

# THIRD-ORDER EIGENVALUE PROBLEMS

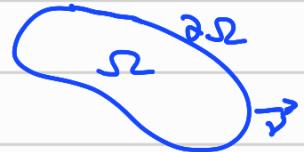
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## BULK-BOUNDARY EIGENVALUE PROBLEMS

$$E[u] = \frac{\int_{\Omega} |\Delta u|^2 dx}{\int_{\Omega} u^2 + \gamma \int_{\partial\Omega} u^2}$$

$0 \leq \gamma \leq \infty$



$$\mathcal{H} = \{u \in H^2 / \partial_\nu u = 0\}$$



$$\left\{ \begin{array}{l} Lu := \Delta^2 u = \lambda u \quad \text{in } \Omega \\ Nu := \partial_\nu u = 0 \quad \text{on } \partial\Omega \\ Bu := -\partial_\nu \Delta u = \gamma \lambda u \quad \text{on } \partial\Omega \end{array} \right.$$

### Historical remarks:

$$\begin{cases} \Delta^2 u = 0 & \text{in } \Omega \\ \partial_\nu u = 0 & \text{on } \partial\Omega \\ -\partial_\nu \Delta u = \gamma u & \text{on } \partial\Omega \end{cases}$$

Kuttler-Sigillito 60's.  
(Biharmonic Steklov eigenvalue problems)

$\gamma = \infty$

Linearization of the Cahn-Hilliard equation  
with dynamic boundary conditions.

[Knopf-Signori, Knopf-Lam-Liu-Metzger, ...]

## FUNCTIONAL ANALYSIS:

$$\mathcal{H} = \{u \in H^2(\Omega) / \partial_\nu u = 0\}$$

$$\begin{cases} \Delta^2 u = f \text{ in } \Omega \\ \partial_\nu u = 0 \text{ on } \partial\Omega \\ -\partial_\nu \Delta u = g \text{ on } \partial\Omega \end{cases}$$

$$L_\gamma^2 : \|u\|^2 = \int_{\Omega} u^2 + \gamma \int_{\partial\Omega} u^2$$

$$\text{Fredholm condition: } \int_{\Omega} f + \gamma \int_{\partial\Omega} g = 0.$$

- Self-adjoint (Agmon-Douglis-Nirenberg)
- Good elliptic estimates
- $\exists \quad 0 = \lambda_1 \leq \lambda_2 \leq \dots \xrightarrow{\infty} \text{eigenvalues}$

## CONVERGENCE TO THE LIMITING CASES

Theorem:

$$\lambda_k(\gamma) \rightarrow \lambda_k(0) \text{ as } \gamma \rightarrow 0$$

$$\gamma \cdot \lambda_k(\gamma) \rightarrow \lambda_k(\infty) \text{ as } \gamma \rightarrow \infty.$$

Eigenfunctions converge (up to rescaling).

Proof:

- monotonicity in  $\gamma$
- Min-max principle

# EIGENVALUE PROBLEM IN A BALL:

**Theorem 4.1.** Fix  $\gamma \geq 0$ . The positive eigenvalues of

$$(4.3) \quad \begin{cases} \Delta^2 u = \lambda^4 u & \text{in } \mathbb{B}_1, \\ \partial_\nu u = 0 & \text{on } \partial\mathbb{B}_1, \\ \partial_\nu(\Delta u) = -\gamma\lambda^4 u & \text{on } \partial\mathbb{B}_1, \end{cases}$$

are all and only the solutions of the equation

$$(4.4) \quad 2j'_\ell(\lambda_\ell)i'_\ell(\lambda_\ell) + \gamma\lambda_\ell(j'_\ell(\lambda_\ell)i_\ell(\lambda_\ell) - i'_\ell(\lambda_\ell)j_\ell(\lambda_\ell)) = 0,$$

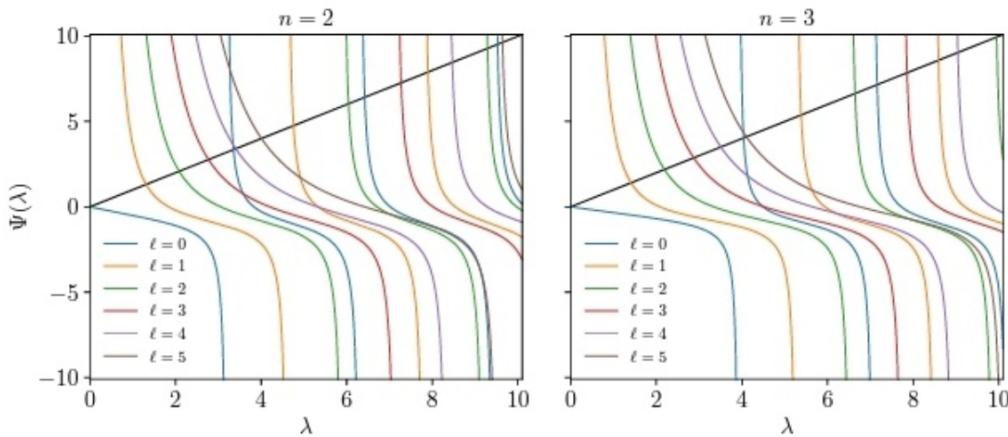
for some  $\ell \in \mathbb{N}$ . For any such solution, the associated eigenfunctions take the form

