

Hadronic Vacuum Polarization: An unblinded window on the g-2 mystery

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challenges in Lattice field theory
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Introduction

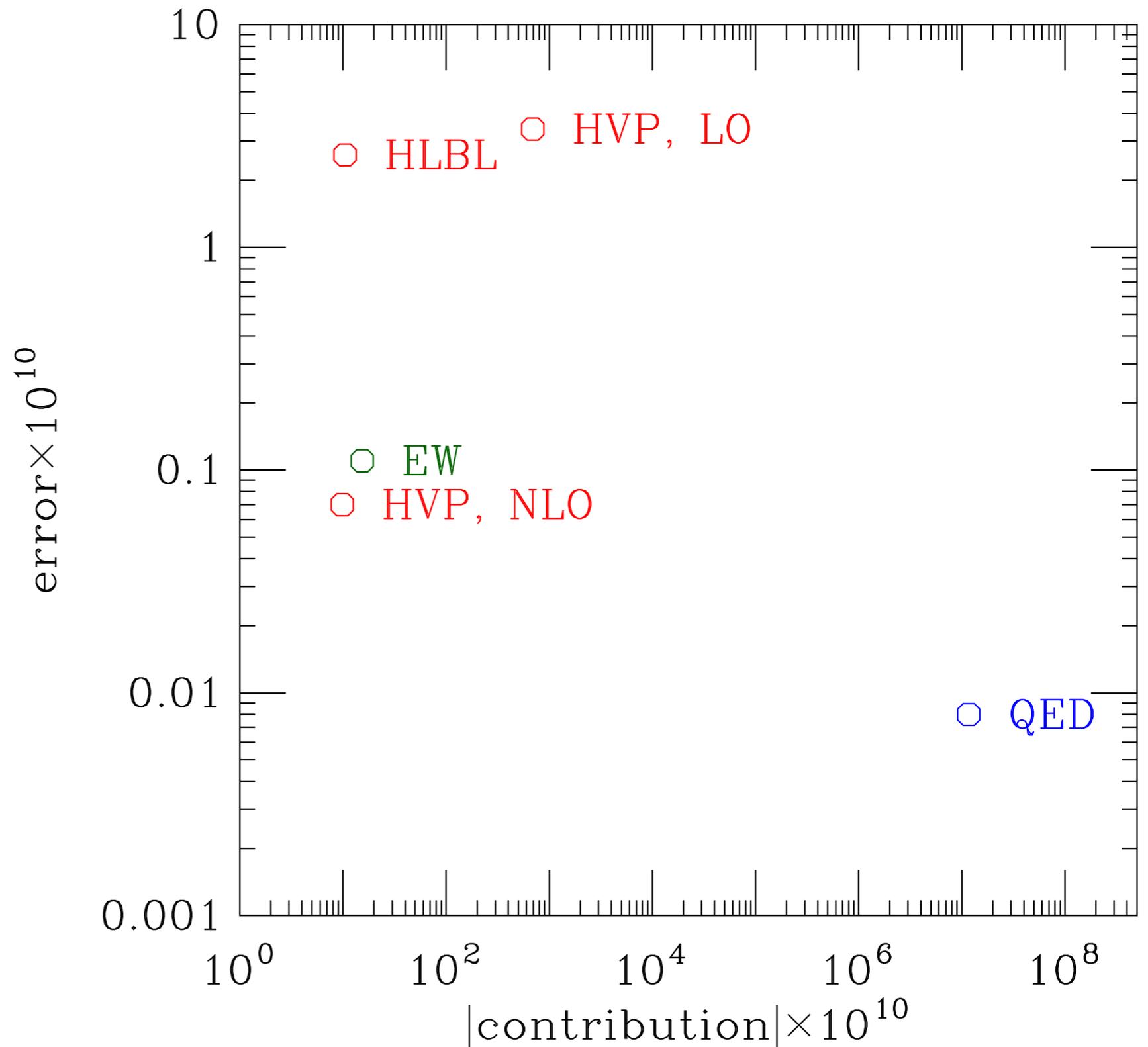
- ◆ Anomalous magnetic moments of electron and muon are two of the most precisely measured quantities in physics
- ◆ E821 at BNL published its final value for the muon in 2006
- ◆ FNAL E989 announced its initial result in April, 2021
 - spectacular agreement with E821
 - continues to run
 - New experiment E34 planned at J-PARC
- ◆ There is $\approx 4.2 \sigma$ difference between data driven standard model (SM) calculation and experiment
- ◆ BMWc 2021 value lies between SM value and experiment
- ◆ It is important to improve the precision of other lattice QCD calculations

Theory Overview

- ◆ SM contributions come from QED (electron & muon), electroweak contributions, and hadronic contributions that involve quarks
 - all forces save gravity contribute
- ◆ Current situation summarized by Muon $g-2$ Theory Initiative
 - T. Aoyama *et al.*, Phys. Rept. 887 (2020), 2006.04822 [hep-ph]
- ◆ Next plot shows how the hadronic corrections dominate the error

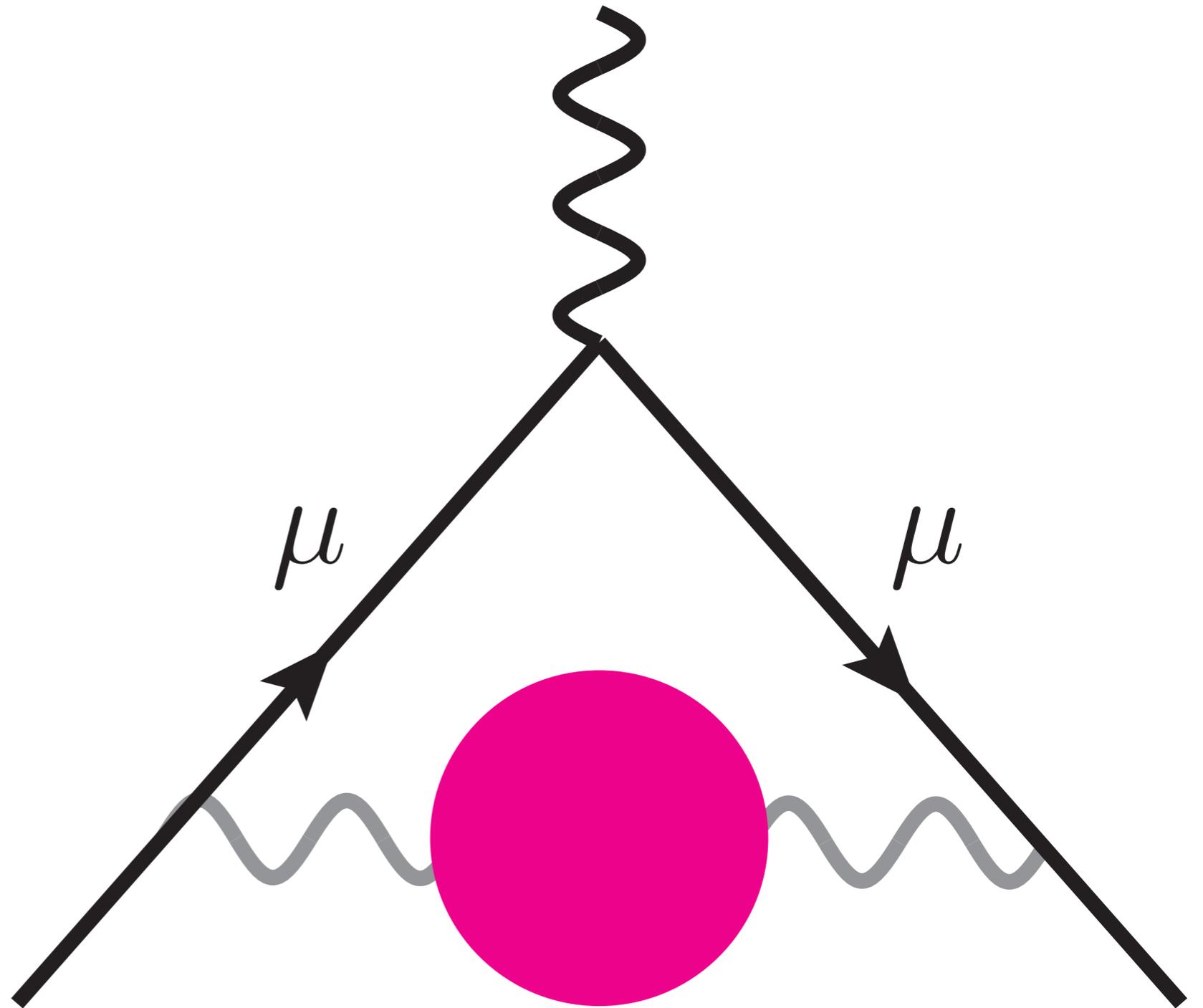
Error vs. Contribution

- QED in blue has very small error
- Electroweak in green has larger error, but small contribution
- Hadronic contributions are all in red
 - LO Hadronic vacuum polarization largest error and 2nd largest contribution
 - HLBL 2nd largest error
- This talk on LO HVP



Hadronic Vacuum Polarization

- HVP diagram looks like 2 loop QED diagram, except that red blob includes all QCD corrections to the quark loop, and there are also ‘disconnected’ contributions with two quark loops that can exchange gluons
- Contribution written as integral over 4-momentum-squared



HVP Calculation

◆ Hadronic part of the current-current two-point function must be integrated over the loop momentum. Rest of diagram uses known photon and muon propagators.

