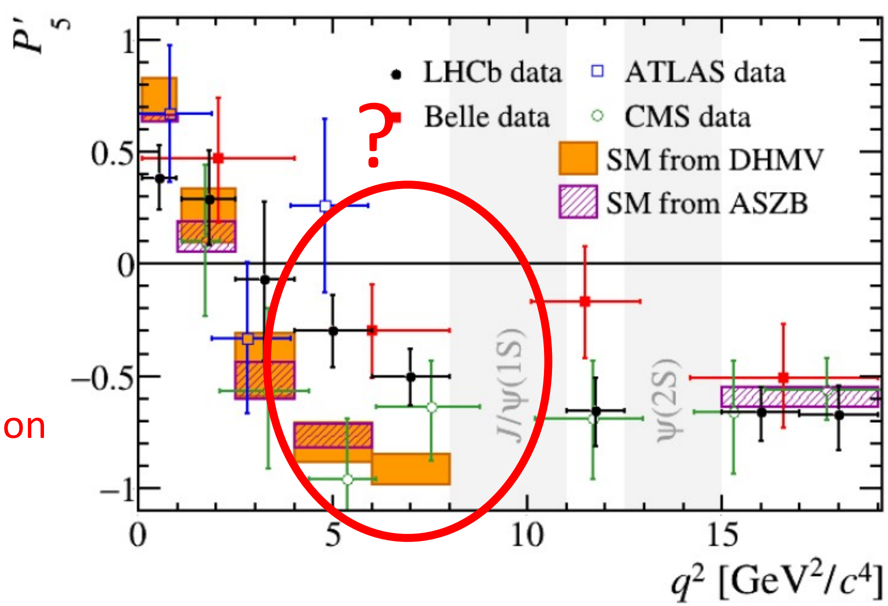
- These observables are also affected by hadronic uncertainties
- A new set of "optimized observables", with form factor cancellations can be defined: [Descotes-Genon et al, JHEP 05 (2013) 137]

$$P'_{i=4,5,6,8} = \frac{S_{j=4,5,7,8}}{\sqrt{F_L(1-F_L)}}$$

- These observable are functions of q^2 and the Wilson coefficients $\mathbf{C_i}$

Example: P'_5

3σ local deviation



- JHEP 02 (2016) 104
- PRL 118 (2017) 111801
- ATLAS-CONF-2017-023
- CMS-PAS-BPH-15-008

Rare decays: lepton flavour universality

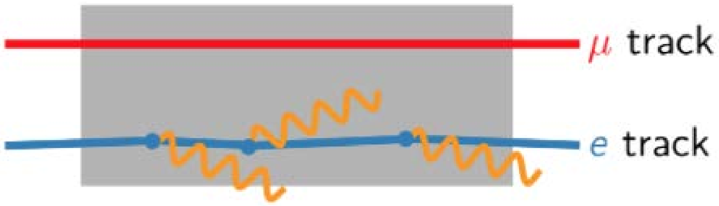
- In the SM all leptons are expected to behave in the same way

Test of lepton universality:

$$R_K = \frac{\mathcal{B}(B^+ \to K^+\mu^+\mu^-)}{\mathcal{B}(B^+ \to K^+e^+e^-)}$$ $= 1.000 + O(m_\mu^2/m_b^2)$

- Precise theory prediction due to cancellation of hadronic form factor uncertainties

- Challenge: bremsstrahlung by electrons

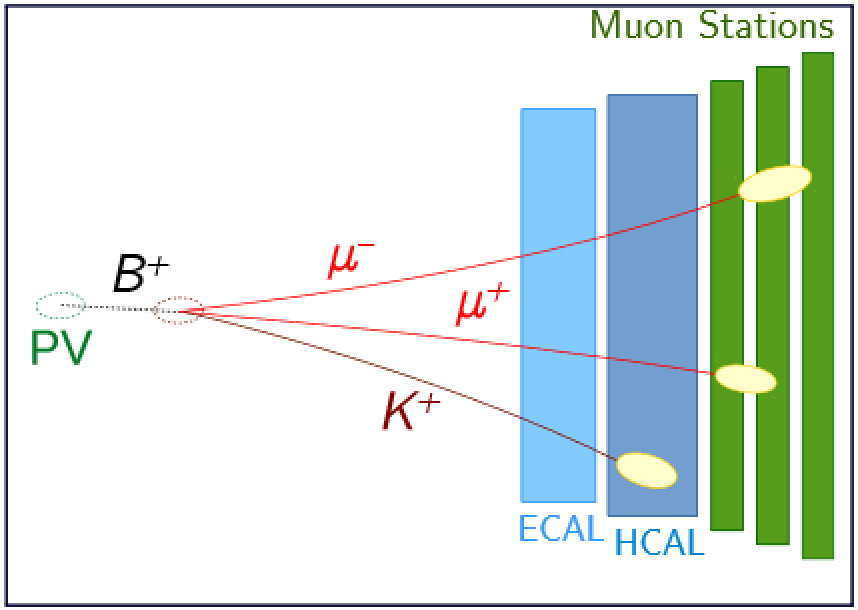


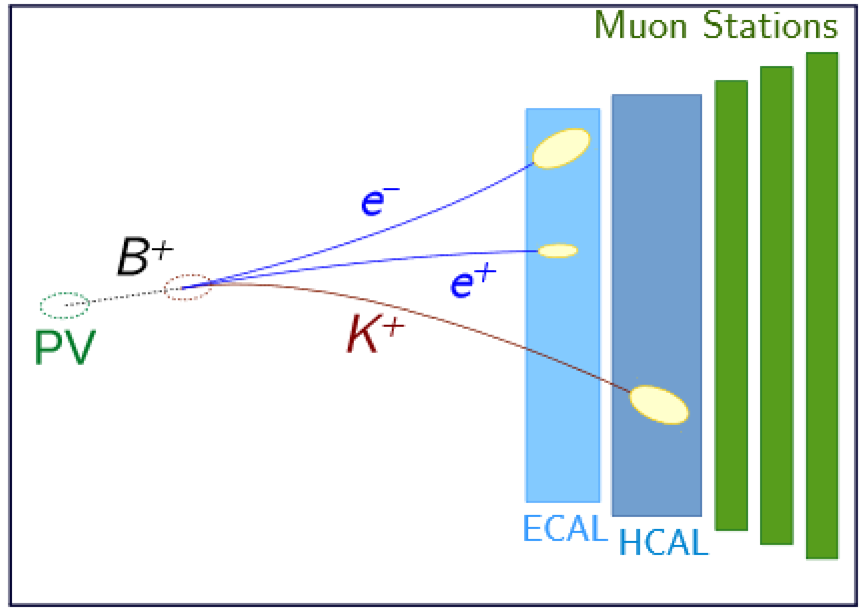
- Experimentally, we perform a double ratio to cancel systematic uncertainties

