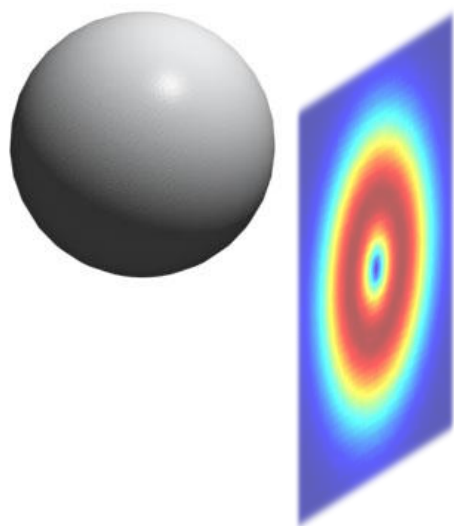
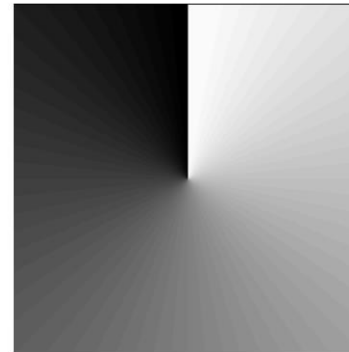
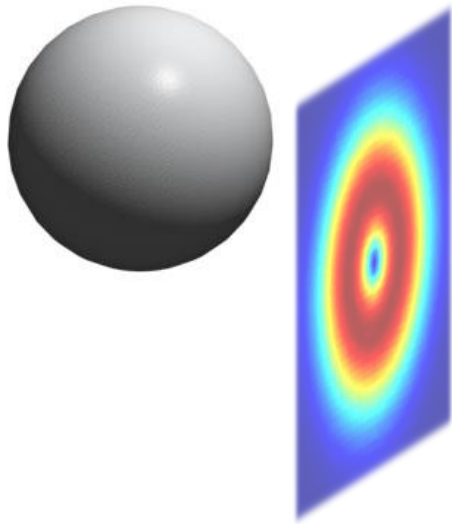
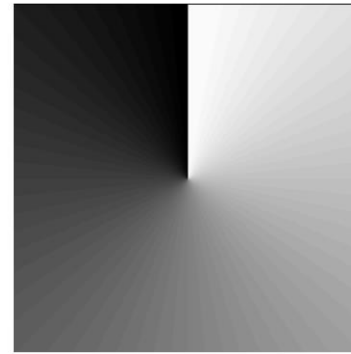
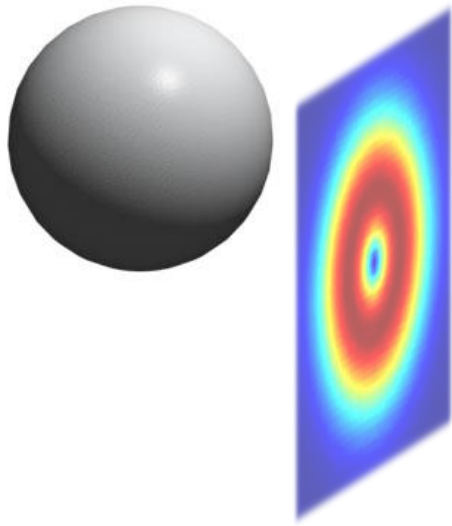




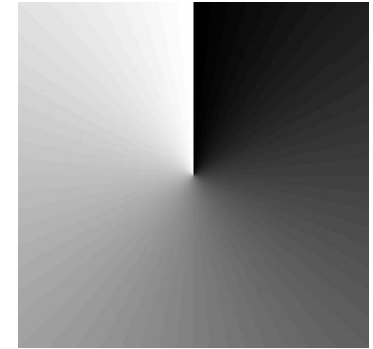


$$q = 1$$

- ✓ Q₁: Does the scattering change if the topological charge $\mathbf{q}$ changes sign?

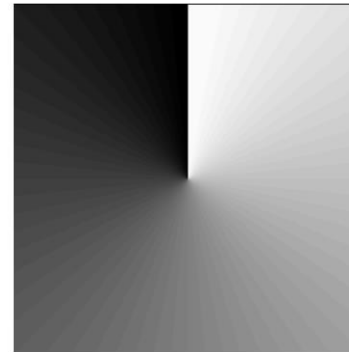
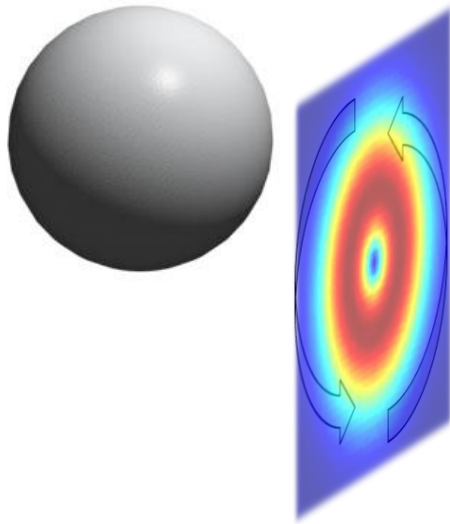


$q = 1$



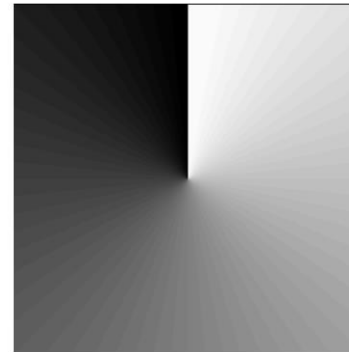
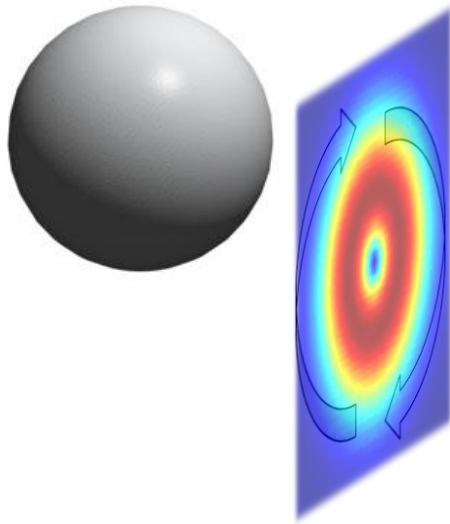
$q = -1$

- ✓ Q_1 : Does the scattering change if the topological charge q changes sign?



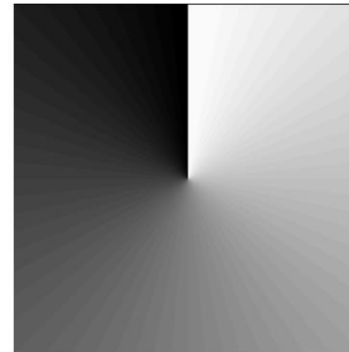
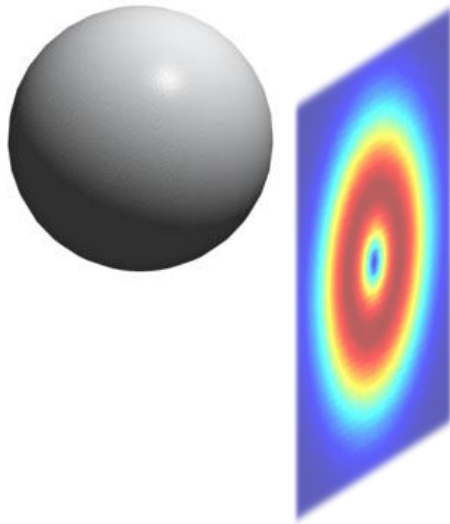
$$q = 1$$

- ✓ Q_1 : Does the scattering change if the topological charge q changes sign?
- ✓ Q_2 : Given a topological charge q , does the sense of circular polarization matter?



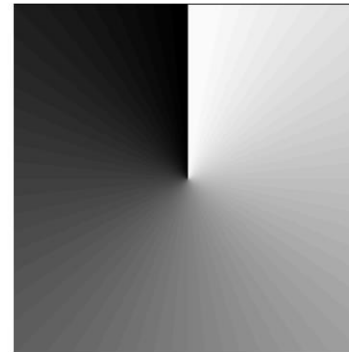
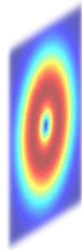
$$q = 1$$

- ✓ Q_1 : Does the scattering change if the topological charge q changes sign?
- ✓ Q_2 : Given a topological charge q , does the sense of circular polarization matter?



$$q = 1$$

- ✓ Q_1 : Does the scattering change if the topological charge q changes sign?
- ✓ Q_2 : Given a topological charge q , does the sense of circular polarization matter?



$$q = 1$$

- ✓ Q_1 : Does the scattering change if the topological charge q changes sign?
- ✓ Q_2 : Given a topological charge q , does the sense of circular polarization matter?
- ✓ Q_3 : Does the focusing regime change the answers to Q_1 and Q_2 ?
Symmetries of the problem don't change



Tailoring Mie scattering with Helicity and Angular Momentum

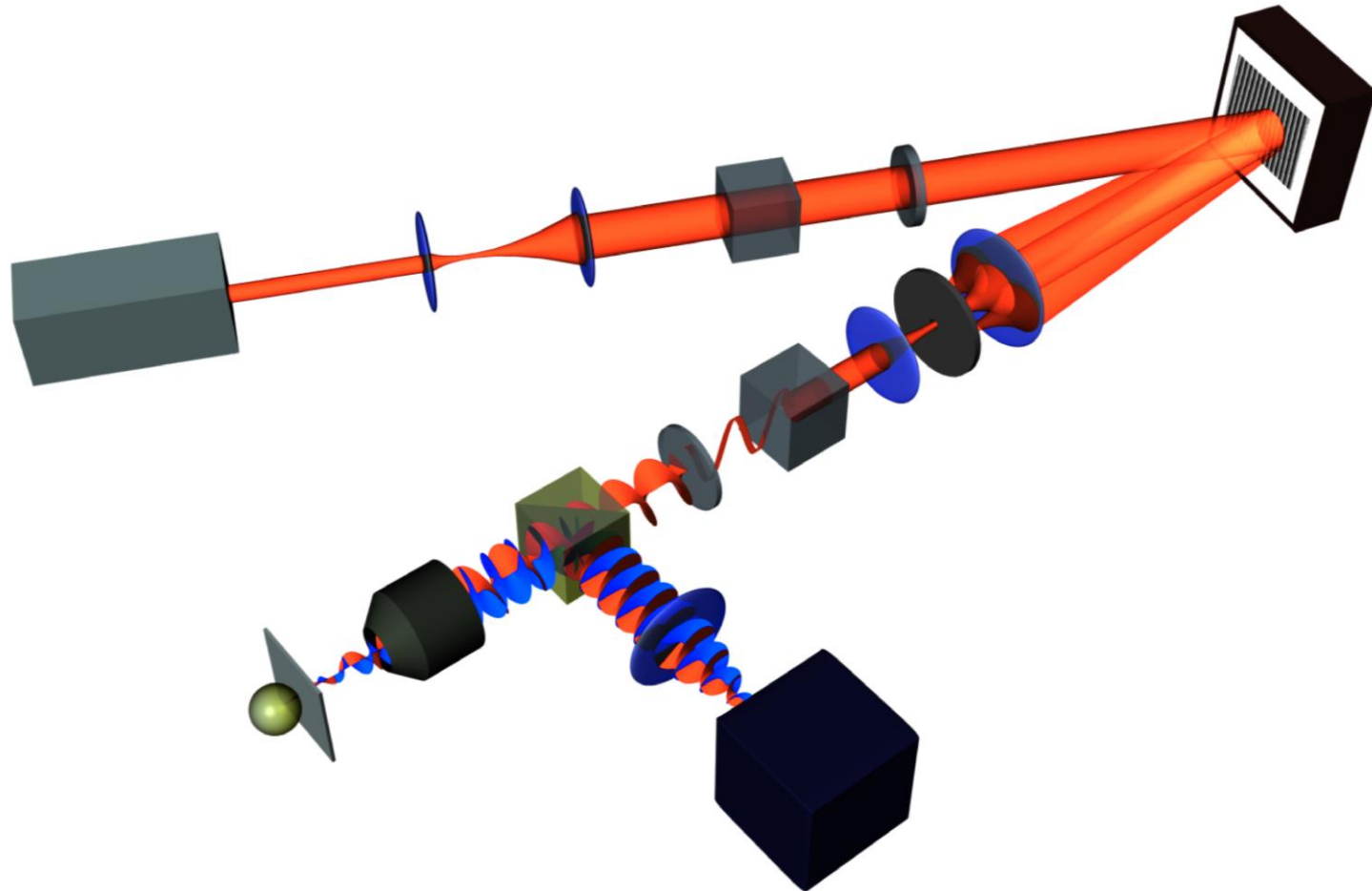
Xavier Zambrana-Puyalto



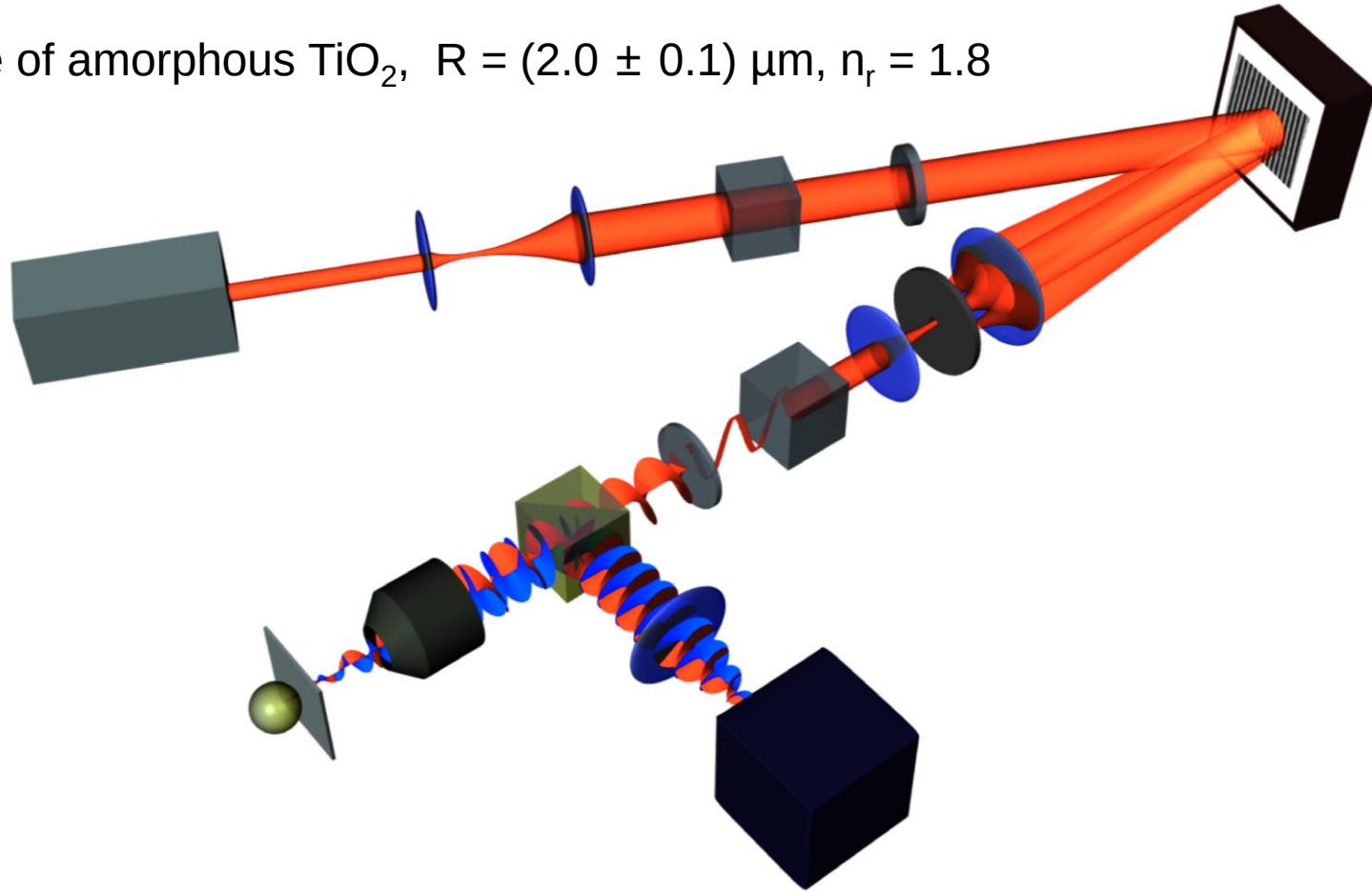
This project has received funding from the European Union's Horizon 2020 research and innovation programme under the Marie Skłodowska-Curie grant agreement No 795838



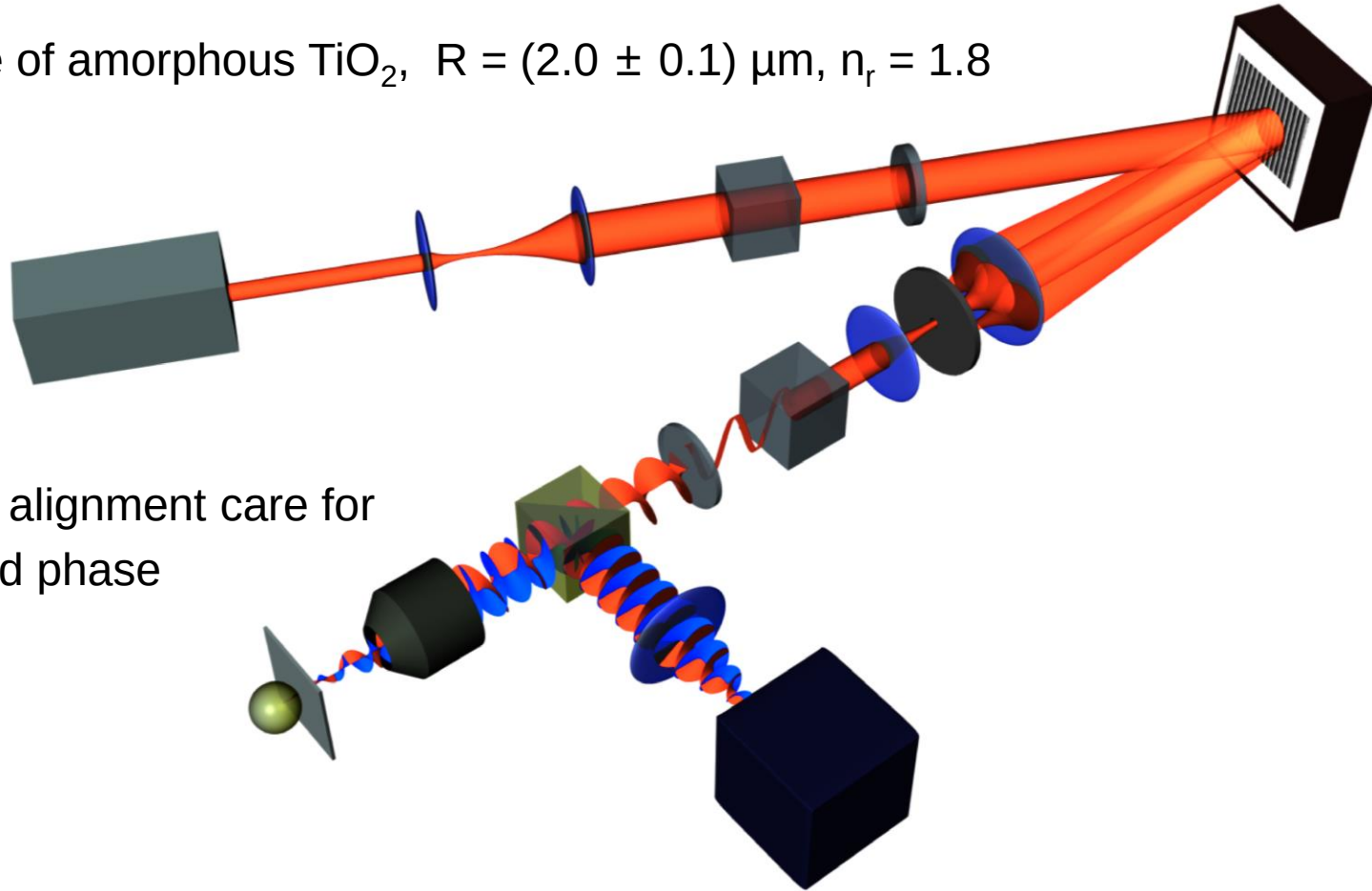
ISTITUTO
ITALIANO DI
TECNOLOGIA



- ✓ Particles made of amorphous TiO_2 , $R = (2.0 \pm 0.1) \mu\text{m}$, $n_r = 1.8$