◆ Two approaches:

◆ Using dispersion relations and optical theorem, can convert integral to one involving

$$R = \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma(e^+e^- \rightarrow \mu^+\mu^-)$$

- This approach relies on careful use of experimental data and is currently the most precise method.

◆ *Ab initio* lattice QCD non-perturbatively calculates the current-current two-point function for Euclidean time or spacelike Q^2 .

- Challenge is to get accurate values at large Euclidean time or small Q^2 .

Lowest Order HVP

- ◆ HVP is calculated as sum of several contributions: light quark connected, strange connected, ..., light disconnected, ..., strong isospin breaking, electromagnetic, etc.
- ◆ $\alpha_{\mu}^{ll}(\text{conn.})$ light quark connected is biggest contribution, by far
- ◆ FNAL/HPQCD/MILC: PRD **101**, 034512 (2020), 1902.04223 [hep-lat]
 - briefly recap

Lattice Ensembles

- ◆ In 2020, we used $N_f=2+1+1$ HISQ ensembles from the MILC collaboration with physical light quark masses

$\approx a$ (fm)	$am_l^{\text{sea}}/am_s^{\text{sea}}/am_c^{\text{sea}}$	w_0/a	M_{π_5} (MeV)	$(L/a)^3 \times (T/a)$	$N_{\text{conf.}}$
0.15	0.00235/0.0647/0.831	1.13670(50)	133.04(70)	$32^3 \times 48$	997
0.15	0.002426/0.0673/0.8447	1.13215(35)	134.73(71)	$32^3 \times 48$	9362
0.12	0.00184/0.0507/0.628	1.41490(60)	132.73(70)	$48^3 \times 64$	998
0.09	0.00120/0.0363/0.432	1.95180(70)	128.34(68)	$64^3 \times 96$	1557
0.06	0.0008/0.022/0.260	3.0170(23)	134.95(72)	$96^3 \times 192$	1230

- ◆ Have retuned 0.12 fm and added statistics for current analysis. Still adding configurations for 0.06 fm.

$\approx a$ [fm]	$N_s^3 \times N_t$	$am_l^{\text{sea}}/am_s^{\text{sea}}/am_c^{\text{sea}}$	w_0/a	M_{π_5} (MeV)	$N_{\text{conf.}}$	N_{wall}
0.15	$32^3 \times 48$	0.002426 / 0.0673 / 0.8447	1.13215(35)	134.73(71)	9362	48
0.12	$48^3 \times 64$	0.001907 / 0.05252 / 0.6382	1.41110(59)	134.86(71)	9011	64
0.09	$64^3 \times 96$	0.00120 / 0.0363 / 0.432	1.95180(70)	128.34(68)	5384	48
0.06	$96^3 \times 128$	0.0008 / 0.022 / 0.260	3.0170(23)	134.95(72)	2621	24

Blinding

- ◆ To avoid confirmation bias in analysis, correlators are all blinded by multiplication by an unknown factor.
- ◆ Once all aspects of analysis are completed, the collaboration will decide to unblind and actual result will be available.
- ◆ Collaboration looked at many variations on the analysis and decided on Tuesday, September 13 that we were ready to unblind.
- ◆ This is the first presentation of our unblinded results for the Euclidean time windows W and W_2

Windows Analysis

- ◆ The statistical noise at large Euclidean time is challenging
 - RBC/UKQCD suggested using windows to achieve higher precision and allow better comparison of different calculations
 - PRL 121, 022003 (2018)
 - FNAL/HPQCD/MILC recently advocated one-sided windows with longer time extent than SD defined in PRL above.
 - 2207.04765 [hep-lat] (use such windows as part of this study)
- ◆ We have considered multiple windows and concentrate on just two here

$$\Theta(t, t_0, t_1, \Delta) = \frac{1}{2} \left[\tanh\left(\frac{t - t_0}{\Delta}\right) - \tanh\left(\frac{t - t_1}{\Delta}\right) \right]$$

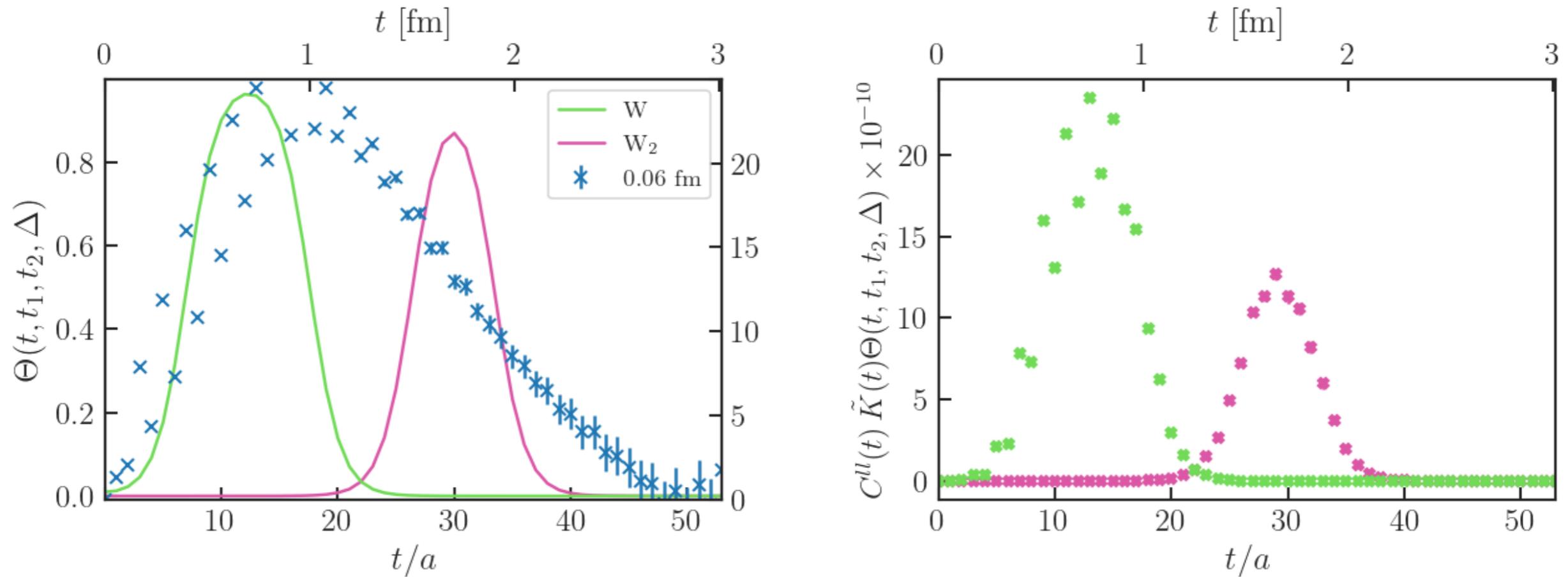
Windows Considered

- ◆ We fix $\Delta = 0.15$ fm.
- ◆ For the one-sided (O.S.), $t_1 = 1, 1.5, 2, 3$.

label	$[t_0, t_1]$ fm
a_μ^{SD}	$[0, 0.4]$
a_μ^{W}	$[0.4, 1]$
$a_\mu^{\text{W}_2}$	$[1.5, 1.9]$
$a_\mu^{\text{O.S.}}(t_1)$	$[0, t_1]$

- ◆ Here, we only present W and W_2 (Aubin *et al.* 2204.12256 [hep-lat])
- ◆ Each window has its own blinding factor, so can unblind independently.