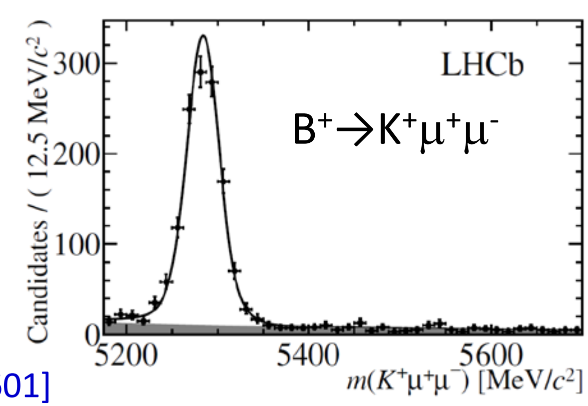
$$R_K = \frac{\mathcal{B}(B^+ \to K^+\mu^+\mu^-)}{\mathcal{B}(B^+ \to K^+J/\psi(\mu^+\mu^-))} \Bigg/ \frac{\mathcal{B}(B^+ \to K^+e^+e^-)}{\mathcal{B}(B^+ \to K^+J/\psi(e^+e^-))}$$

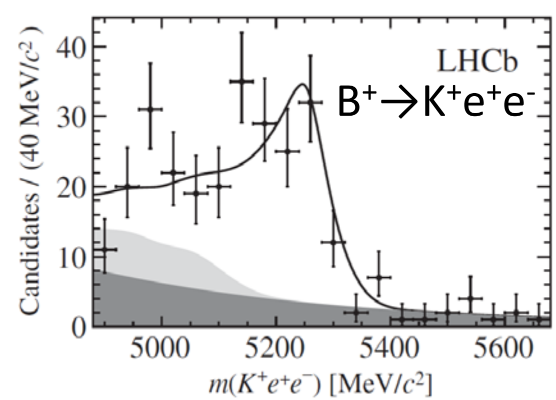
Rare decays: lepton flavour universality

B mass vesus q^2 for $B^+ \to K^+ \ell^+ \ell^-$





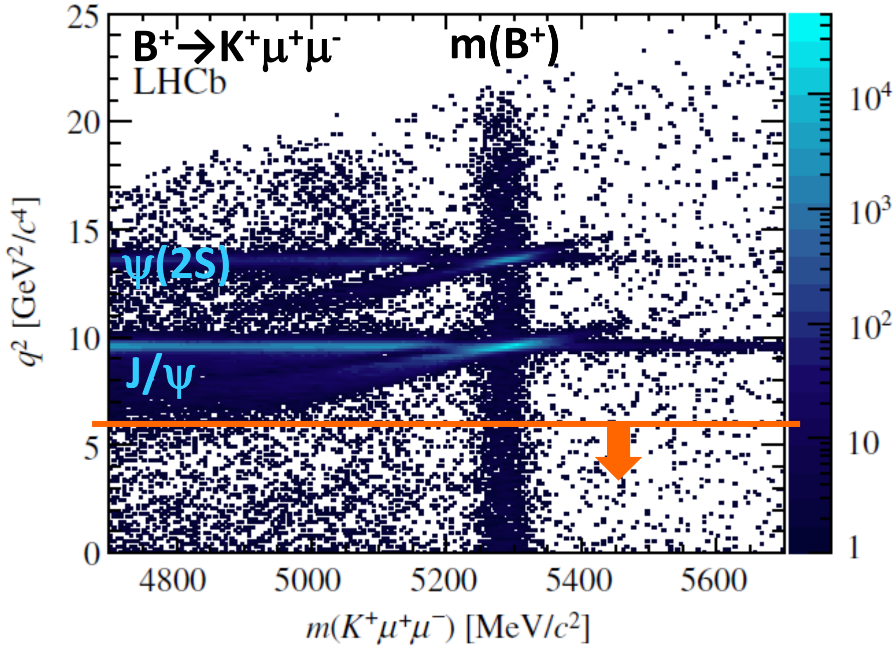


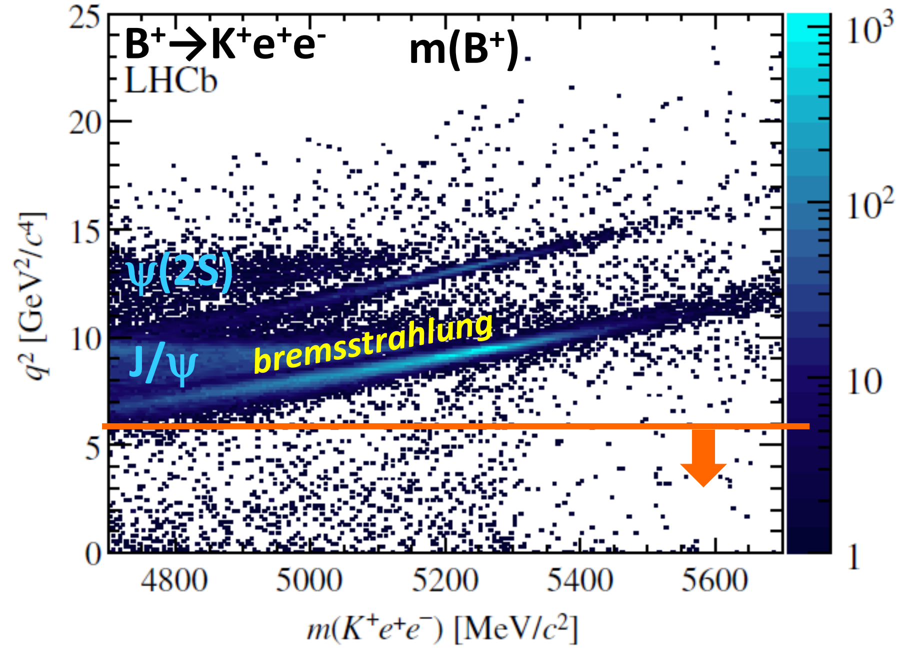


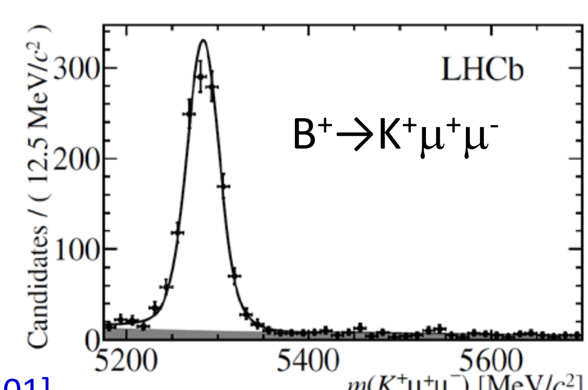
[PRL 113 (2014) 151601]