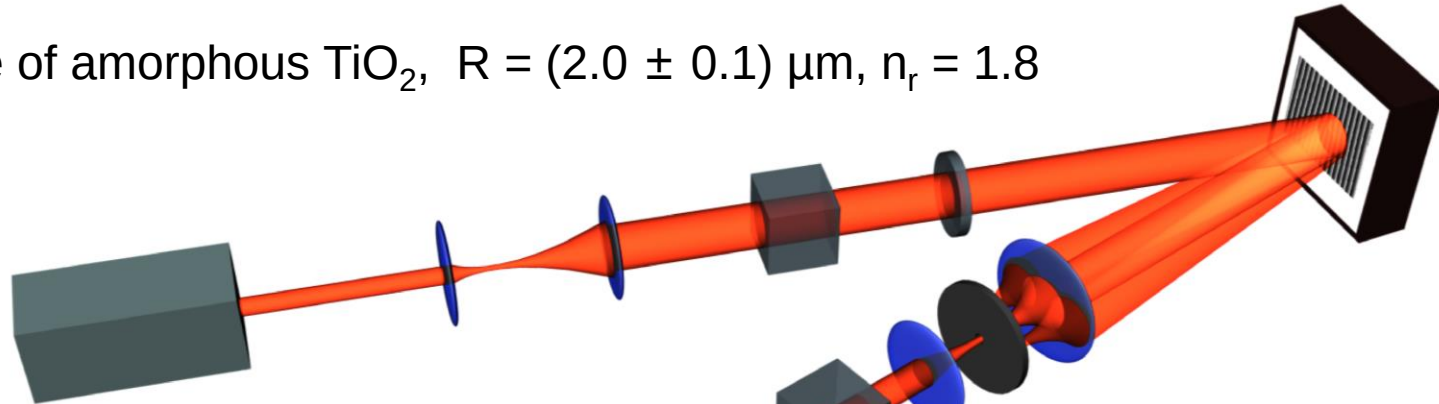
- ✓ Particles made of amorphous TiO_2 , $R = (2.0 \pm 0.1) \mu\text{m}$, $n_r = 1.8$



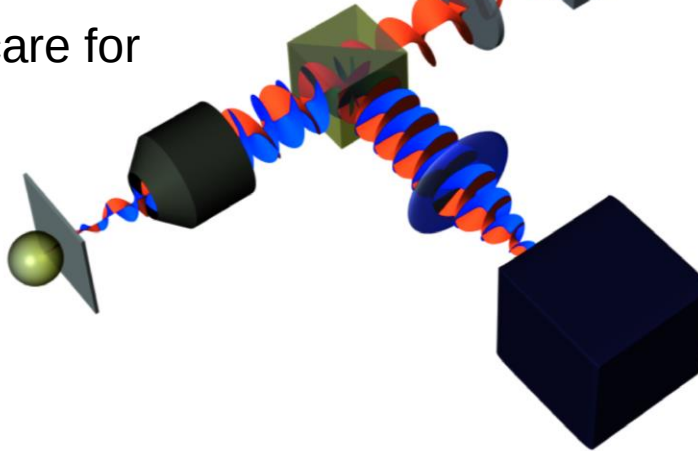
- ✓ λ scans. Great alignment care for polarization and phase

Experimental details

- ✓ Particles made of amorphous TiO_2 , $R = (2.0 \pm 0.1) \mu\text{m}$, $n_r = 1.8$



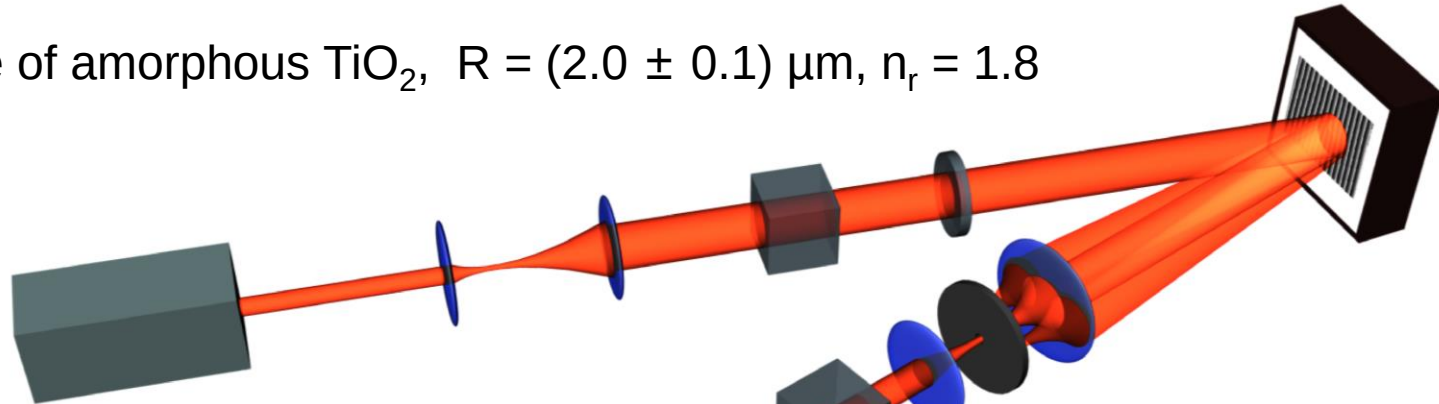
- ✓ λ scans. Great alignment care for polarization and phase



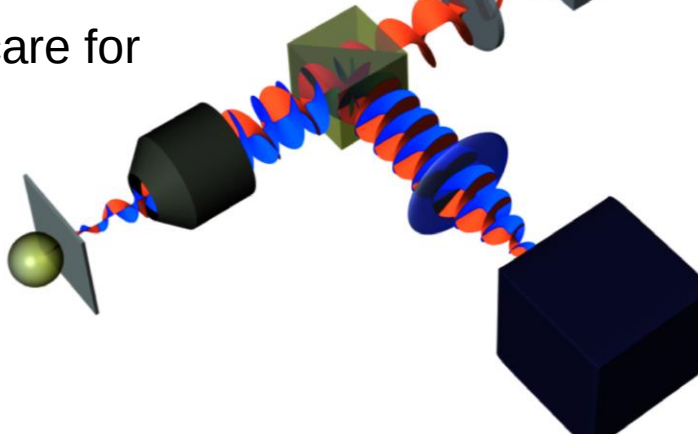
- ✓ Blue/red: two helicities

Experimental details

- ✓ Particles made of amorphous TiO_2 , $R = (2.0 \pm 0.1) \mu\text{m}$, $n_r = 1.8$



- ✓ λ scans. Great alignment care for polarization and phase

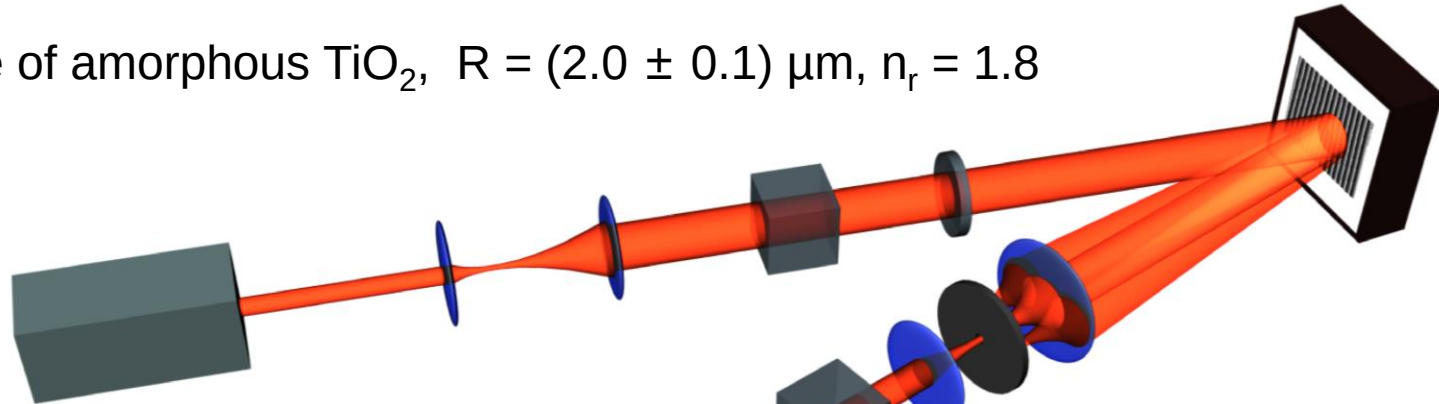


- ✓ Blue/red: two helicities
- ✓ Backscattering dimensionless figure of merit

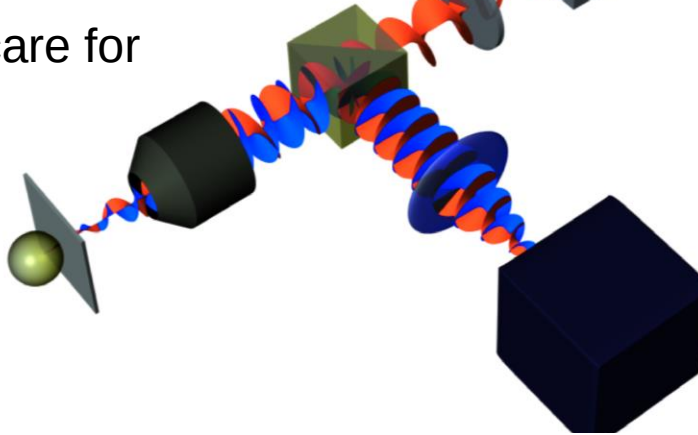
$$I = \frac{P_p - P_0}{P_s - P_0}$$

Experimental details

- ✓ Particles made of amorphous TiO_2 , $R = (2.0 \pm 0.1) \mu\text{m}$, $n_r = 1.8$



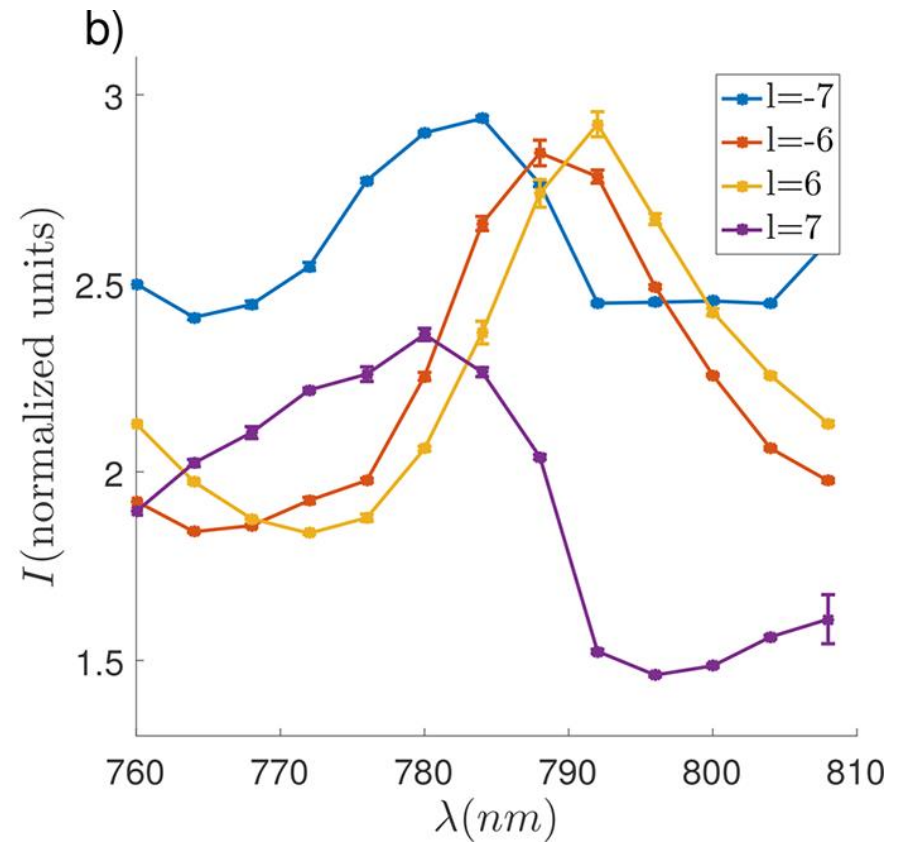
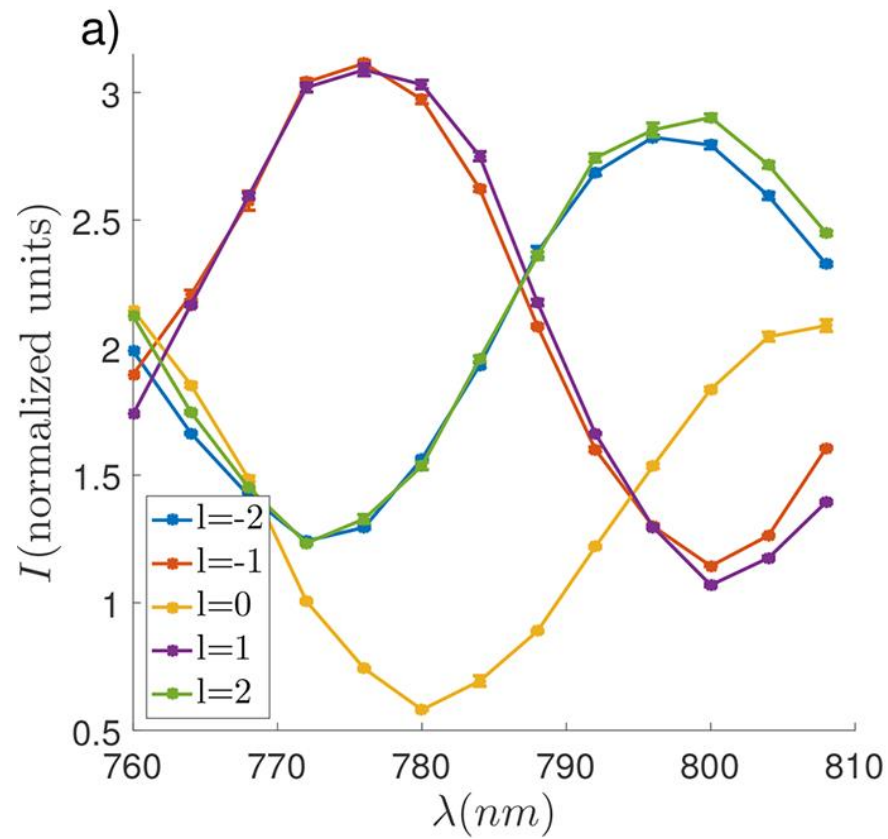
- ✓ λ scans. Great alignment care for polarization and phase



- ✓ Blue/red: two helicities
- ✓ Backscattering dimensionless figure of merit
- ✓ Mirror and cylindrically symmetric

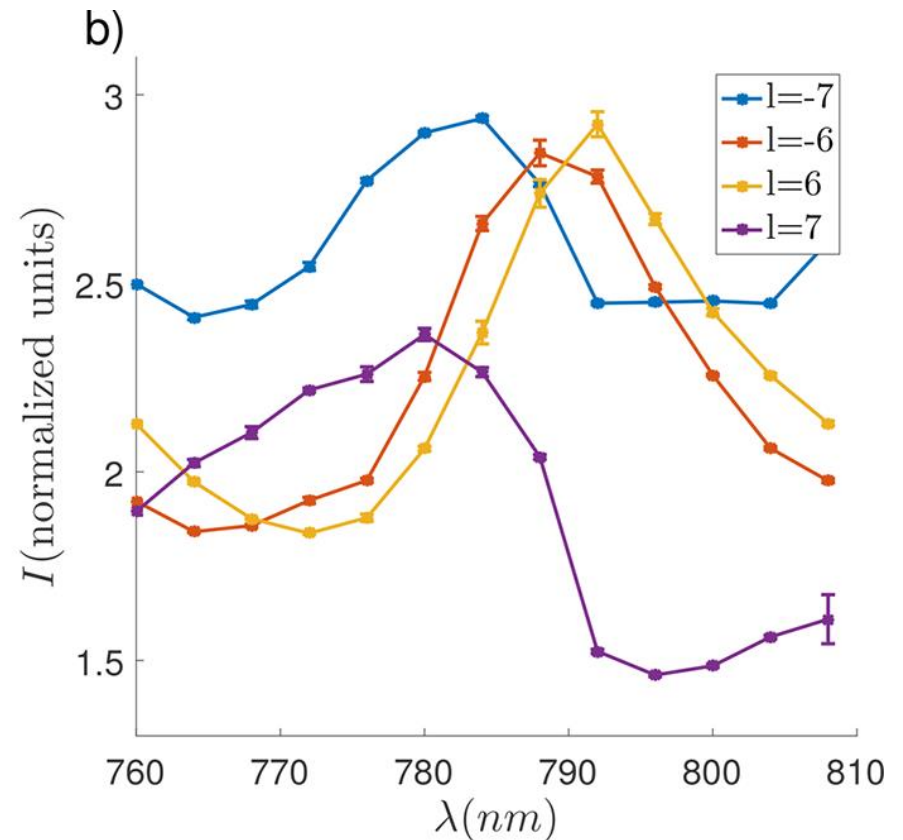
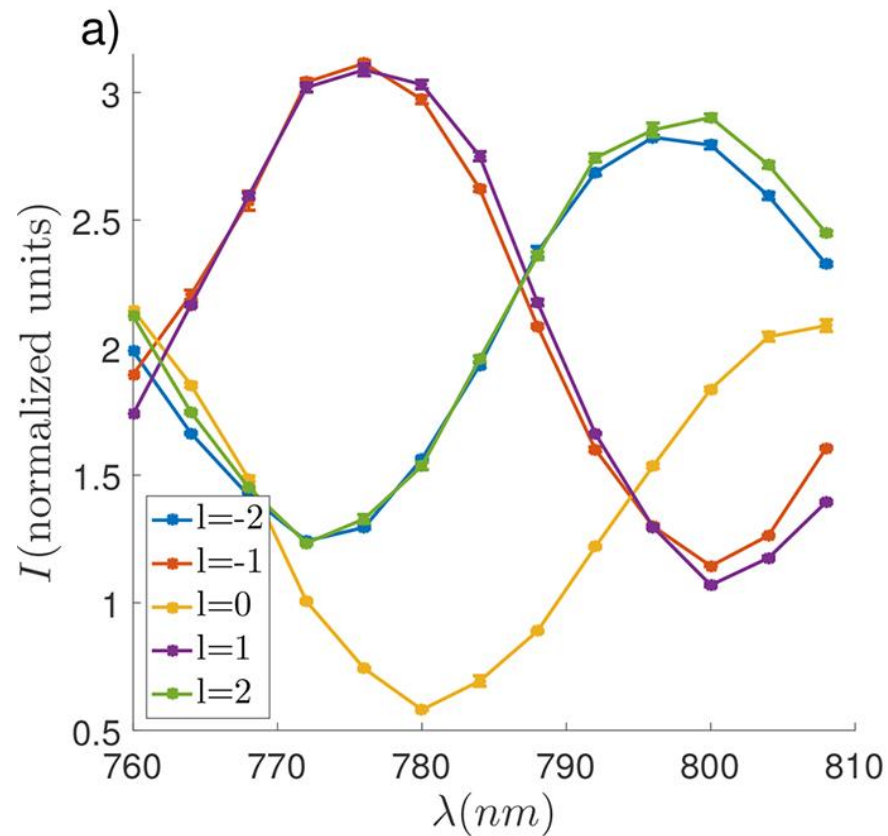
$$I = \frac{P_p - P_0}{P_s - P_0}$$

Experimental results (I)



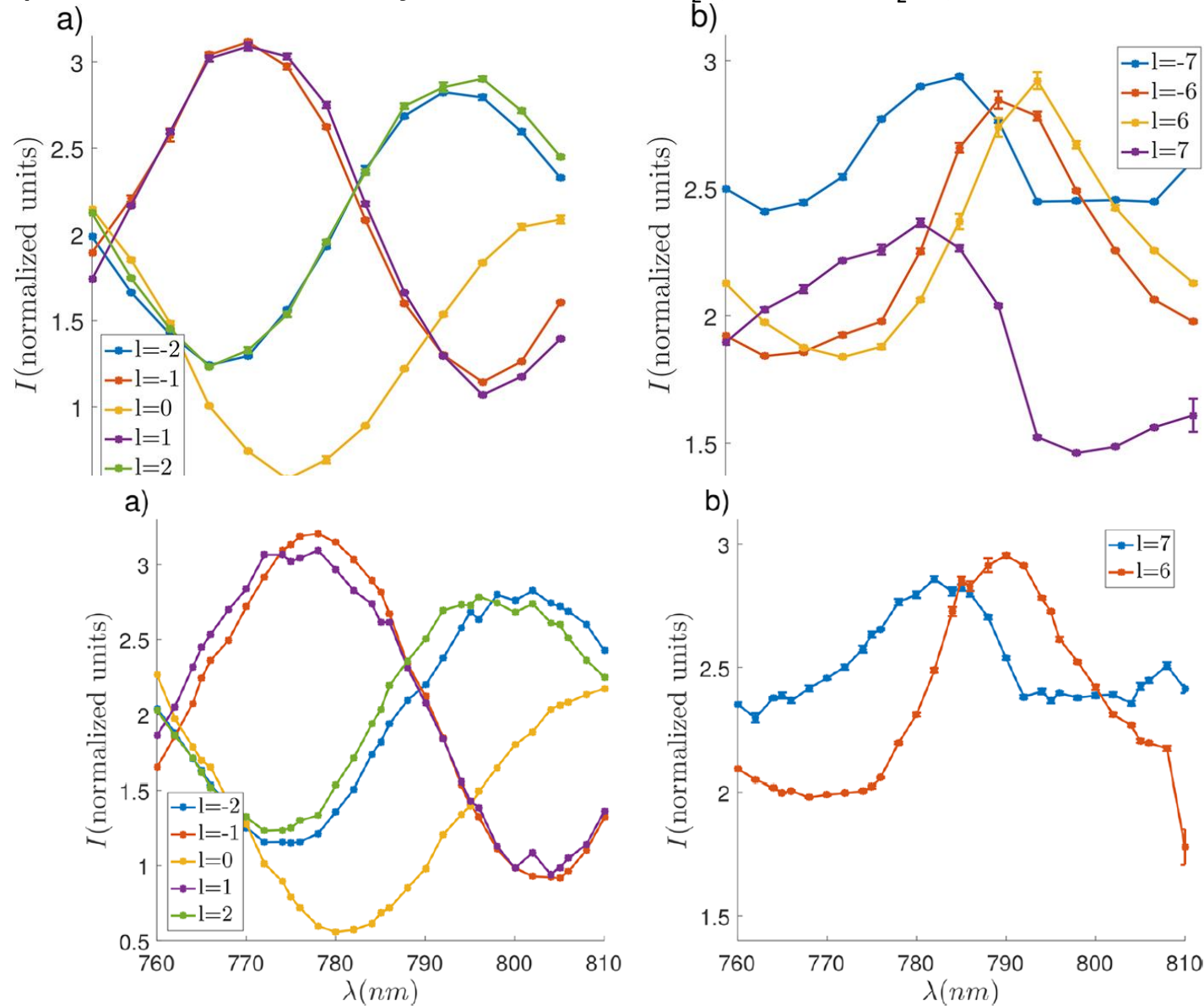
Experimental results (I)