Effect of Window



- ◆ Left: a_μ integrand in blue; W window factor in green; W_2 in red
- ◆ Right: integrand after multiplication by window factor
- ◆ note effect of staggering on W

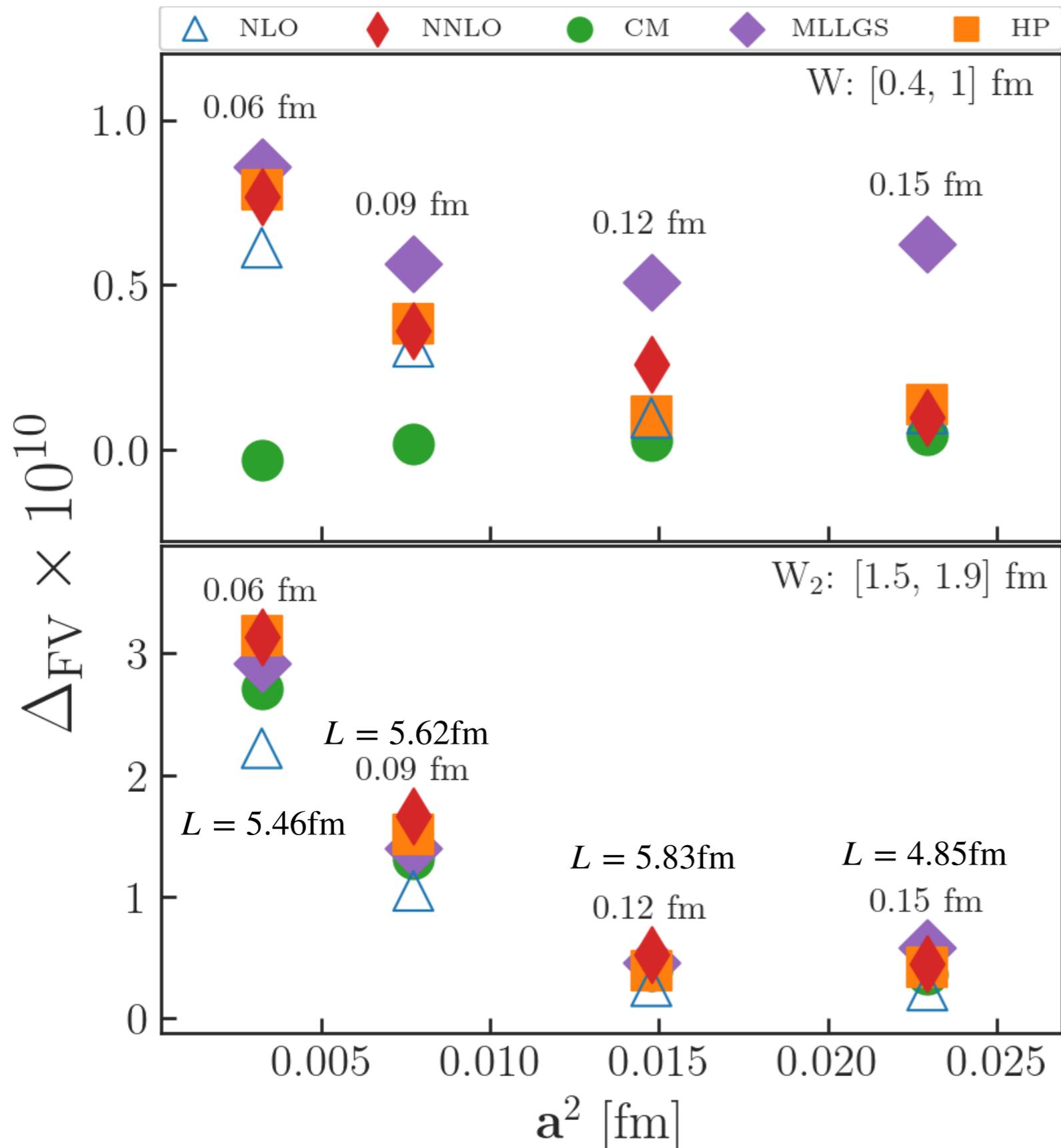
Corrections

- ◆ Three corrections are applied: volume, mass mistuning, and taste breaking. (Latter is optional, see below.)
- ◆ $a_\mu(L_\infty, m_{\pi_{phys}}) = a_\mu(L_{latt}, m_{\pi_{latt,\xi_1}}, \dots, m_{\pi_{latt,\xi_{16}}}) + \Delta_{FV} + \Delta_{m_\pi} + \Delta_{TB}$
- ◆ $\Delta_{FV} = a_\mu(L_\infty, m_{\pi_{latt,\xi_1}}, \dots, m_{\pi_{latt,\xi_{16}}}) - a_\mu(L_{latt}, m_{\pi_{latt,\xi_1}}, \dots, m_{\pi_{latt,\xi_{16}}})$
- ◆ $\Delta_{m_\pi} = a_\mu(L_\infty, m_{\pi_{phys,\xi_1}}, \dots, m_{\pi_{phys,\xi_{16}}}) - a_\mu(L_\infty, m_{\pi_{latt,\xi_1}}, \dots, m_{\pi_{latt,\xi_{16}}})$
- ◆ $\Delta_{TB} = a_\mu(L_\infty, m_{\pi_{phys}}, \dots, m_{\pi_{phys}}) - a_\mu(L_\infty, m_{\pi_{phys,\xi_1}}, \dots, m_{\pi_{phys,\xi_{16}}})$
- ◆ Correction terms calculated on each ensemble using several models

Correction Models

- ◆ We consider several models
 - Chiral Perturbation Theory (ChiPT NLO, NNLO)
 - Meyer-Lellouch-Lüscher-Gournaris-Sakurai (MLLGS)
 - Chiral Model (CM, and CM' variation)
 - Hansen and Patella (HP)
 - last is used only for finite volume correction
- ◆ We also try neglecting Δ_{TB} at each lattice spacing and allowing continuum limit to eliminate taste breaking
- ◆ Don't need to use the same model for all correction terms.
 - many, many variations

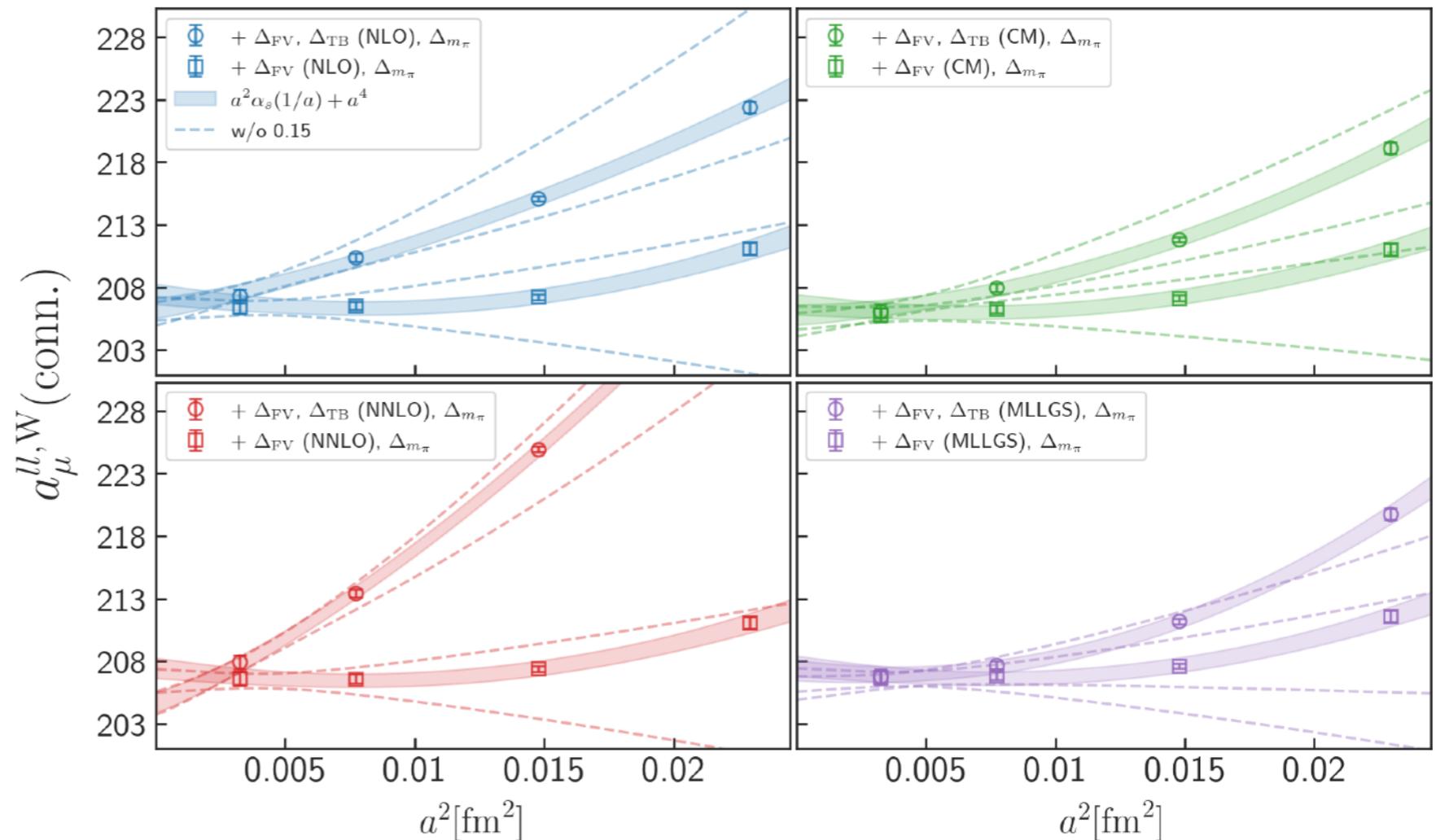
Finite Volume Correction



- FV correction for W (top) and W_2 (bottom) windows, shows much better consistency for the window at larger time advocated by Aubin et al.
- FV correction is so small at smallest volume (coarsest ensemble) because taste breaking is larger there.