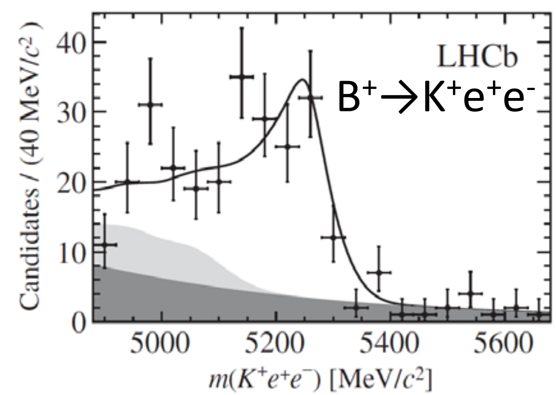
Rare decays: lepton flavour universality

B mass vesus q^2 for $B^+ \to K^+ \ell^+ \ell^-$





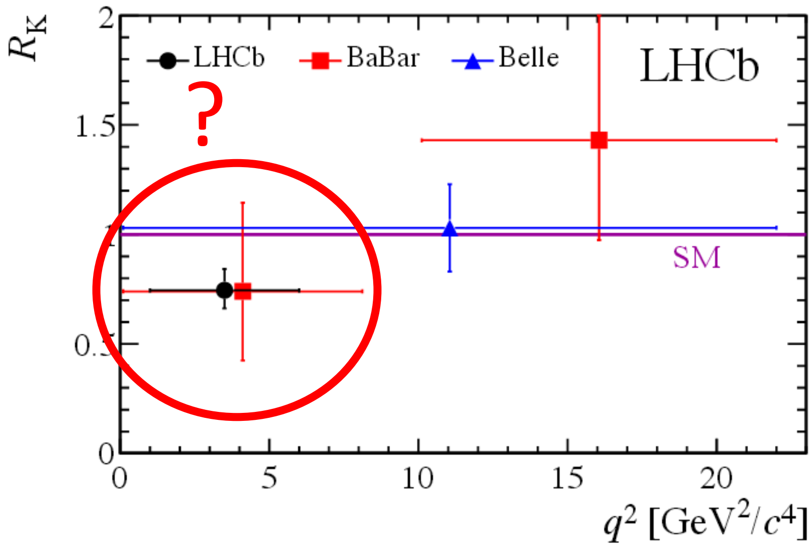


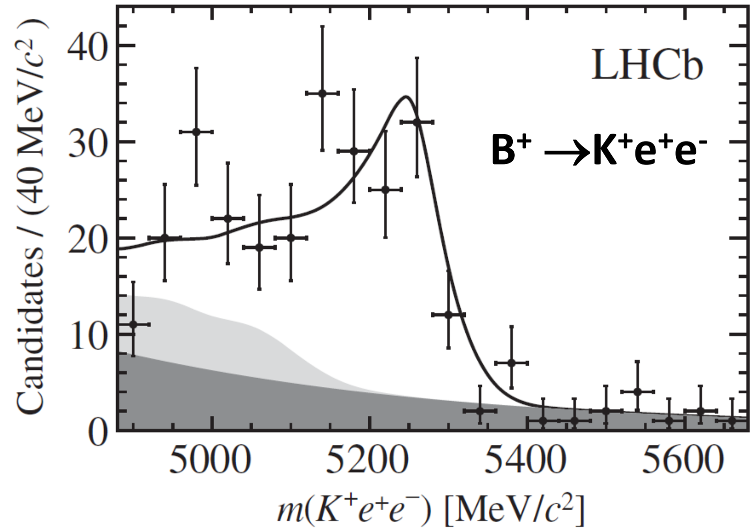


[PRL 113 (2014) 151601]

!? Rare decays: lepton flavour universality

Results with Run1 data: [PRL 113 (2014) 151601]





1 GeV < q^2 < 6 GeV

$$R_K = 0.745^{+0.090}_{-0.074}\,(\text{stat}) \pm 0.036\,(\text{syst})$$

$\rightarrow$ Consistent, but lower, than the SM at **2.6σ**

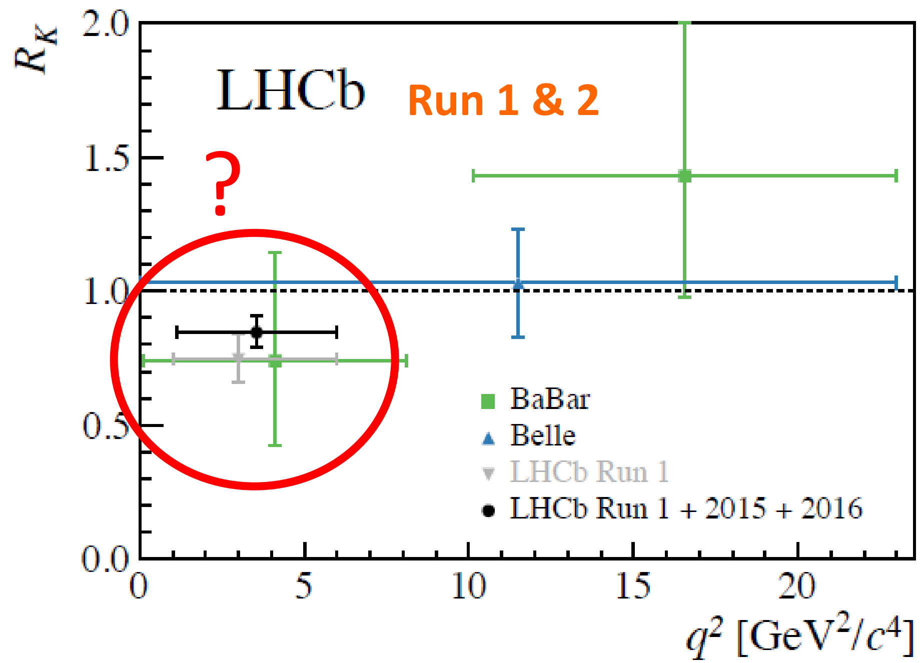
Rare decays: lepton flavour universality

New results (Moriond 2019):

Including partial sample of Run2 (2fb^{-1})

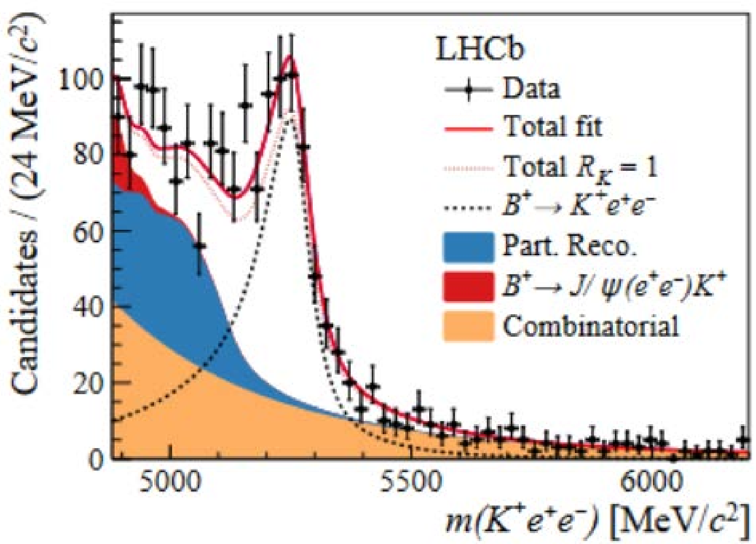
[*PRL* 122 (2019) 191801]

