

A scaling hypothesis for PEPS

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Overview

Introduction: PEPS for 2-D quantum systems and 3-D stat-mech

Variational optimization of PEPS fixed points

PEPS for critical systems: a scaling hypothesis

Outlook

Bram Vanhecke, Juraj Hasik, Frank Verstraete, LV
arXiv:2102.03143

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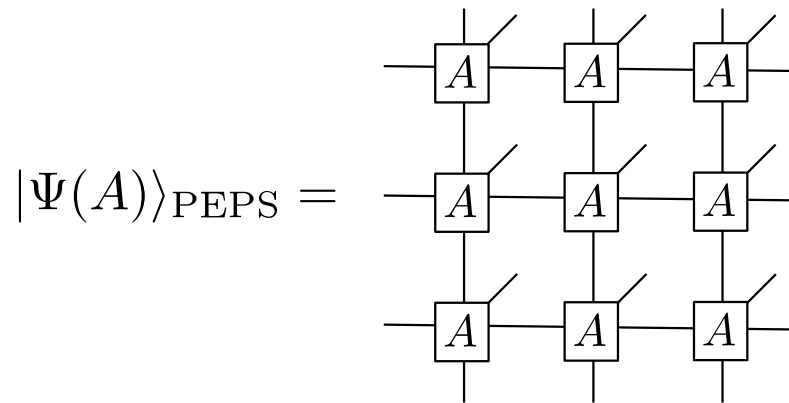
PEPS for critical systems: a scaling hypothesis

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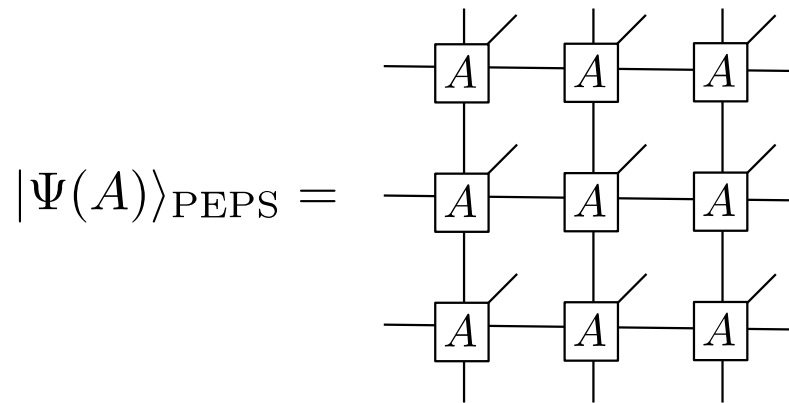
Projected entangled-pair states (PEPS) are variational wavefunctions for 2-D lattice systems



- formulated directly in the thermodynamic limit
- parametrized by a single local tensor (larger unit cells possible)

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Projected entangled-pair states (PEPS) are variational wavefunctions for 2-D lattice systems

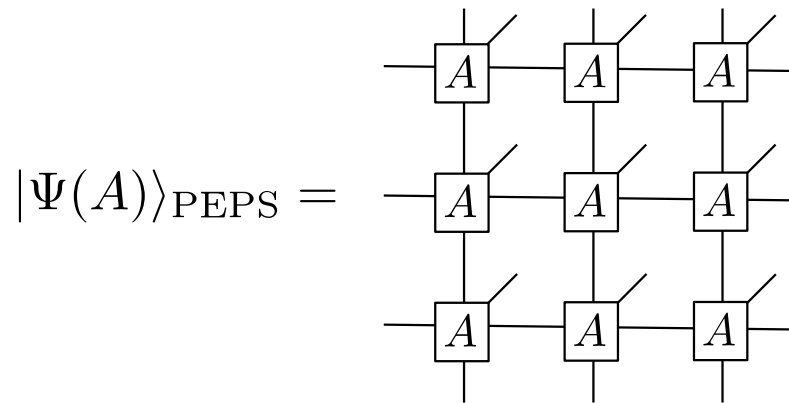


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- parametrized by a single local tensor (larger unit cells possible)

- ground-state wavefunctions for 2-D quantum lattice systems (spins, bosons, electrons)

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Projected entangled-pair states (PEPS) are variational wavefunctions for 2-D lattice systems

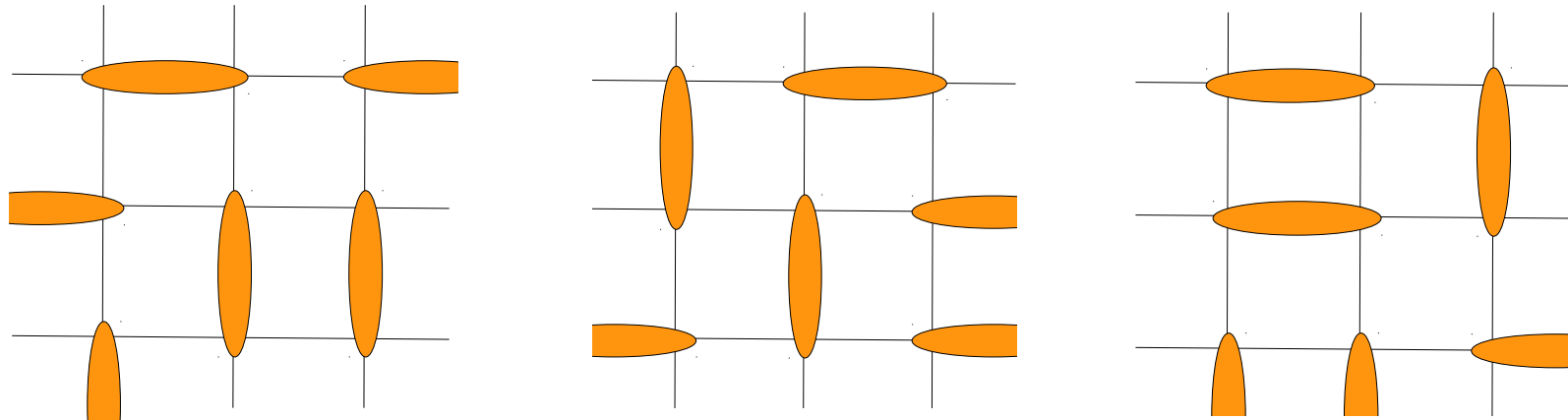


- formulated directly in the thermodynamic limit
- parametrized by a single local tensor (larger unit cells possible)