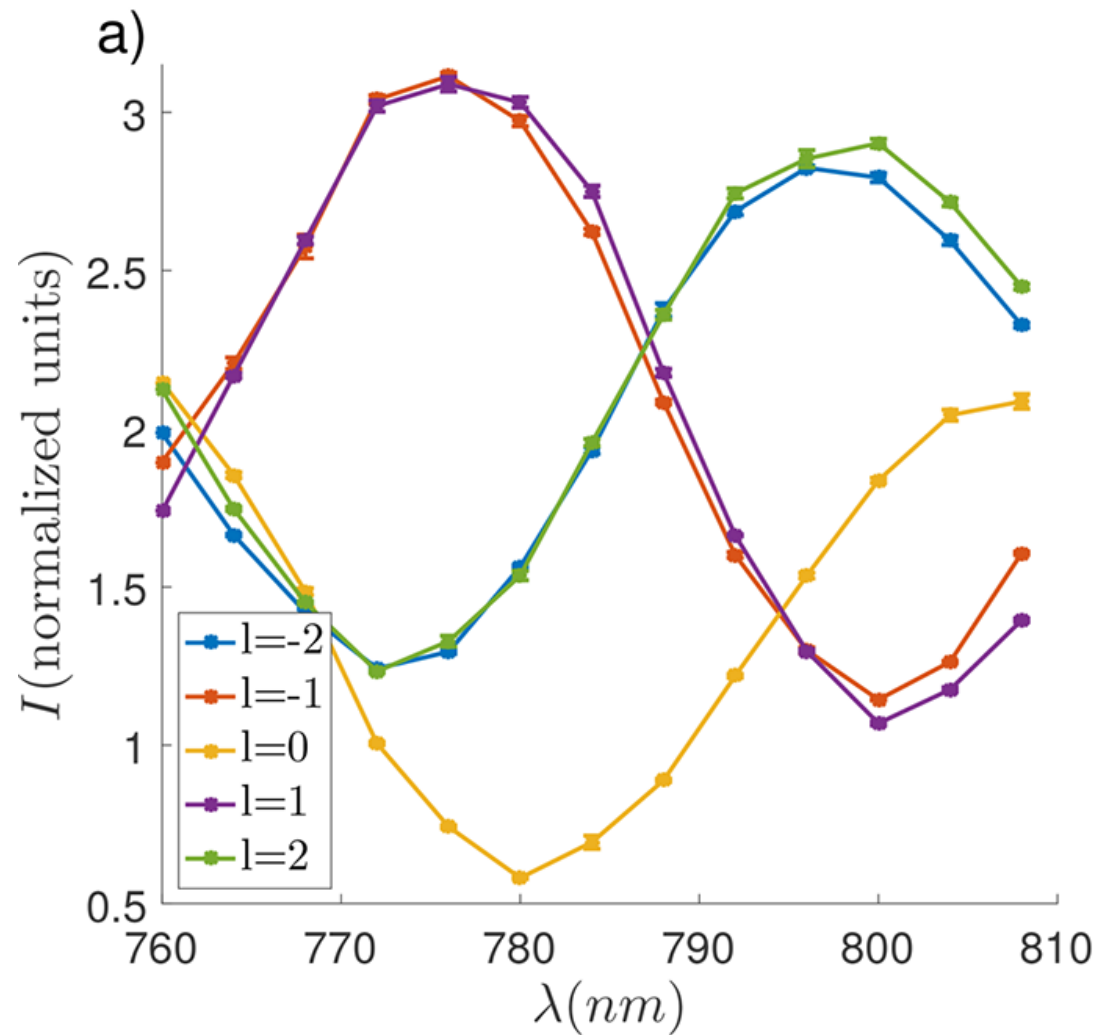
✓ Experiment has some symmetries. $(\lambda, m_z, p) = (\lambda, -m_z, -p)$



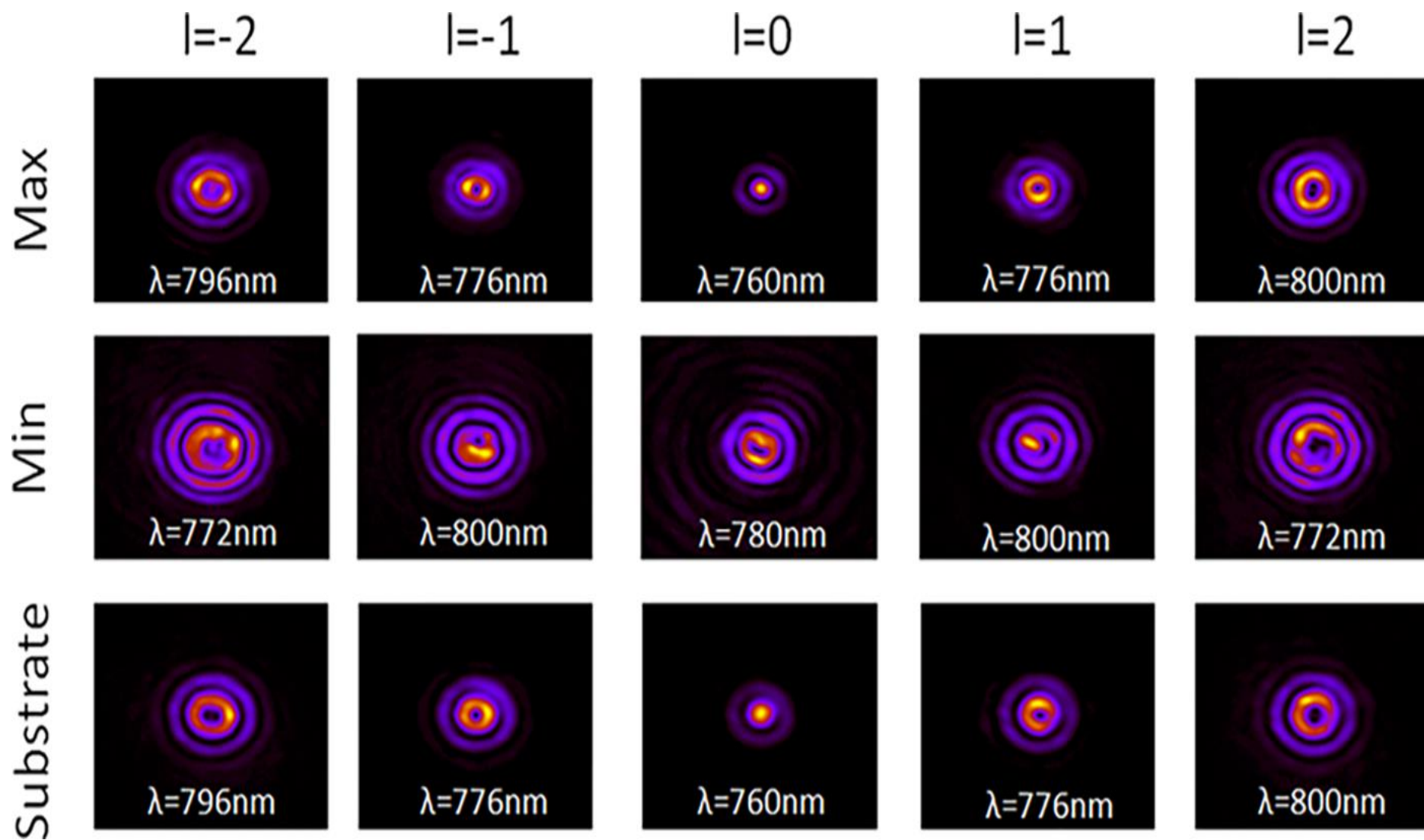
Experimental results (I)

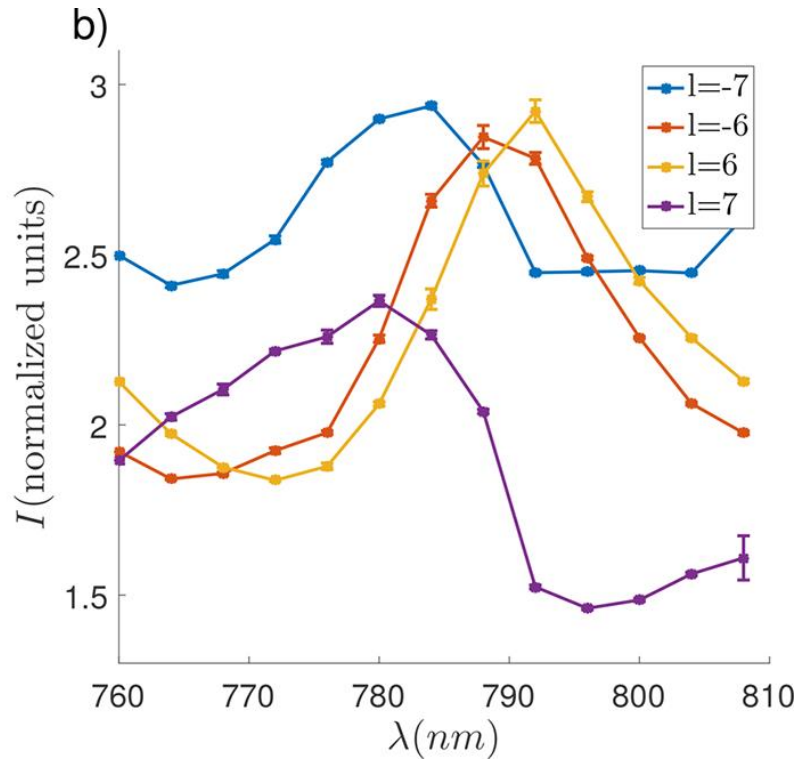
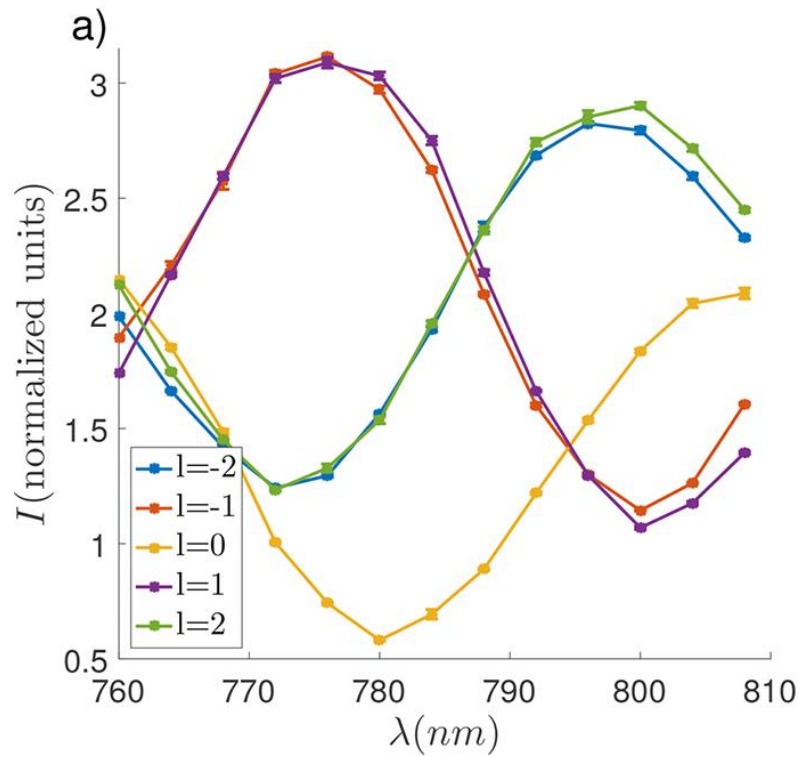
✓ Experiment has some symmetries. $(\lambda, m_z, \rho) = (\lambda, -m_z, -\rho)$

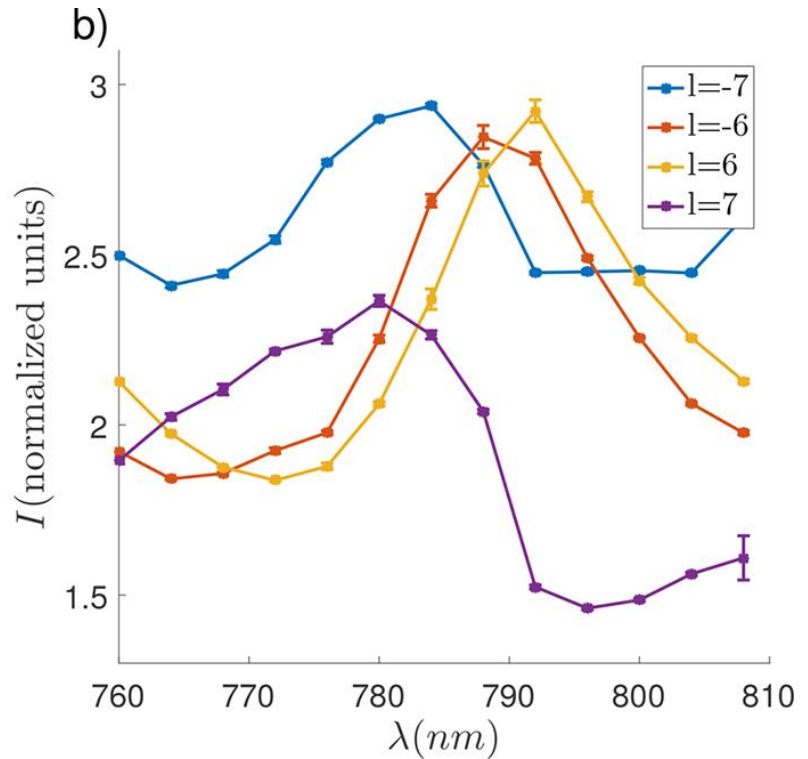
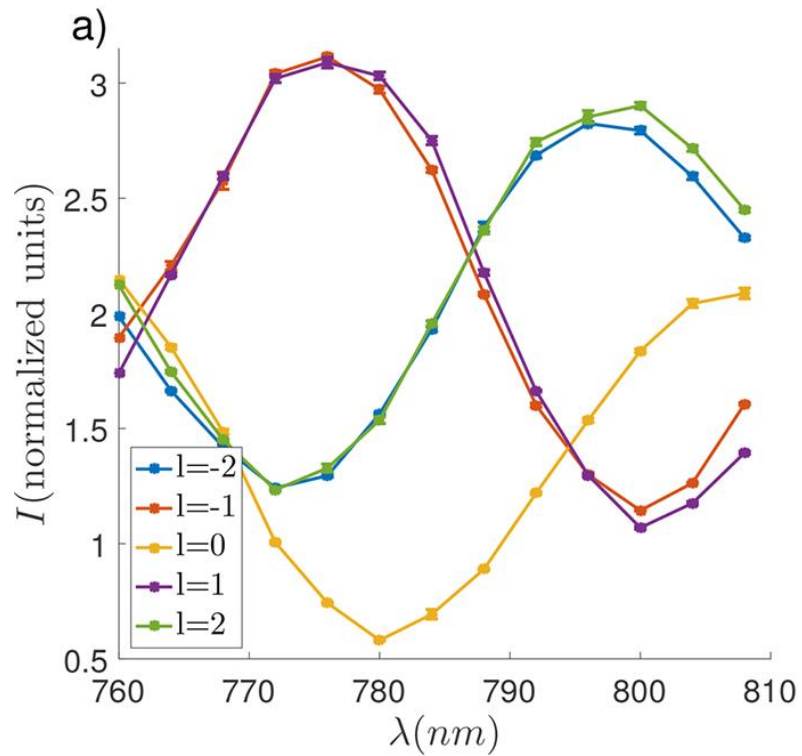




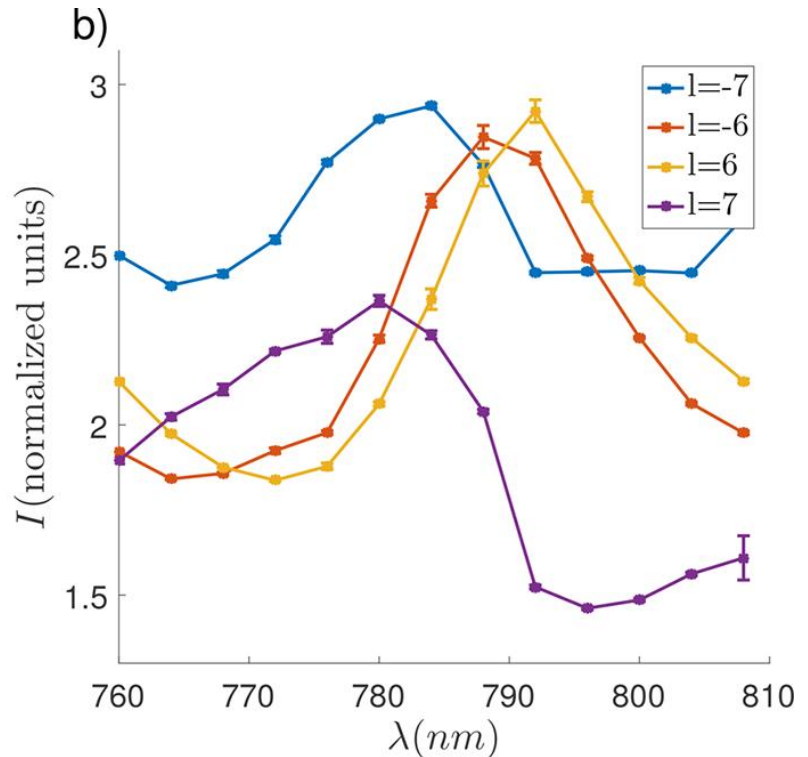
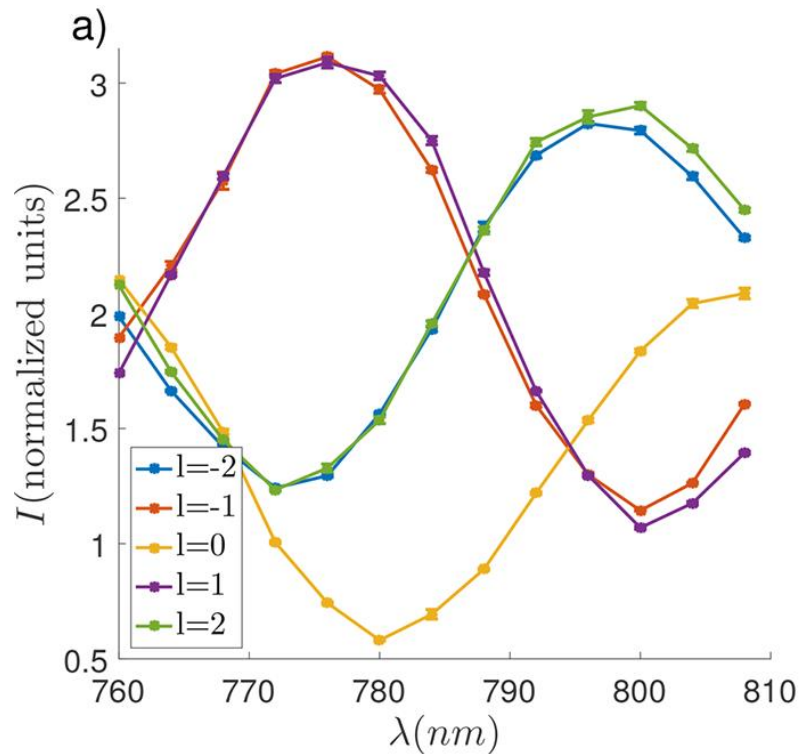
Experimental results (II)