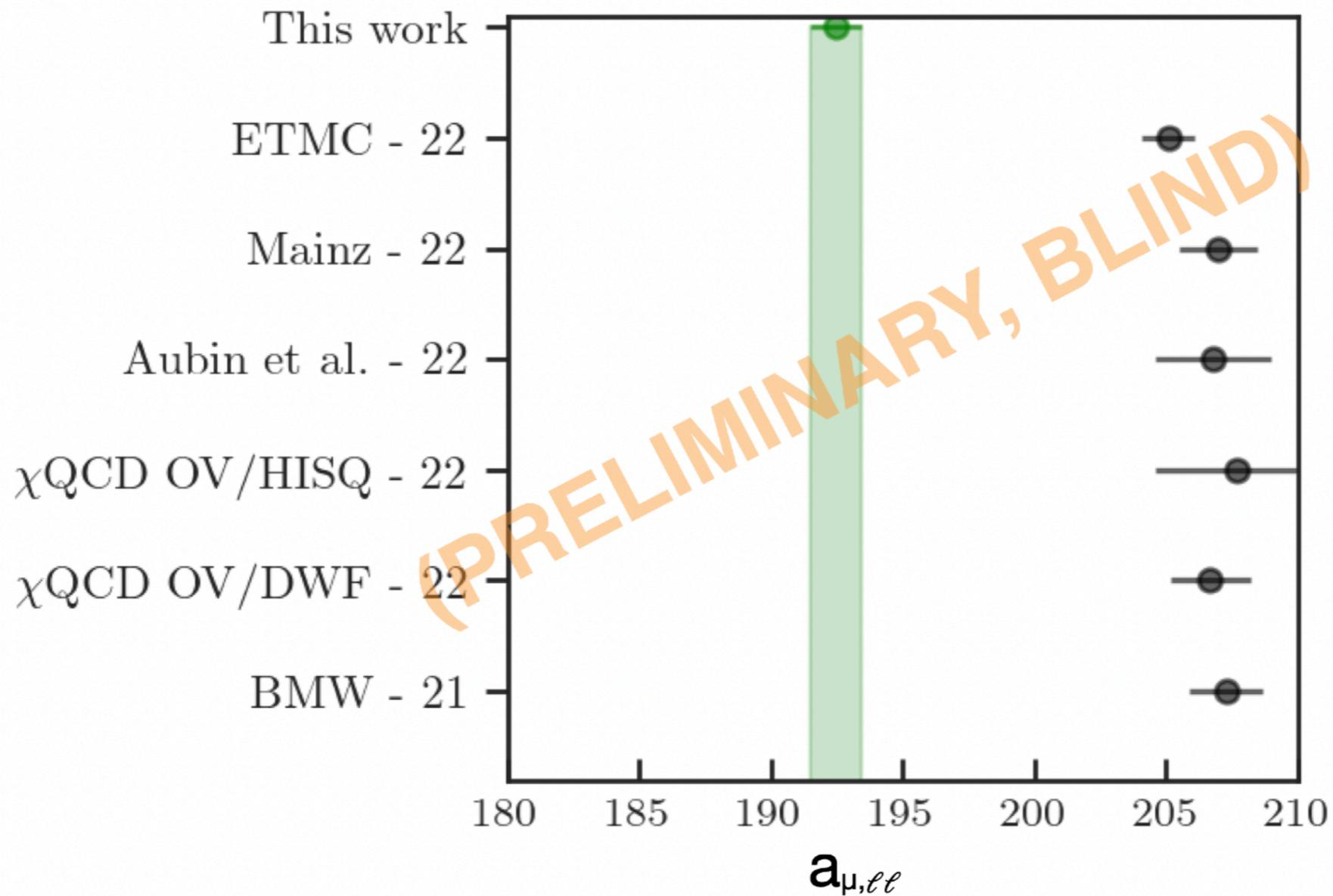
To Correct TB or Not?

- We can allow continuum limit to remove taste breaking or remove on each ensemble.
- We see some differences as $a \rightarrow 0$ depending on model & whether we include coarsest ensemble.
- Blinded result in range 190–195



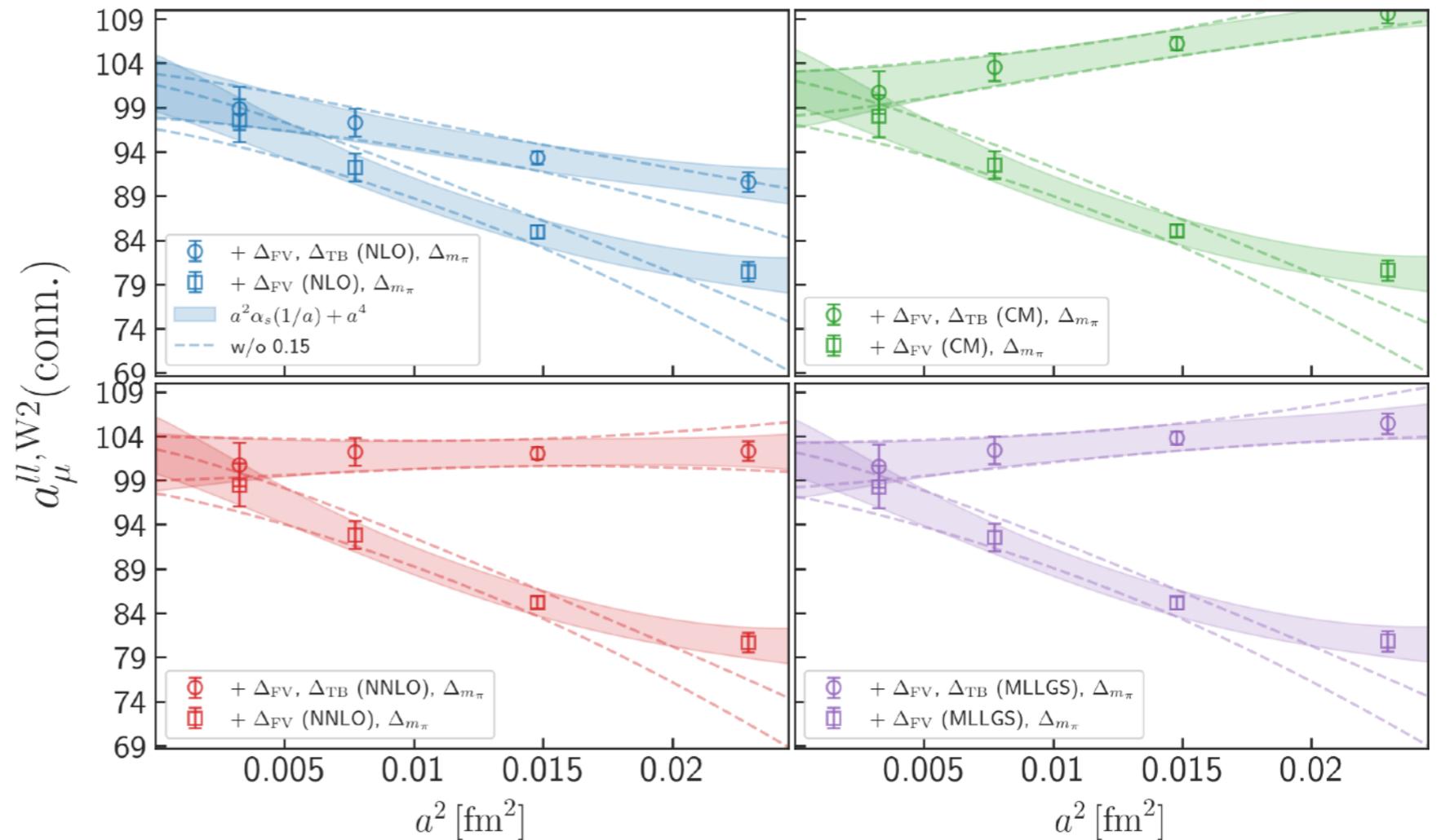
W window

Blinded Result for W



To Correct TB or Not? II

- Lattice spacing dependence is quite different for window at larger time.
- Model corrections can differ quite a bit, but as $a \rightarrow 0$ results are more consistent, than in previous case.
- Error is also larger.