- ground-state wavefunctions for 2-D quantum lattice systems (spins, bosons, electrons)
- transfer-matrix fixed points for 3-D stat-mech systems

Introduction

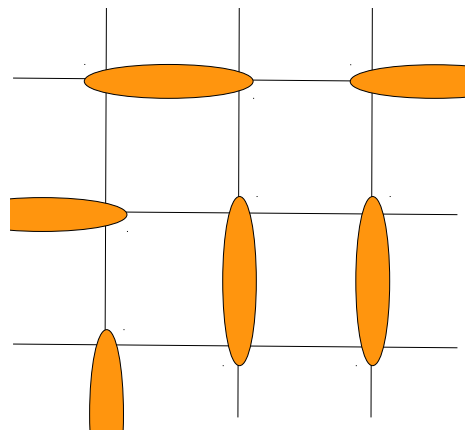
Example: dimer-covering problem



→ how does the number of dimer coverings scale with system size?

Introduction

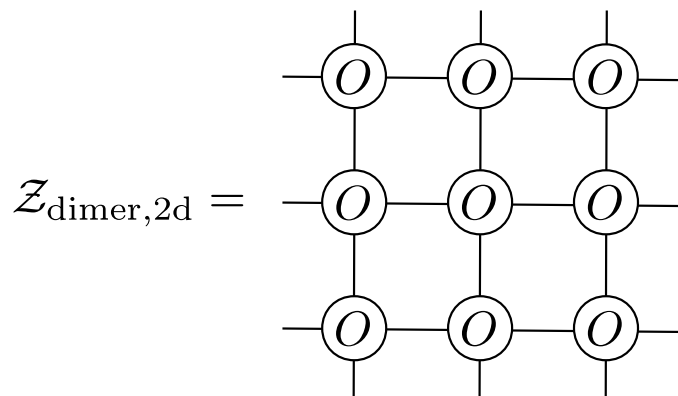
Example: dimer-covering problem



	0	1	0	0	1
0			0	0	
1	0	0	0	0	0
0	0	1	0	1	0
1	0	0	0		

dimer rule on
the vertices

Tensor network that sums all allowed configurations

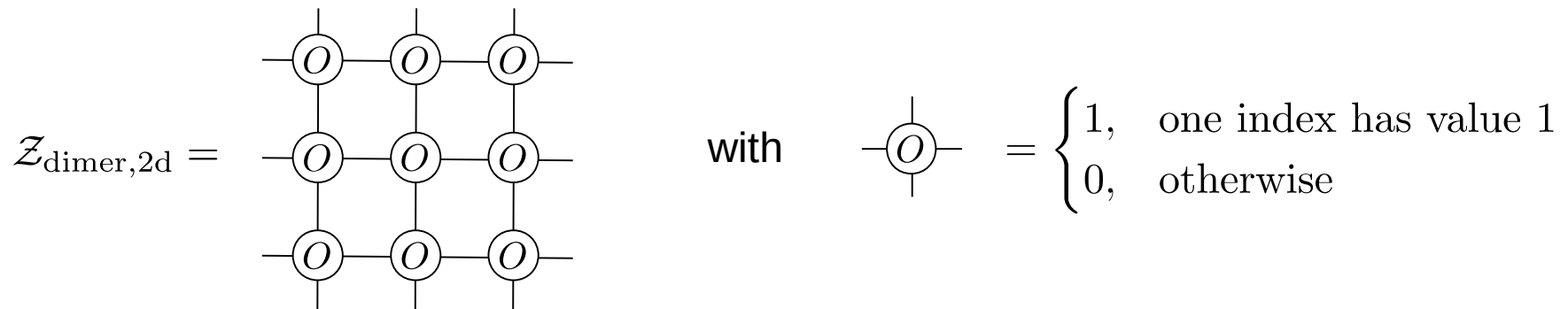


with $\text{---} \bigcirc \text{---} = \begin{cases} 1, & \text{one index has value 1} \\ 0, & \text{otherwise} \end{cases}$

→ how does this number scale with system size?

Introduction

Example: dimer-covering problem



Contraction of infinite network gives extremely precise results
(this is a critical model)

exact $S = 0.2915609040$

Baxter's CTMRG (D=6, IBM-360) $S = 0.29146021876$

boundary MPS (D=250, laptop) $S = 0.2915608913$

→ variational results!

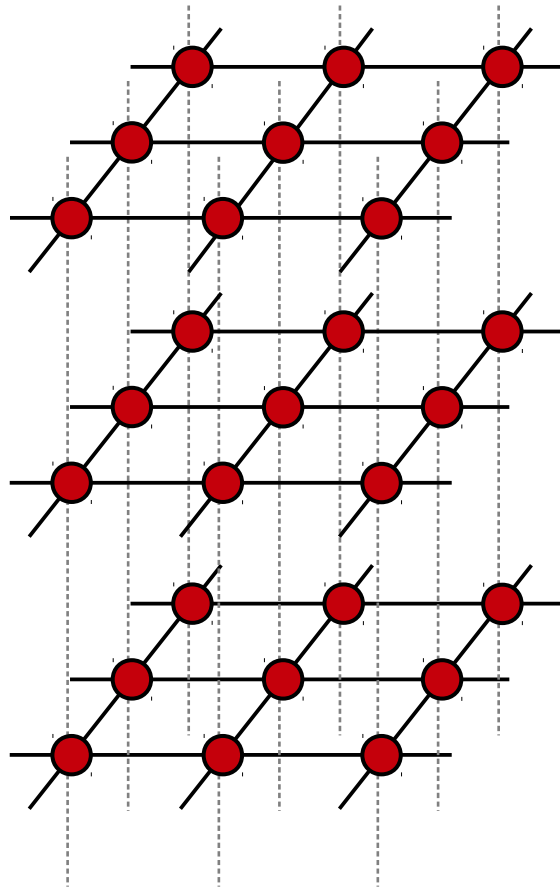
TABLE I. The values of κ/s obtained from the $r = 1, 2, \dots, 6$ approximations for $s = 1.0, 4.0, 10.0$, and ∞ , together with the values obtained by geometric extrapolation from the 4, 5, 6 results.