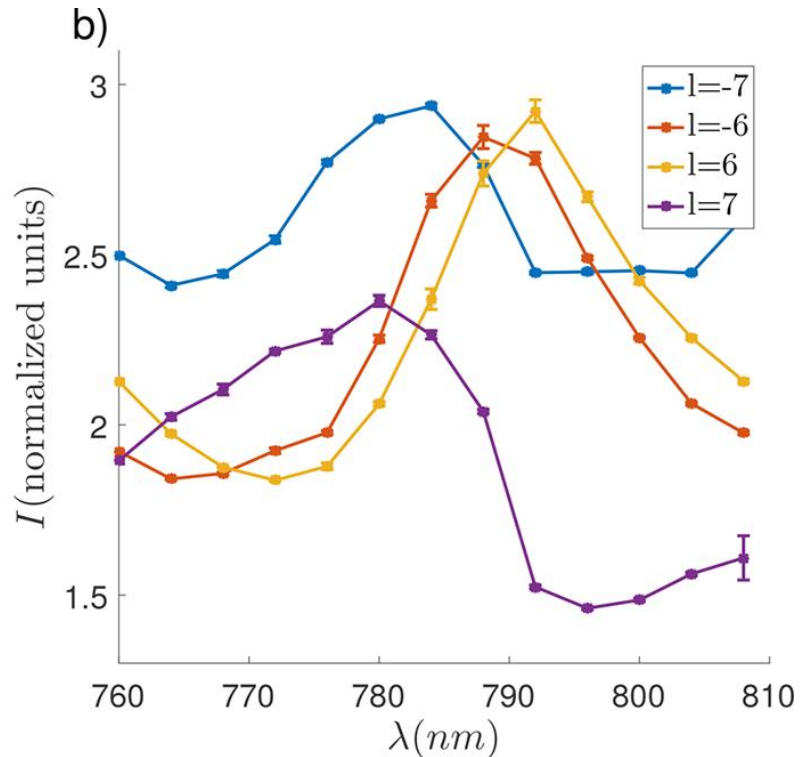
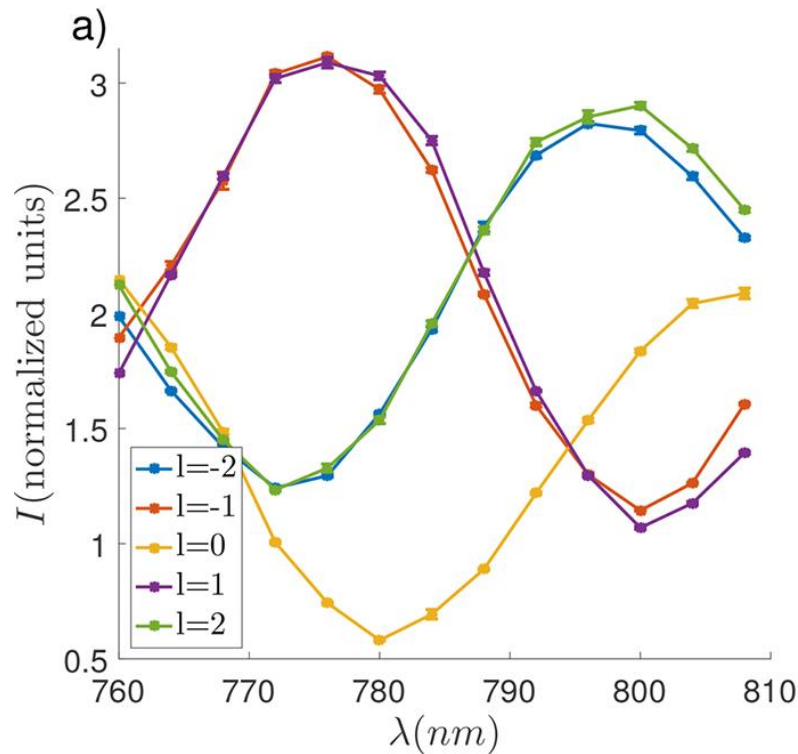




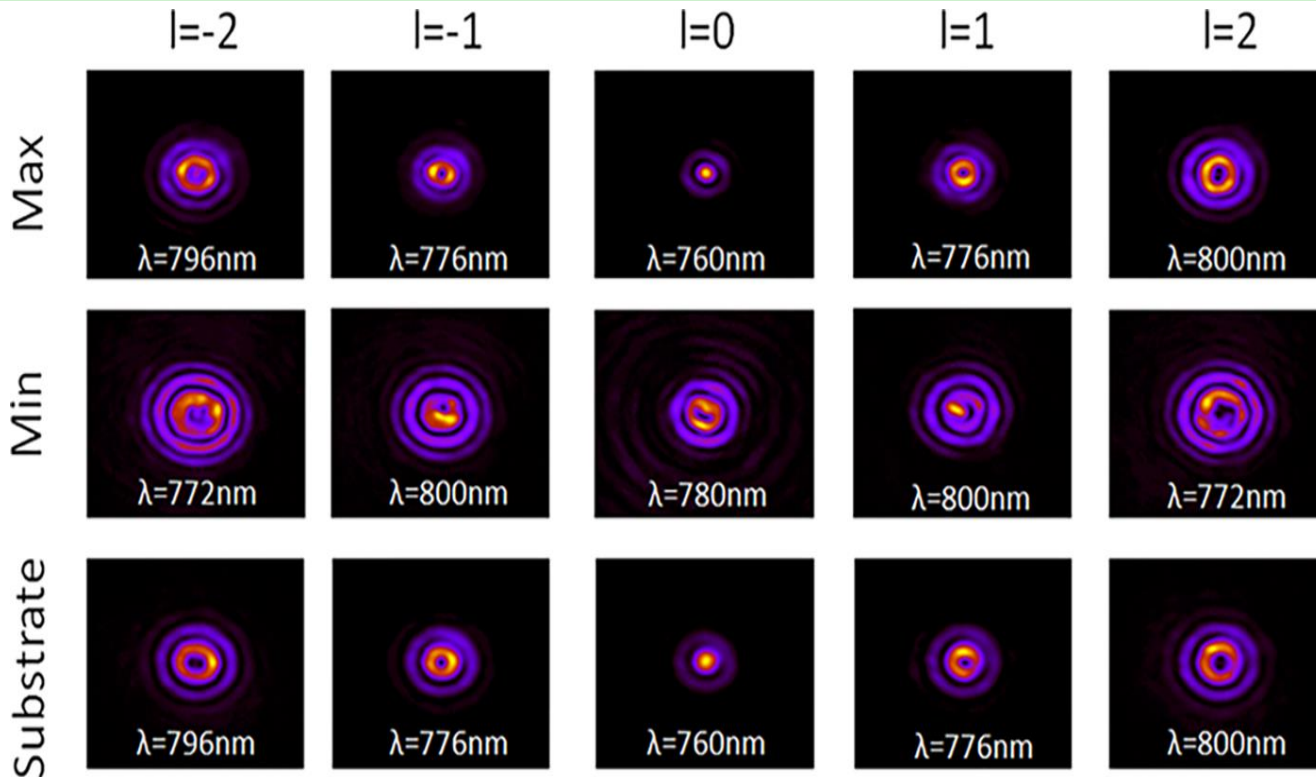
✓ 1) Why do these resonances move around?



- ✓ 1) Why do these resonances move around?
- ✓ 2) Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't?

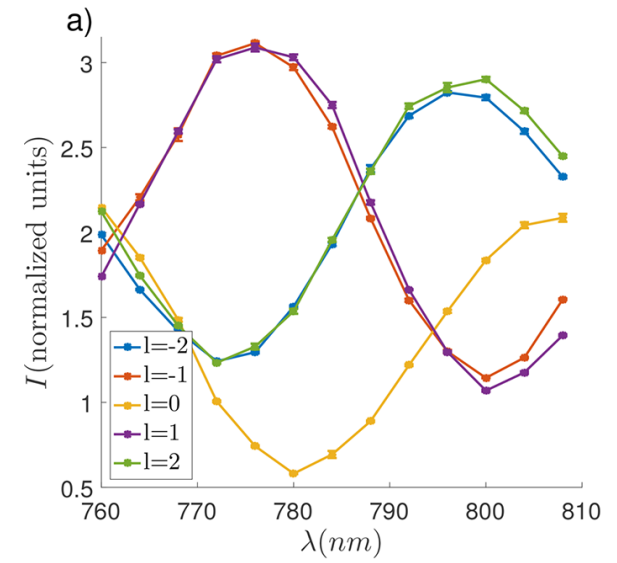


- ✓ 1) Why do these resonances move around?
- ✓ 2) Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't?
- ✓ 3) Why do $l = 0$ and $l = 7$ have different linewidths?



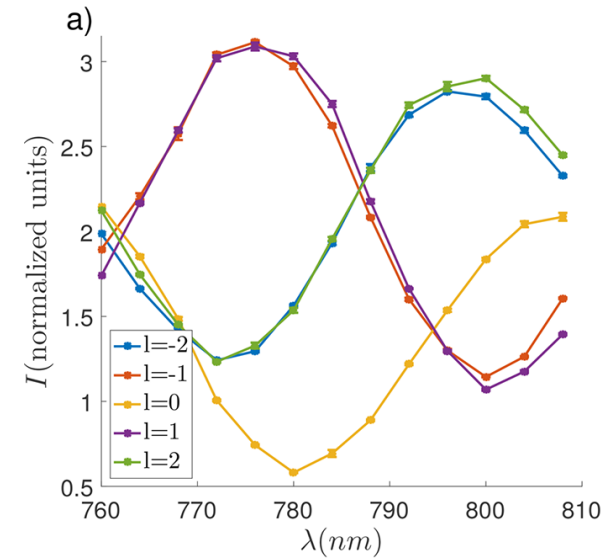
- ✓ 1) Why do these resonances move around?
- ✓ 2) Why do $|l| = \pm 1$ behave equally, while $|l| = \pm 7$ don't?
- ✓ 3) Why do $|l| = 0$ and $|l| = 7$ have different linewidths?
- ✓ 4) Why do we see different intensity patterns?

Why do these resonances move around?



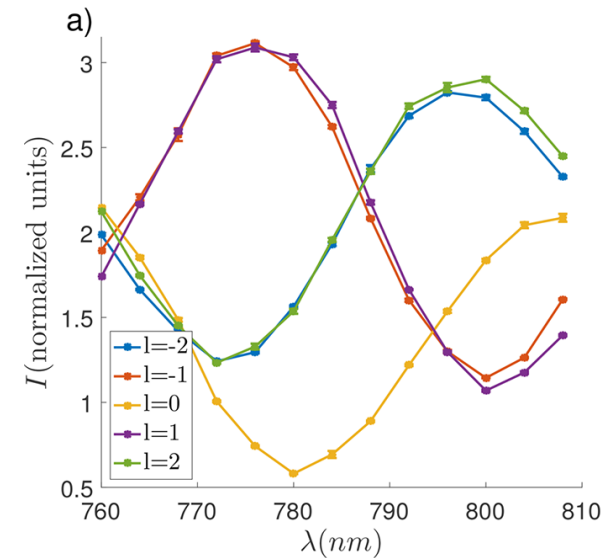
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination



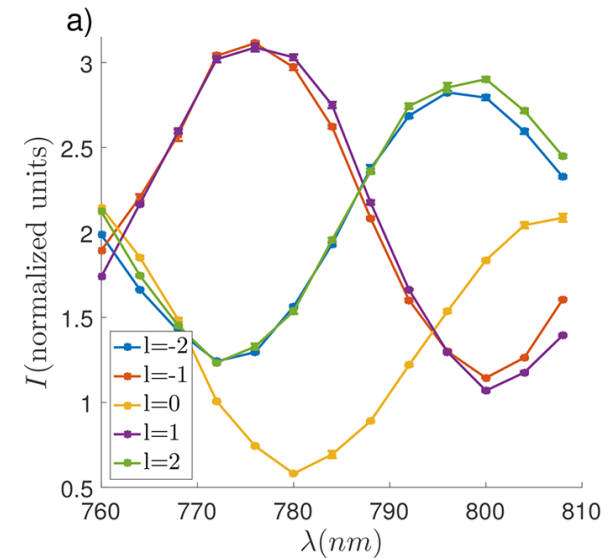
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes



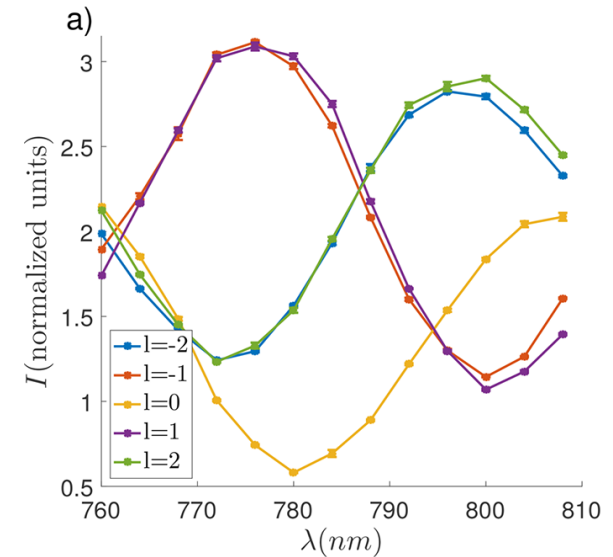
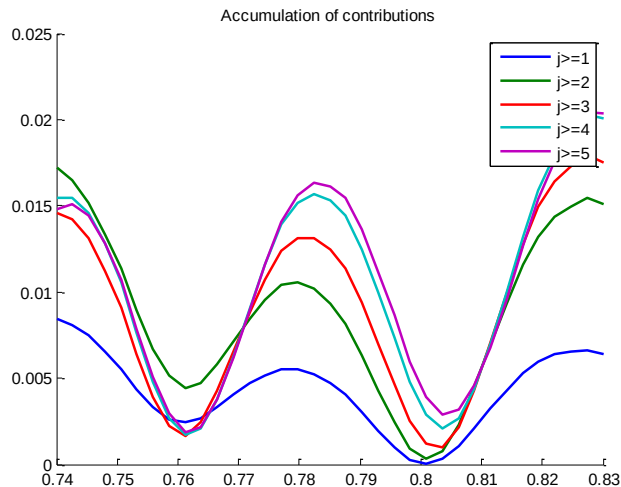
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory



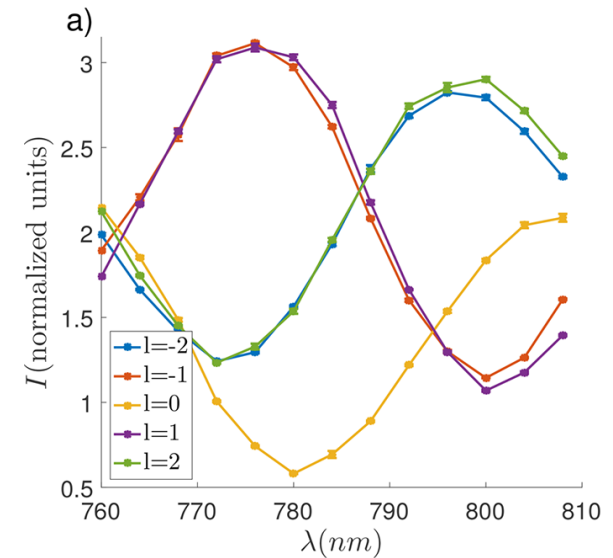
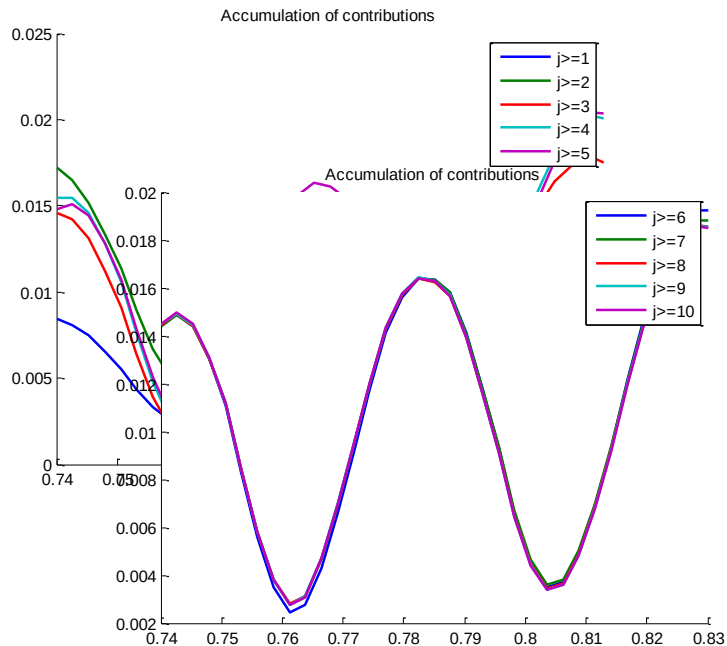
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory



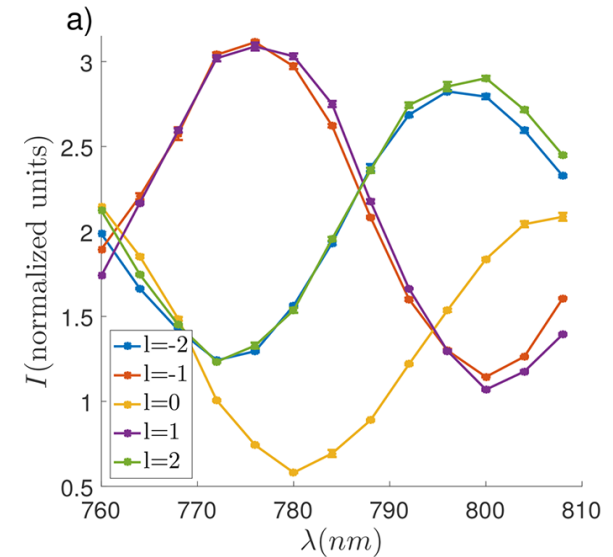
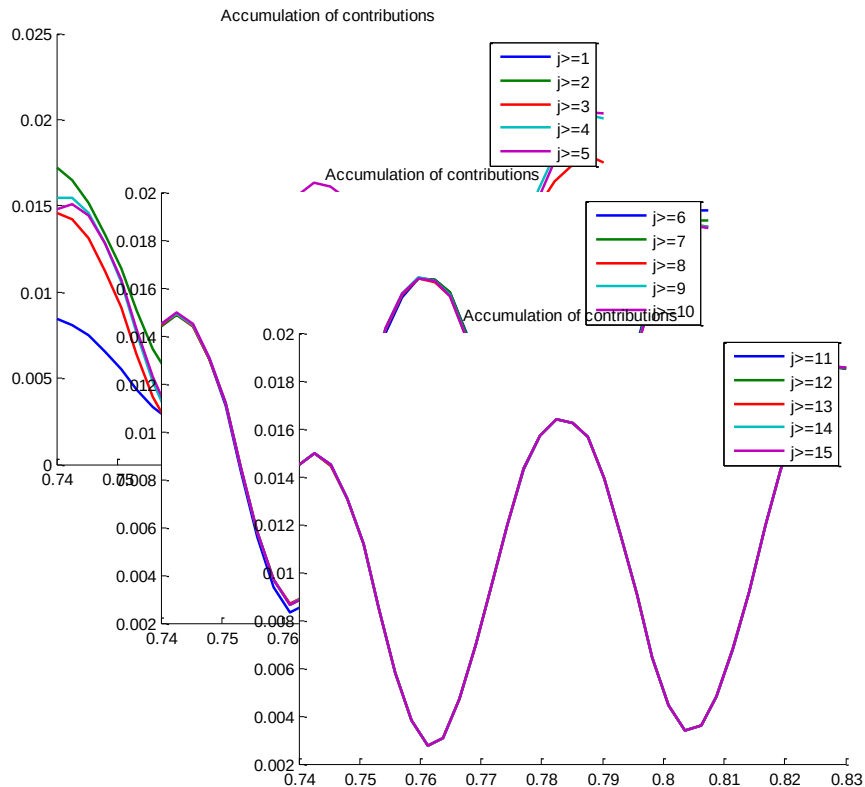
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory



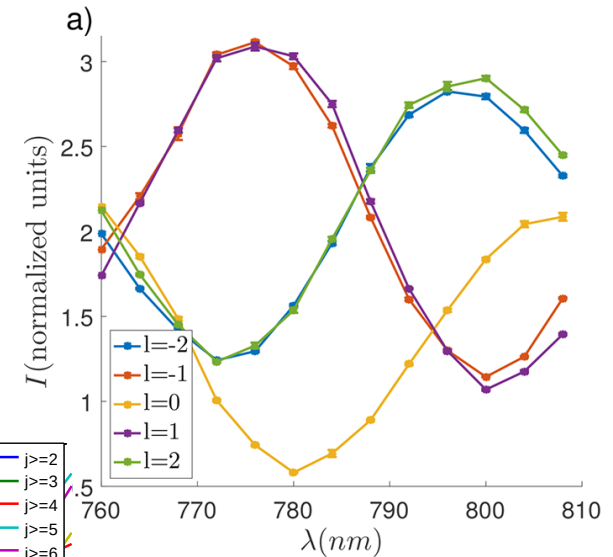
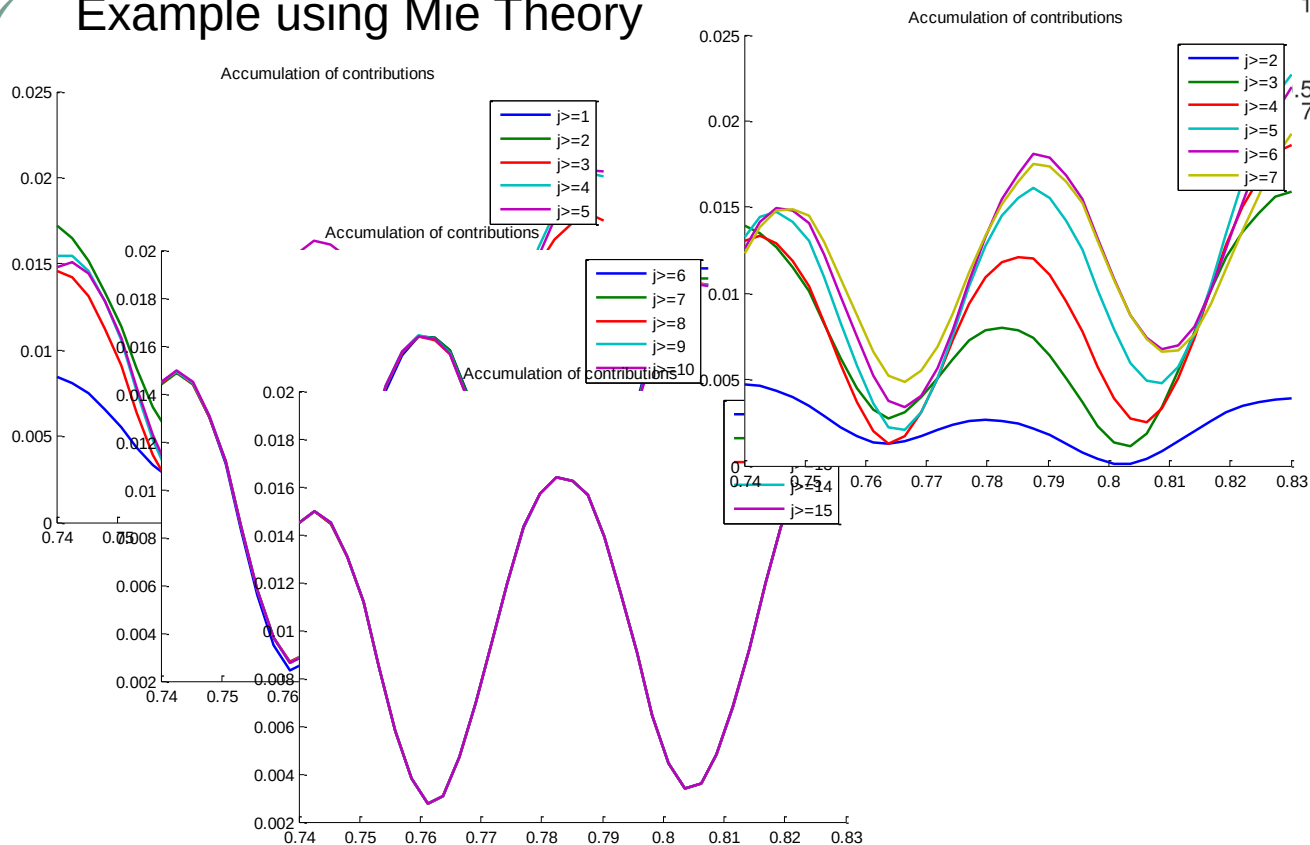
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory



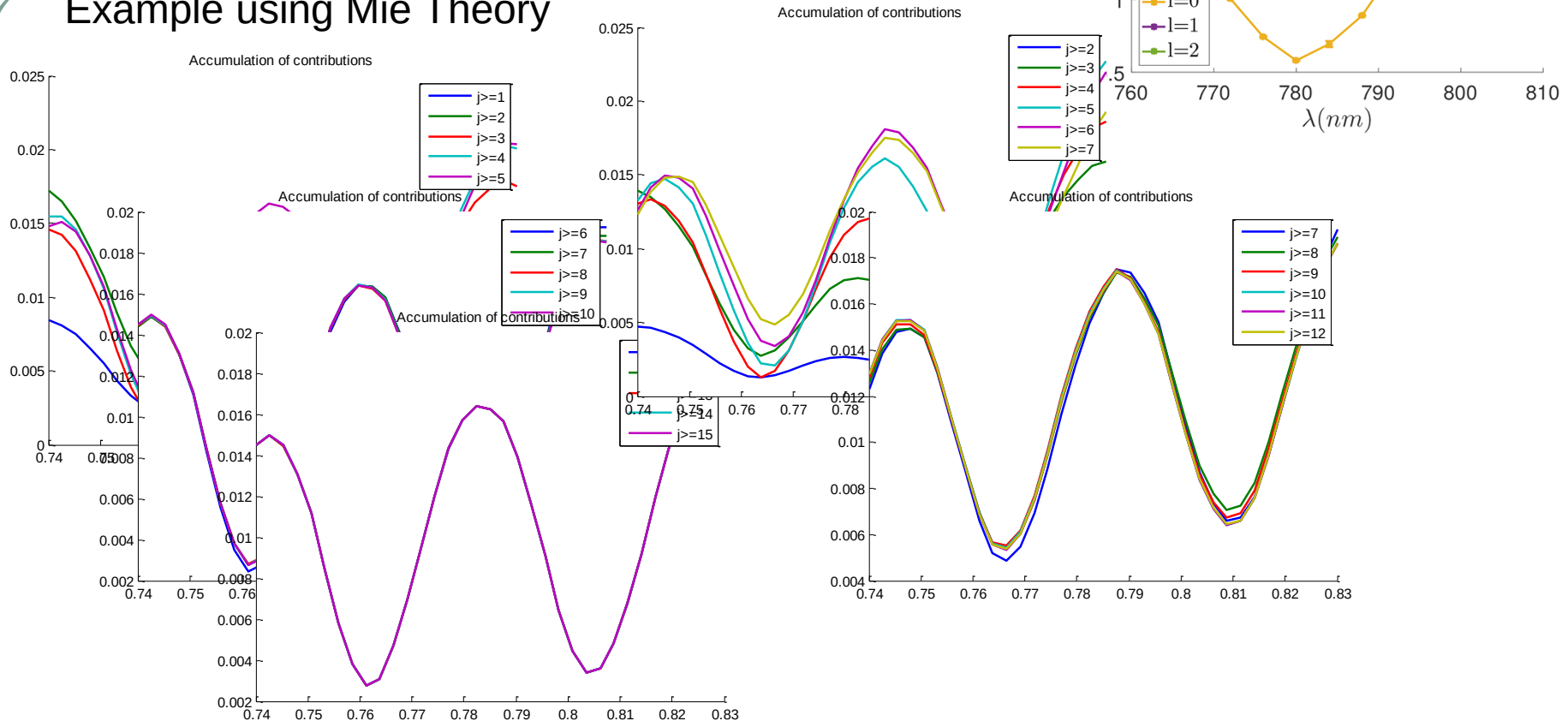
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them ‘power oscillations’. They are due to the superposition of many modes
- ✓ Example using Mie Theory



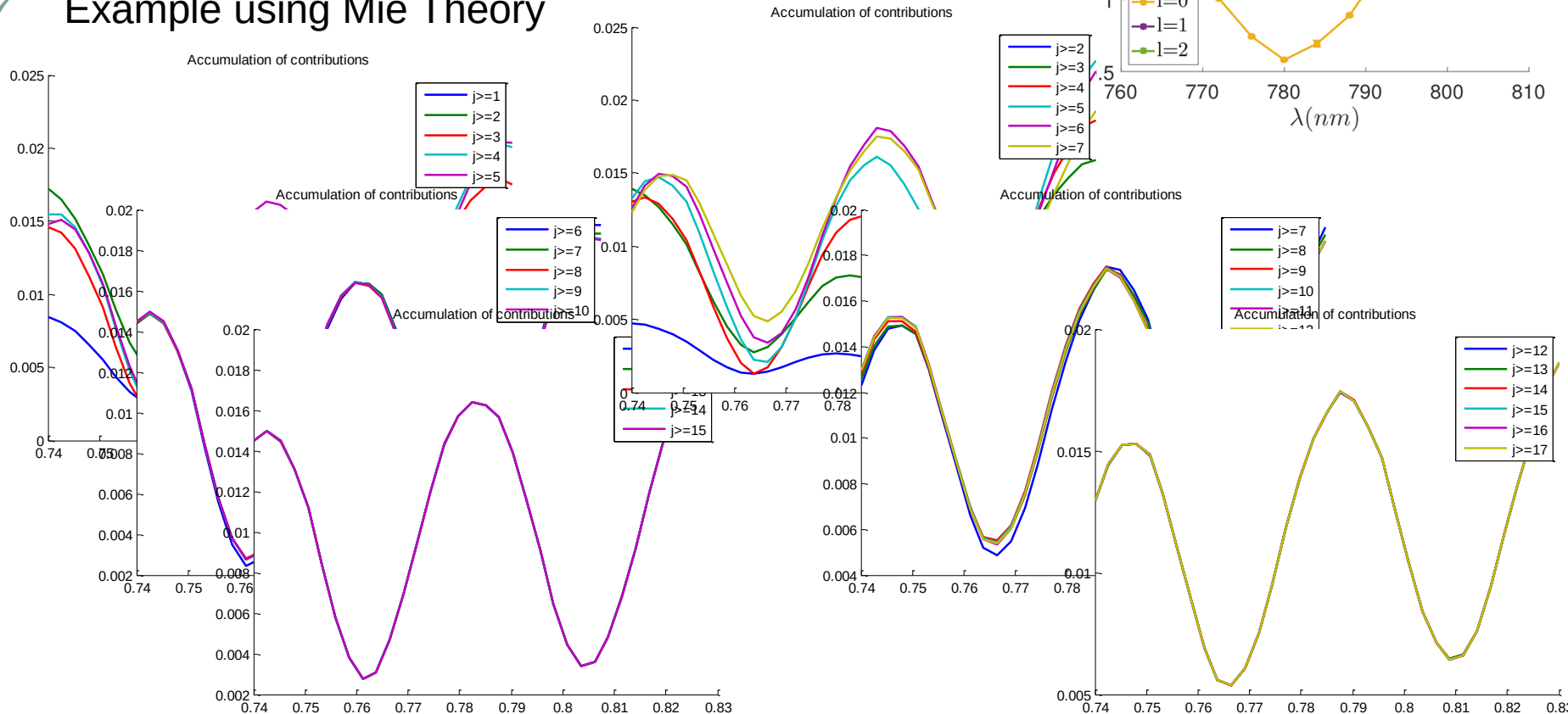
Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory

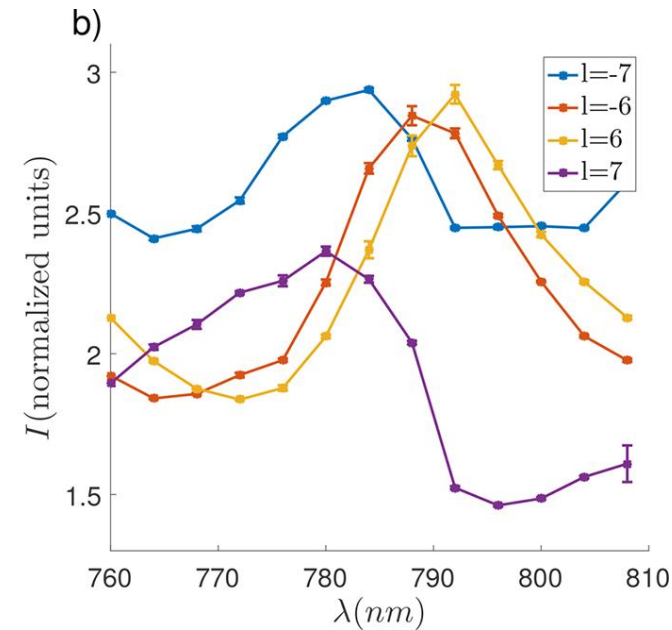
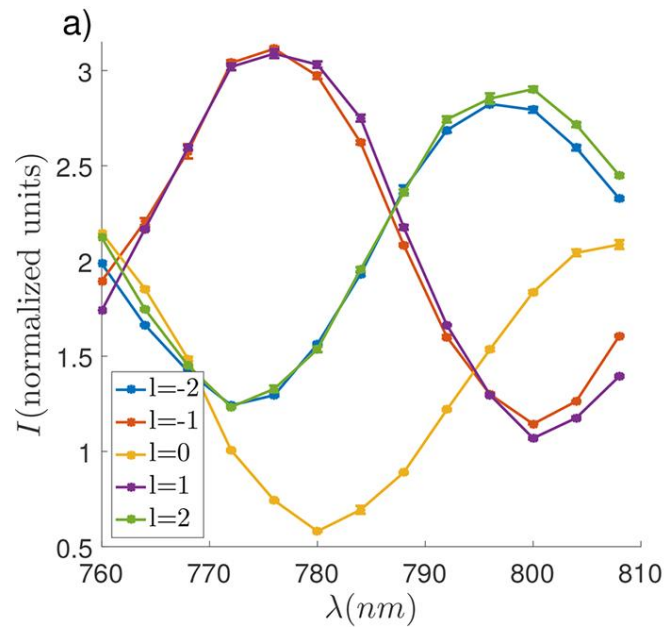


Why do these resonances move around?

- ✓ Resonances are associated to eigenmodes. Their spectral position cannot be tailored with the illumination
- ✓ You can call them 'power oscillations'. They are due to the superposition of many modes
- ✓ Example using Mie Theory



Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (I)



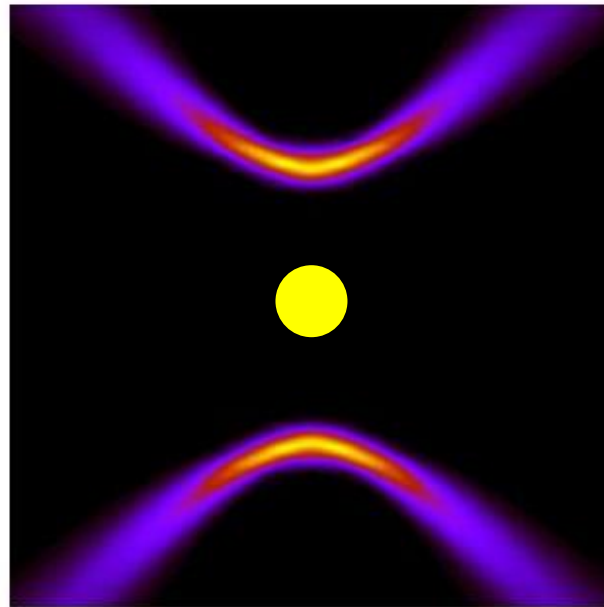
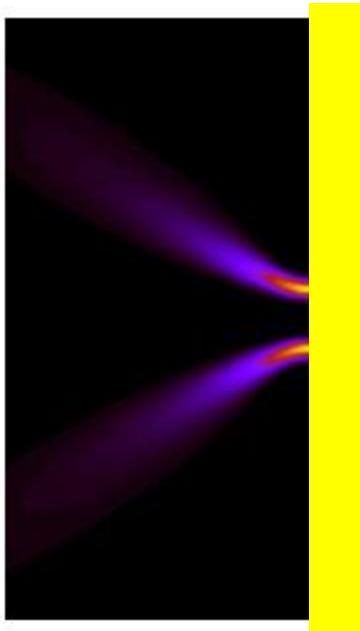


Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (I)

- ✓ Intuitive picture: Interesting phenomena happen when the size of the scatterer and the beam width are similar

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (I)

- ✓ Intuitive picture: Interesting phenomena happen when the size of the scatterer and the beam width are similar
- ✓ 2 extreme cases:
 - i) Beam width is much smaller than the scatterer
 - ii) Beam width is much larger than the scatterer





Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)



Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)



Rigorous picture:

Simulation



Mie Theory

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

✓ Rigorous picture: Simulation | Mie Theory

✓ Mie Theory: Given a particle (R, n_r), then

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

- ✓ Rigorous picture: Simulation | Mie Theory
- ✓ **Mie Theory**: Given a particle (R, n_r), then
 - ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

- ✓ Rigorous picture: Simulation | Mie Theory
- ✓ **Mie Theory**: Given a particle (R, n_r), then
 - ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$
 - ✓ The 'fine details' (narrow resonances) about scattering are found in the higher order modes

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

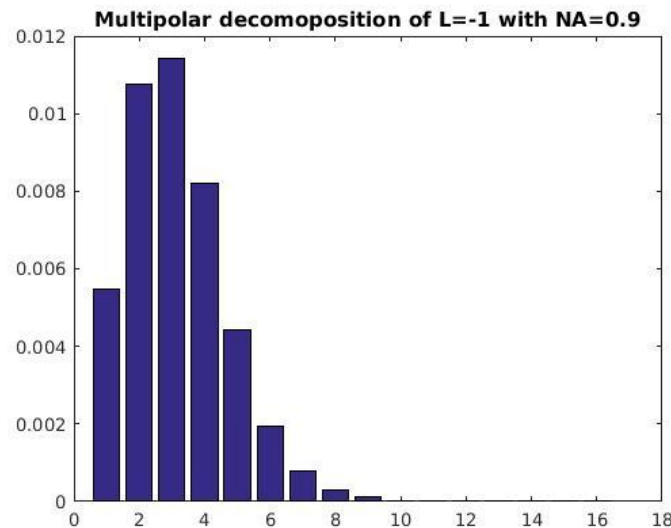
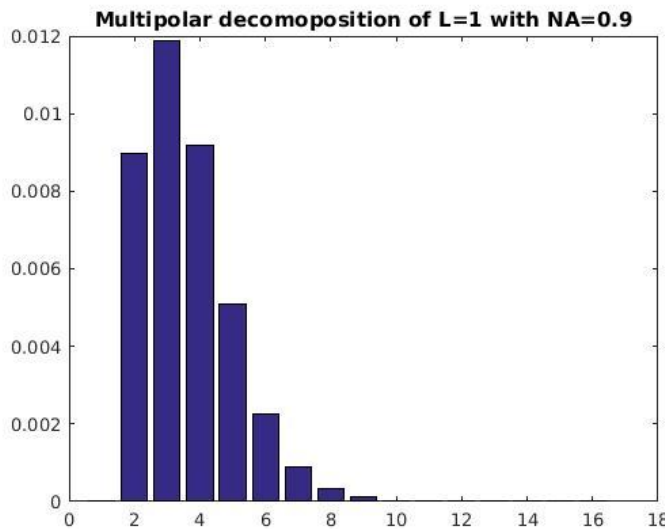
- ✓ Rigorous picture: Simulation | Mie Theory
- ✓ **Mie Theory**: Given a particle (R, n_r), then
 - ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$
 - ✓ The 'fine details' (narrow resonances) about scattering are found in the higher order modes
 - ✓ What's the difference between the multipolar decomposition for $l = \pm 1$?

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

✓ Rigorous picture: Simulation | Mie Theory

✓ **Mie Theory**: Given a particle (R, n_r), then

- ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ The 'fine details' (narrow resonances) about scattering are found in the higher order modes
- ✓ What's the difference between the multipolar decomposition for $l = \pm 1$?

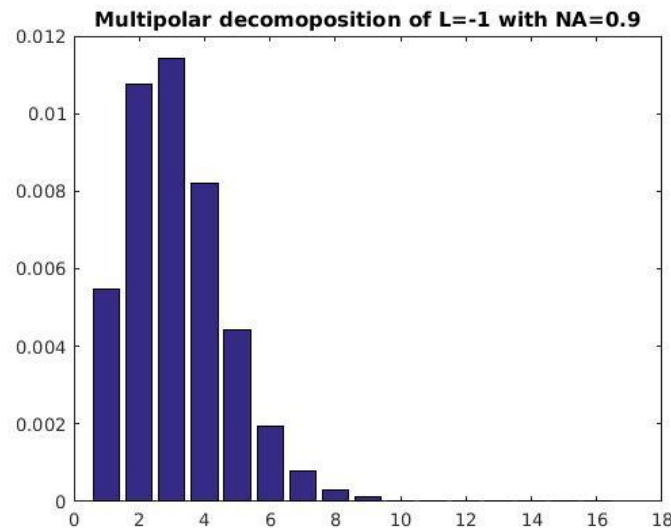
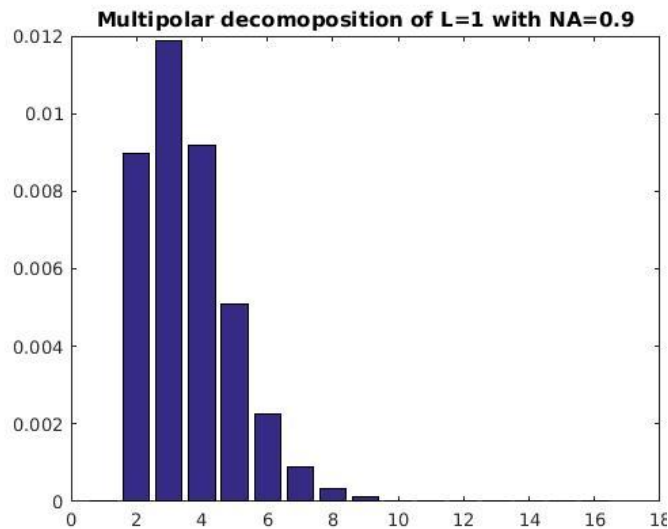


Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

✓ Rigorous picture: Simulation | Mie Theory

✓ **Mie Theory**: Given a particle (R, n_r), then

- ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ The 'fine details' (narrow resonances) about scattering are found in the higher order modes
- ✓ What's the difference between the multipolar decomposition for $l = \pm 1$?



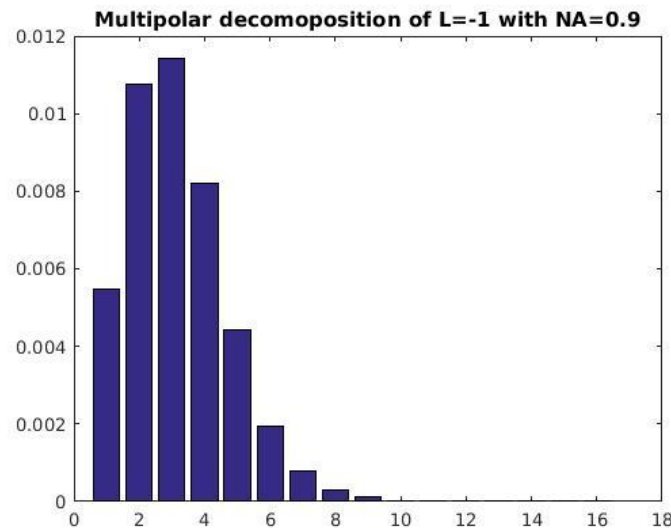
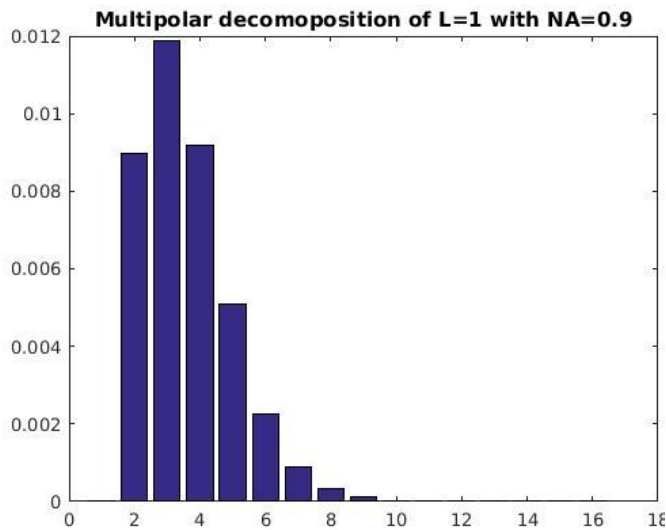
Differences in
 $j=1,2,3$ then it's
almost equal!