W_2 window

Bayesian Model Averaging

- ◆ Introduced by Jay and Neil, PRD **103**, 114502 (2021).
- ◆ Useful when considering multiple models (or parameter values like t_{\min} in fits).

$$\text{pr}(M | D) \equiv \exp \left[-\frac{1}{2} \left(\chi_{\text{aug}}^2(\mathbf{a}^*) + 2k + 2N_{\text{cut}} \right) \right]$$

gives the weight of each model in the average.

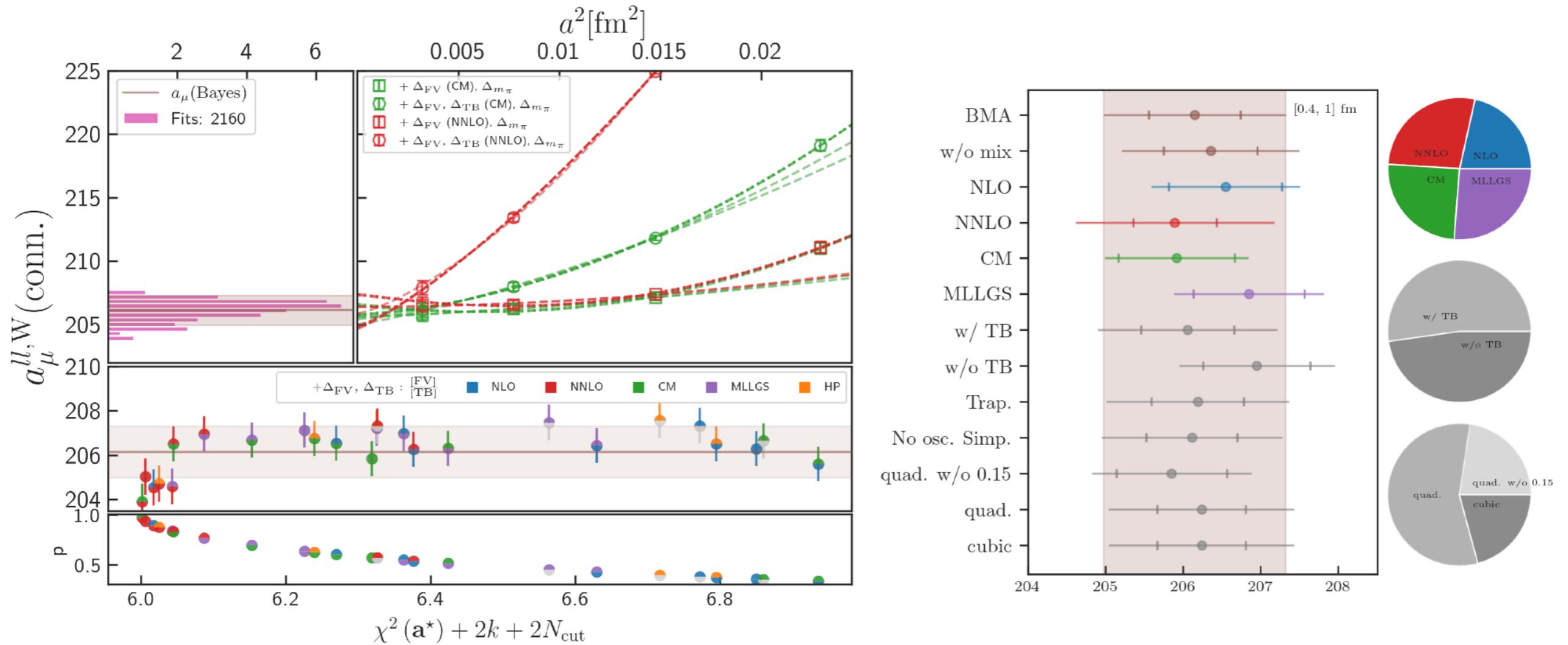
$$\langle a_{\mu} \rangle = \sum_i \langle a_{\mu} \rangle_i \text{pr}(M_i | D)$$

is the average over the models.

Bayesian Model Averaging II

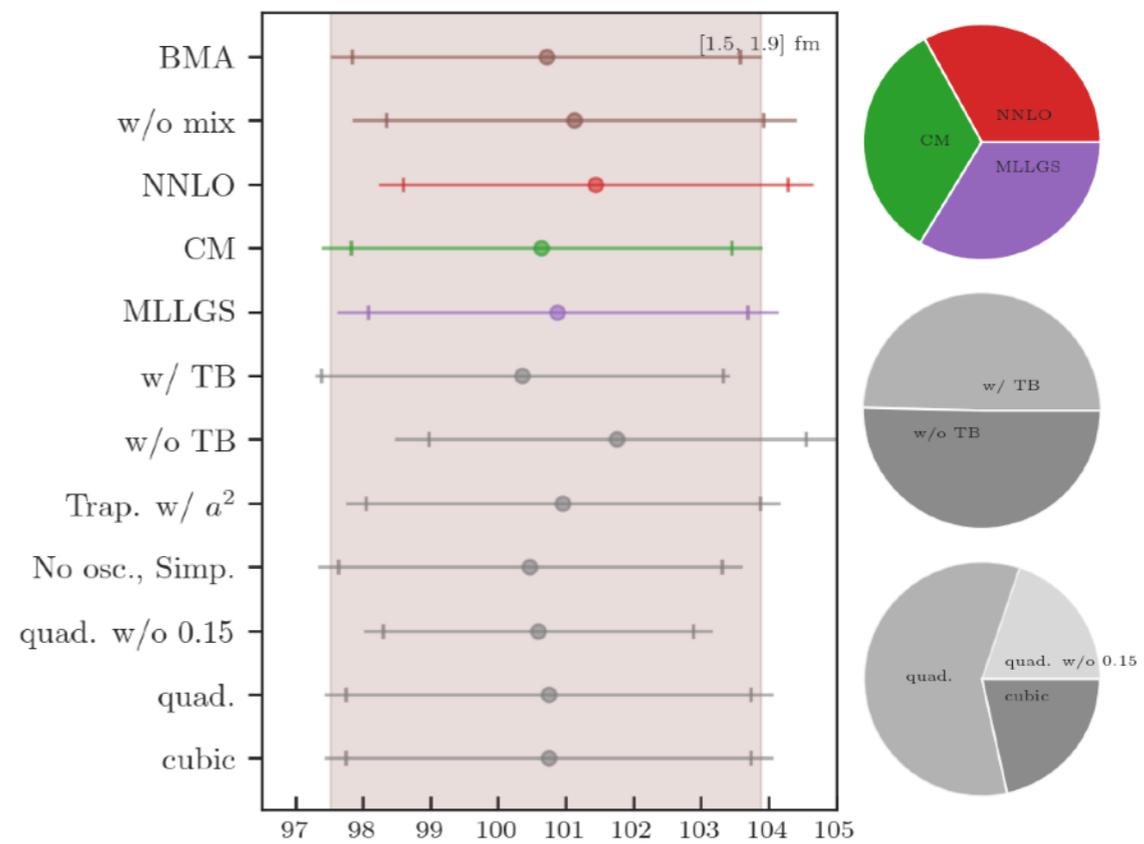
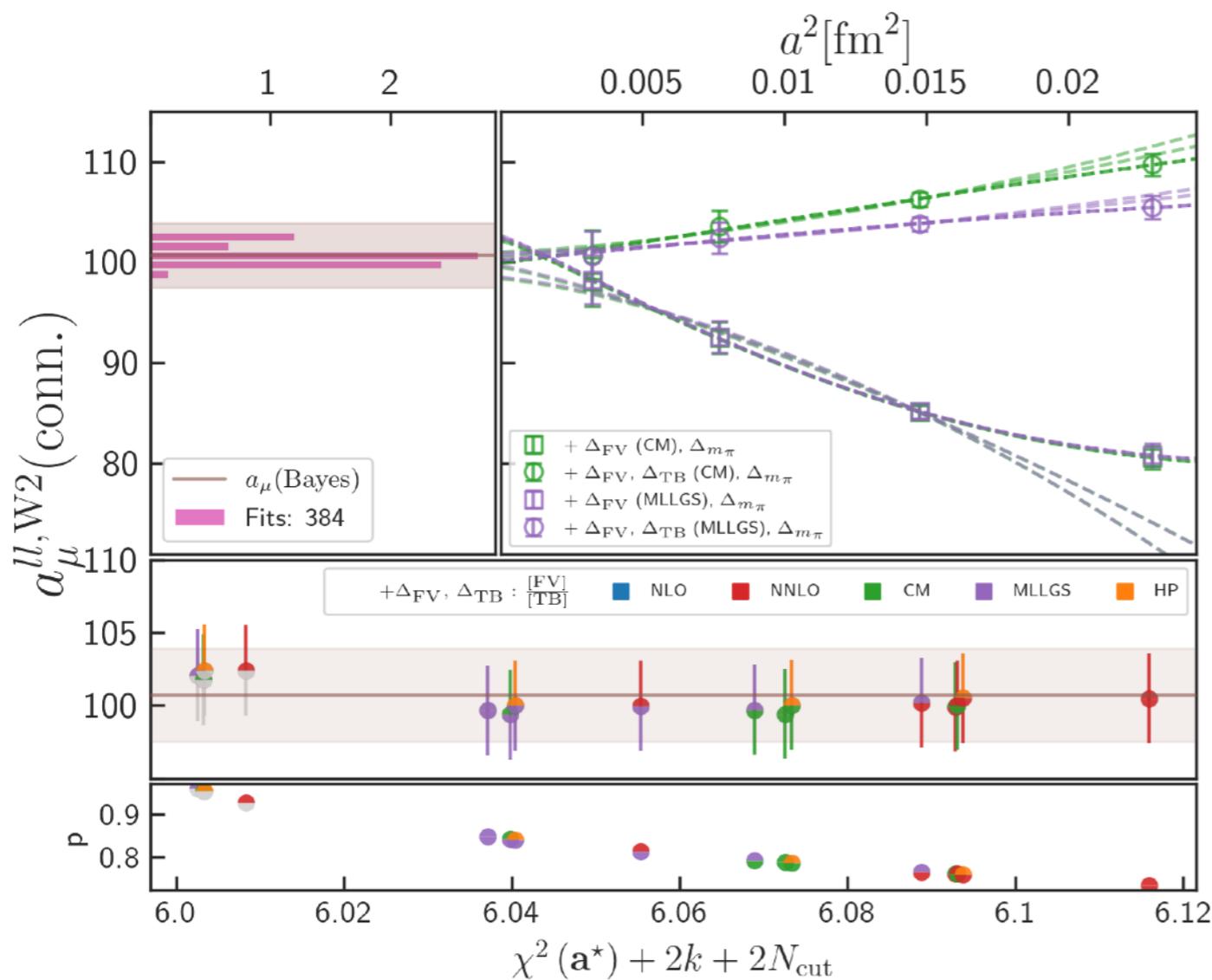
- ◆ Many variations in how the fit is done:
 - choice of model for each correction FV, mistuning, TB
 - also no taste breaking correction
 - apply corrections to a reduced region of time
 - remove opposite parity contributions to vector-correlator that come from using staggered quarks
 - dropping coarsest ensemble
 - variations in the number of powers of a^2 and α_s in continuum fit
 - inclusion of sea-quark mistuning term