r	κ/s			
	$s = 1.0$	$s = 4.0$	$s = 10.0$	$s = \infty$
1	1.937416664	1.444670083	1.356095932	1.299038106
2	1.940215341	1.460590906	1.381143005	1.335033348
3	1.940215344	1.460623453	1.381458447	1.337338271
4	1.940215351	1.460629381	1.381506501	1.337984697
5	1.940215351	1.460629397	1.381508315	1.338250017
6	1.940215351	1.460629398	1.381508512	1.338380390
Extrapolated	1.940215351	1.460629398	1.381508536	1.338506344

Introduction

Example: dimer-covering problem

$$\mathcal{Z}_{\text{dimer},3\text{d}} =$$



with the tensor

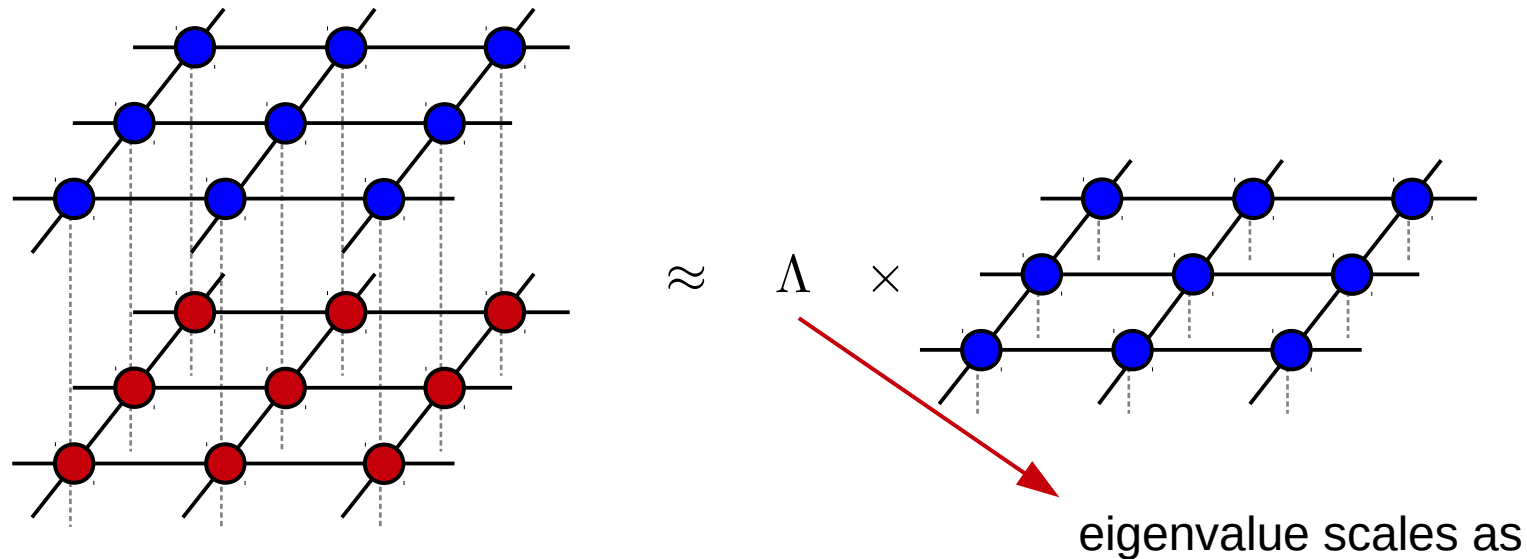
$$\text{tensor symbol} = \begin{cases} 1, & \text{one index has value 2} \\ 0, & \text{otherwise} \end{cases}$$

→ plane-to-plane transfer matrix

Introduction

Example: dimer-covering problem

→ approximate the fixed point of the transfer matrix as a PEPS



$$\Lambda = \lambda^{N_x N_y}$$

→ how do we optimize PEPS fixed points, and how do extrapolate data?

such that entropy density is given by

$$S = \log \lambda$$

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Variational optimization

An optimal PEPS approximation is characterized by the variational principle

→ for quantum Hamiltonians

$$A = \arg \min_A \left(\frac{\langle \Psi(A) | \mathcal{H} | \Psi(A) \rangle}{\langle \Psi(A) | \Psi(A) \rangle} \right)$$

gradient-based methods were developed that optimize the variational cost function directly

Corboz, PRB 94, 035133 (2016)

LV, Haegeman, Corboz, Verstraete, PRB 94, 155123 (2016)