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)

✓ Rigorous picture: Simulation | Mie Theory

✓ **Mie Theory:** Given a particle (R, n_r), then

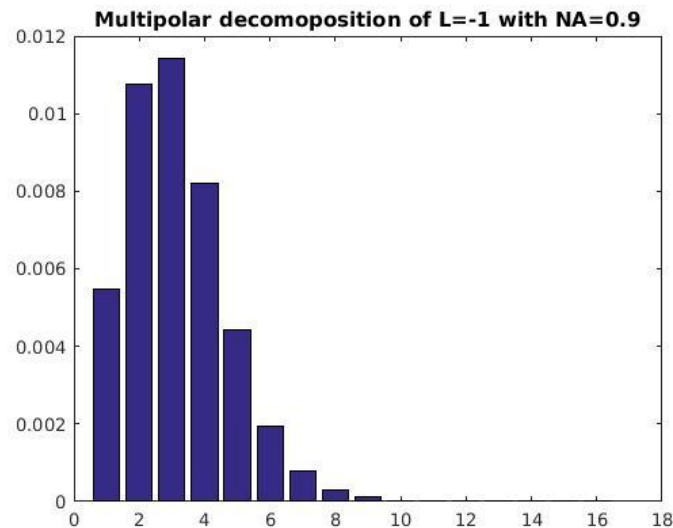
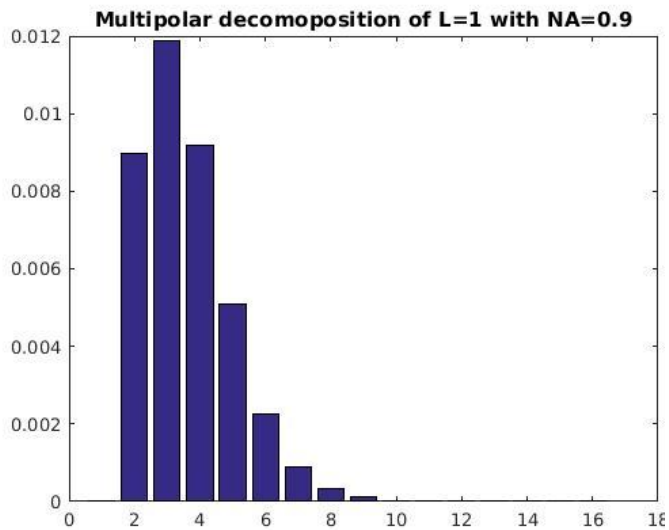
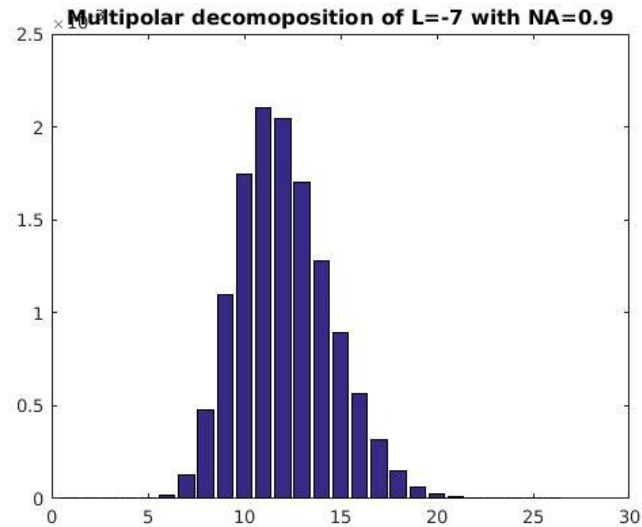
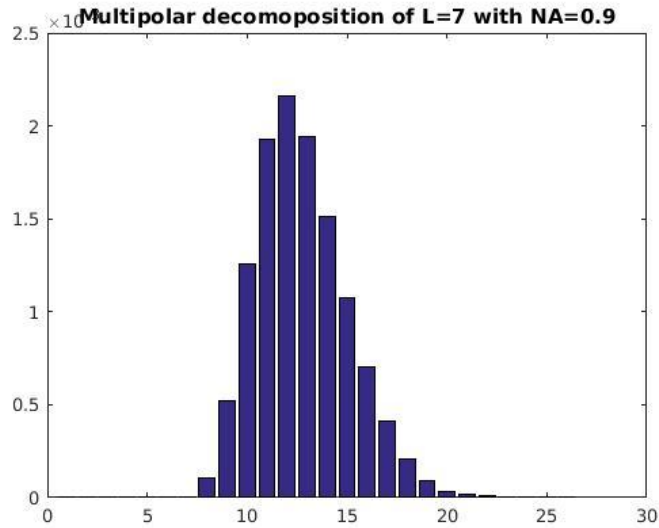
- ✓ The scattering is composed of j^* modes, where $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ The 'fine details' (narrow resonances) about scattering are found in the higher order modes
- ✓ What's the difference between the multipolar decomposition for $l = \pm 1$?



Differences in
 $j=1,2,3$ then it's
almost equal!

✓ And in between $l = \pm 7$?

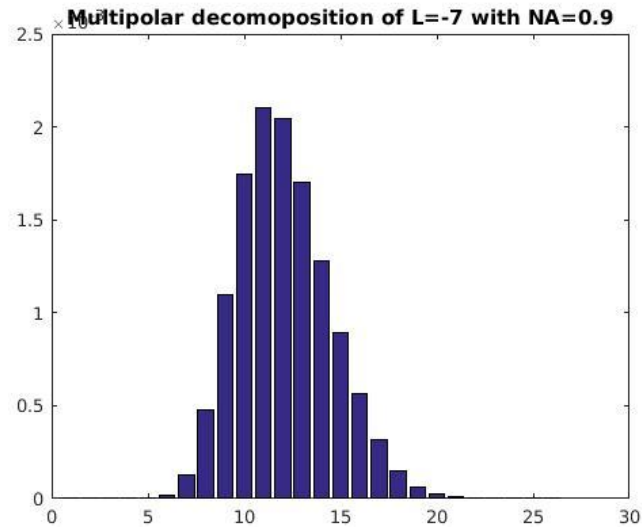
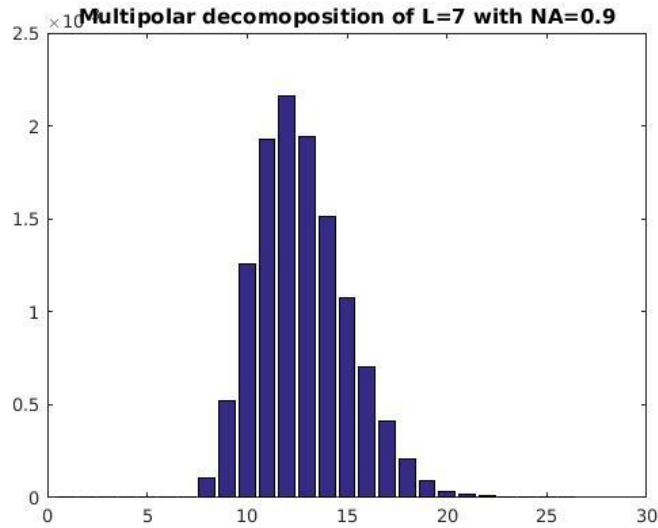
Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)



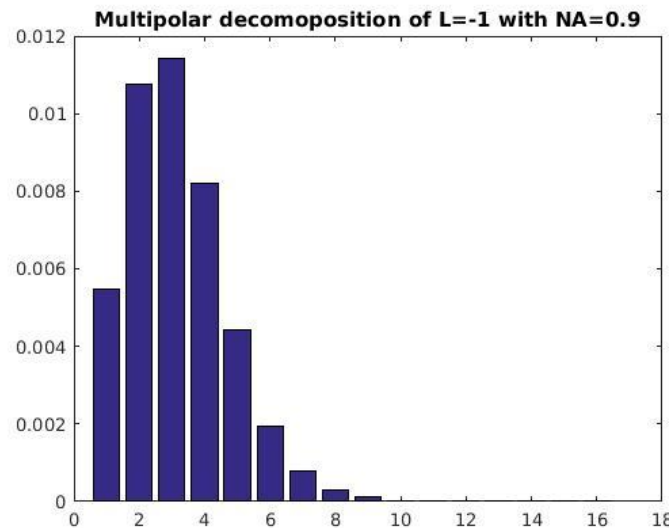
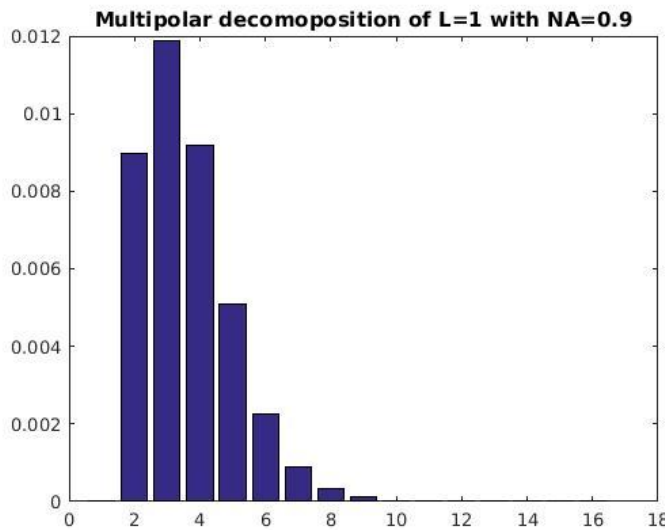
Differences in $j=1,2,3$ then it's almost equal!

✓ And in between $l = \pm 7$?

Why do $l = \pm 1$ behave equally, while $l = \pm 7$ don't? (II)



The decomposition is different in modes where $j \sim j^*$



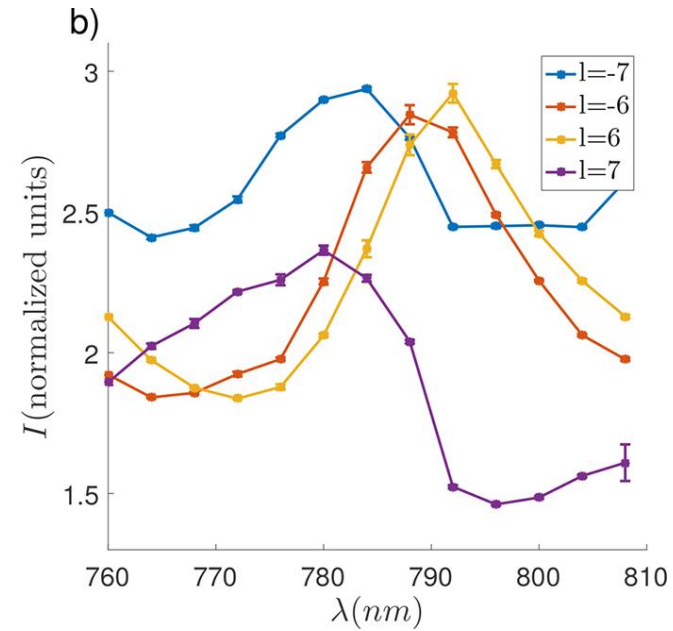
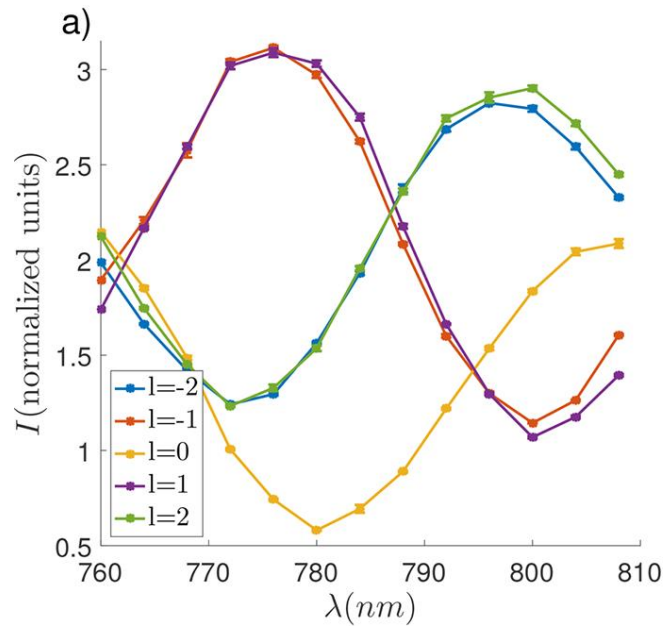
Differences in $j=1,2,3$ then it's almost equal!

✓ And in between $l = \pm 7$?



Why do $l = 0$ and $l = 7$ have different linewidths?

Why do $l = 0$ and $l = 7$ have different linewidths?



Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

Why do $l = 0$ and $l = 7$ have different linewidths?

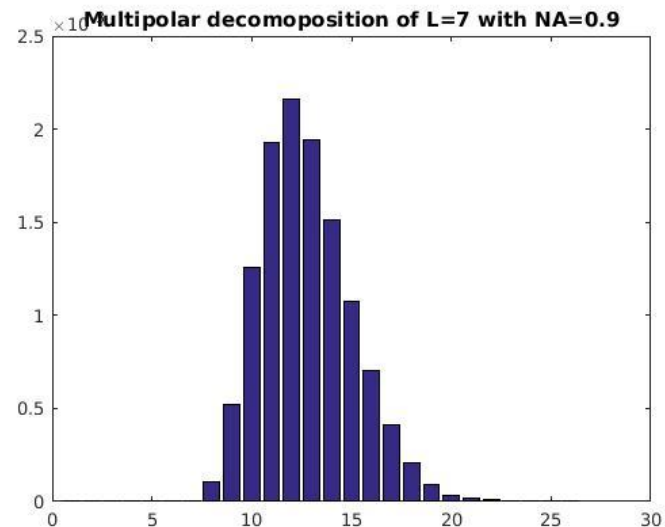
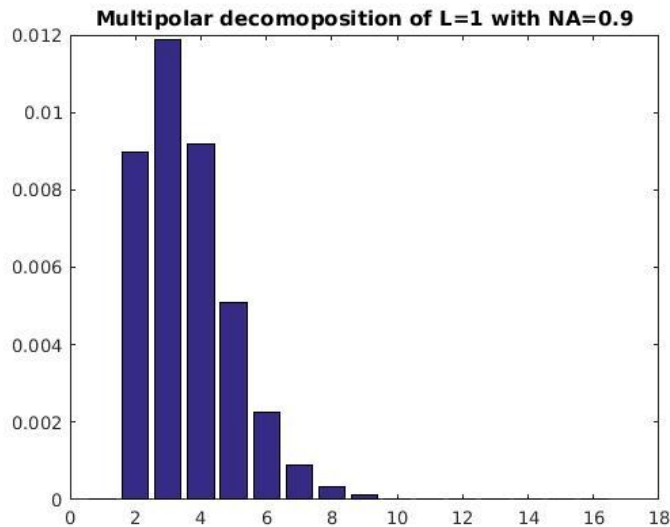
$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

- ✓ Angular momentum suppresses the scattering from the first $j = |m_z|$ modes

Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

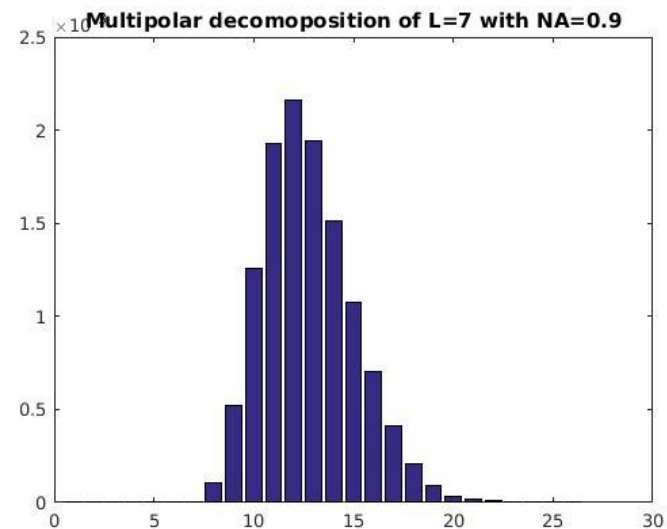
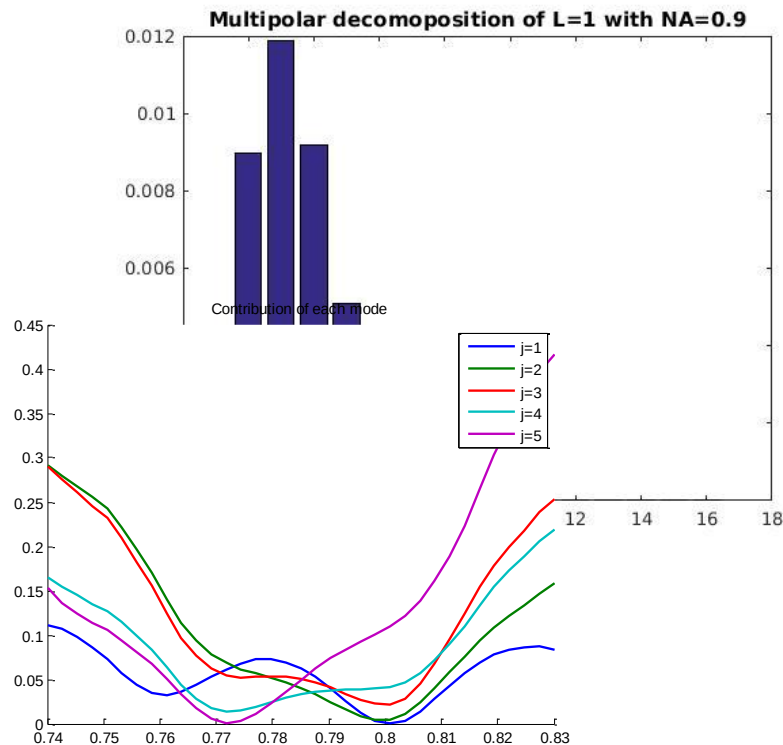
- ✓ Angular momentum suppresses the scattering from the first $j = |m_z|$ modes



Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

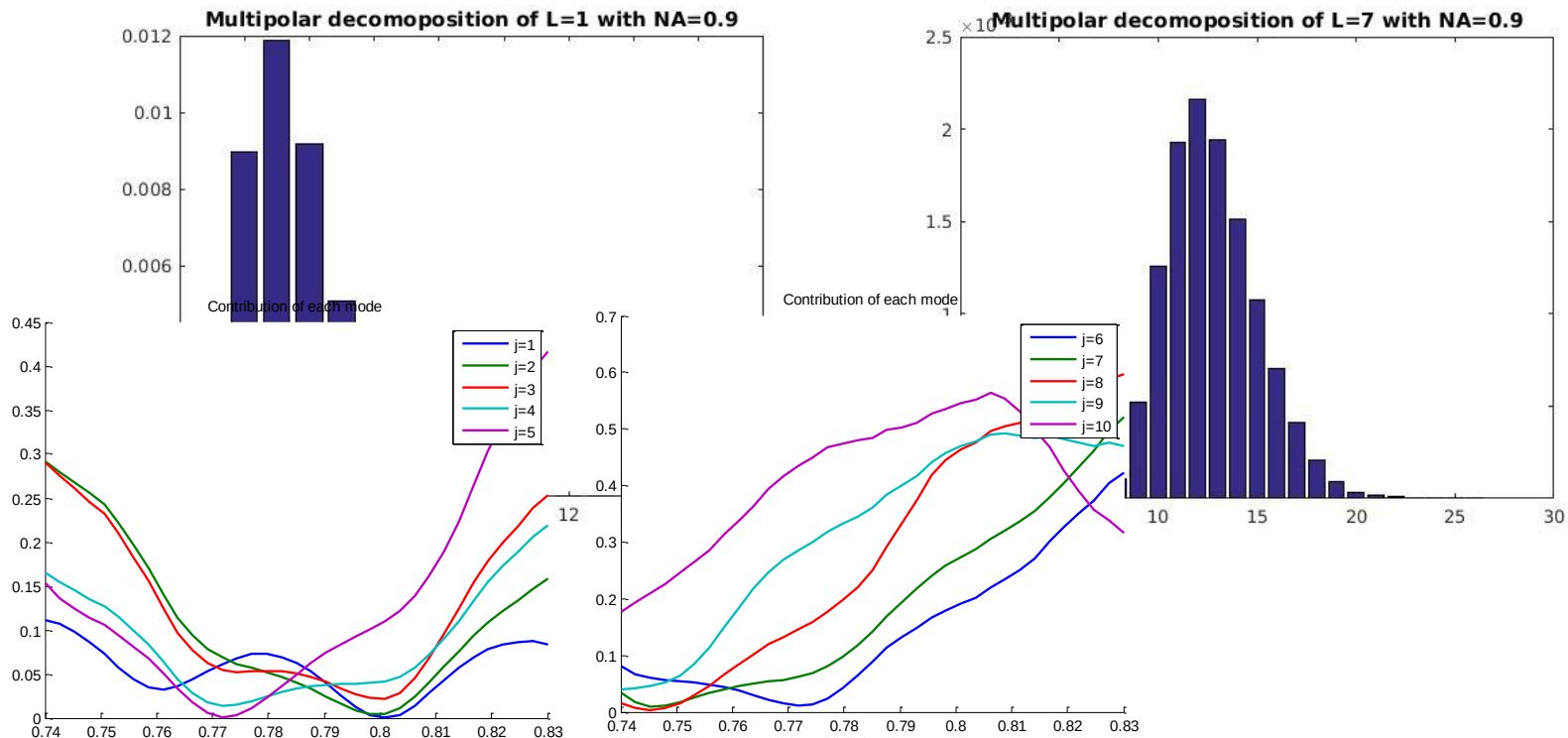
- ✓ Angular momentum suppresses the scattering from the first $j = |m_z|$ modes



Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

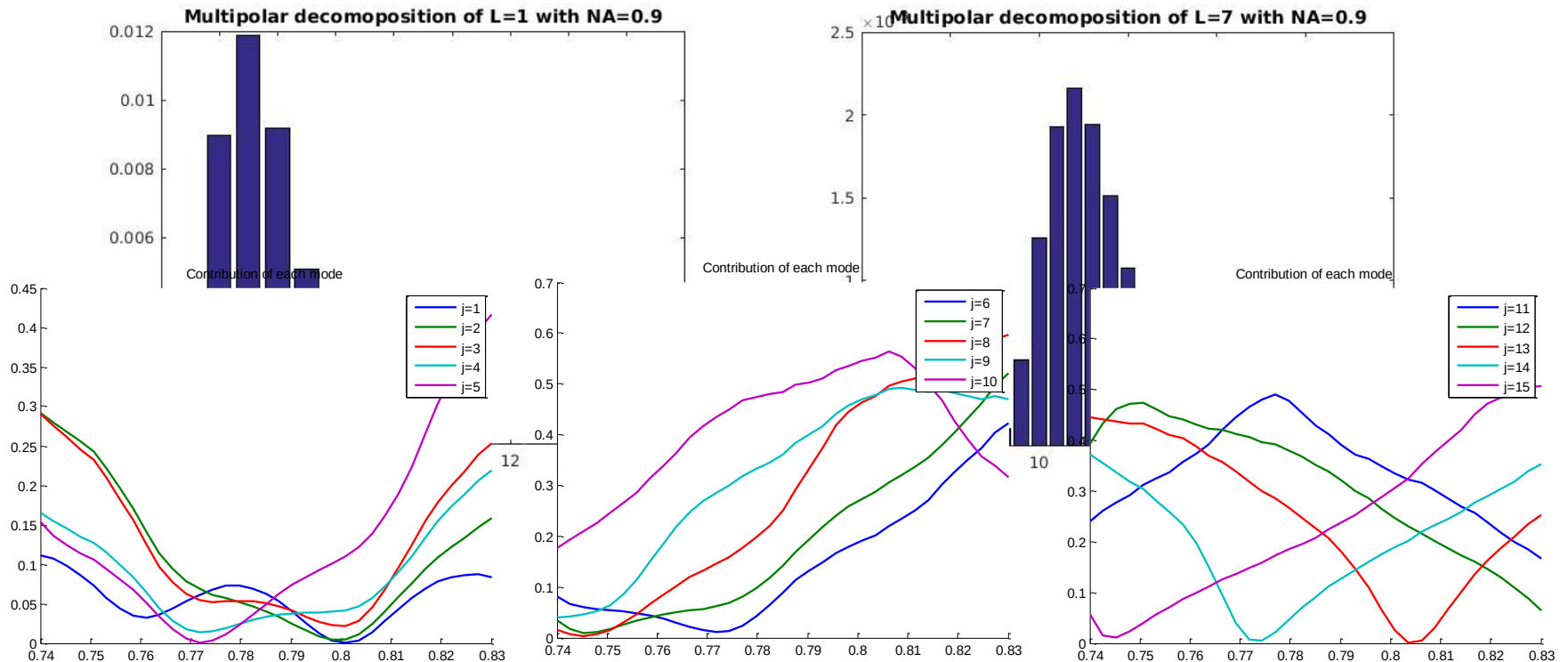
✓ Angular momentum suppresses the scattering from the first $j = |m_z|$ modes



Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

✓ Angular momentum suppresses the scattering from the first $j = |m_z|$ modes

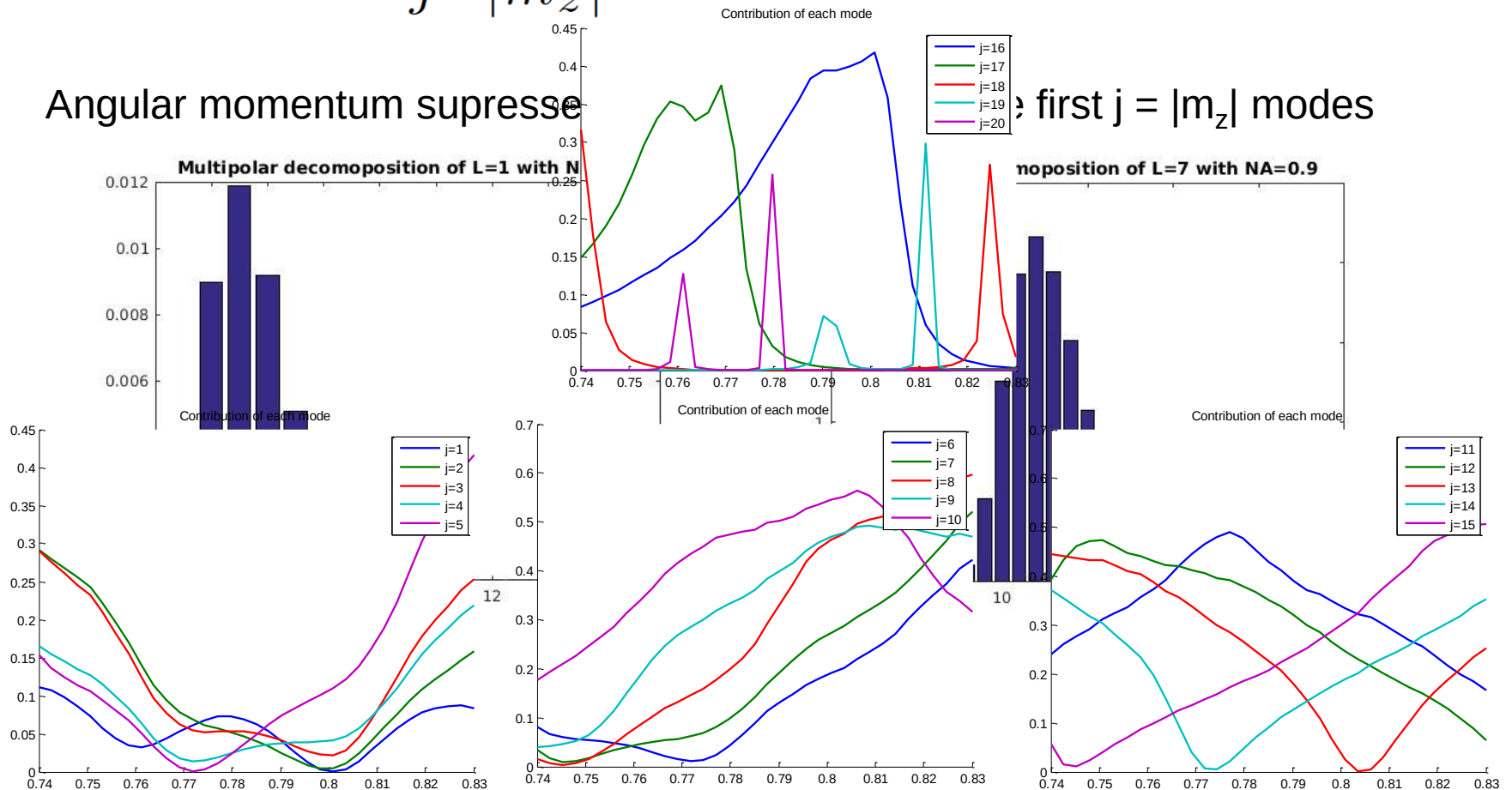


Why do $l = 0$ and $l = 7$ have different linewidths?

$$C_s \propto \sum_{j=|m_z|}^{\infty} |C_{jm_z}|^2 (|a_j|^2 + |b_j|^2)$$

Angular momentum suppression

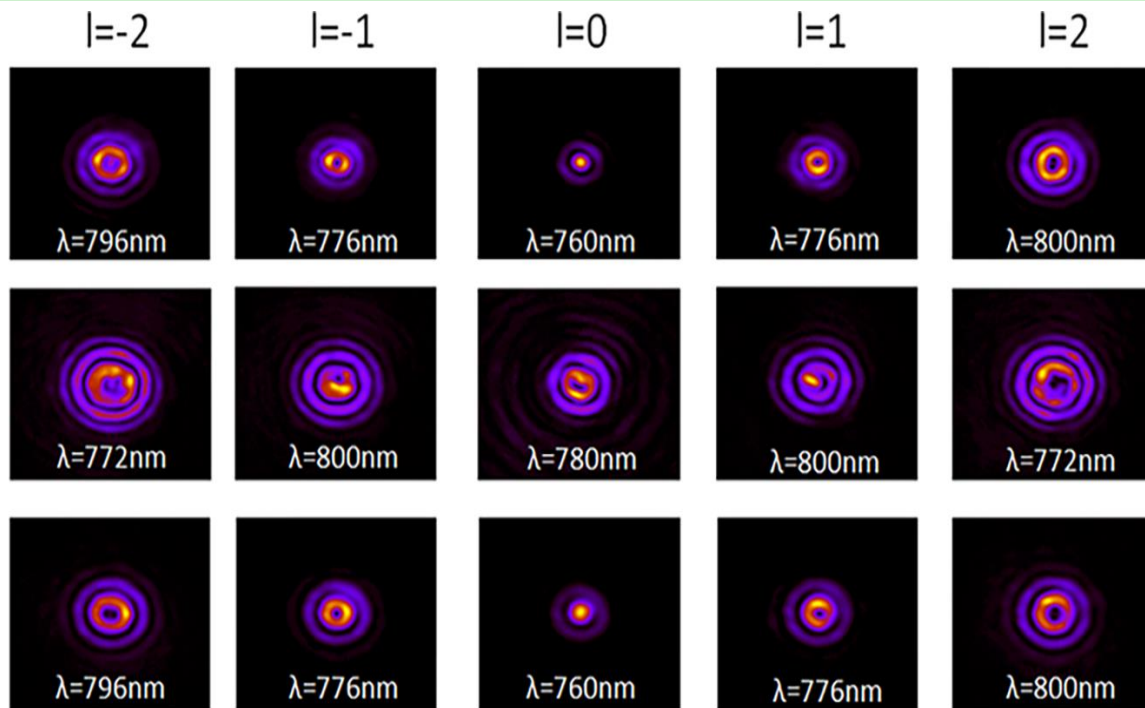
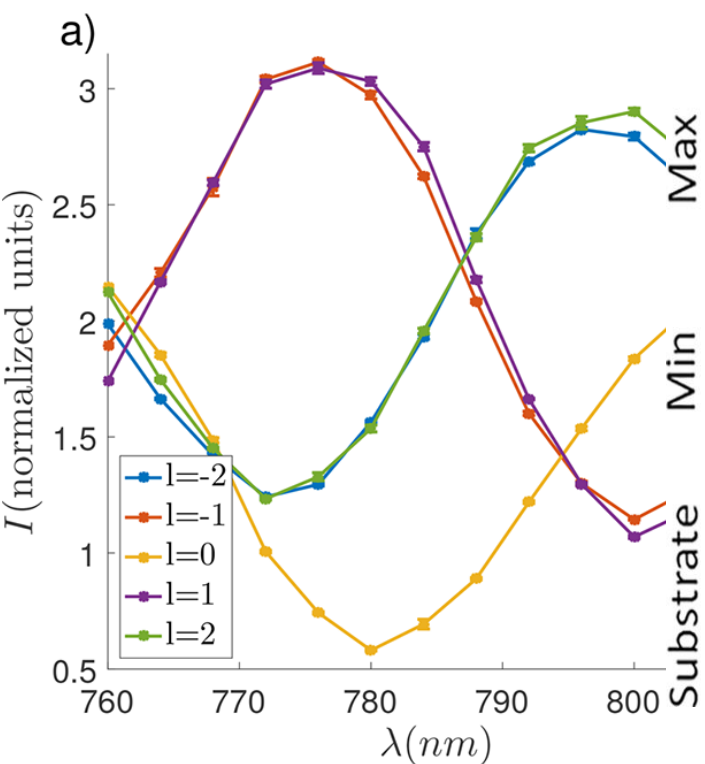
first $j = |m_z|$ modes



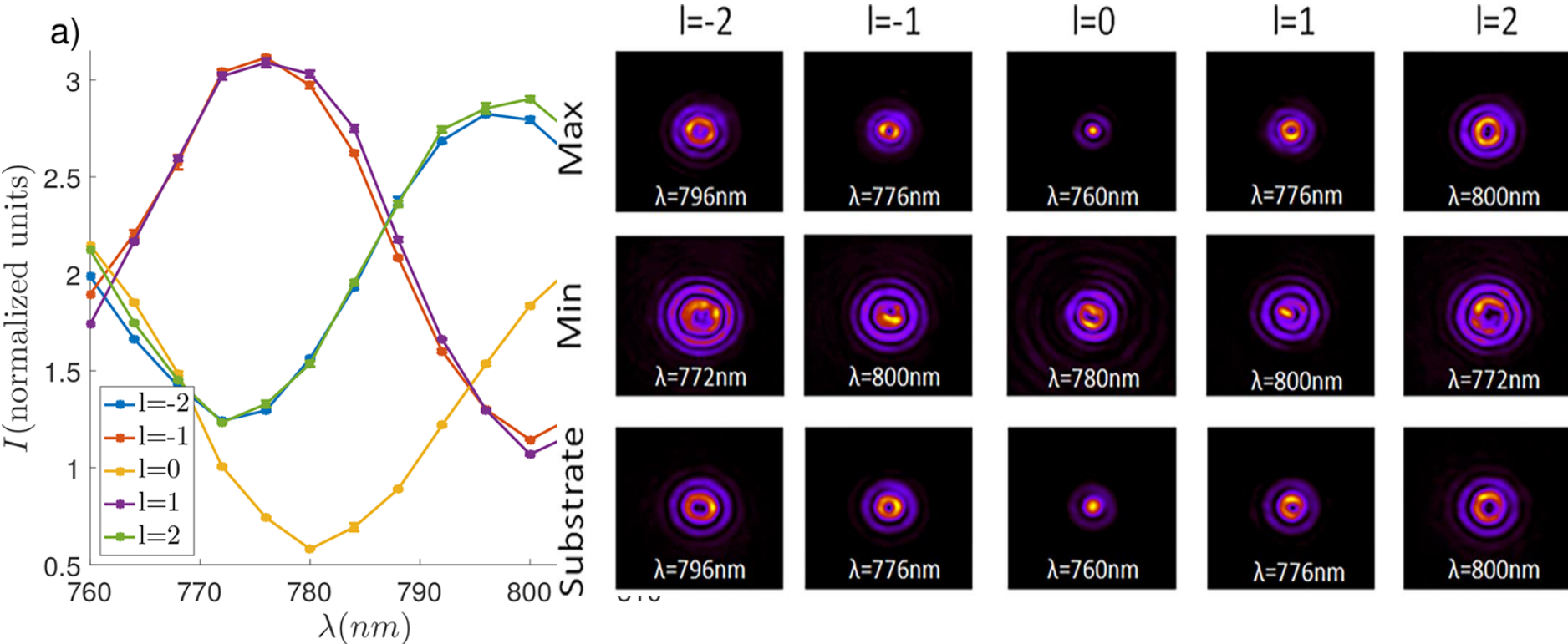


Why do we see different intensity patterns? (I)

Why do we see different intensity patterns? (I)

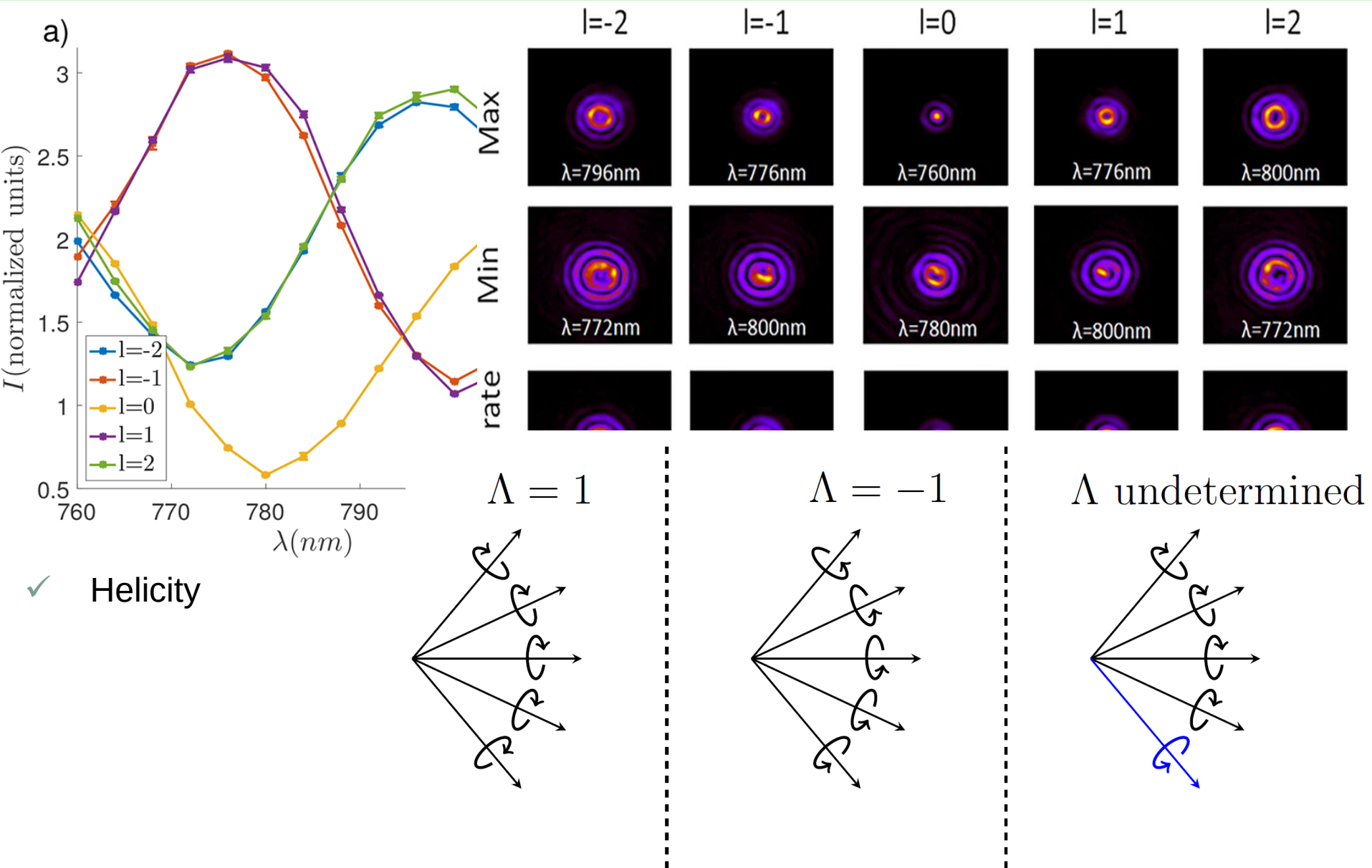


Why do we see different intensity patterns? (I)

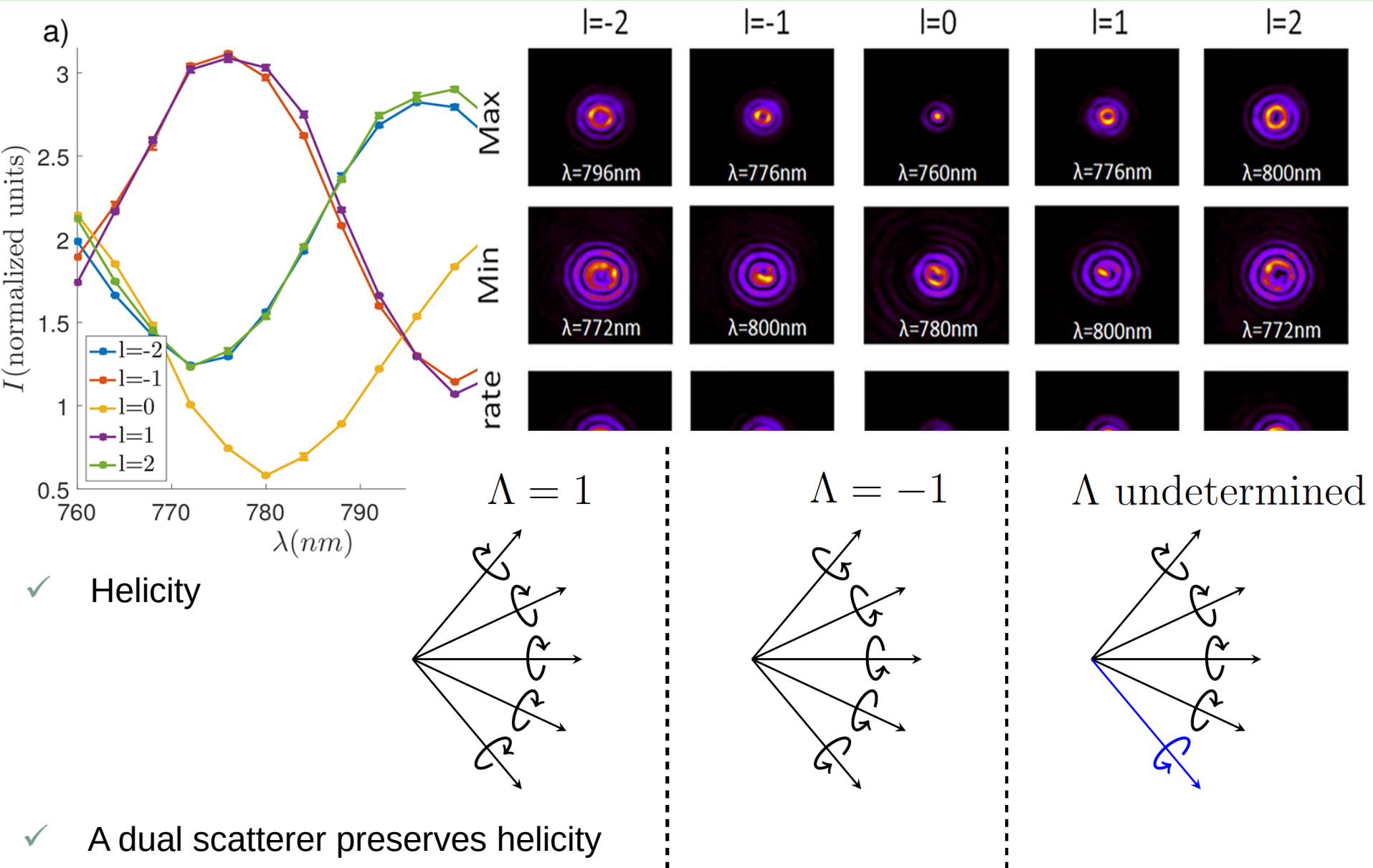


✓ Helicity

Why do we see different intensity patterns? (I)