With improved reconstruction and re-optimized analysed estrategy



1.1 GeV < q^2 < 6 GeV

$$R_K = 0.846 \,^{+0.060}_{-0.054}\,(\text{stat.}) \,^{+0.016}_{-0.014}\,(\text{syst.})$$



→ Still consistent, lower, than the SM at **2.5σ**

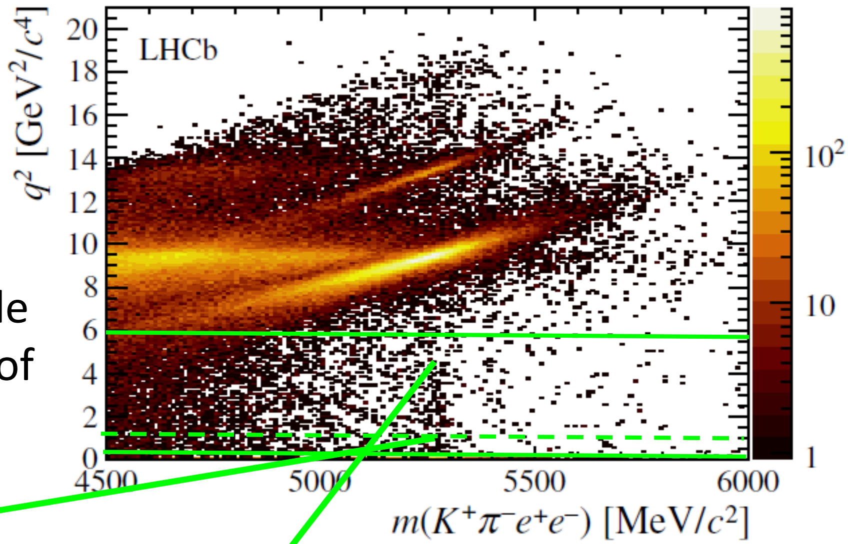
Not confirmed, not ruled out…

Rare decays: lepton flavour universality

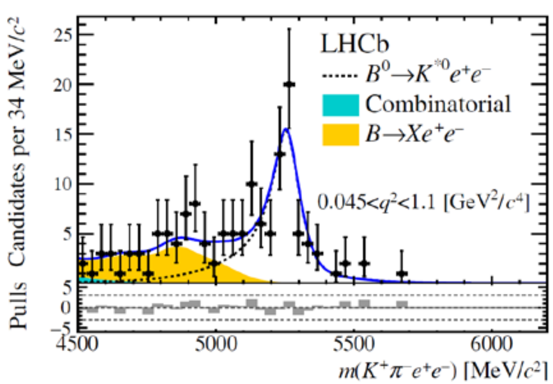
- Measurement in the $B \to K^{*}\mu^{+}\mu^{-}$ channel, R_{K*}:

[JHEP 08 (2017) 055]

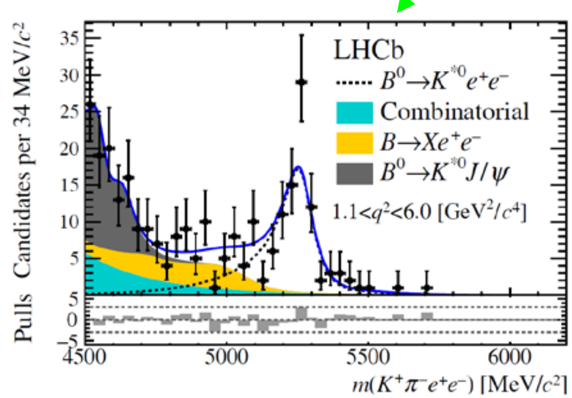
$$\mathcal{R}_{K^{*0}} = \frac{\mathcal{B}(B^0 \to K^{*0}\mu^+\mu^-)}{\mathcal{B}(B^0 \to K^{*0}e^+e^-)}$$



- Computed in two bins of q^2
  - [0.045, 1.1 GeV^2] avoiding the photon pole
  - [1.1, 6.0 GeV^2] avoiding the radiative tail of J/ψ modes



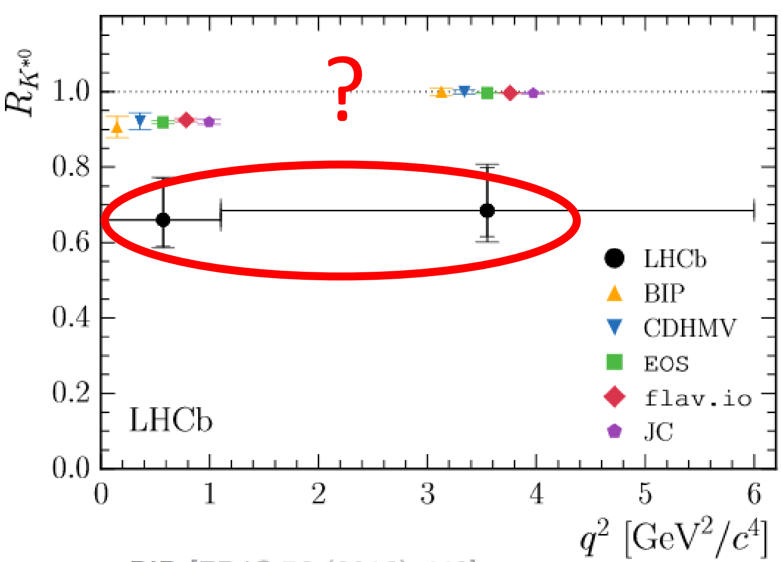
0.045 GeV < q^2 < 1.1 GeV

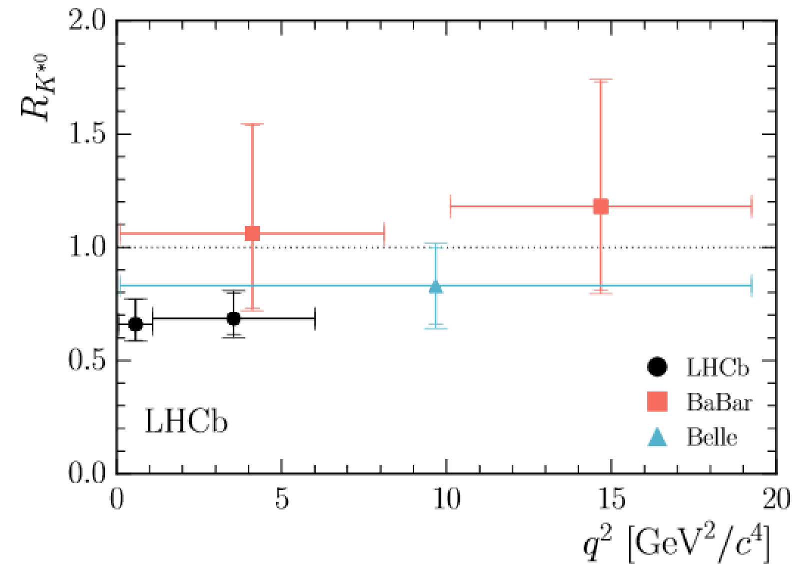


1.1 GeV < q^2 < 6 GeV

Rare decays: lepton flavour universality

- Results: [JHEP 08 (2017) 055]