BMA for W



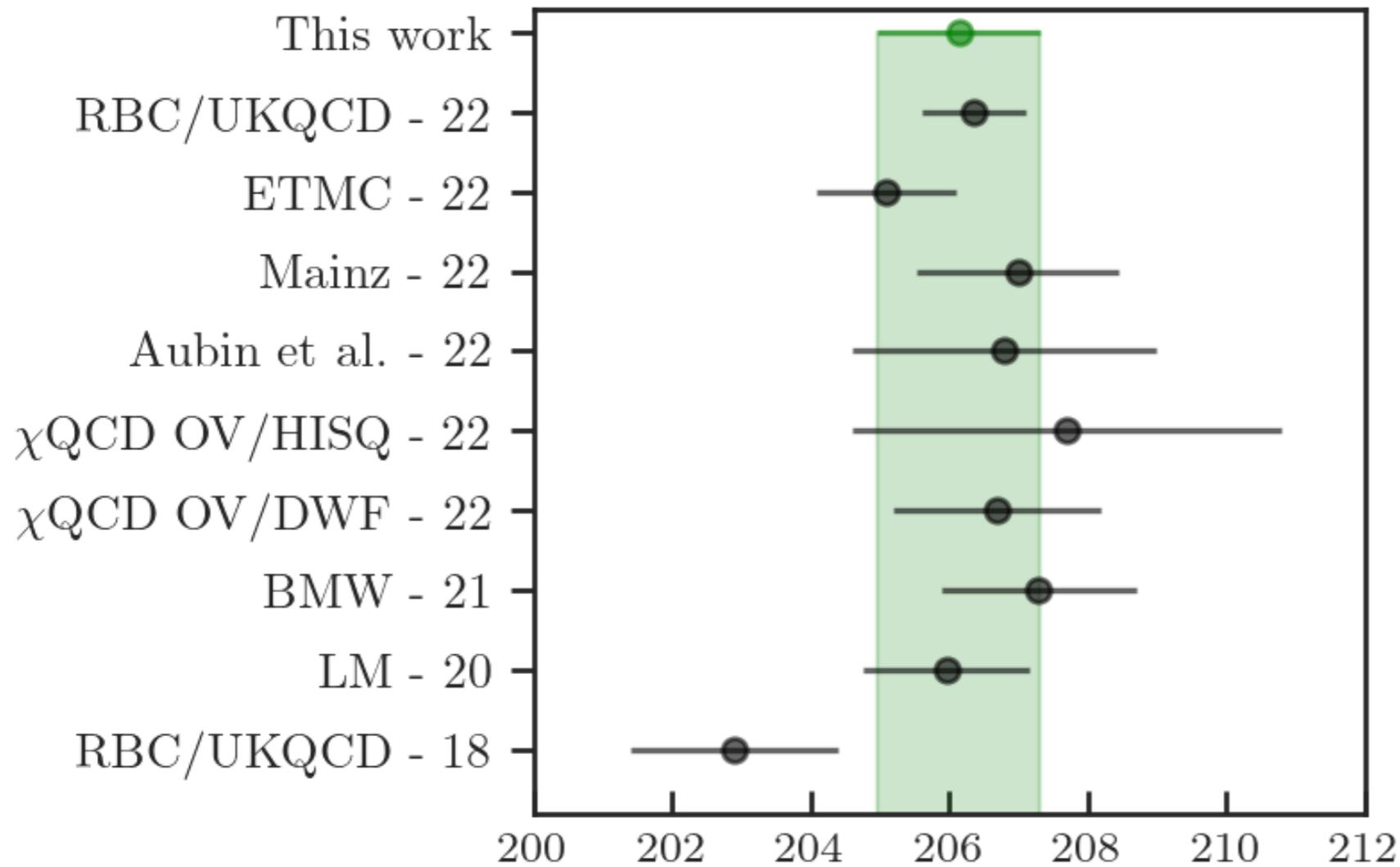
- (L) Four panels show many aspects of the various fits: histogram of 2,160 fits; examples of fits using CM and NLO chiral perturbation theory; 50 best fits; p-value for data contribution to χ^2 .
- (R) Model average using only subsets of the models.

BMA for W_2



- Similar to previous slide but for the window suggested by Aubin et al.

Unblinded Result for $a_{\mu ll}^W(\text{conn.})$



- $206.1(1.2) \times 10^{10}$
- Our result is in excellent agreement with recent results.
- Our error is not quite as small as RBC/UKQCD-22, but comparable to best of the rest.