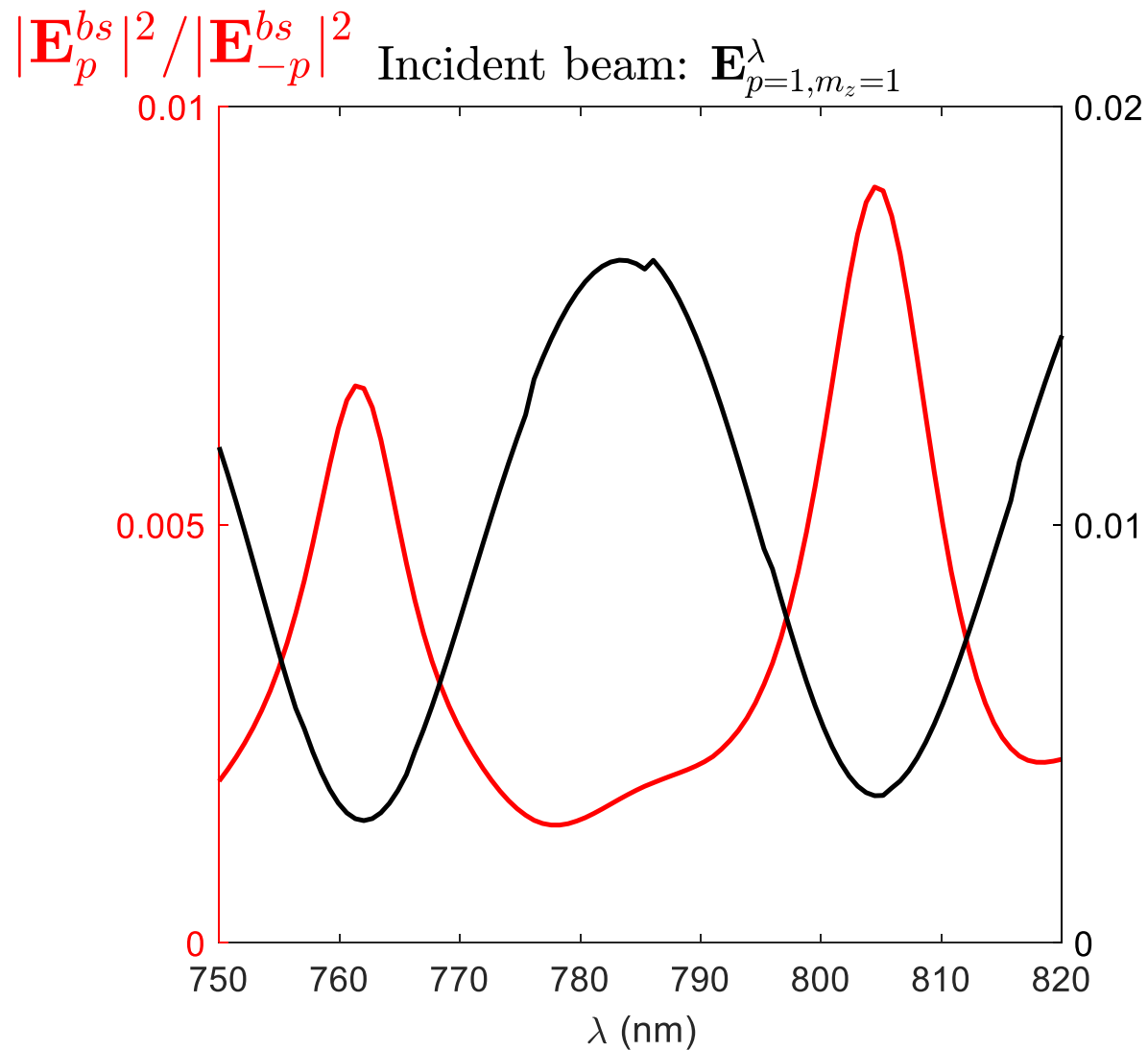
Why do we see different intensity patterns? (I)



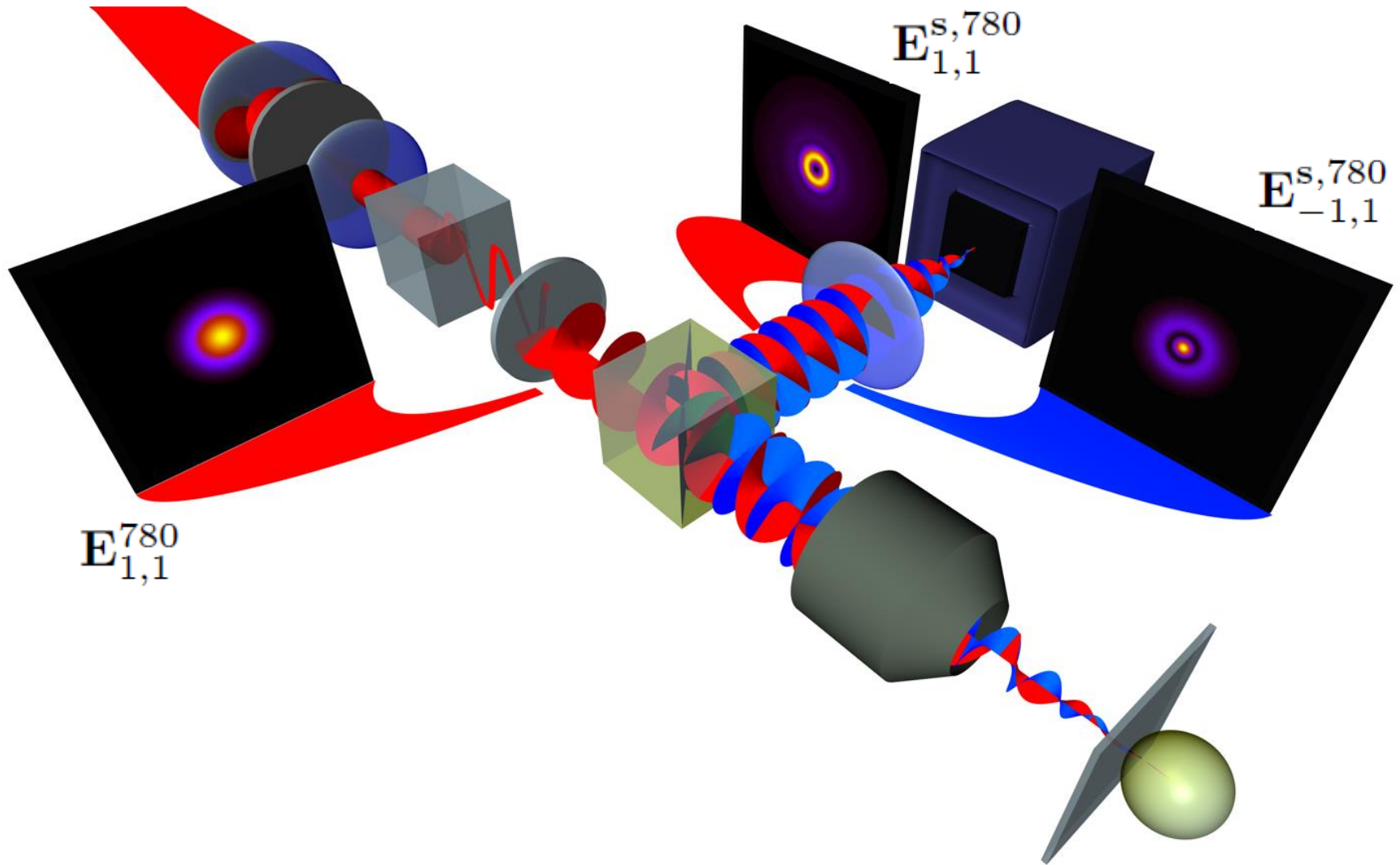


Why do we see different intensity patterns? (II)

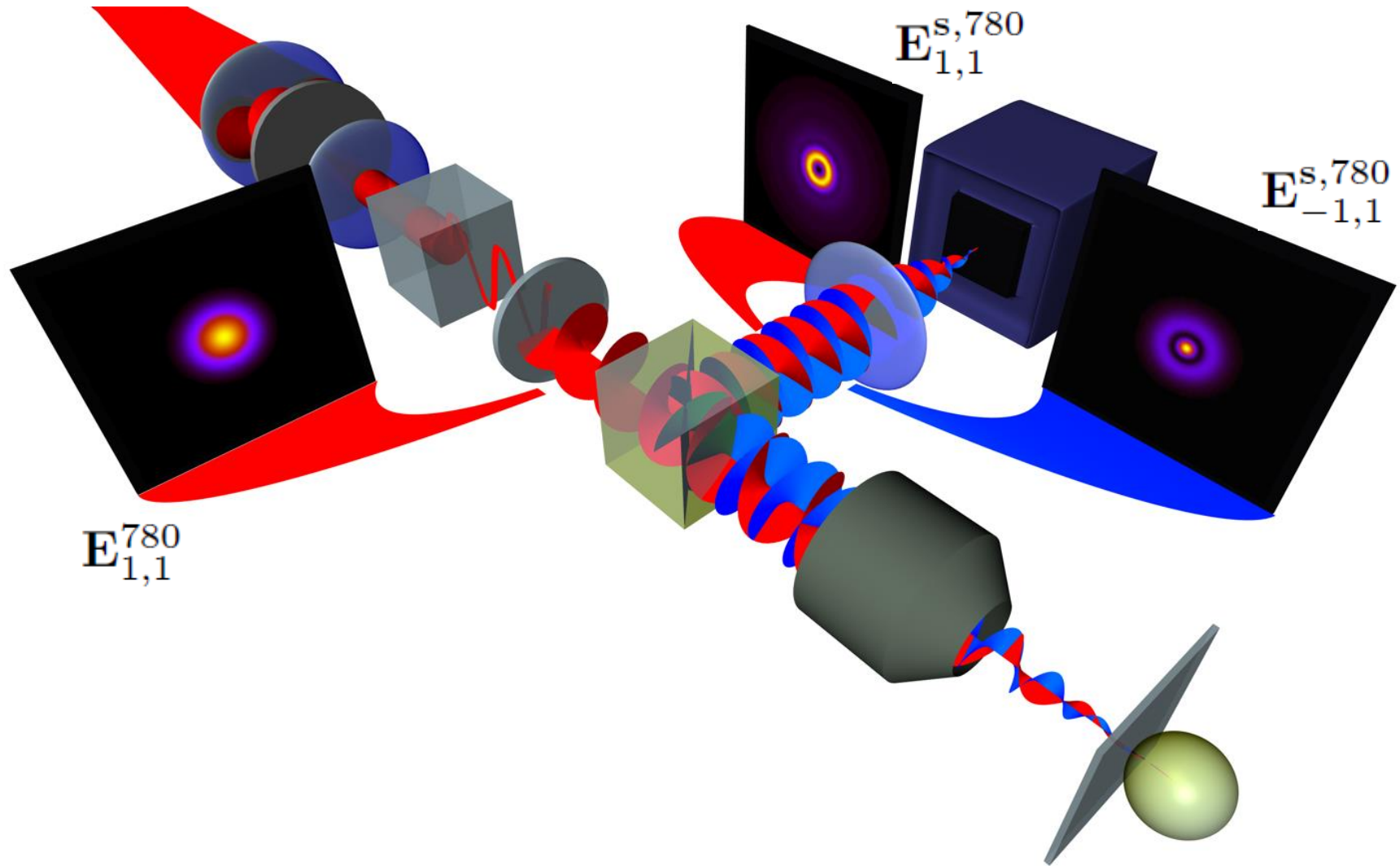
Why do we see different intensity patterns? (II)



Why do we see different intensity patterns? (II)



Why do we see different intensity patterns? (II)



- ✓ Multipoles with opposite helicity behave similarly in opposite semispaces

- ✓ Helicity and Angular momentum play a crucial role in light scattering

- ✓ Helicity and Angular momentum play a crucial role in light scattering
- ✓ Increasing the Angular Momentum of the illumination can spectrally tune and sharpen the scattered power

- ✓ Helicity and Angular momentum play a crucial role in light scattering
- ✓ Increasing the Angular Momentum of the illumination can spectrally tune and sharpen the scattered power
- ✓ Interesting effects arise when we tailor the illumination so that we excite the multipolar modes $j^* \approx 2\pi R/\lambda \approx 16$

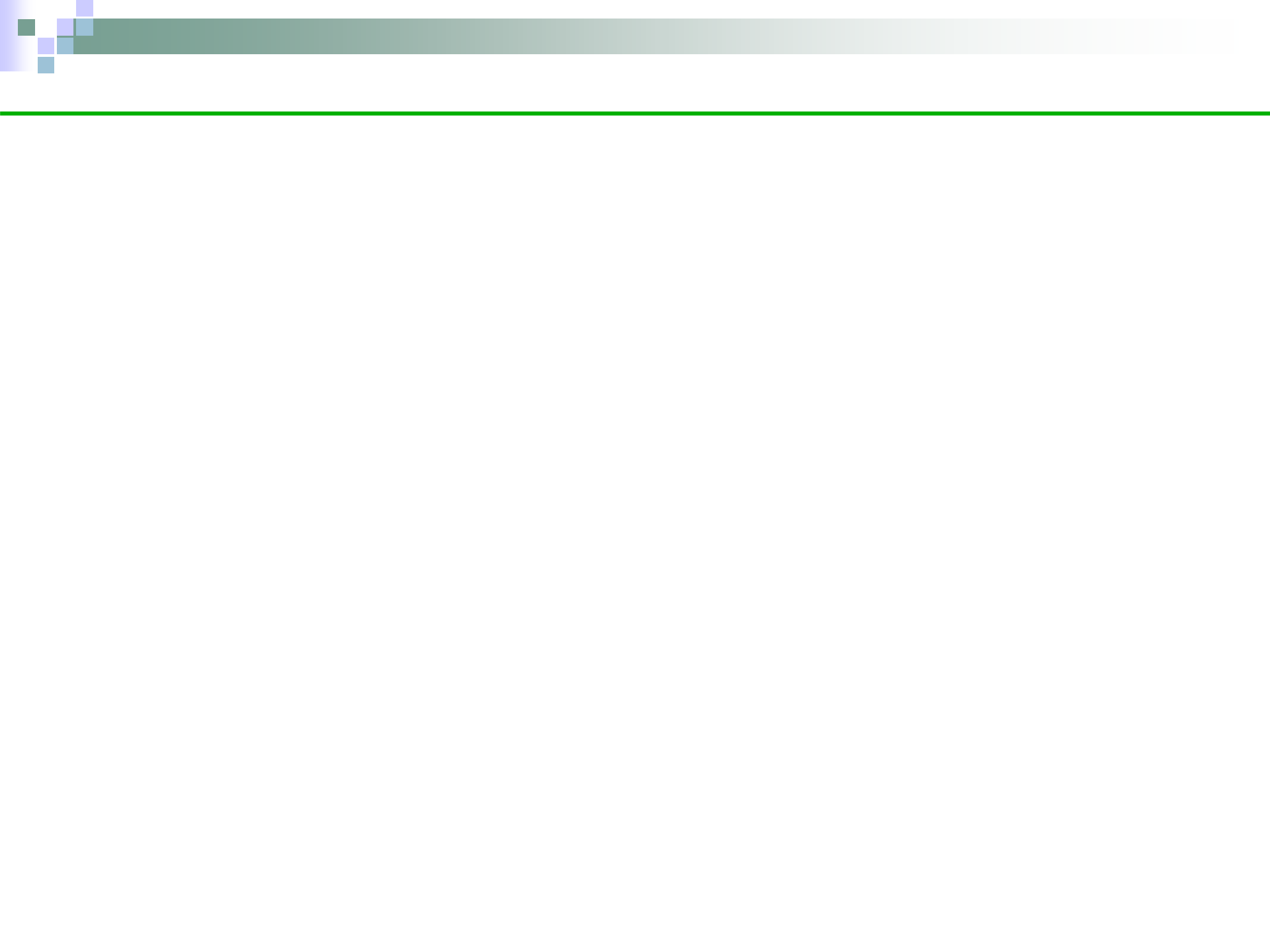
- ✓ Helicity and Angular momentum play a crucial role in light scattering
- ✓ Increasing the Angular Momentum of the illumination can spectrally tune and sharpen the scattered power
- ✓ Interesting effects arise when we tailor the illumination so that we excite the multipolar modes $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ Q_1 : Does the scattering change if the topological charge $\mathbf{q}$ changes sign?
- ✓ A_1 : It depends

- ✓ Helicity and Angular momentum play a crucial role in light scattering
- ✓ Increasing the Angular Momentum of the illumination can spectrally tune and sharpen the scattered power
- ✓ Interesting effects arise when we tailor the illumination so that we excite the multipolar modes $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ Q_1 : Does the scattering change if the topological charge $\mathbf{q}$ changes sign?
- ✓ A_1 : It depends
- ✓ Q_2 : Given a topological charge $\mathbf{q}$, does the sense of circular polarization matter?
- ✓ A_2 : It depends

- ✓ Helicity and Angular momentum play a crucial role in light scattering
- ✓ Increasing the Angular Momentum of the illumination can spectrally tune and sharpen the scattered power
- ✓ Interesting effects arise when we tailor the illumination so that we excite the multipolar modes $j^* \approx 2\pi R/\lambda \approx 16$
- ✓ Q_1 : Does the scattering change if the topological charge $\mathbf{q}$ changes sign?
- ✓ A_1 : It depends
- ✓ Q_2 : Given a topological charge $\mathbf{q}$, does the sense of circular polarization matter?
- ✓ A_2 : It depends
- ✓ Q_3 : Does the focusing regime change the answers to Q_1 and Q_2 ?
Symmetries of the problem don't change
- ✓ A_3 : Yes

- 
- ✓ X. Zambrana-Puyalto, X. Vidal, P. Wozniak, P. Banzer, and G. Molina-Terriza, “Tailoring Multipolar Mie Scattering with Helicity and Angular Momentum”, ACS Photonics **5**, 2936-2944 (2018)

- ✓ X. Zambrana-Puyalto, X. Vidal, P. Wozniak, P. Banzer, and G. Molina-Terriza, "Tailoring Multipolar Mie Scattering with Helicity and Angular Momentum", *ACS Photonics* **5**, 2936-2944 (2018)



- 
-
- ✓ Take the expression of the incident field

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formalism

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formalism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formalism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formalism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

$$I_q^{L/R} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\mathbf{E}_{+1/-1,q}^{\text{t}}|^2 dx dy,$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formalism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

$$I_q^{L/R} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\mathbf{E}_{+1/-1,q}^{\text{t}}|^2 dx dy, \longrightarrow$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

$$I_q^{L/R} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\mathbf{E}_{+1/-1,q}^{\text{t}}|^2 dx dy, \longrightarrow I_q^{L/R} = I_{-q}^{R/L}$$

- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

$$I_q^{L/R} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\mathbf{E}_{+1/-1,q}^{\text{t}}|^2 dx dy, \longrightarrow I_q^{L/R} = I_{-q}^{R/L}$$


- ✓ Take the expression of the incident field

$$\mathbf{E}_{p,q}^{\text{in}} = A \rho^q e^{(iq\phi + ikz)} e^{(-\rho^2/w_0)} e^{-i\omega t} \hat{\mathbf{e}}_p$$

- ✓ Compute the transmitted field through the hole with a convolution using the Green dyadic formulism

$$\mathbf{E}_{p,q}^{\text{t}} = \overline{\mathbf{S}}\{\mathbf{E}_{p,q}^{\text{in}}\}$$

- ✓ Relate this field with its mirror symmetric counterpart

$$\mathbf{E}_{p,q}^{\text{t}} = \exp(i\alpha) M_{\{\hat{\mathbf{z}}\}} \mathbf{E}_{-p,-q}^{\text{t}}$$

- ✓ Anti-commutation relations

$$M_{\{\hat{\mathbf{z}}\}}^\dagger \Lambda M_{\{\hat{\mathbf{z}}\}} = -\Lambda M_{\{\hat{\mathbf{z}}\}}^\dagger J_z M_{\{\hat{\mathbf{z}}\}} = -J_z$$

- ✓ Q_3 can be found using the definition of CD_q and the transmitted intensity

$$I_q^{L/R} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |\mathbf{E}_{+1/-1,q}^{\text{t}}|^2 dx dy, \xrightarrow{\quad\quad\quad} I_q^{L/R} = I_{-q}^{R/L}$$

$$CD_q = -\overleftarrow{CD_{-q}}$$