- BIP [EPJC 76 (2016) 440]
- CDHMV [JHEP 04 (2017) 016]
- EOS [PRD 95 (2017) 035029]
- flav.io [EPJC 77 (2017) 377]
- JC [PRD 93 (2016) 014028]

- LHCb, JHEP 08 (2017) 055
- BaBar, PRD 86(2012) 032012
- Belle, PRL (2009) 171801

Low q^2 [0.045-1.1 GeV^2]: SM▼ = 0.922(22)

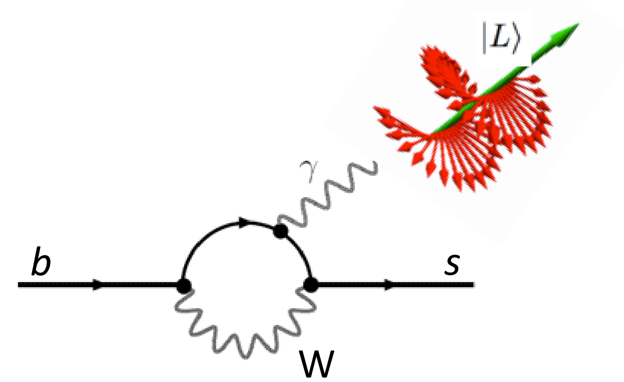
$$\boxed{R_{K^{*0}} = 0.66\ {}^{+\,0.11}_{-\,0.07}\,(\mathrm{stat}) \pm 0.03\,(\mathrm{syst})}$$

Central q^2: [1.1-6 GeV^2]: SM▼ = 1.000(6)

$$\boxed{R_{K^{*0}} = 0.69\ {}^{+\,0.11}_{-\,0.07}\,(\mathrm{stat}) \pm 0.05\,(\mathrm{syst})}$$

→ Consistent, but lower than the SM at **2.1-2.3σ** (low q^2) and **2.4-2.5σ** (central q^2)

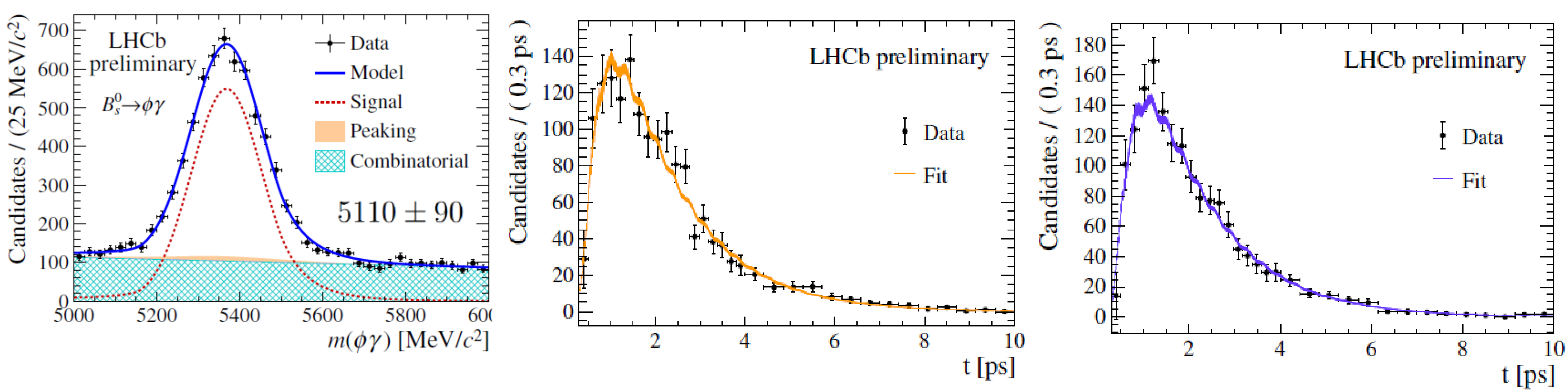
Rare decays: radiative



New results (Moriond 2019): [PRL 123 (2019) 081802]

- Time dependent distribution for $B_s \to \phi\gamma$ is sensitive to the photon polarization (photon is left-handed polarized in $b \to s$ transitions)

$$\Gamma_{B,\overline{B}}(t) = \mathcal{B}_0 e^{-\Gamma t}[\cosh(\frac{\Delta\Gamma}{2}t) - A^{\Delta} \sinh(\frac{\Delta\Gamma}{2}t) \pm C \cos(\Delta m\ t) \mp S \sin(\Delta m\ t)]$$



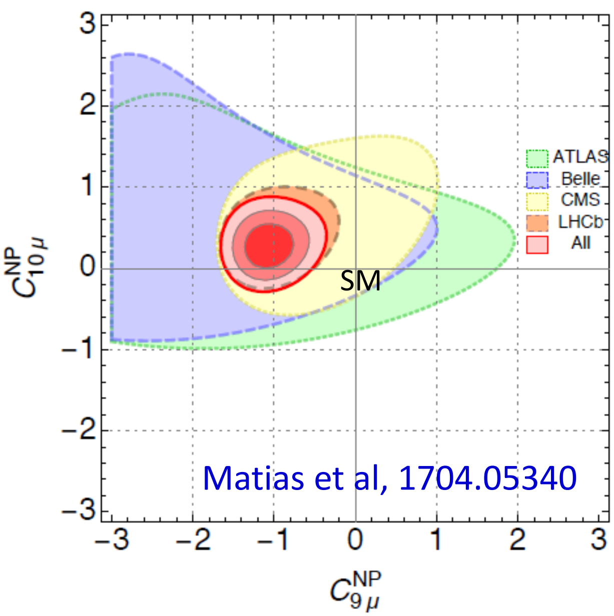
$$S_{\phi\gamma} = 0.43 \pm 0.30 \pm 0.11$$