Error Budget for W

- ◆ Continuum extrapolation is dominant source of error
- ◆ Scale setting, current renormalization, and finite volume are all close in size

Source	$a_\mu^{ll,W}$ (conn.) %
Continuum extrapolation ($a \rightarrow 0$, Δ_{TB})	0.37
Scale setting (w_0 (fm), w_0/a)	0.24
Current renormalization (Z_V)	0.25
Monte Carlo statistics	0.17
Finite-volume correction (Δ_{FV})	0.21
Pion-mass adjustment (Δ_{m_π})	0.09
Total	0.57%

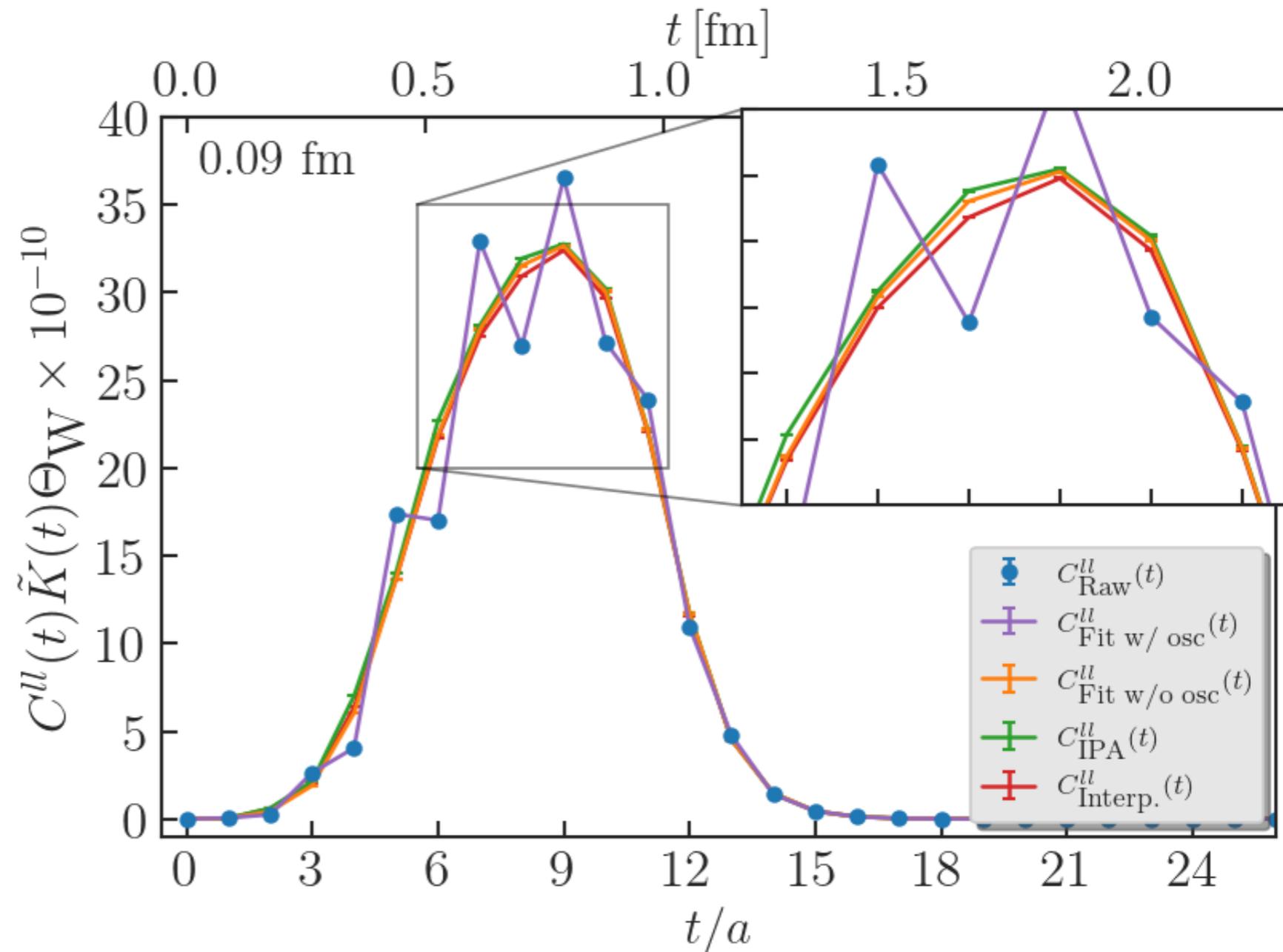
Error Budget for W_2

- ◆ This window is at larger time, so limited statistical precision is dominant source of error.
- ◆ Scale setting, continuum extrapolation, and mass adjustment are also considerable.

Source	$a_{\mu}^{ll, W_2}(\text{conn.})$ %
Monte Carlo statistics	2.42
Scale setting (w_0 (fm), w_0/a)	1.28
Continuum extrapolation ($a \rightarrow 0$, Δ_{TB})	1.03
Pion-mass adjustment ($\Delta_{m_{\pi}}$)	0.92
Finite-volume correction (Δ_{FV})	0.29
Current renormalization (Z_V)	0.22
Total	3.16%

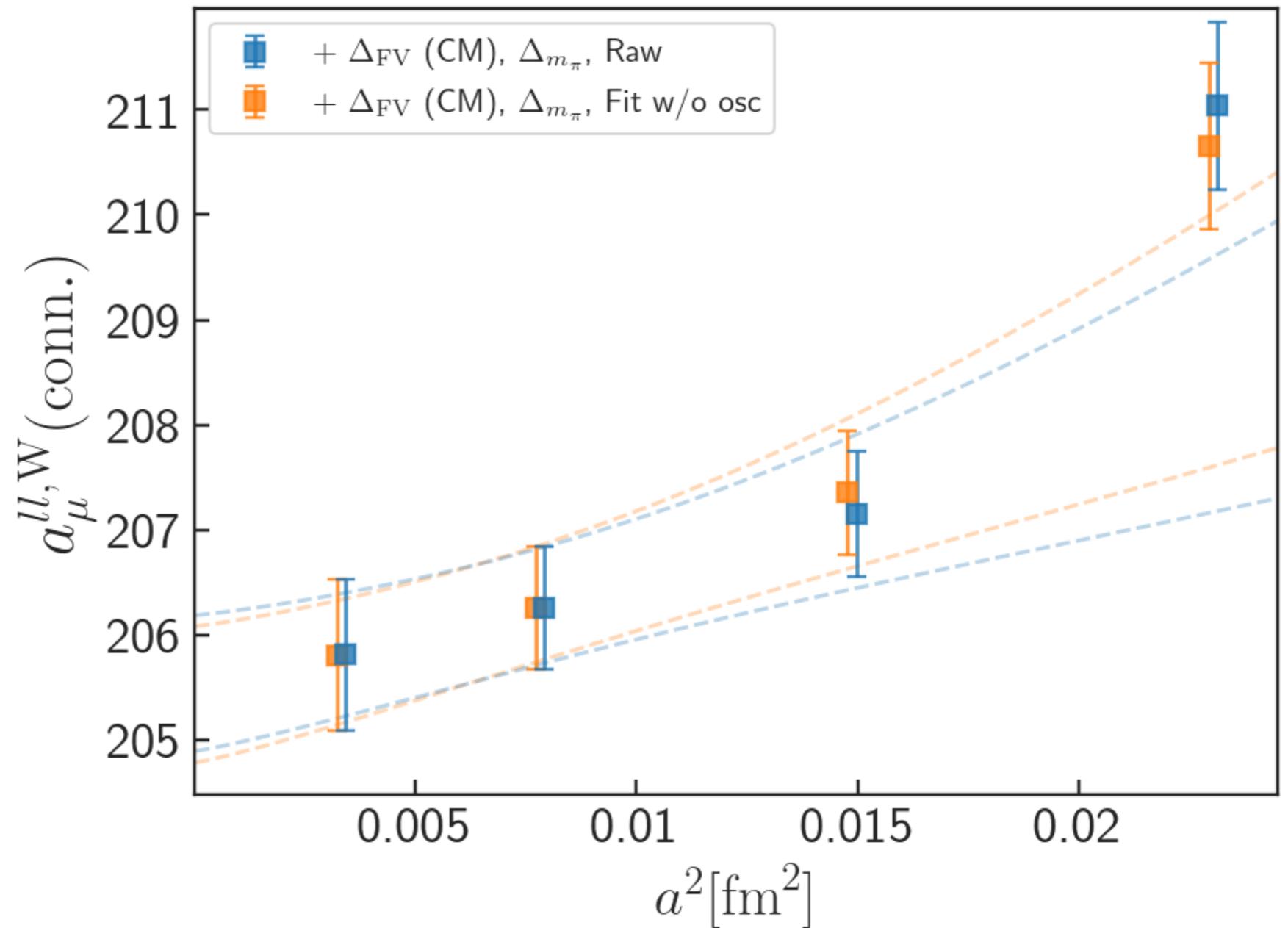
Staggered Opposite Parity

- Well known that staggered hadronic propagators often couple to opposite parity state resulting in $(-1)^t$ terms
- We have explored some ways to eliminate these terms
- Fit and eliminate the opposite parity part