→ for transfer matrices

$$A = \arg \min_A \left(-\log \frac{\langle \Psi(A) | \mathcal{T} | \Psi(A) \rangle}{\langle \Psi(A) | \Psi(A) \rangle} \right)$$

similar gradient-based method for optimization

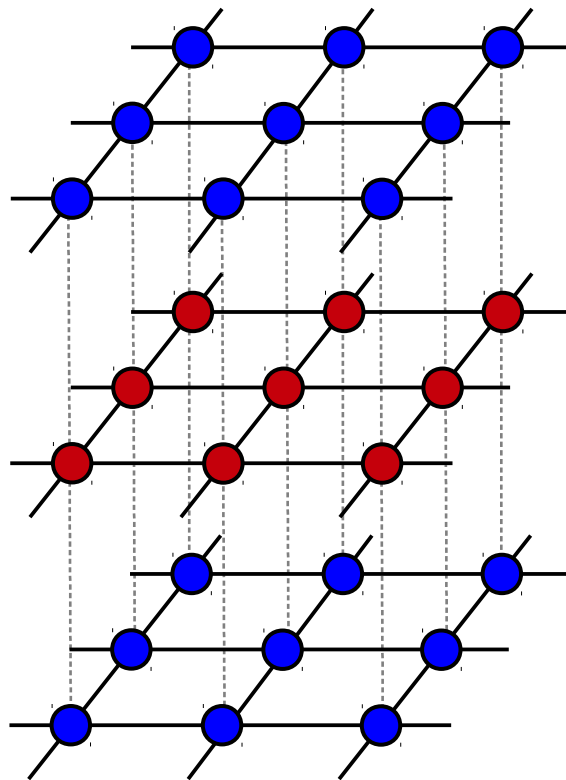
Nishino et al, Nucl. Phys. B 575, 504 (2000)

LV, Vanhecke, Verstraete, PRE 98, 042145 (2018)

Variational optimization

Direct variational optimization of the free energy

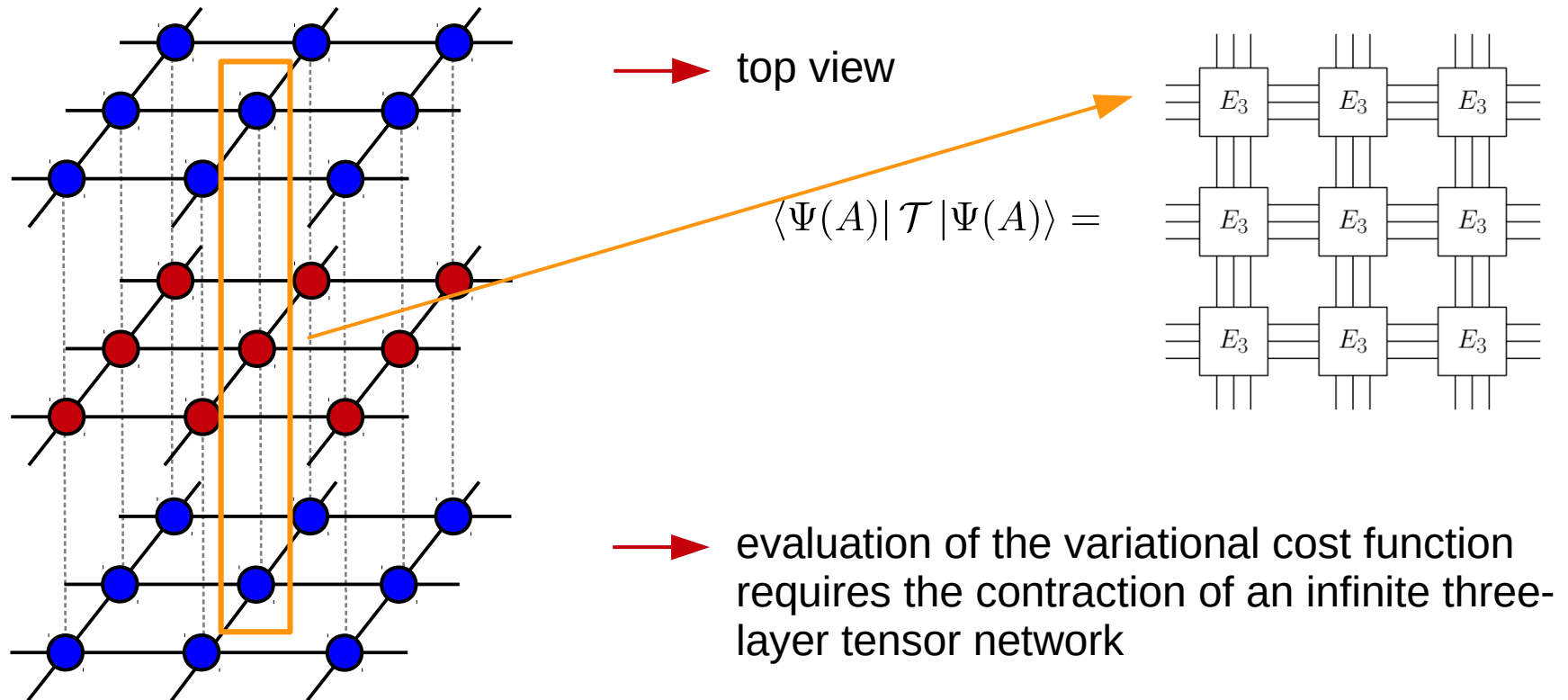
$$A = \arg \min_A \left(-\log \frac{\langle \Psi(A) | \mathcal{T} | \Psi(A) \rangle}{\langle \Psi(A) | \Psi(A) \rangle} \right)$$



Variational optimization

Direct variational optimization of the free energy

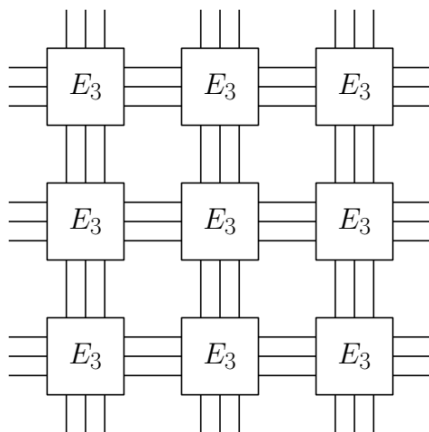
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Variational optimization