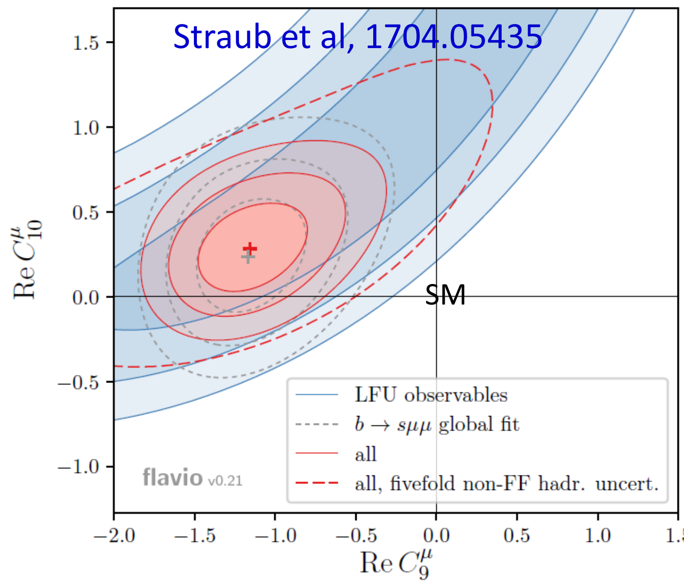
$$C_{\phi\gamma} = 0.11 \pm 0.29 \pm 0.11$$

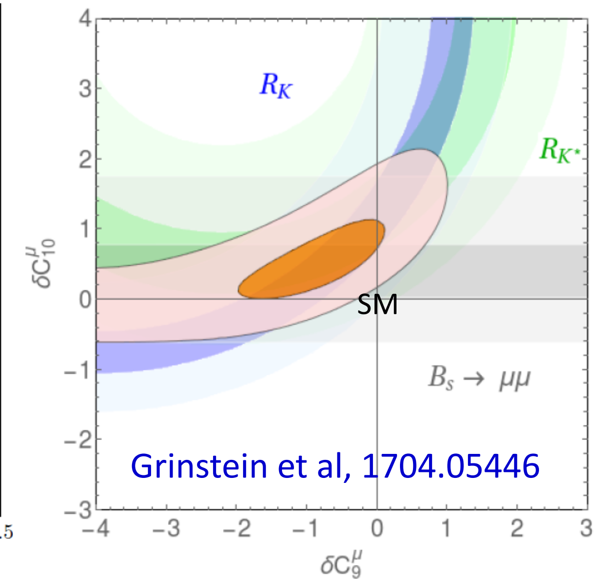
$$\mathcal{A}^{\Delta}_{\phi\gamma} = -0.67\,^{+0.37}_{-0.41} \pm 0.17$$

Rare decays: global fits

Global fits (some cases with more than 100 observables)



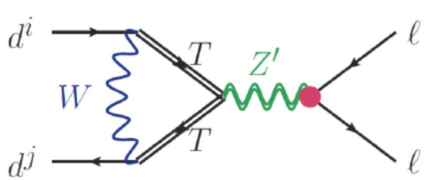


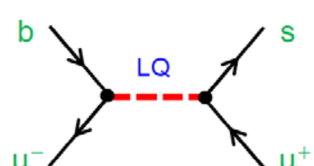


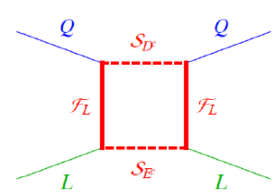
New Physics hypothesis preferred over SM by more than 4 - 5σ

Main effect on the $C_{9\mu}$ coefficient: **4.27^{SM} -1.1^{NP}**

Triggered models with Z', leptoquarks (LQ), new fermions and scalars….



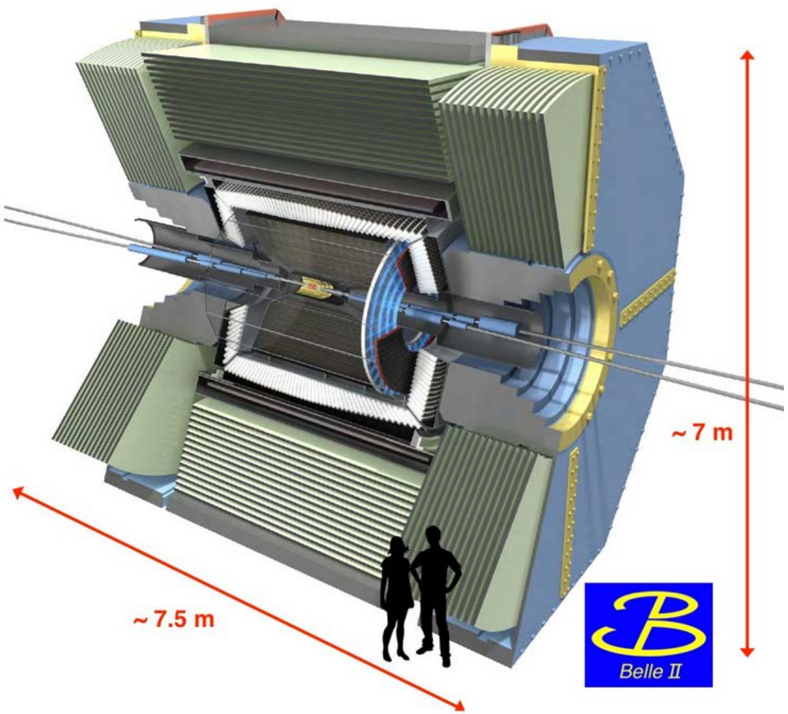




Future plans

Belle II

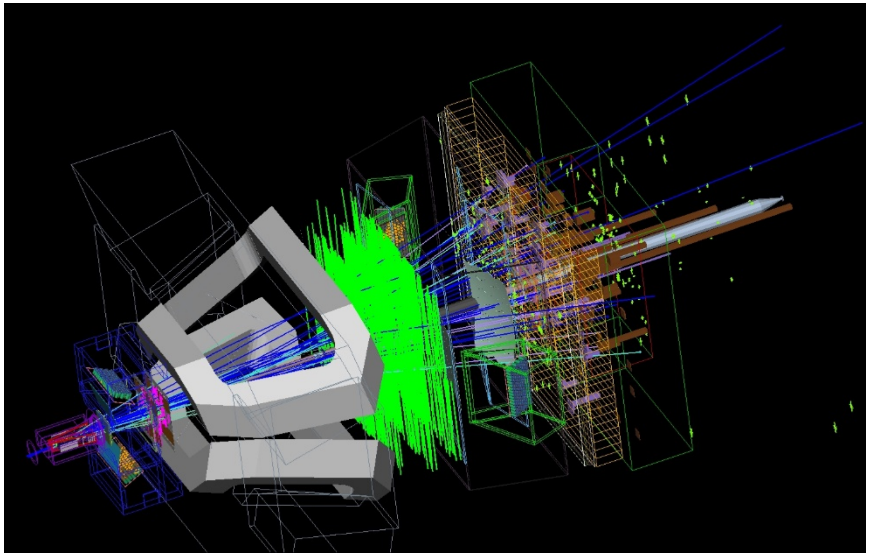
e+e- asymmetric collider (SuperKEK) at the $\Upsilon(4S)$ energy in Tsukuba, Japan



Just started !