$$u_\ell(r)\mathcal{Y}_\ell(\theta)$$

for some spherical harmonic  $\mathcal{Y}_\ell$  of order  $\ell$ , where

$$(4.5) \quad u_\ell(r) = i'_\ell(\lambda_\ell)j_\ell(\lambda_\ell r) - j'_\ell(\lambda_\ell)i_\ell(\lambda_\ell r).$$

## The fundamental mode in the ball:



Theorem: The smallest non-zero eigenvalue corresponds to the mode  $\ell = 1$ .

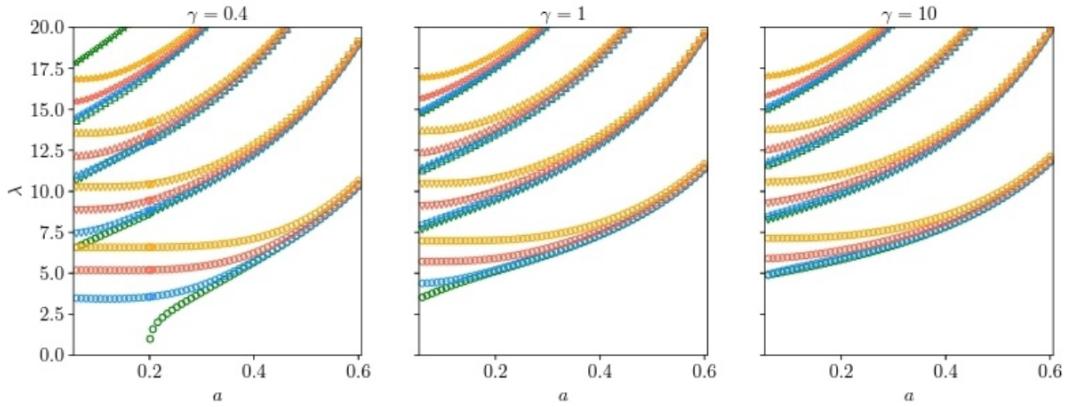
Proof: Inspired by [Chasman] on bilaplacian eigenvalues.

# The eigenvalue problem in the annulus: $\mathcal{A} = \{a < r < 1\}$

$$\begin{pmatrix} j_\ell'(\lambda a) & y_\ell'(\lambda a) & i_\ell'(\lambda a) & k_\ell'(\lambda a) \\ j_\ell'(\lambda) & y_\ell'(\lambda) & i_\ell'(\lambda) & k_\ell'(\lambda) \\ j_\ell'(\lambda a) - \gamma \lambda j_\ell(\lambda a) & y_\ell'(\lambda a) - \gamma \lambda y_\ell(\lambda a) & -i_\ell'(\lambda a) - \gamma \lambda i_\ell(\lambda a) & -k_\ell'(\lambda a) - \gamma \lambda k_\ell(\lambda a) \\ j_\ell'(\lambda) + \gamma \lambda j_\ell(\lambda) & y_\ell'(\lambda) + \gamma \lambda y_\ell(\lambda) & -i_\ell'(\lambda) + \gamma \lambda i_\ell(\lambda) & -k_\ell'(\lambda) + \gamma \lambda k_\ell(\lambda) \end{pmatrix}$$

$\circ$	$\ell = 0, k = 1$	$\circ$	$\ell = 1, k = 1$	$\circ$	$\ell = 2, k = 1$	$\circ$	$\ell = 3, k = 1$
$\nabla$	$\ell = 0, k = 2$	$\nabla$	$\ell = 1, k = 2$	$\nabla$	$\ell = 2, k = 2$	$\nabla$	$\ell = 3, k = 2$
$\Delta$	$\ell = 0, k = 3$	$\Delta$	$\ell = 1, k = 3$	$\Delta$	$\ell = 2, k = 3$	$\Delta$	$\ell = 3, k = 3$
$\star$	$\ell = 0, k = 4$	$\star$	$\ell = 1, k = 4$	$\star$	$\ell = 2, k = 4$	$\star$	$\ell = 3, k = 4$

$n = 2$



**Proposition 5.3.** Assume that  $n \geq 2$  and  $\gamma \in (0, \frac{1}{n})$ . There exists  $a_* \in (0, 1)$  such that, or  $a_* < a < 1$ , the eigenfunction corresponding to the lowest non-zero eigenvalue is radially symmetric, while if  $0 < a < a_*$ , the eigenfunction is not radially symmetric but corresponds to the mode  $\ell = 1$ . For  $n = 2$ ,  $a_* = 1 - 2\gamma$ .