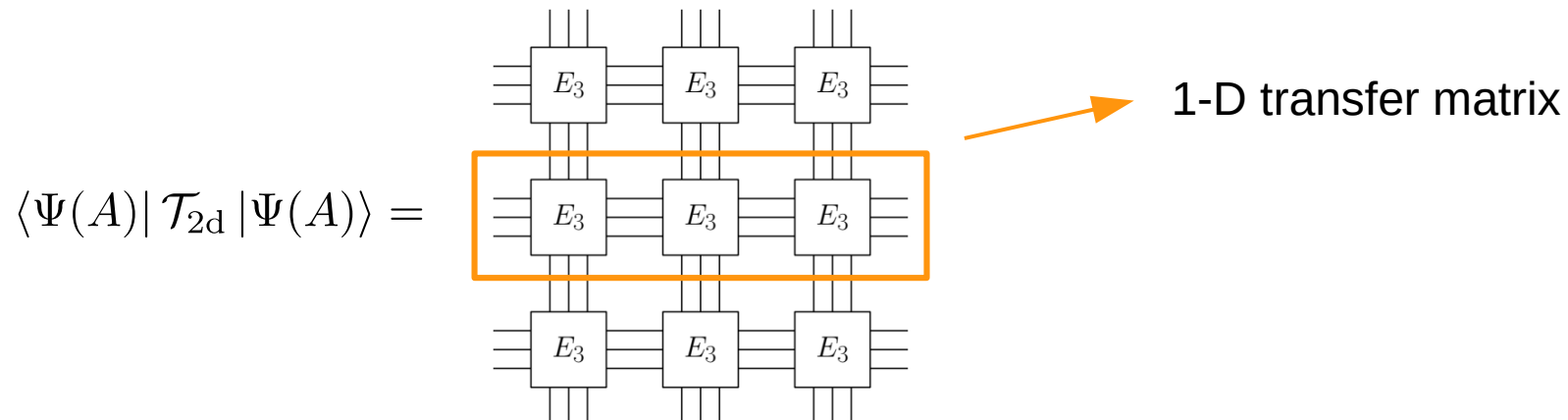
Evaluation of the variational cost function requires the contraction of an infinite three-layer tensor network

$$\langle \Psi(A) | \mathcal{T}_{2d} | \Psi(A) \rangle =$$



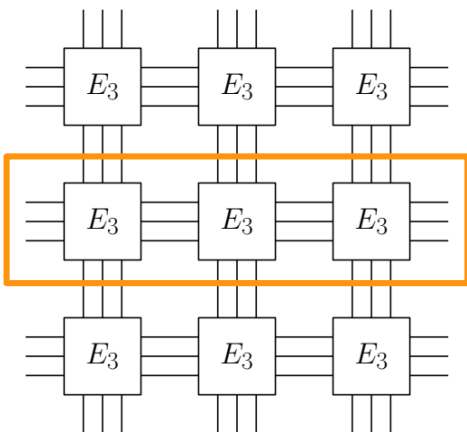
Variational optimization

Evaluation of the variational cost function requires the contraction of an infinite three-layer tensor network



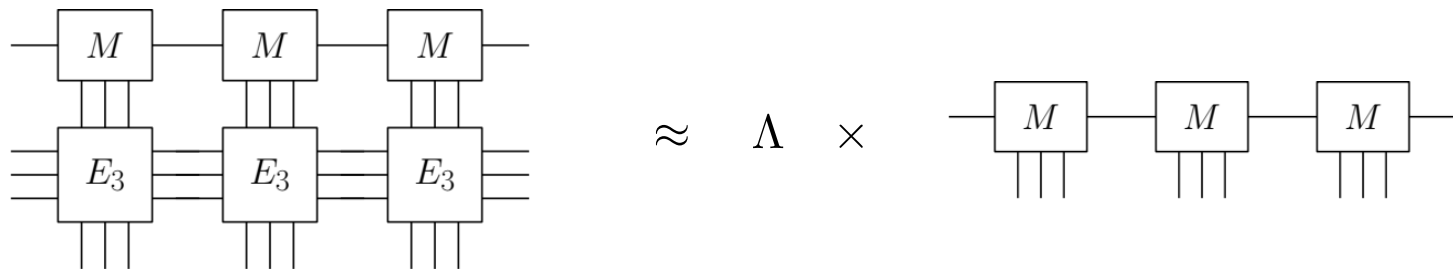
Variational optimization

Evaluation of the variational cost function requires the contraction of an infinite three-layer tensor network

$$\langle \Psi(A) | \mathcal{T}_{2d} | \Psi(A) \rangle =$$


1-D transfer matrix

→ approximate the boundary of the 1-D transfer matrix as an MPS



$$\approx \Lambda \times$$

→ apply the variational principle for the boundary contraction

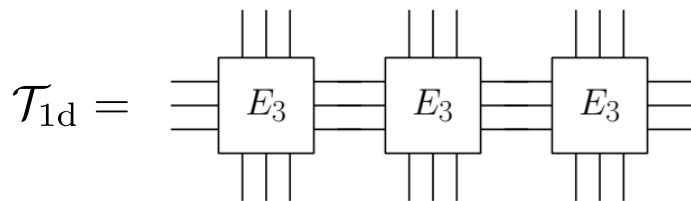
$$M = \arg \min_M \left(-\log \frac{\langle \Psi(M) | \mathcal{T}_{1d} | \Psi(M) \rangle}{\langle \Psi(M) | \Psi(M) \rangle} \right)$$

Variational optimization

Apply the variational principle for the boundary contraction

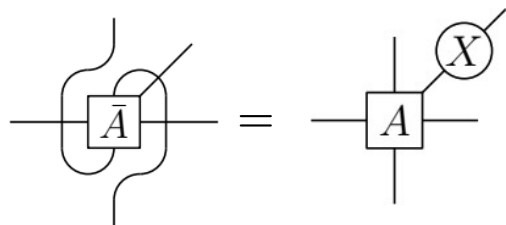
$$M = \arg \min_M \left(-\log \frac{\langle \Psi(M) | \mathcal{T}_{1d} | \Psi(M) \rangle}{\langle \Psi(M) | \Psi(M) \rangle} \right)$$

with

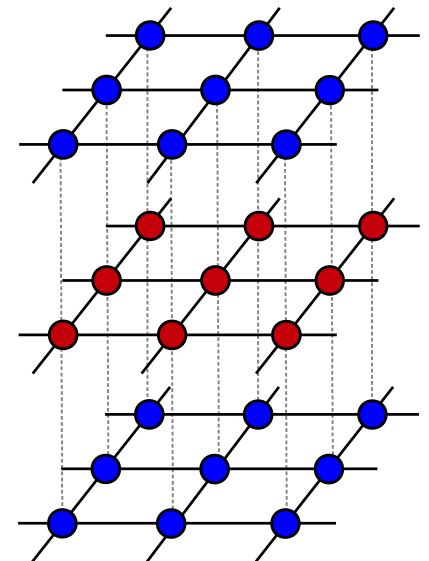


→ the transfer matrix has to be hermitian!