LHCb Upgrade

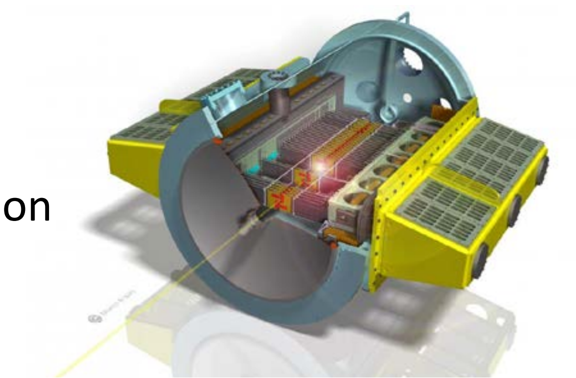
Improved trigger, improved detector



Starting in 2021

Future plans

- At present we did not find evidence for New Phyiscs, but some “anomalies” (let’s cross fingers!!)



- Standard Model deviations are expected to be small
- Most of the measurements are limited by the statistical precision

	LHC era		High-lumi LHC era		
	2010-2012	2015-2018	2020-2022	2025-2028	2030+
ATLAS & CMS	25 fb^{-1}	100 fb^{-1}	300 fb^{-1}	$\rightarrow$	3000 fb^{-1}
LHCb	3 fb^{-1}	8 fb^{-1}	23 fb^{-1}	46 fb^{-1}	100 fb^{-1}
Belle II		0.5 ab^{-1}	25 ab^{-1}	50 ab^{-1}	-

Remember that we have 10^{11} $b\bar{b}$ pairs/fb !

(At Belle II: 10^9 $B\bar{B}$ pairs/ab)

Future plans

Complementarity:

LHCb

ATLAS & CMS

- Rare decays: $B_{d,s} \to \mu\mu$
- B_s system
- b-baryons
- Spectroscopy
- CKM phases (β, γ)
- Gluonic penguins
- EW penguins
- Charm physics
- Semileptonics: Mixing, A_{SL}
- Semileptonics: V_{xb}
- $B \to \tau\nu$, $D\tau\mu$,
- $B \to K^{*}\nu\nu$
- τ-physics

Some only LHCb, some only Belle II

On-going

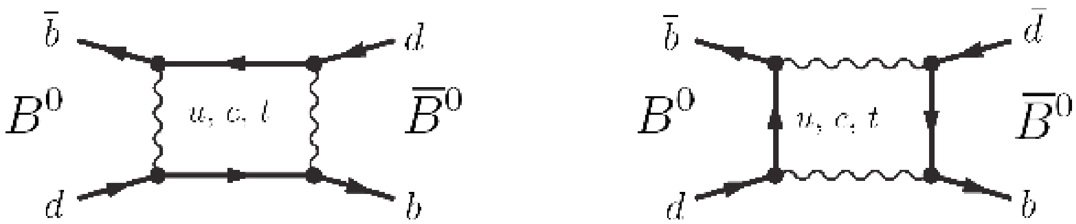
Belle II

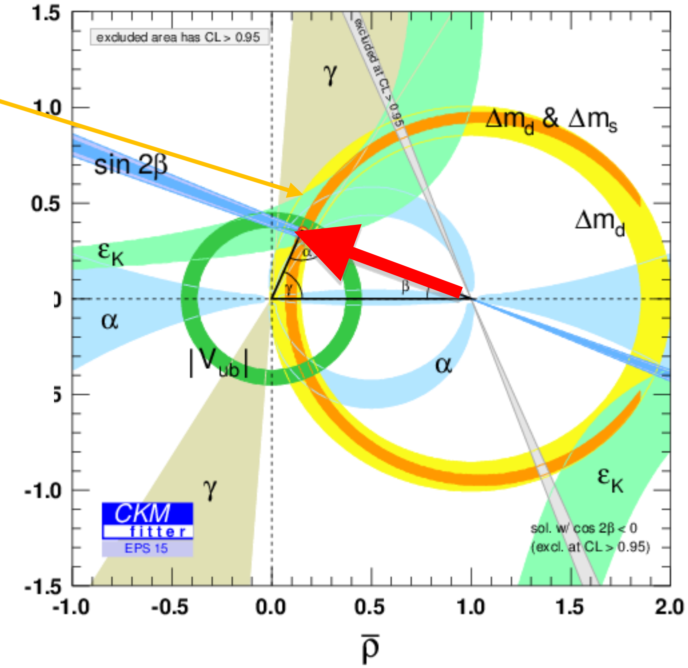
Let's hope!

B mixing and CP violation

Measuring the other side of the CKM triangle

- Through the mass difference Δm_d, Δm_s in the box diagrams:



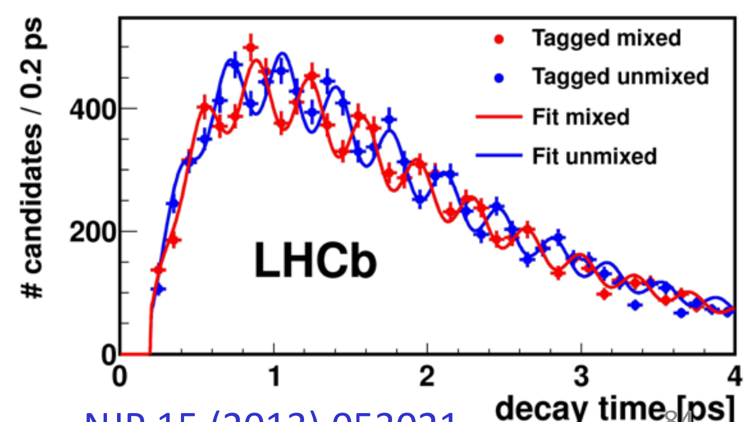


The time-dependent asymmetry A(t) of $B_{d,s}$ events that changed flavour (mixed) or not (unmixed) is:

$$\boxed{A_{mix}(t)\frac{N^{unmix}(t)-N^{mix}(t)}{N^{unmix}(t)+N^{mix}(t)} = \cos\Delta mt}$$

with experimental effects:

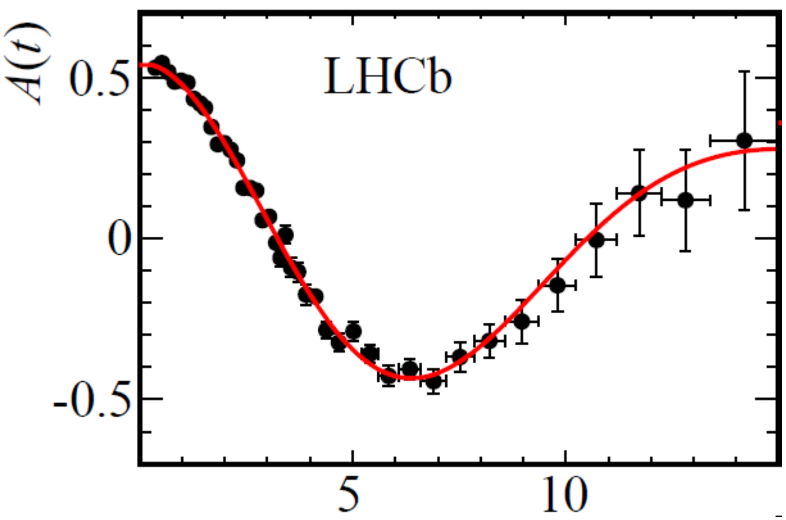
$$=\{(1-2\omega)\times\cos\Delta m\Delta t\}\otimes R(\Delta t)$$



NJP 15 (2013) 053021

B mixing and CP violation

Measuring the other side of the CKM triangle

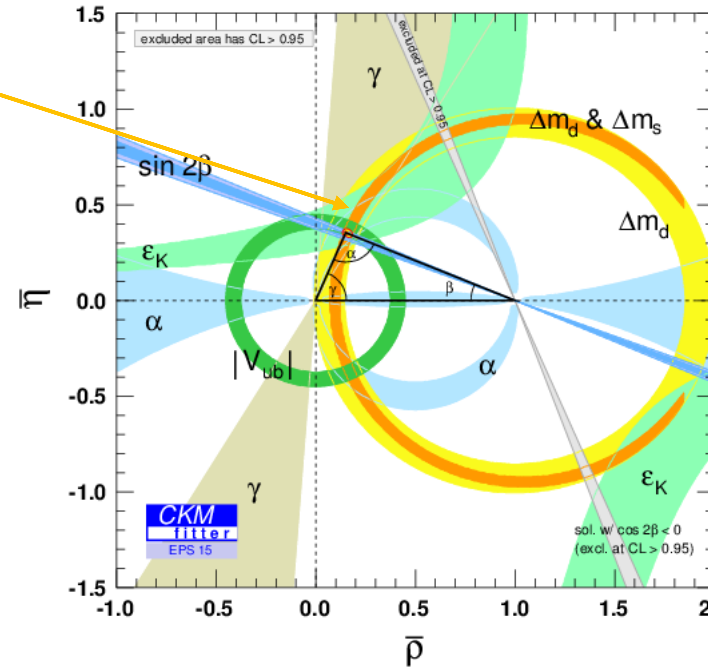


(best tagging category)

[EPJ C (2016) 76]

$\Delta m_d = (0.5050 \pm 0.0021 \pm 0.0010)\ \mathrm{ps}^{-1}$

$\Delta m_s = (17.768 \pm 0.023 \pm 0.006)\ \mathrm{ps}^{-1}$



$$\frac{\Delta m_d}{\Delta m_s} = \frac{m_{B_d} f_{B_d}^2 \hat{B}_{B_d}}{m_{B_s} f_{B_s}^2 \hat{B}_{B_s}} \frac{|V_{td}|^2}{|V_{ts}|^2}$$

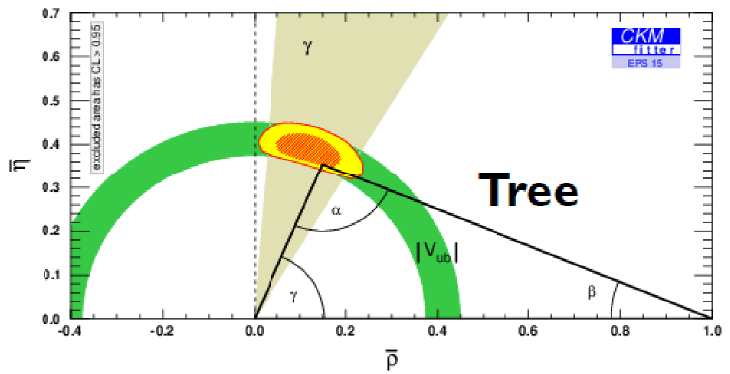
error from theory

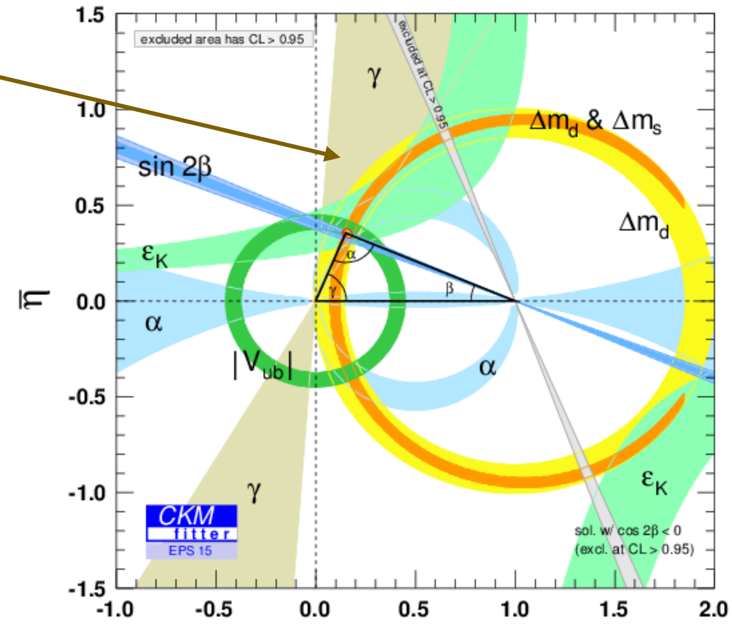
$$|V_{td}/V_{ts}| = 0.216 \pm 0.001 \pm 0.011$$

B mixing and CP violation

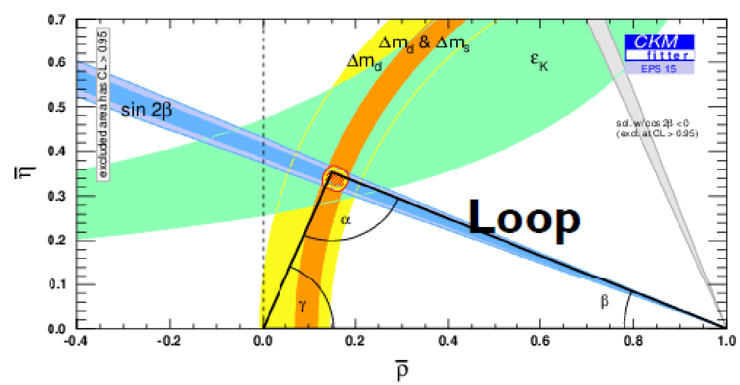
Measuring the γ angle of the CKM triangle

Key to determine the apex from tree level processes, with V_{ub} is the "reference clock"





To be compared with constraints from loop processes:



$\gamma = 72.2^{+5.3}_{-5.8}$