## EIGENVALUE PROBLEMS IN CONFORMAL GEOMETRY

Flat case:  $\Omega \subset \mathbb{R}^n$

$$Lu = \Delta^2 u$$

$\Omega$

$$Nu = \partial_\nu u$$

$\partial\Omega$

$$Bu = -\partial_\nu \Delta u$$

Conformal version:  $(M^n, g)$

$$Lu = \Delta_g^2 u + \operatorname{div}_g(4A_g - 2J_g)u$$

$$Nu = \partial_\nu u$$

$$Bu = -\partial_\nu \Delta u + \dots$$

$$B_g^3 u = -\eta \Delta u - 2\tilde{\Delta} \eta u + 2\langle H_0, \tilde{D}^2 u \rangle - \frac{2}{3} H \tilde{\Delta} u + \frac{2}{3} \langle \tilde{\nabla} H, \tilde{\nabla} u \rangle + \left( -\frac{1}{3} H^2 - 2A(\eta, \eta) + 2\tilde{J} + \frac{1}{2} |H_0|^2 \right) \eta u,$$

(Chang-Qing Case, ...).

## GEOMETRIC INTEREST

- Conformal invariance:

$$\tilde{L} = e^{-4f} L$$

$$\tilde{B} = e^{-3f} B$$

$$\tilde{g} = e^{2f} g$$

- Chern-Gauss-Bonnet formula in dim = 4:

$$8\pi^2 \chi(M) = \int_M \left( \frac{|W|^2_g}{4} + Q_g \right) dv_g + \int_{\Sigma} \left( T_g - \frac{2}{3} \operatorname{tr} H_0^3 \right) d\sigma_h$$

conformal invariant

L

B

umbilic boundary

- Generalize 2-dim geometry to 4-dim

## CLASSIFICATION RESULTS:

$M^4$  lcf, umbilic boundary

- M simply connected,  $\partial M$  one component  $\Rightarrow M \cong$  half-sphere

- $\chi(M) = 1, \gamma(M) > 0 \Rightarrow M \cong$  half-sphere [Raoulot]

- Theorem: M simply connected,  $\partial M$  two boundary components,  $R > 0, Q > 0 \Rightarrow M \cong$  annulus  $\{r \leq r \leq 1\}$

conformal modulus

## BACK TO THE EIGENVALUE PROBLEM:

$$\begin{cases} Lu=0 & \text{in } M \\ Nu=0 & \text{on } \partial M \\ Bu=\lambda u & \text{on } \partial M \end{cases}$$

**Theorem 1.4.** We get the following bounds for  $\lambda_1^g > 0$ :

- a. If  $M$  is conformally equivalent to a half-sphere  $\mathbb{S}_+^4$ ,

$$\lambda_1^g \text{Vol}(\Sigma) \leq 24\pi^2.$$

and it is attained at a flat disk.

- b. If  $M$  is conformally equivalent to an annulus  $\mathcal{A}_\rho$  (with boundaries  $\Sigma_1, \Sigma_\rho$ ),

$$(1.24) \quad \lambda_1^g \text{Vol}(\Sigma) \leq c\left(\rho, \frac{\text{Vol}(\Sigma_\rho)}{\text{Vol}(\Sigma_1)}\right),$$

where this is a constant can be explicitly computed.

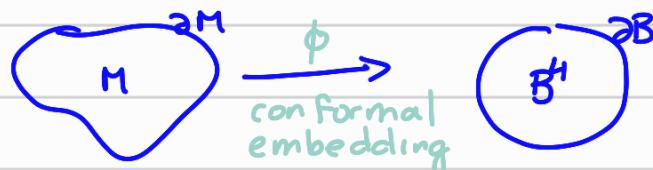
In addition, there is  $\rho^* > 0$  such that for  $\rho \leq \rho^*$  the bound is sharp.

### Proof:

- Variational characterization of eigenvalues

$$\lambda_1 = \min_{\substack{\|u\|=0 \\ Nu=0}} \frac{\int_M (\Delta u)^2 + \dots}{\int_{\partial M} u^2} \quad ) \text{conformally invariant quantity}$$

- Use conformal properties to transform to a model geometry



- Use as test functions:  $U_i = (u_i^{(n)} Y_i) \circ \phi$ ,  
 $i=0,1,\dots,4$

$\uparrow$   
eigenfunctions for  $\lambda_1$   
on  $S^3$

- Hersch trick (calibration).

- Lengthy calculations in model cases.  $\begin{matrix} \text{ball} \\ \text{annulus} \end{matrix}$

- Bifurcation in annulus!