Impose symmetry constraints on the PEPS tensors



→ implies that transfer matrix is hermitian, and we can use variational algorithms



Variational optimization

Hierarchy of variational principles

→ boundary PEPS $A = \arg \min_A \left(-\log \frac{\langle \Psi(A) | \mathcal{T}_{2d} | \Psi(A) \rangle}{\langle \Psi(A) | \Psi(A) \rangle} \right)$

control parameter: PEPS bond dimension D

→ boundary MPS $M = \arg \min_M \left(-\log \frac{\langle \Psi(M) | \mathcal{T}_{1d} | \Psi(M) \rangle}{\langle \Psi(M) | \Psi(M) \rangle} \right)$

control parameter: MPS bond dimension χ

We can characterize a variational optimum for given values of (D, χ)

→ independent of contraction or optimization algorithm!

Variational optimization

In practice, optimizing both variational principles can be done

→ boundary PEPS: gradient optimization (e.g., BLGS)

$$g = -\frac{2}{\lambda(A, \bar{A})} \left(\begin{array}{c} \text{Diagram 1: } G_l \text{ --- } \begin{array}{c} M_c \\ | \\ E_2 \\ | \\ \bar{M}_c \end{array} \text{ --- } G_r \end{array} \right)^{-1} \left(\begin{array}{c} \text{Diagram 2: } G_l \text{ --- } \begin{array}{c} M_c \\ | \\ e_3 \\ | \\ M_c \end{array} \text{ --- } G_r \\ - \lambda(A, \bar{A}) \begin{array}{c} \text{Diagram 3: } G_l \text{ --- } \begin{array}{c} M_c \\ | \\ e_2 \\ | \\ \bar{M}_c \end{array} \text{ --- } G_r \end{array} \right)$$

“Hellmann-Feynman theorem”: thanks to variational principle for the environment, the differentials of the environment tensors vanish

→ boundary MPS: vumps algorithm finds variational optimum

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PEPS for critical systems

Finite-size scaling in Monte-Carlo or exact diagonalization

Finite-entanglement scaling in MPS

for certain bond dimension, use effective
correlation length as characteristic length scale

Nishino, Okunishi, Kikuchi, Phys. Lett. A 213, 69 (1996)

Tagliacozzo, de Oliveira, Iblisdir, Latorre,
Phys. Rev. B 78, 024410 (2008)

Pollmann, Mukerjee, Turner, Moore,
PRL 102, 25570 (2009)

For PEPS, situation is more complicated

algorithms give different results

- contraction: CTMRG, boundary-MPS, TRG and variations
- PEPS optimization: simple update, full update, variational optimization

two control parameters: two-step extrapolation?

high computational cost to get data points

PEPS for critical systems

Recently, a two-step procedure was proposed for extrapolating data

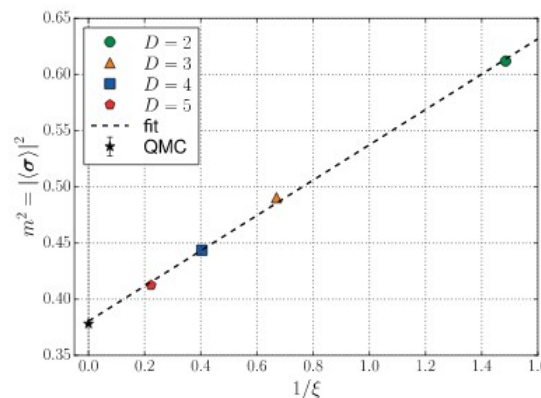
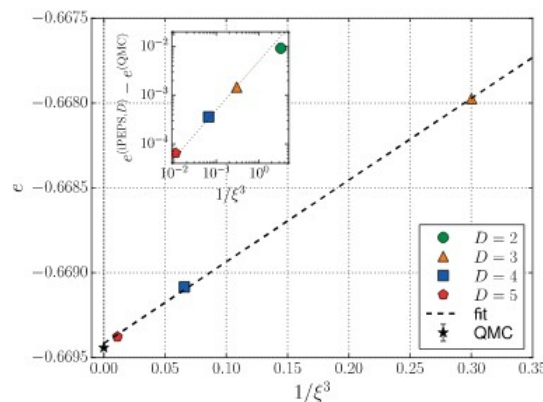
use variational optimization of PEPS tensors,
with χ large enough to get convergence

find PEPS correlation length by extrapolating
environment bond dimension

$$\xi_D = \lim_{\chi \rightarrow \infty} \xi_{\text{boundary-MPS}}$$

Rader, Läuchli, PRX 8, 031030 (2018)

Corboz, Czarnecki, Kapteijns, Tagliacozzo,
PRX 8, 060310 (2018)



(taken from Rader & Läuchli)

→ needs very large χ to get a single data point

PEPS for critical systems

A scaling hypothesis for PEPS

→ place finite-D and finite- χ on an equal footing

PEPS for critical systems

A scaling hypothesis for PEPS

→ place finite-D and finite- χ on an equal footing

Procedure

characterize optimum by variational principles
for boundary PEPS *and* boundary MPS

extract correlation length from boundary MPS,
and formulate scaling hypothesis

$$O(D, \chi) = O(\xi_{\{D, \chi\}})$$

→ single effective length scale that models both finite-D and finite- χ effects

This approach is a lot more efficient

- one can get meaningful data for small χ
- a lot more data points for extrapolations