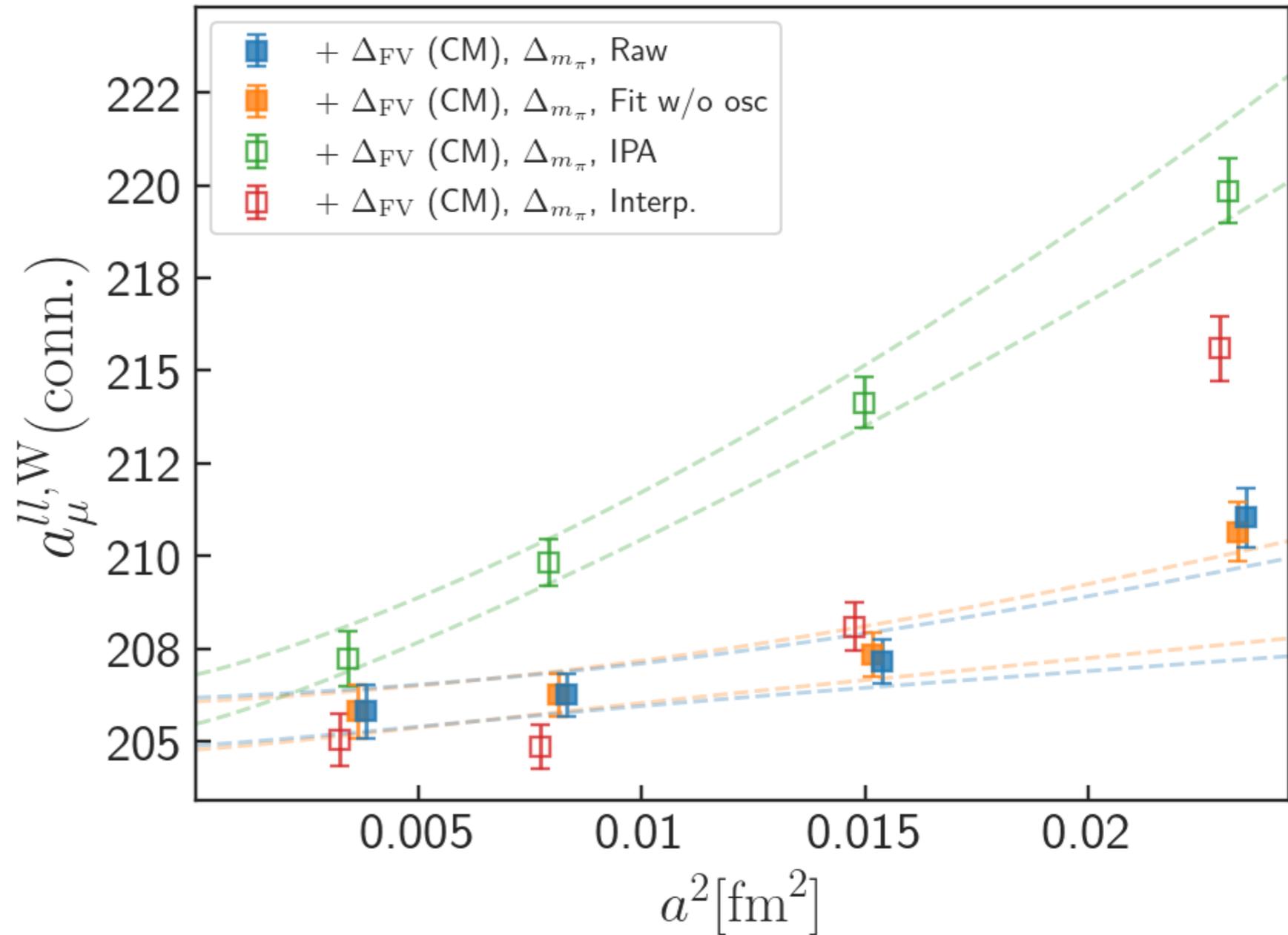
Staggered Opposite Parity II

- Removing the opposite parity contribution explicitly makes little difference
- Here we use a chiral model for the finite volume correction.



Staggered Opposite Parity III

- Improved parity averaging (IPA) and interpolation of even only and odd only times are not useful approaches



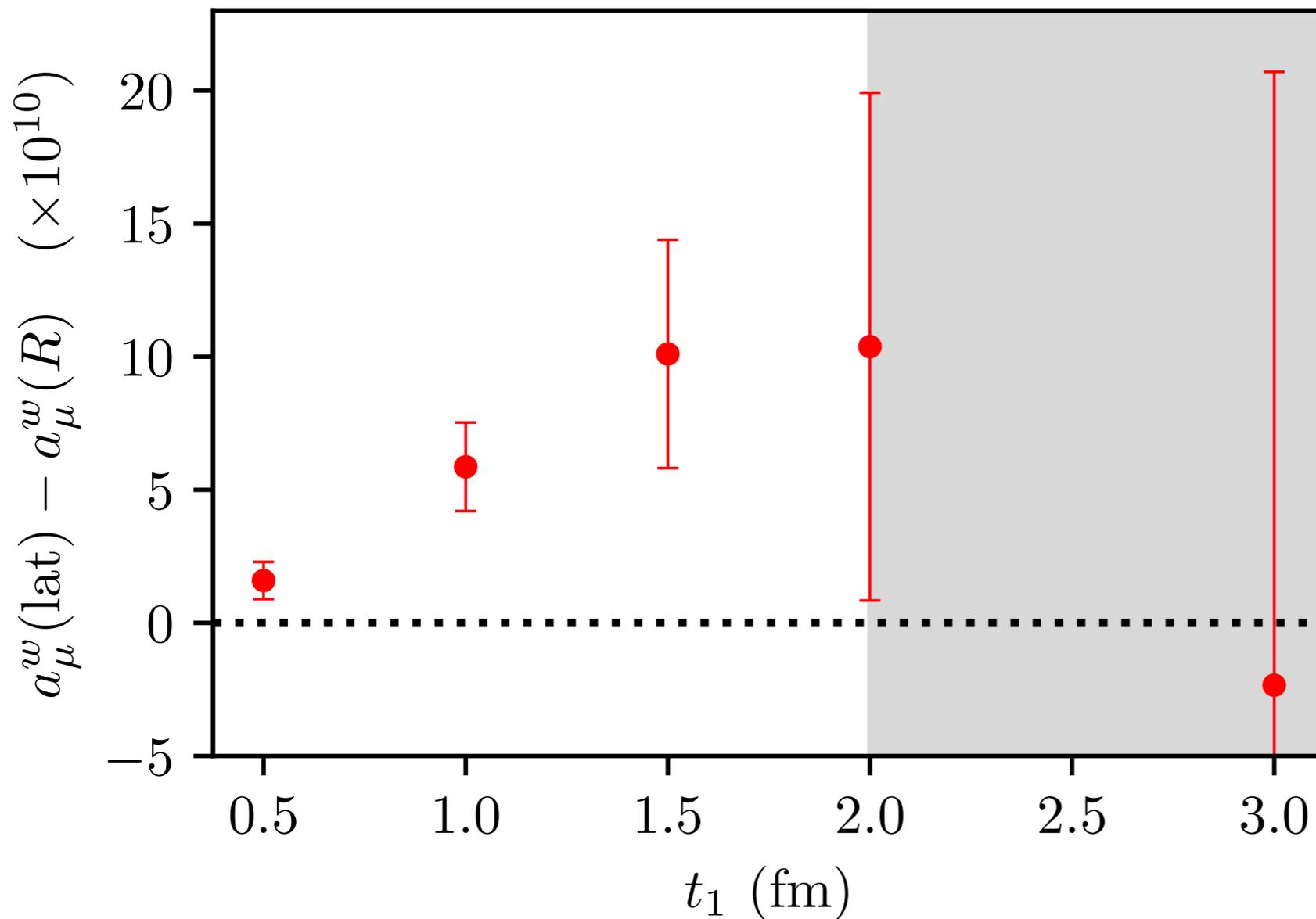
Towards a Complete Calculation

- ◆ Ultimate goal is a_μ , so we need:
 - better scale setting
 - extending range of ensembles with gauge flow data
 - Ω baryon mass (Yin Lin)
 - better statistical accuracy at large time
 - Michael Lynch's Lattice '22 poster on low-mode improvement
 - Shaun Lahert's work on two pion states (not presented here)
 - now analyzing 0.12 fm ensemble
 - strong isospin breaking
 - Curtis Peterson's analysis (not presented here)
 - electromagnetic corrections
 - Gaurav Ray's work was presented at Lattice 2022

Conclusions

- ◆ Contributions to a_μ from various windows in Euclidean time provide valuable benchmarks for lattice QCD calculations on the way to complete HVP calculation
- ◆ This is our first announced window result
 - We expect a paper on arXiv in a couple of weeks
- ◆ The lattice community needs to continue to work hard on the full set of hadronic contributions to a_μ
 - The tension between the data driven (dispersive) approach and lattice QCD is of critical importance and must be resolved or explained

One sided windows



- Difference between lattice and R-ratio determination for various one-sided windows.
- From 2207.04765, using data from 2020.
- We have analyzed several windows with our updated data set