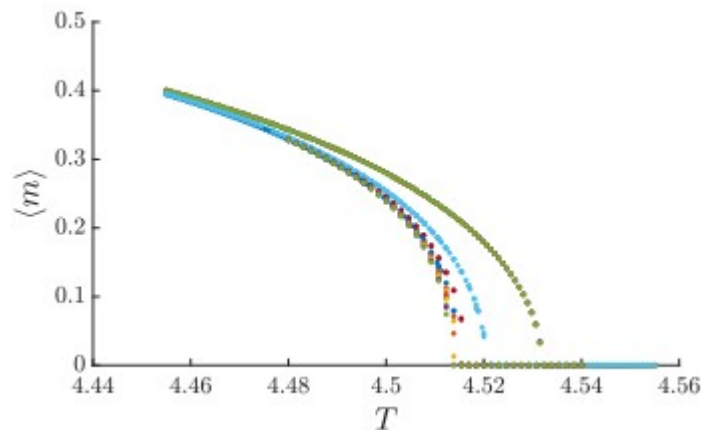
PEPS for critical systems

Application: classical 3-D Ising model

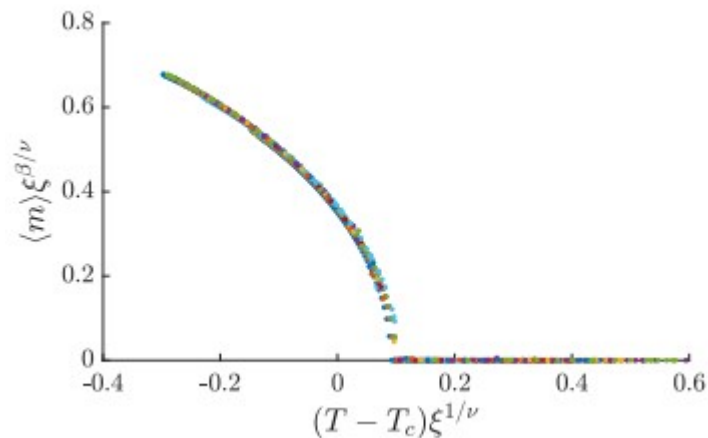
PEPS for critical systems

Application: classical 3-D Ising model

→ generate a lot of data points for different (D, χ)



→ apply scaling hypothesis to get data collapse



optimize data collapse
for critical temperature

$$T_c = 4.51170$$

current world record for TN,
dixit Nishino

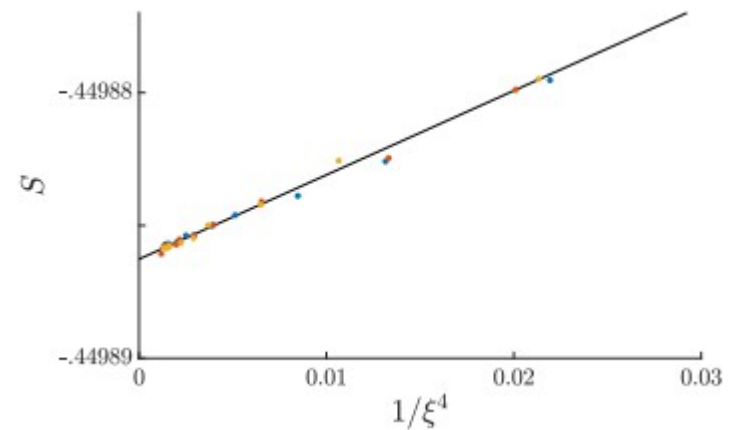
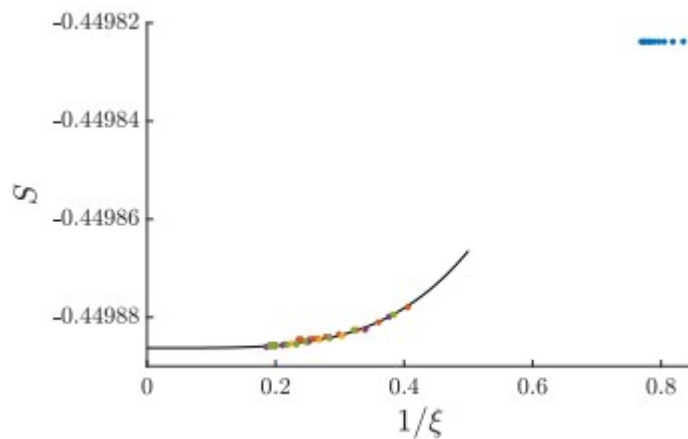
PEPS for critical systems

Application: 3-D dimer covering problem

PEPS for critical systems

Application: 3-D dimer covering problem

→ generate a lot of data points for different (D, χ) and plot entropy vs inverse correlation length



We get unprecedented precision on the dimer entropy

$$S = 0.449886267(25)$$

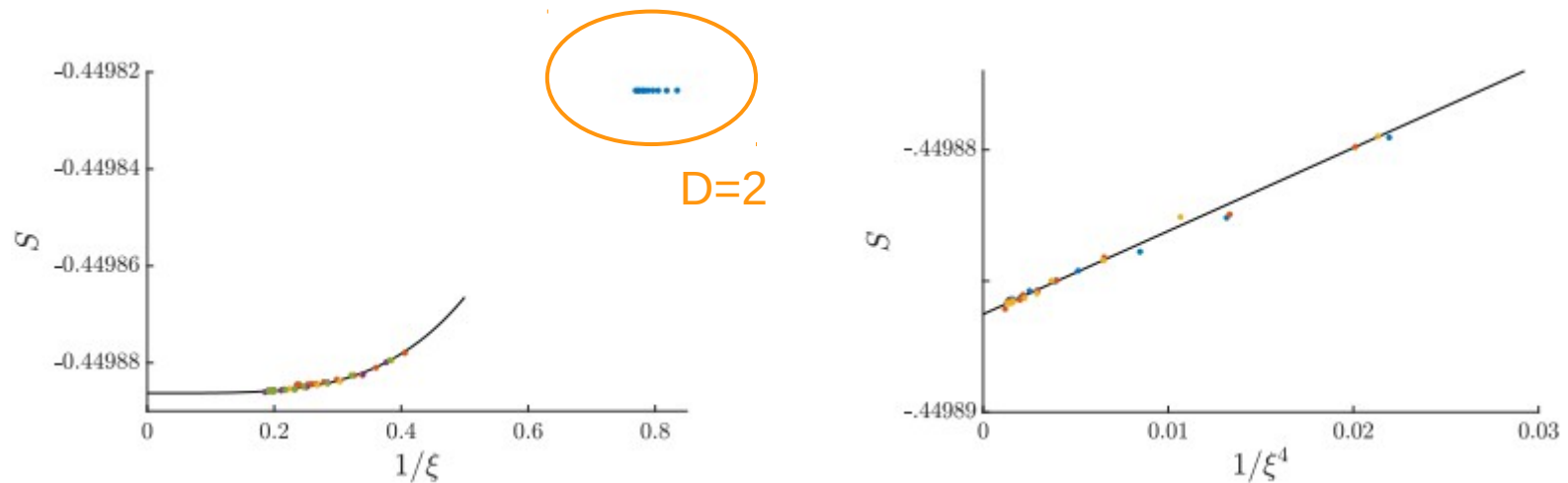
compare to best MC result

$$S = 0.4466(6)$$

PEPS for critical systems

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compare to best MC result

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PEPS for critical systems

We can do the same procedure for 2-D quantum systems

PEPS for critical systems

We can do the same procedure for 2-D quantum systems

optimize PEPS tensor for variational energy,
evaluated with fixed environment bond dimension

use backwards differentiation for the cost function

$$e(A)_{\{D,\chi\}} = \frac{\langle \Psi(A) | \mathcal{H} | \Psi(A) \rangle_{\chi}}{\langle \Psi(A) | \Psi(A) \rangle_{\chi}}$$

Liao, Liu, Wang, Xiang, PRX 9, 031041 (2019)
Hasik, Poilblanc, Becca, SciPost Phys. 10, 012 (2021)

open-source software package for backwards
differentiation of CTMRG contraction

J. Hasik and G. B. Mbeng, “peps-torch” (2020)

PEPS for critical systems

Application: 2-D quantum Heisenberg model

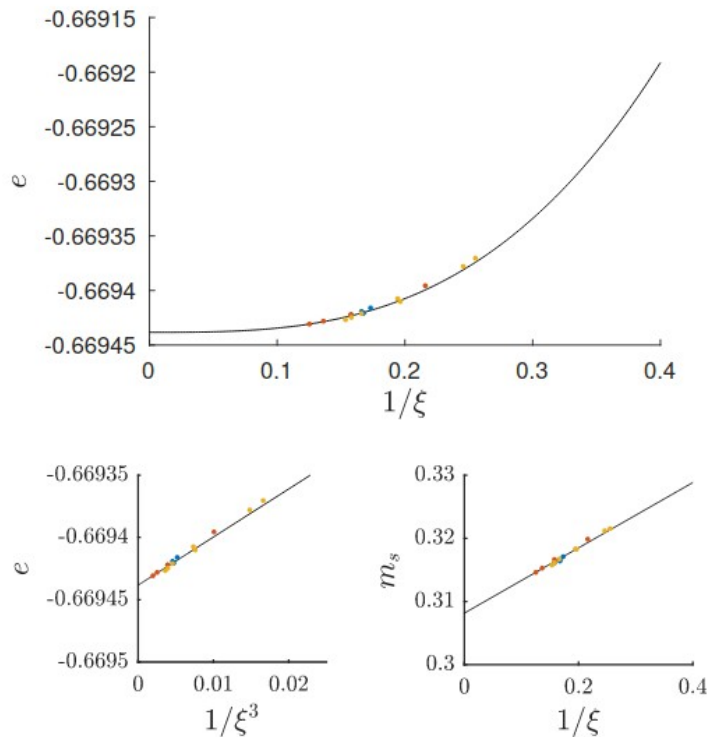


FIG. 3. Data for the 2-D Heisenberg model for $D = (6, 7, 8)$ and selected χ from 17 up to 200 for $D = 6, 7$ and 147 for $D = 8$. Top: energy vs. inverse correlation length of the environment MPS. Bottom-left: energy vs. the cube of the inverse correlation length. Bottom-right: magnetization vs that same inverse correlation length. Fits are performed on all the data and are shown in black.

optimize PEPS for different (D, χ)

energy/magnetisation vs inverse correlation length nicely collapses on a single curve

get extrapolated values

$$e = -0.6694401(10)$$

$$m_s = 0.30771(31)$$

compare to MC values

$$e = -0.669437(5)$$

$$m_s = 0.30743(1)$$

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Accurate PEPS simulations of (3+0)-D or (2+1)-D critical theories are possible

- optimize at finite (D, χ) through variational principles

- extract a single effective length scale from
optimized PEPS and environment

- scaling hypothesis

Outlook

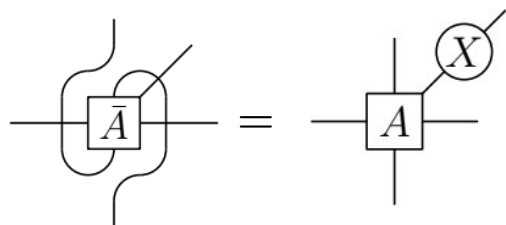
Accurate PEPS simulations of (3+0)-D critical theories are possible

optimize at finite (D, χ) through variational principles

extract a single effective length scale from
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scaling hypothesis

Symmetry constraints on PEPS tensor



what ground states do not permit such a parametrization?

More exotic critical points

What about (3+1)-D systems?



3 control parameters, 3-step hierarchy of variational principles, single effective length scale?

Thank you!